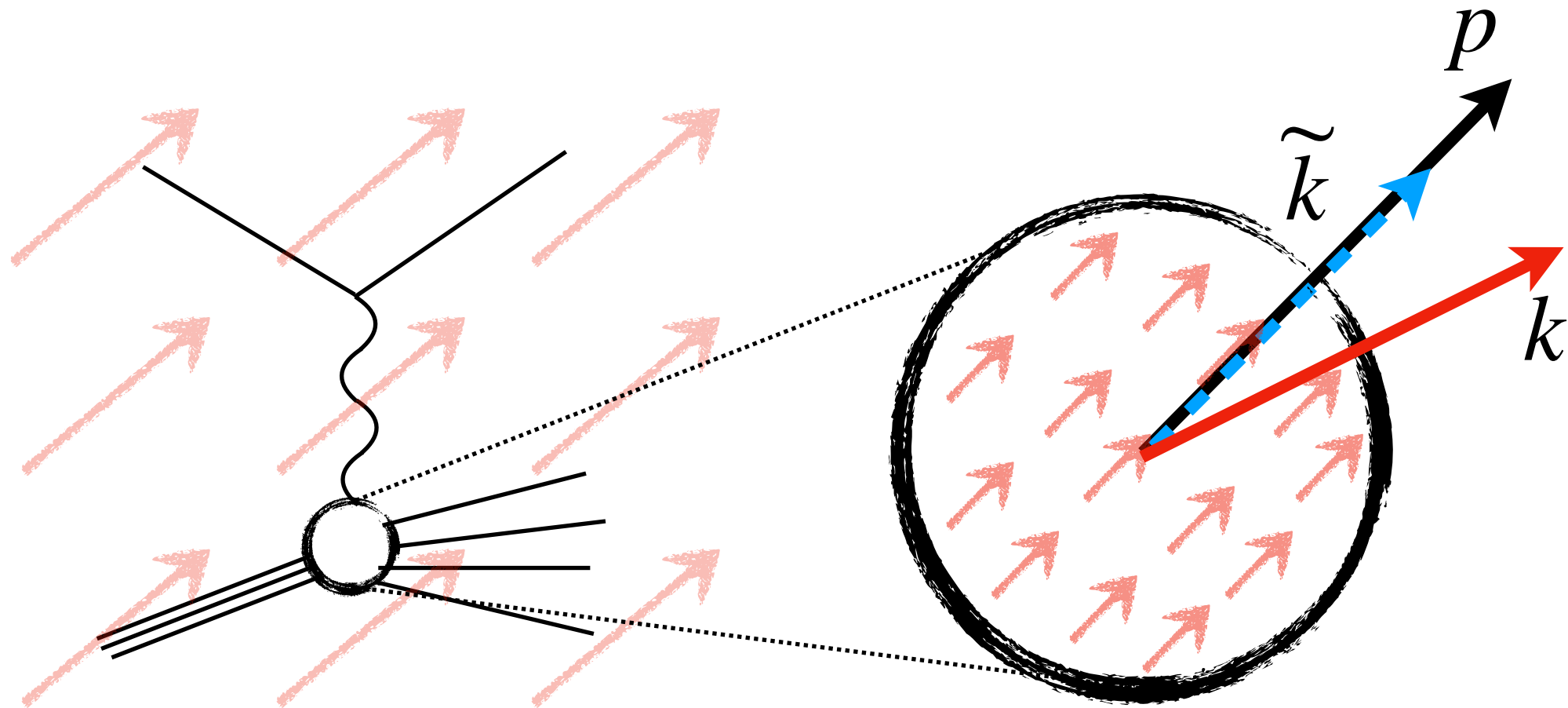


Lorentz violation analysis with ZEUS data



Nathan Sherrill
Indiana University
ZEUS meeting, January 2020

In collaboration with Enrico Lunghi



<http://www.indiana.edu/~iucss/>



Talk overview

What is Lorentz violation?

How to search for Lorentz violation?

