

# Prospects for detecting the axion annihilation signal using continuous gravitational-wave searches

Sylvia J. Zhu (DESY, AEI Hannover)

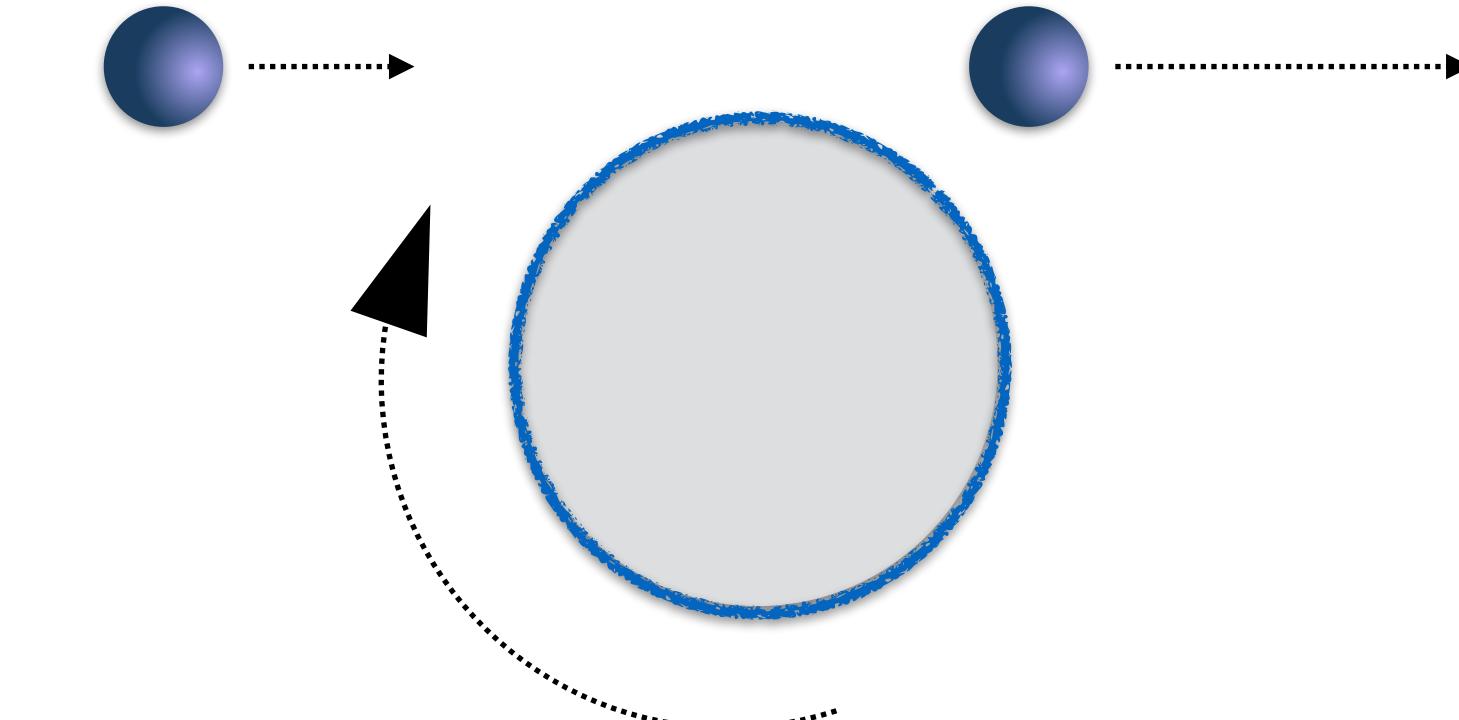
[arXiv:2003.03359](https://www.aei.mpg.de/continuouswaves/arxiv200303359)  
<https://www.aei.mpg.de/continuouswaves/arxiv200303359>

# Superradiance instability

Particles around a black hole increase in number by extracting black hole spin

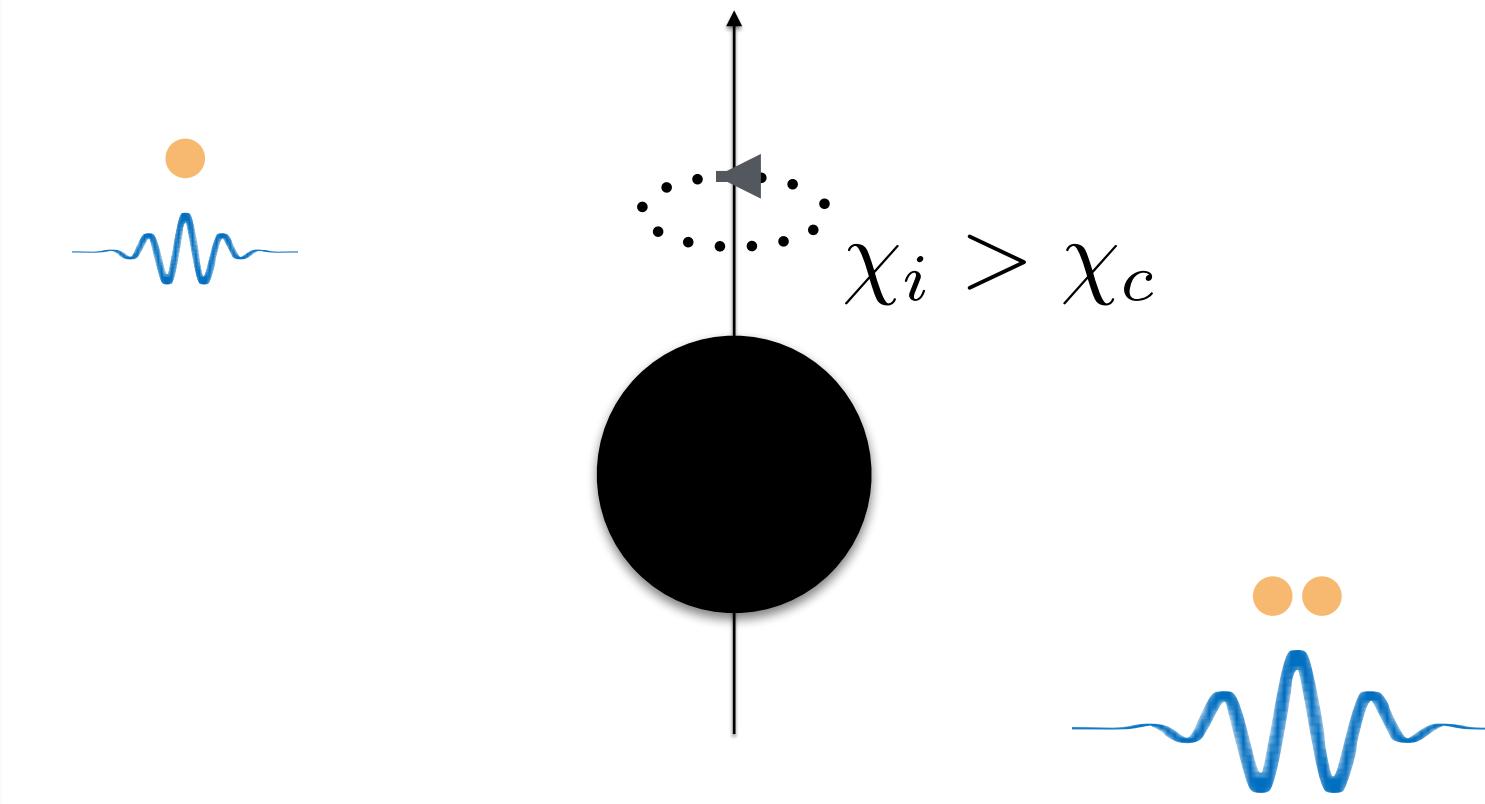
classical analog

$$v_{\text{ball}} < v_{\text{cylinder}}$$

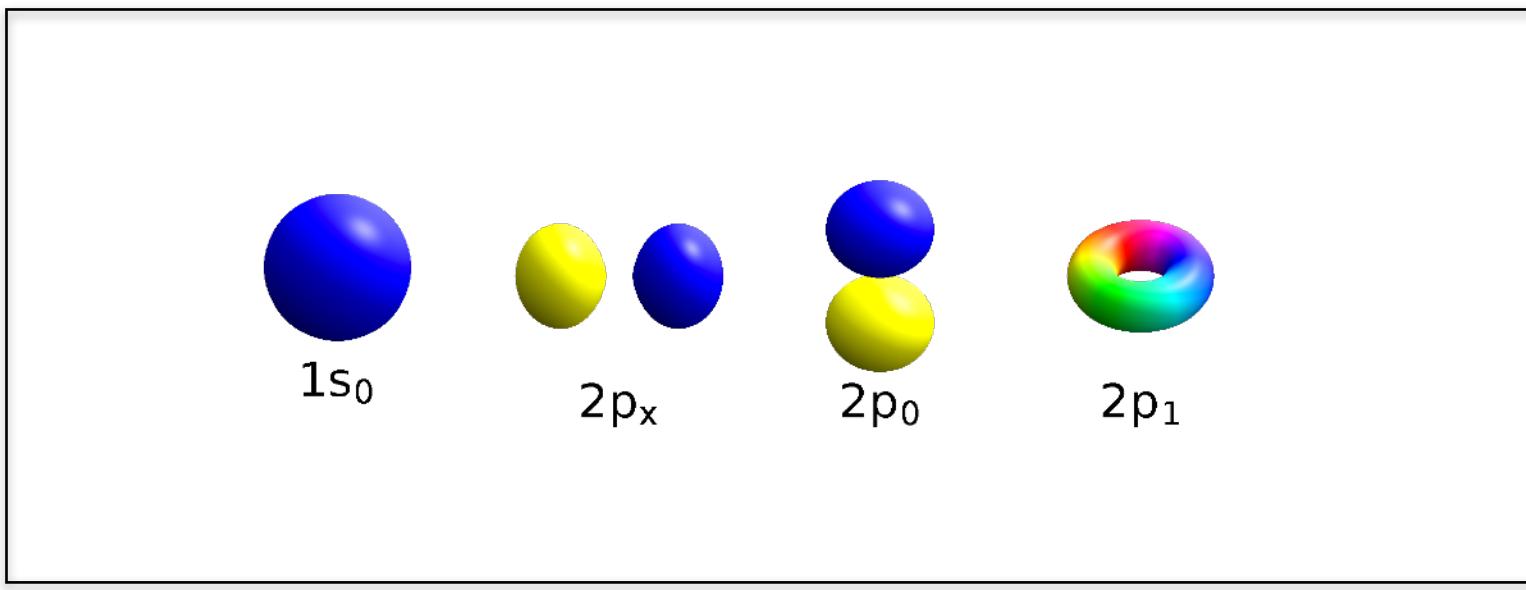
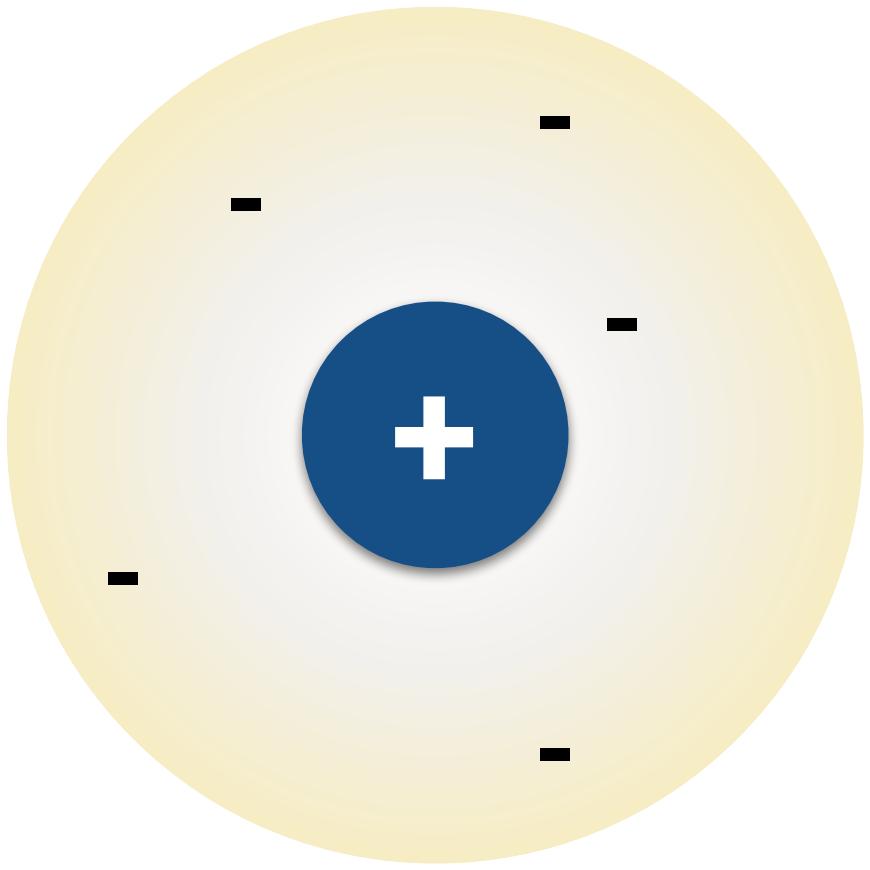


superradiance

$$\Omega_b < \Omega_{\text{BH}}$$

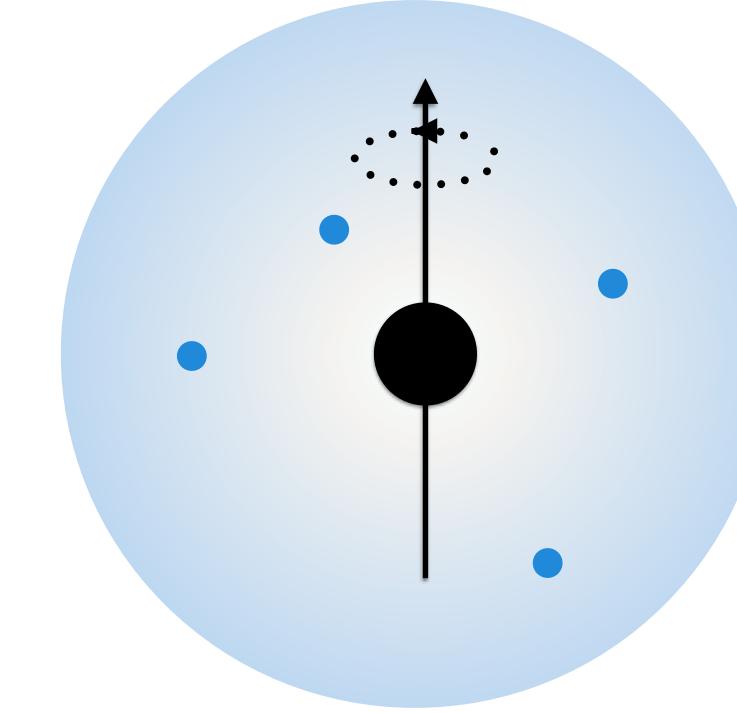


# Gravitational “atom”

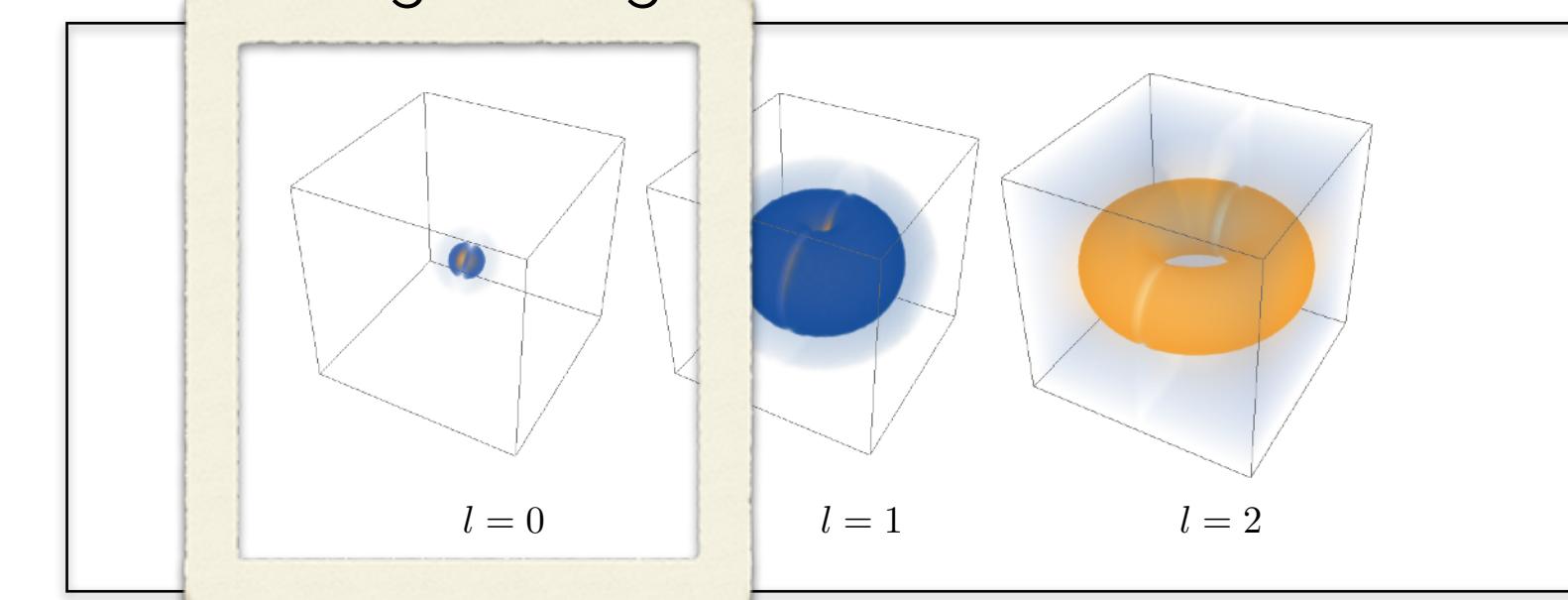


[Wikimedia Commons]

$$\alpha \propto q_e q_{\text{nuc}}$$



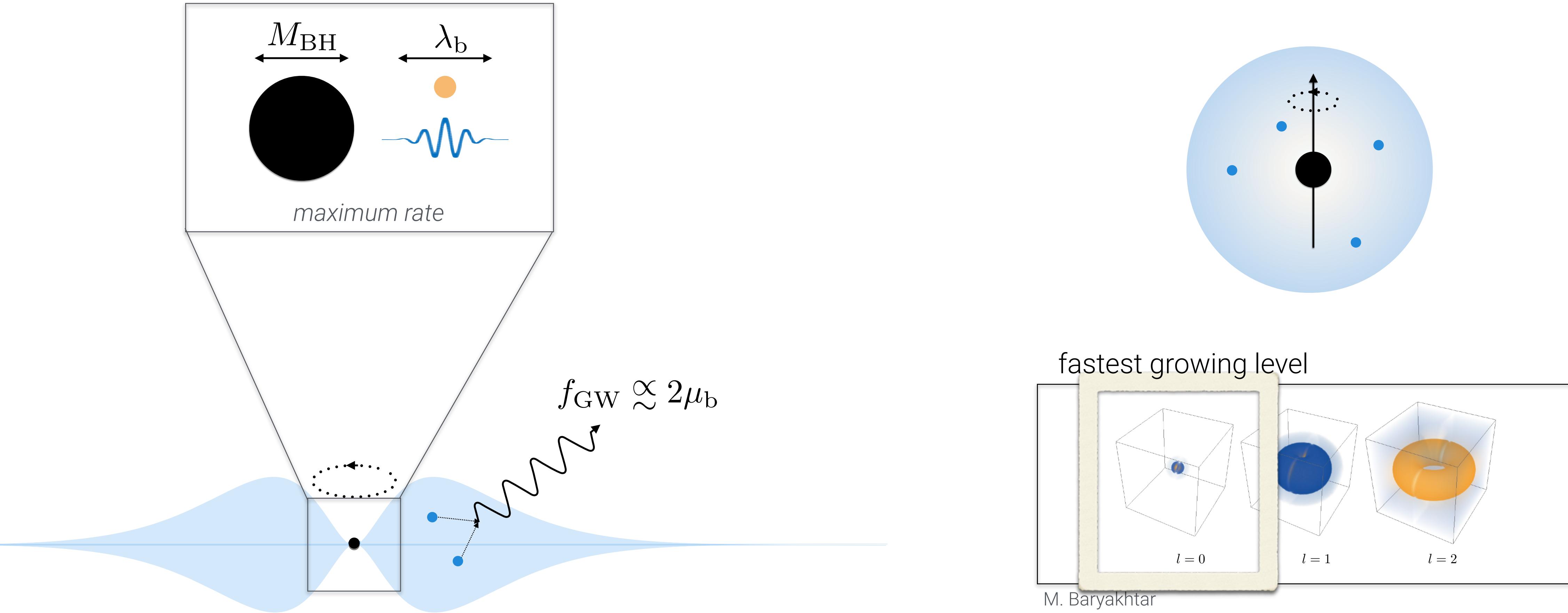
fastest growing level



M. Baryakhtar

$$\alpha \propto \mu_b M_{\text{BH}}$$

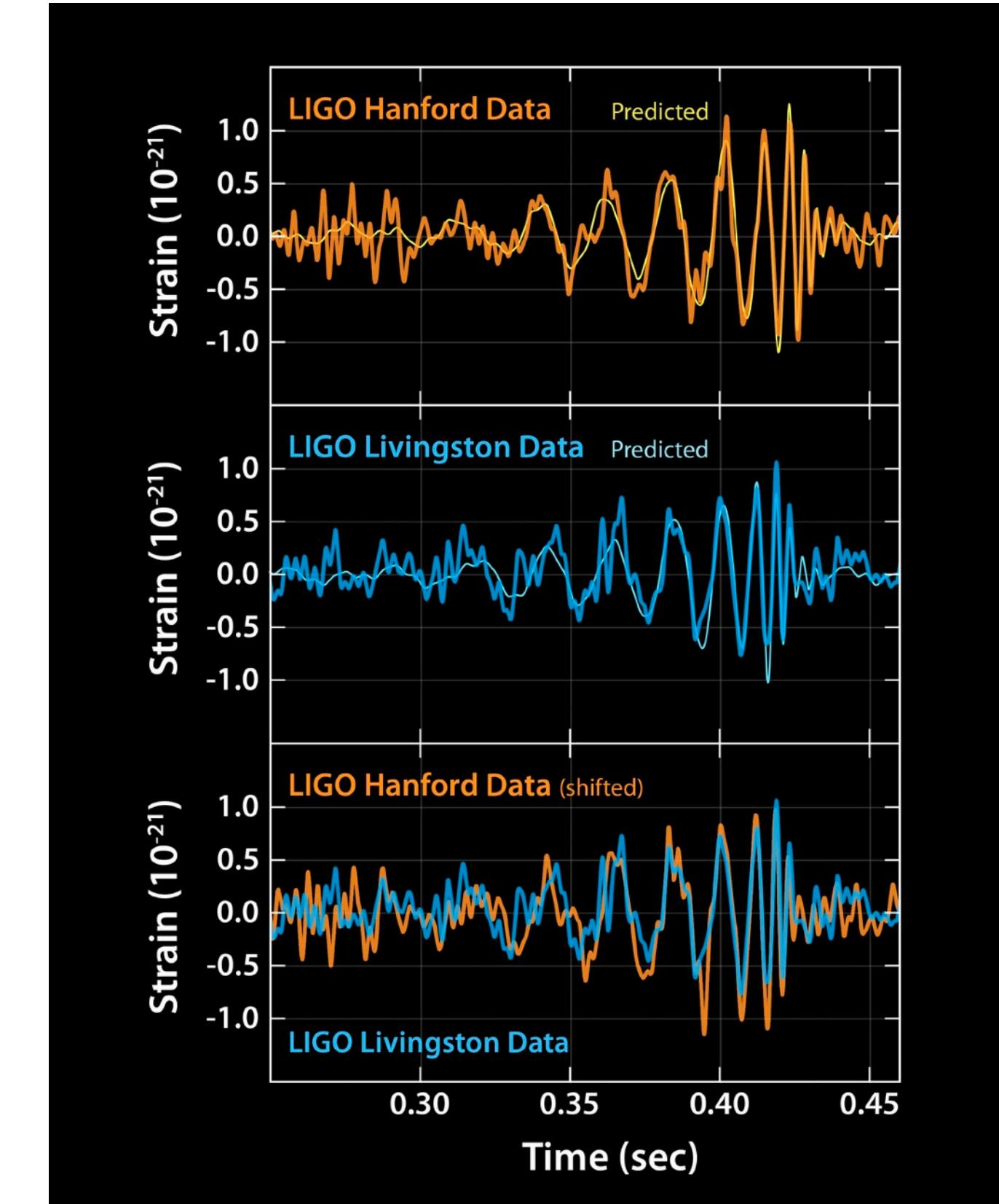
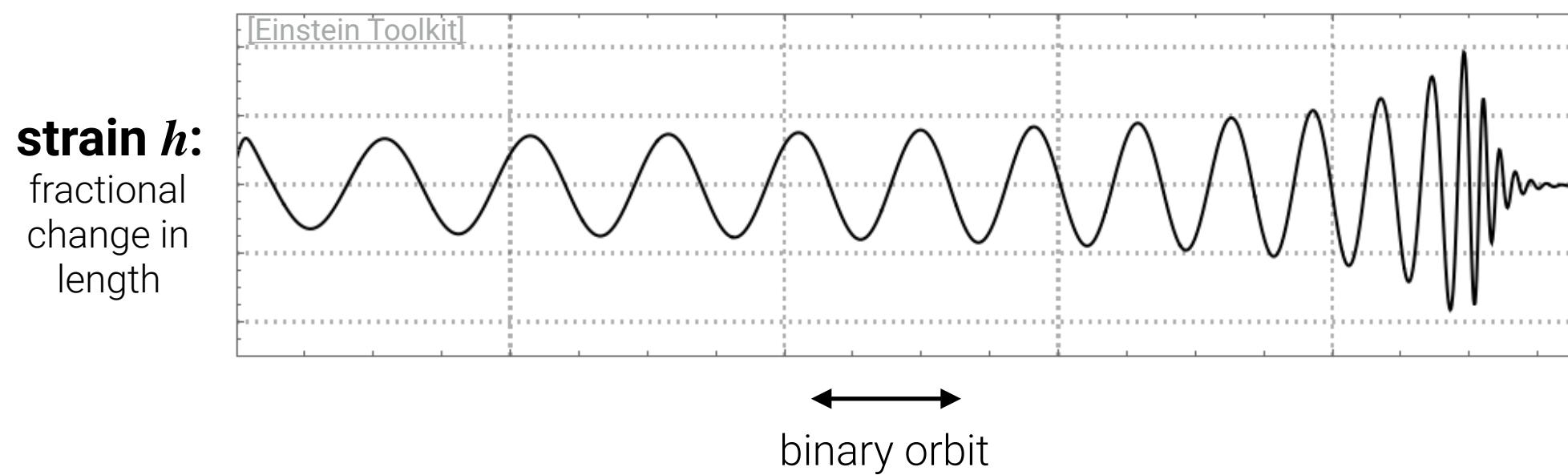
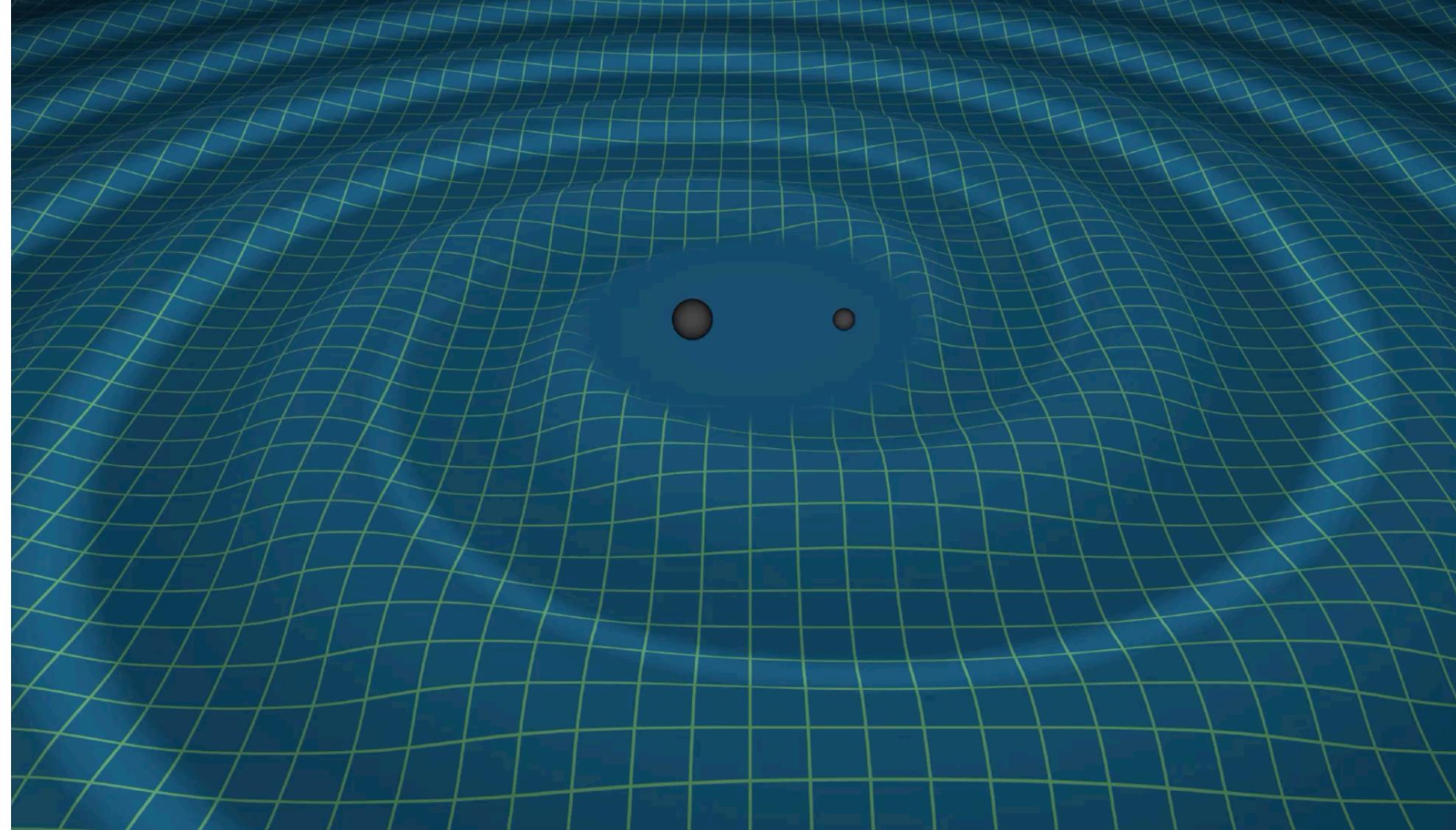
# Gravitational “atom” sources continuous gravitational waves



$$\alpha \propto \mu_b M_{\text{BH}}$$

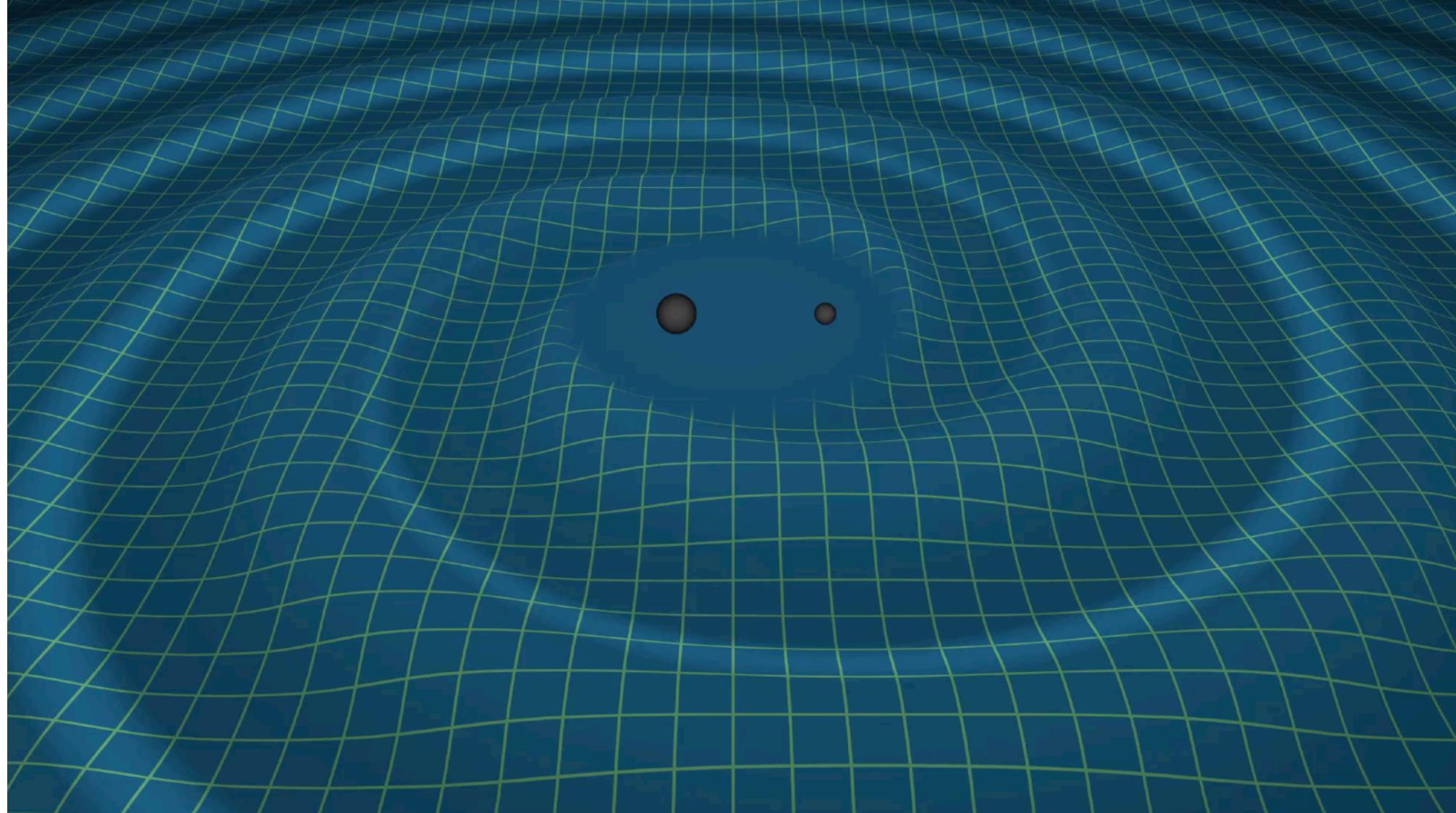
# Some types of gravitational-wave signals

