Two remarkable features of Composite Dark Matter

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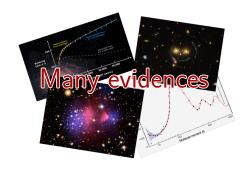




DESY seminar, 10th February 2020

Based on 1905.13244 with A. Davoli, A. De Simone, D. Marzocca

Dark matter problem



But what is it? Not yet known

Many possible candidates →

→ Weakly Interacting Massive Particle (WIMP)

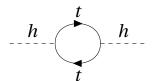
Naturalness/Hierarchy

Naturalness principle

A guide on the size of couplings: $[g] = \alpha \rightarrow g \approx \Lambda^{\alpha}$

$$[\lambda] = 0 \rightarrow |\lambda| \approx 1$$
, $[m_h^2] = 2 \rightarrow m_h^2 \approx m_{Pl}^2$

Hierarchy problem



h Loop gives contribution δm_h^2 : $\delta m_h^2 \propto \Lambda_{NP}^2 \gg m_h^2$

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Open issues

Dark Matter

 $\Omega_{\rm DM} h^2 \approx 0.1198$



Naturalness

 $m_h \ll m_{\rm Pl}$



Composite Higgs

New particles \rightarrow DM candidate? New scale $f: m_h \leq f \ll m_{\rm Pl}$

Two remarkable features of Composite Dark Matter

CH paradigm (see e.g. 1105.5403)

There is a strong sector at $m_* = g_* f$

$$\mathscr{G} \xrightarrow{\mathsf{condensate}} \mathscr{H} \longrightarrow \mathsf{Particles} \; \mathsf{generated}$$

Explicit:

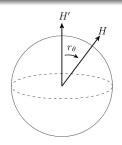
$$\mathscr{Q} \xrightarrow{\mathsf{Yukawa}, \mathsf{gauge}} \mathscr{H} \Longrightarrow \mathsf{Masses} \mathsf{generated}$$

Where is the SM?

$$SM \subset SO(4)_c \subset \mathcal{H}' \neq \mathcal{H}$$

The two groups are misaligned!

CH paradigm (see e.g. 1105.5403)



Two different vacua ↓
Two different field expansions

Related by r_{θ} : $\sin \theta = v/f$

Gauge basis T related to \mathcal{H}' : $\langle h \rangle = v$

Physical basis T_{θ} related to \mathcal{H} : $\langle h \rangle = 0$

$$T_{\theta} = r_{\theta} T r_{\theta}^{-1}$$

Generators

$$\mathscr{G} = SO(7), \ \mathscr{H} = SO(5) \times SO(2)$$

$$T_{L,R} \sim \text{SO}(4)_c, \ T_5 \sim \frac{\text{SO}(5)}{\text{SO}(4)_c}, \ T_2 \sim \text{SO}(2), \ \hat{T} \sim \frac{\text{SO}(7)}{\text{SO}(5) \times \text{SO}(2)}$$

$$(T, \hat{T}) = \left(egin{array}{c|c} T_{L,R} & T_5 & \hat{T}_1 & \hat{T}_2 \ \hline T_5 & \hat{T}_1 & \hat{T}_2 \ \hline \hat{T}_1 & \hat{T}_2 & T_2 \ \hline \hat{T}_1 & \hat{T}_2 & T_2 \end{array}
ight).$$

 $SO(7) \rightarrow SO(5) \times SO(2) \implies 10$ NGBs generated in a (5,2)

Our coset and pNGBs

	Field	SO(4)'	C_2	P_7
2HDM {	$\overline{\phi_1}$	4	+	+
	ϕ_2	4	_	+
singlets $\left\{ \right.$	η	1	+	_
	κ	1	_	_

Symmetries

 C_2 : no sizable vacuum for $\phi_2 \rightarrow$ safe custodial breaking

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 P_7 : no $\eta \bar{q} q$ coupling \rightarrow stability of DM

Particle identification

$$\phi_1 = \begin{pmatrix} G_1 \\ G_2 \\ G_3 \\ h \end{pmatrix}, \qquad \phi_2 = \begin{pmatrix} \dfrac{-i}{\sqrt{2}}(H_+ - H_-) \\ \dfrac{1}{\sqrt{2}}(H_+ + H_-) \\ H_0 \\ A_0 \end{pmatrix}$$

