



Standard Model Electroweak Processes at LHC (and ILC)



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Outline

- Snapshot of the Electroweak Standard Model
- Electroweak Processes at LHC
- LHC Detectors
 - Calibration and Performance with Standard Model Physics
- Electroweak Measurements at the LHC
 - W Mass
 - Gauge Boson Couplings
 - Electroweak Mixing Angle
- Beyond LHC: ILC
- Summary



Snapshot of the Electroweak Standard Model

- **Electroweak Standard Model:**

- developed by Glashow, Salam and Weinberg in the 1960s
- unified gauge theory of QED and weak interactions → gauge group $SU(2)_L \times U(1)_Y$



NP 1979

$$\begin{pmatrix} \nu_e \\ e^- \\ e_R^- \\ \nu_{e,R} \end{pmatrix}_L \quad \begin{pmatrix} \nu_\mu \\ \mu^- \\ \mu_R^- \\ \nu_{\mu,R} \end{pmatrix}_L \quad \begin{pmatrix} \nu_\tau \\ \tau^- \\ \tau_R^- \\ \nu_{\tau,R} \end{pmatrix}_L$$

$$\begin{pmatrix} u \\ d \\ u_R \\ d_R \end{pmatrix}_L \quad \begin{pmatrix} c \\ s \\ c_R \\ s_R \end{pmatrix}_L \quad \begin{pmatrix} t \\ b \\ t_R \\ b_R \end{pmatrix}_L$$

- matter particles are fermions

- leptons and quarks
- left-handed doublets
- right-handed singlets

- interactions are mediated by gauge bosons

- $W^\pm$ and Z bosons, photons
- mixing of neutral fields
- weak mixing angle θ_w
- non-abelian theory:

$$\begin{pmatrix} A \\ Z \end{pmatrix} = \begin{pmatrix} \cos \theta_w & \sin \theta_w \\ -\sin \theta_w & \cos \theta_w \end{pmatrix} \begin{pmatrix} B \\ W_3 \end{pmatrix}$$

- gauge bosons can couple to themselves

- **Higgs mechanism provides masses for fermions and bosons** (lecture on Thursday)



Electroweak Measurements

- masses of gauge bosons: M_W, M_Z
- masses of fermions: the most important is the mass of the top quark m_t
(lecture on Thursday)

- important prediction of the Standard Model: $\frac{M_W}{M_Z} = \cos \theta_w$ or $\rho = \frac{M_W}{M_Z \cos \theta_w} = 1$
(to lowest order)

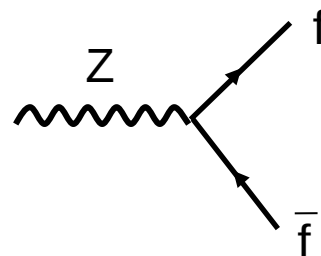
- couplings of fermions to gauge bosons

$$-ig \frac{1}{2} (\mathbf{c}_V \gamma^\mu - \mathbf{c}_A \gamma^\mu \gamma^5)$$

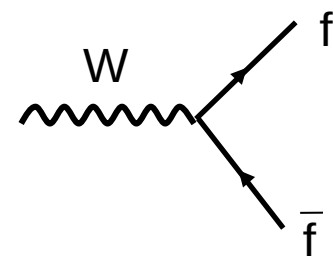
- couplings to Z- and W-bosons involve θ_w :

$$c_A^{f,Z} = T_3 = \pm \frac{1}{2}, 0$$

$$c_V^{f,Z} = T_3 - 2Q \sin^2 \theta_w$$

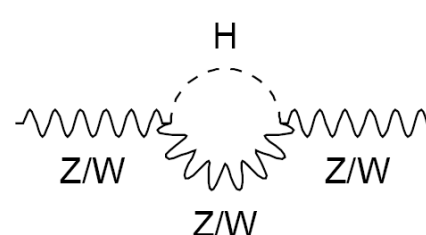
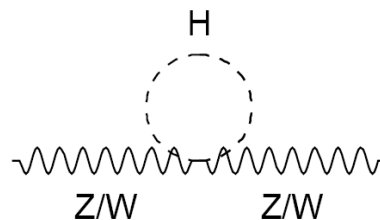
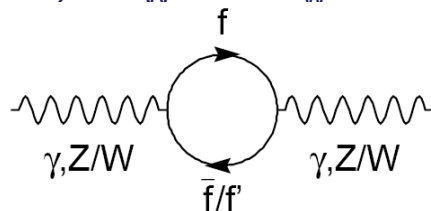


$$g_z = \frac{e}{\sin \theta_w \cos \theta_w}$$



$$g_w = \sqrt{2} \frac{e}{\sin \theta_w}$$

- consequence: measurement of masses and couplings test Standard Model relations
→ $M_Z, M_W, \sin^2 \theta_w$ are precisely measured → radiative corrections need to be included

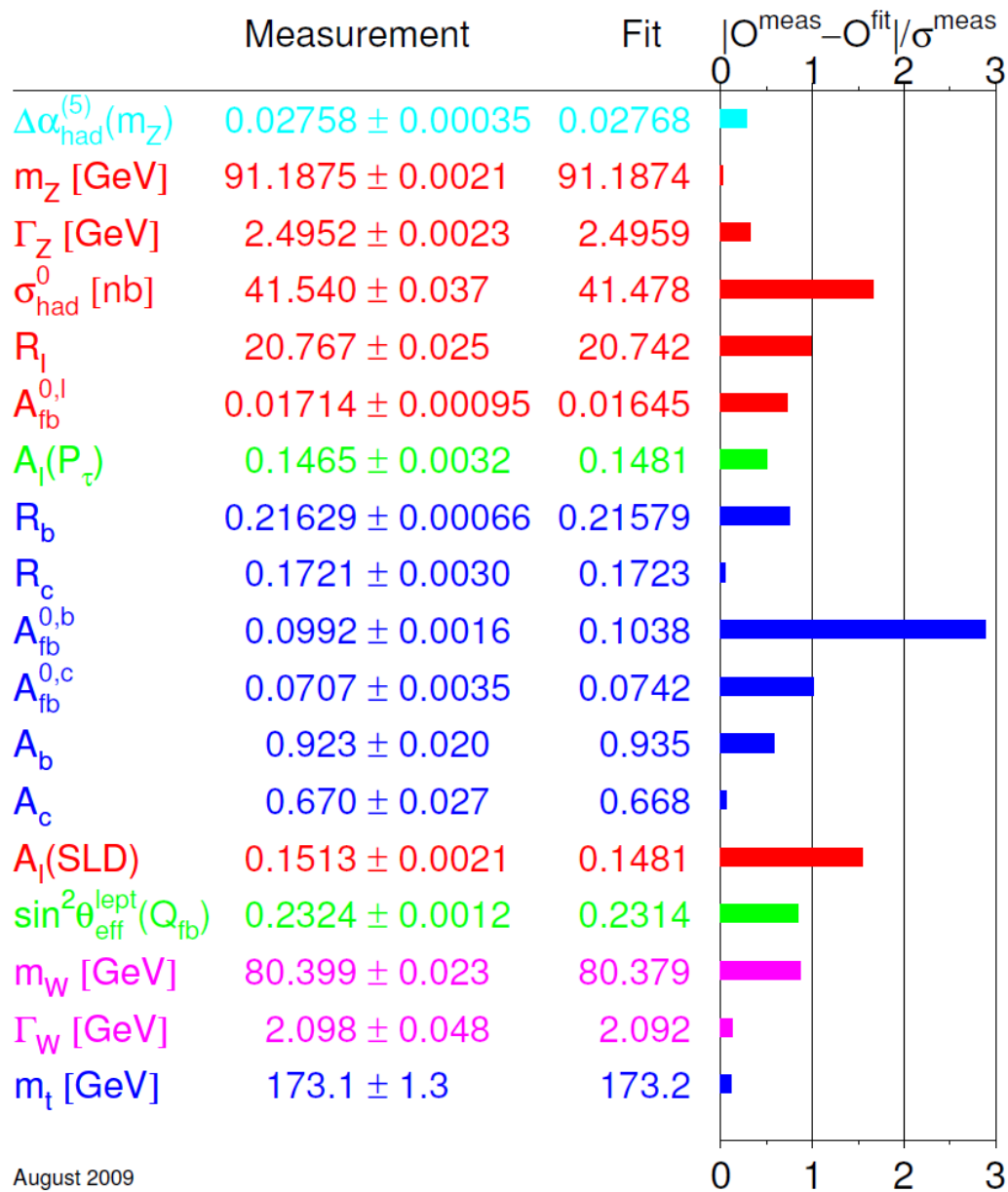


$$\sim m_t^2$$

$$\sim \log (M_H/M_W)$$



Test of the Electroweak Standard Model



- bars show difference between measurement and theory
- relatively good agreement of measurements with predictions
- but room for new physics (SUSY, ...)

August 2009



Physics at the LHC

- the proton-proton collider LHC and the 4 experiments:



- number of physics events:

$$N = L \cdot \sigma$$

L = luminosity

σ = cross-section

- LHC luminosities and centre-of-mass energies

- 2009:

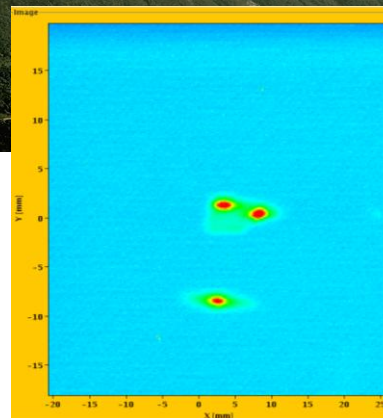
- 900 GeV injection energy
- $9 \mu\text{b}^{-1}$

- 2010/2011

- 7 TeV max. energy
- up to 1 fb^{-1}

- later

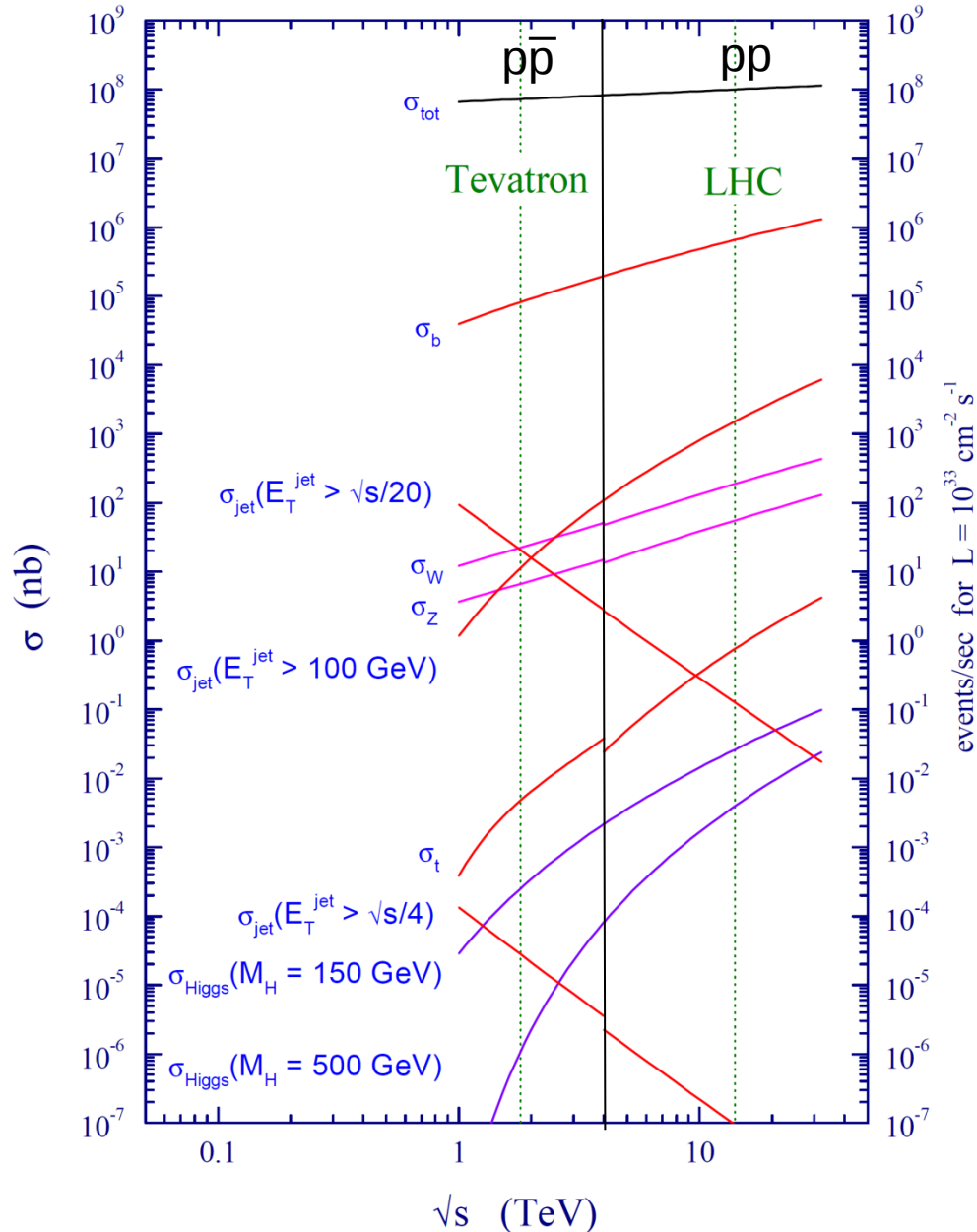
- 13-14 TeV ultimate energy
- $50\text{-}70 \text{ fb}^{-1}$ until 2016
- 3000 fb^{-1} until 2030





Overview of Production Rates at LHC

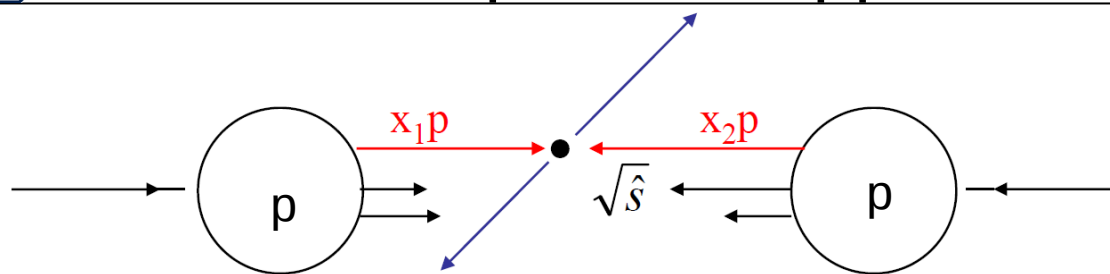
- typical production cross-sections at Tevatron and LHC:



- inelastic pp scattering dominates
→ "minimum bias"
- QCD processes have large cross-sections → 2-jet events, 3-jet events, ...
- many events with W and Z bosons
- top-pair production
- new particles: Higgs, SUSY, ..
- 13 orders of magnitude



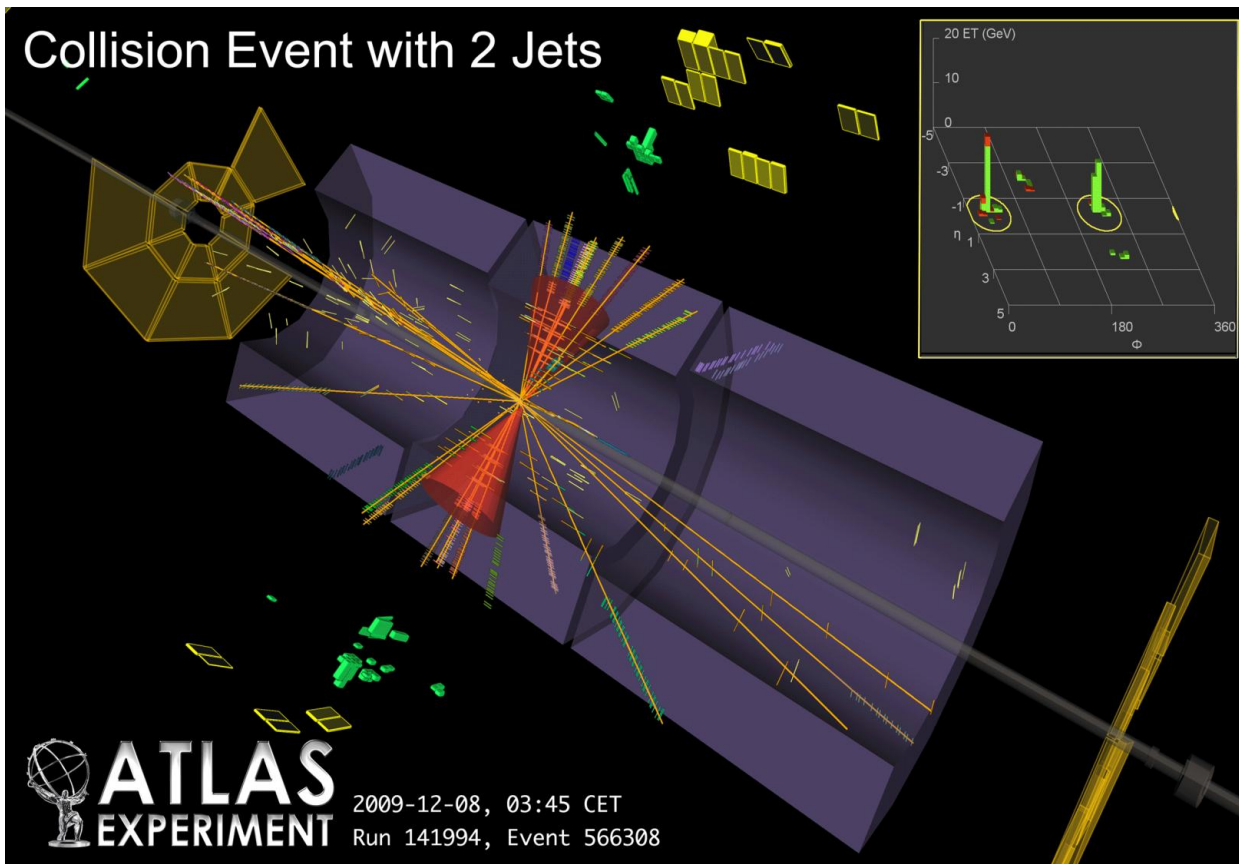
Properties of pp-Collision Events



- hard scattering of partons inside proton:
- gluons and quarks carry momentum fractions x_1, x_2
- in the event picture:
 $qq \rightarrow qq, gg \rightarrow gg, gq \rightarrow gq$
- QCD events are more frequent than EW events

- many hadrons produced from proton remnants:
"underlying event"
- the typical LHC event has no hard scattering:
"minimum bias" $\rightarrow$ pile-up
- strategy: identify the interesting physics with high- p_T leptons, photons and jets or missing transverse energy E_T^{miss}