transient signals

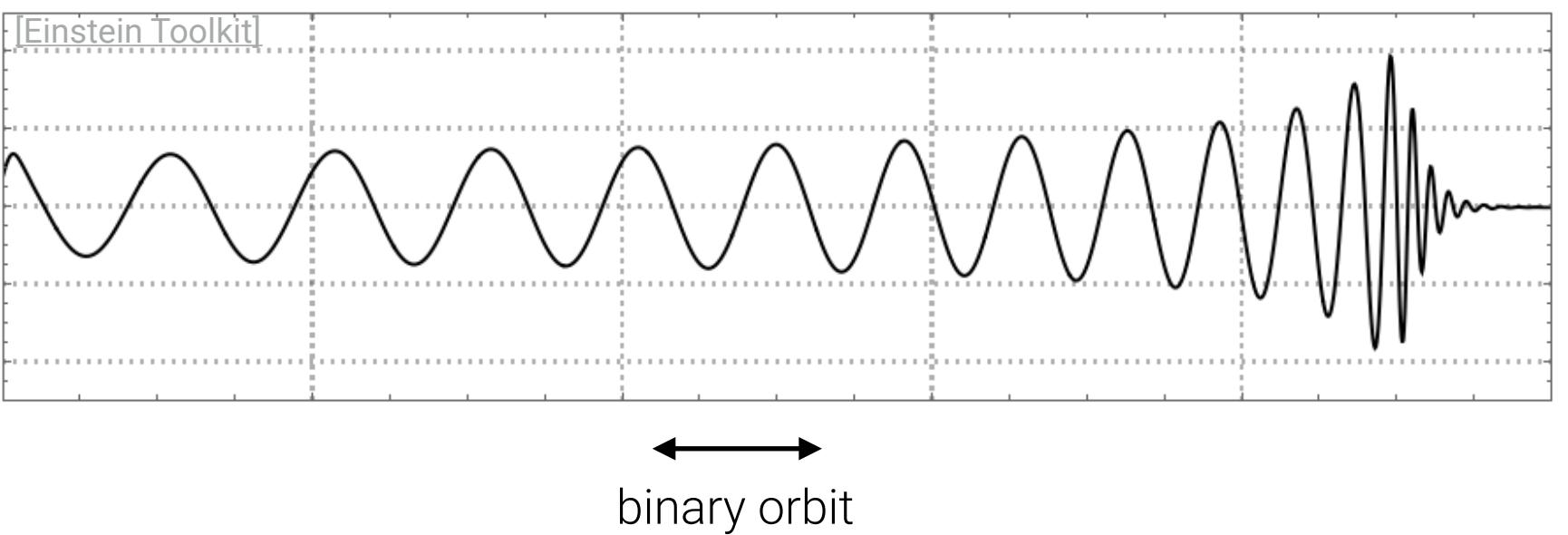


# Some types of gravitational-wave signals

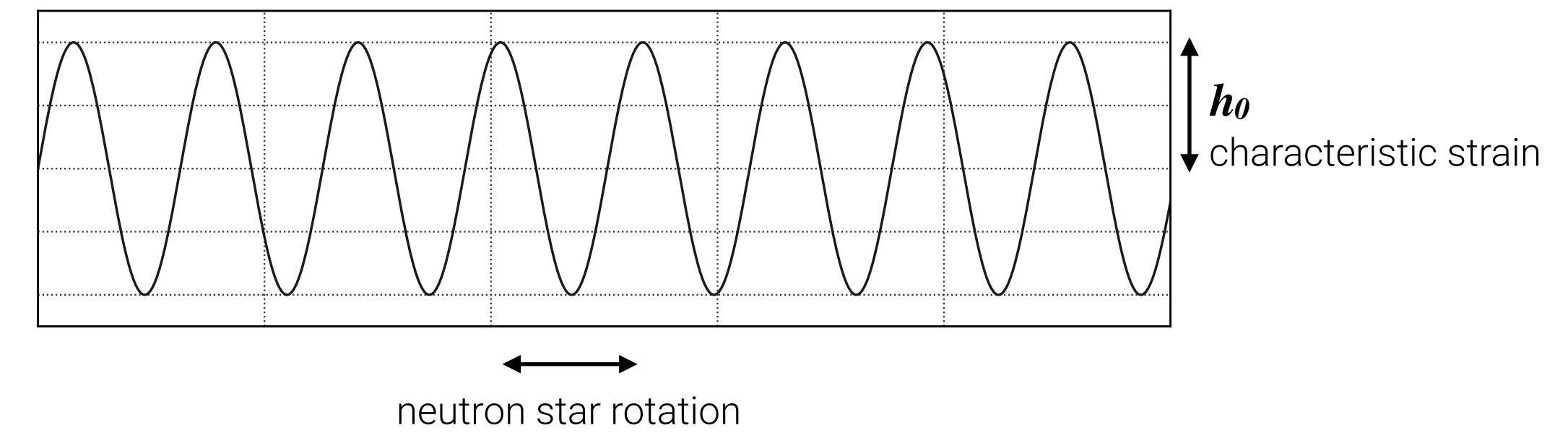
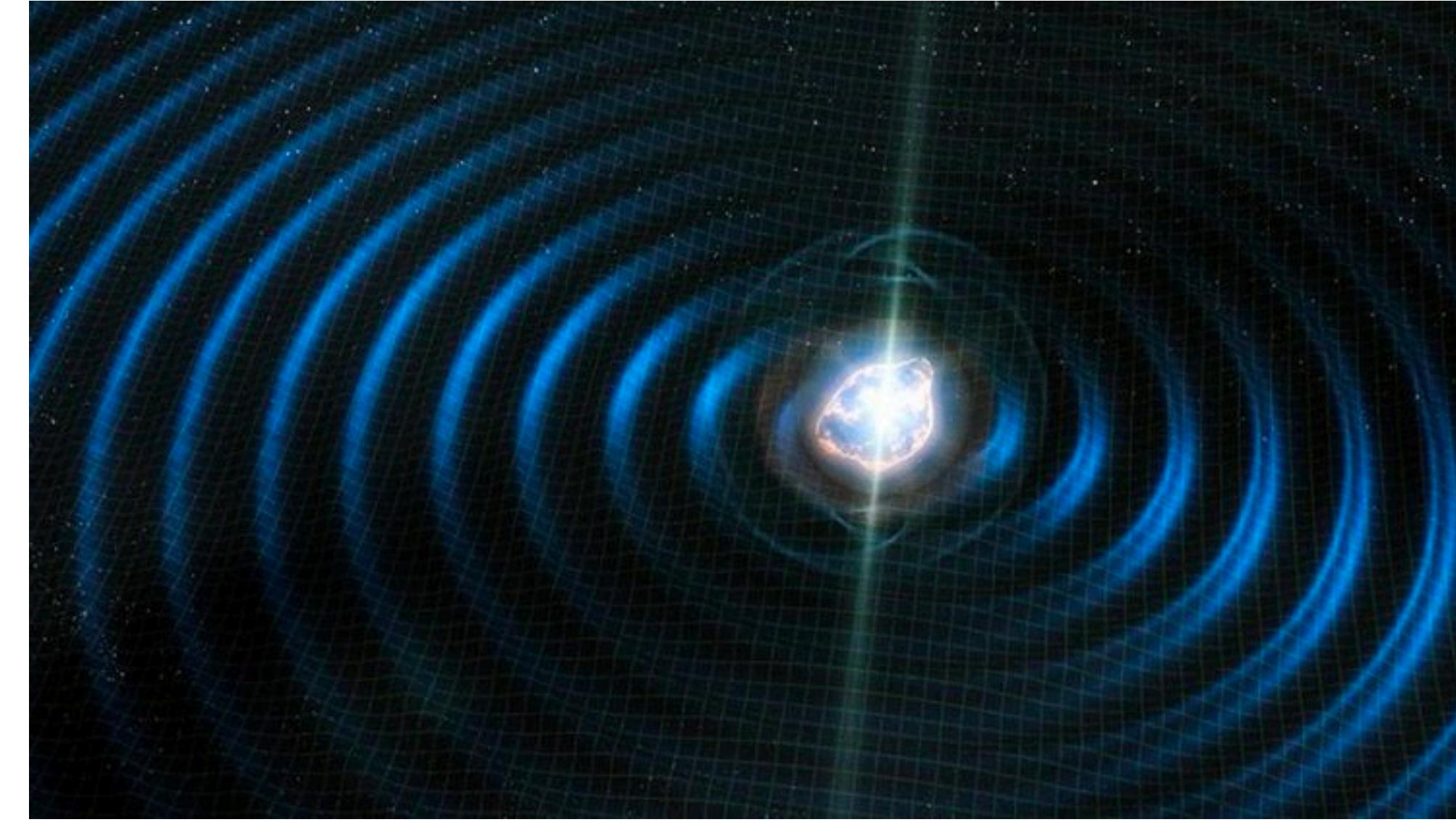
transient signals



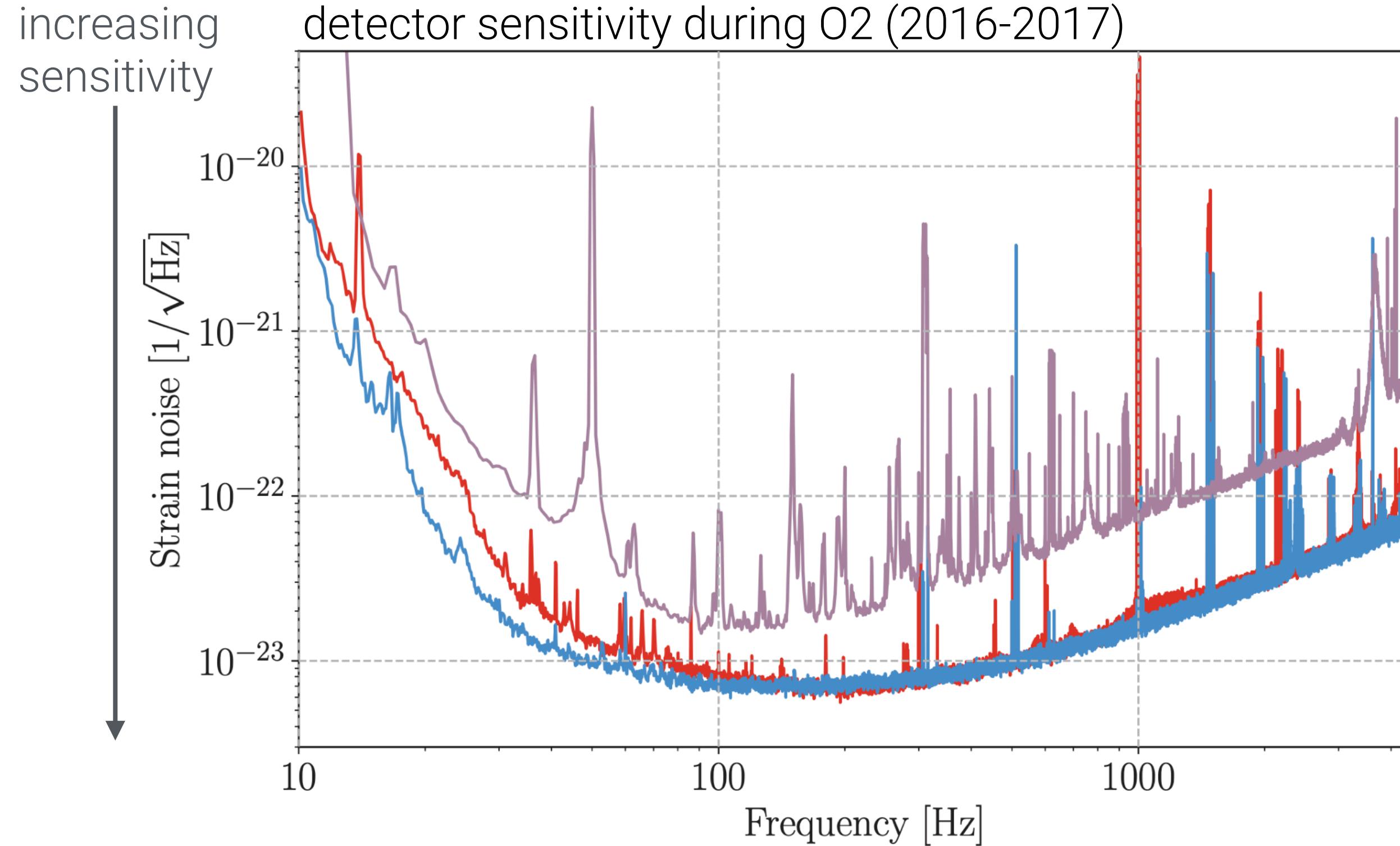
**strain  $h$ :**  
fractional change in length



continuous signals

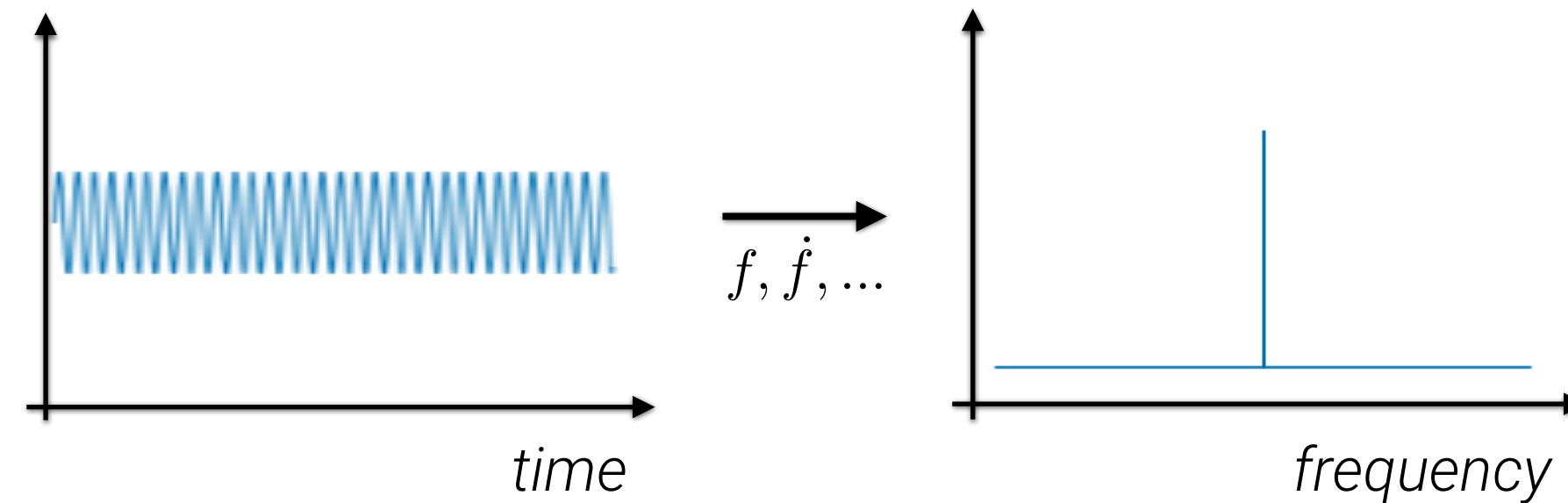
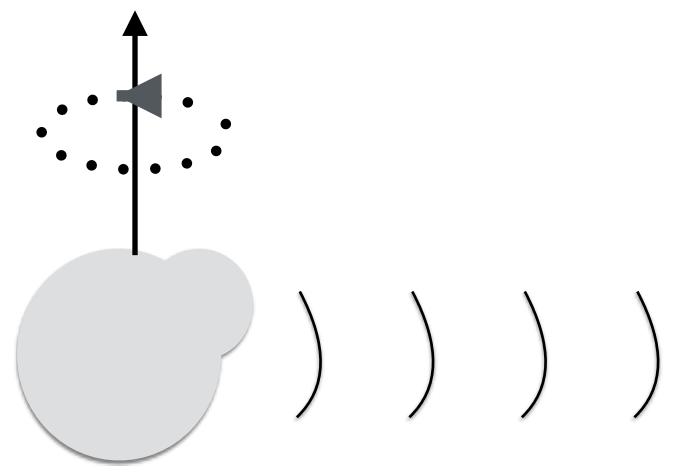


# Gravitational-wave interferometers

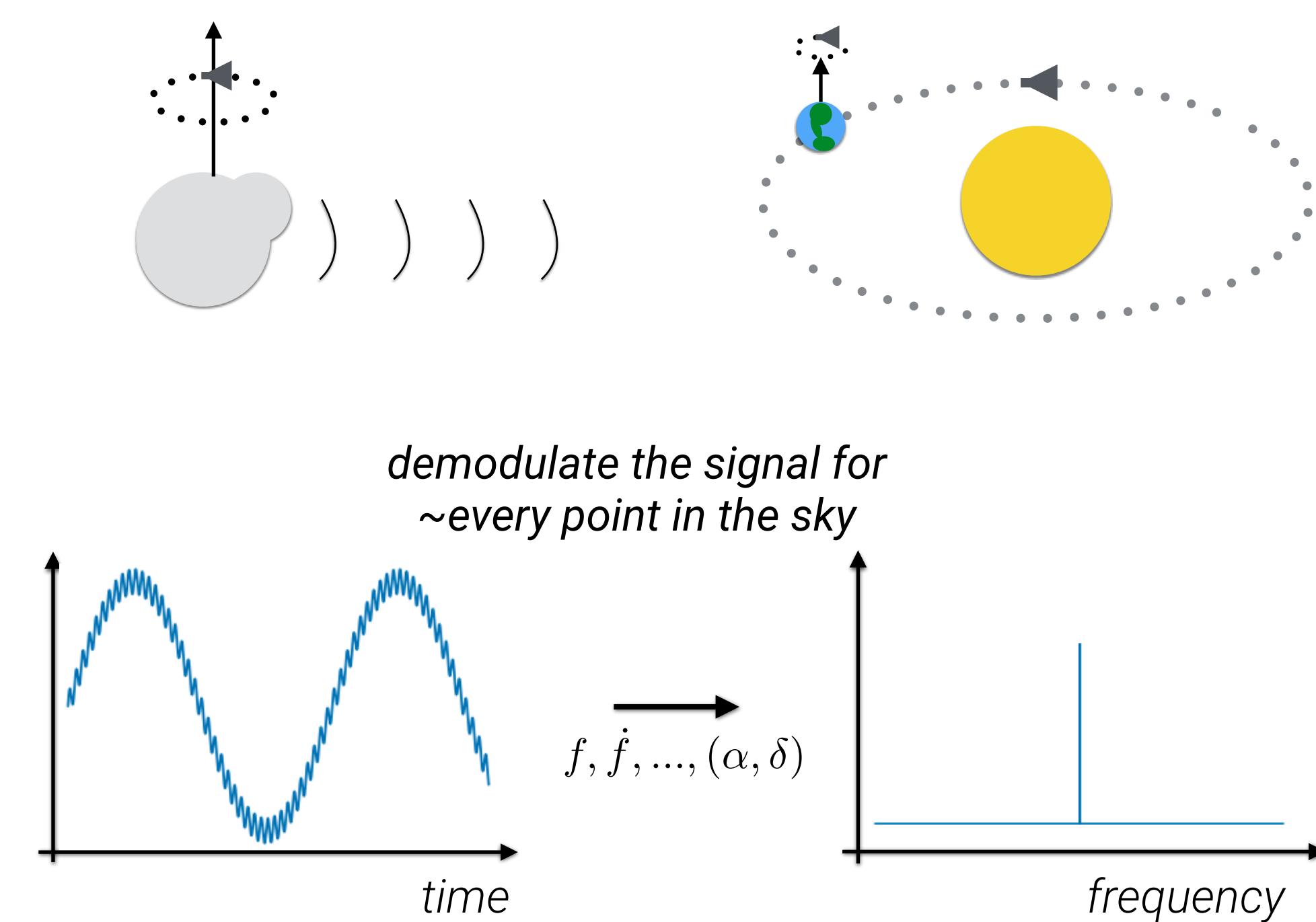
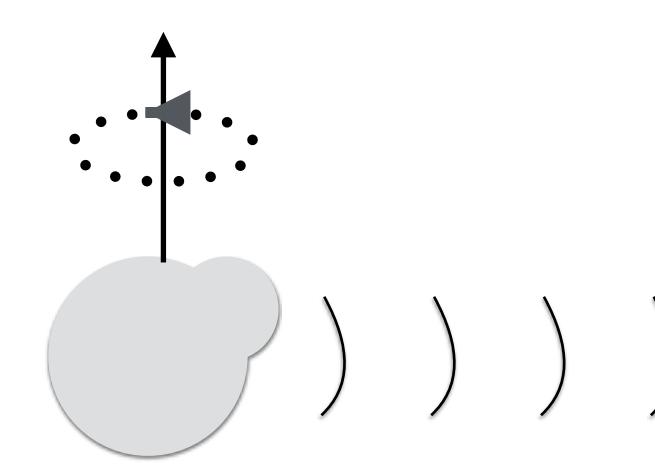


# Searching for continuous gravitational waves

the waveform at the **source**:



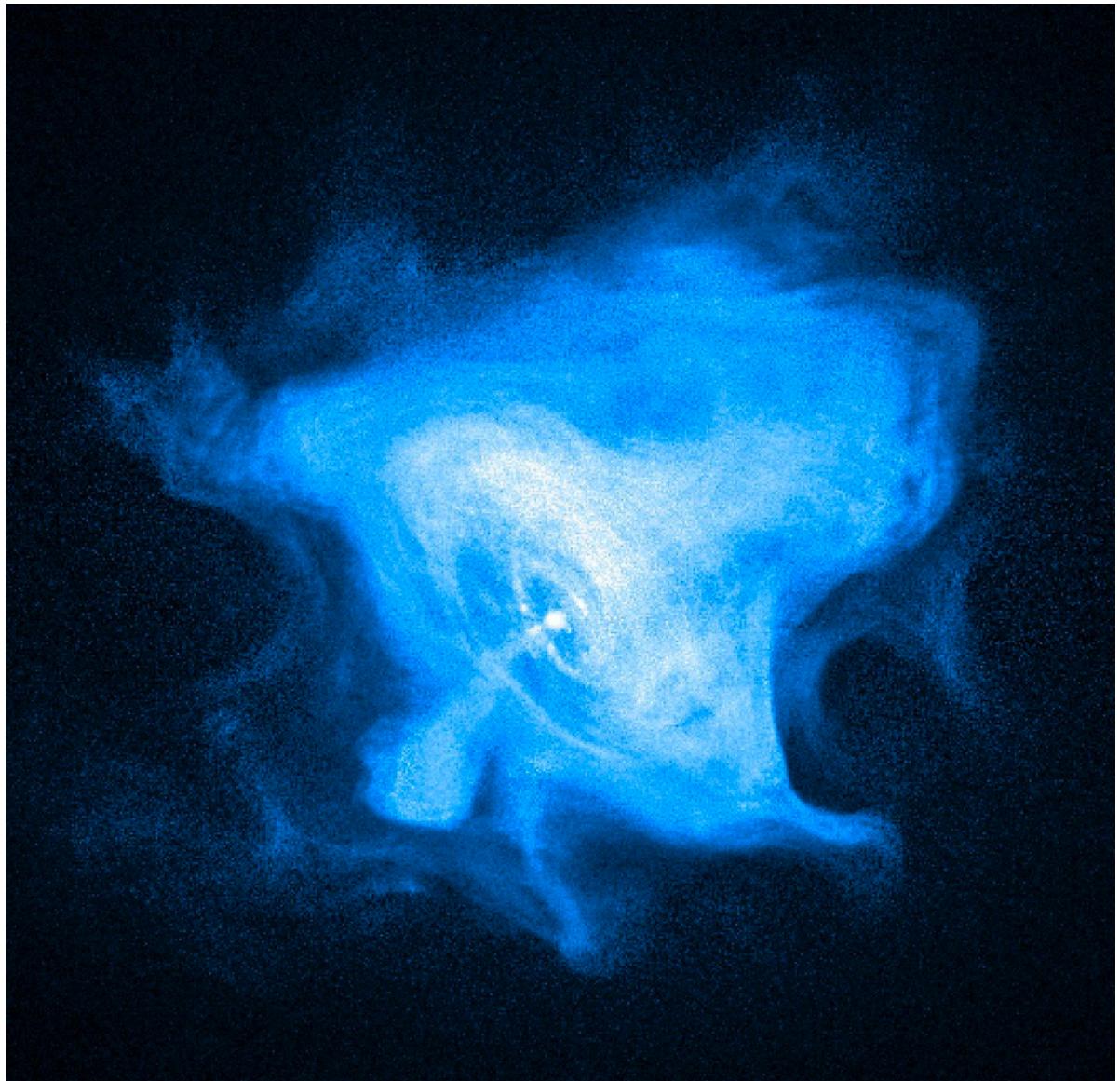
the waveform at the **detector**:



# Types of searches

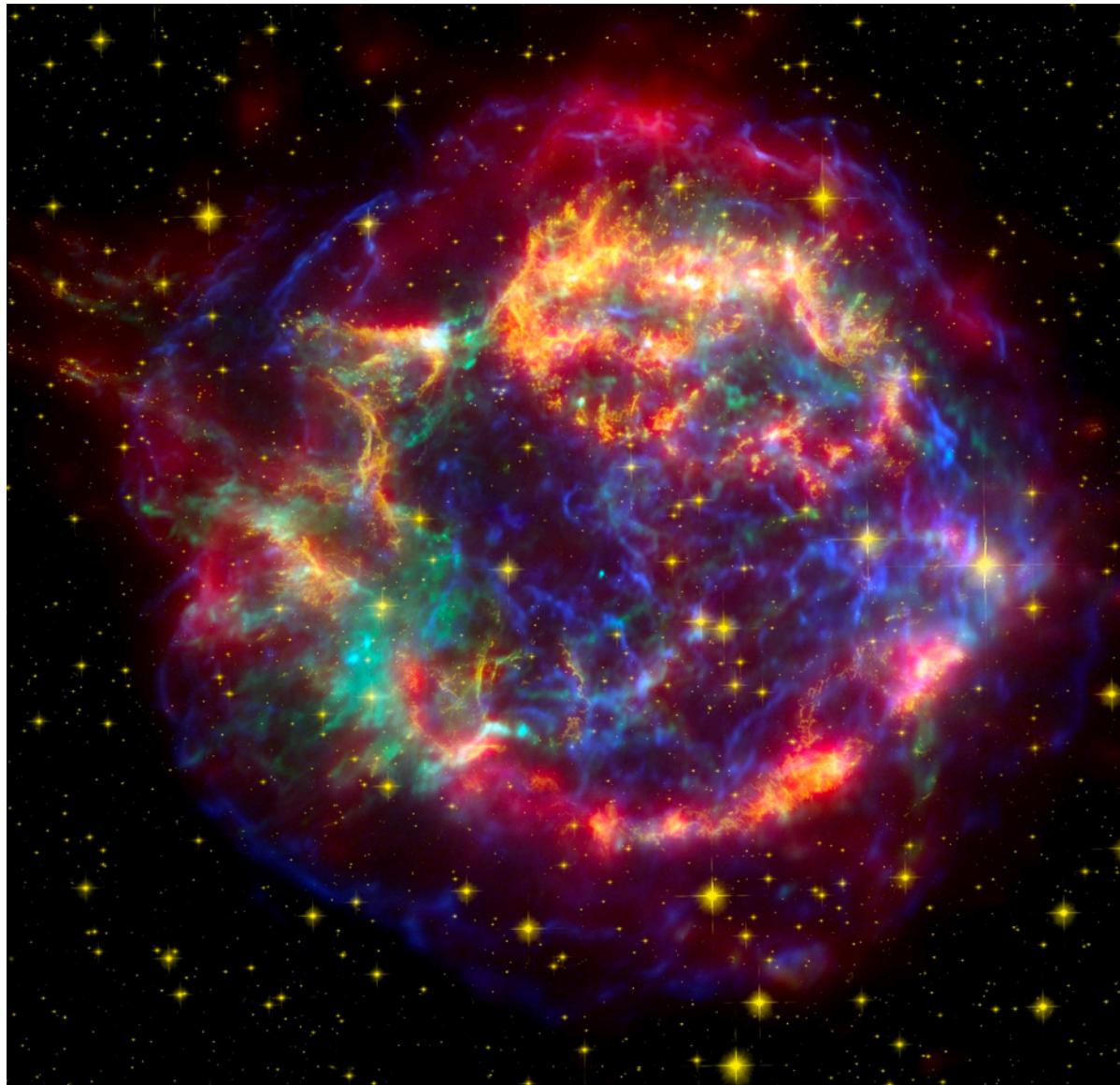
## targeted search

known pulsars



## directed search

known neutron stars



## all-sky search

minimal assumptions

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computational cost

search sensitivity

# Types of searches

## targeted search

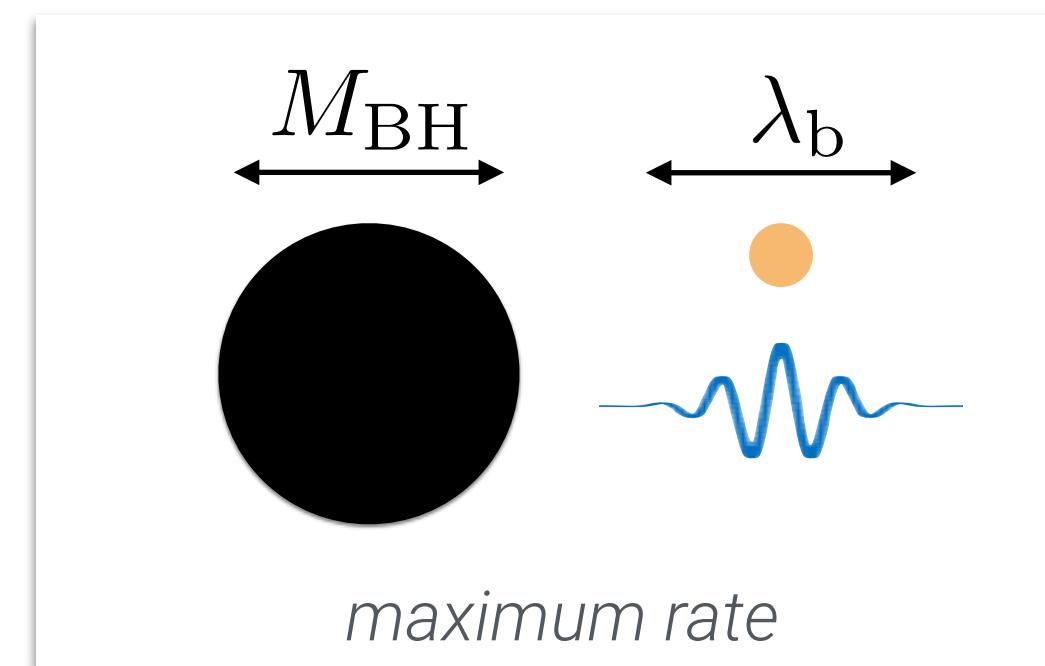
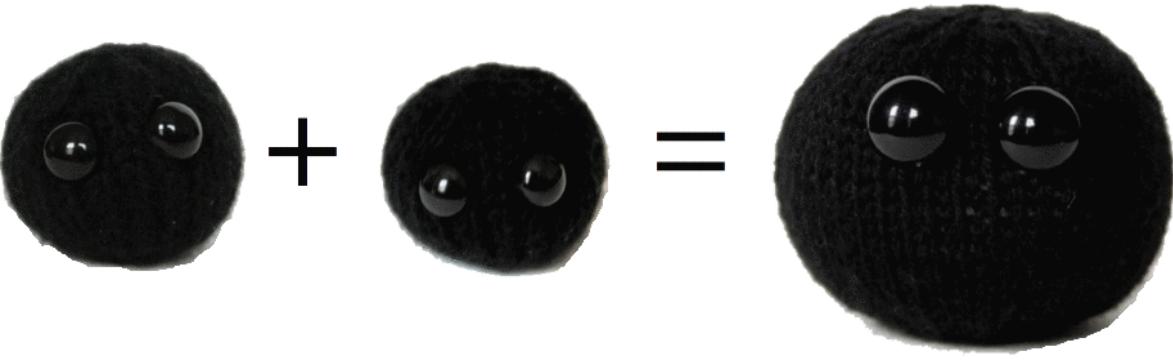
known black hole  
+  
known boson mass



## directed search

known black hole  
+  
unknown boson mass

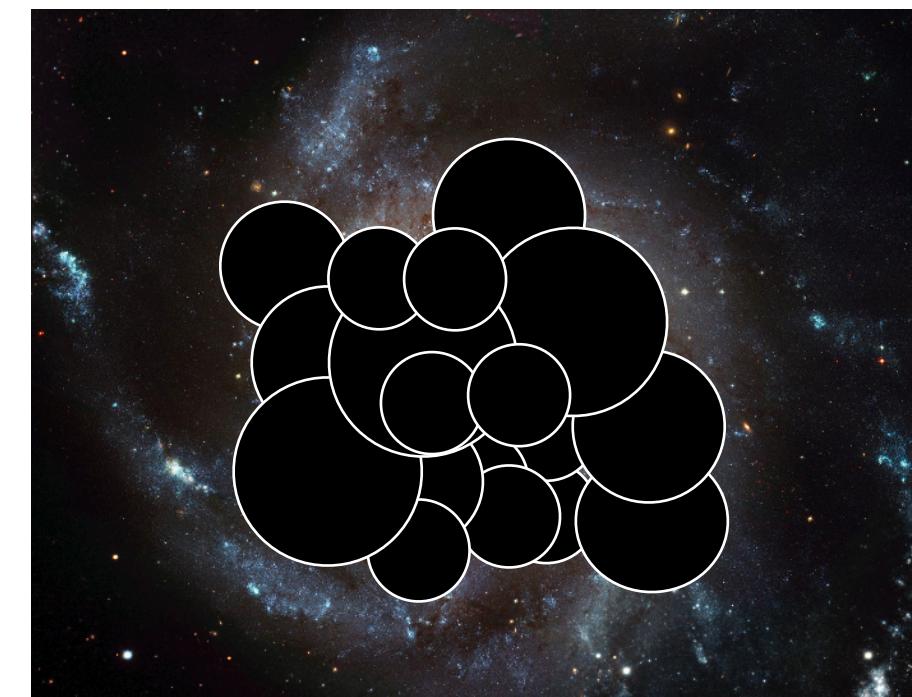
e.g., [Isi et al., PRD 99 \(2019\)](#)