Why that H_0 and A_0 identification? (spoilers ahead)

$$\mathcal{L} \supset -\frac{m_q}{v} \bar{q} q(k_q h + k_{H_0 q} H_0) - i \frac{m_q}{v^2} \bar{q} \gamma^5 q(g_{\eta \kappa q} \eta \kappa + \tilde{g} h A_0)$$

CCWZ construction

First part of the model: CCWZ Lagrangian

$$\vec{\Phi} = U(\Pi) \vec{F}$$
, where $U \rightarrow g \cdot U \cdot h^{-1}$

We define a symbol: $d_{\mu} \equiv i \sum_{I} \text{Tr}[U^{-1} D_{\mu} U \hat{T}_{I}] \hat{T}_{I}$

 D_{μ} contains the SM gauge bosons!

Simplest invariant object:
$$\mathscr{L}_{\Pi}^{(2)} = \frac{f^2}{4} d_{\mu} d^{\mu}$$

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CCWZ construction

We work in the physical basis!

$$U_{\theta} = \exp\left(i\frac{\sqrt{2}}{f}\Pi_{I}\hat{T}_{I}^{\theta}\right)$$

 d_{μ} contains SM particles $\left(W_{\mu},Z_{\mu},h\right)$ and new particles

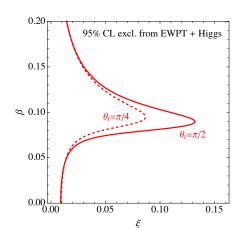
- $\Rightarrow m_Z, m_W$ generated
- ⇒ SM Higgs interactions modified
- ⇒ totally new interactions with the new pNGBs

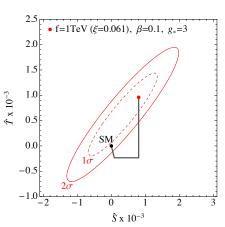
VEV structure

$$\begin{split} \sin\theta_1 &= \sqrt{\xi}\cos\beta, \quad \sin\theta_2 = \sqrt{\xi}\sin\beta \\ & \qquad \qquad \downarrow \\ m_W^2 &= \frac{g^2v^2}{4}, \quad m_Z^2 = \frac{v^2(g^2+g'^2)}{4} \left[1 - \xi\left(1 - \cos(4\beta)\right)\right] \\ & \qquad \qquad \downarrow \\ (\Delta\hat{T})_{\rm 2HDM} &= \frac{\xi}{4}(1 - \cos4\beta) \approx 2\xi\beta^2 \end{split}$$

What is the limit from EWPT?

EWPT





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Partial compositeness

Linear coupling to strong sector:

$$\boldsymbol{\mathscr{L}}_{\mathrm{int}}^f = \bar{q}_L^\alpha \boldsymbol{\mathscr{Y}}_L^{\alpha T} \boldsymbol{\mathscr{O}}_L + \bar{t}_R \boldsymbol{\mathscr{Y}}_R^T \boldsymbol{\mathscr{O}}_R + \mathrm{h.c.}$$

We want \mathcal{H} invariants:

$$\bar{\mathscr{Y}}_L^{\alpha} \equiv \left(r_{\theta}^{-1} U_{\theta}^{\dagger} \mathscr{Y}_L\right)^{\alpha}, \quad \bar{\mathscr{Y}}_R \equiv r_{\theta}^{-1} U_{\theta}^{\dagger} \mathscr{Y}_R$$

Now we can give mass:

$$\mathcal{L}_{t} = c_{t} \frac{m_{*}}{g_{*}^{2}} \bar{q}_{L}^{\alpha} \left(\bar{\mathcal{Y}}_{L}^{\alpha}\right)^{I} \left(\bar{\mathcal{Y}}_{R}\right)_{I}^{\dagger} t_{R}$$

Partial compositeness

Next choice, fermion embedding:

$$\mathbf{7}_{\frac{2}{3}} = \mathbf{2}_{\frac{7}{6}} \oplus \mathbf{2}_{\frac{1}{6}} \oplus \mathbf{1}_{\frac{2}{3}} \oplus \mathbf{1}_{\frac{2}{3}} \oplus \mathbf{1}_{\frac{2}{3}}$$