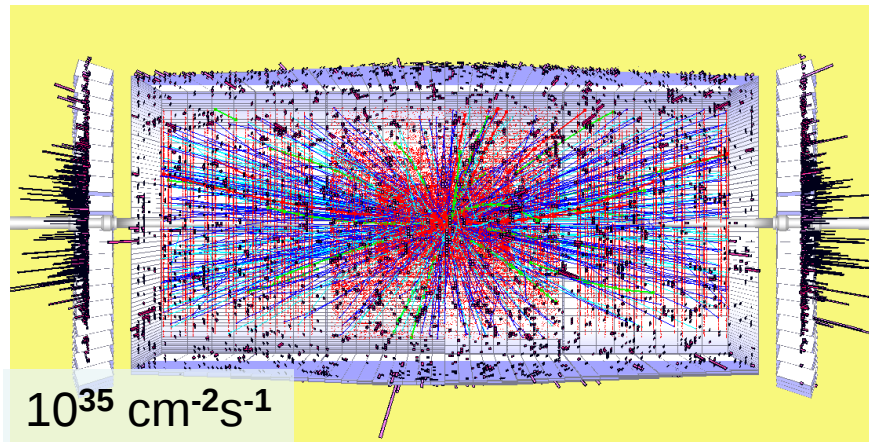
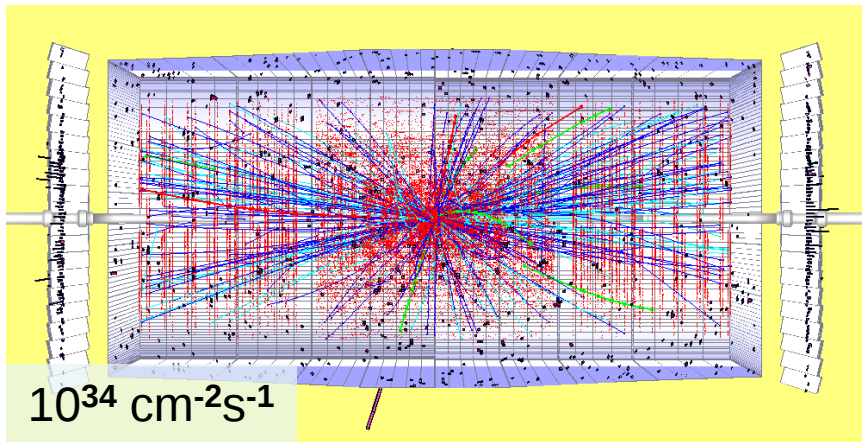
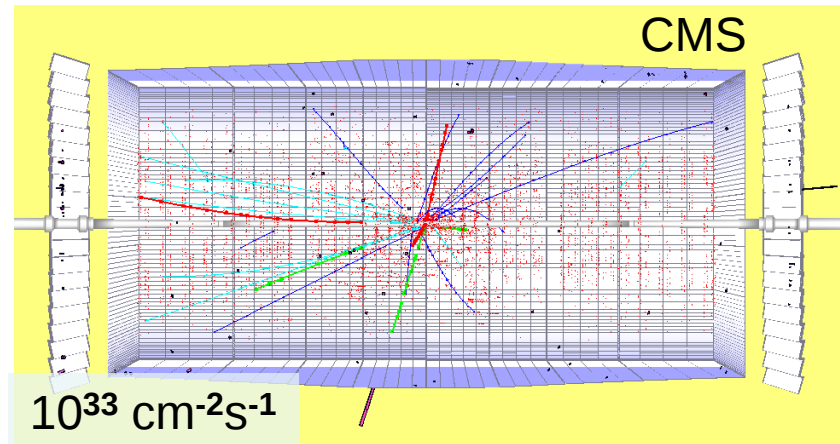
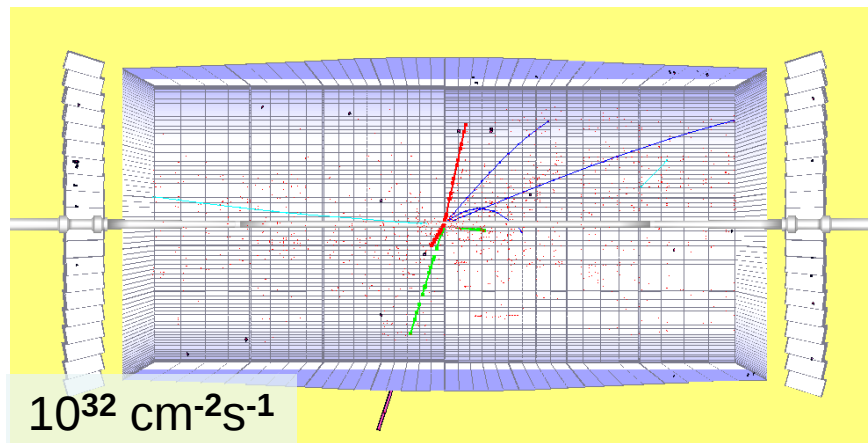
- a 2-jet event at 900 GeV in ATLAS:





Effect of Event Pile-Up

- number of simultaneous pp reactions increase with luminosity:



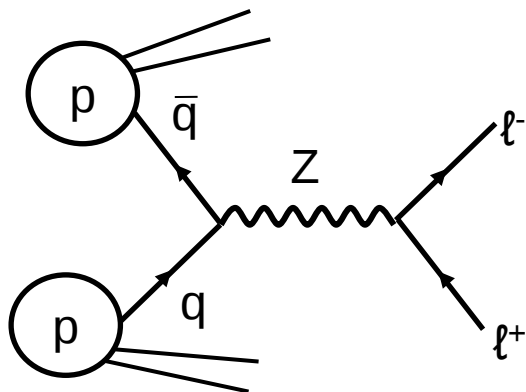
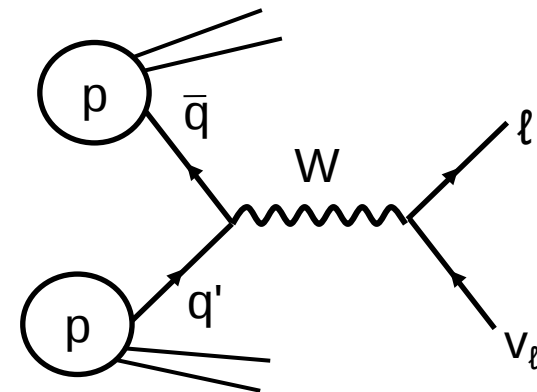
- at nominal LHC luminosity of $10^{34} \text{ cm}^{-2}\text{s}^{-1}$ about 25 events per bunch crossing



Electroweak Physics at the LHC

- W boson production:

- cross-section measurement in $W \rightarrow e\nu, \mu\nu, \tau\nu$ channels (hadronic final states are difficult to select at LHC)
- measurement of the W boson mass
- background to search channels (W+jets)



- Z boson production:

- cross-section measurement in $Z \rightarrow e\bar{e}, \mu\bar{\mu}, \tau\bar{\tau}$ channels
- calibration of the detector
- $\sin^2\theta_w$ from forw.-backw. asymmetry in $Z \rightarrow e\bar{e}$ events
- measurement of systematic uncertainties for other processes, background to search channels (Z+jets)

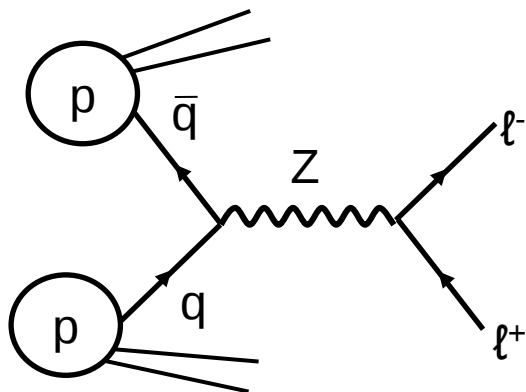
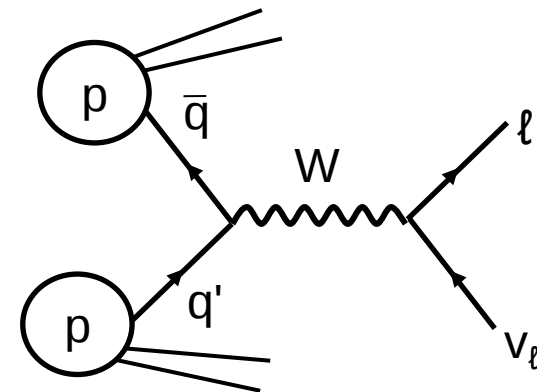
Process	$\sigma_{\text{Lo}} @ 7 \text{ TeV}$ [pb]	$\sigma_{\text{Lo}} @ 14 \text{ TeV}$ [pb]	Evts (7 TeV) in 200 pb ⁻¹	Ratio 7 TeV/14 TeV
QCD pt>100 GeV	3.2 E+05	1.4 E+06	64 M	0.2
Z inclusive	2.4 E+04	5.7 E+04	5 M (3% $Z \rightarrow \ell\bar{\ell}$)	0.4
W inclusive	9.5 E+04	2.1 E+05	19 M (11% $W \rightarrow \ell\nu$)	0.5
top pairs	8.4 E+01	4.7 E+02	17000	0.2
H (m=150 GeV)	4.0	16.0	800	0.3



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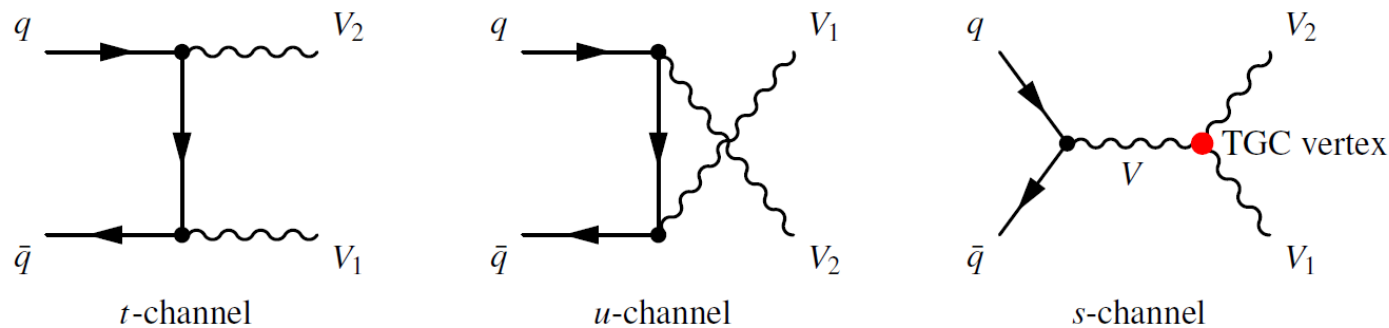
- cross-section measurement in $Z \rightarrow ee, \mu\mu, \tau\tau$ channels
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Di-boson Production at the LHC

- measurement of W -pairs, Z -pairs, WZ , $W\gamma$, and $Z\gamma$ final states



- determination of production cross-sections
- measurement of the couplings between gauge bosons $\rightarrow$ prediction of the SM
 - triple-gauge boson vertex
- examples of cross-sections:

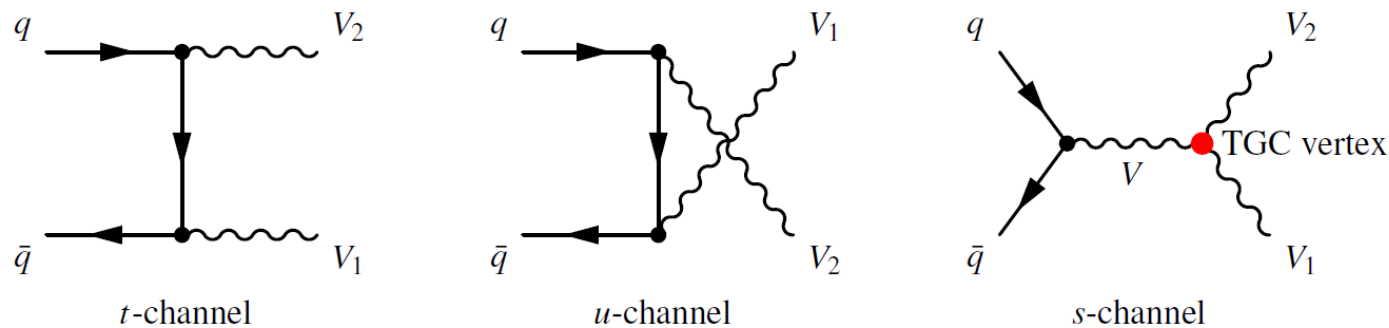
Diboson mode	Conditions	$\sqrt{s} = 14 \text{ TeV}$ $\sigma [pb]$
W^+W^- [15]	W -boson width included	111.6
$W^\pm Z$ [15]	Z and W on mass shell	47.8
ZZ [15]	Z 's on mass shell	14.8
$W^\pm \gamma$ [16]	$E_T^\gamma > 7 \text{ GeV}$, $\Delta R(\ell, \gamma) > 0.7$	451
$Z\gamma$ [17]	$E_T^\gamma > 7 \text{ GeV}$, $\Delta R(\ell, \gamma) > 0.7$	219

factor 100-1000
smaller than
 W, Z production



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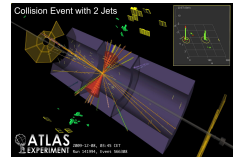
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ATLAS and CMS Detectors



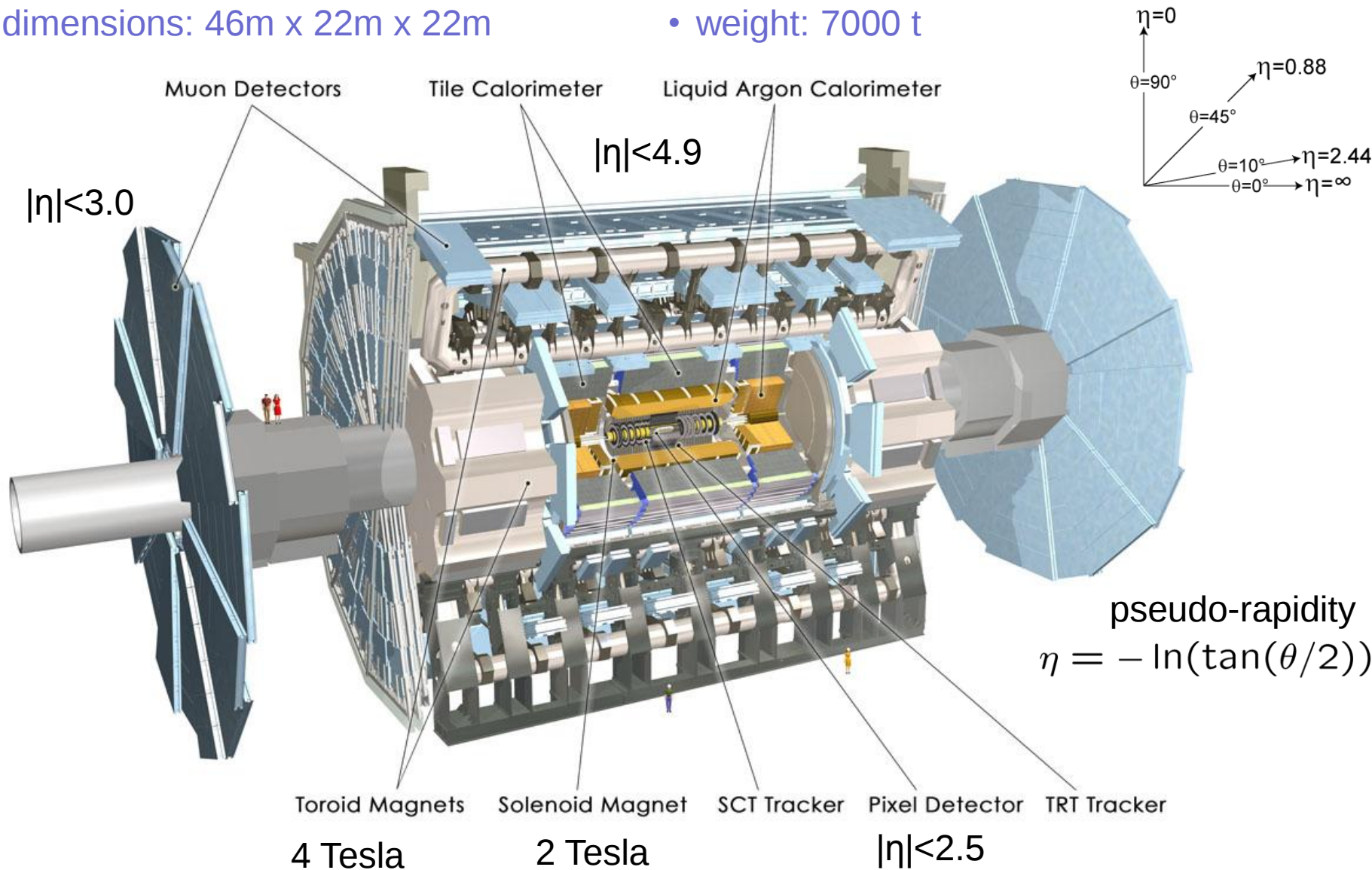
- challenges for LHC detectors:
 - high collision rate: 25 ns proton bunch-crossing
 - up to 25 minimum-bias events/crossing → pile-up
 - radiation hardness → detectors and electronics tolerate 10^{15} particles/cm²
 - high granularity → isolate interesting leptons/photons/jets from pile-up
 - very good angular coverage → measure missing energy (neutrinos, SUSY)
 - fast trigger with high background suppression rate (~ 1 accepted event/ 10^6)
- performance goals:
 - get high efficiency for particle detection
 - suppress background from QCD jets

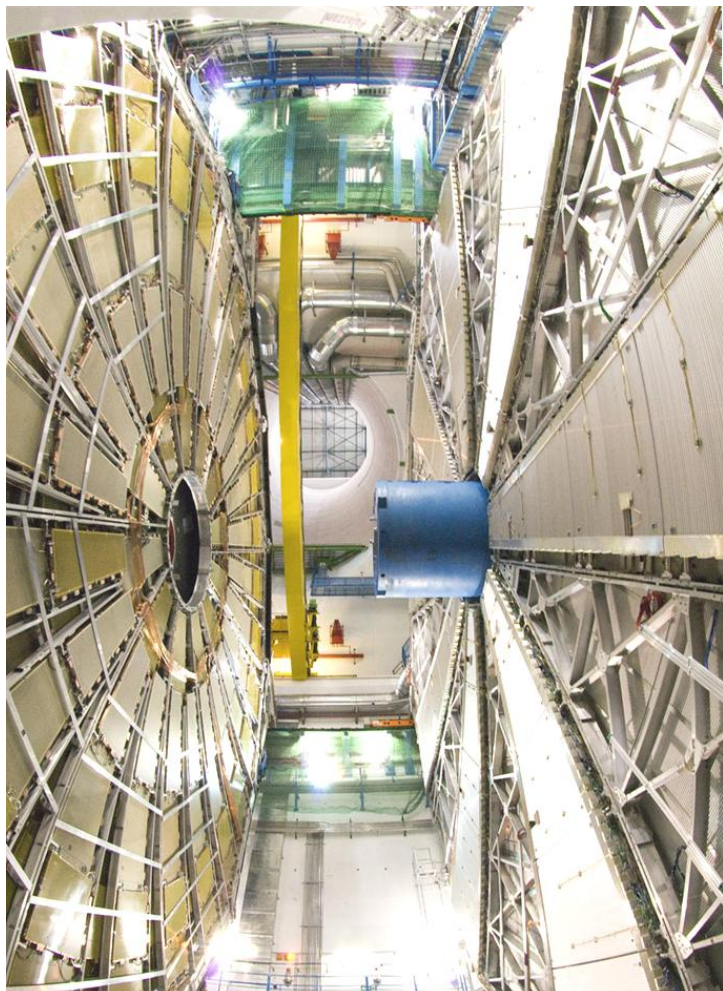
Object	Efficiency	BG suppr.	Motivation	K. Desch
electrons	>70%	100000	Z,W,top,H $\rightarrow 4\ell$,...	
photons	80%	1000	H $\rightarrow \gamma\gamma$	
muons	>97% for $p_T > 1\text{GeV}$		H $\rightarrow 4\ell$,..., z, w	
tau leptons	50%	100	H/A $\rightarrow \tau\tau$, SUSY,...	
b-jets	50%	100	top, H $\rightarrow bb$, SUSY,...	