## all-sky search

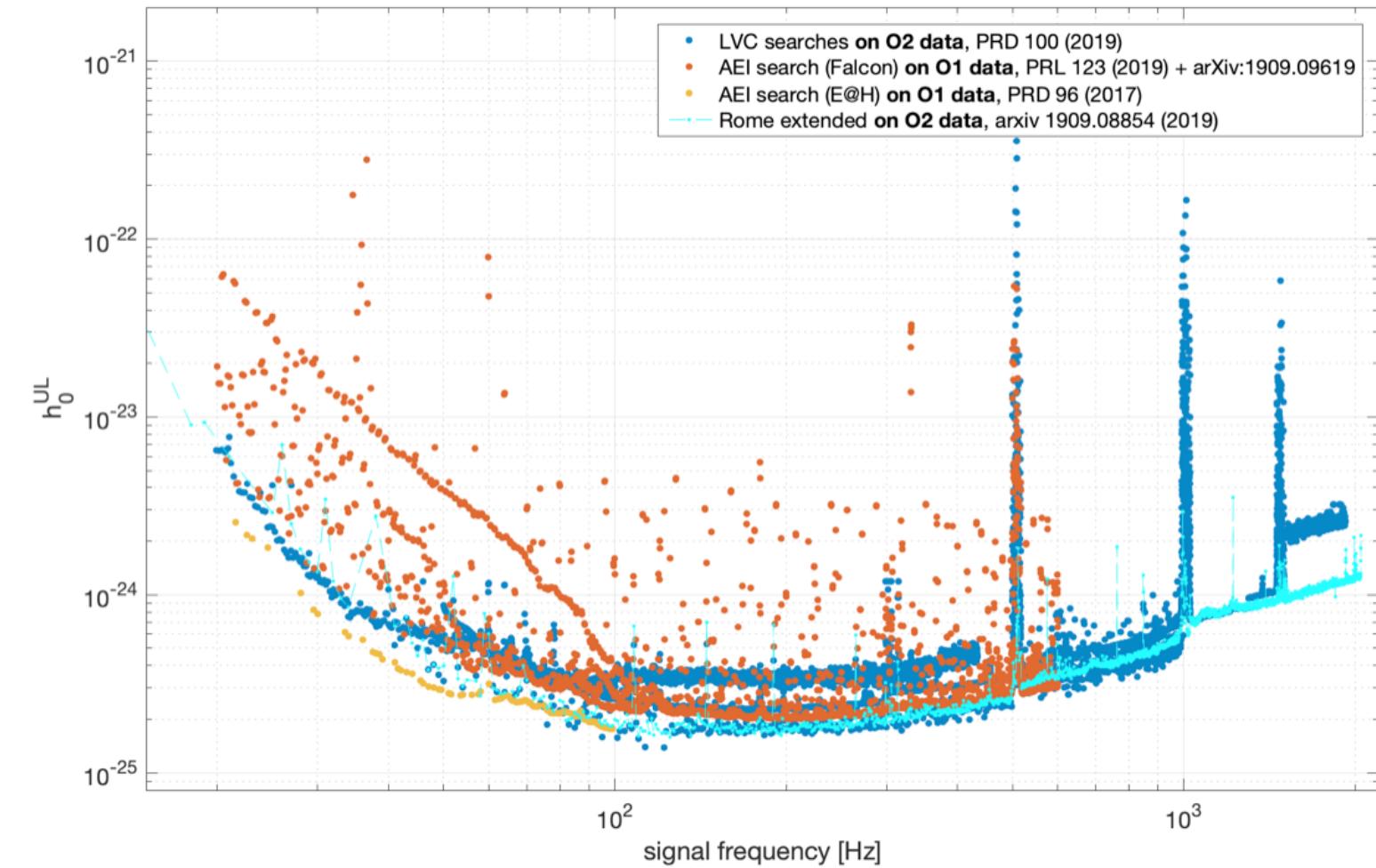
minimal assumptions on  
black hole positions or  
boson mass

e.g., [Palomba et al., PRL 123 \(2019\)](#)



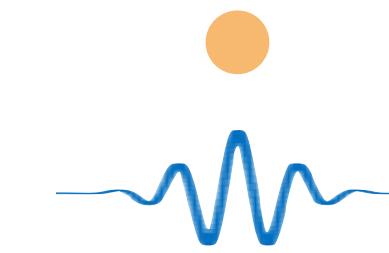
# Interpreting null results (upper limits) – boson clouds

upper limit on gravitational-wave signal strength

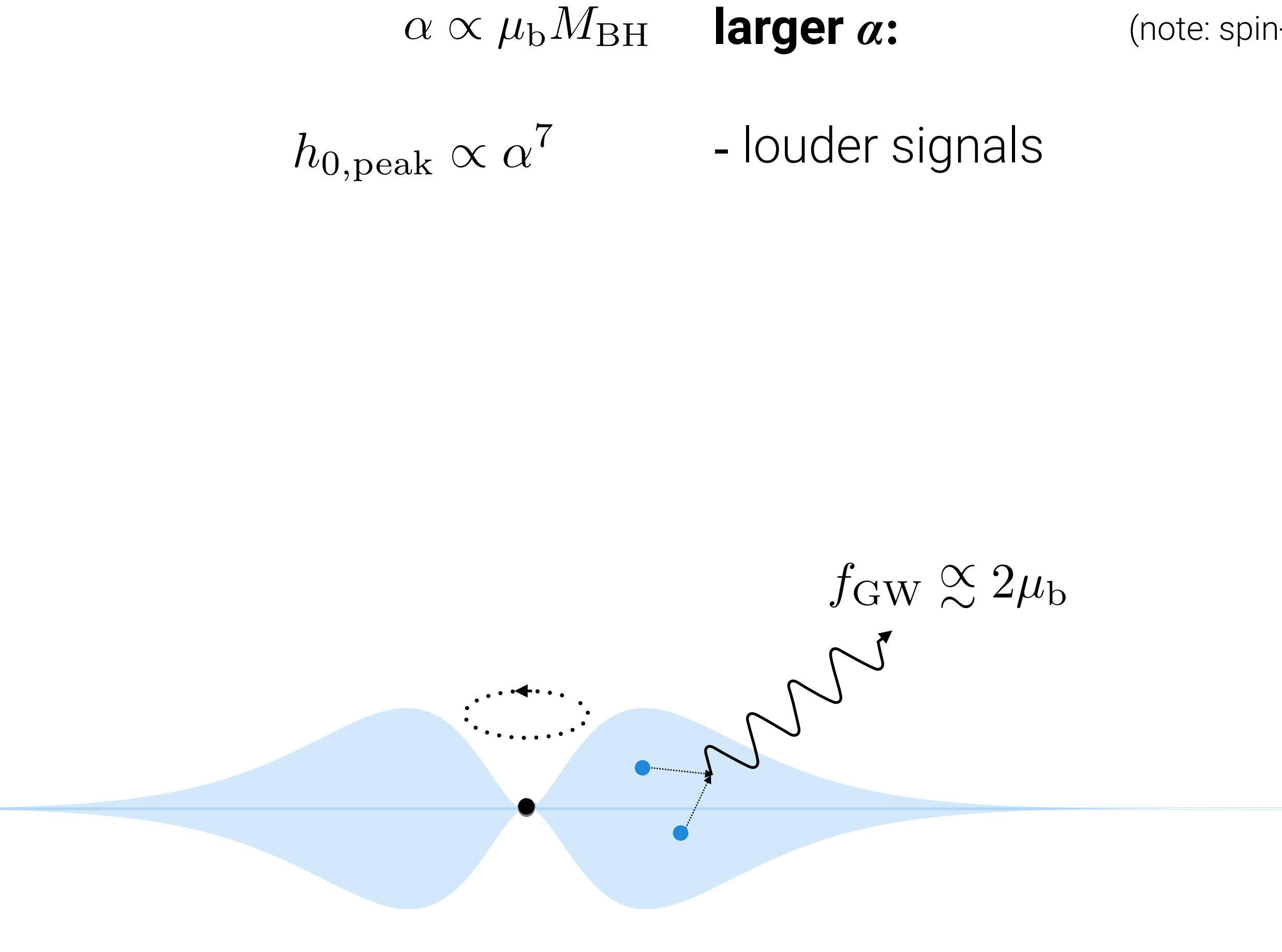


?

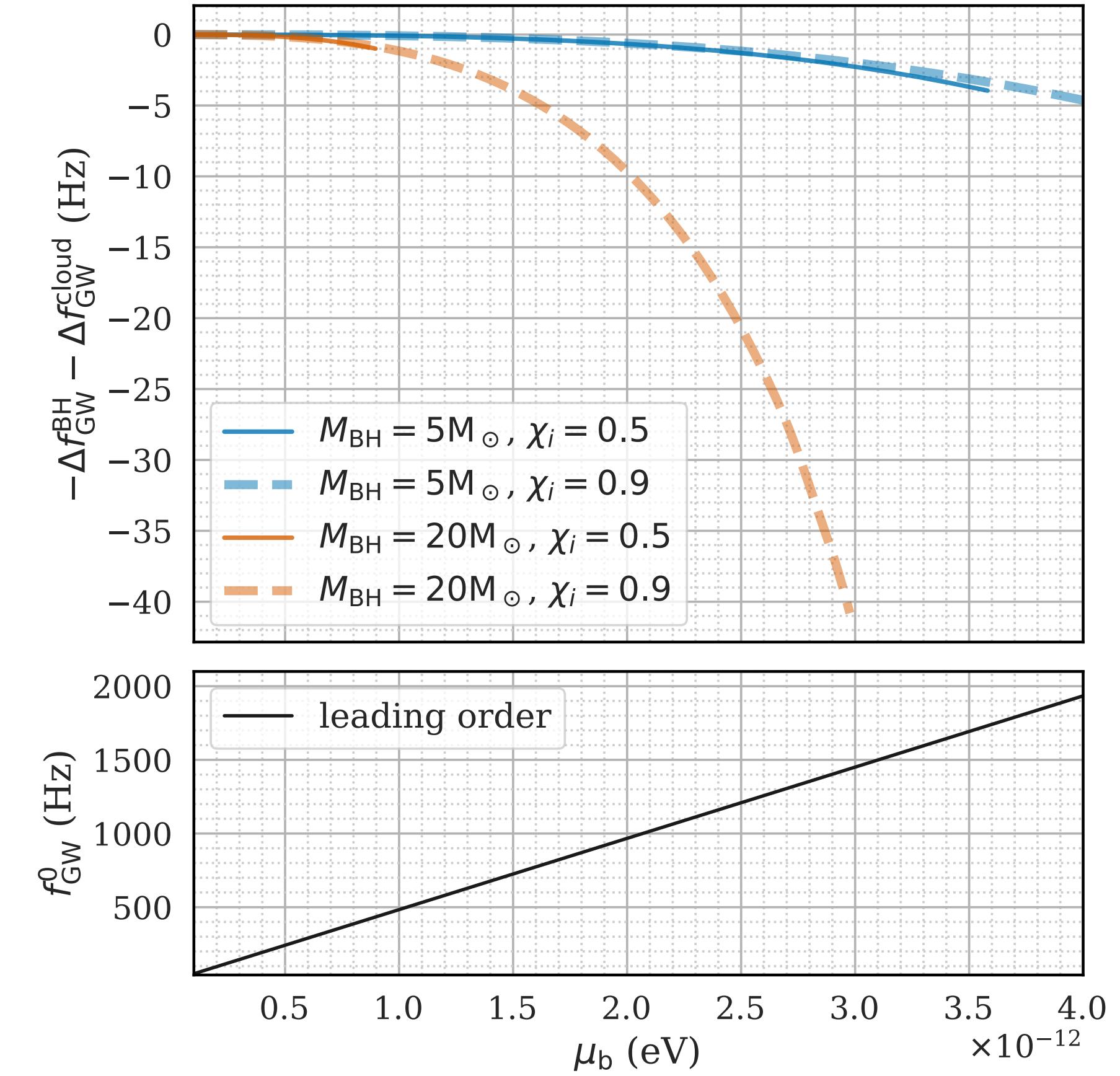
range of disfavored boson masses



# Gravitational “atom” sources continuous gravitational waves

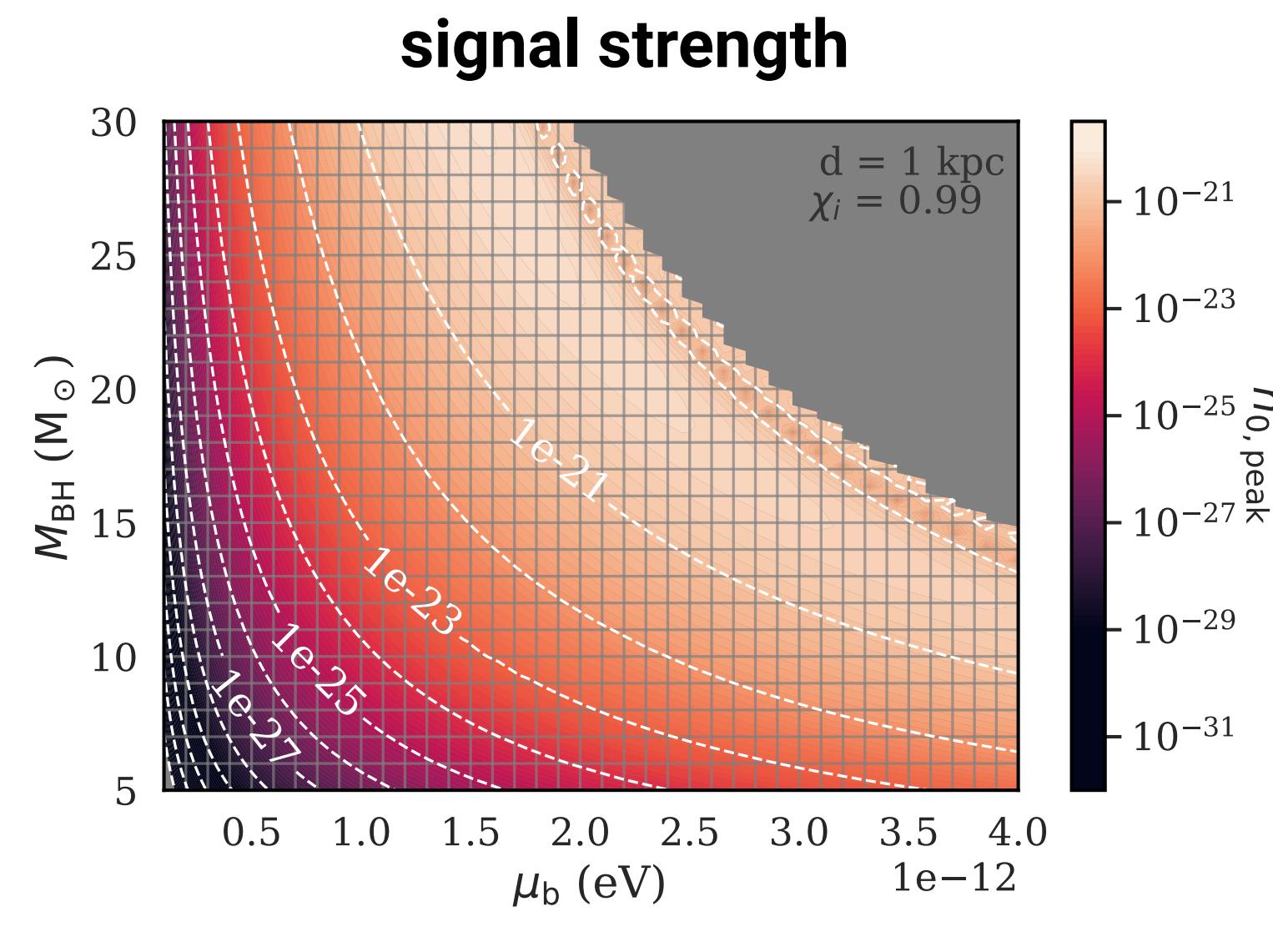


(note: spin-0 particles)

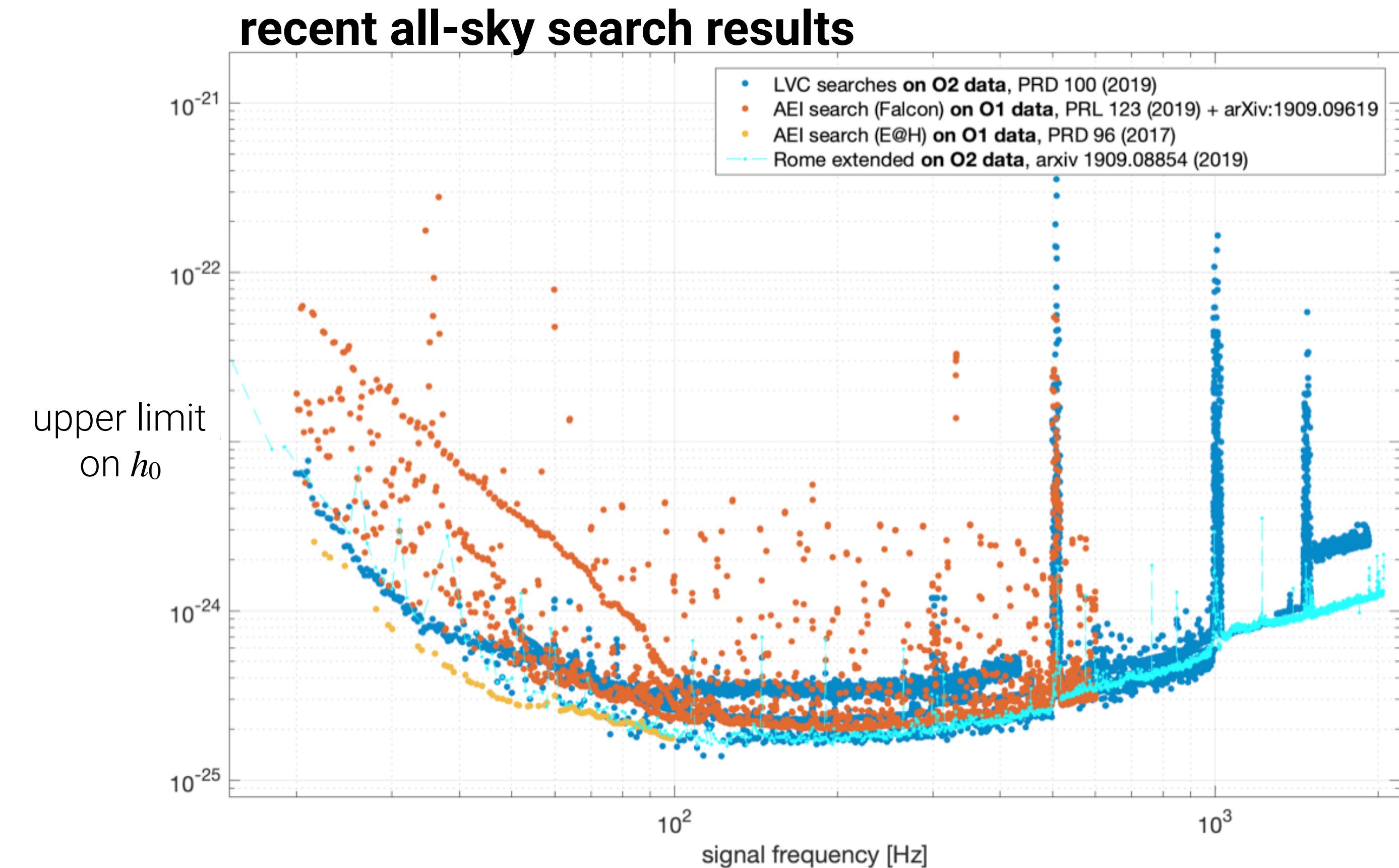


# What kind of signals do we expect?

$\alpha \propto \mu_b M_{\text{BH}}$  determines most of the signal properties



$$h_{0,\text{peak}} \propto \alpha^7$$



# Gravitational “atom” sources continuous gravitational waves

$$\alpha \propto \mu_b M_{\text{BH}}$$

**larger  $\alpha$ :**

(note: spin-0 particles)

$$h_{0,\text{peak}} \propto \alpha^7$$

- louder signals, but

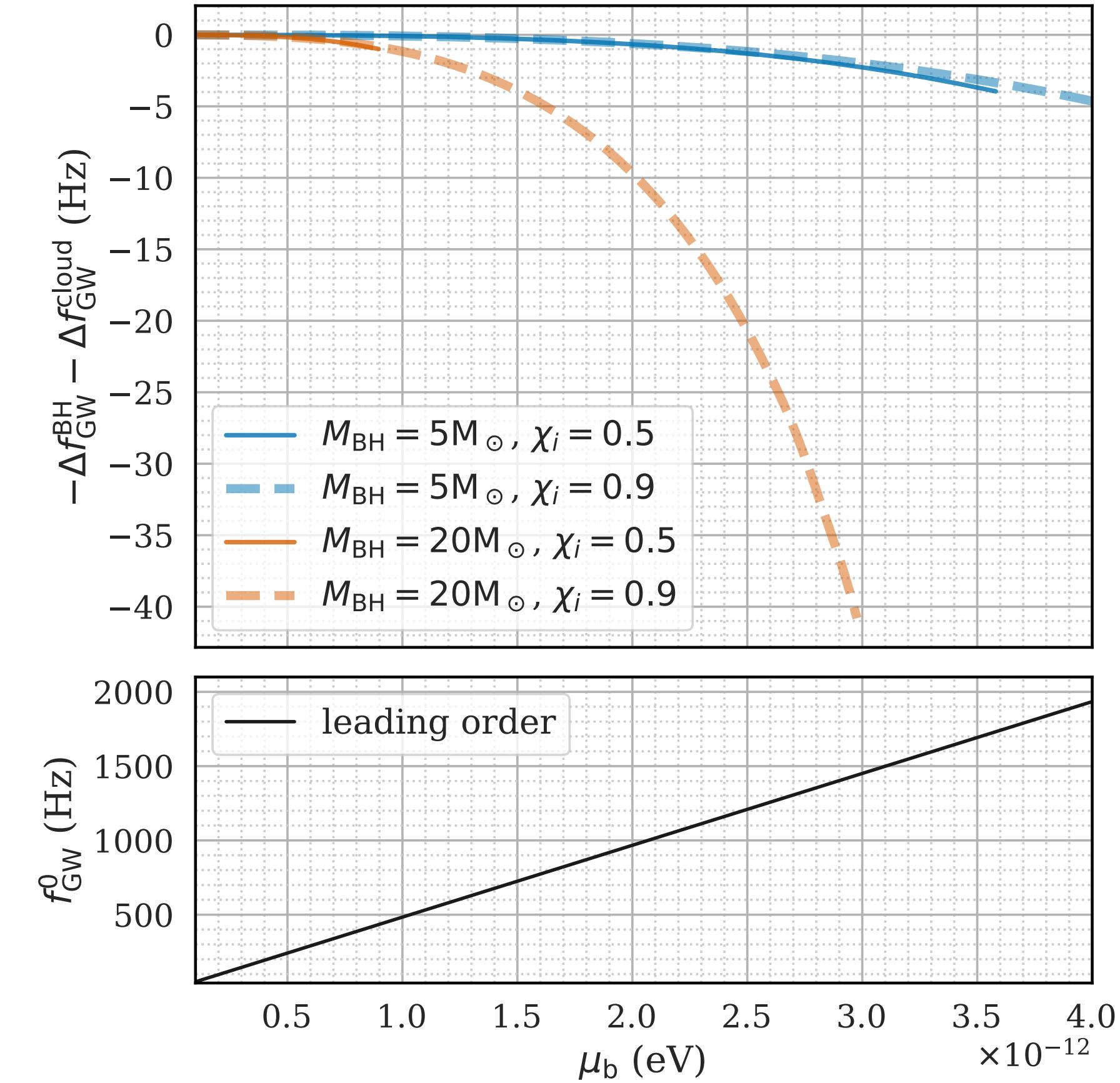
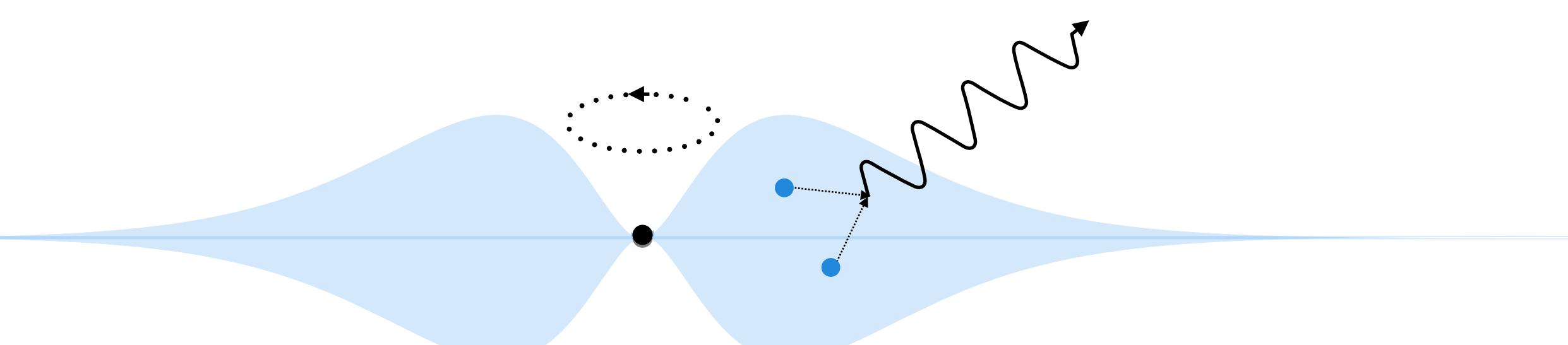
$$\tau_{\text{GW}} \propto \alpha^{-15}$$

- shorter signal timescales

$$\chi_c \approx \frac{4\alpha}{1 + 4\alpha^2}$$

- requires faster spinning BHs

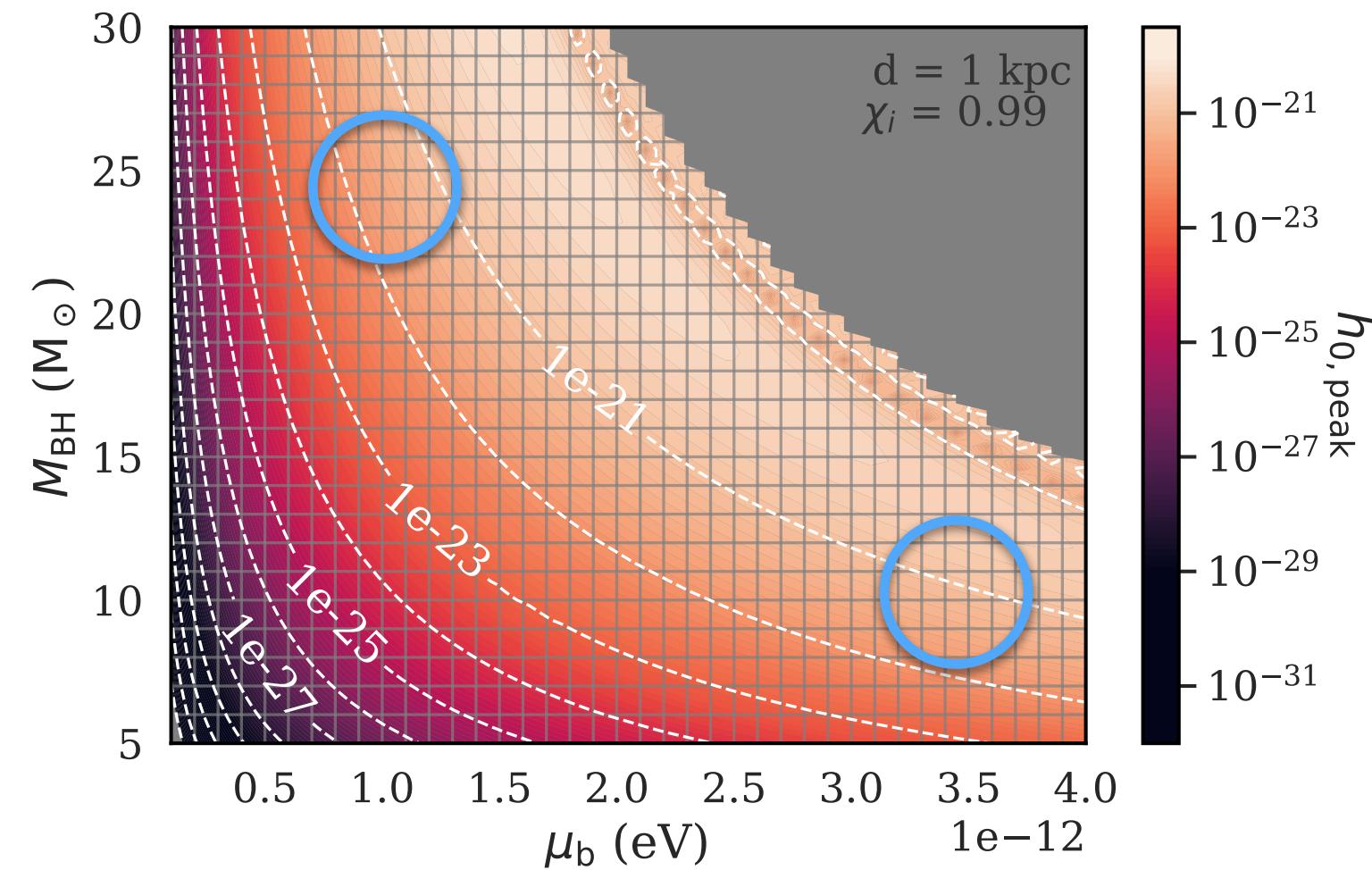
$$f_{\text{GW}} \propto 2\mu_b$$



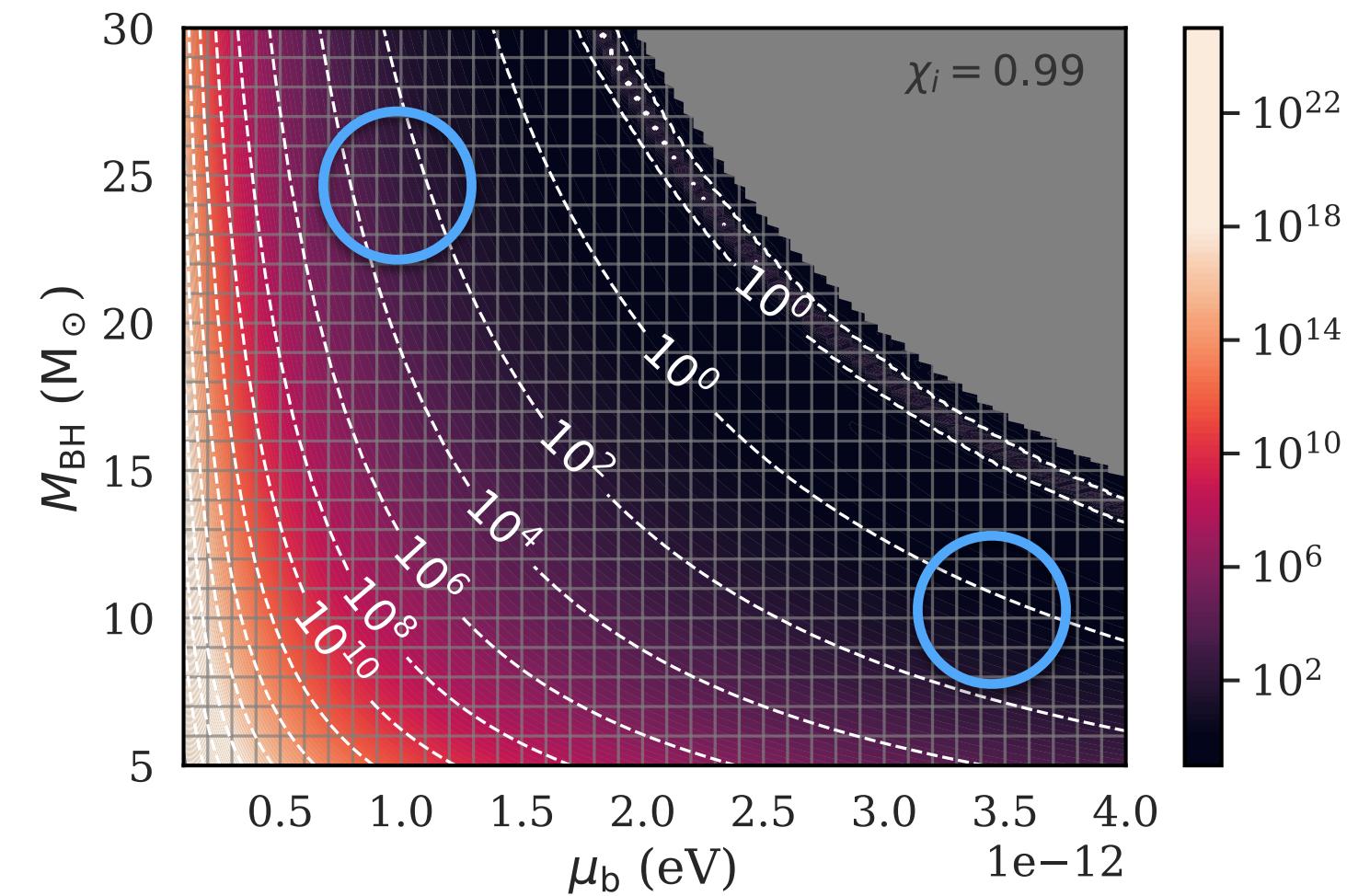
# What kind of signals do we expect?

