First Structures in Fuzzy Dark Matter

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Göttingen FDM-workshop

Fuzzy Dark Matter (FDM) motivations and tensions

- ► New probes for FDM:
- Full-physics cosmological simulations (Mocz et al., 2019; Mocz et al., 2020)
 - ► First galaxies uniquely probe physical nature of dark matter
 - Test with James Webb Space Telescope (JWST)
- ► Large axion mass limit (Mocz et al., 2018)
 - ► connection between FDM & CDM
- ► Adding in the strong-CP scale to FDM
- Further probes: small-scale astrophysical features of FDM [student works]
 - ► dynamical friction (Lancaster, Giovanetti, **Mocz**,..., 2020)
 - dynamical heating (Church, Mocz & Ostriker, 2019)
 - ► soliton + SMBH (Davies & **Mocz**, 2019)

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Jerry Ostriker Columbia



Jesús Zavala Iceland

The nearly century-old dark matter problem is one of the most intriguing mysteries in modern physics. (Zwicky, 1933)

We do not know the nature of 84% of matter in the Universe, yet it is thought to govern cosmic structure and hold galaxies and clusters together.























- Assume DM is a cold, ultralight scalar field (Peebles, 2000; Hu, Barkana & Gruzinov, 2000; Hui et al., 2017)
- ► T = o in early universe, forms a BEC ⇒ macroscopic quantum properties
- 'Quantum Pressure' suppresses gravitational collapse below de Broglie wavelength
 - ▶ Require $m \sim 10^{-22}$ eV to have $\lambda_{dB} \sim 1 \text{ kpc}$ for $\nu \sim 100 \text{ km s}^{-1}$
- Governed by Schrödinger–Poisson

$$i\hbar\frac{\partial\psi}{\partial t} = -\frac{\hbar^2}{2m}\nabla^2\psi + m\nabla\psi, \quad \nabla^2\nabla = 4\pi G(\rho - \overline{\rho}), \quad \rho = |\psi|^2 \qquad (1)$$

Assume DM is a cold, ultralight scalar field (Peebles, 2000; Hu, Barkana & Gruzinov, 2000; Hui et al., 2017)

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Motivation for FDM

Astrophysics

CDM small scale challenges (Primack, 2009)

- deficit of dwarf galaxies / missing satellites problem (Klypin et al., 1999; Moore et al., 1999)
- problem with the abundance of isolated dwarfs (Zavala et al., 2009; Papastergis et al., 2011; Klypin et al., 2015)
- too-big-to-fail problem (Boylan-Kolchin, Bullock & Kaplinghat, 2011, 2012)
- cusp-core problem (Moore, 1994; Flores & Primack, 1994; Gentile et al., 2004; Donato et al., 2009; de Blok, 2010)

Theoretical Physics

- Axions solve the strong CP problem in QCD (Peccei–Quinn theory; $m \sim 10^{-5}-10^{-3}$ eV)
- String-theory compactifications provide class of ultralight axions ($m \sim 10^{-22}$ eV) (Arvanitaki et al., 2010; Bachlechner et al., 2019)

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FDM particle mass constraints from dwarf spheroidals



Particle needs to be light (m ~ 10⁻²² eV) to explain DM-dominated dwarf spheroidals (Fornax, Sculptor) as a pure fuzzy dark matter soliton core (Marsh & Pop, 2015)

FDM particle mass constraints from CMB



- A small boson mass suppresses large k initial DM power spectrum
- need $m \ge 10^{-24}$ eV, otherwise inconsistent with CMB fluctuations

(Hlozek et al., 2015; Hlozek, Marsh & Grin, 2017)

FDM particle mass constraints from Ly– α forest



▶ $m \gtrsim 10^{-21}$ eV, otherwise not enough Mpc-scale power in the Ly- α forest (Armengaud et al., 2017; Iršič et al., 2017)