The ATLAS Detector

• dimensions: 46m x 22m x 22m

• weight: 7000 t

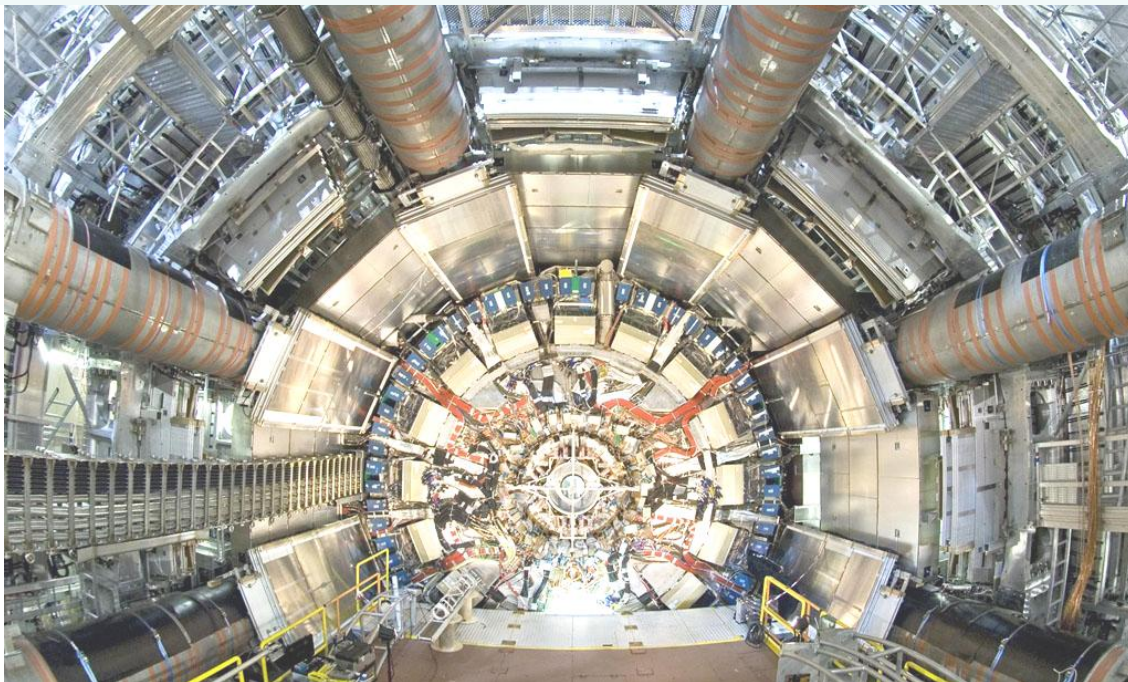




muon spectrometer:
1000 chambers with monitored
drift tubes
500 trigger RPC chambers

Arno Straesser - Star

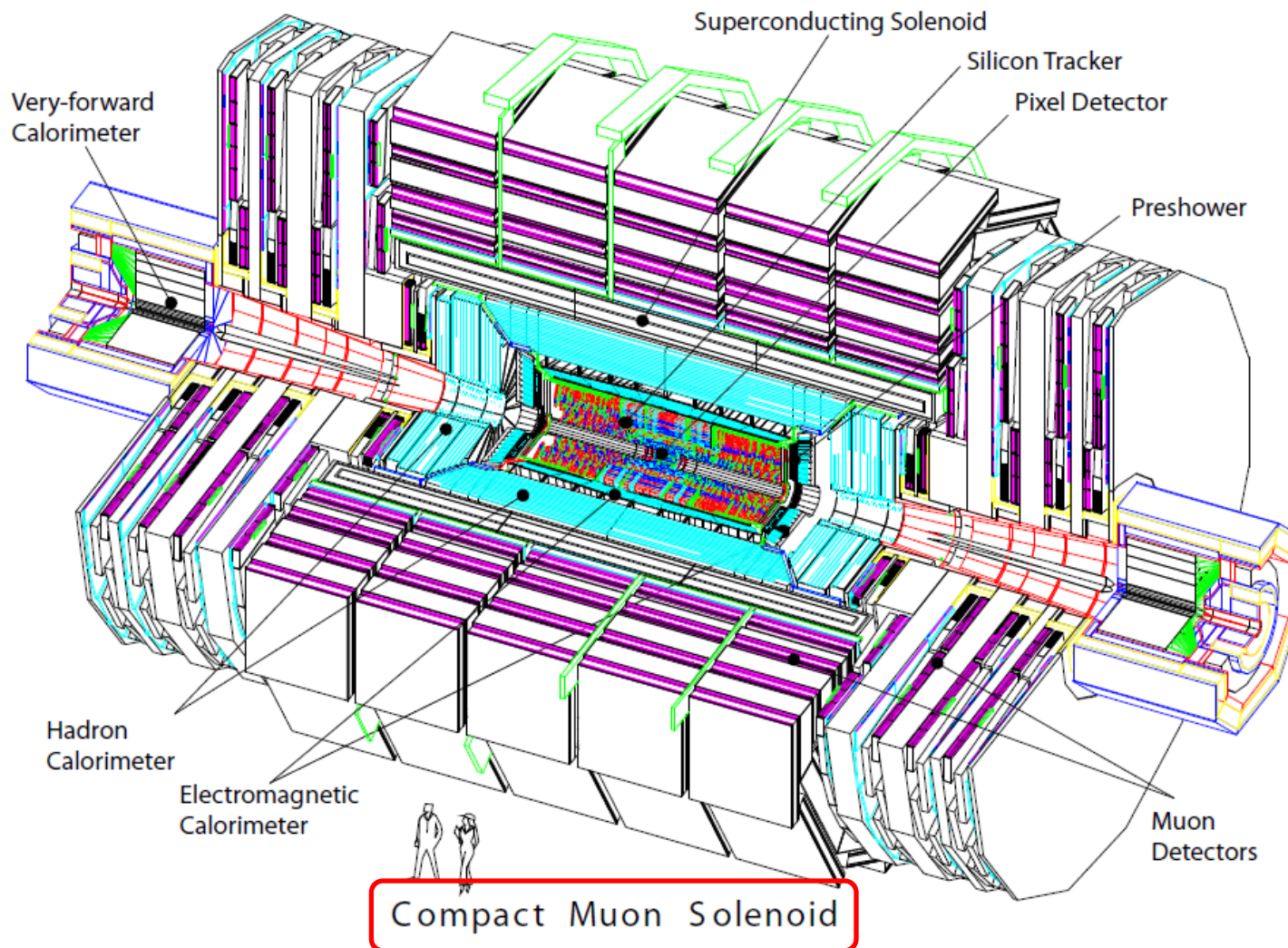
toroid magnet - LAr/Tile calorimeter - silicon tracker



pixel detector:
80 M channels
pixel size $50 \times 400 \mu\text{m}^2$

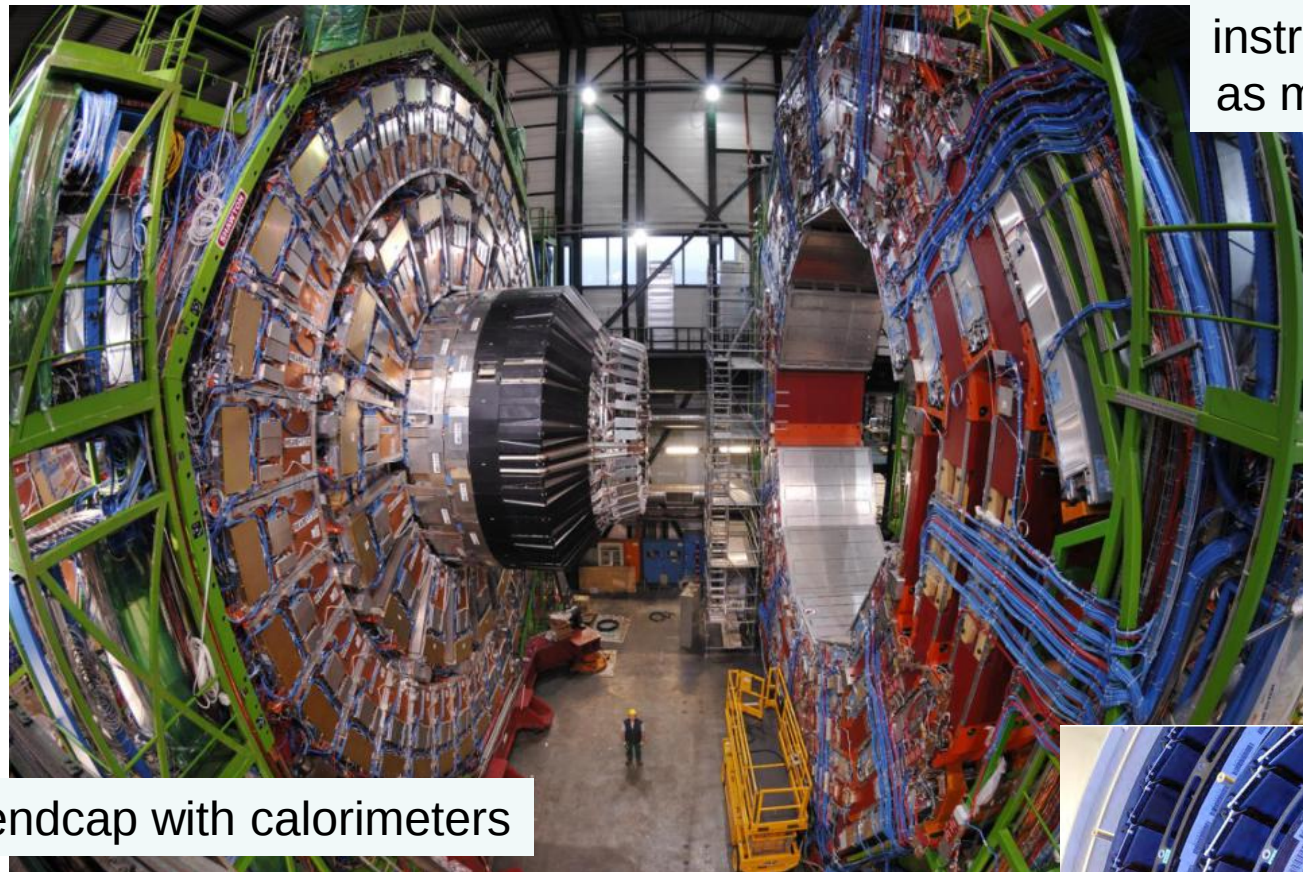
The CMS Detector

- dimensions: 22m x 15m x 15m
- weight: 12500 t



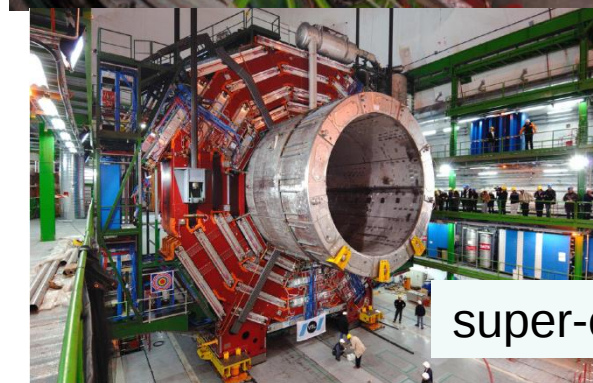
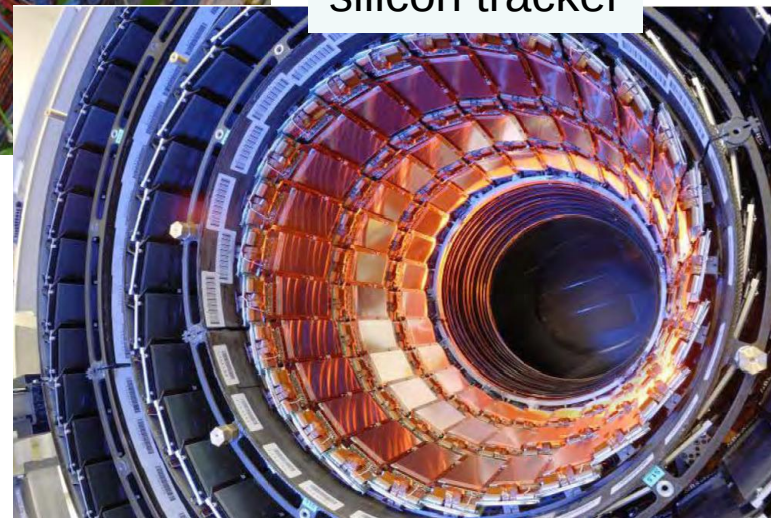


Detector Components of CMS



instrumented iron yoke
as muon spectrometer

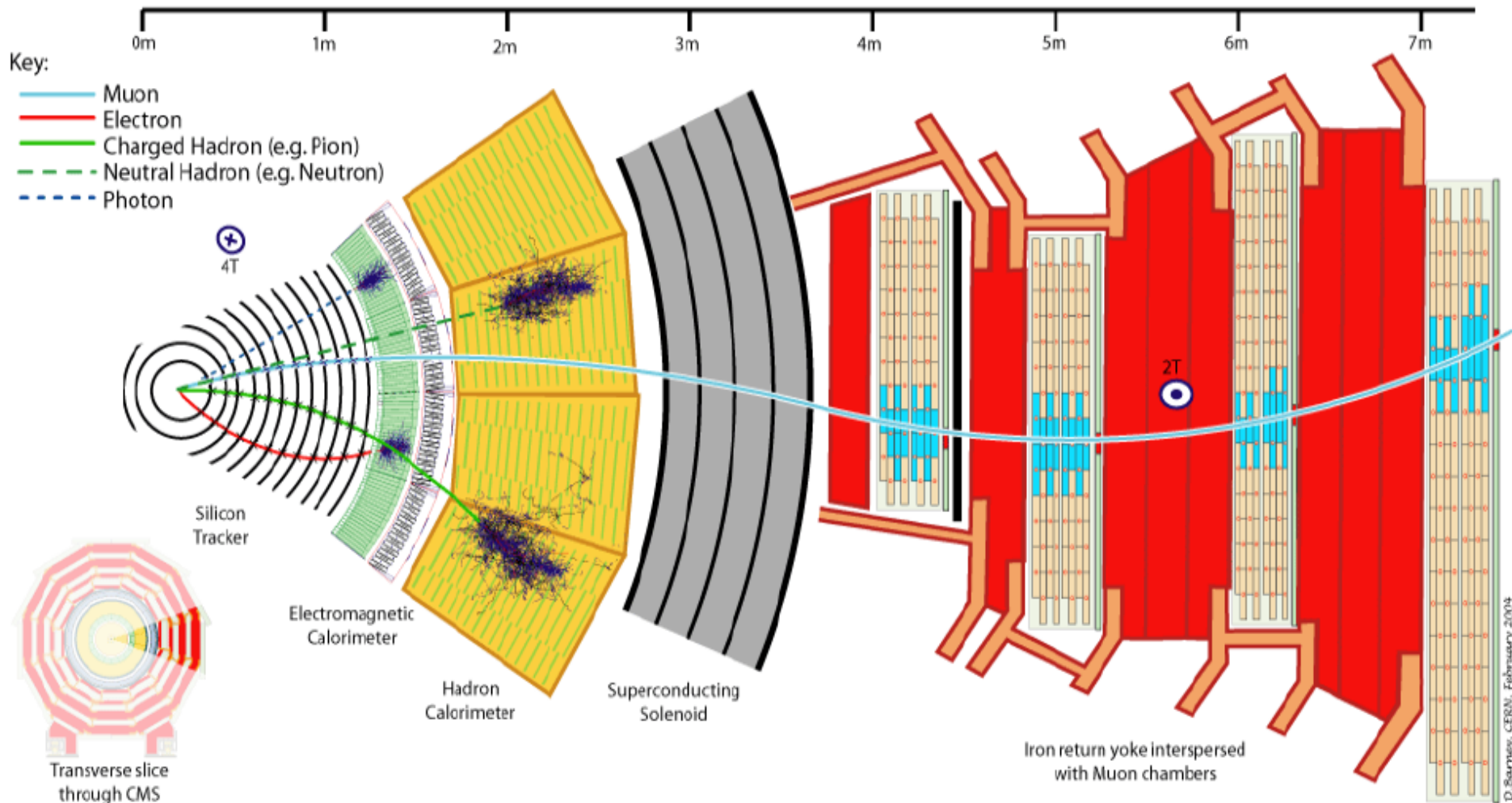
silicon tracker



super-conducting magnet



Particle Identification



Charged
Particles

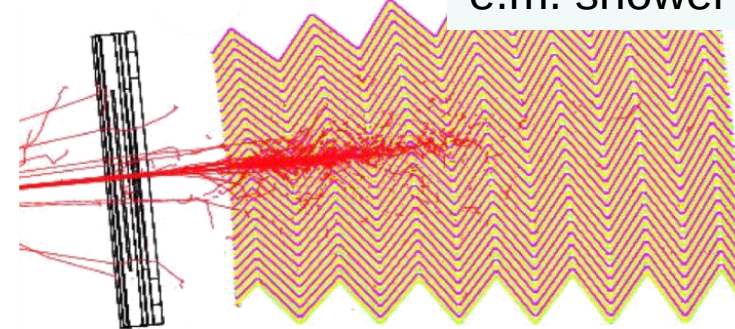
Electrons
Photons

Hadrons

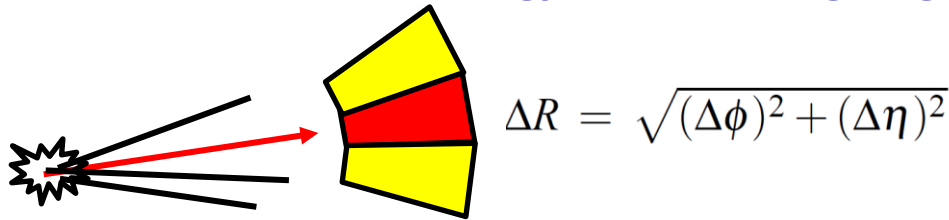
3.8 Tesla
Solenoid

Iron yoke (return flux)
Muons

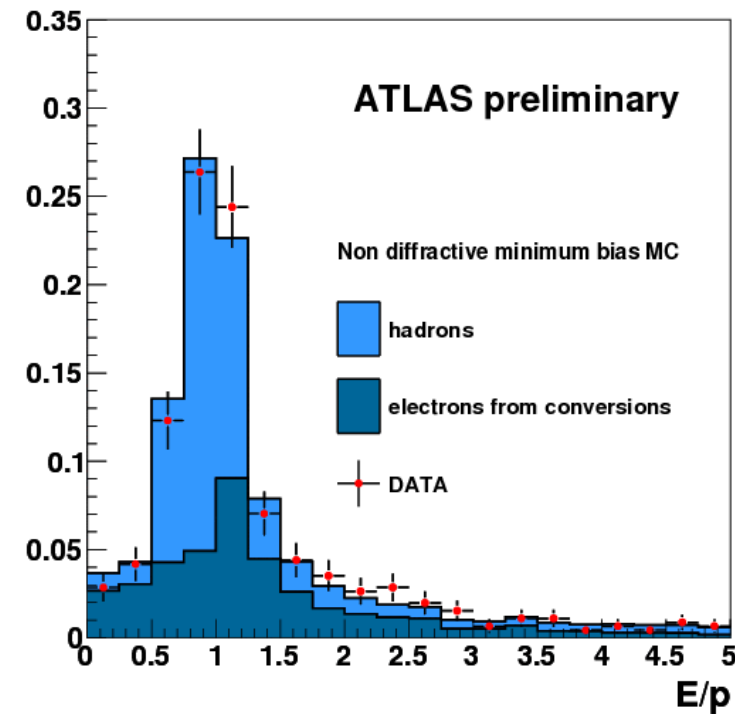
e.m. shower



- e.g. electron:
- energy and shower shape measured in calorimeter
- track matched to energy cluster using angles

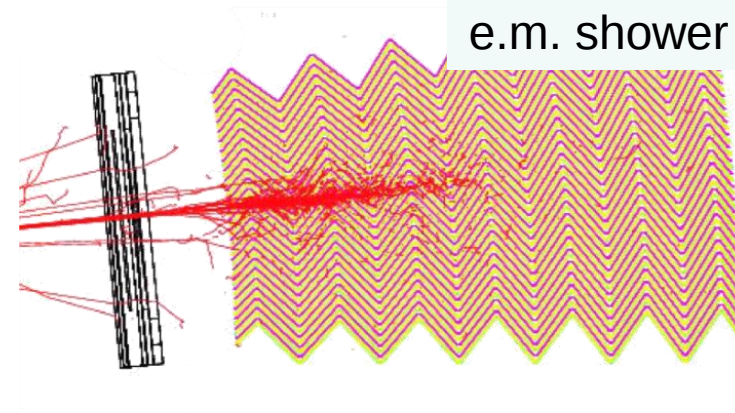


- compare calorimetric energy and momentum measurement in tracker: $E/p \approx 1$
- no leakage of energy into hadronic calorimeter
- combined performance is important

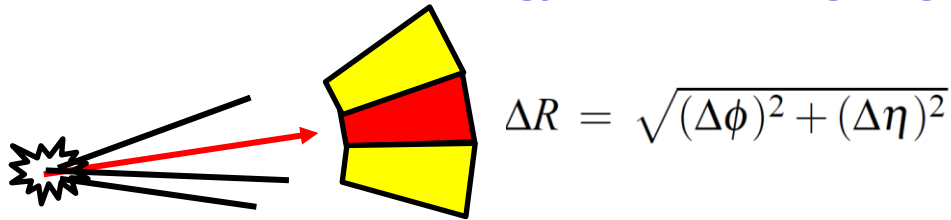


2009 data

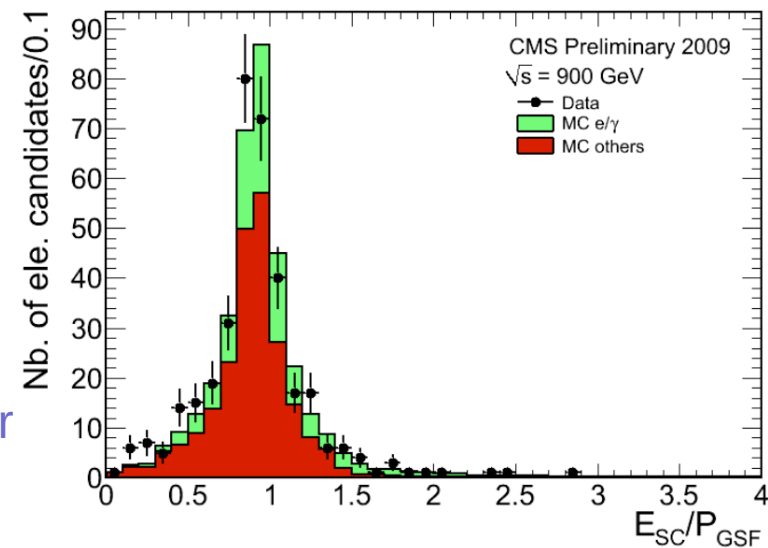
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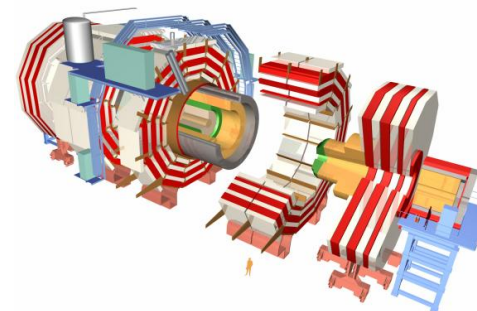
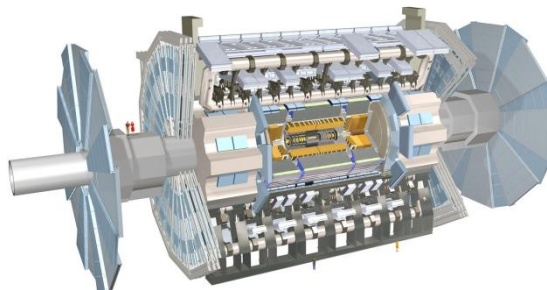
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2009 data