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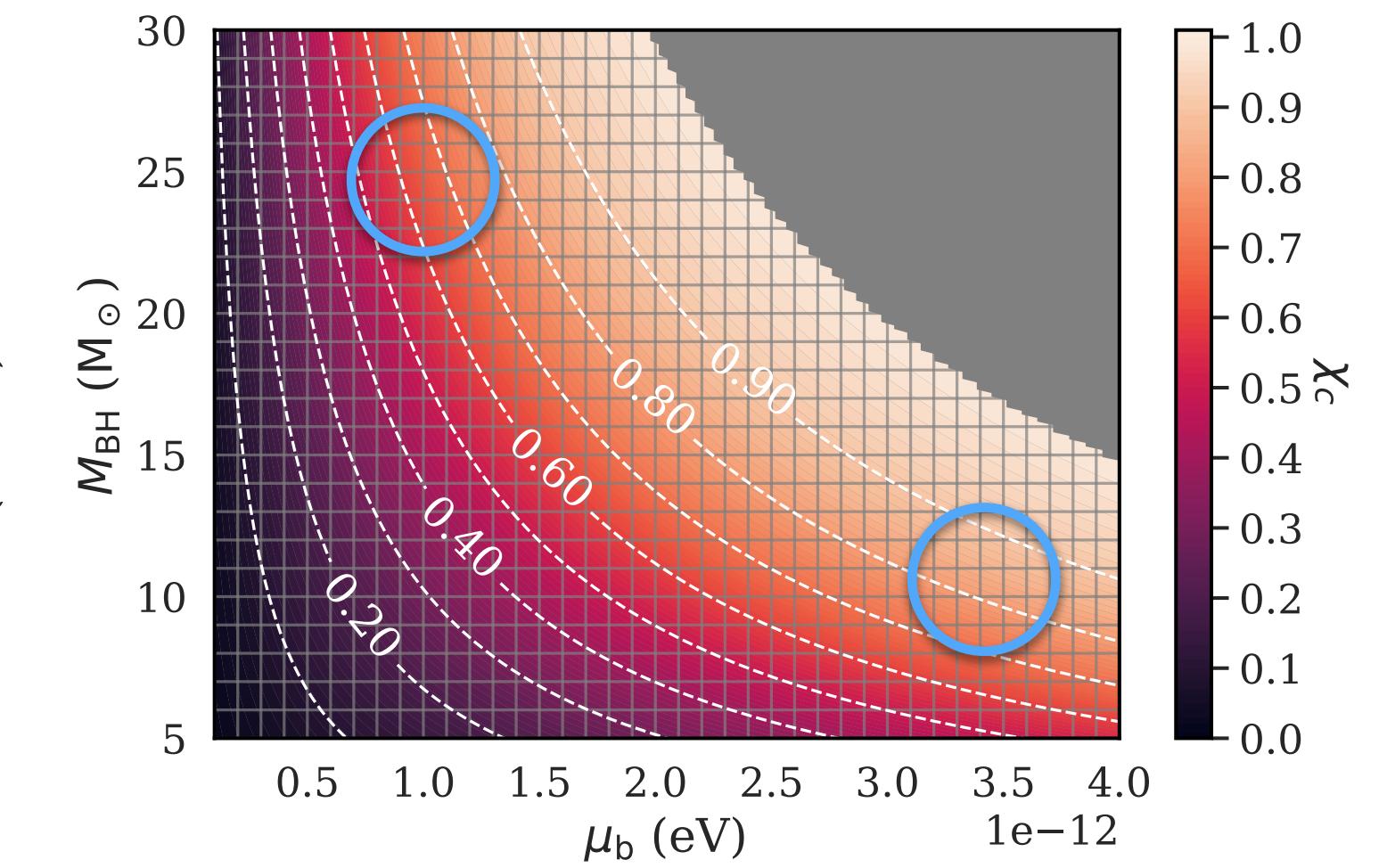
**signal strength**



**signal duration**



**whether the system forms**



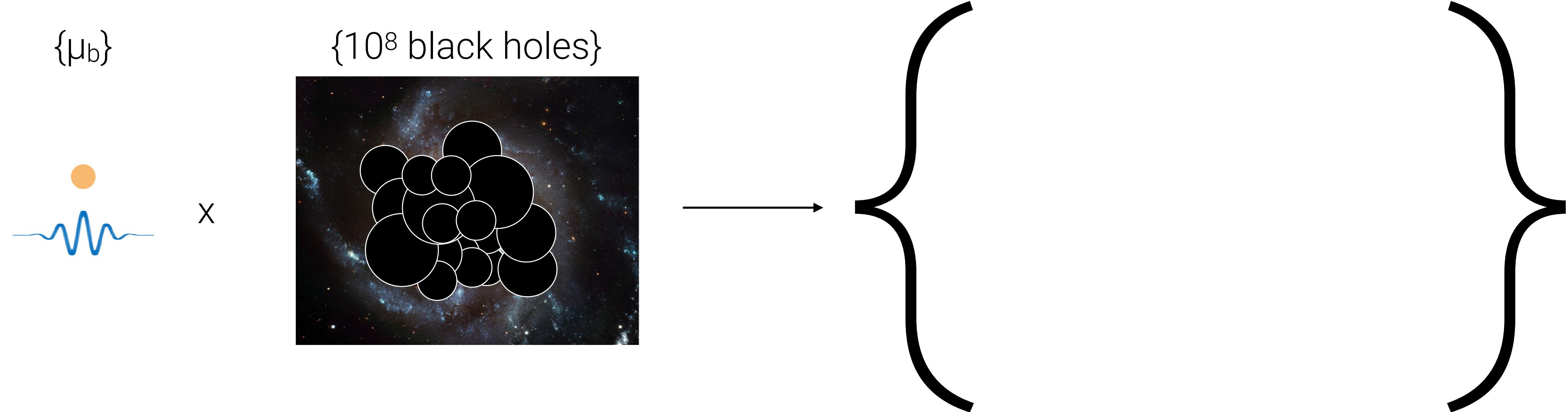
$$h_{0,\text{peak}} \propto \alpha^7$$

$$\tau_{\text{GW}} \propto \alpha^{-15}$$

$$\chi_c \approx \frac{4\alpha}{1 + 4\alpha^2}$$

# Calculate the signal from every black hole

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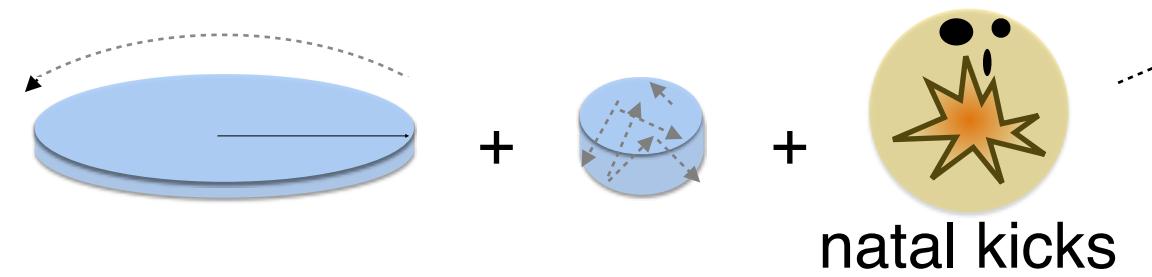
# Simulate $10^8$ isolated black holes in the Galaxy

**galaxy = disk + bulge + halo**

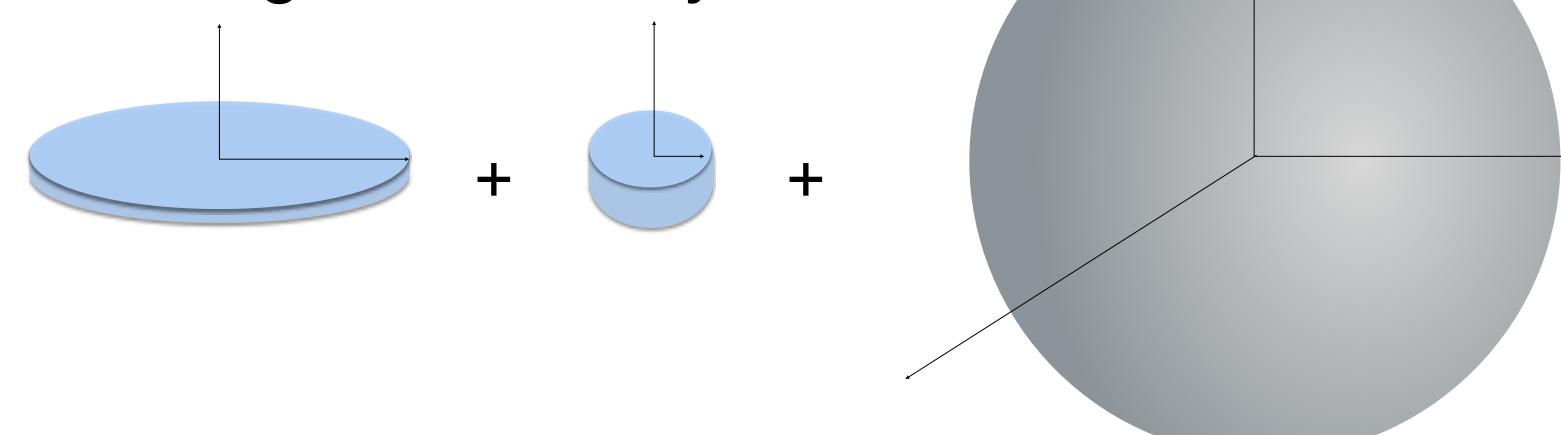
black holes are born:

$$85\% \quad + \quad 15\%$$

with an initial velocity:



and move through the Galaxy



**mass: power-law distribution**

$$M_{\text{BH}} \in [5M_{\odot}, 20M_{\odot}]$$

$$M_{\text{BH}} \in [5M_{\odot}, 30M_{\odot}]$$

**spin: uniform distribution**

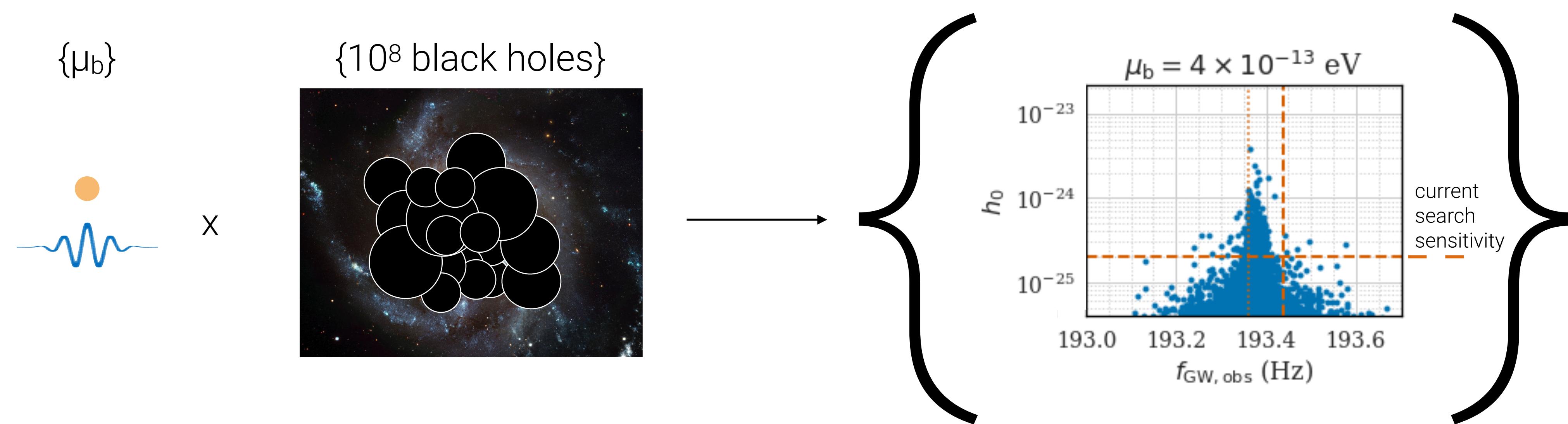
$$\chi_i \in [0, 1]$$

$$\chi_i \in [0, 0.5]$$

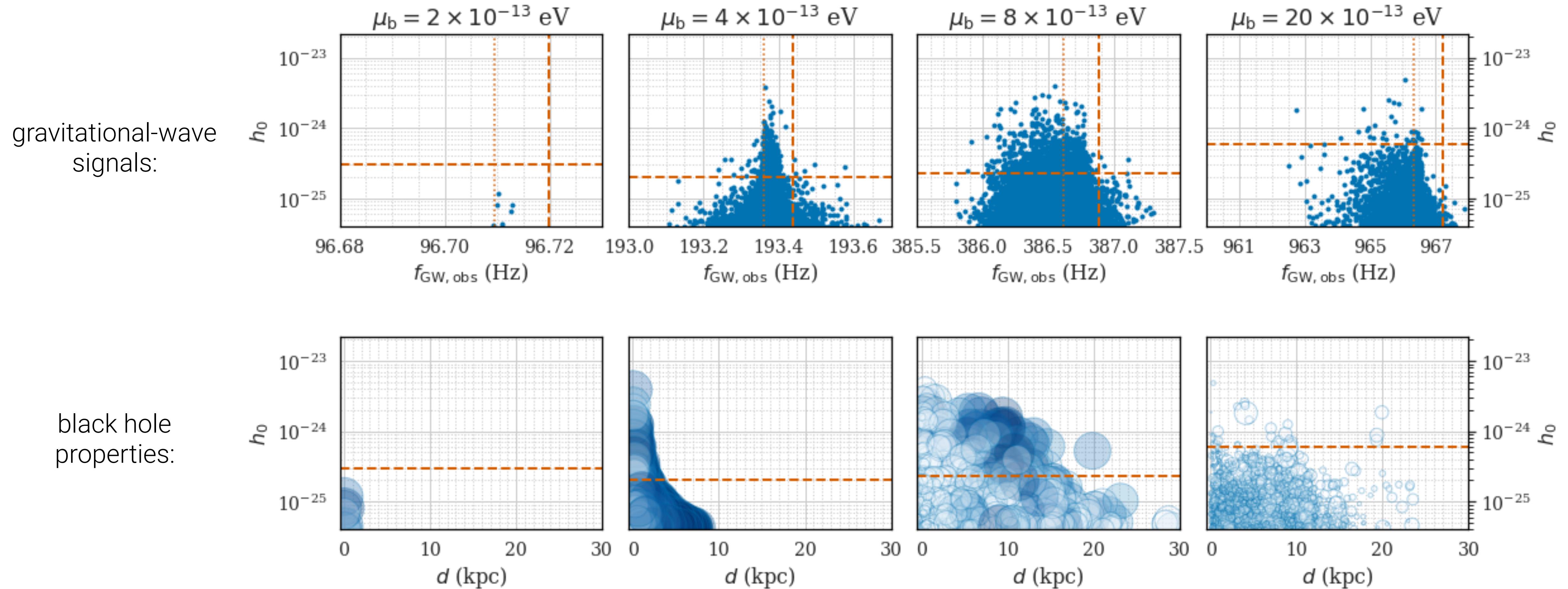
$$\chi_i \in [0, 0.3]$$

[Tsuna, Kawanaka, and Totani, MNRAS 477 (2018)]

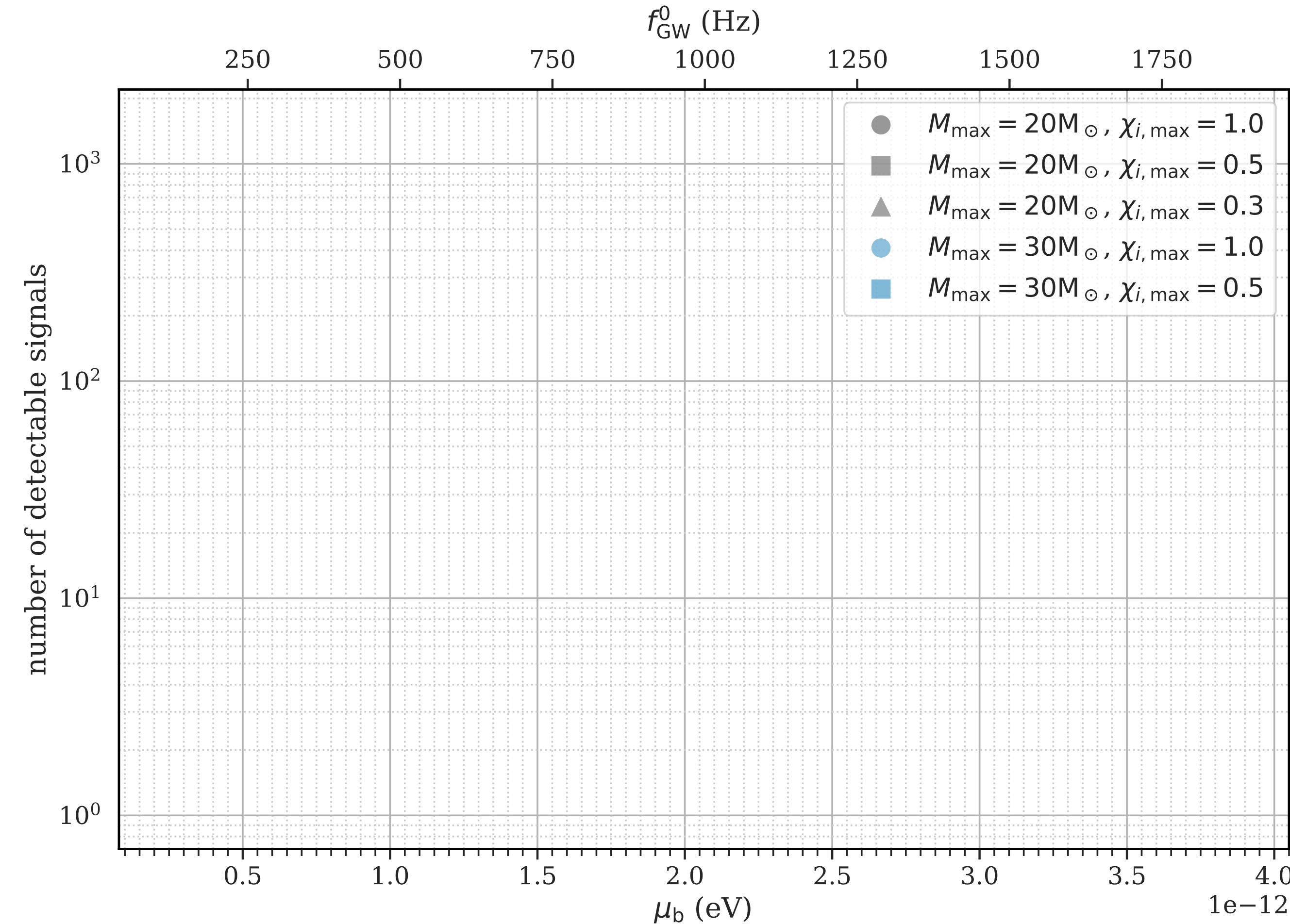
# Calculate the signal from every black hole



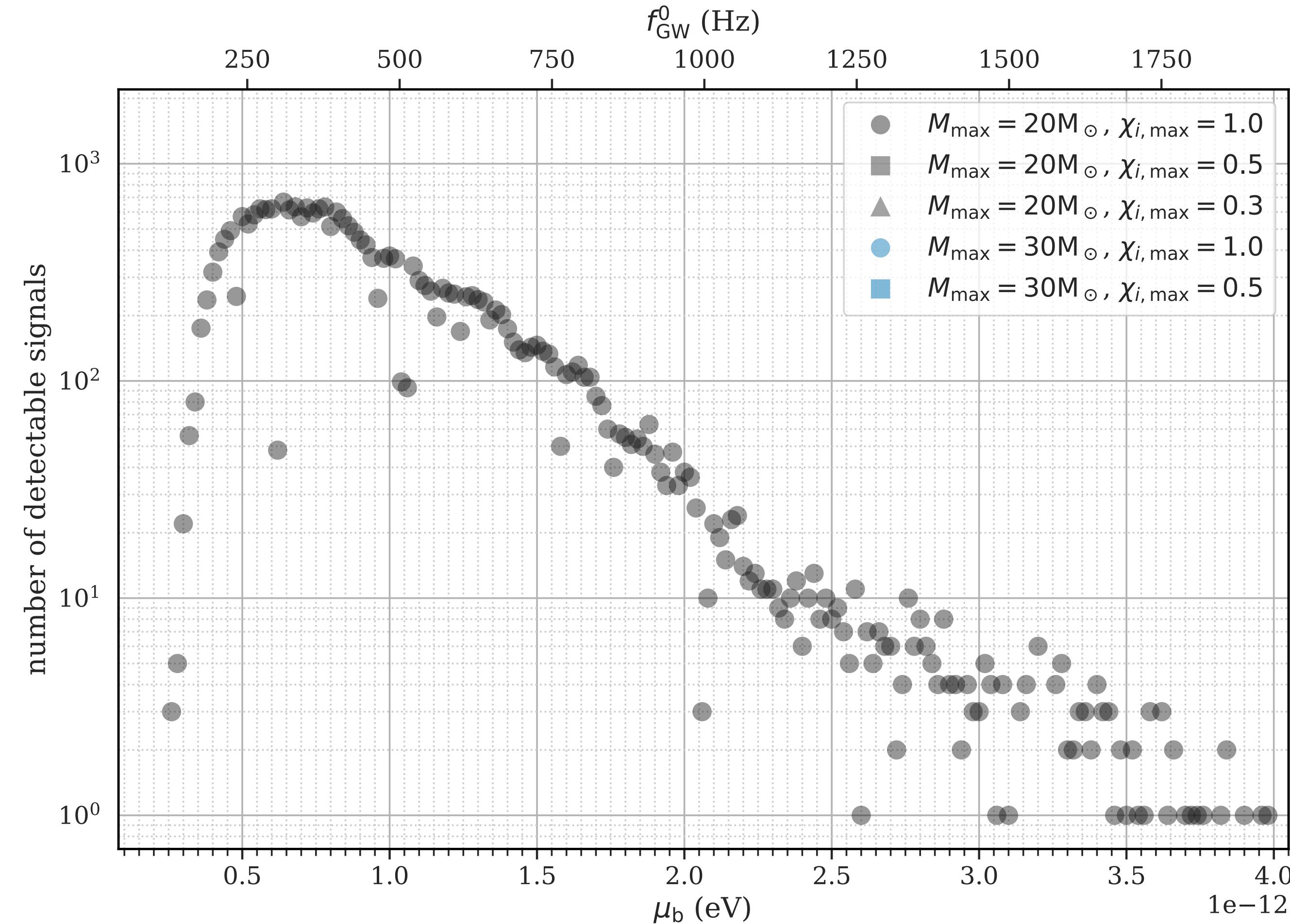
# What do the ensemble signals look like?



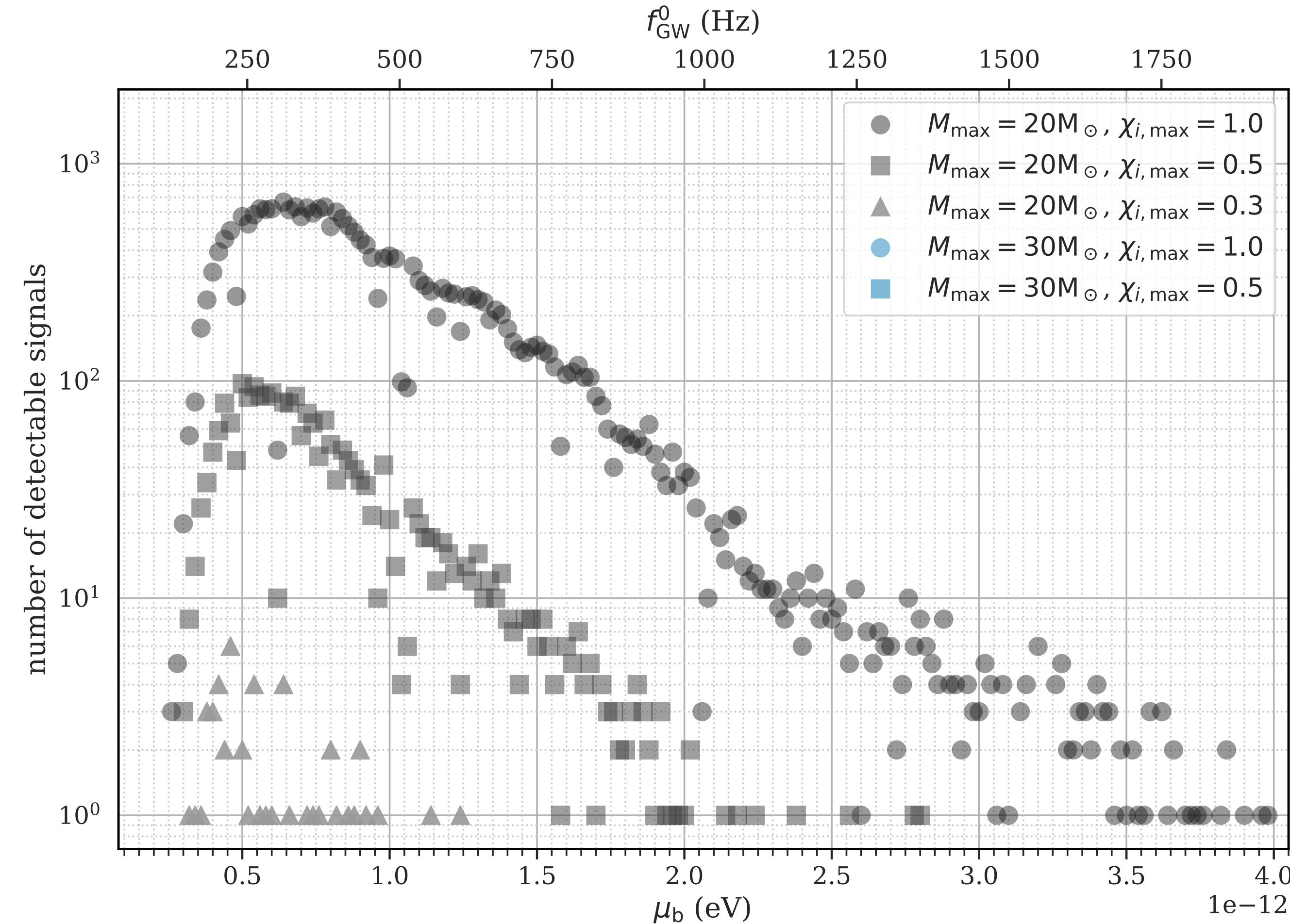
# What are the current detection prospects?



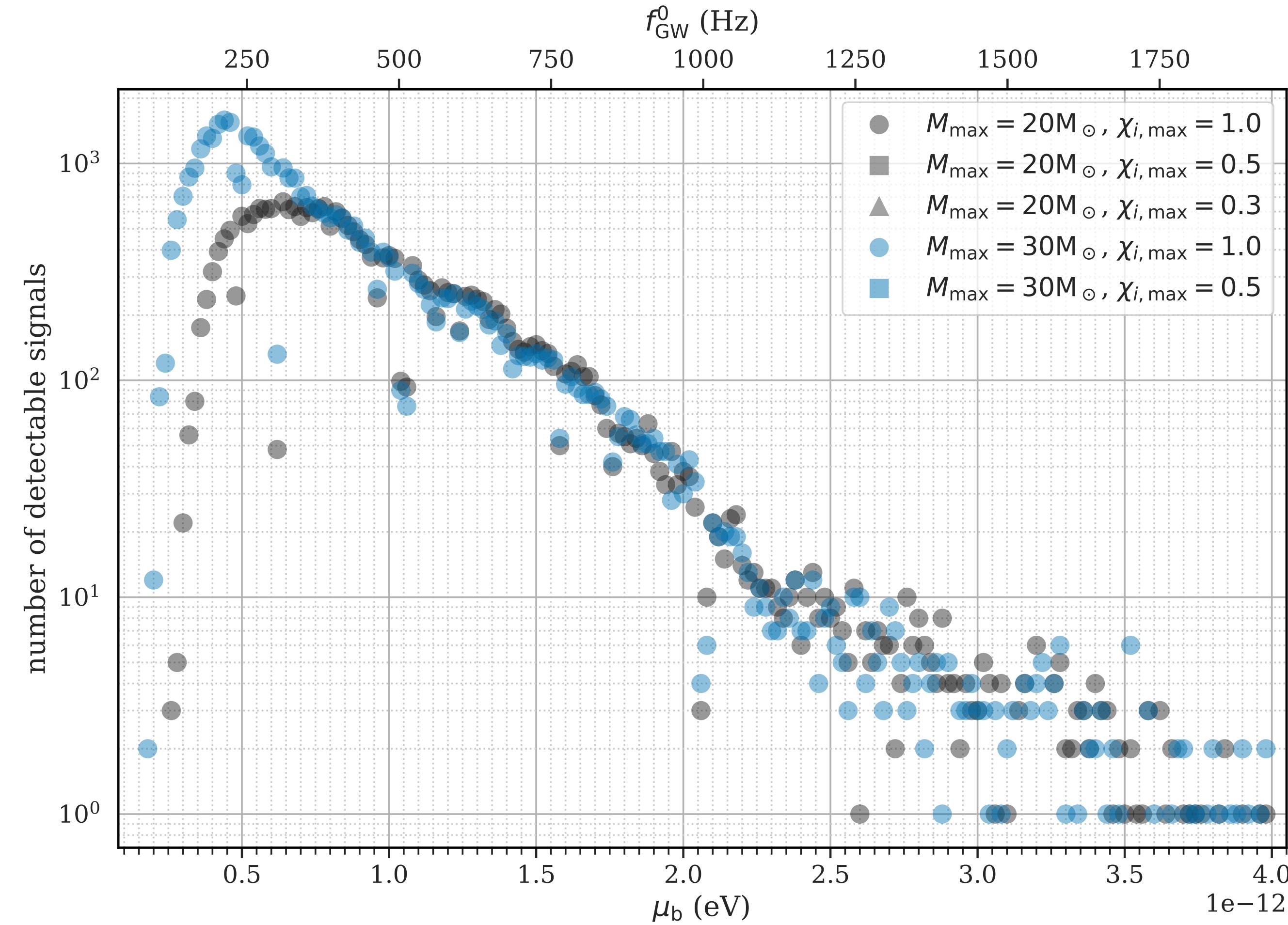
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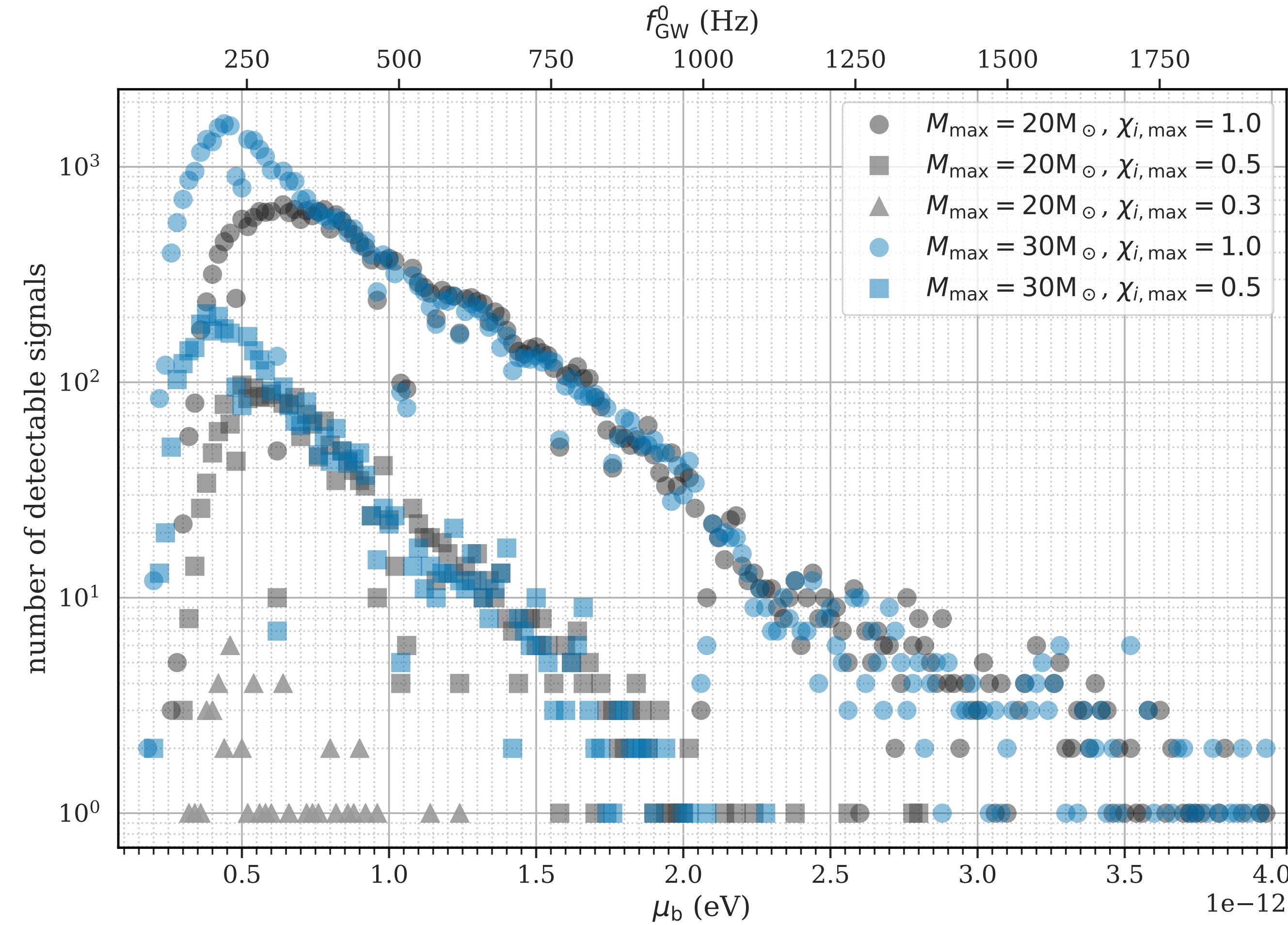
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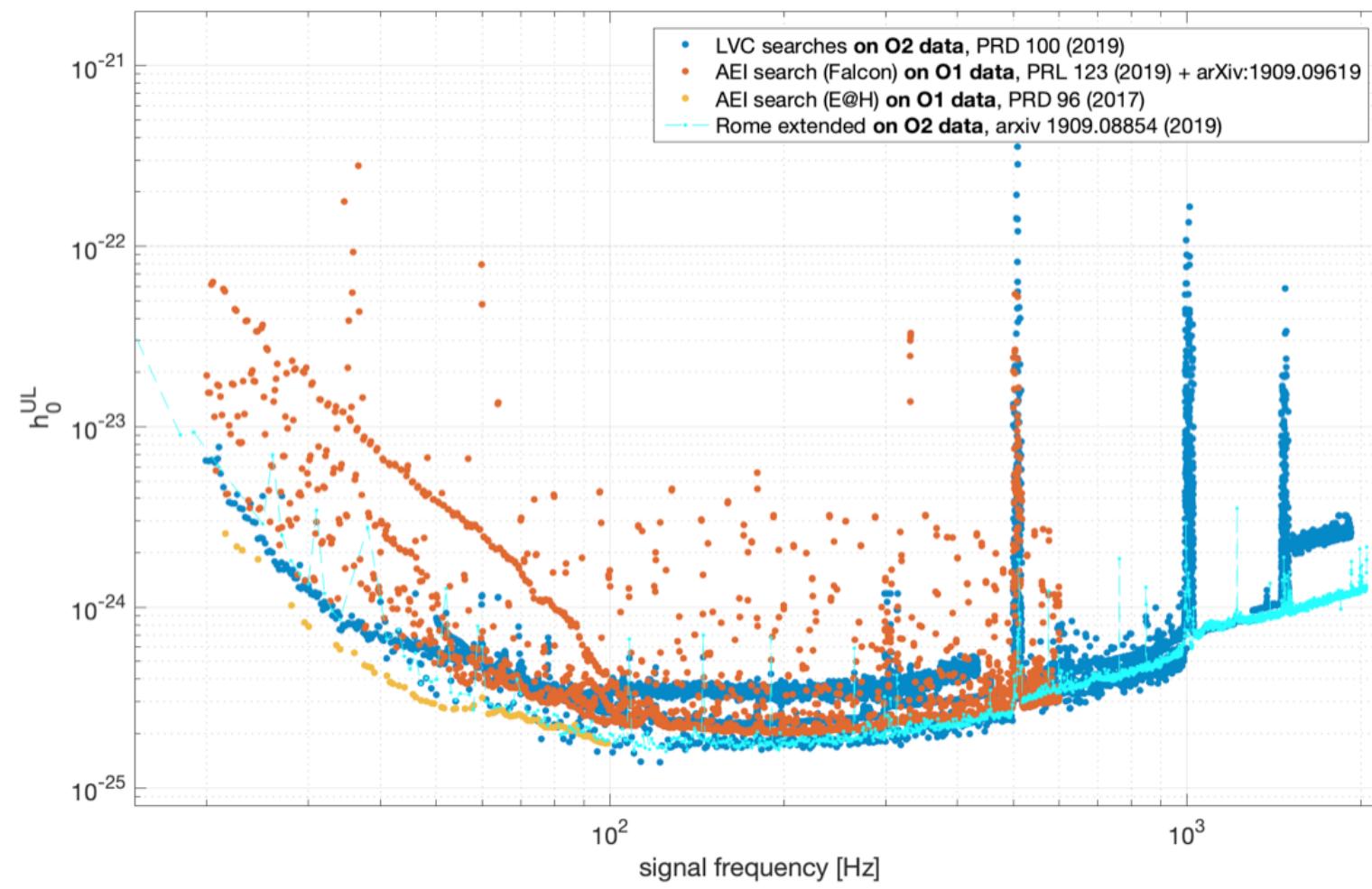