Catch-22

- Catch-22 Problem: moderate tension in setting the particle mass to simultaneously capture large cores and right amount of substructure
- ► Also, simple soliton core model cannot simultaneously explain constant DM surface density (~ 75 M_☉/pc²) inferred from observations of satellite galaxies (Burkert, 2020) (see also Safarzadeh & Spergel (2020)):



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Catch-22

We need simulations to investigate:

- possibility additional nonlinear structure formation?
- modifications to soliton profile in context of cosmology, baryons

Fuzzy Cosmological Simulations

* (Mocz & Succi, 2015), (Mocz et al., 2017), (Mocz et al., 2020), Mocz+ (2019) Phys. Rev. Lett.





Cosmological simulations

- Full-physics (baryons, star formation, feedback) quantum mechanical simulations with quantum wave effects
- Proper initial conditions at z = 127 from AxionCAMB
- ▶ 5 million core hours on Stampede2 and Odyssey



- Limitation: method is memory-expensive (need to resolve kpc interference)
- Restricted to study of first galaxies/structures at z ~ 6, small box size (~ 2 Mpc)

Cosmological simulations – dark matter



Small scales: anatomy of a filament



0.5 Mpc

Summary



JWST Mock Images



Collapse of cylindrical filament



we identified a nonlinear structure formation channel:
 in cosmological context, filaments form first
 cylindrical 'soliton' core unstable to spherical collapse

DM power spectrum



- quantum pressure tensor adds extra suppression of small-scale power
- but we found extra power from interference at 10s kpc by z = 7
- agreement with CDM above 1 Mpc

Conclusions

cosmological first objects summary

- First galaxies uniquely probe the physical nature of dark matter
- Future missions (e.g. JWST) will open an observational window into this emergent world
- Nonlinear structure formation eases Catch-22 problem
- In FDM:
 - Primordial stars form along dense dark matter filaments
 - Dark matter filaments show coherent interference patterns
 - Dark matter filaments develop cylindrical soliton-like cores which are unstable under gravity and collapse into spherical solitons
 - Gas & stars trace cored dark matter profile

Numerical Method: (**Mocz** et al., 2017)

2nd-order unitary spectral leap-frog scheme. 'Kick-drift-kick'
Calculate potential:

$$V = \operatorname{ifft} \left[-\operatorname{fft} \left[4\pi G(\rho - \overline{\rho}) \right] / k^2 \right]$$
(2)

Potential half-step 'kick':

$$\psi \leftarrow \exp\left[-i(\Delta t/2)(m/\hbar)V\right]\psi \tag{3}$$

► Full 'drift' (kinetic) step in Fourier-space:

$$\hat{\psi} = \operatorname{fft}[\psi]$$
 (4)

$$\hat{\psi} \leftarrow \exp\left[-i\Delta t(\hbar/m)k^2/2\right]\hat{\psi}$$

$$\psi \leftarrow \operatorname{ifft}\left[\hat{\psi}\right]$$
(6)

Another 'kick'

(*unitary algorithm adaptable to quantum computers)



Numerical Method: Gas (Mocz et al., 2014; Mocz,

Vogelsberger & Hernquist, 2014; Mocz et al., 2015, 2016; Mocz, 2017)

Moving Mesh Magnetohydrodynamics





Schrödinger/Vlasov–Poisson correspondence

Do the 3D Schrödinger equations encode collisionless dynamics (6D)?

$$i\hbar \frac{\partial \psi}{\partial t} = -\frac{\hbar^2}{2m} \nabla^2 \psi + m V \psi$$

$$\iff (?)$$

$$\frac{\partial f}{\partial t} + \mathbf{v} \cdot \frac{\partial f}{\partial \mathbf{x}} - \nabla V \cdot \frac{\partial f}{\partial \mathbf{v}} = \mathbf{0}$$
(8)

In Mocz et al. (2018) we explore the limiting behavior of large boson mass (e.g., QCD axion)