Performance Comparison

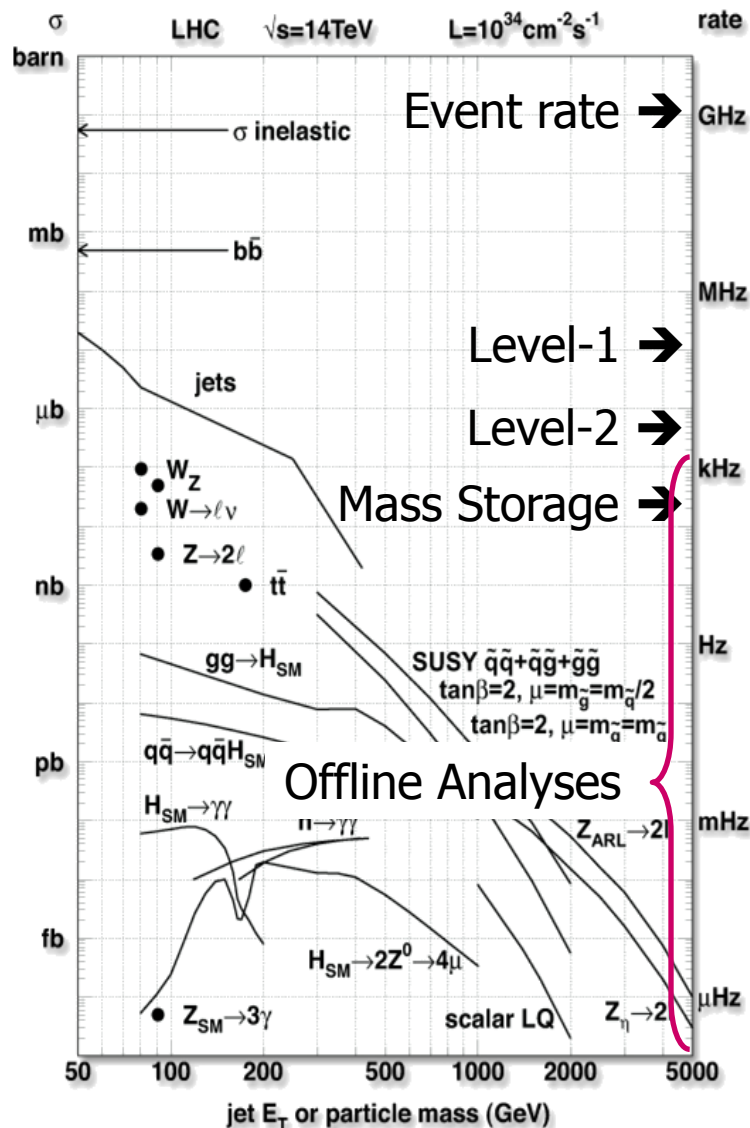


	ATLAS		CMS	
B-Field	2 T (solenoid)	4 T (toroid)	3.8 T (solenoid)	
Tracking	Si (strips+pixel) + Gas (TRT)	$\sigma(p_T)/p_T$ at 100 GeV: 3.8 %	Si (strips+pixel)	$\sigma(p_T)/p_T$ at 100 GeV: 1.5 %
ECAL	Pb - LAr (high granularity)	$\sigma(E)/E = 9\%/\sqrt{E} + 0.7\%$	PbWO ₄ crystals (high E resolution)	$\sigma(E)/E = 3\%/\sqrt{E} + 0.25\%$
HCAL	Fe – scintillator (10 λ)	$\sigma(E)/E$ (ECAL+HCAL)= $70\%/\sqrt{E} + 3.3\%$	Brass – scintillator (7 λ)	$\sigma(E)/E$ (ECAL+HCAL)= $70\%/\sqrt{E} + 8\%$
Muon Spectrometer	ion. chambers, air-core magnet	$\sigma(p_T)/p_T$ at 1TeV: 7 %	ion. chambers, instrumented iron	$\sigma(p_T)/p_T$ at 1TeV: 5 %

- D. Froidevaux, P. Sphicas, General-purpose Detectors for the Large Hadron Collider, Annu. Rev. Nucl. Part. Sci. 2006. 56:375–440



From Detector To Physics: The Trigger



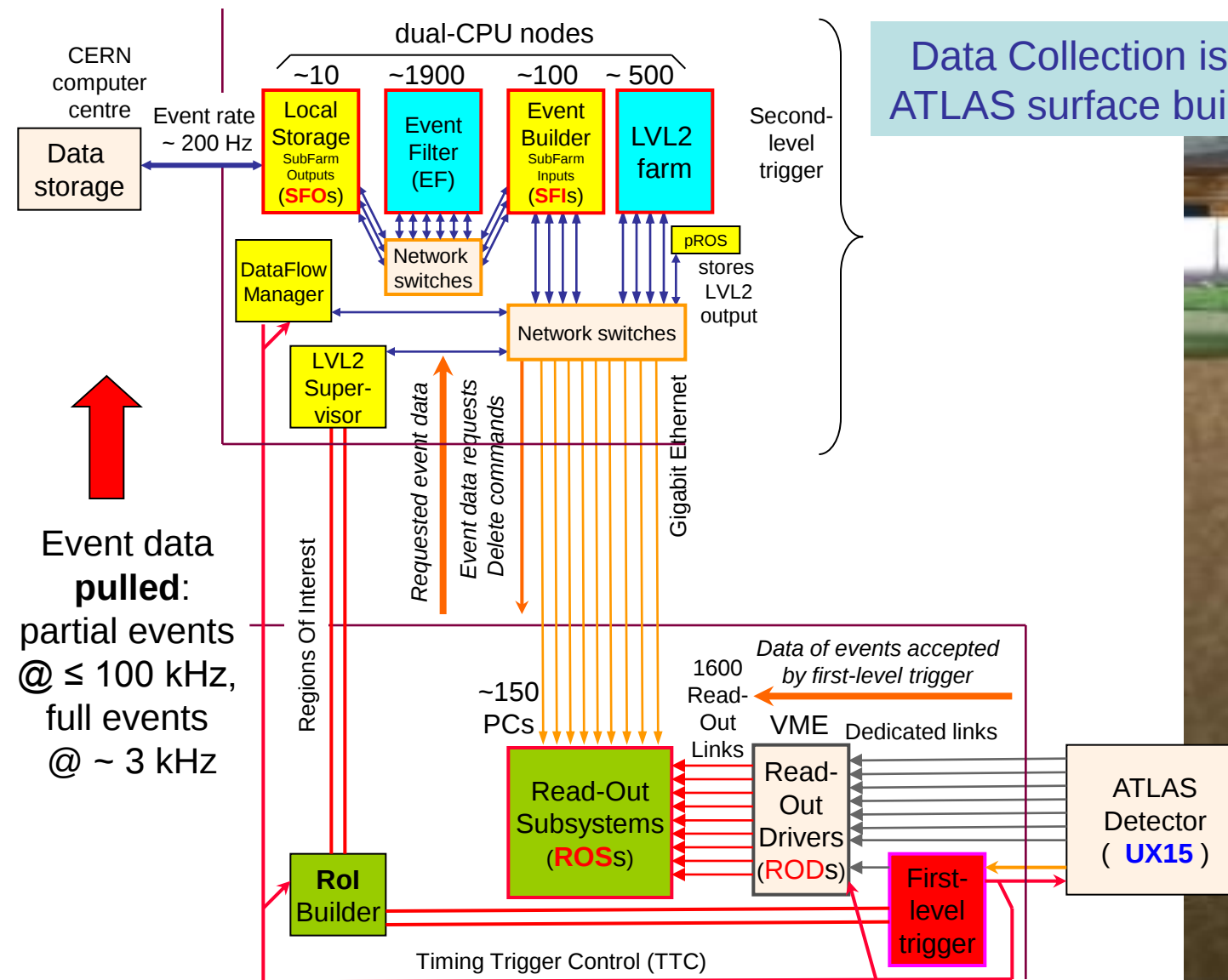
- Bunch crossing 40 MHz
- σ total 70 mb
- Event rate ~ 1 GHz
- Number of events/BC ~ 25
- Number of particles/event ~ 1500
- Event size ~ 1.5 MB
- Mass storage rate ~ 200 Hz

- Need to have trigger of high performance
 - 6 orders of rate reduction
 - complex events and 140 M channels

- Level-1: hardware based at 40 MHz
 - Level-2: software based at 100 kHz
 - Level-3: event filter at 3 kHz
 - Storage at 200 Hz
- } PC farms



ATLAS Trigger & Data Acquisition at Point-1



Data Collection is situated in ATLAS surface building (SDX1)

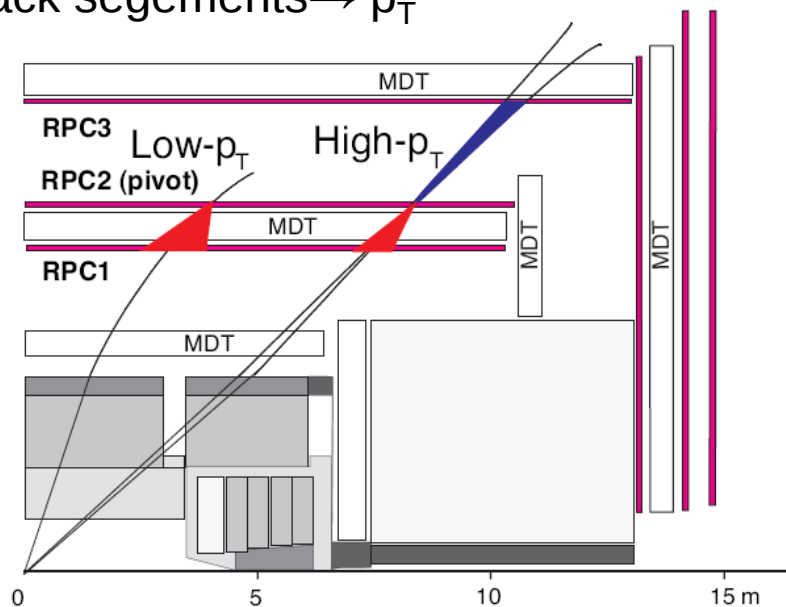


Event data **pushed** @ ≤ 100 kHz,
1600 fragments of ~ 1 kByte each
Arno Straesser - Standard Model Electroweak Processes at LHC



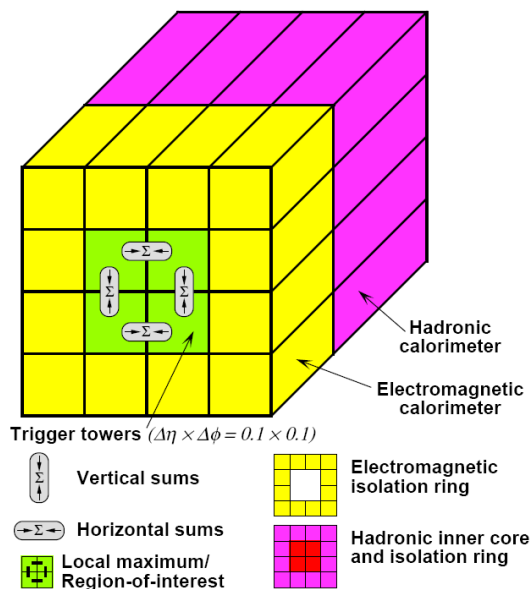
Muon and Calorimeter Level-1 Triggers

- muon trigger detectors: RPC = resistive plate chamber
- pre-defined track segments $\rightarrow p_T$



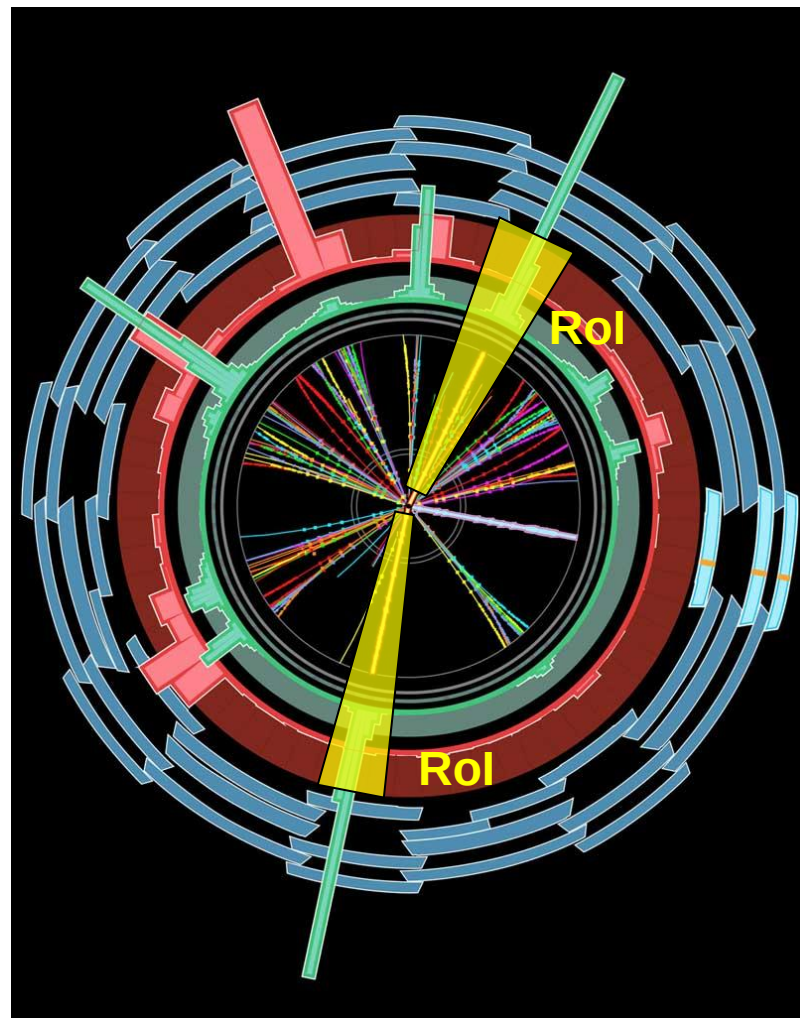
muon trigger receiver

- calorimeter: "trigger-tower"
- energy sums in pre-defined (η, ϕ) regions



calo trigger receiver

- early event rejection in each stage
- L1 muon trigger:
low/high p_T signatures
- L1 calo trigger:
energy sums in calorimeter towers
- L1 Region-of-Interest $\rightarrow$ L2
- high level trigger (software):
 - L2:
 - local reconstruction within RoI
 - tracking information available
 - event filter (EF):
 - access to fully reconstructed event
- CMS trigger is conceptually similar

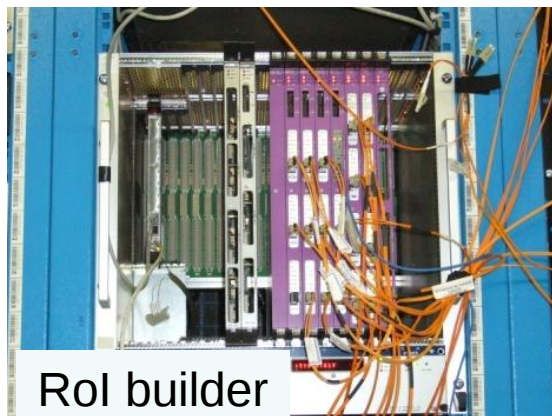




High Level Trigger

- trigger hardware:

read-out buffer PCs



RoI builder



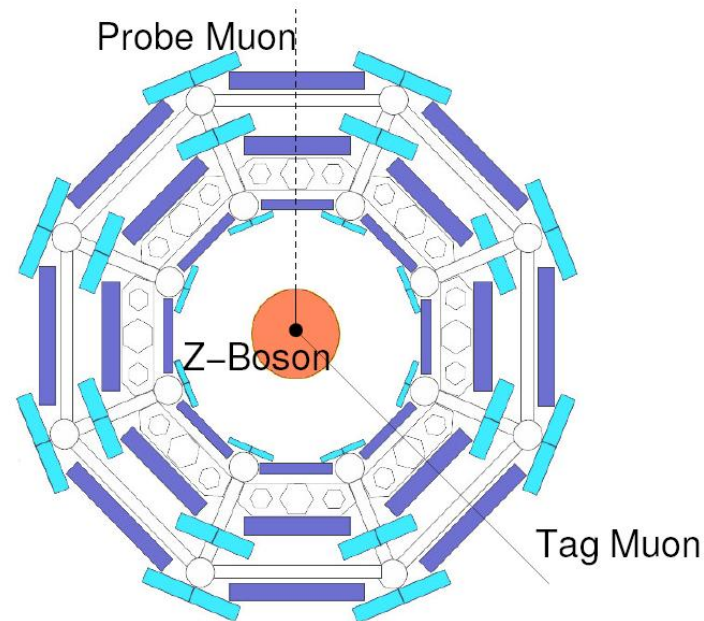
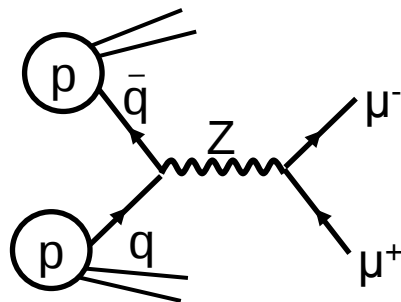
event builder PCs



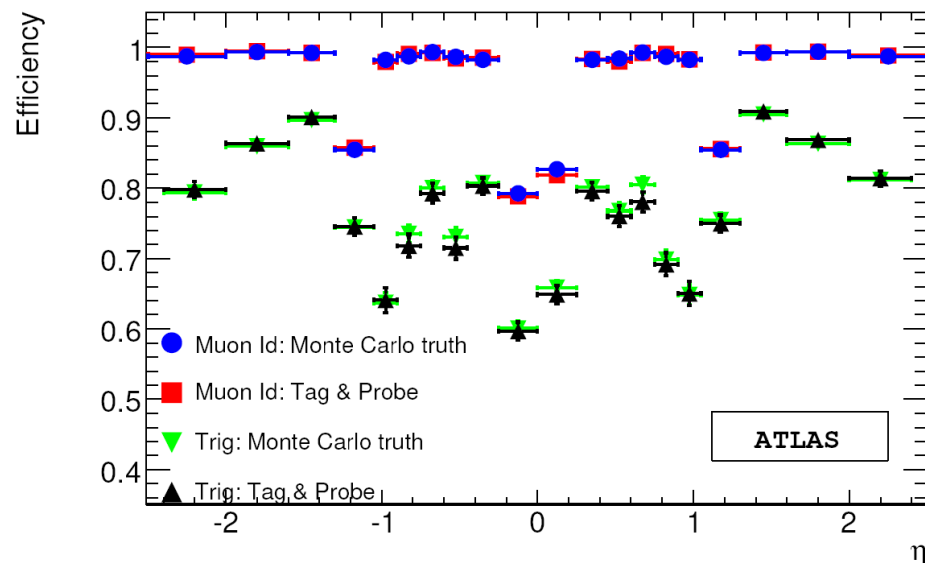
HLT farm



farm output



- commonly applied method to measure performance in data: → **tag & probe**
- example: measure muon identification efficiency in $Z \rightarrow \mu\mu$ events
 - identify triggered and well measured muon: → **tag**
 - use Z decay kinematics to find the 2nd reconstructed muon: → **probe**
 - check if 2nd muon passed selection → **muon identification efficiency**
 - similar for muon trigger and muon reconstruction efficiency
- many more examples:
 - $Z \rightarrow ee$, $J/\psi \rightarrow ee$, ...