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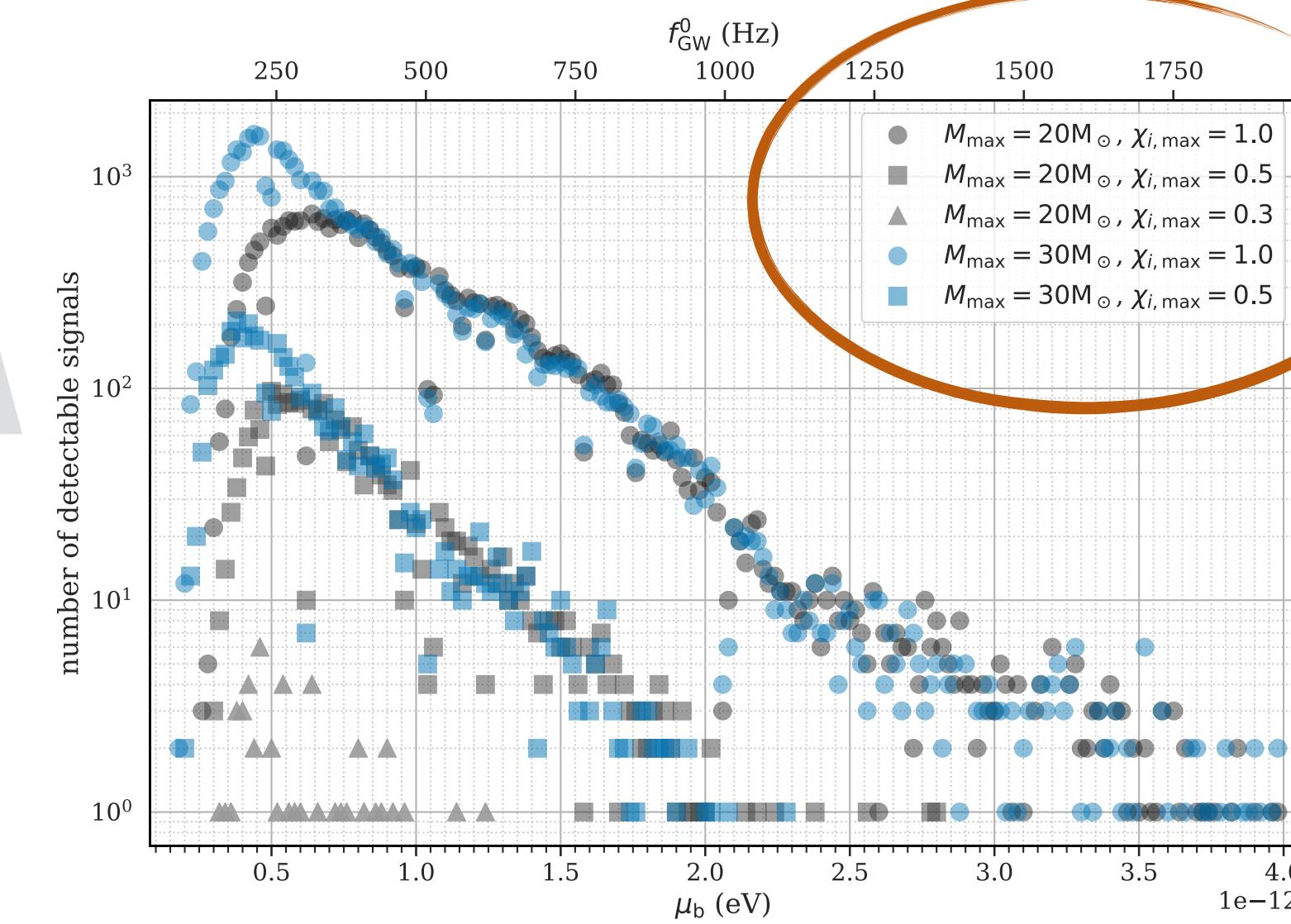


# Interpreting null results (upper limits) – boson clouds

upper limit on gravitational-wave signal strength



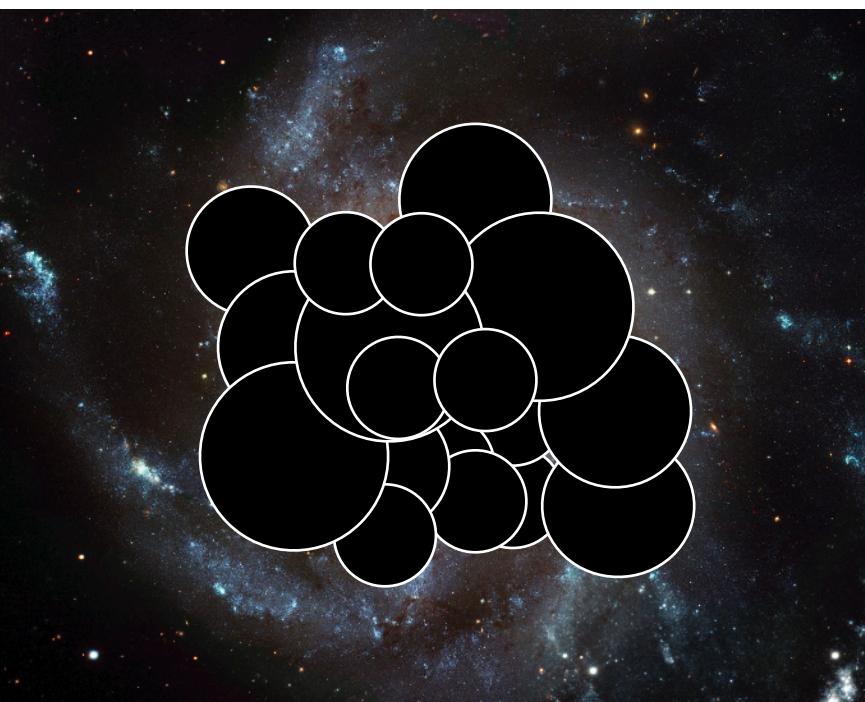
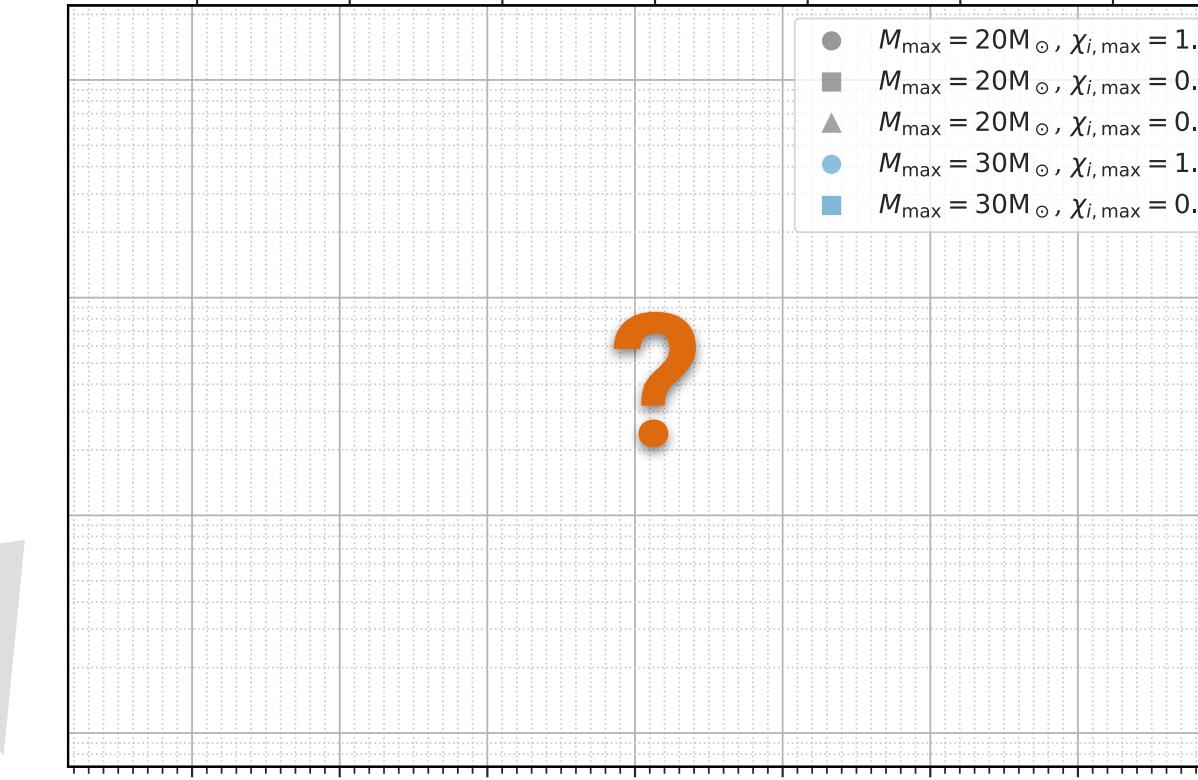
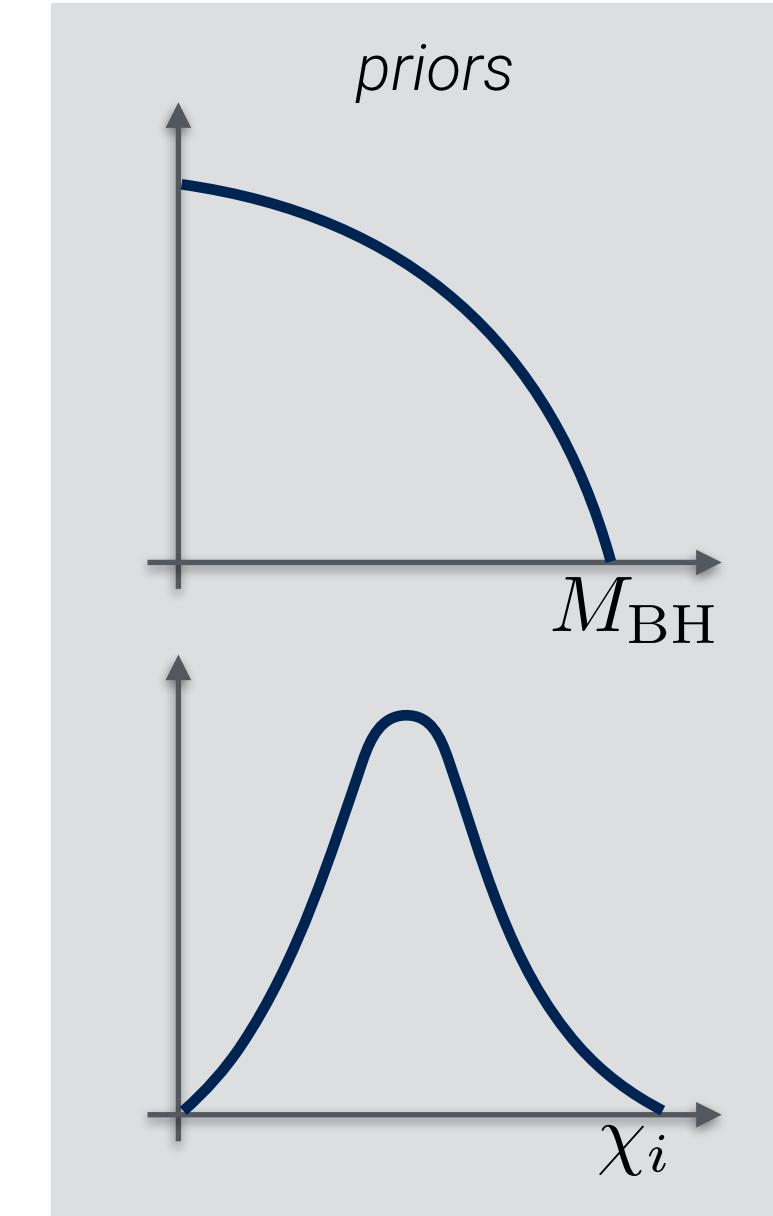
expectations on number of detectable signals



range of disfavored boson masses

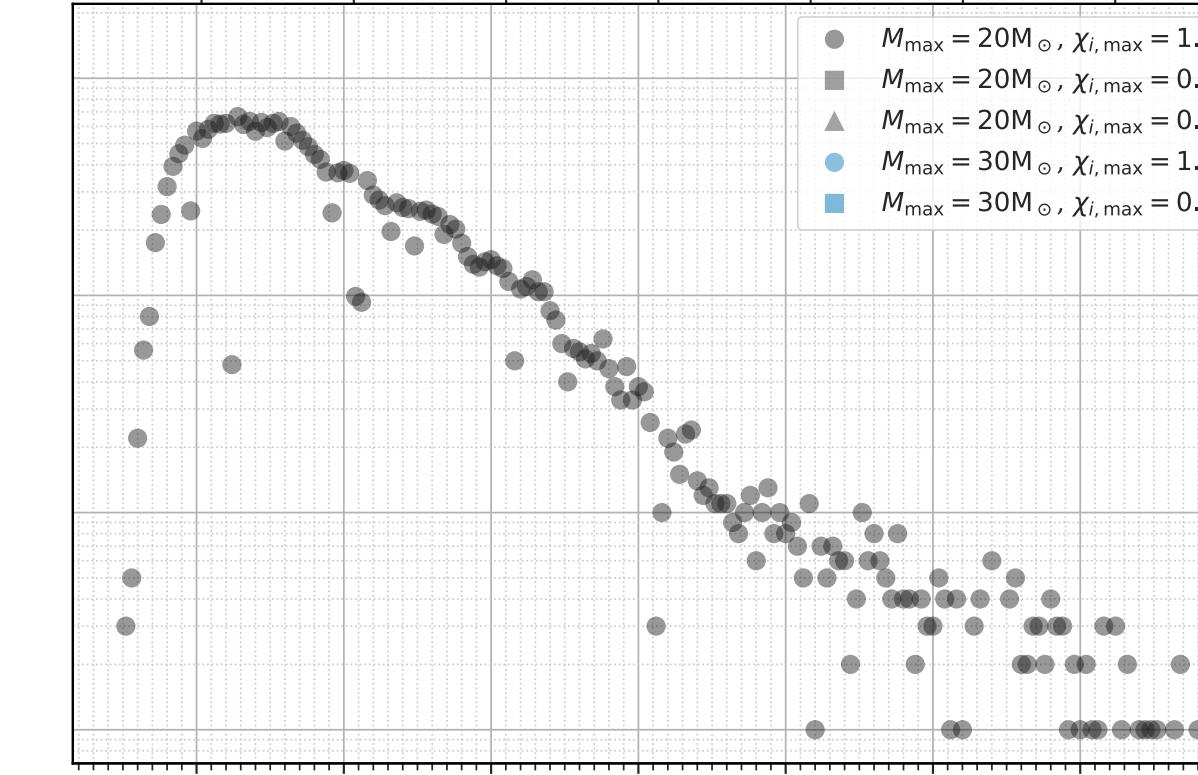
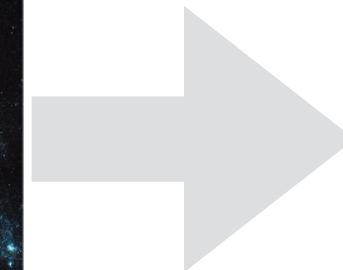
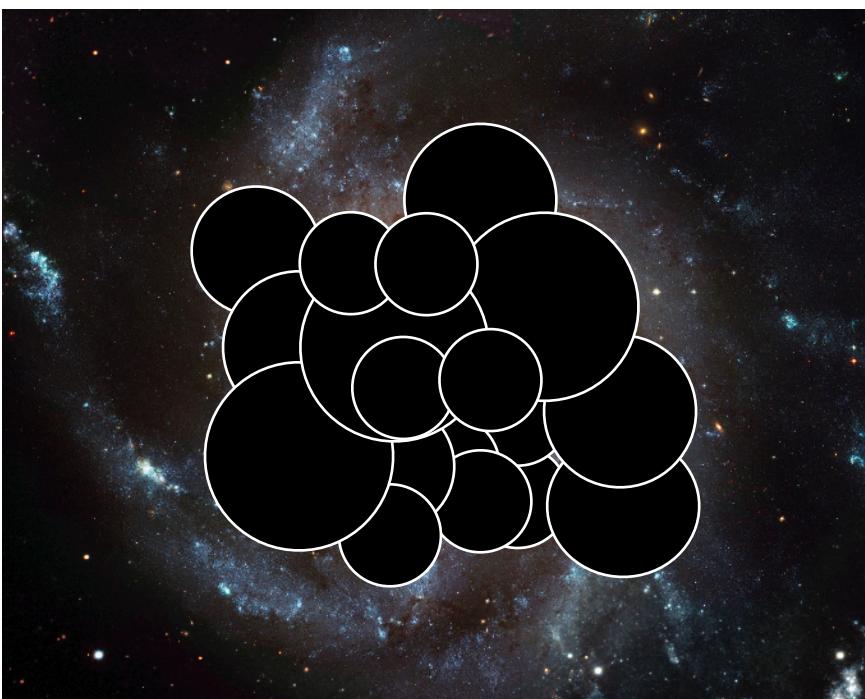
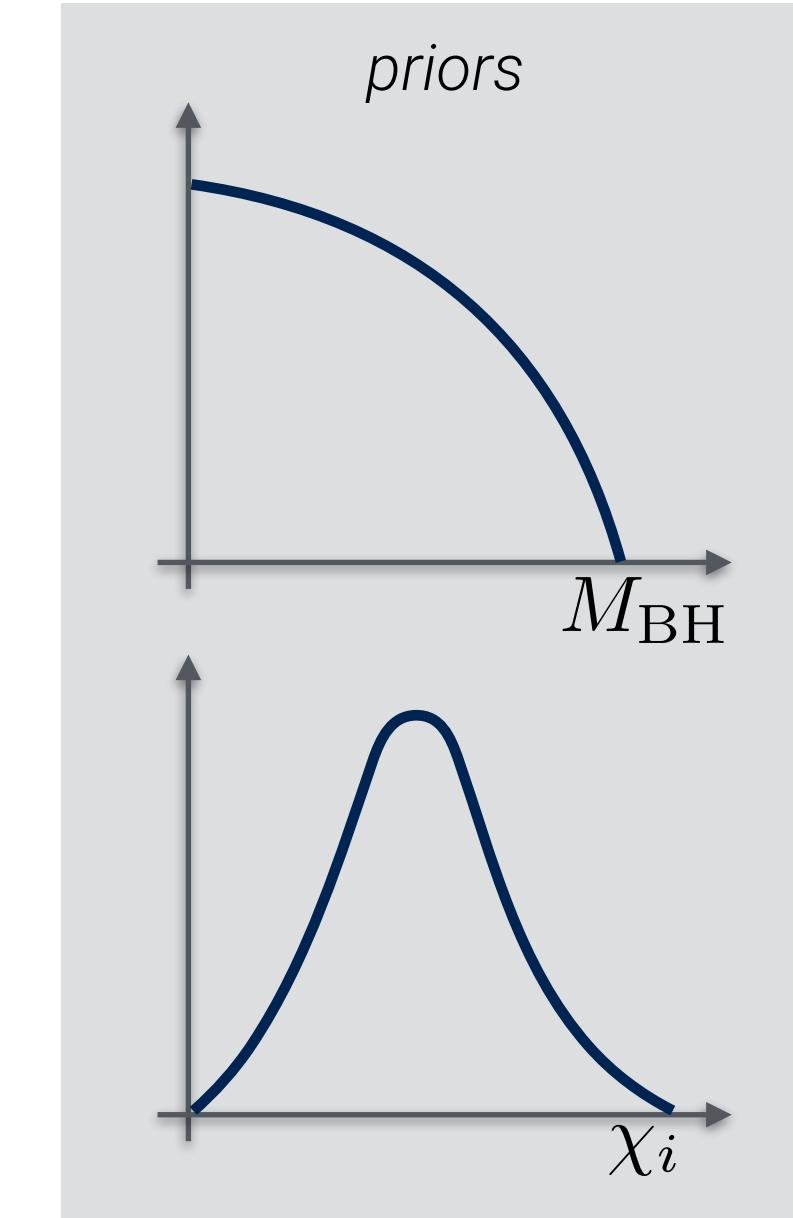
# Interpreting null results (upper limits) – boson clouds

(speculating)

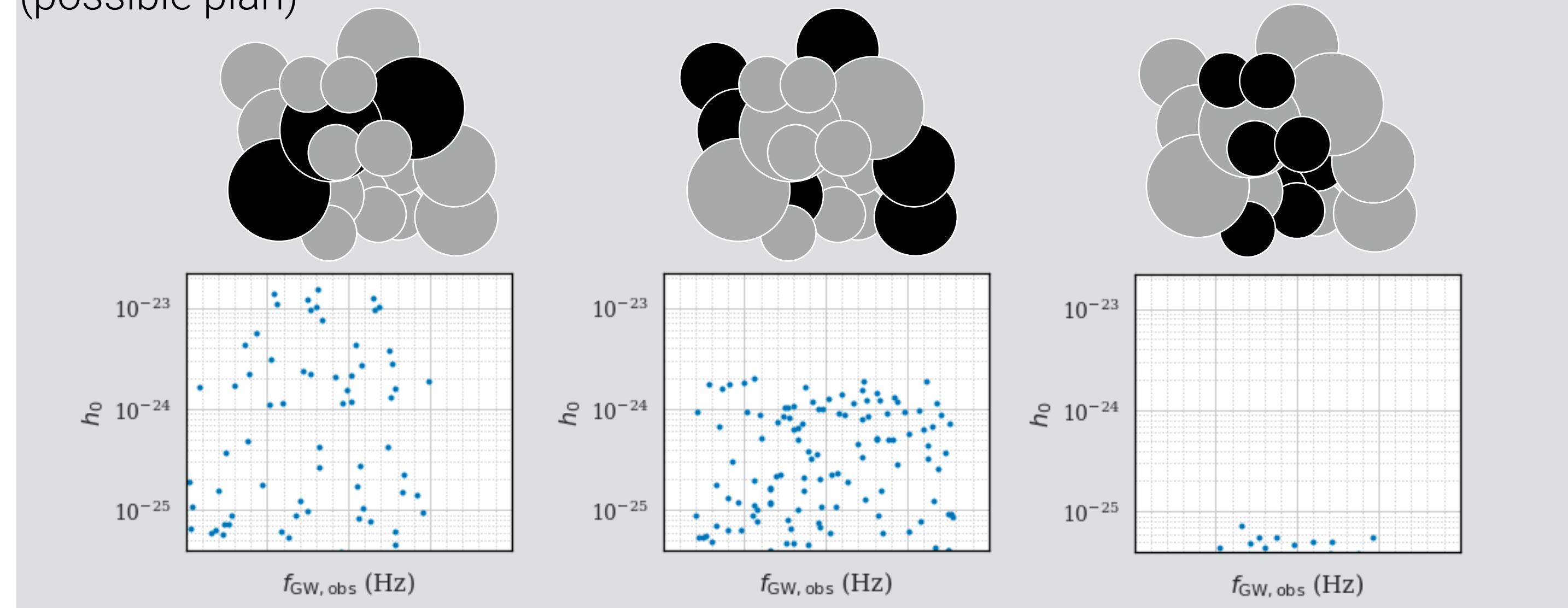


# Interpreting null results (upper limits) – boson clouds

(speculating)



(possible plan)



# Thanks

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[arXiv:2003.03359](https://arxiv.org/abs/0303.03359)

<https://www.aei.mpg.de/continuouswaves/arxiv200303359>

# extra slides

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# Recent searches using GWs

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Some recent searches (non exhaustive list!!)

- Directed search for signal from a single BH: [Isi et al., PRD 99 \(2019\)](#); [Ling et al., PRD 101 \(2020\)](#)
- Search for the stochastic background: [Tsukada et al., PRD 99 \(2019\)](#)
- All-sky search: [Palomba et al., PRL 123 \(2019\)](#)
- Indirect constraints from spinning black holes: [Ng et al., arXiv:1908.02312](#)

# The ensemble signal

2 bosons annihilate  $\rightarrow$  gravitational waves

black holes are moving towards or away from us

all\* black holes will have radiating clouds  
(i.e., all black holes with the right properties)

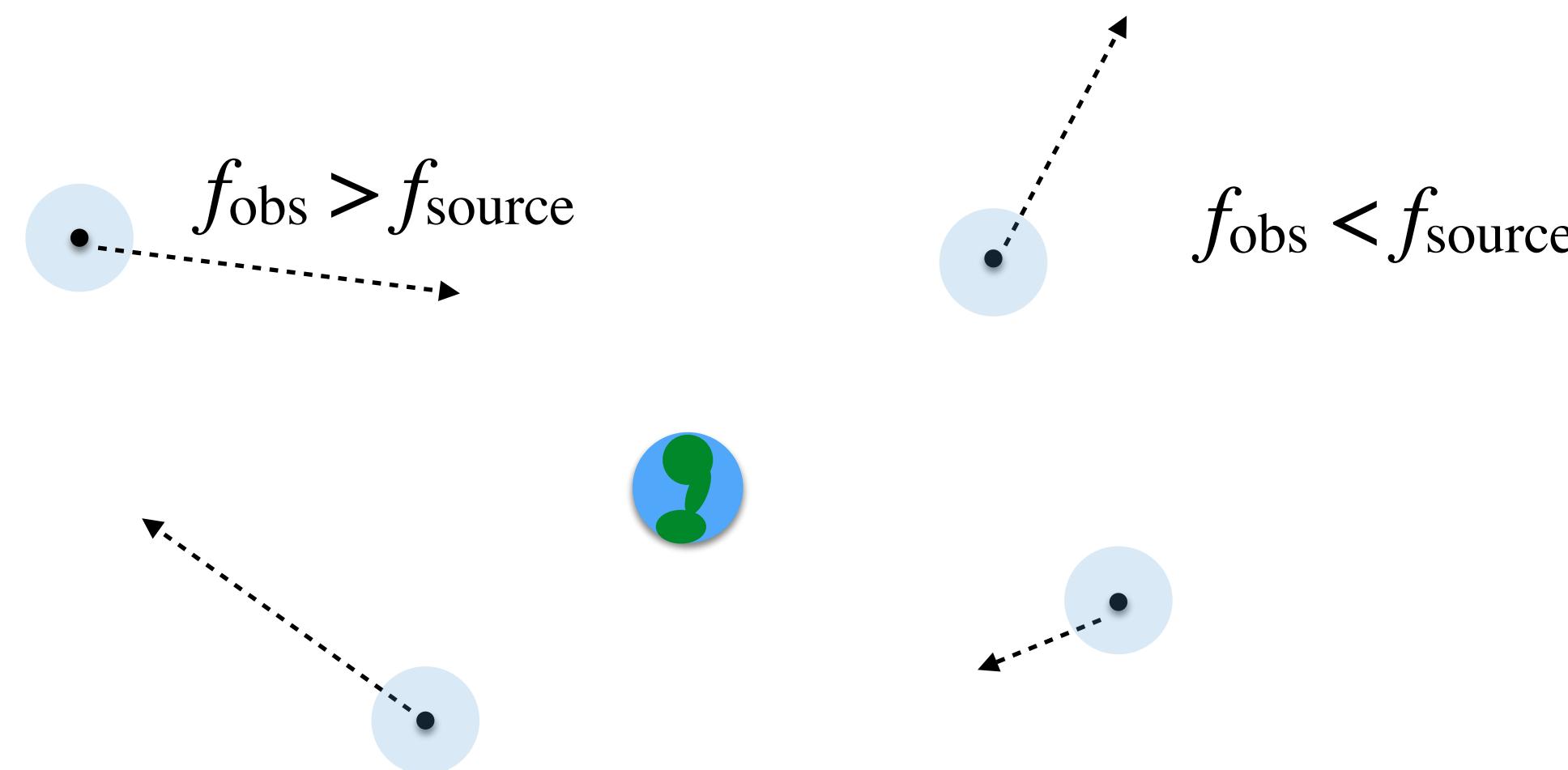
**inherent signal frequency** depends almost entirely on **boson mass**

$$f_{\text{source}}$$

**observed signal frequency** will be **Doppler shifted** due to this motion

$$f_{\text{obs}}$$

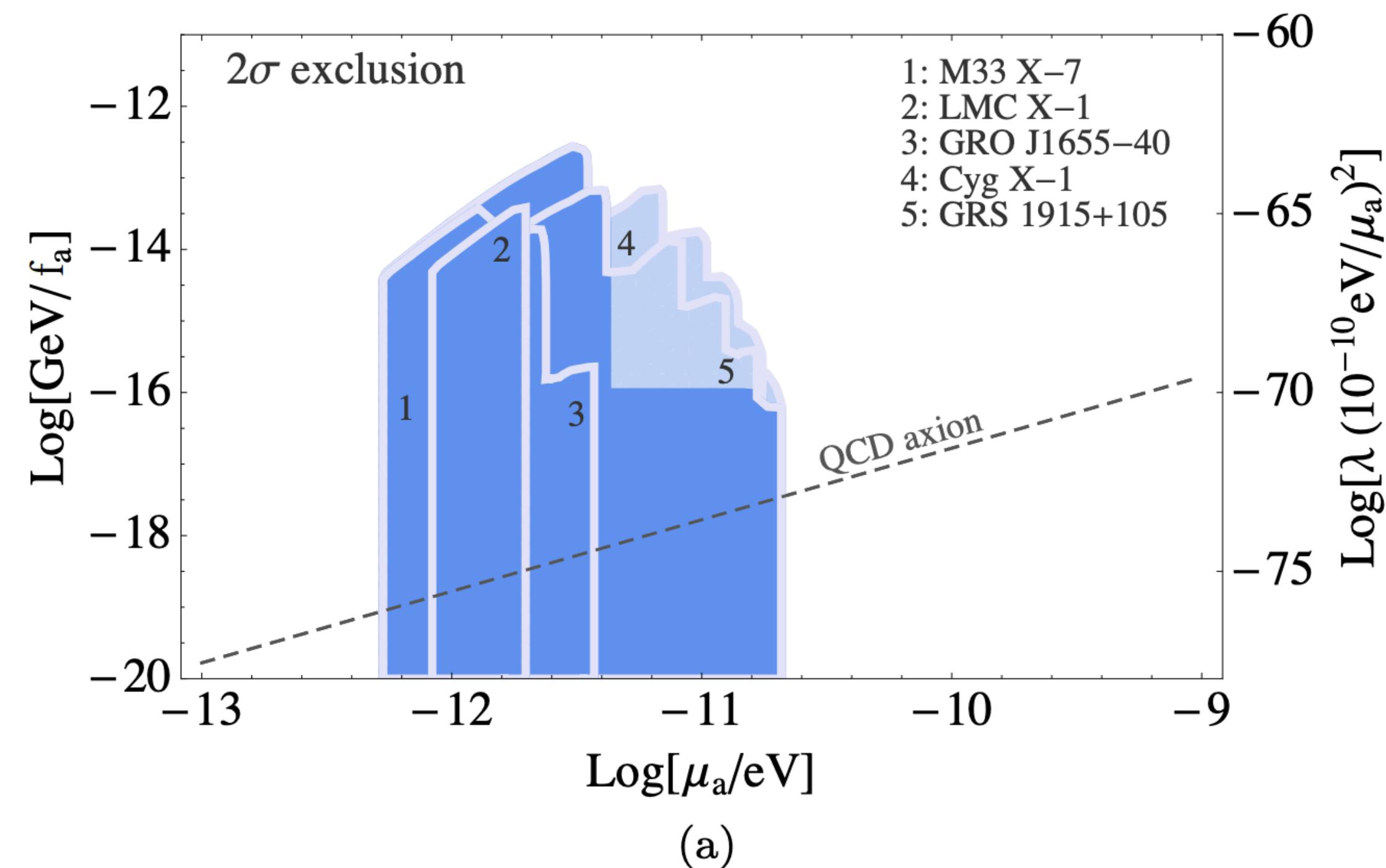
we should be looking for the **ensemble signal** produced by all the boson clouds near us