Schrödinger/Vlasov–Poisson correspondence



3D wave function can encode 6D distribution function: $\psi(\mathbf{x}) \propto \sum_{\mathbf{v}} \sqrt{f(\mathbf{x}, \mathbf{v})} e^{im\mathbf{x}\cdot\mathbf{v}/\hbar + 2\pi i\phi_{\text{rand},\mathbf{v}}} d^3v$

Schrödinger/Vlasov–Poisson correspondence





Conclusions

SP-VP correspondence summary

- ► classical limit for the gravitational potential V recovered as $\mathcal{O}(m^{-2})$ (\Rightarrow forces as $\mathcal{O}(m^{-1})$), while density has $\mathcal{O}(1)$ interference patterns on scale of λ_{dB}
- soliton cores regularize caustic singularities
- fuzzy halos are NFW-like with a soliton core

Future Work

Larger statistical simulations samples

- need approximate particle-based simulation methods
 Fuzzy Dark Matter
- Simultaneously constrain particle mass & strong-CP symmetry-breaking scale using a variety of techniques
 - Lyman- α
 - nonlinear halo mass function
 - dynamics in a soliton core
 - soliton black hole interaction
 - stellar streams
 - dynamical heating/timing problems
 - gravitational lensing

Return to Catch-22

- Catch-22 Problem: moderate tension in setting the particle mass to simultaneously capture large cores and right amount of substructure
- ► Resolutions:
 - Dark matter is not fuzzy
 - FDM wave interference
 - Substructure forms below exponential cutoff (nonlinear effects)
 - Baryonic physics/feedback (connection between galaxies and their dark matter halos (Wechsler & Tinker, 2018))
 - Attractive self-interaction that arises from strong-CP symmetry-breaking scale (Desjacques, Kehagias & Riotto, 2018)

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Strong-CP symmetry-breaking scale

Starting point

$$S[\phi] = \int d^4x \sqrt{-g} \left[\frac{1}{2} (\partial \phi)^2 - \left(m^2 f^2 \right) \left(1 - \cos \frac{\phi}{f} \right) \right]$$
(9)

 $f \simeq 10^{17} \text{ GeV}$ is the decay constant, $m^2/f^2 = 10^{-96} \leftarrow tiny!$

⇒ Gross-Pitaevskii-Poisson equations in expanding universe

$$i\hbar\left(\frac{\partial}{\partial t} + \frac{3}{2}H\right)\psi = -\frac{\hbar^2}{2m}\nabla^2\psi + mV\psi - \frac{4\pi\hbar^2a_s}{m^2}|\psi|^2\psi \qquad (10)$$
$$\nabla^2 V = 4\pi G(\rho - \overline{\rho}) \qquad (11)$$

 a_s is effective *s*-scattering cross section

Strong-CP symmetry-breaking scale

Phase-transition: dilute to dense solitons above a critical mass



increasing self-interaction \rightarrow



Work in progress with undergrad Noah Notis

Strong-CP symmetry-breaking scale



FDM: $a_s = 0$ $a_s = 1 \cdot 10^{-75}$ cm



 $a_{\rm s} = 2 \cdot 10^{-75} {\rm \ cm}$



(not fully-converged @prelim res.)



relative power-spectrum

Small-scale features of FDM Student work highlights

FDM dynamical friction Lachlan Lancaster, Cara Giovanetti,

Mocz+ 2019





Applications:

- satellite infall / timing problem
- final parsec problem



gravity vs quantum pressure

- gravity vs velocity dispersion
 - random walk
- gravity vs mixed



Small-scale features of FDM Student work highlights

Soliton + SMBH

Elliot Davies & Mocz (2020)



#AAS235 Honolulu, HI



► BH makes soliton more compact, ⇒ hydrogen Bohr solution



Small-scale features of FDM Student work highlights

FDM dynamical heating



 FDM interference fluctuations heat Milky Way old stellar disk



(Mocz et al., 2017)

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