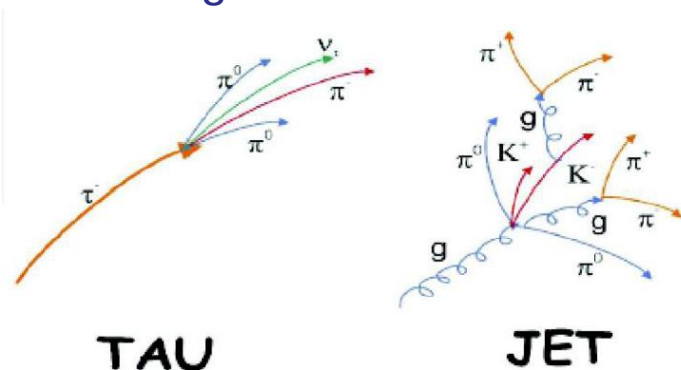
Detector Performance with Z Bosons

- identification of hadronic tau leptons
 - narrow hadronic jet
 - 1 or 3 particle tracks
- tau may look like hadronic jets from quarks and gluons

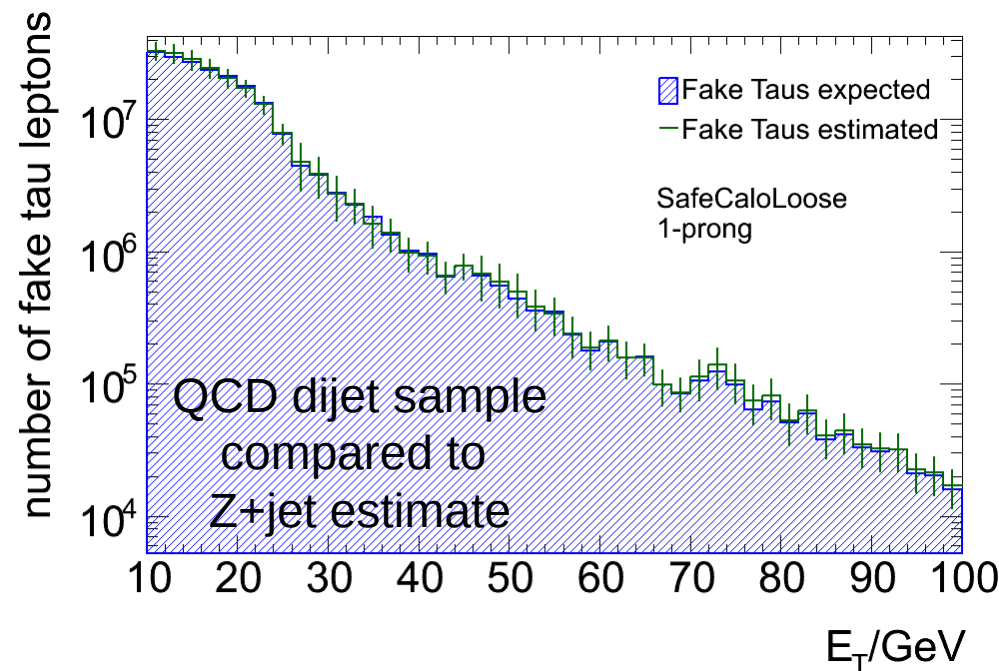
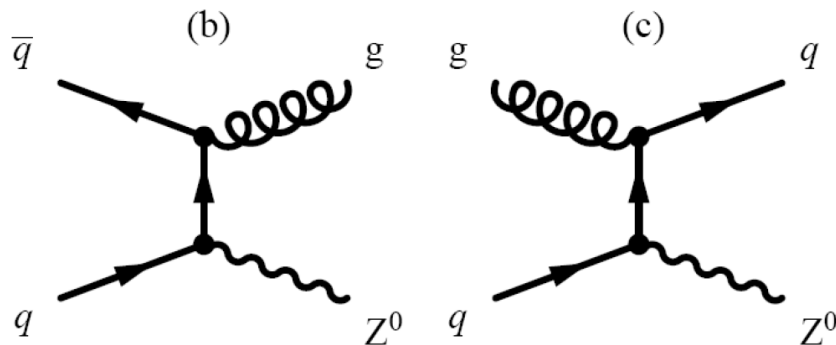
τ^- decay mode	BR %	
$e^- \bar{\nu}_e \bar{\nu}_\tau$	17.85	leptonic
$\mu^- \bar{\nu}_\mu \bar{\nu}_\tau$	17.36	
$\pi^- \bar{\nu}_\tau$	10.91	hadronic
$\pi^- \pi^0 \bar{\nu}_\tau$	25.52	1-prong
$\pi^- 2\pi^0 \bar{\nu}_\tau$	9.27	
$\pi^- 3\pi^0 \bar{\nu}_\tau$	1.04	
$K^- \bar{\nu}_\tau + \text{Neutrals}$	1.57	
$\pi^- \pi^+ \pi^- \bar{\nu}_\tau$	9.32	hadronic
$\pi^- \pi^+ \pi^- \pi^0 \bar{\nu}_\tau$	4.61	3-prong
$K^- \pi^+ \pi^- \bar{\nu}_\tau + \text{Neutrals}$	0.48	

72%

23%



- $Z + \text{jet} \rightarrow \ell\ell + \text{jet}$ events help to estimate how often a jet is misidentified as a tau



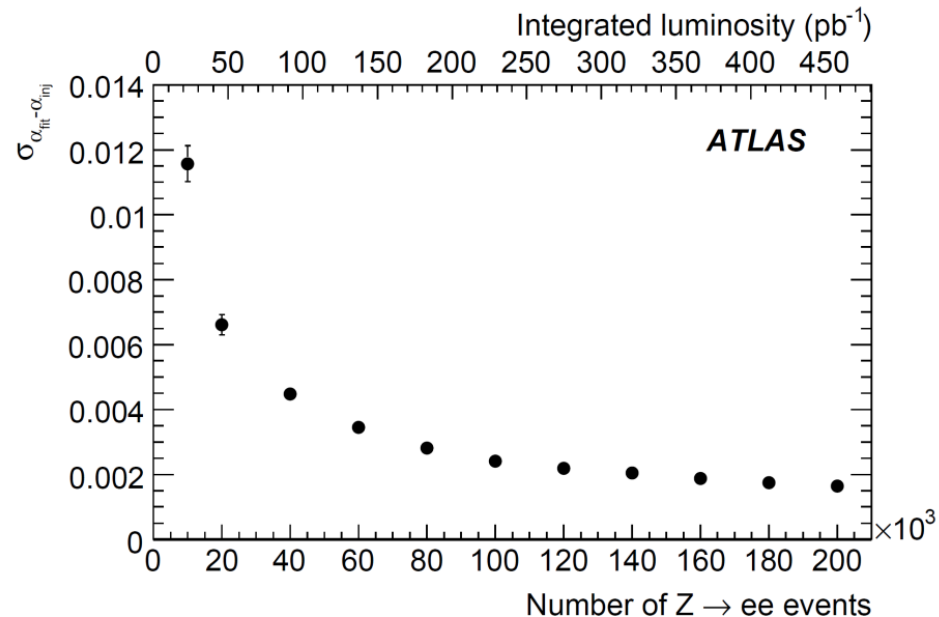
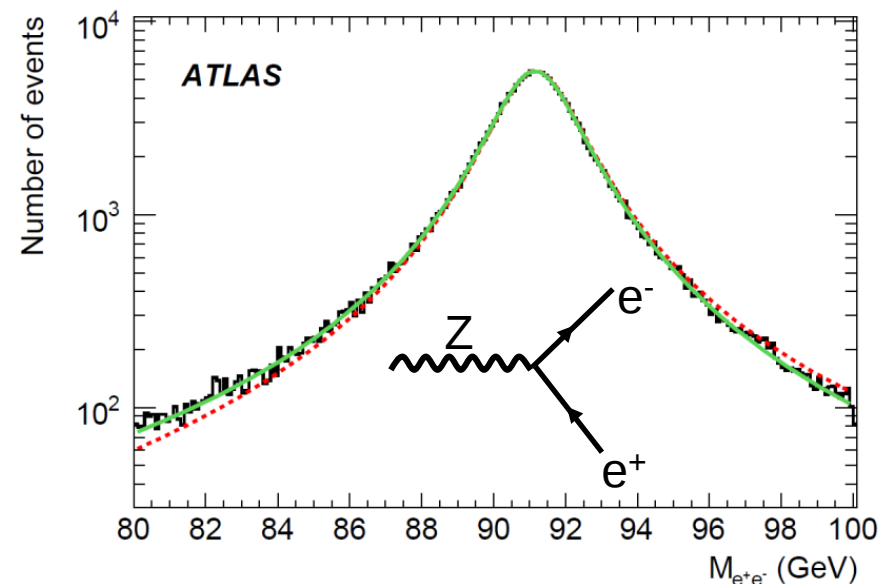


Detector Calibration with Z Bosons

- reconstruct pairs of electrons and positrons in the detector
- measure the invariant mass of electron and positron M_{ee}
- the mass spectrum should peak at the Z boson mass M_Z
- line-shape modelled with Breit-Wigner distribution:

$$BW(M) \sim \frac{M^2}{(M^2 - M_Z^2)^2 + \Gamma_Z^2 M^4 / M_Z^2}$$

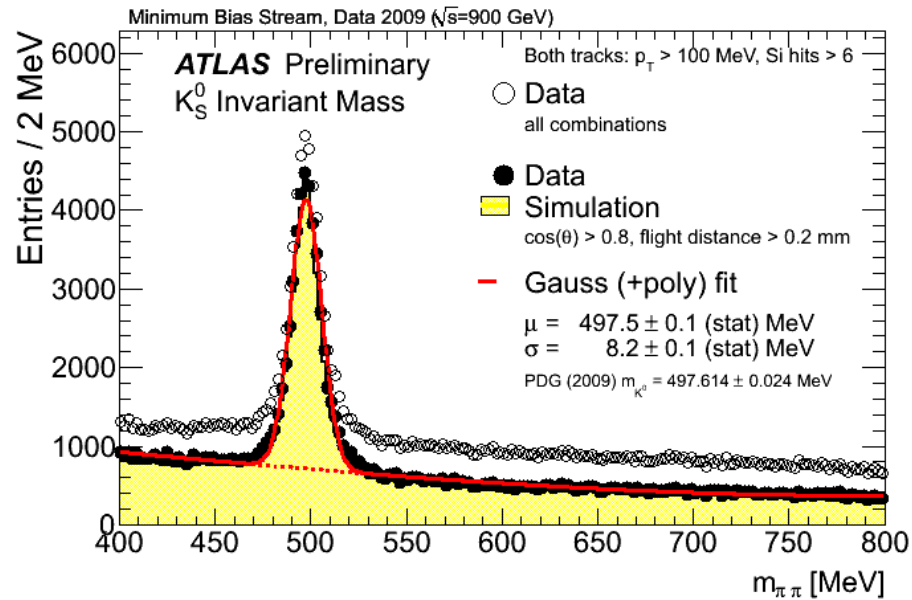
- use the Z mass constraint from LEP
 $M_Z = 91.1875$ GeV to calibrate the energy scale of the calorimeter
- e.g. constant term of the energy resolution:
140k $Z \rightarrow ee$ events needed to reach 0.2%





Detector Performance with First Physics Data

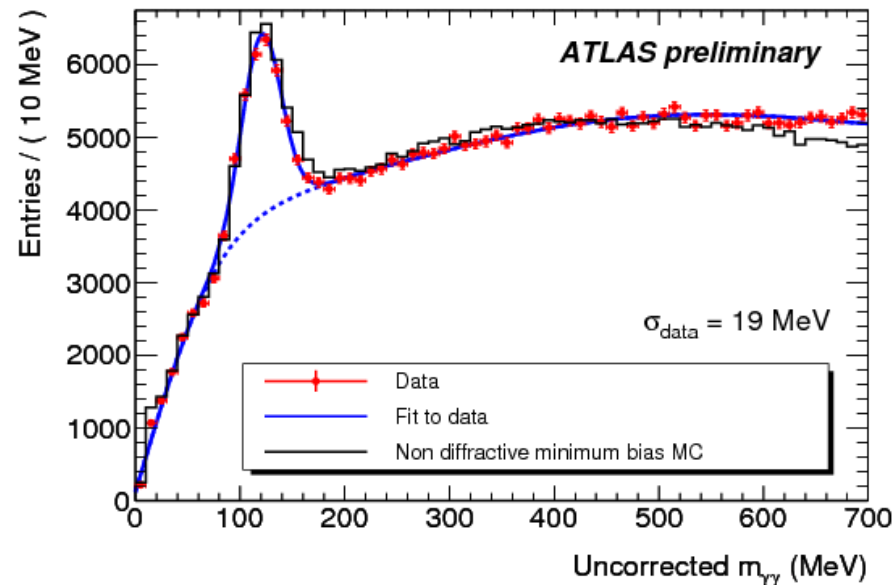
- tracking performance
using K_S mesons:



- calorimeter performance
with $\pi^0 \rightarrow \gamma\gamma$

peak value agrees with pion
mass to 1% (1.5 MeV)

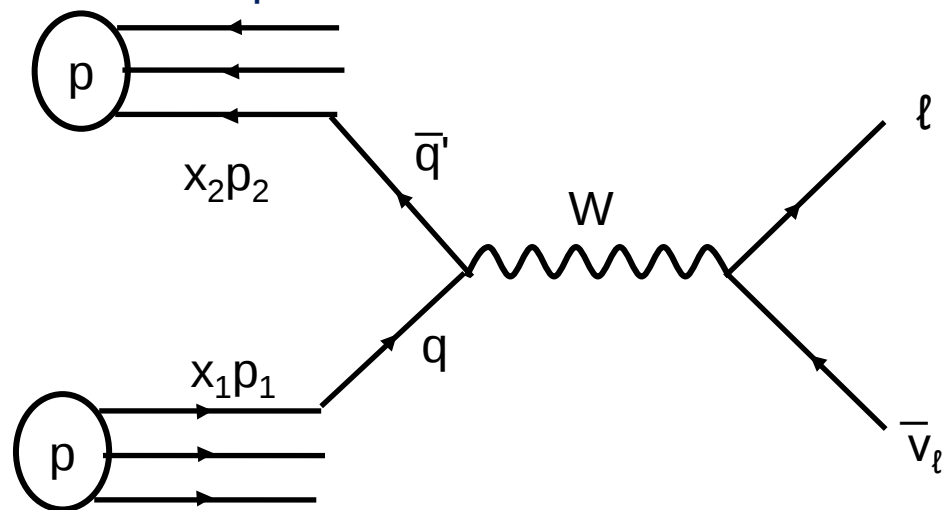
resolution is 19 MeV $\sim$ 2%



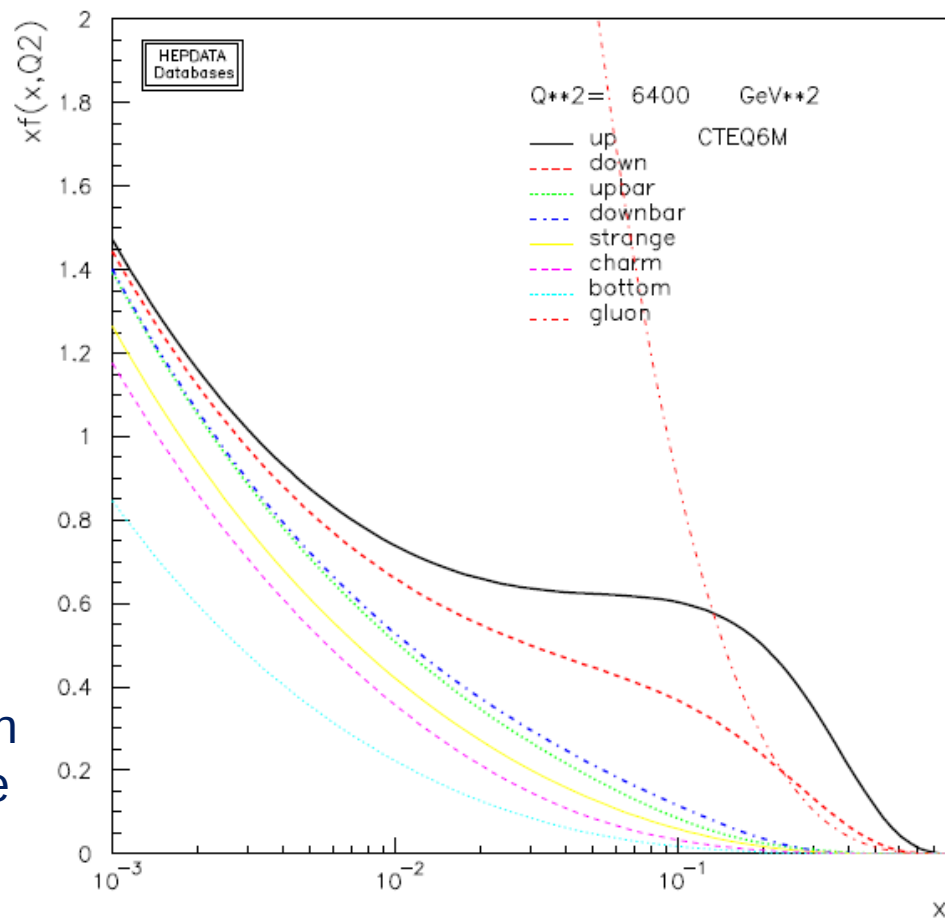


A more detailed example: W production

- W boson production to lowest order:



- quarks carry momentum fractions x_1, x_2
- the probability to find a parton i with a given energy fraction x in a proton is given by the parton distribution function $f_i(x, Q^2)$

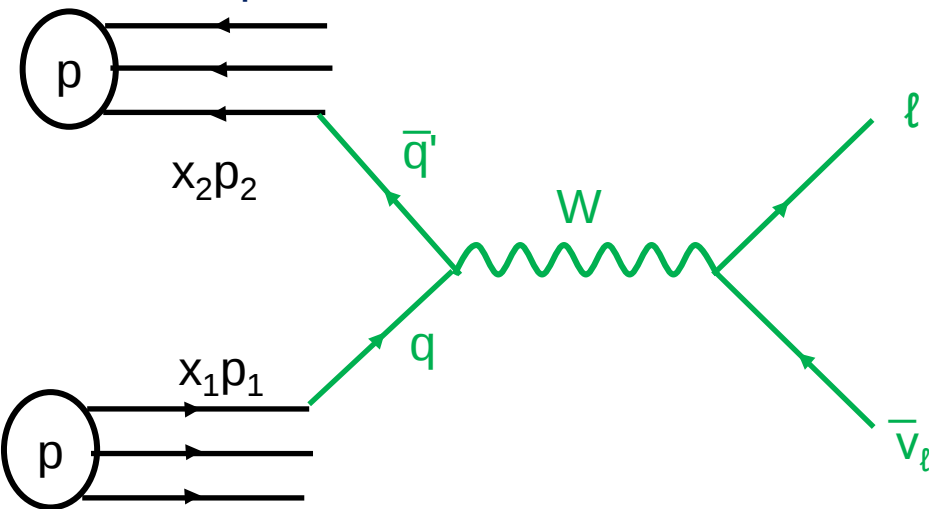


- it is easy to calculate that the mass of the object produced (e.g. a W) is given by $\sqrt{\hat{s}} = \sqrt{x_1 x_2 s}$ with s = centre-of-mass energy squared