# Mass and self coupling

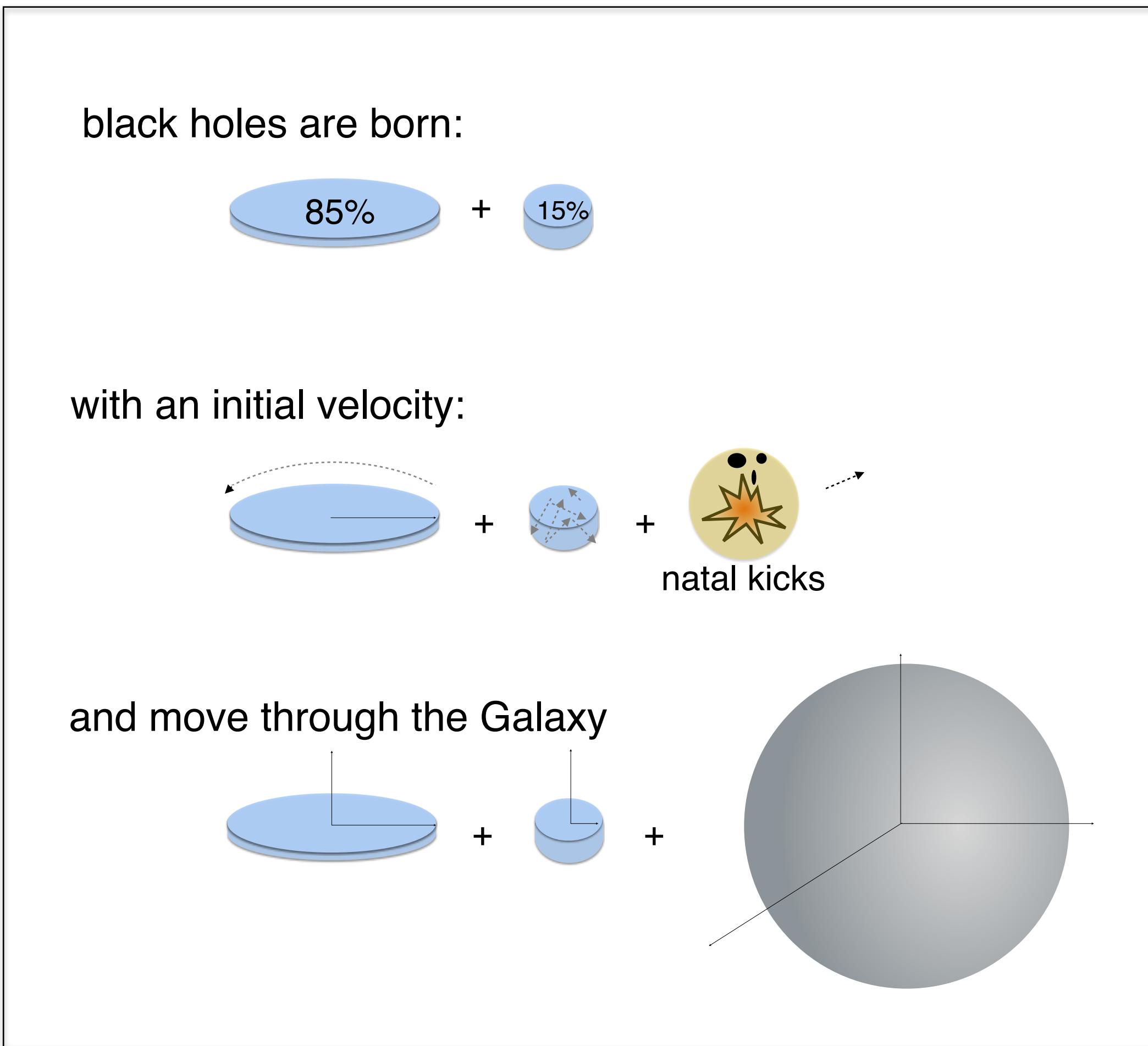
We are interested in particles for which the self coupling is not so strong that cloud collapses as a bosenova.

e.g., the regions here are determined based on observations of spinning black holes

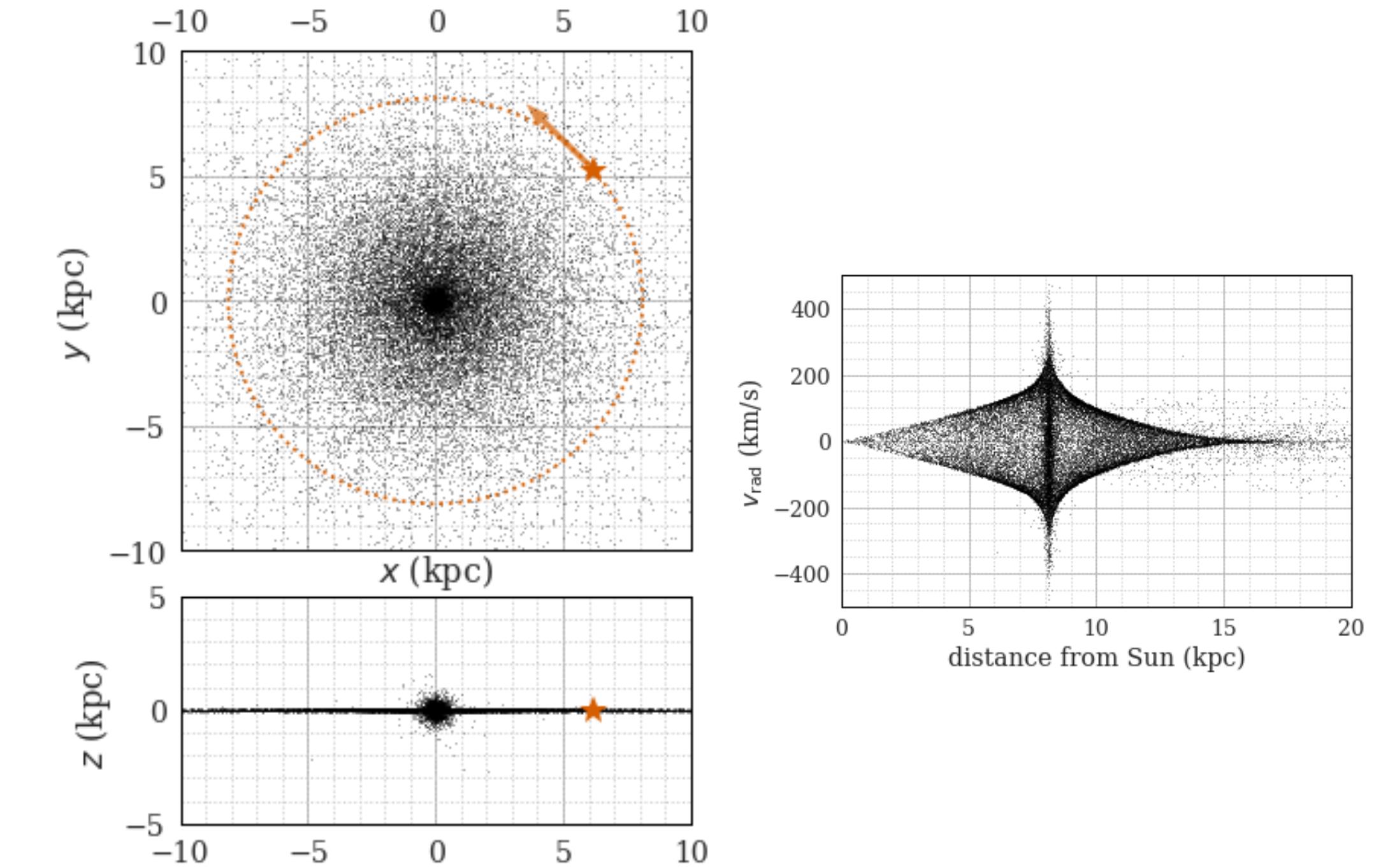


# $10^8$ isolated black holes in the Galaxy

galaxy = disk + bulge + halo (see D. Tsuna et al. 2018)

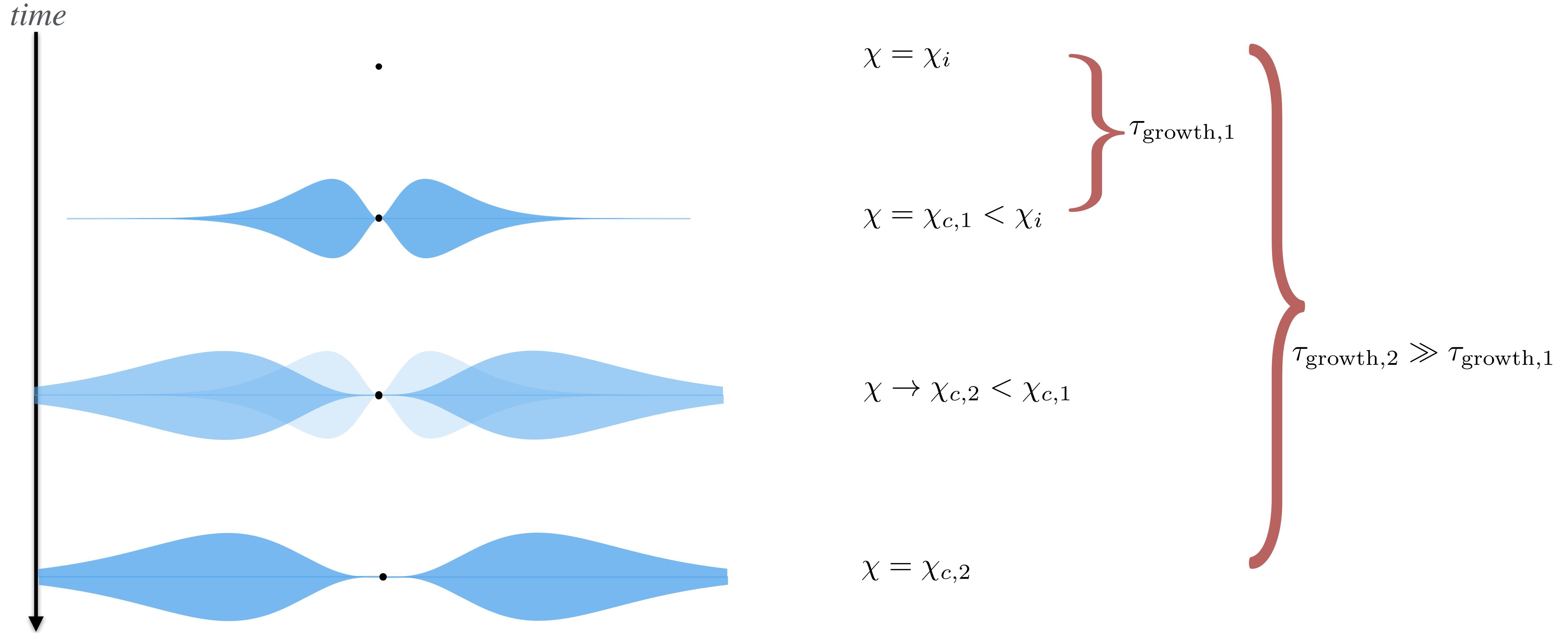


final black hole spatial properties

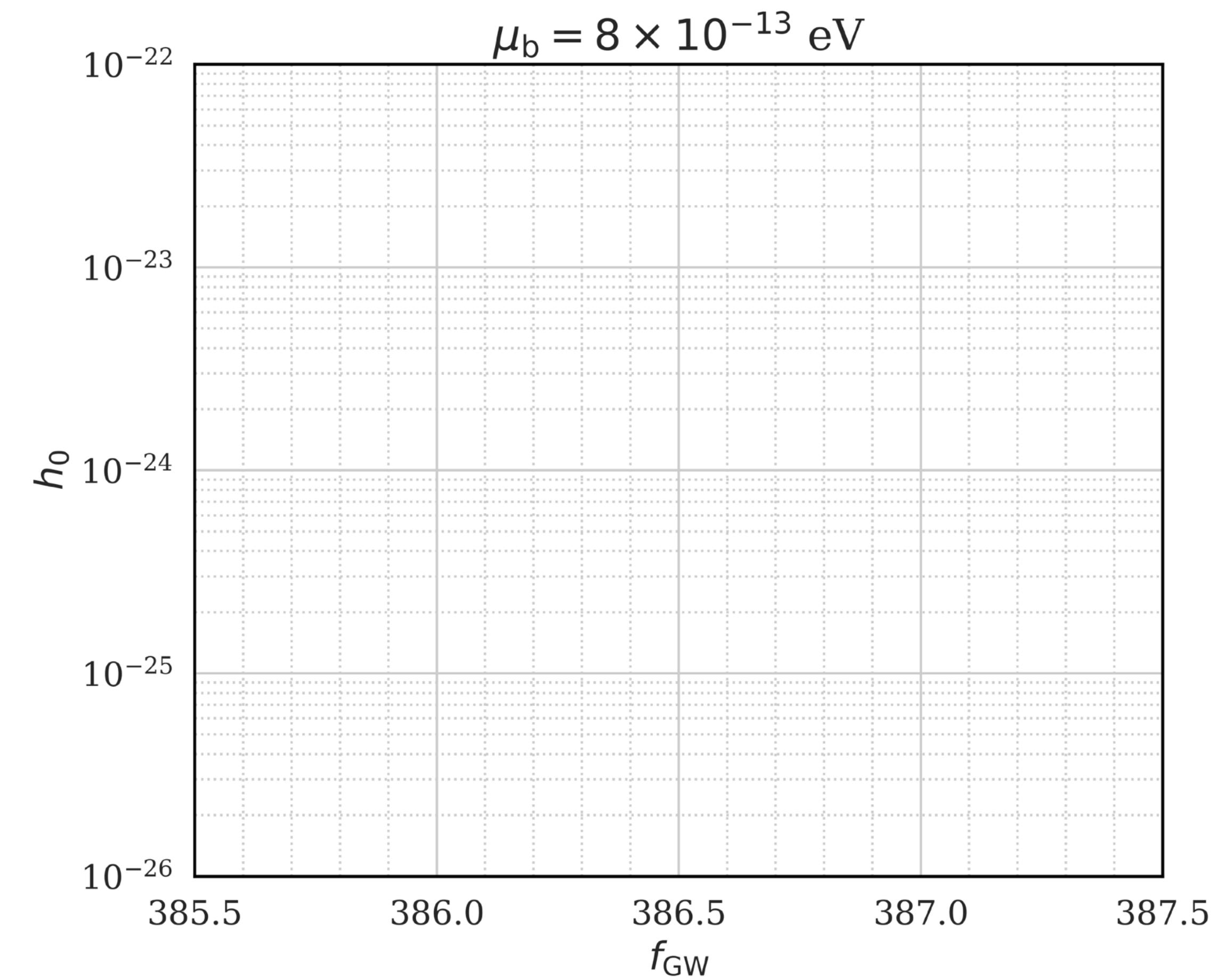
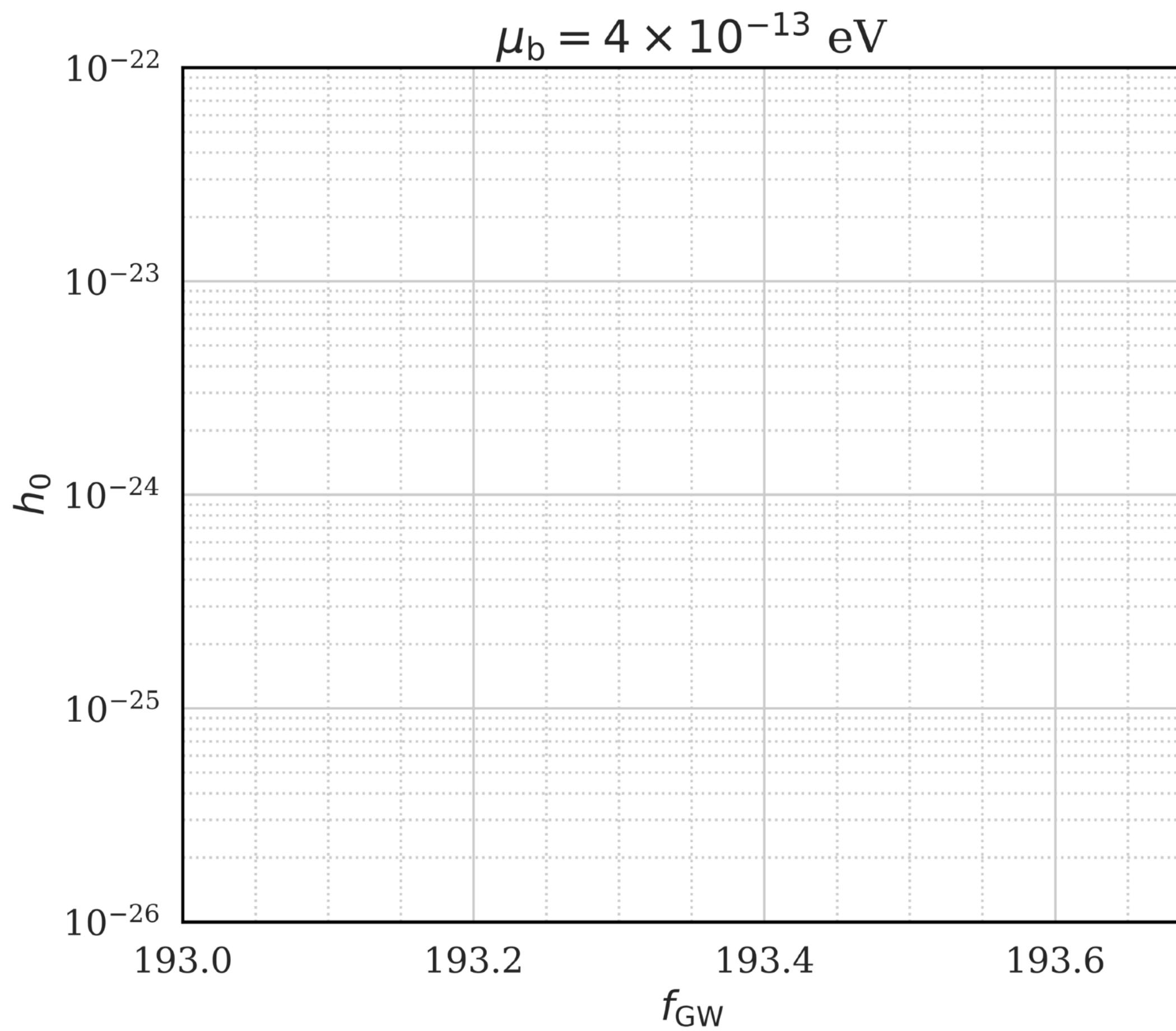


# formation of the second cloud level

When the second level fully forms, the first level falls back into the black hole and emission ceases.

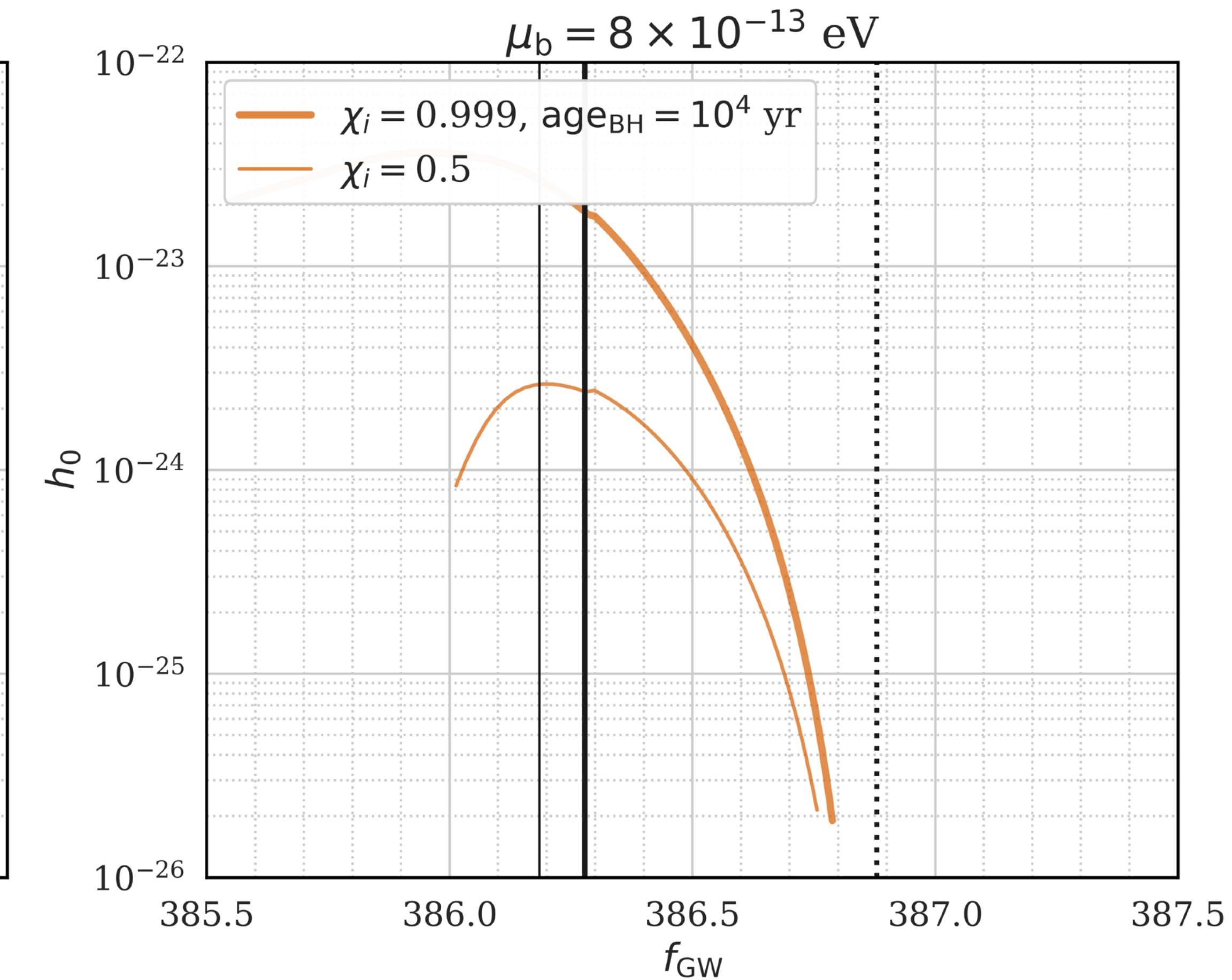
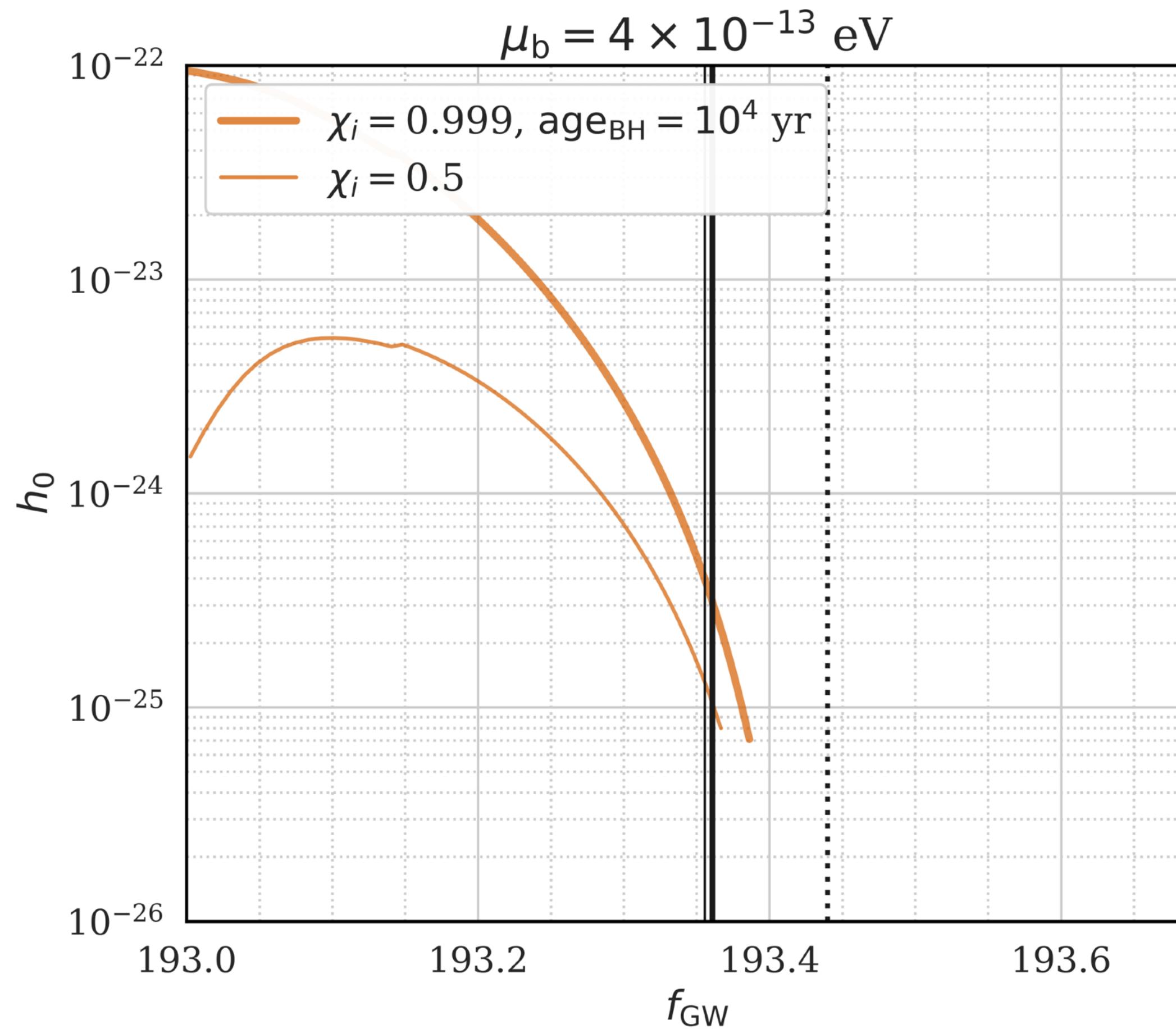


# What do the ensemble signals look like?



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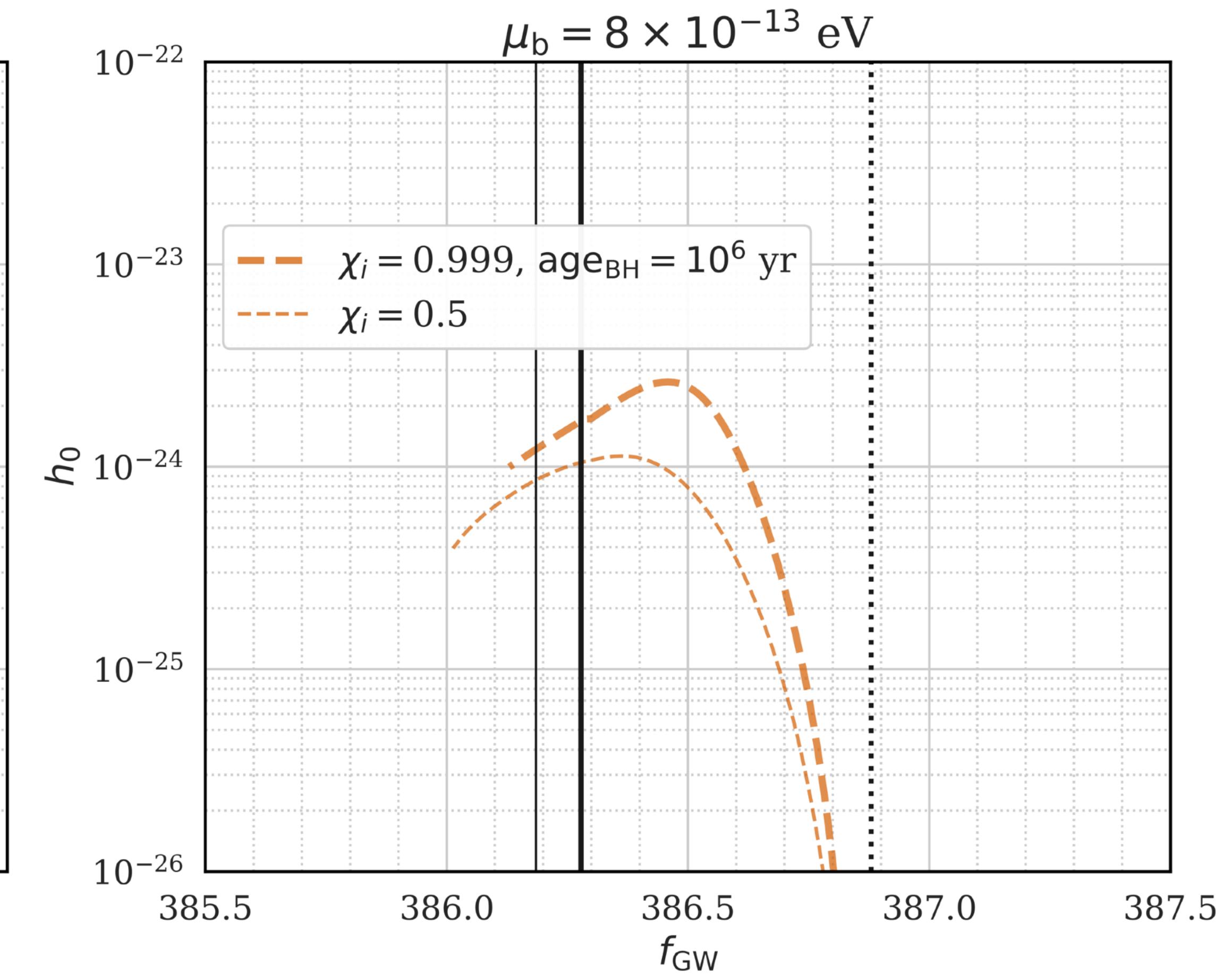
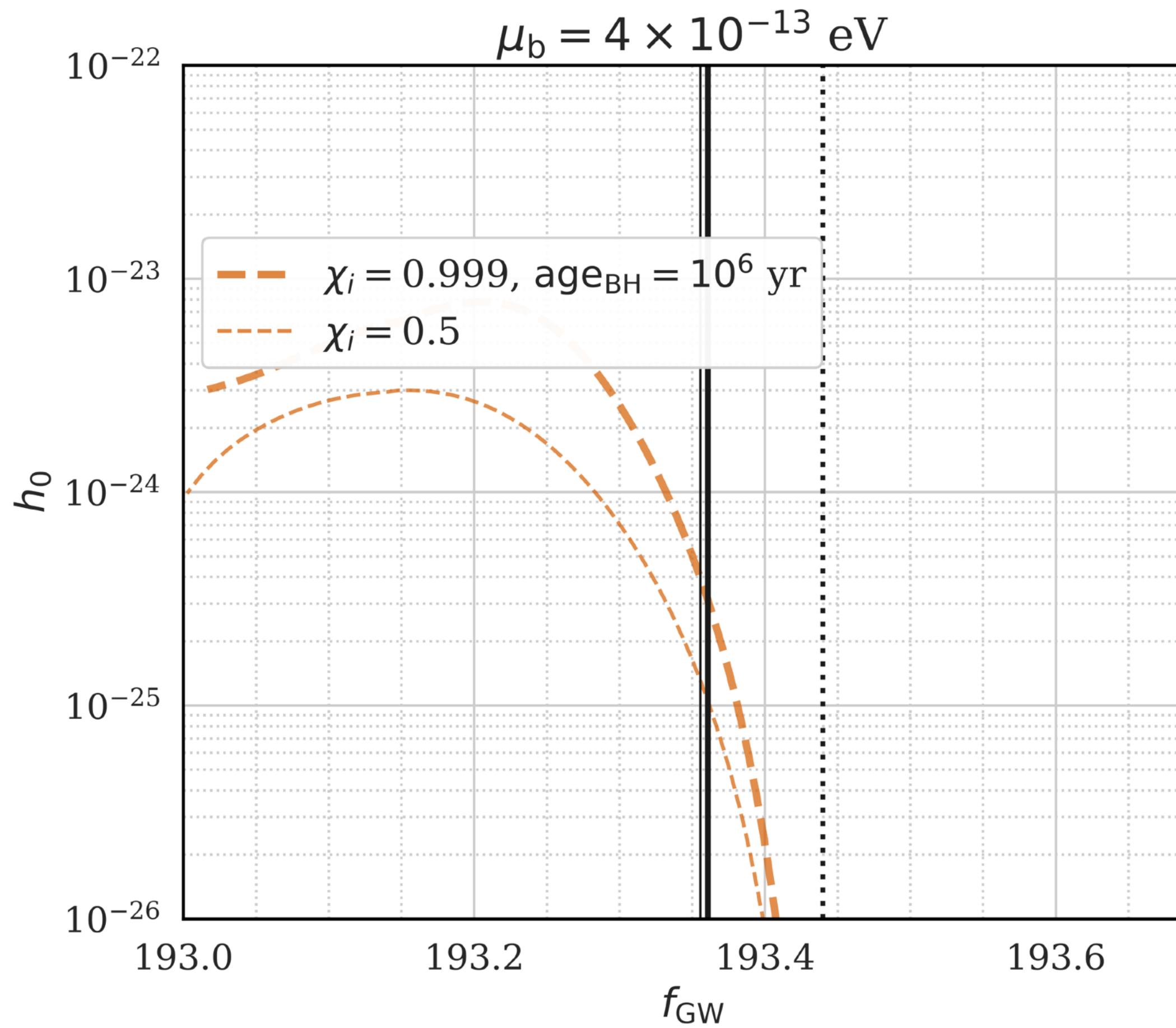
The orange curves define the envelopes of the potential signals over a wide range of boson masses.



The black solid vertical lines show  $\min f_{\text{GW}}$  given  $M_{\text{BH}} = 20M_{\odot}$ ; the dotted line shows the rest mass frequency.