Effects on high-energy hadrons

Application: deep inelastic scattering

Estimates for colliders

An analysis with real data

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Based on: arXiv:1911.04002; PRD **98**, 115018 (2018); PLB **769**, 272 (2017)

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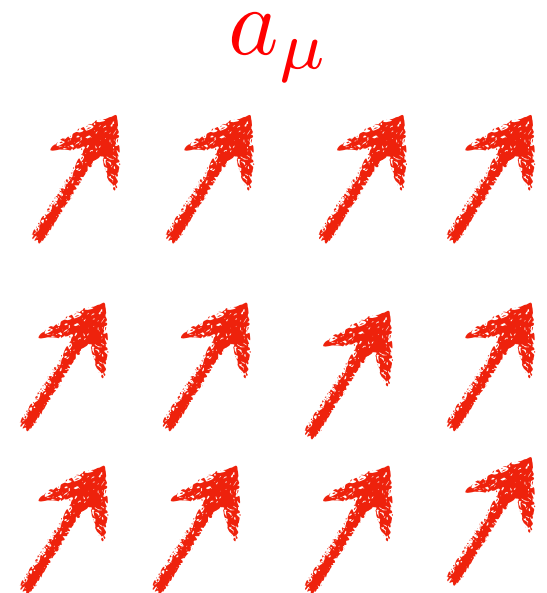
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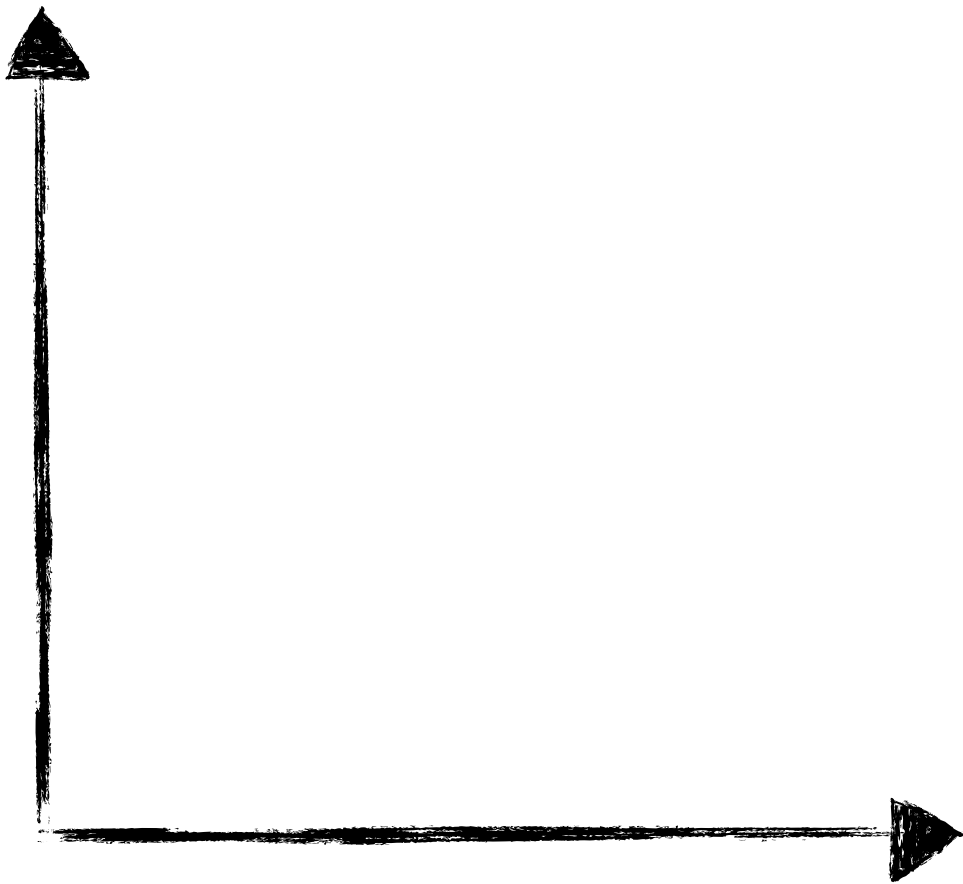
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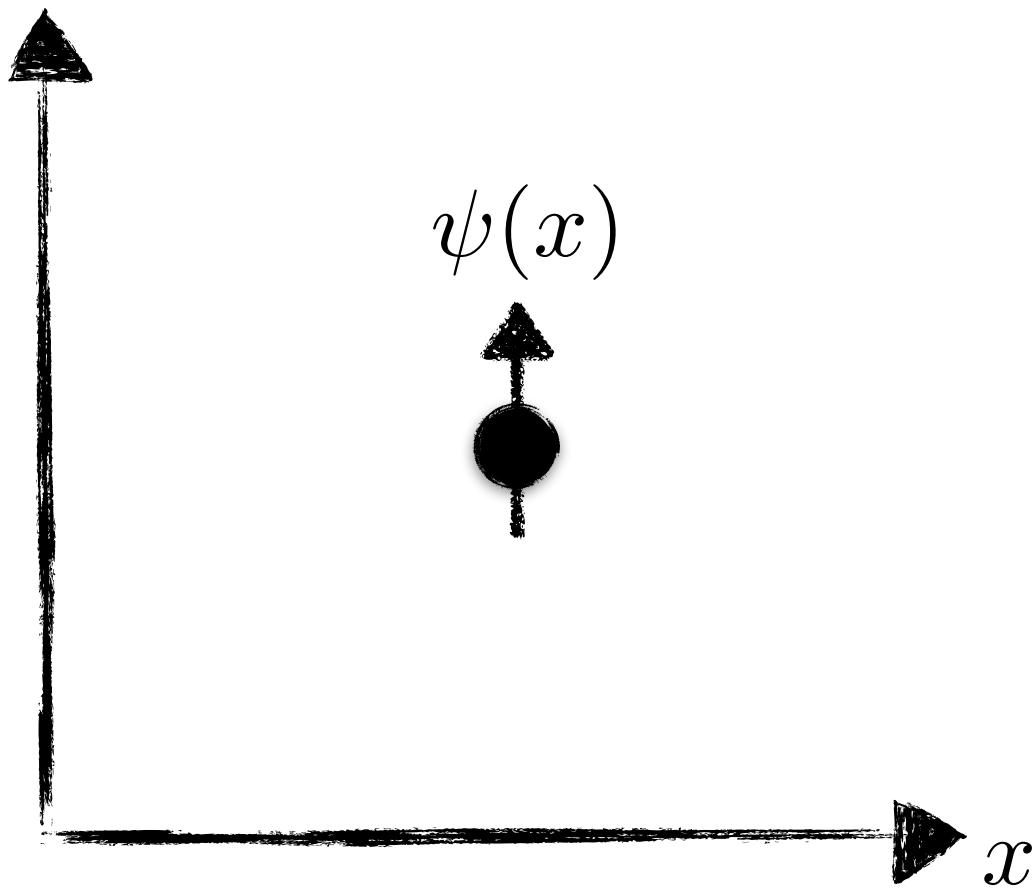