Do you remember P_7 ? It forbids a 5th component

$$\mathcal{Y}_{L} = \frac{y_{L}}{\sqrt{2}} \begin{pmatrix} 0 & 0 & i & 1 & 0 & 0 & 0 \\ i & -1 & 0 & 0 & 0 & 0 & 0 \end{pmatrix}$$

$$\mathcal{Y}_R = y_R \begin{pmatrix} 0 & 0 & 0 & 0 & \cos \theta_t & i \sin \theta_t \end{pmatrix}$$

$$Y_t \approx \frac{c_t}{g_*} \frac{y_L y_R}{g_*} (\sqrt{1 - \xi} \cos \theta_t + \beta \sin \theta_t)$$

Potential construction

Without getting into the details of the invariant construction...

Fermions

Chirality
$$\rightarrow$$
 New objects $\bar{\Delta}_L = \bar{\mathscr{Y}}_L^* \bar{\mathscr{Y}}_L$, $\bar{\Delta}_R = \bar{\mathscr{Y}}_R^* \bar{\mathscr{Y}}_R$

$$\mathscr{I}_1 = \bar{\Delta}_L^{ii}$$
, $\mathscr{I}_3 = \bar{\Delta}_L^{ij} \bar{\Delta}_L^{ji}$ and so on

Gauge

$$\mathcal{G}^\alpha = g \, T_L^\alpha, \, \mathcal{G}'^\alpha = g' \, T_R^3 \quad \to \quad \bar{\mathcal{G}} \equiv r_\theta^{-1} U_\theta^\dagger \mathcal{G} \, U_\theta \, r_\theta \, .$$

Make invariants with $21 = (10,1) \oplus (5,2) \oplus (1,1)$

Embedding + symmetries $\rightarrow V = \sum_{i} c_{i} \mathcal{I}_{i}, c_{i} \approx \mathcal{O}(1)$

Potential

 \mathscr{I}_4 breaks $C_2 \Longrightarrow h$, H_0 can both acquire a VEV!

We want $\xi \approx 0.05$, $\beta \approx 0.1$, $m_h = 125 \,\text{GeV}$

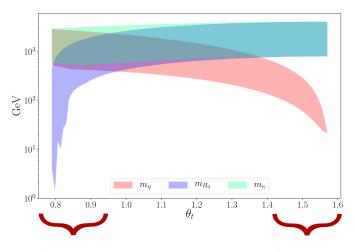
Where is the fine tuning?

$$\xi \approx c_3 + \frac{g_*^2}{y_L^2} c_1$$

$$\tan \beta \approx \frac{2y_L^2}{g_*^2} \frac{c_4}{c_2} \tan(2\theta_t)$$

$$m_h^2 \approx \frac{N_c g_*^2}{8\pi^2} m_t^2 \left(2\frac{y_L^2}{y_R^2} c_3 + \frac{y_R^2}{y_L^2} c_5 (3 + 4\cos(2\theta_t) + \cos(4\theta_t)) \right)$$

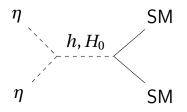
Spectrum



 $\theta_t \approx \pi/4$ Non-thermal relic

 $\theta_t \approx \pi/2$ Standard freeze-out

Relic density



$$SM = b, W, Z, h, t$$

 $\bigwedge W^*, Z^*$ off-shell decays can be relevant!

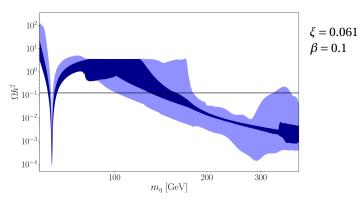
$$\Omega h^2 = \frac{0.03}{\int_{x_F}^{\infty} dx \frac{\sqrt{g_*}}{x^2} \frac{\langle \sigma v \rangle_{\text{eff}}}{1 \, \text{pb}}}, \quad \text{with } x_F \sim 25$$

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Relic density $\theta_t \lesssim \pi/2$

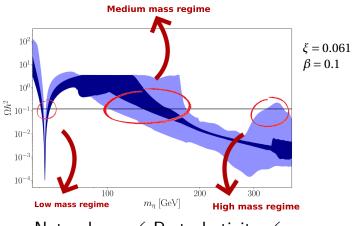
All processes included, no experimental limits (for now)