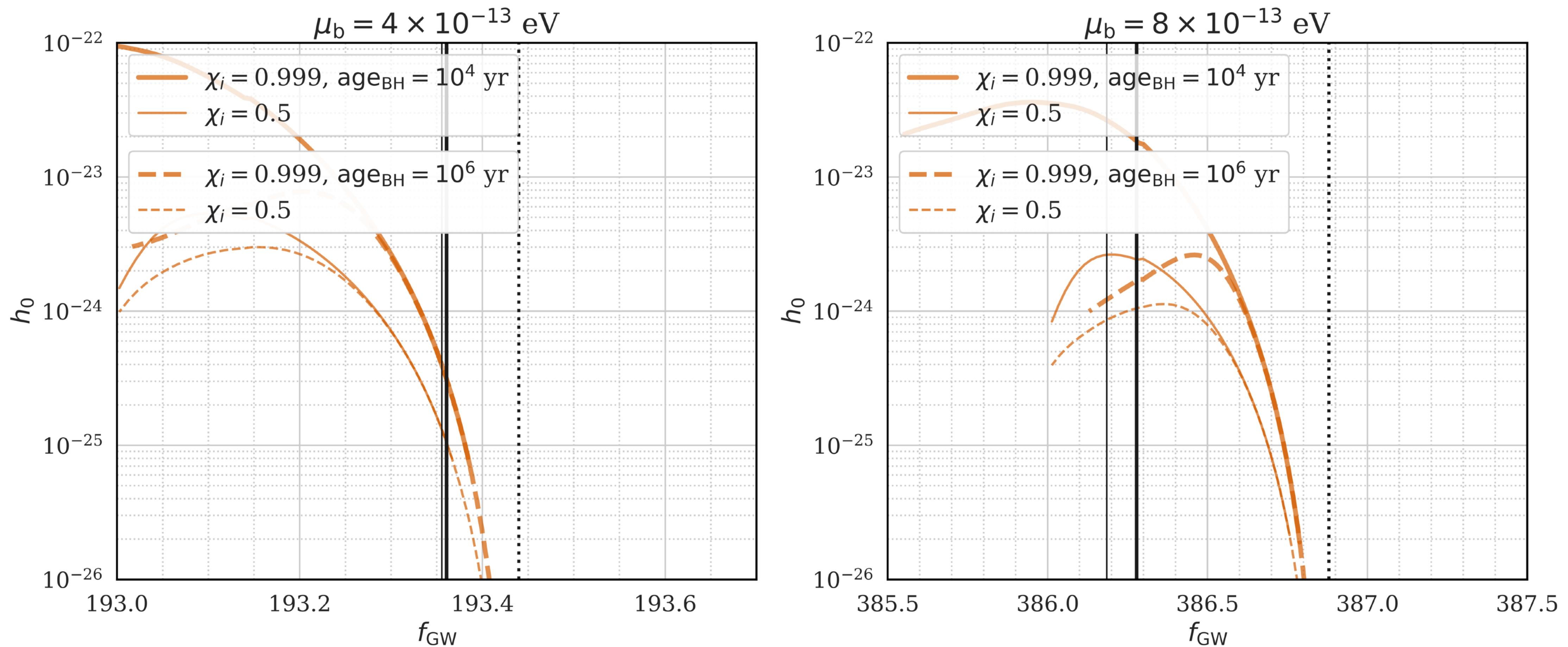
# What do the ensemble signals look like?

Signals get quieter as black holes age => envelopes drop



The black solid vertical lines show  $\min f_{\text{GW}}$  given  $M_{\text{BH}} = 20M_{\odot}$ ; the dotted line shows the rest mass frequency.

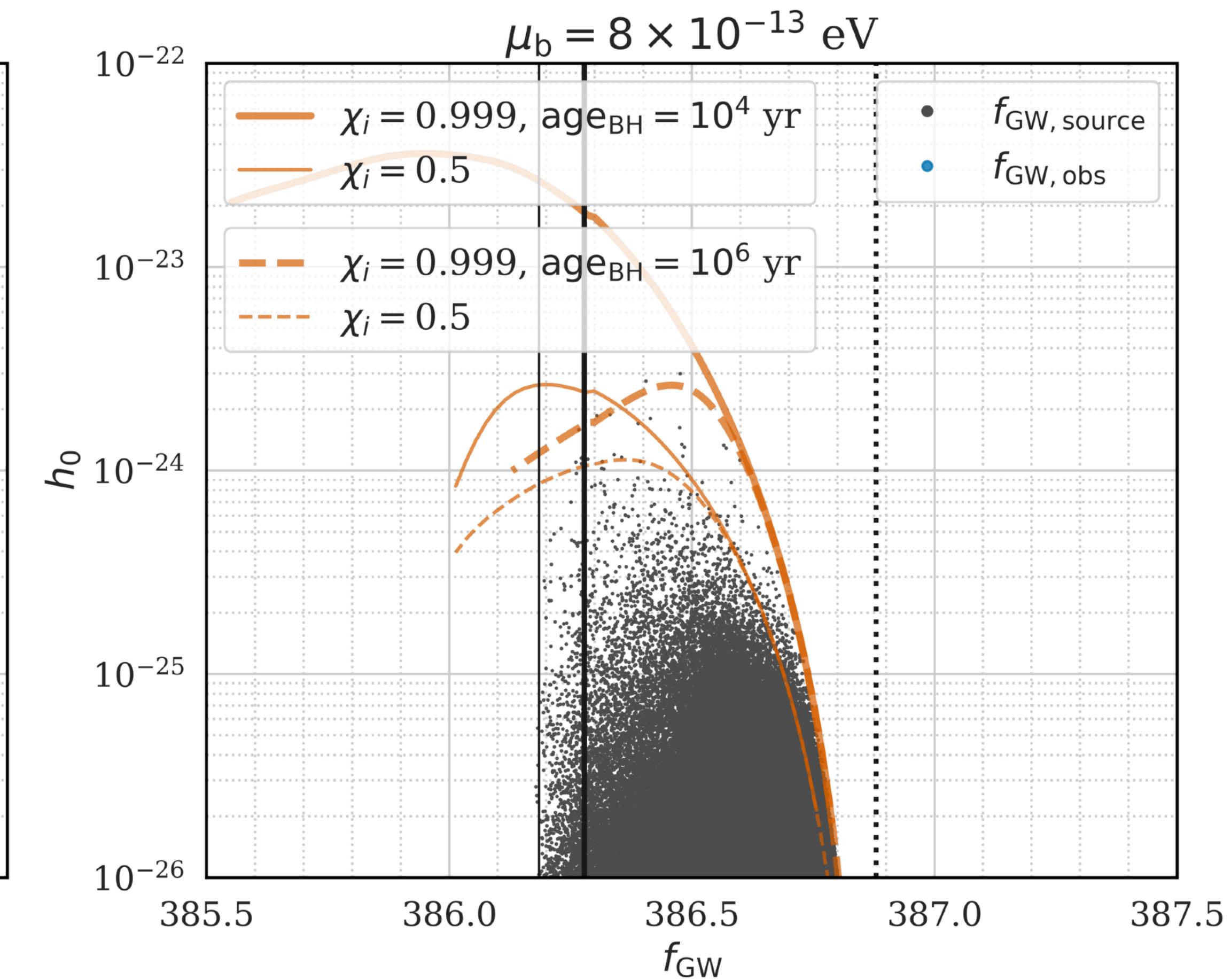
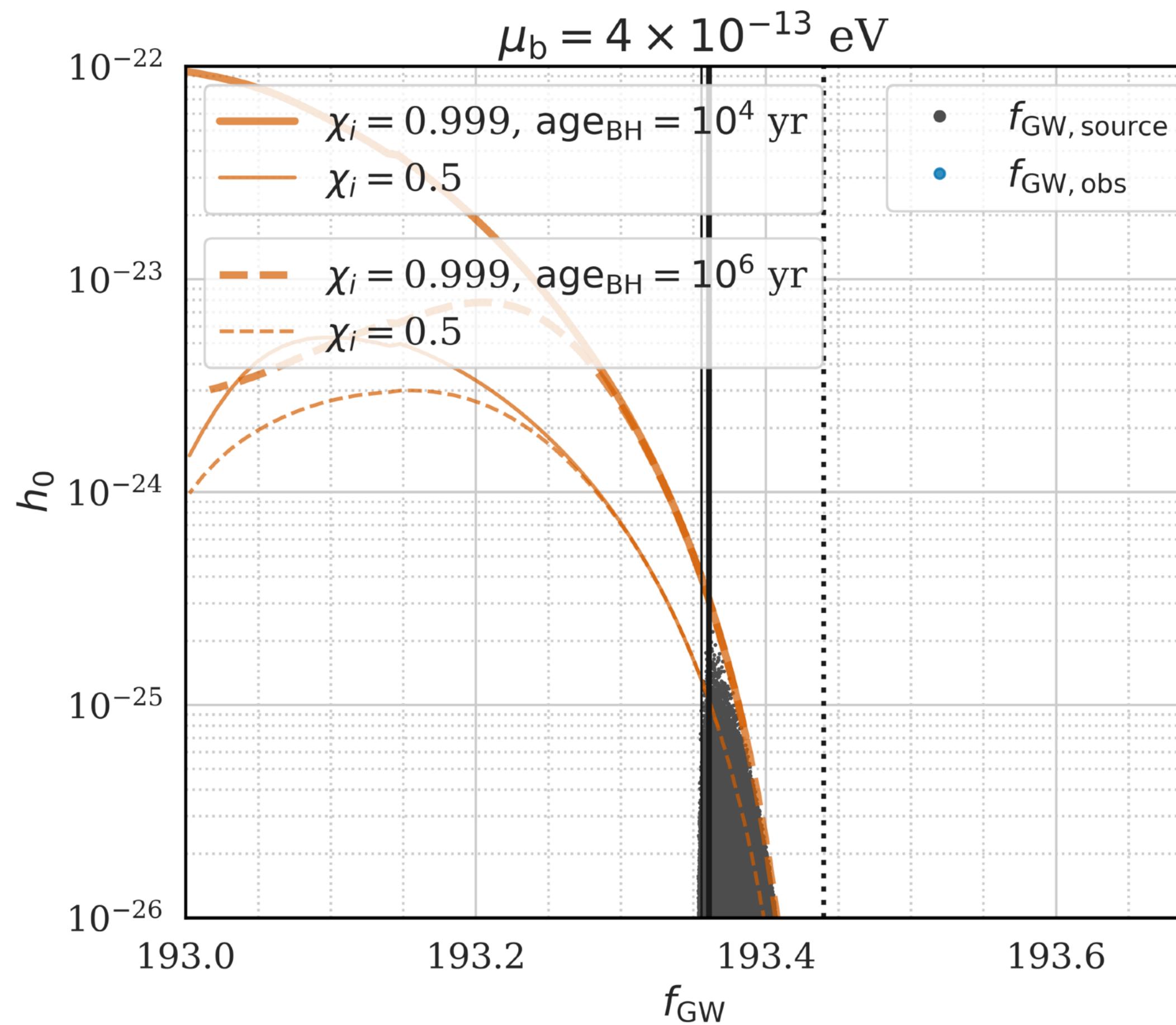
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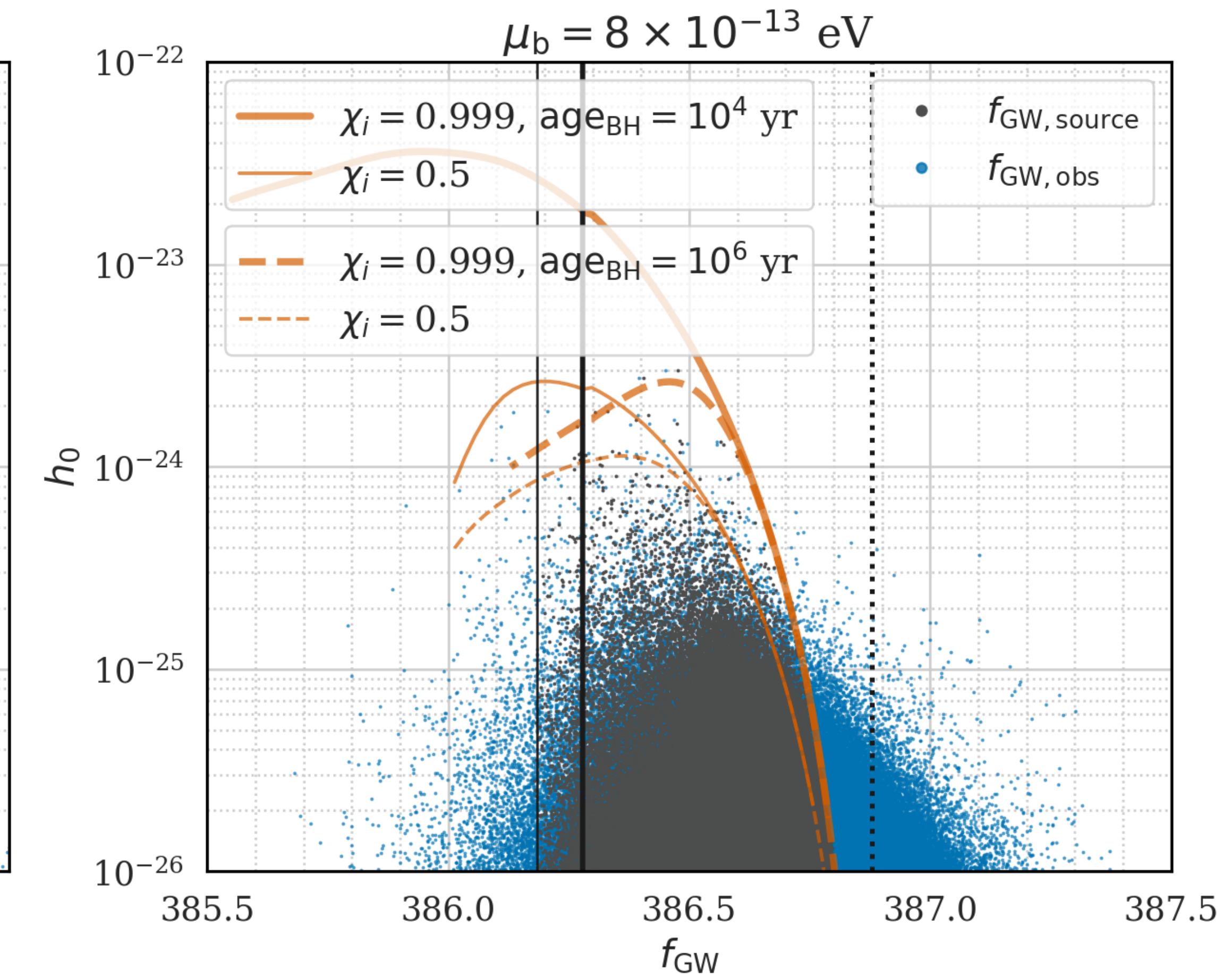
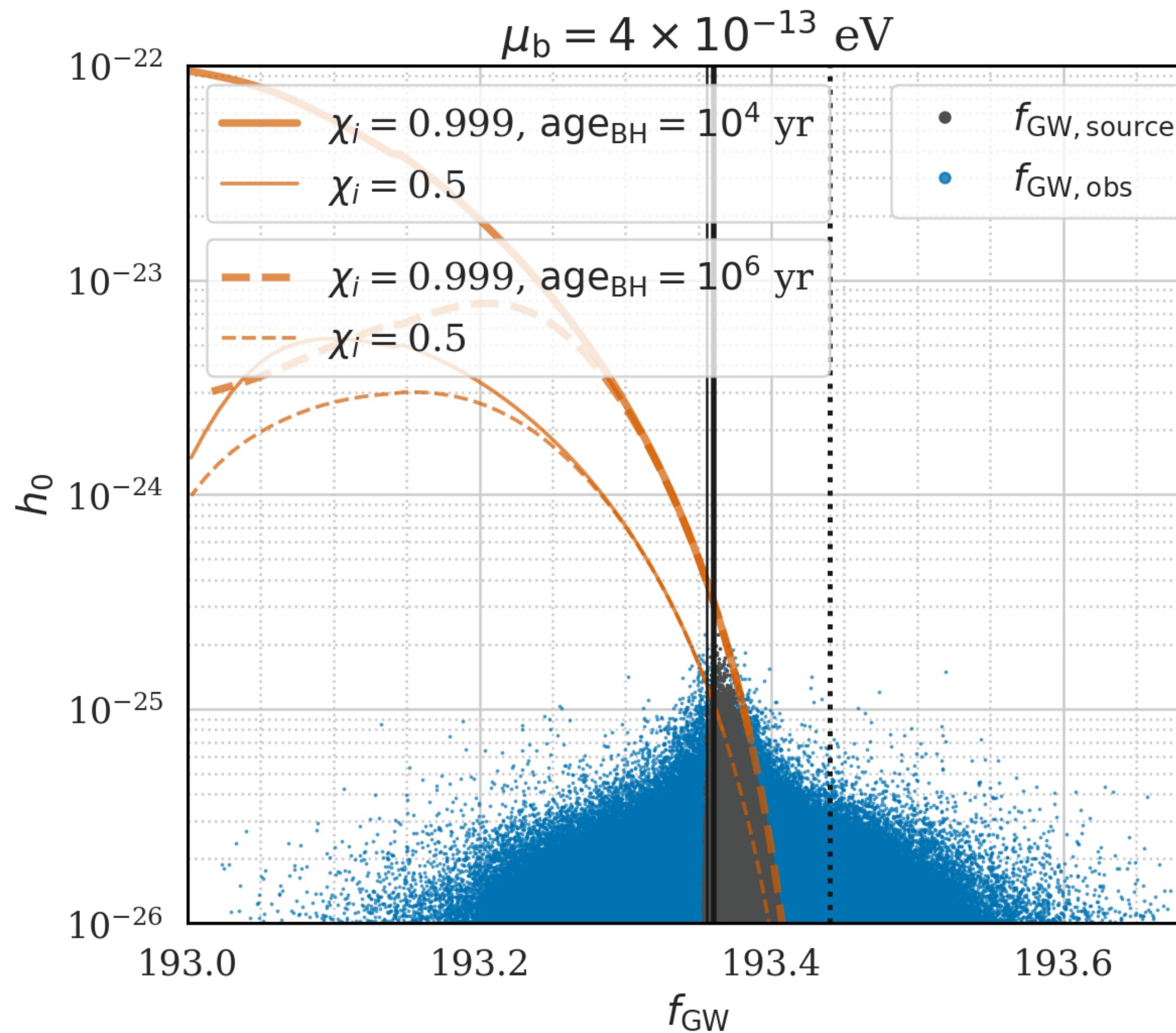
The signals w/ their source frequencies fill in the regions defined by the envelopes and vertical lines.



The black solid vertical lines show  $\min f_{\text{GW}}$  given  $M_{\text{BH}} = 20M_\odot$ ; the dotted line shows the rest mass frequency.

# What do the ensemble signals look like?

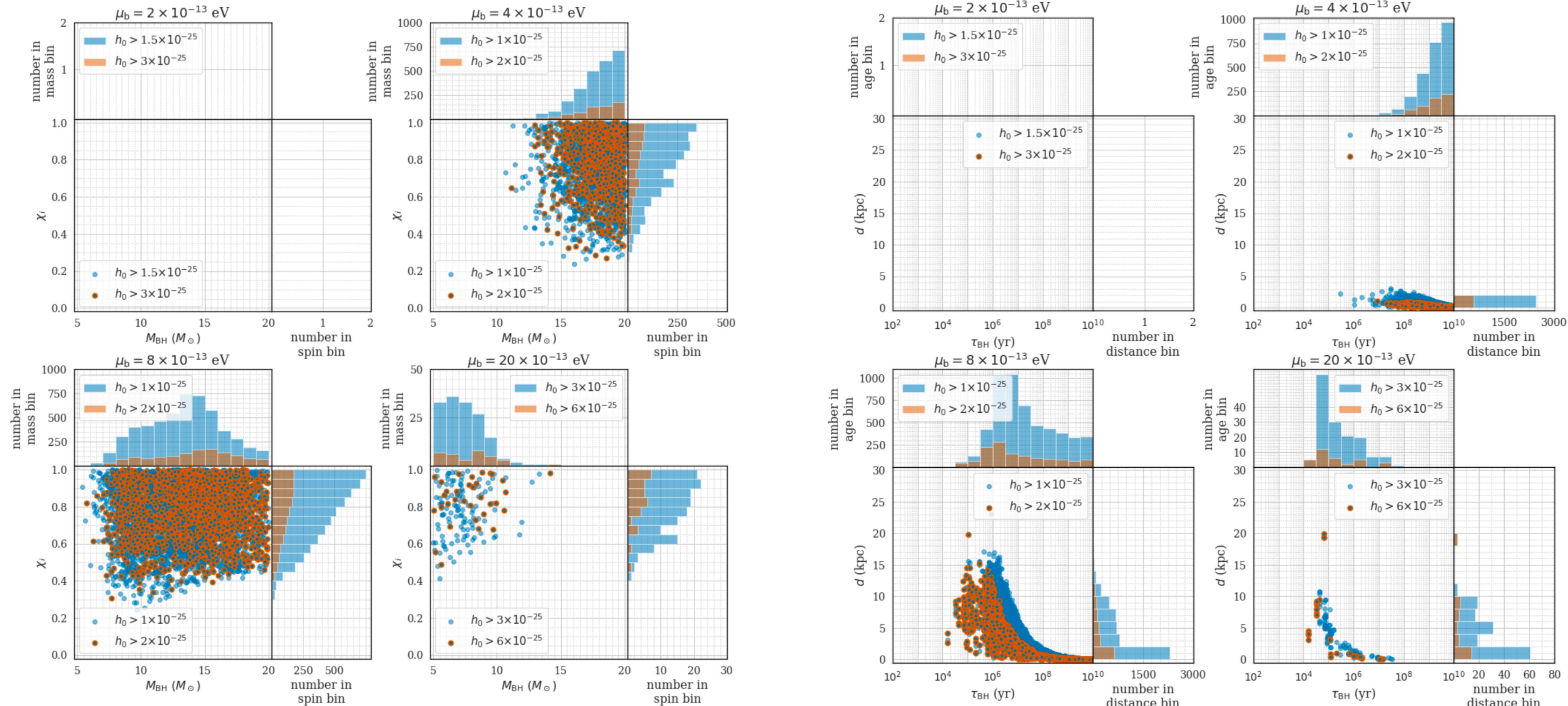
The black hole velocities add a large Doppler shift and "smear out" the distribution.



The black solid vertical lines show  $\min f_{\text{GW}}$  given  $M_{\text{BH}} = 20M_\odot$ ; the dotted line shows the rest mass frequency.

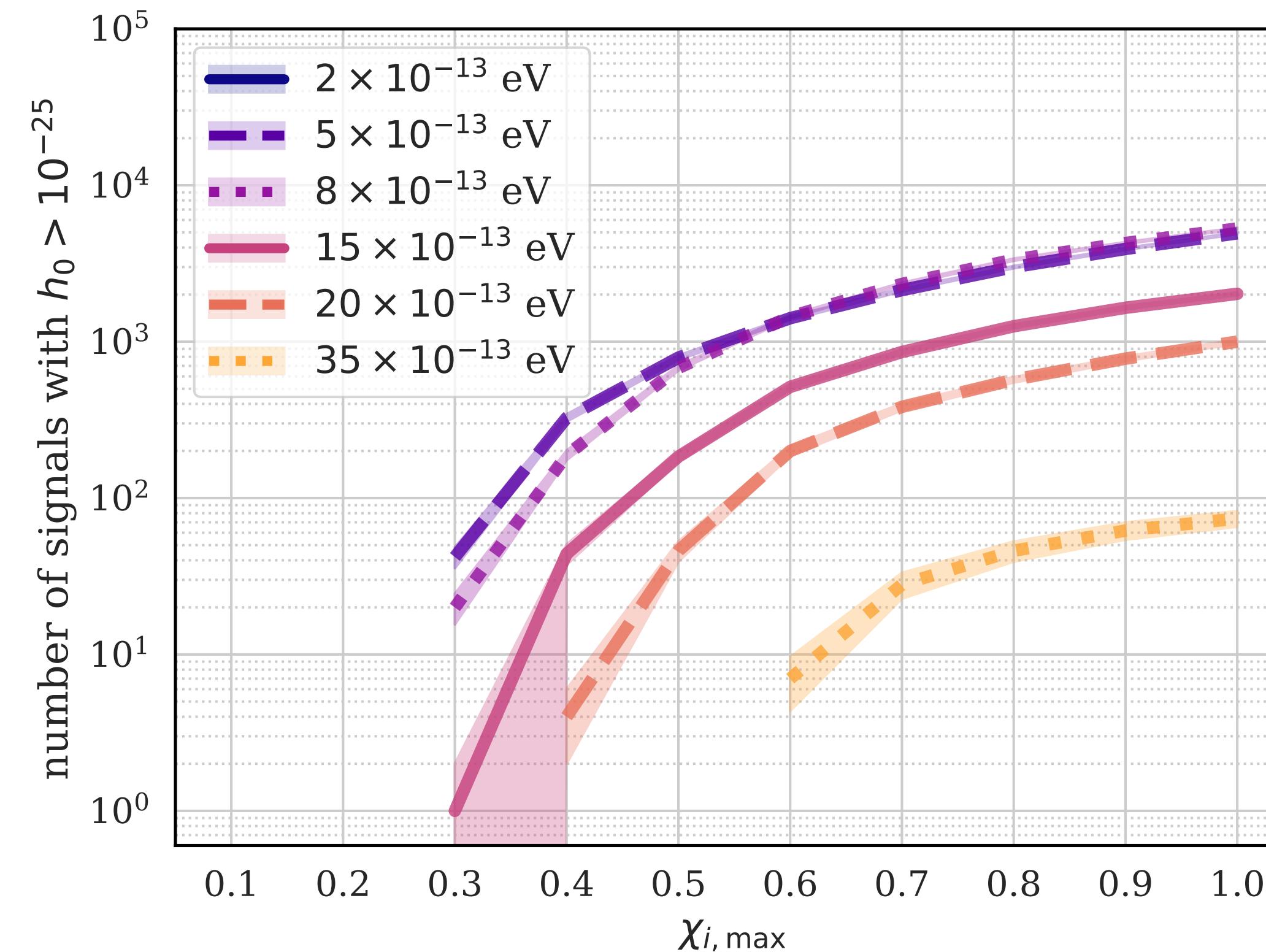
# Which black holes produce ~detectable signals?

The “relevant” black holes depend on the boson mass, through both the signal properties and the detector sensitivity



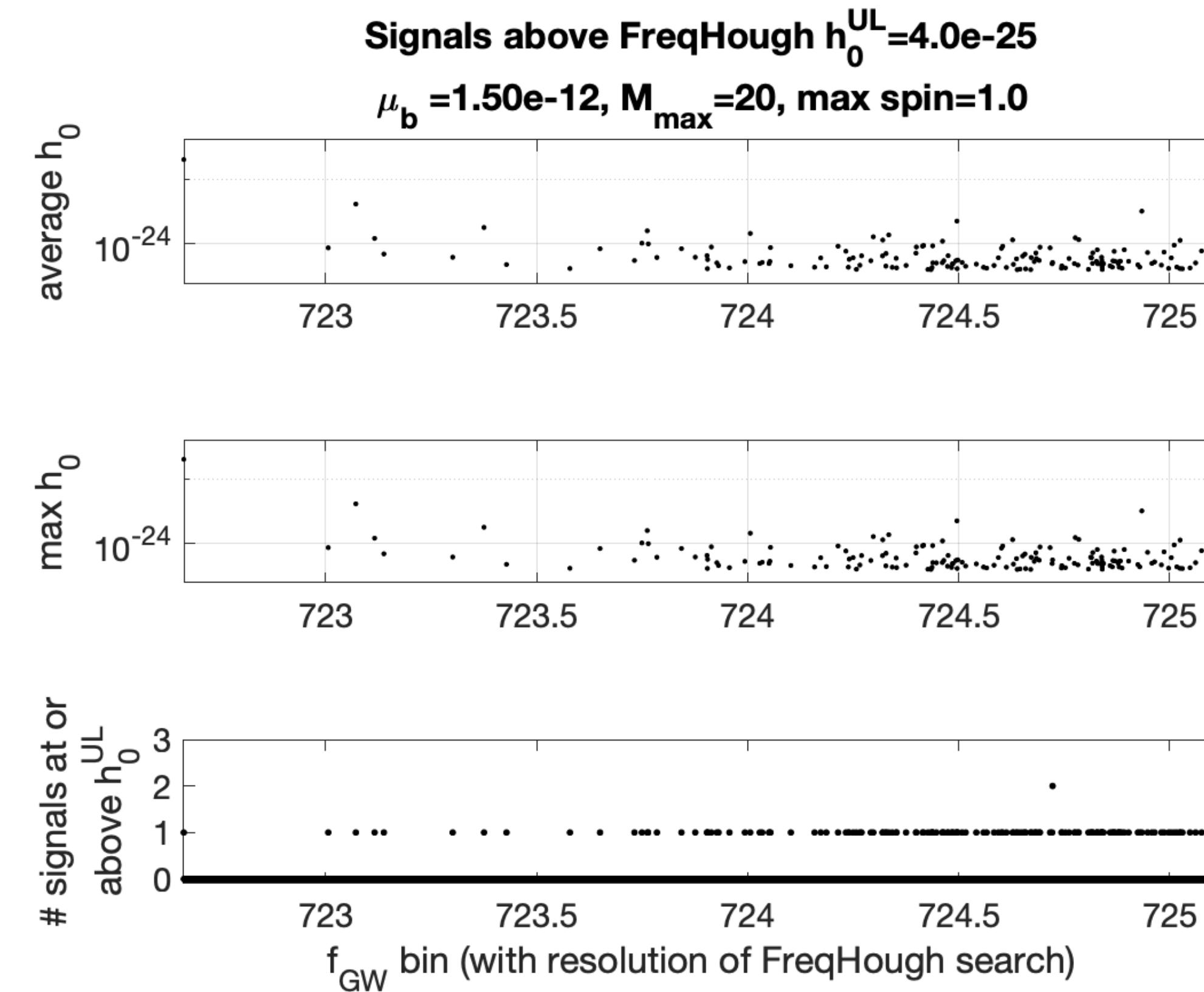
# Varying spin

For our BH population, initial spins of  $>0.3$  are needed to produce  $\sim$ detectable signals



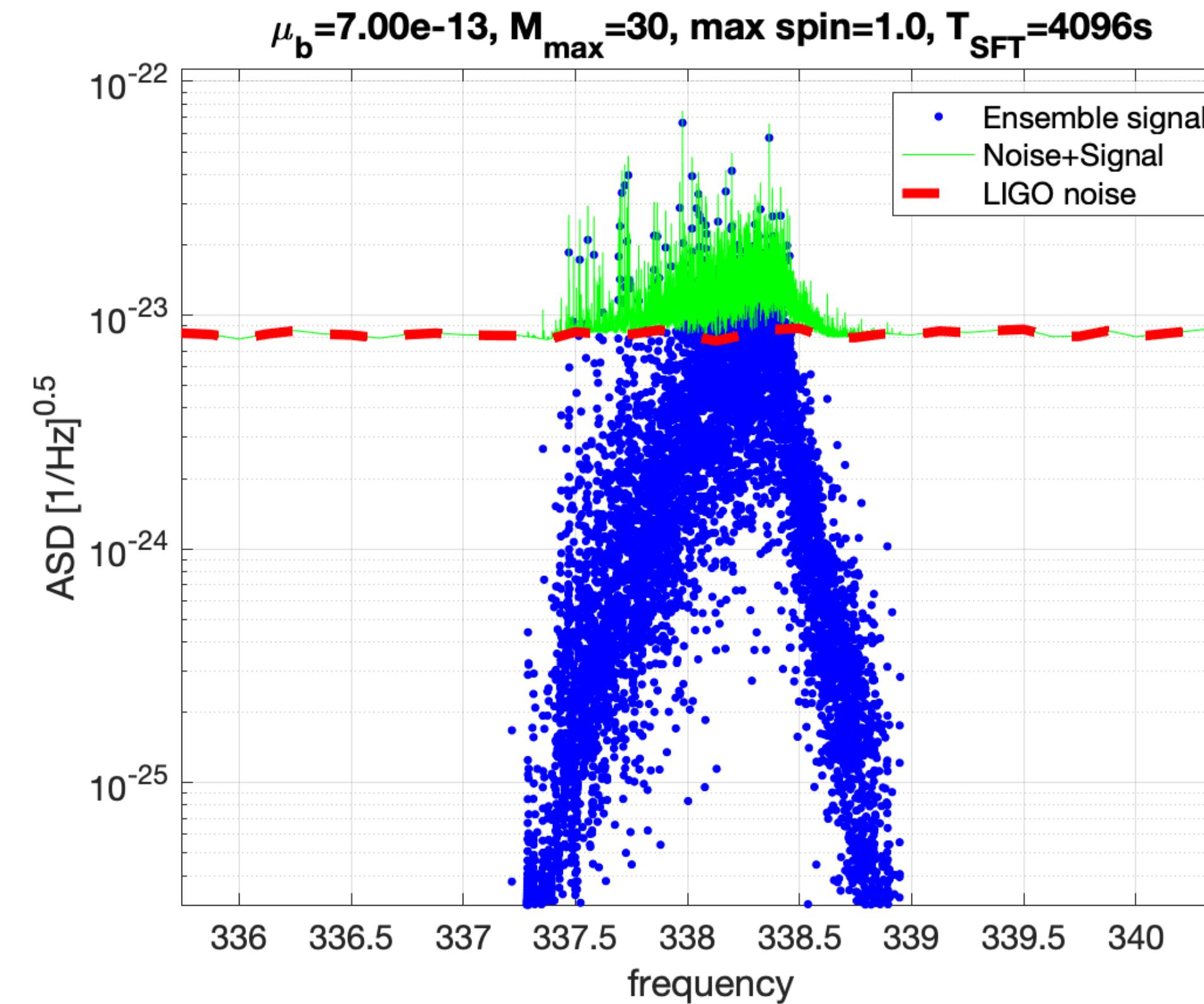
# Concerns about signal density

Searches can turn up multiple signals per frequency bin => signal confusion



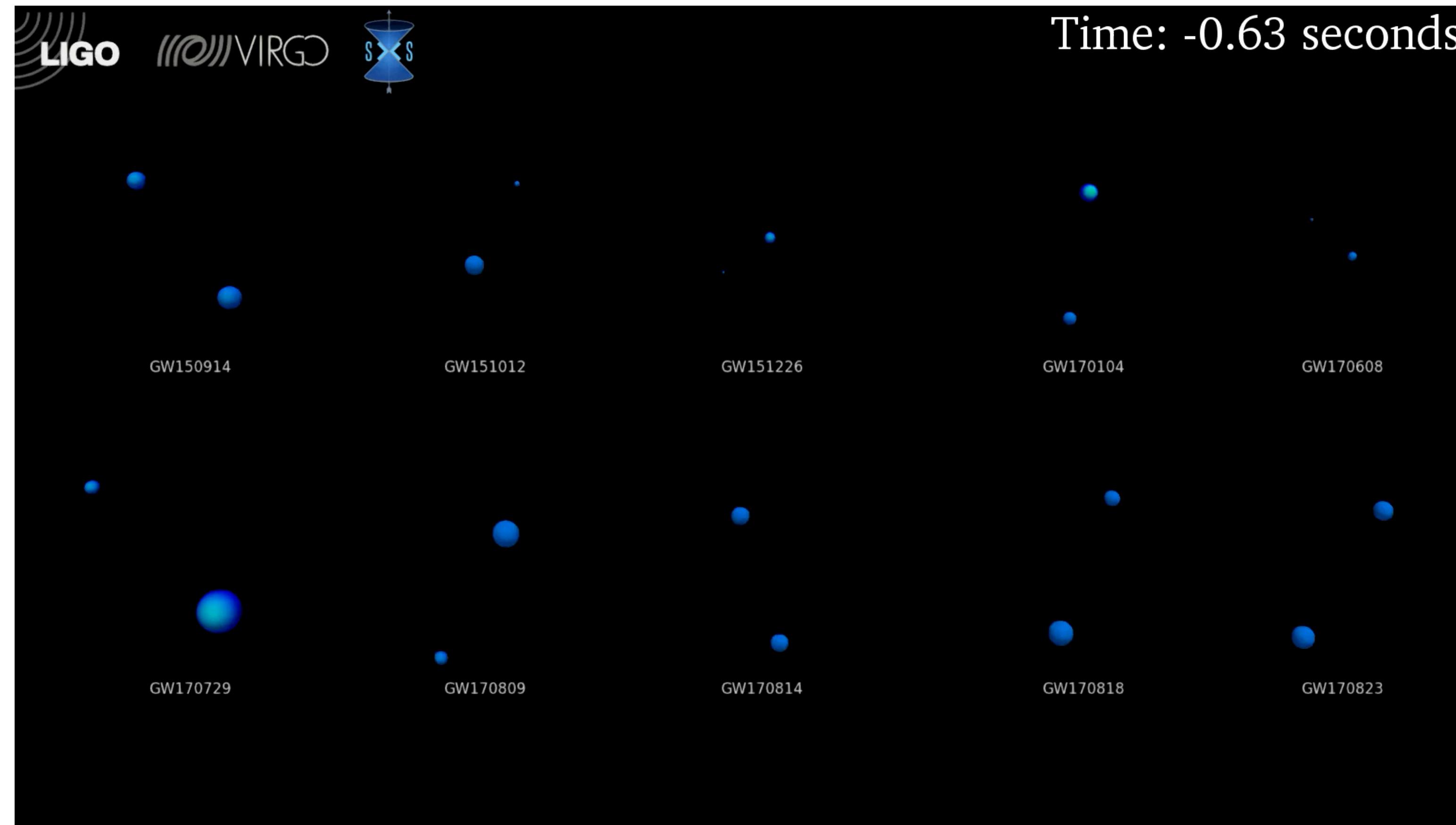
# Concerns about signal density

A dense signal background can affect the noise floor beyond what typical CW searches expect