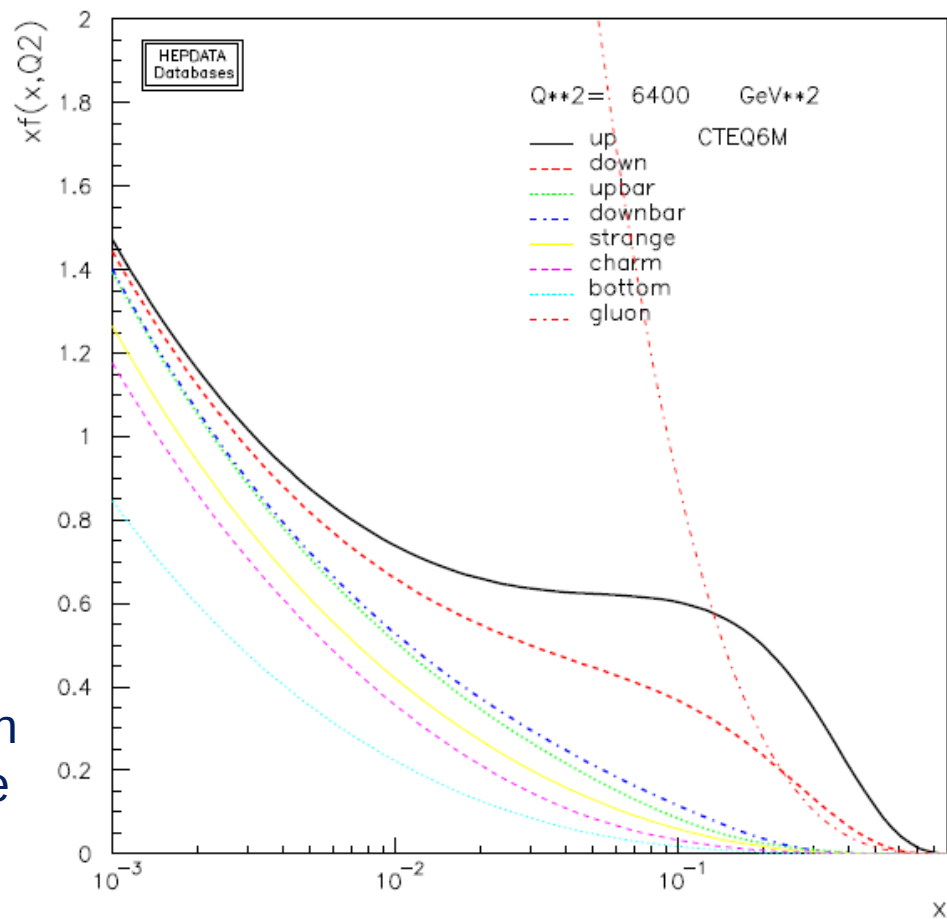


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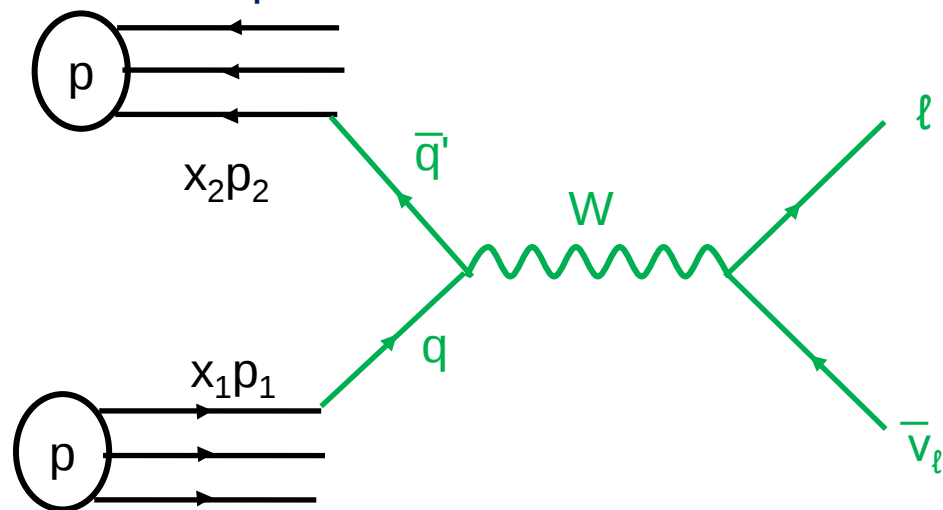
- it is easy to calculate that the mass of the object produced (e.g. a W) is given by $\sqrt{\hat{s}} = \sqrt{x_1 x_2 s}$ with s = centre-of-mass energy squared
- the production cross section can then be written as

$$\sigma_{ij}(\hat{s})$$

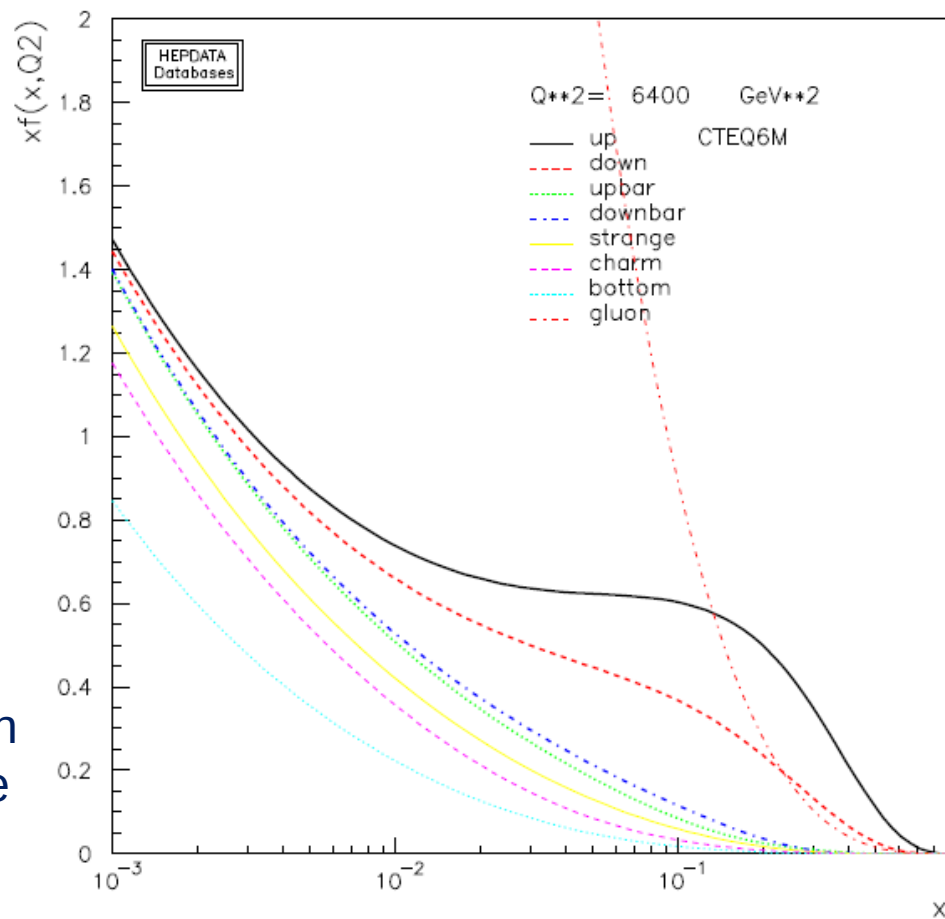


A more detailed example: W production

- W boson production to lowest order:



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- the probability to find a parton i with a given energy fraction x in a proton is given by the **parton distribution function** $f_i(x, Q^2)$



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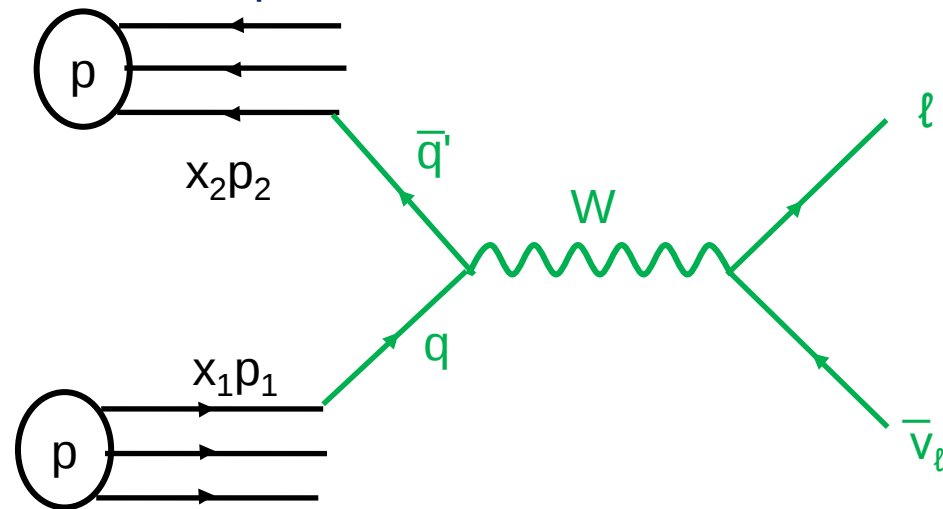
$$\sigma_{ij}(\hat{s})$$

$$\delta(\hat{s} - x_1 x_2 s)$$

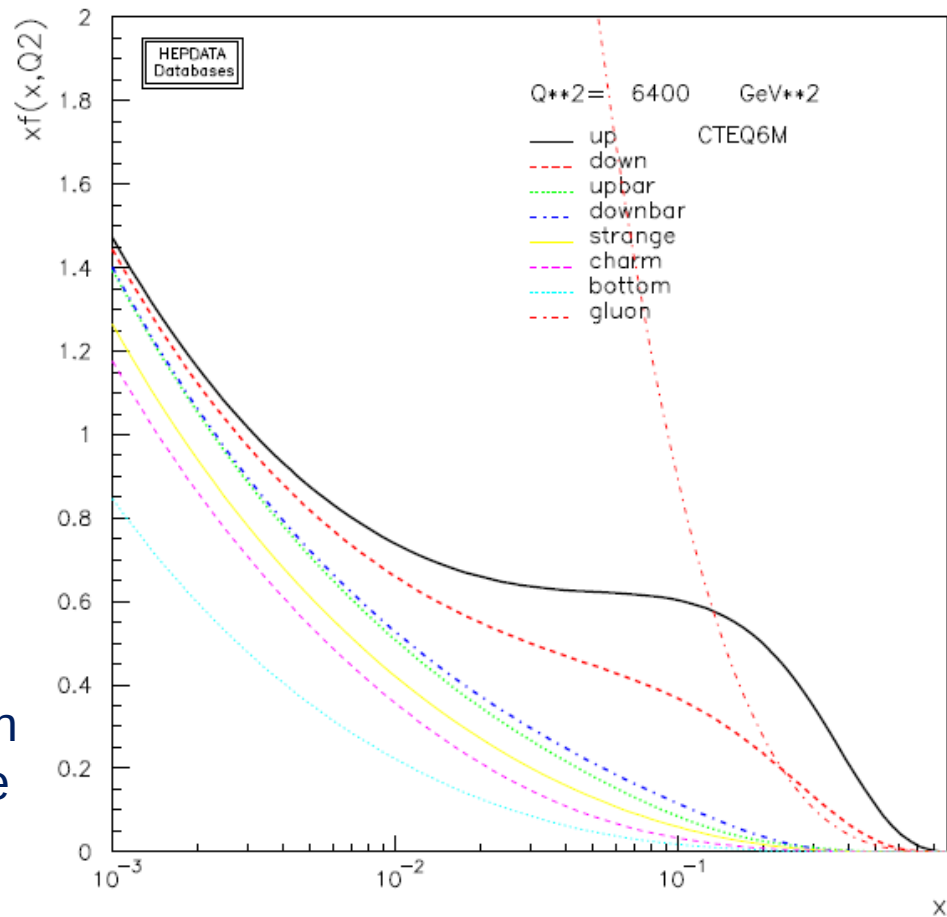


A more detailed example: W production

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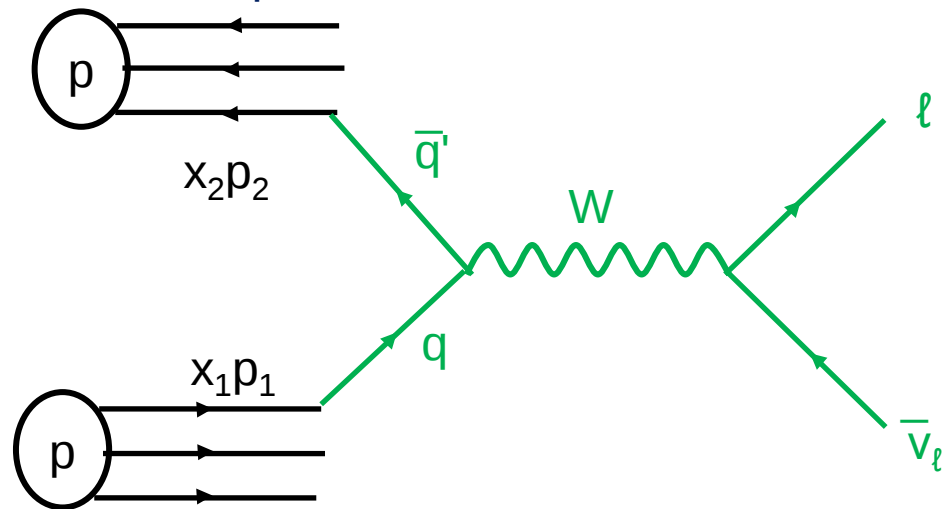
- it is easy to calculate that the mass of the object produced (e.g. a W) is given by $\sqrt{\hat{s}} = \sqrt{x_1 x_2 s}$ with s = centre-of-mass energy squared
- the production cross section can then be written as

$$\sigma_{ij}(\hat{s}) \int_0^1 \int_0^1 f_i(x_1) f_j(x_2) \delta(\hat{s} - x_1 x_2 s) dx_1 dx_2$$

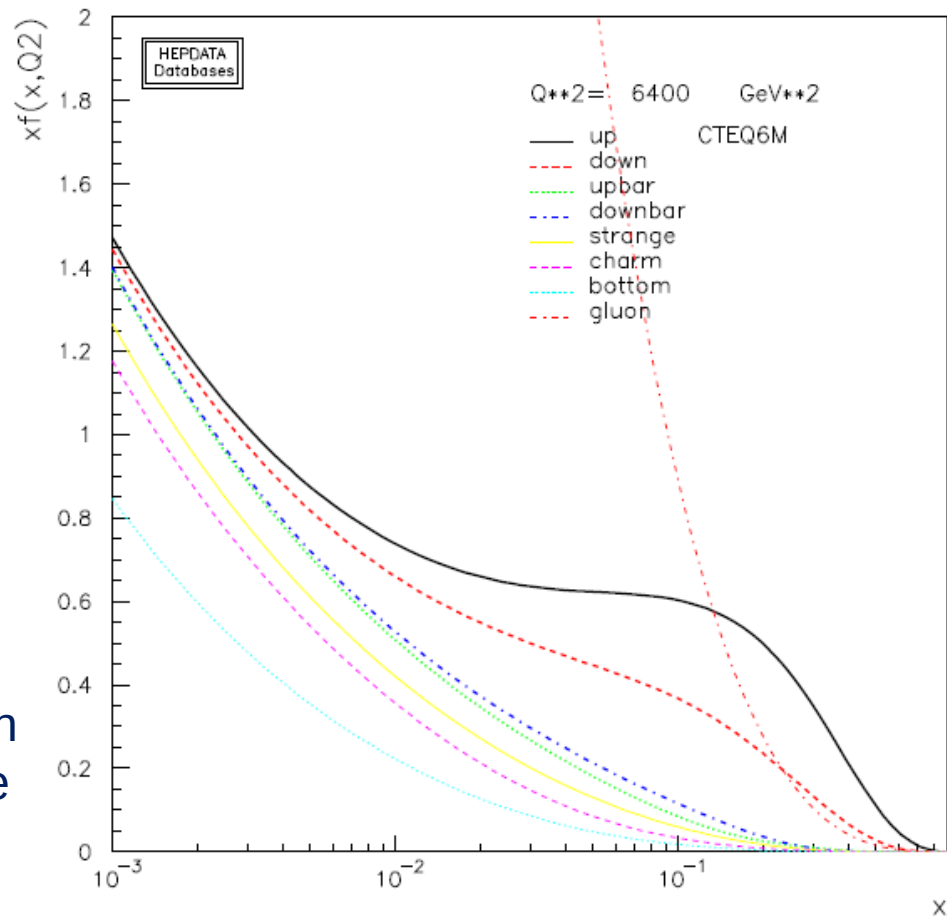


A more detailed example: W production

- W boson production to lowest order:



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- the probability to find a parton i with a given energy fraction x in a proton is given by the parton distribution function $f_i(x, Q^2)$



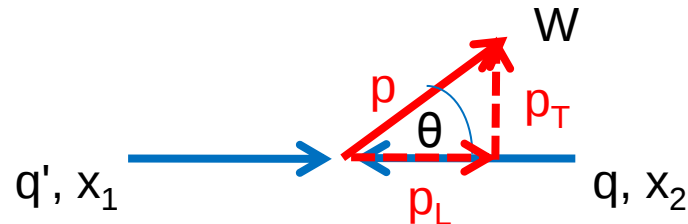
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- the production cross section can then be written as

$$\sigma = \sum_{i,j} \int \sigma_{ij}(\hat{s}) \int_0^1 \int_0^1 f_i(x_1) f_j(x_2) \delta(\hat{s} - x_1 x_2 s) dx_1 dx_2 d\hat{s}$$

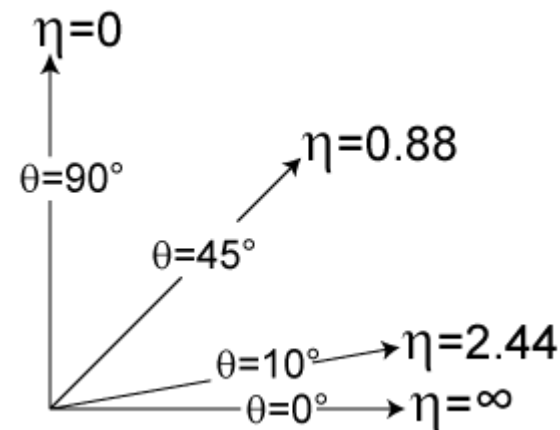


W production kinematics

- At hadron machines the scattering direction is usually given in terms of pseudo-rapidity



$$\eta = \frac{1}{2} \ln \frac{|\vec{p}| + p_L}{|\vec{p}| - p_L} = -\ln \tan \left(\frac{\theta}{2} \right)$$



- in a simplified picture x_1 , x_2 are related to the pseudorapidity of the boson with mass M :

$$x_1 \approx \frac{M}{\sqrt{s}} e^{+\eta} \quad x_2 \approx \frac{M}{\sqrt{s}} e^{-\eta}$$

- examples: $\eta = 0$ then $x_1 = x_2 \approx 0.01$
 $\eta = 0.88$ then $x_1 \approx 0.03$ and $x_2 \approx 0.005$
 $\eta = 3.0$ then $x_1 \approx 0.23$ and $x_2 \approx 0.0006$
for $M=80.4$ GeV and $\sqrt{s}=7$ TeV
- at different scattering directions, different regions of the parton distribution functions are probed



W differential cross-section

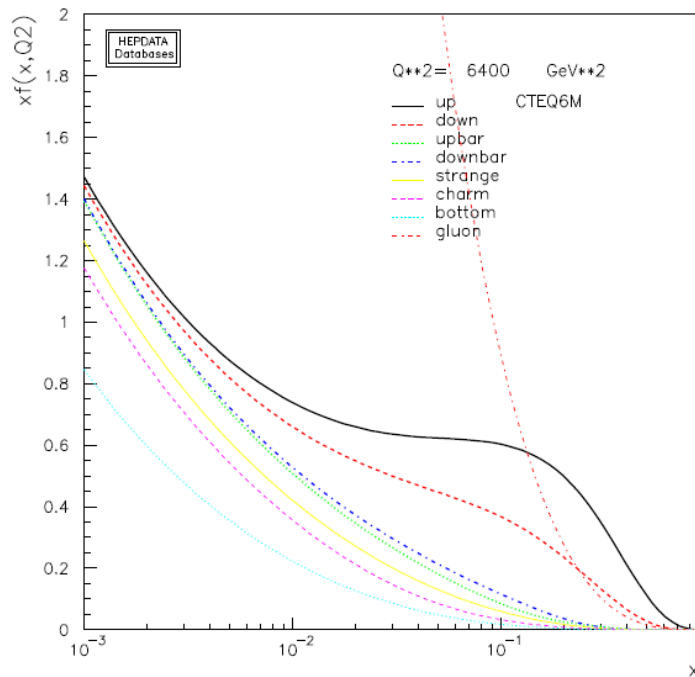
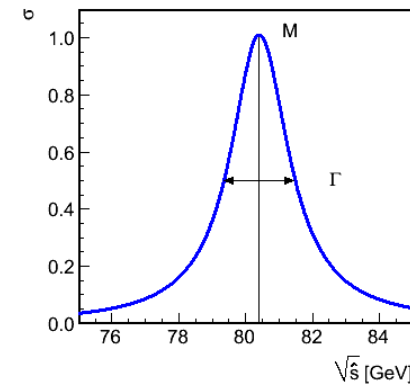
- If you assume that the W boson resonance has zero width ($\Gamma_W \approx 0$)

$$\hat{\sigma}_{ij}(\hat{s}) = \sigma_{ij} \delta(\hat{s} - M^2) M^2 \quad \sigma_{ij} = \sigma(ij \rightarrow W)$$

one can perform the integrals easily and one obtains:

$$\frac{d\sigma}{d\eta}(pp \rightarrow X) \approx \sum_{i,j} \sigma_{ij} \frac{M^2}{s} f_i\left(\frac{M}{\sqrt{s}} e^{+\eta}\right) f_j\left(\frac{M}{\sqrt{s}} e^{-\eta}\right)$$

- consequence: you can calculate the diff. cross-section by simply multiplying the pdf's (and by summing over the possible intermediate states ij)



plots and numerical tables:

<http://durpdg.dur.ac.uk/HEPDATA/PDF>

try it yourself!

(works also for Higgs at LO, ...)