Naturalness ✓ Perturbativity ✓

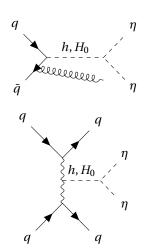
Relic density $\theta_t \lesssim \pi/2$

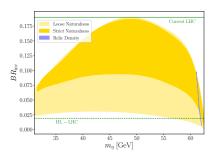
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Naturalness ✓ Perturbativity ✓

LHC

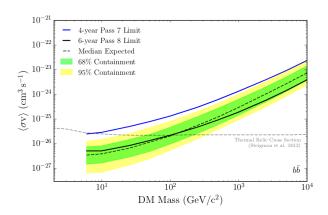




$$\Gamma_{h \to \eta \eta} = \frac{g_{\eta h}^2}{32\pi m_h} v^2 \sqrt{1 - \frac{4m_{\eta}^2}{m_h^2}}$$

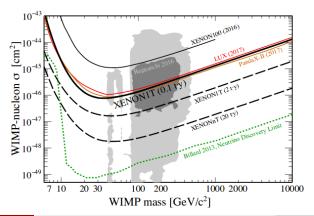
Indirect detection

- \Rightarrow We use γ -ray data
- \Rightarrow Overestimate $\sigma_{b\bar{b}} = \sigma_{\text{tot}}$
- ⇒ Processes are not p-wave suppressed



Direct detection

- ⇒ We use XENON result
- ⇒ Energy exchanged of order keV
- ⇒ Processes are not p-wave suppressed



Experimental constraints

- LHC: (CMS, 1809.05937) MJ, VBF weaker than BR_{inv} and not relevant at higher masses BR_{inv} constraining in the HL projection

ID: (FERMI-LAT, 1503.02641)
 dSPh data: we are safe

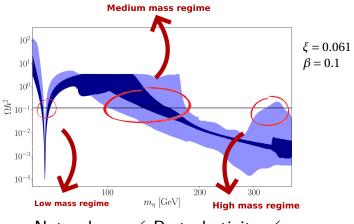
- DD: (XENON1T, 1705.06655)

$$a_{q} = \frac{1}{2} \left[\frac{g_{q}}{v^{2}} - \left(k_{q} \frac{g_{\eta h}}{m_{h}^{2}} - k_{H_{0}q} \frac{g_{\eta H_{0}}}{m_{H_{0}}^{2}} \right) \right]$$

cancellation possible, increasing coefficient tuning to evade

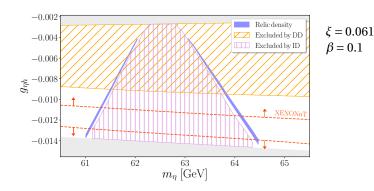
Relic density $\theta_t \lesssim \pi/2$ (again)

All processes included, no experimental limits (for now)



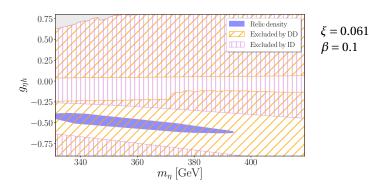
Naturalness ✓ Perturbativity ✓

Low mass regime (benchmark)



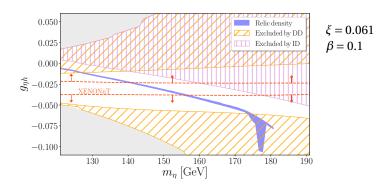
Limited by DD, ID safe by a factor $1.1 \div 1.2$ Available mass range $\lesssim 1\,\text{GeV}$ wide

High mass regime (benchmark)



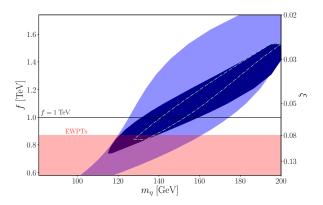
Excluded by DD up to $\xi \approx 0.01$ Mass range defined by $H_0 - h$ cancellation

Medium mass regime (benchmark)



Limited by DD, insensitive to other searches Available mass range $130\,\text{GeV} \div 175\,\text{GeV}$

Medium mass regime (scan)



Relic density ✓ Direct detection ✓

Lowest fine tuning set by EWPTs

XENONnT will reduce, but not exclude, the model parameter space

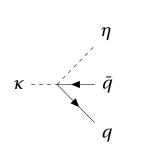
Relic density $\theta_t \gtrsim \pi/4$

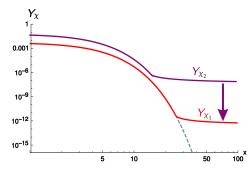
Let's go to the other θ_t regime

 $\theta_{\it t} \gtrsim \frac{\pi}{4}$ has many consequences!