# sources of transient gravitational waves

A changing mass quadrupole produces gravitational waves



# Continuous gravitational waves

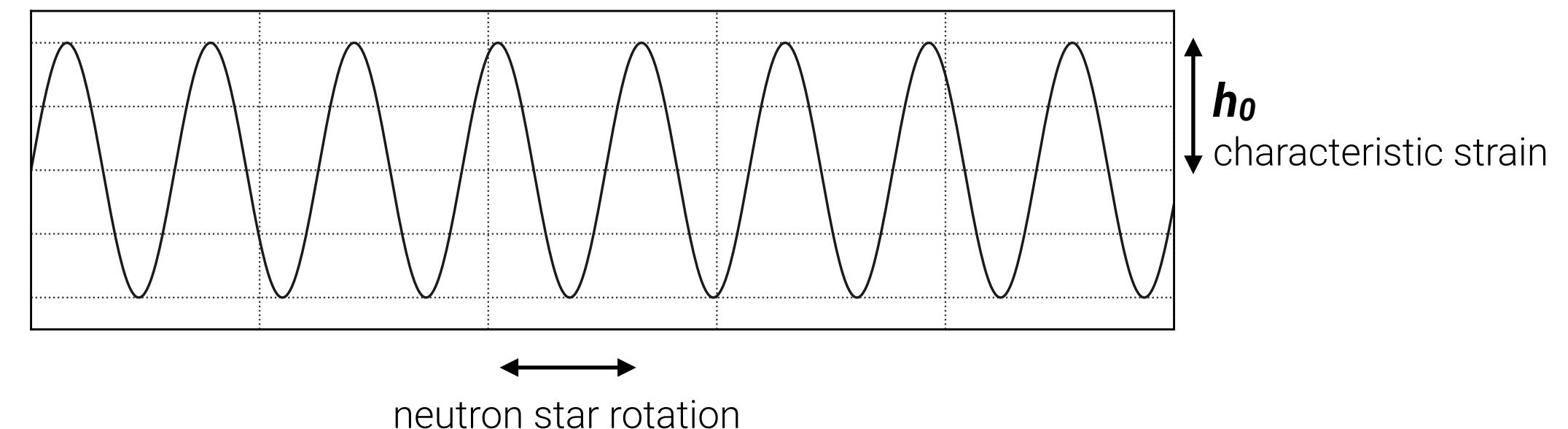
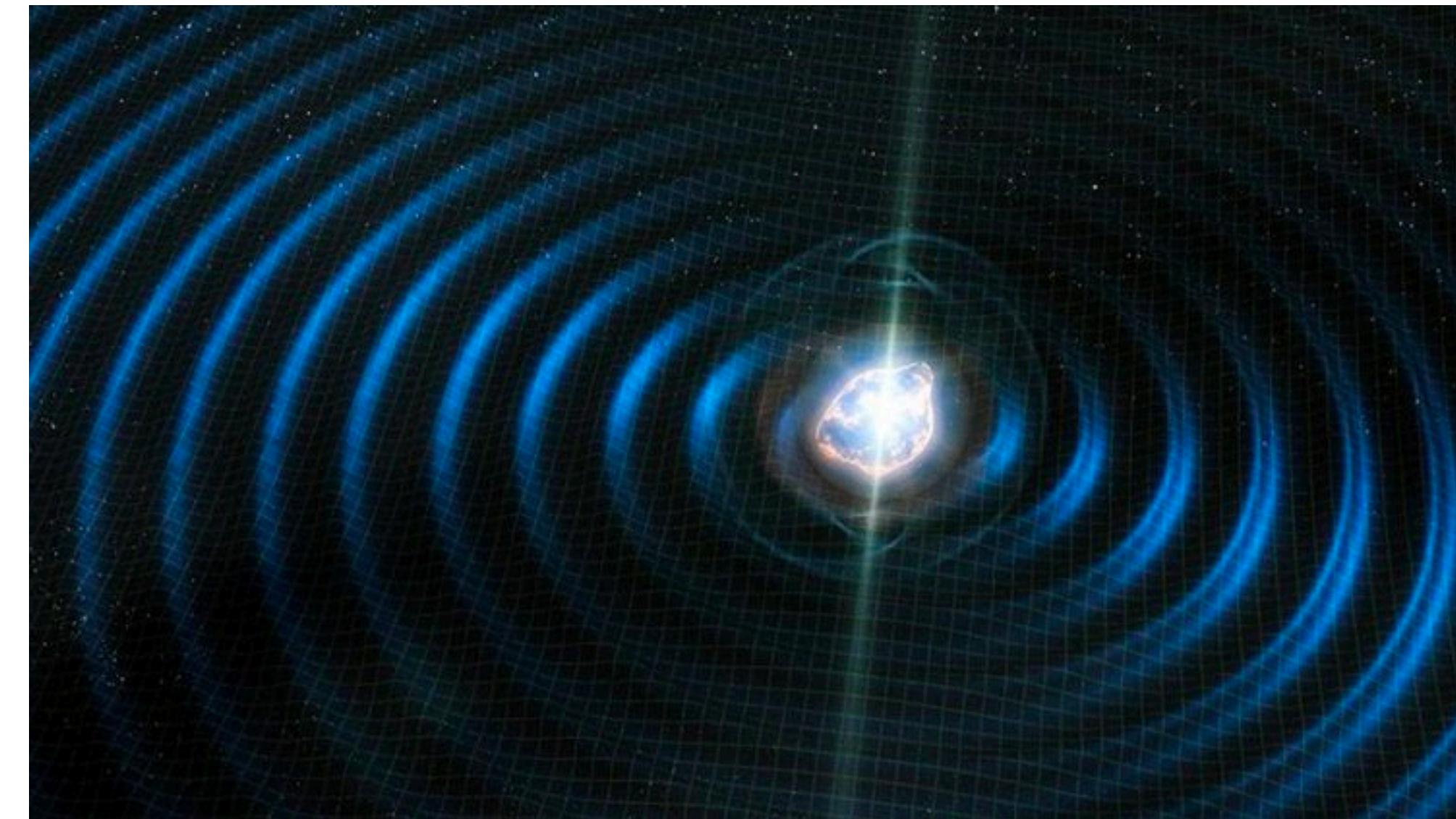
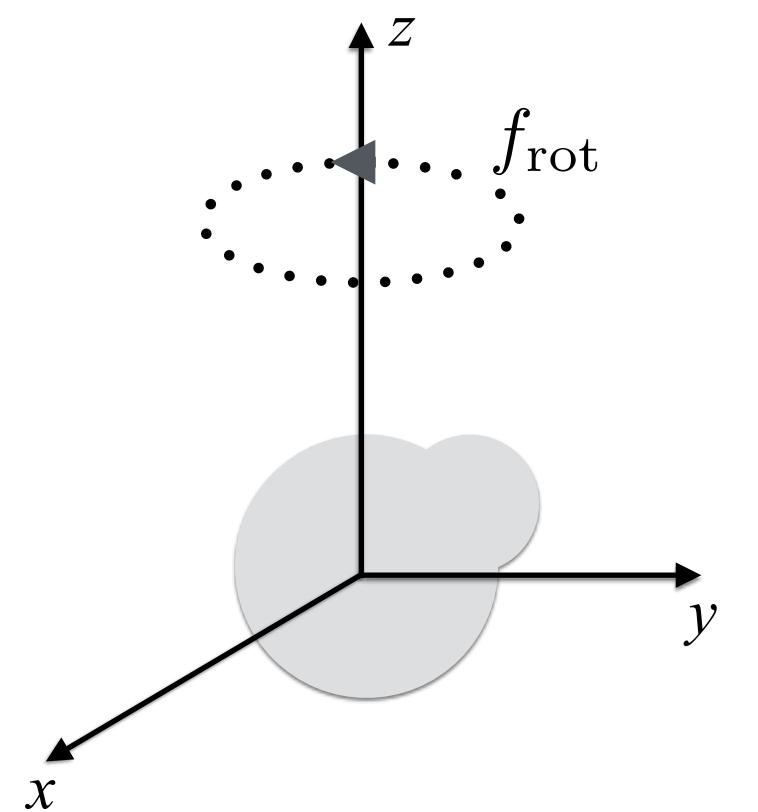
A changing mass quadrupole produces gravitational waves

$$h_0 = \frac{4\pi G I_{zz}}{c^4} \frac{f_{\text{GW}}^2 \epsilon}{d}$$

2x rotation frequency  $f_{\text{rot}}$

ellipticity  $\frac{|I_{xx} - I_{yy}|}{I_{zz}}$

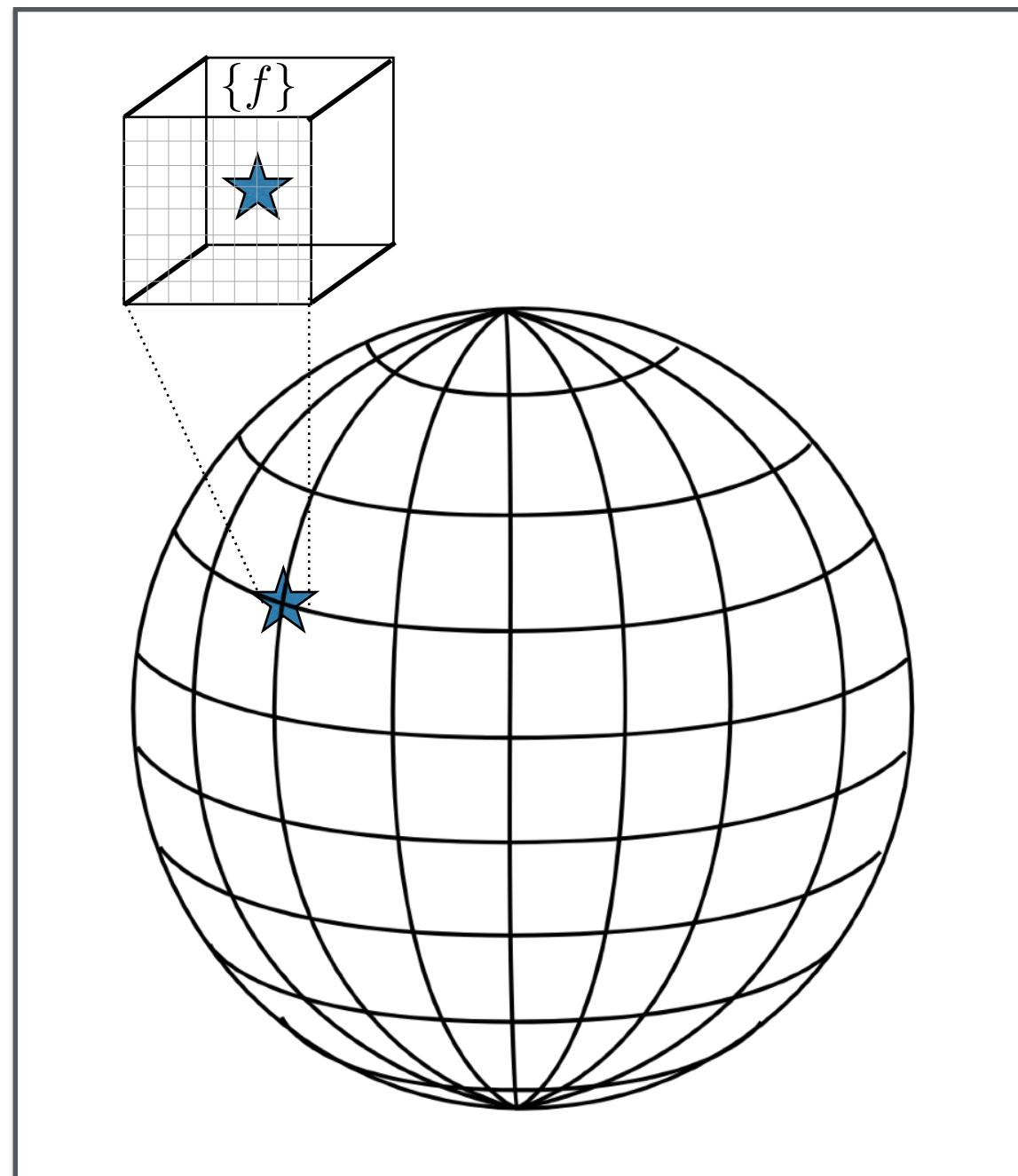
distance from detector



# Types of searches

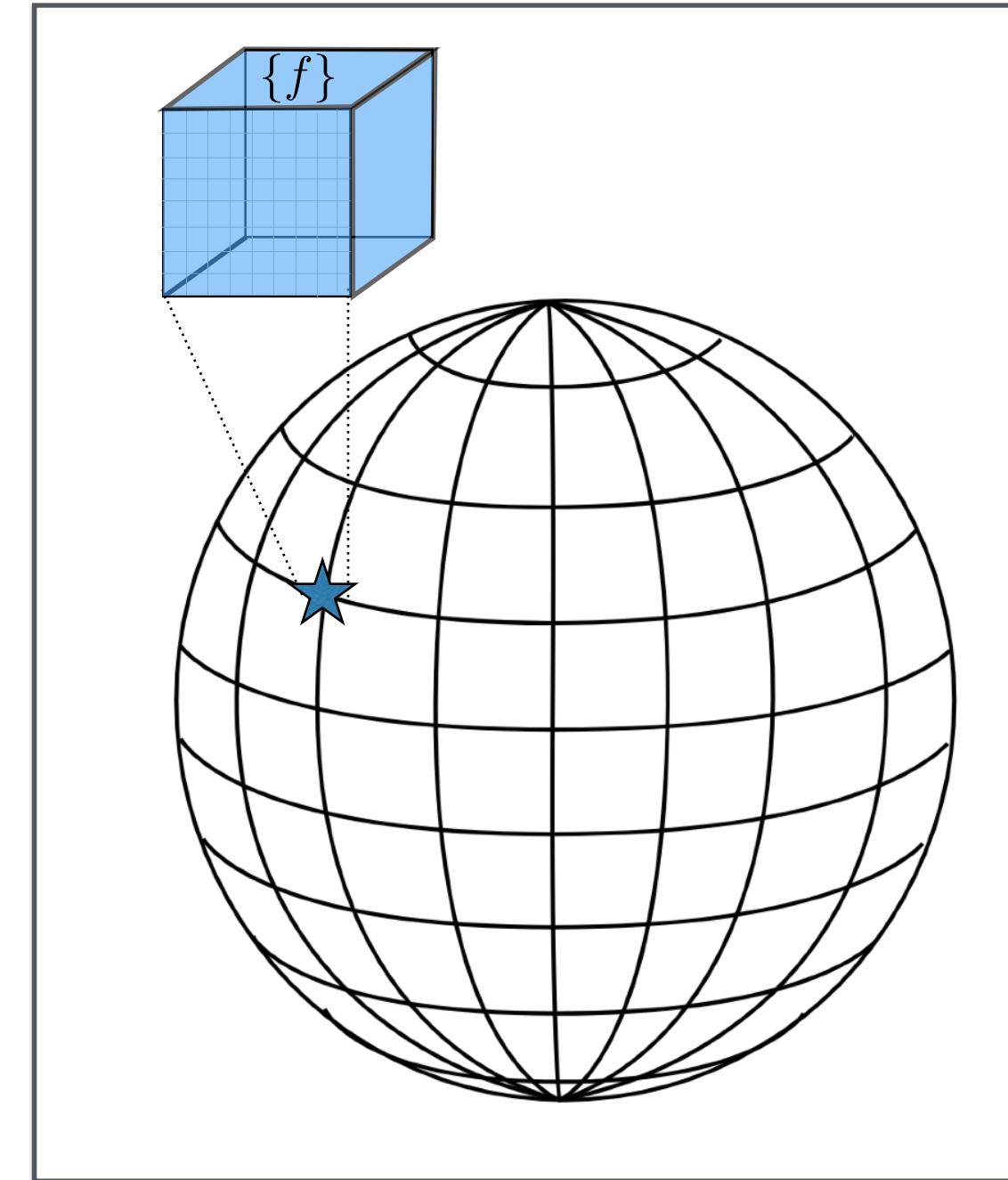
## targeted search

known pulsars



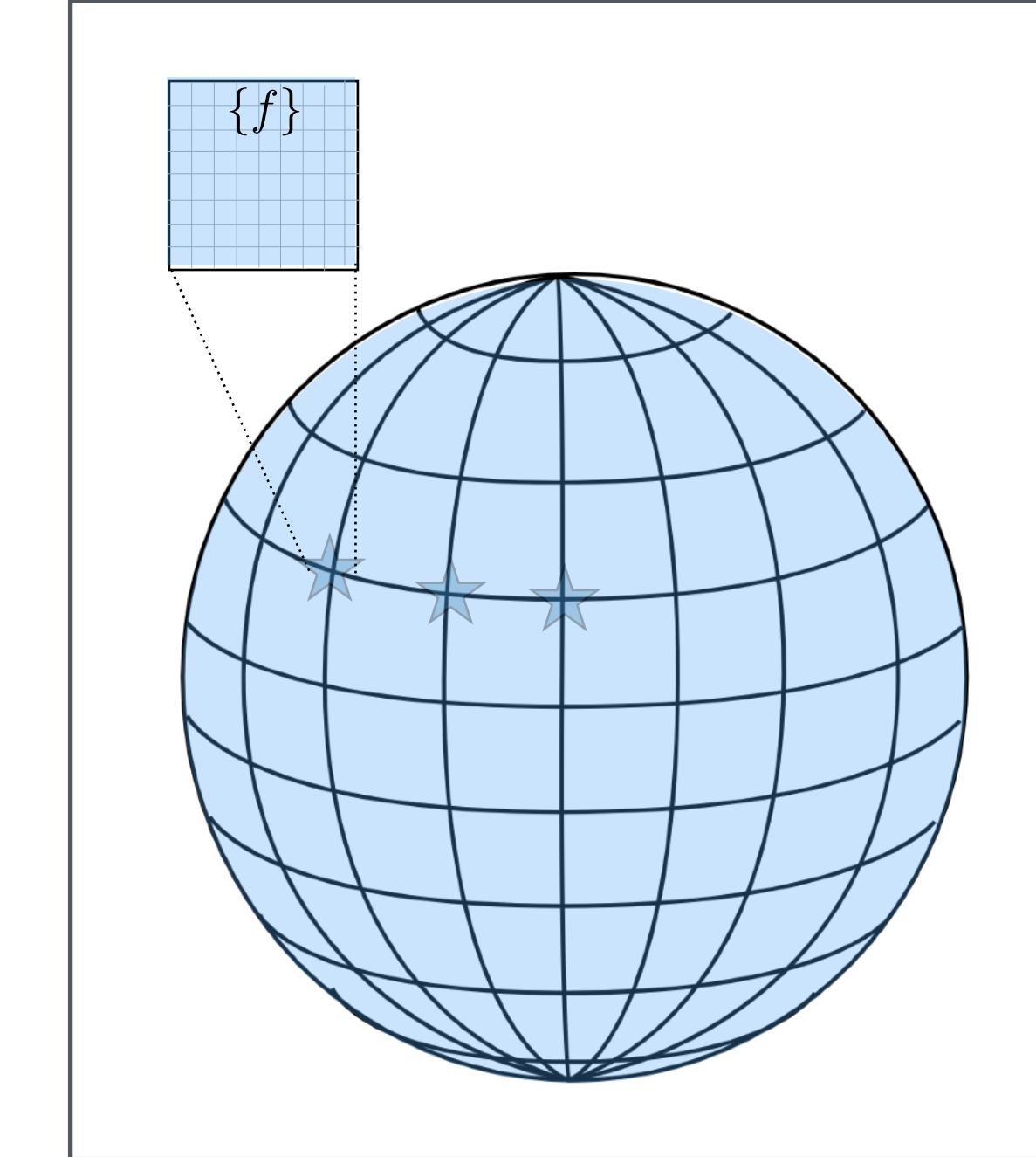
## directed search

known neutron stars



## all-sky search

minimal assumptions



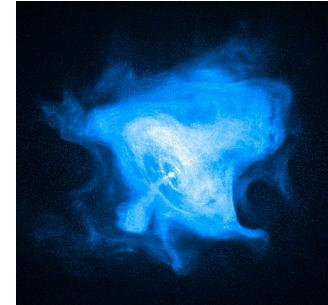
computational cost



search sensitivity



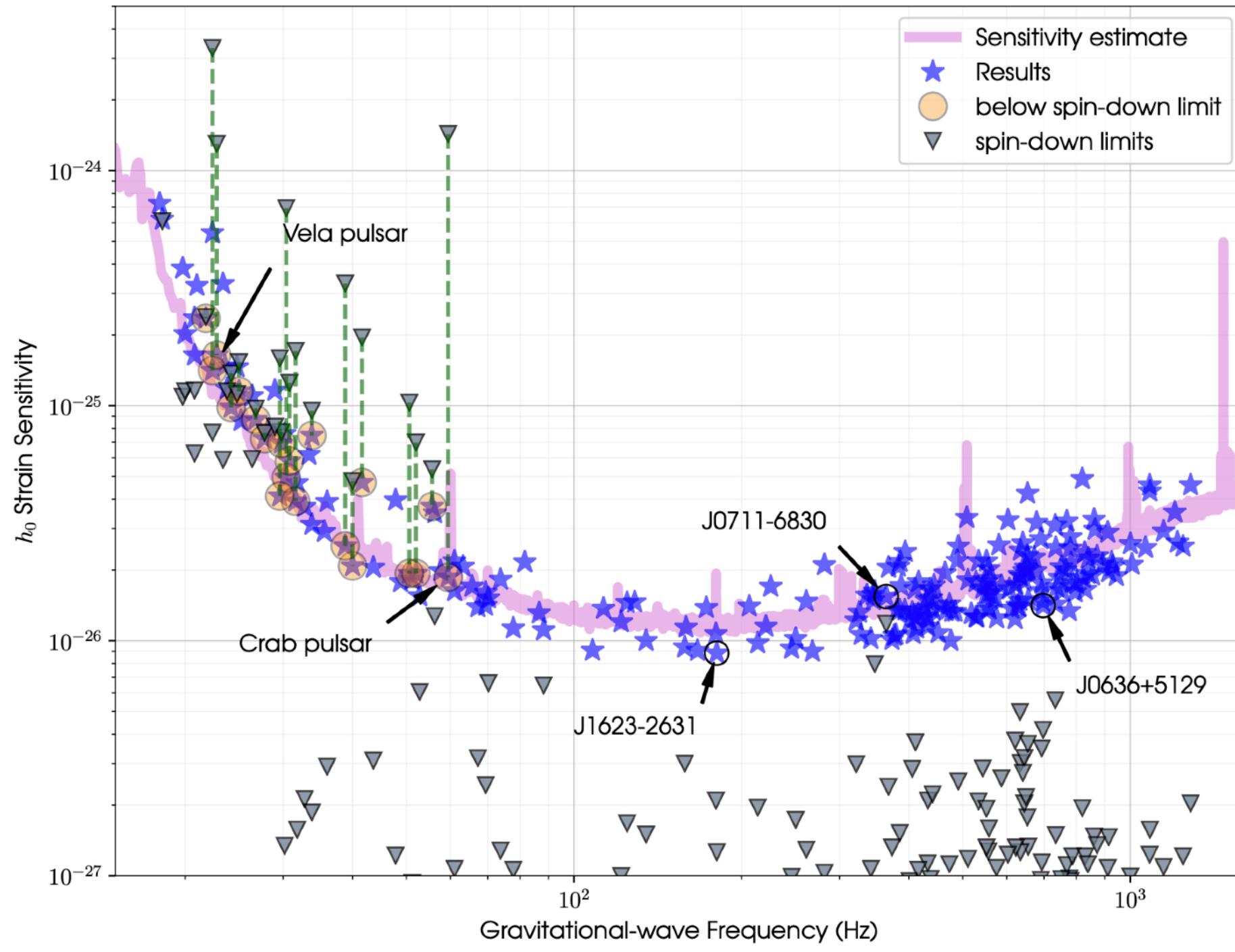
# Interpreting null results (upper limits) – neutron stars



$$h_0 = \frac{4\pi G I_{zz}}{c^4} \frac{f_{\text{GW}}^2 \epsilon}{d}$$

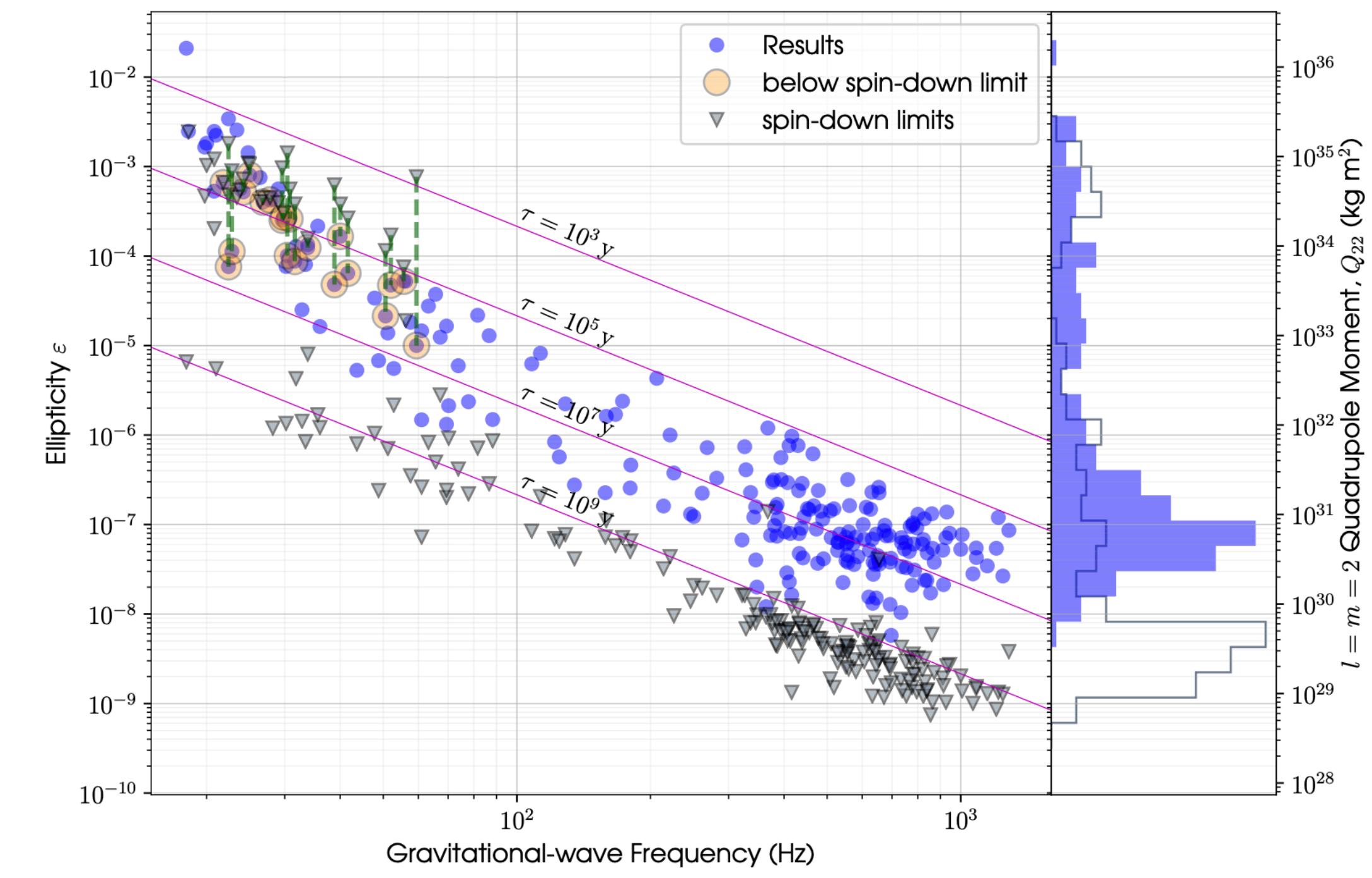
2x rotation frequency  $f_{\text{rot}}$   
 ellipticity  $\frac{|I_{xx} - I_{yy}|}{I_{zz}}$   
 distance from detector

upper limit on gravitational-wave signal strength

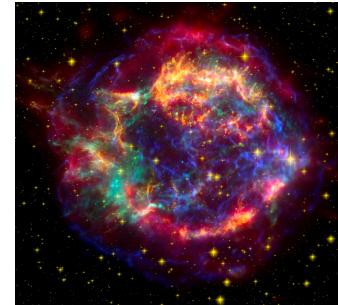


directly translates to

upper limit on neutron star ellipticity (~max mountain height)



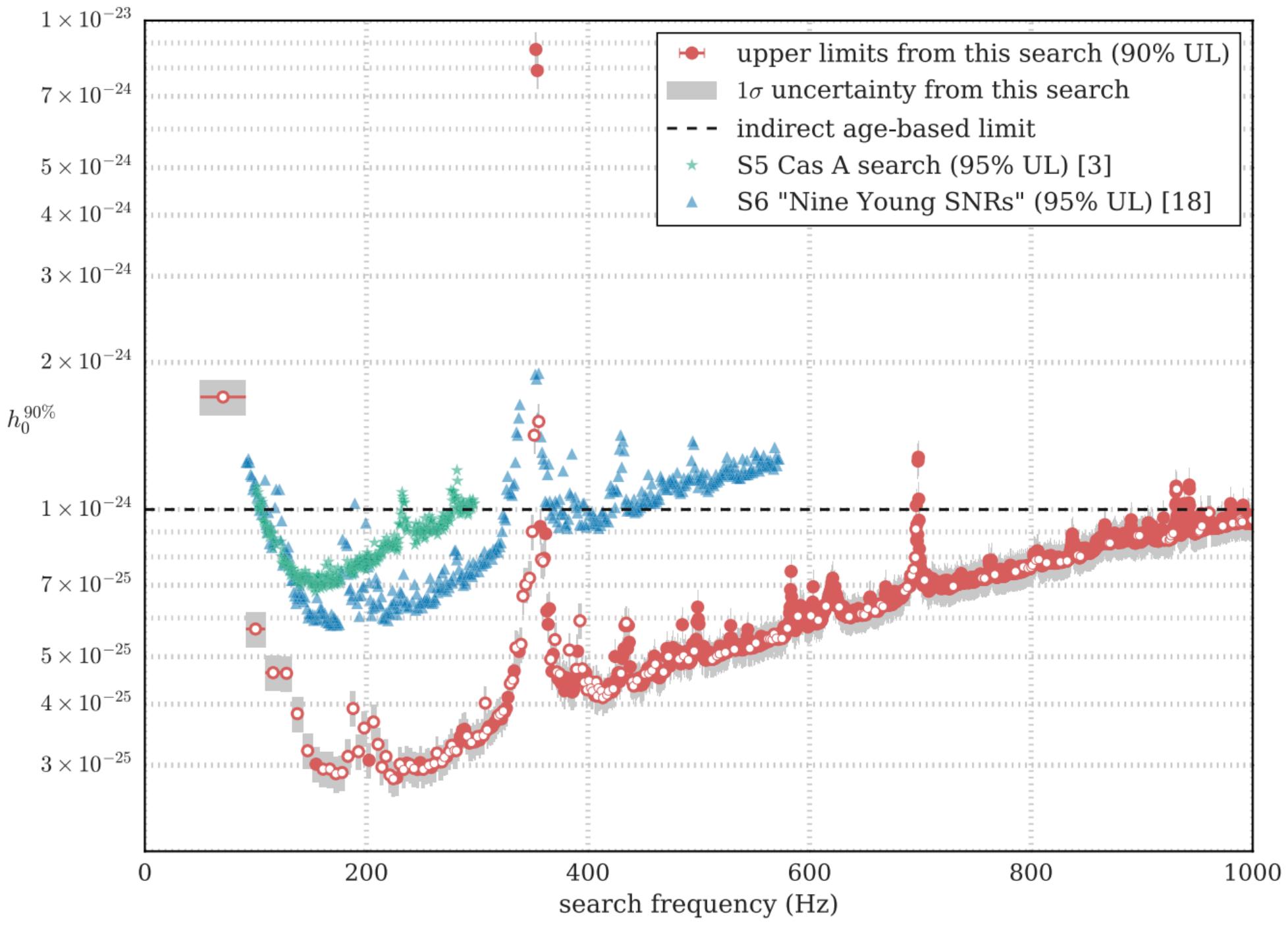
# Interpreting null results (upper limits) – neutron stars



$$h_0 = \frac{4\pi G I_{zz}}{c^4} \frac{f_{\text{GW}}^2 \epsilon}{d}$$

2x rotation frequency  $f_{\text{rot}}$   
ellipticity  $\frac{|I_{xx} - I_{yy}|}{I_{zz}}$   
distance from detector

upper limit on gravitational-wave signal strength



translates to,  
assuming  $f_{\text{GW}}$ ,

upper limit on neutron star ellipticity (~max mountain height)