W differential cross-section

$$\frac{d\sigma}{d\eta}(pp \rightarrow X) \approx \sum_{i,j} \sigma_{ij} \frac{M^2}{s} f_i\left(\frac{M}{\sqrt{s}} e^{+\eta}\right) f_j\left(\frac{M}{\sqrt{s}} e^{-\eta}\right)$$

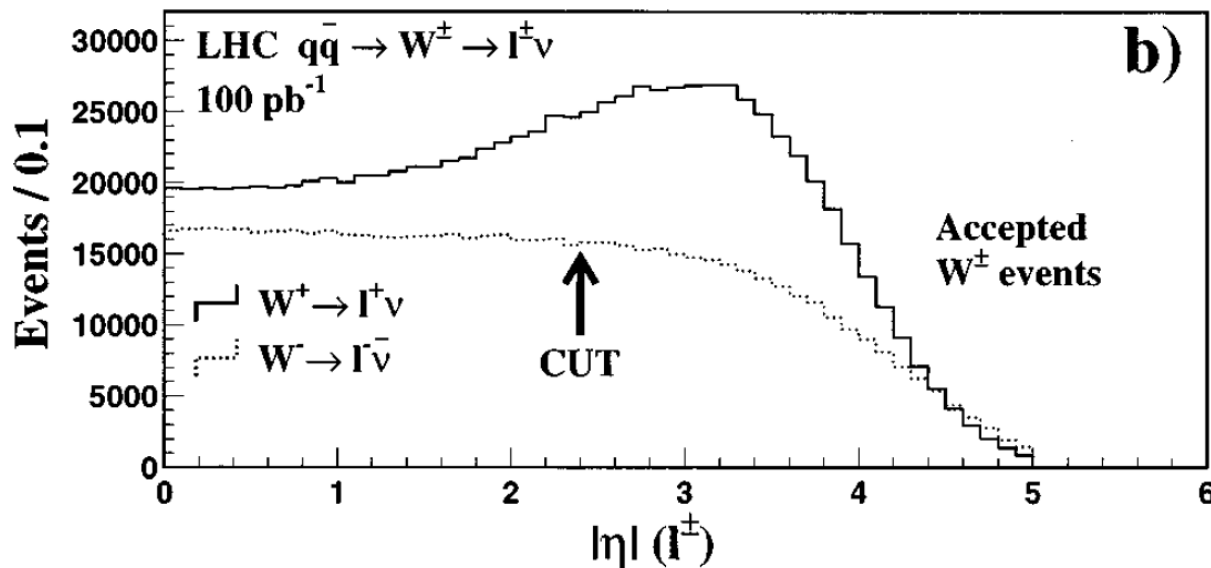
- example:

$$u + \bar{d} \rightarrow W^+$$

$$\bar{u} + d \rightarrow W^-$$

$$\sigma(u + \bar{d} \rightarrow W^+) = \sigma(\bar{u} + d \rightarrow W^-) \equiv \sigma_W$$

- anti-u and anti-d pdf's are about the same (sea-quarks)
- at $\eta = 3.0$ ($x_1 \approx 0.23$, $x_2 \approx 0.0006$)
 - the u- and d-pdf differ by a factor of 2 (valence quarks)
 - W^+ production is enhanced with respect to W^-



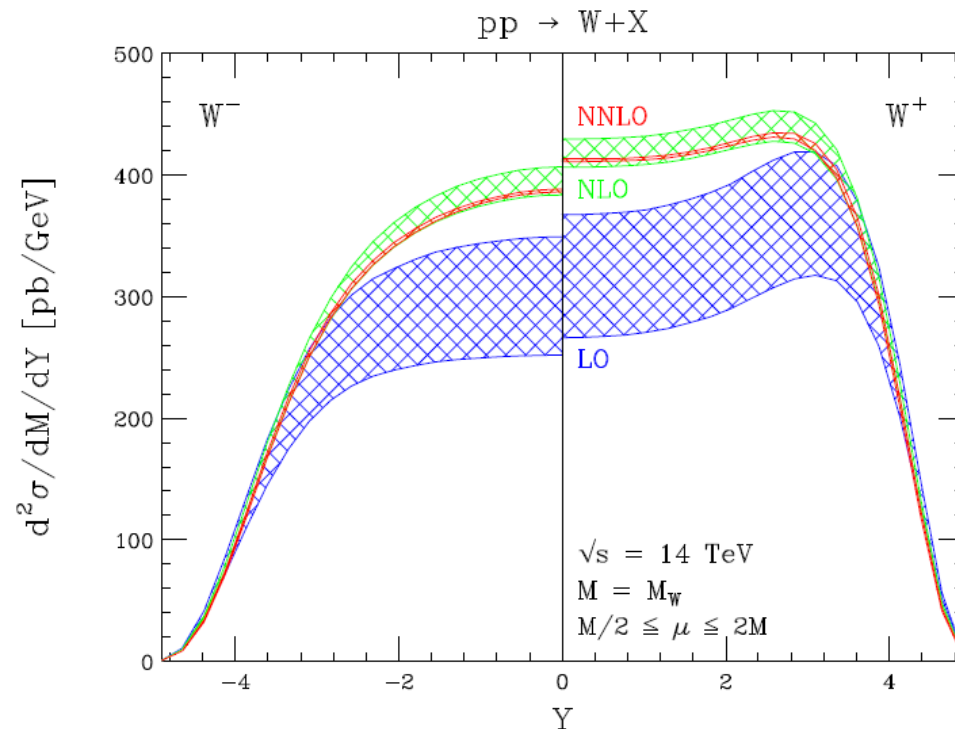
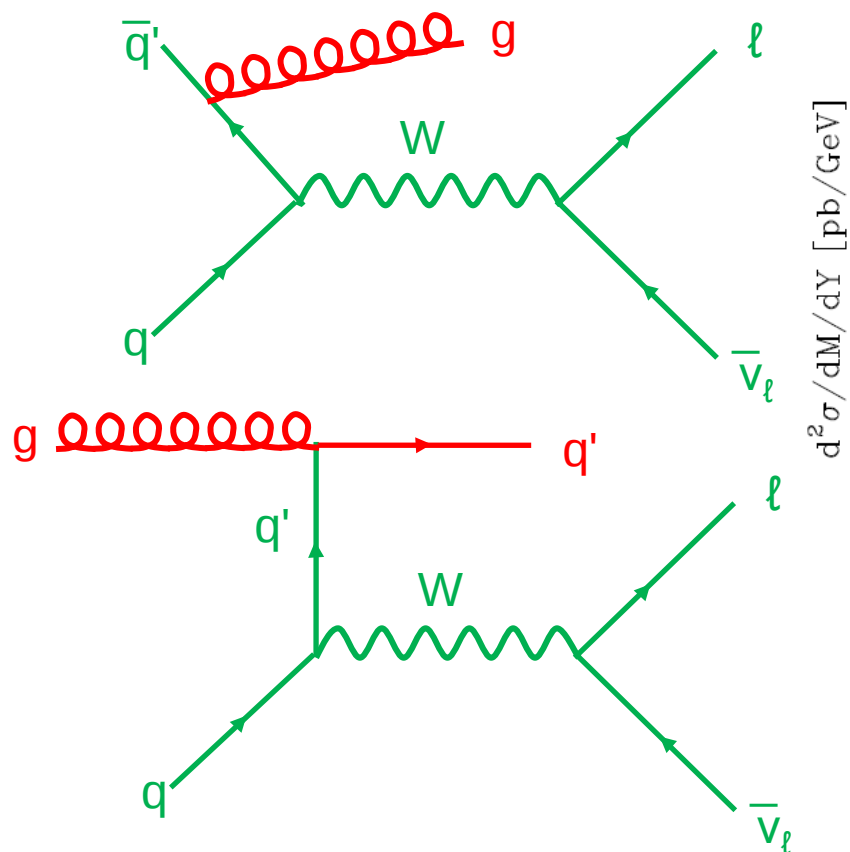
- possibility to test/measure pdf's

M. Dittmar et al., Phys. Rev. D 56 (1997) 7284



W+jet Production

- there are correction to the lowest order description, e.g. due to gluon radiation
- corrections are next-to-leading order (NLO) in α_s or even NNLO, N³LO, ...
→ W+jet production



- theoretical uncertainties are typically reduced when going to higher order calculations



W and Z production cross-sections

- W selection:

- lepton $p_T > 25$ GeV
- $E_T^{\text{miss}} > 25$ GeV $\rightarrow$ neutrino
- W mass would be

$$M_W \approx \sqrt{E^\ell E^\nu (1 - \cos \alpha(\ell, \nu))}$$

but E^ν and angle can not be measured
 $\rightarrow$ transverse mass: $M_T > 40$ GeV

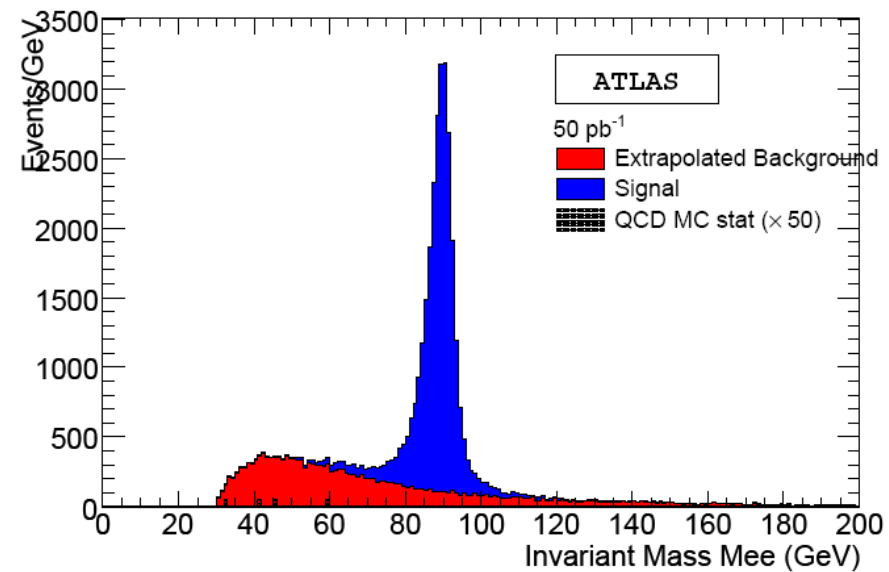
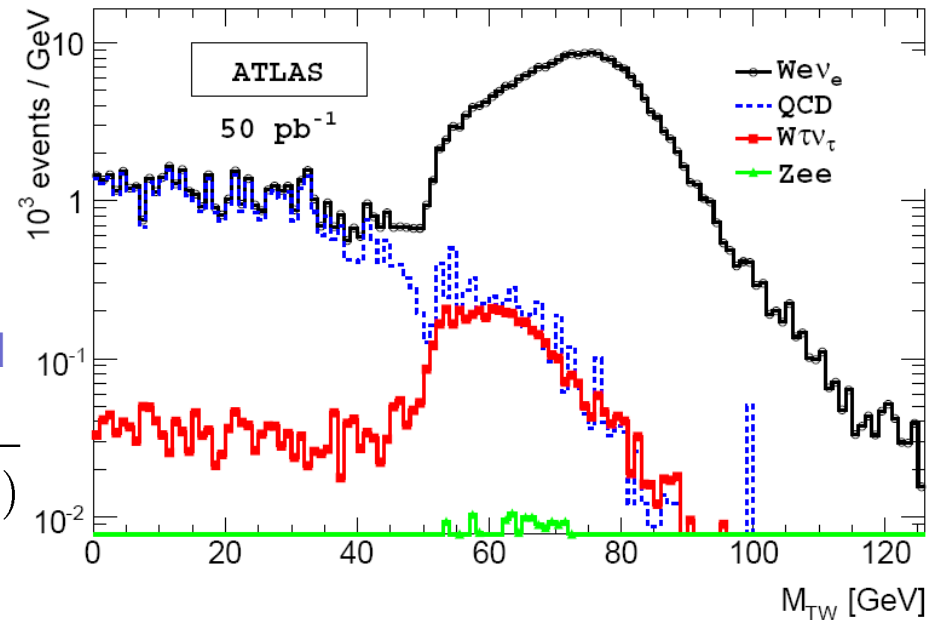
$$M_T = \sqrt{p_T^\ell p_T^{\text{miss}} (1 - \cos \Delta\phi(p_T^\ell, p_T^{\text{miss}}))}$$

- Z selection:

- lepton $p_T > 15$ GeV
- $80 < M_{\ell\ell} < 100$ GeV
- isolated electrons and muons

- background at 10% level (QCD)

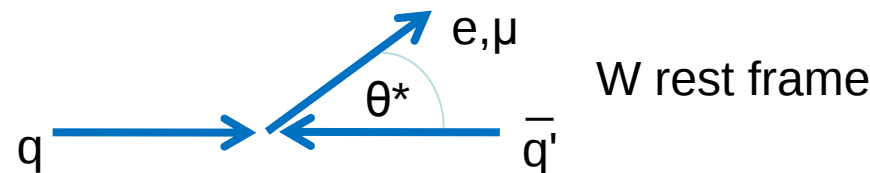
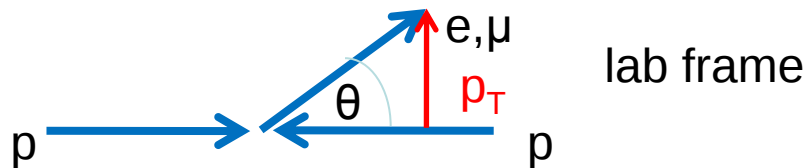
- stat. uncertainty $\sim 1/\sqrt{N}$ will be very small
- systematics will dominate:
background uncertainties, luminosity and
theory (PDF, gluon radiation, underlying
event,...)





Measurement of M_W

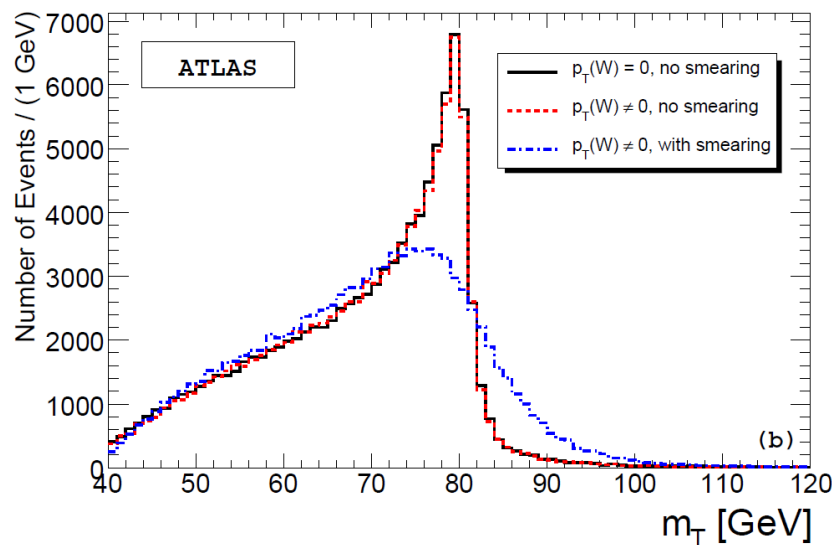
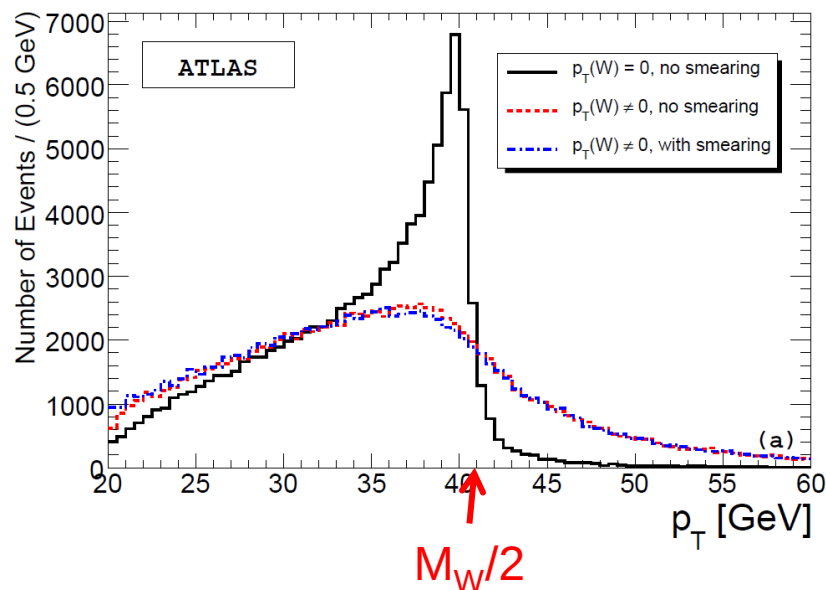
- kinematics of the $W \rightarrow \ell \nu_\ell$ decay $\rightarrow$ only the lepton is measured in the detector:



- differential decay rate in the lab frame can be traced back to decay rate in rest frame:

$$\frac{dN}{dp_T} = \frac{2p_T}{M_W} \frac{1}{\sqrt{\left(\frac{M_W}{2}\right)^2 - p_T^2}} \frac{dN}{d\cos\theta^*}$$

- Jacobian factor is singular at $p_T = M_W/2 \rightarrow$ Mass peak in p_T spectrum



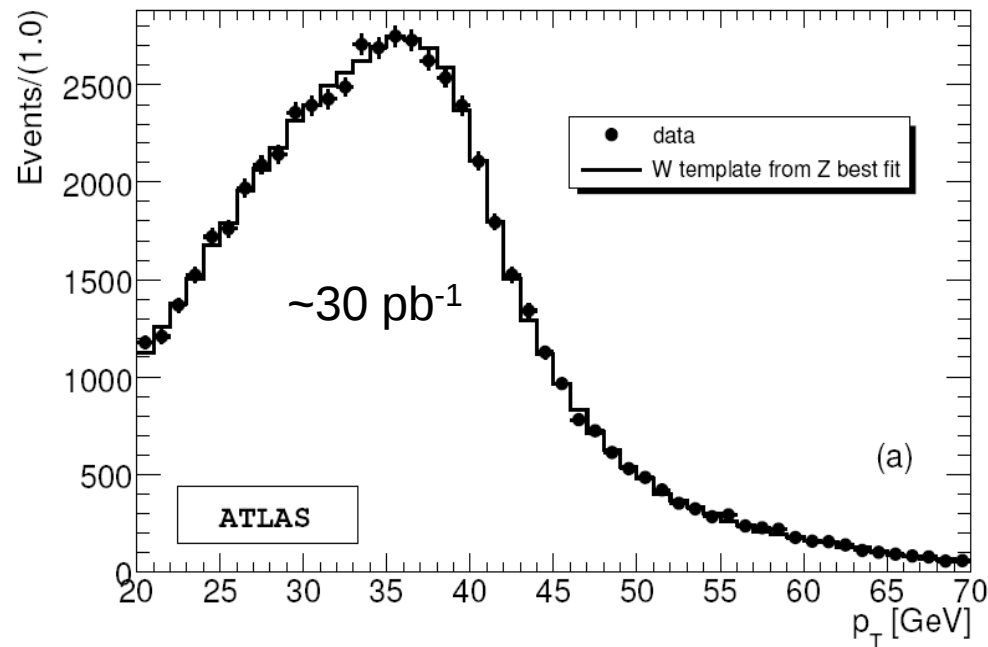
- p_T of W is non-zero $\rightarrow$ changes spectrum

$$M_T = \sqrt{p_T^\ell p_T^{miss} (1 - \cos \Delta\phi(p_T^\ell, p_T^{miss}))}$$



Measurement of M_W

- fit spectra of lepton p_T and $M_T = \sqrt{p_T^\ell p_T^{miss}(1 - \cos \Delta\phi(p_T^\ell, p_T^{miss}))}$
- lepton energy scale from Z events
- adjust Monte Carlo samples with different M_W to fit best to data spectrum
- most sensitive region is upper edge of the Jacobian peak (~ 40 GeV)

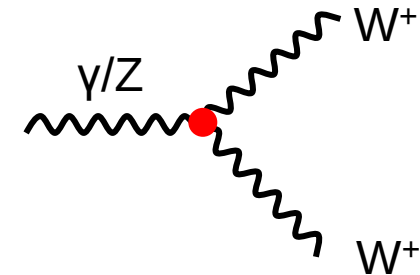


- statistical uncertainty in $\sim 30 \text{ pb}^{-1}$ is ~ 60 MeV in $W \rightarrow \mu\nu$ and $W \rightarrow e\nu$ channel
- systematics: lepton energy scale 110 MeV and 200 MeV due to W recoil, p_T^W
- "warm-up" for the ultimate 15 MeV precision $\rightarrow 10\text{-}20 \text{ fb}^{-1}$
- W mass precision today: 23 MeV (LEP+Tevatron)



Couplings between Gauge Bosons

- non-abelian nature of SM predicts
→ triple gauge boson couplings (TGC)



- coupling strength in the SM:

$$g_{WW\gamma} = e \quad g_{WWZ} = e \cot \theta_w$$

- measure these couplings and test for anomalous contributions
- effective Lagrangian is used to model these (V= γ or Z):

$$L/g_{WWV} = i \boxed{g_1^V} (W_{\mu\nu}^* W^\mu V^\nu - W_{\mu\nu} W^{*\mu} V^\nu) + i \boxed{\kappa^V} W_\mu^* W_\nu V^{\mu\nu} + \frac{\boxed{\lambda^V}}{M_W^2} W_{\rho\mu}^* W_\nu^\mu V^{\nu\rho}$$

- in the SM: $g_1^\gamma = g_1^Z = 1$ $\kappa^\gamma = \kappa^Z = 1$ $\lambda^\gamma = \lambda^Z = 0$
- at LHC the di-boson decays to electrons and muons will be measured:

$$\left. \begin{aligned} pp &\rightarrow ZW \rightarrow \ell\ell \ell\nu \\ pp &\rightarrow WW \rightarrow \ell\nu \ell\nu \\ pp &\rightarrow ZZ \rightarrow \ell\ell \ell\ell, \ell\ell \nu\nu \\ pp &\rightarrow Z\gamma \rightarrow \ell\ell \gamma \\ pp &\rightarrow W\gamma \rightarrow \ell\nu \gamma \end{aligned} \right\}$$