- \Rightarrow η becomes heavy (~ TeV) $m_{\kappa} > m_{\eta} > m_{H_0}$, but η still stable
- ⇒ Particles become close in mass, so coannihilations and long-lived particles
- \Rightarrow many natural suppressions become unnatural, because $\sin(2\theta_t) \sim 1$

Relic density $\theta_t \gtrsim \pi/4$





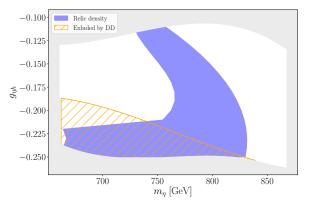
Two contributions: $\Omega_{\rm DM} h^2 = \Omega_{\eta} h^2 + \frac{m_{\eta}}{m_{\kappa}} \Omega_{\kappa} h^2$

Late time decay: $t_{FO} < \tau_{\kappa} < t_{BBN}$

$$\Gamma_{\kappa} \stackrel{\propto}{\sim} \frac{\Delta m^5 m_q^2}{m_k^2 v^4} \rightarrow 20 \, \text{GeV} \lesssim \Delta m \lesssim 50 \, \text{GeV}$$

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Non thermal



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 $\xi = 0.01$ $\beta = 0.2$

Evade DD $\rightarrow \xi \approx 0.01$ needed $\Omega_{\kappa} \approx \Omega_{\eta}$: same order contribution

Conclusions

- CH model based on SO(7)/SO(5)xSO(2) delivers a viable DM candidate consistent with all exp. results;
- the level of fine tuning required is $f \gtrsim 0.8 \, \text{TeV}$ (set by EWPT);
- DM can be produced non-thermally via decays of an heavier pNGB (but with higher fine tuning).

Thank you!

BACKUP

$$\xi \approx \frac{2N_c y_L^4 c_{(2,0)}^{(1)} + g_*^2 \left(N_c y_L^2 c_{(1,0)}^{(1)} - 3g^2 c_g^{(1)} - g'^2 c_{g'}^{(1)} \right)}{N_c y_L^4 c_{(2,0)}^{(1)}}$$

$$\tan \beta \approx \frac{N_c c_{(1,1)}^{(1)} y_L^2 y_R^2 \sin 2\theta_t}{2g_*^2 (g'^2 c_{g'}^{(1)} + 2N_c y_R^2 c_{(0,1)}^{(1)} \cos 2\theta_t)}$$

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$$a_{q} = \frac{1}{2} \left[\frac{g_{q}}{v^{2}} - \left(k_{q} \frac{g_{\eta h}}{m_{h}^{2}} - k_{H_{0} q} \frac{g_{\eta H_{0}}}{m_{H_{0}}^{2}} \right) \right]$$

$$g_q \approx -2\xi \frac{\cos\beta\cos\theta_q}{\cos(\beta-\alpha_q\theta_q)}$$

$$k_q = 1 + \mathcal{O}(\xi)$$

$$g_{\eta h} \approx -\frac{g_*^2}{8\pi^2} \cos^2\theta_t \left[2\tilde{c}_y^{(5)} - \tilde{c}_y^{(7)} + 2\tilde{c}_y^{(8)} + 2\cos(2\theta_t)(\tilde{c}_y^{(5)} + \tilde{c}_y^{(8)}) \right]$$

Fixing coefficients ≠ **Tuning coefficients**

Fixing coefficient (with tuning)

$$\xi = c_1 + c_2 = 0.01 \rightarrow c_1 = 1.01, c_2 = -1$$

Change
$$c_2$$
: $\delta c_2 = +5\% \rightarrow \delta \xi = +500\%$

Fixing coefficient (without tuning)

$$m_h = c_1 + c_2 = 2 \rightarrow c_1 = 1, c_2 = 1$$

Change
$$c_2$$
: $\delta c_2 = +5\% \rightarrow \delta m_h = +2.5\%$

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