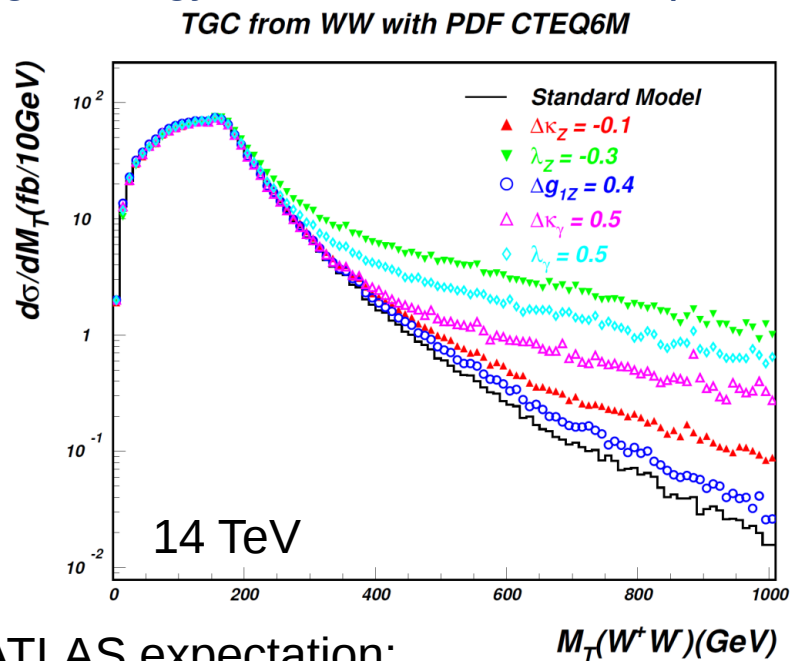
high p_T leptons and photons
Z mass and transverse mass constraints

- anomalous couplings between ZZZ and ZZγ are studied, too → do not exist in SM

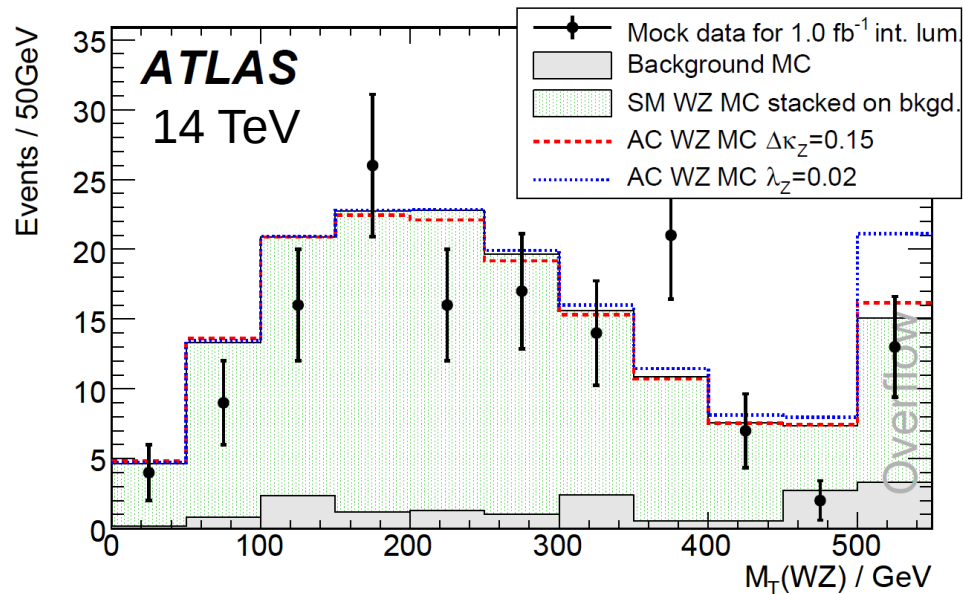
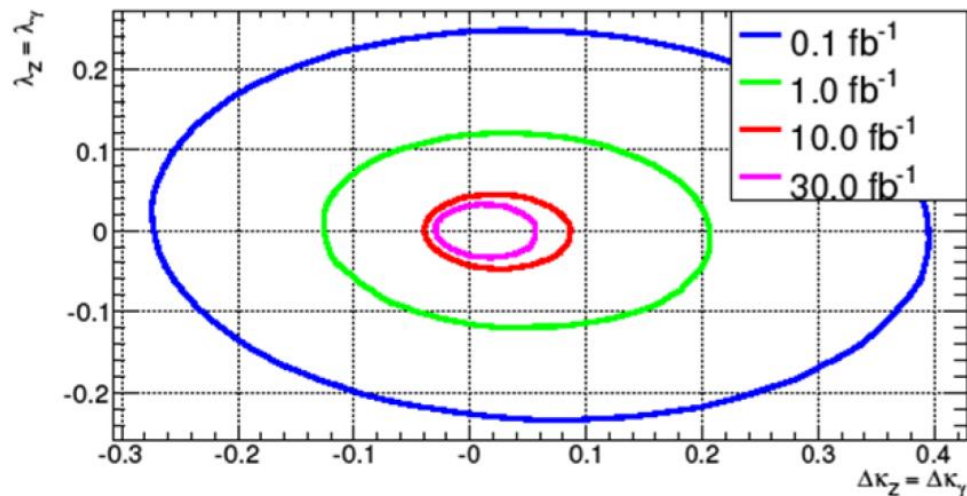


Gauge boson self-couplings

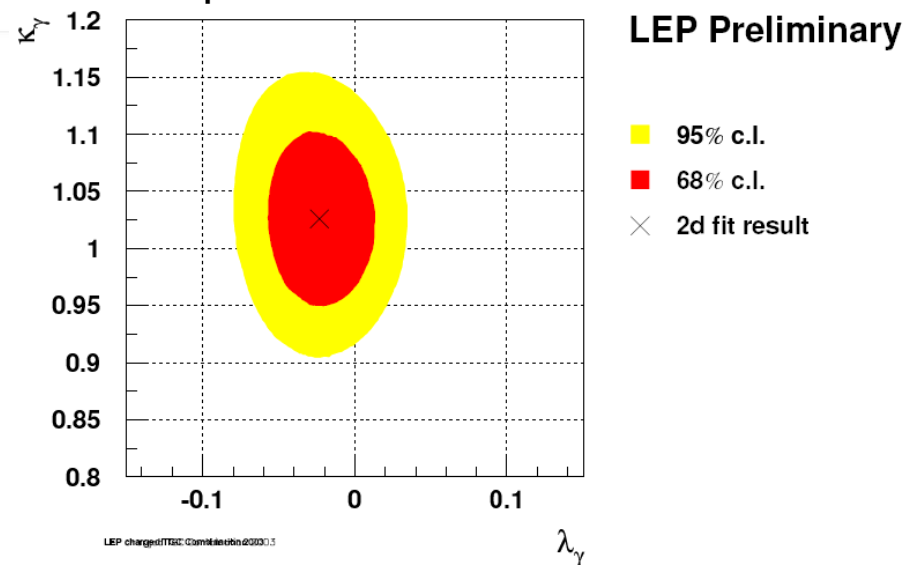
- high energy reach of LHC will help to improve the coupling measurements:



ATLAS expectation:



compared to LEP:





Weak Mixing Angle at the LHC

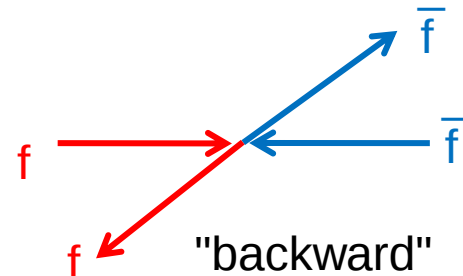
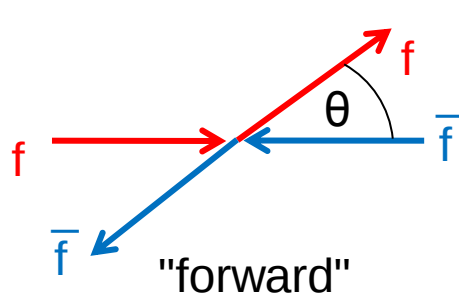
- Z boson production at LHC is the "inverse" of the LEP process

$$\text{LHC: } q\bar{q} \rightarrow Z/\gamma^* \rightarrow e^+e^-$$

$$\text{LEP: } e^+e^- \rightarrow Z/\gamma^* \rightarrow q\bar{q}$$

- in the Z rest frame:
measurement of forward-backward asymmetry

$$A_{FB} = \frac{N_F - N_B}{N_F + N_B}$$



- A_{FB} depends on Z-to-fermion couplings

$$A_{FB} = A_{FB}(c_V^f, c_A^f) = A_{FB}(\sin^2 \theta_w)$$

→ precise measurement of $\sin^2 \theta_w$

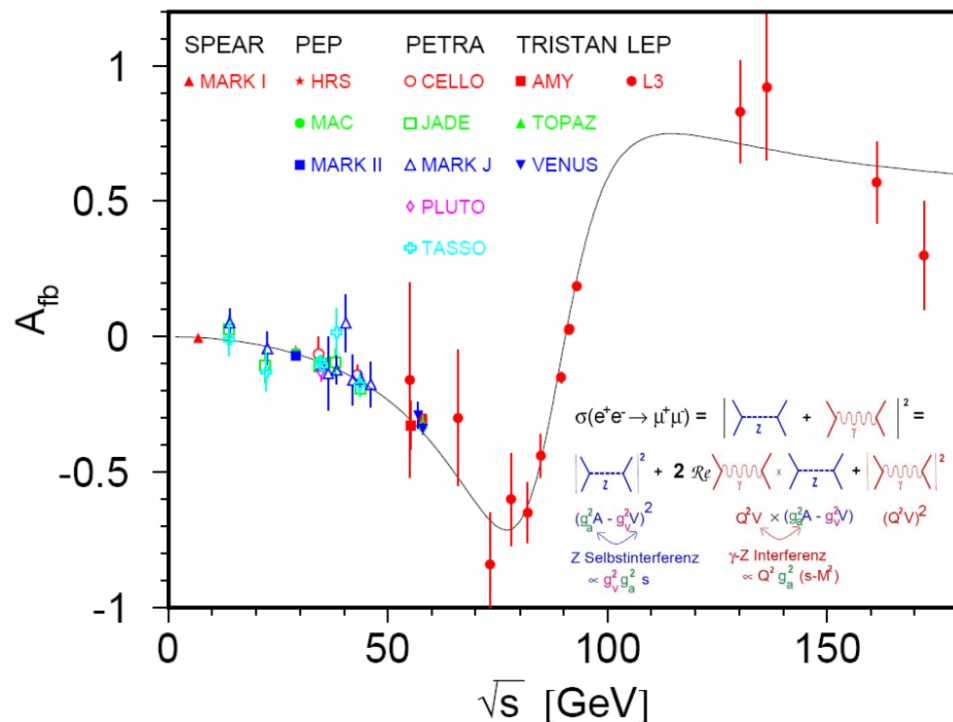
$$\delta \sin^2 \theta_w = 1.6 \cdot 10^{-4}$$

$$\delta M_W / M_W = 2.8 \cdot 10^{-4}$$

$$\delta M_Z / M_Z = 0.2 \cdot 10^{-4}$$

→ test the Standard Model

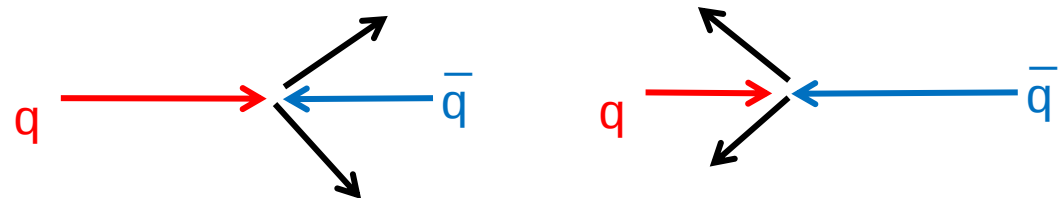
$$\rho = \frac{M_W}{M_Z \cos \theta_w} = 1 + \Delta \rho(\text{rad. corr.})$$



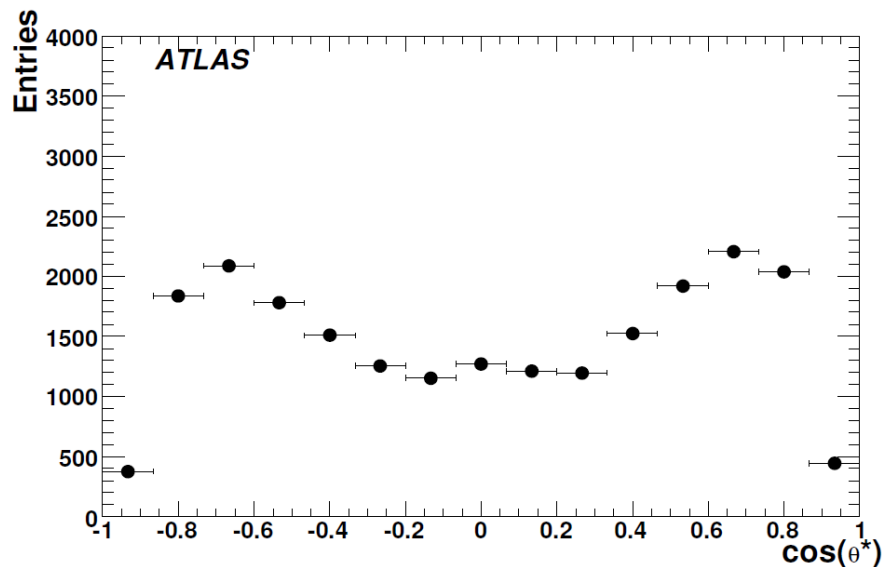
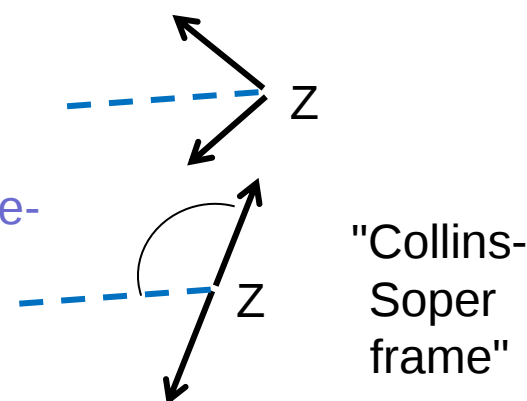
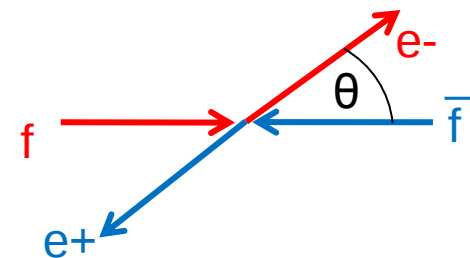


Weak Mixing Angle Measurement

- LEP was a e^+e^- collider with symmetric beam energies
 - automatically in Z rest frame
- quark-quark system at LHC has varying longitudinal momentum $\rightarrow$ depends on share between x_1 and x_2



- a trick has to be done: we know the Z decays to e^+e^-
 - $\rightarrow$ the reference direction is the bi-section between e^+ and e^-
 - $\rightarrow$ go into the e^+e^- rest frame = Z boson rest frame



- good enough approximation to measure the asymmetry

$$A_{\text{FB}} = b(a - \sin^2 \theta_{\text{eff}}^{\ell})$$

- LHC expects a precision of $\delta \sin^2 \theta_w = 3 \cdot 10^{-4}$
- 100 fb^{-1} and includes forward detectors
- compare to LEP: $\delta \sin^2 \theta_w = 1.6 \cdot 10^{-4}$



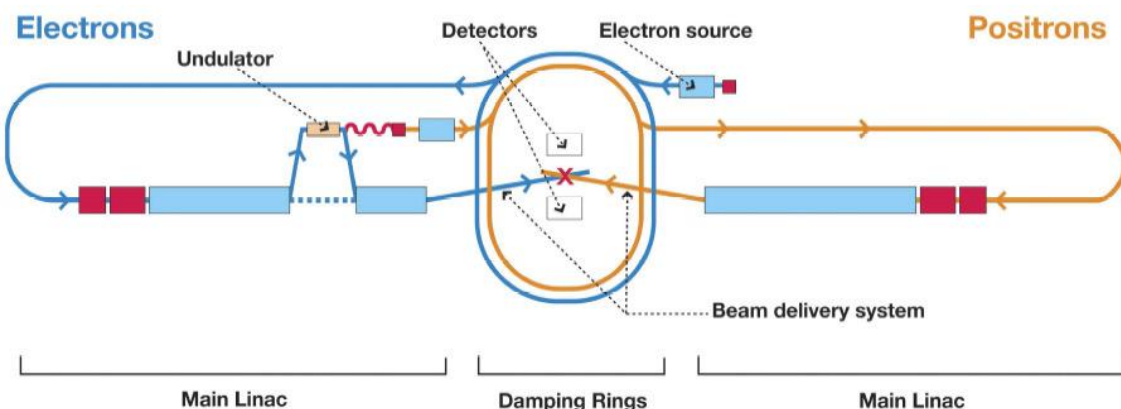
Summary of LHC Electroweak Measurements

- production cross-sections of W and Z boson
- W mass with 15 MeV ultimate precision
- di-boson production WW, WZ, ZZ, ...
- triple gauge boson couplings to precision beyond $\Delta\kappa/\Delta\lambda/\Delta g_1=0.01$
- perform a new measurement of the weak mixing angle
- $\delta\sin^2\theta_w = 3\cdot 10^{-4}$ and maybe better
- need a very good understanding of the ATLAS and CMS detectors
 - known electro-weak processes are a main reference
- need a well-performing LHC machine to deliver good data
- we know that the Standard Model is just an effective theory which will need extensions at very high energies
 - keep your eyes open when looking at invariant mass spectra, etc.



A Quick Look into the Future: the ILC

- International Linear Collider



e+e- machine
500 GeV centre-of-mass energy
luminosity $2 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$
polarization e- 80% e+ 30% (60%)
 $L = 500 \text{ fb}^{-1}$ in 4 years

- detectors better than at LEP: jet measurement, e.m. calorimetry, tracking
- mainly for
 - studying SUSY physics
 - measuring Higgs boson properties
- and...
 - very precise **top mass** determination → top-pair threshold scan $\Delta m_t \approx 100 \text{ MeV}$
 - further improved **TGC measurements**
 - more precise measurement of $\sin^2 \theta_w \rightarrow$ **GigaZ option at Z peak energies**



Backup for Further Questions



ATLAS Detectors at Start-up

Subdetector	Number of Channels	Approximate Operational Fraction
Pixels	80 M	97.9%
SCT Silicon Strips	6.3 M	99.3%
TRT Transition Radiation Tracker	350 k	98.2%
LAr EM Calorimeter	170 k	98.8%
Tile calorimeter	9800	99.2%
Hadronic endcap LAr calorimeter	5600	99.9%
Forward LAr calorimeter	3500	100%
MDT Muon Drift Tubes	350 k	99.7%
CSC Cathode Strip Chambers	31 k	98.4%
RPC Barrel Muon Trigger	370 k	98.5%
TGC Endcap Muon Trigger	320 k	99.4%
LVL1 Calo trigger	7160	99.8%

Presented by J. Kanzaki at KEK in Feb. 19, 2010

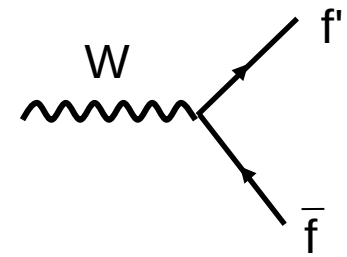
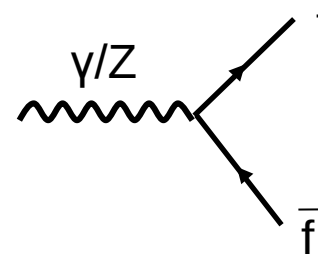


Fermion-Boson Couplings

- couplings of fermions to gauge bosons

$$-ig\frac{1}{2}(c_V\gamma^\mu - c_A\gamma^\mu\gamma^5)$$

- coupling to Z-bosons involves θ_w :



	γ	W	Z
c_V	$2Q$	$T_3 = \pm\frac{1}{2}, 0$	$T_3 - 2Q \sin^2 \theta_w$
c_A	0	T_3	T_3

	γ	W	Z
g	e	$\sqrt{2} \frac{e}{\sin \theta_w}$	$g_z = \frac{e}{\sin \theta_w \cos \theta_w}$



Cross-sections at LHC

LO cross sections in pb, inclusive

E_{cm} [TeV]/ Process	7	10	14	Evts (7 TeV) in 200/pb	Ratio 7/14
QCD $p_t > 100$ GeV	3.2E+05	6.8E+05	1.4E+06	6.4E+07	0.2
Z incl	2.5E+04	3.6E+04	5.7E+04	5.0E+06	0.4
W incl	9.5E+04	1.4E+05	2.1E+05	1.9E+07	0.5
ttbar	8.4E+01	2.2E+02	4.8E+02	1.7E+04	0.2
$H_{(150 \text{ GeV})}$	4.0	8.2	16.0	8.0E+02	0.3

no branching ratios included !

G. Dissertori



Particle Identification

- combined performance is important

	Tracker	ECAL	HCAL	Muon-System	Vertex-Detector
Electrons	++	+++	+	-	-
Photons	+	+++	+	-	-
Muons	+(+)	+	+	+++	(+)
Taus	++	++	++	++	++
B-Jets	++	+	+	+	+++
Jets	+	+++	+++	(+)	-
Missing E_T	+	+++	+++	+	-

K. Desch