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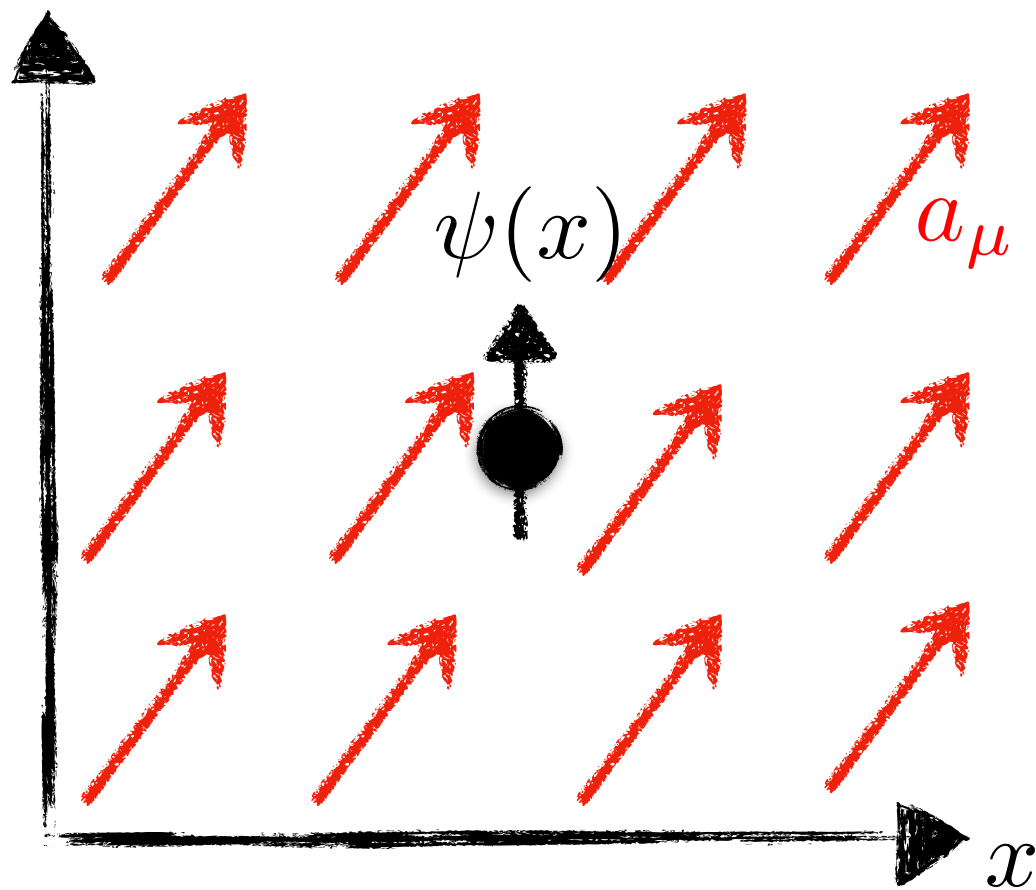
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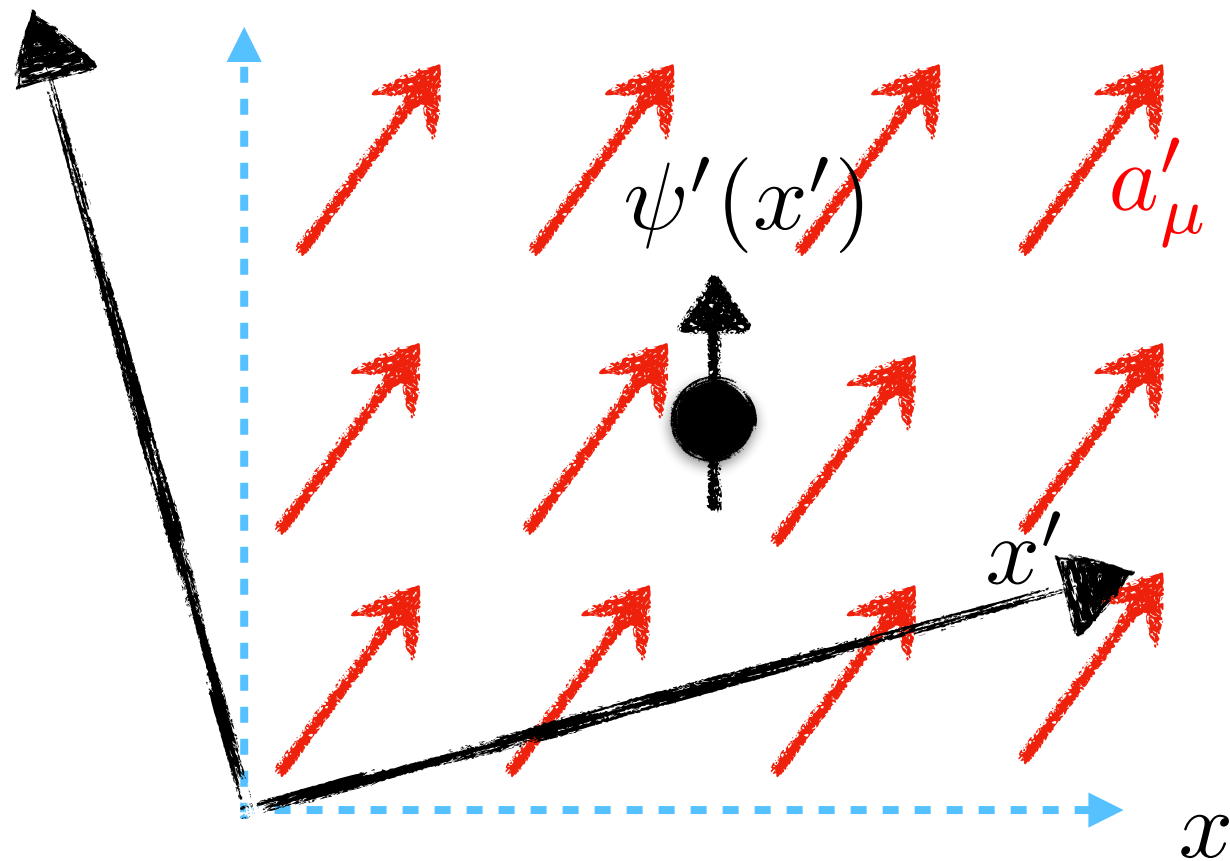
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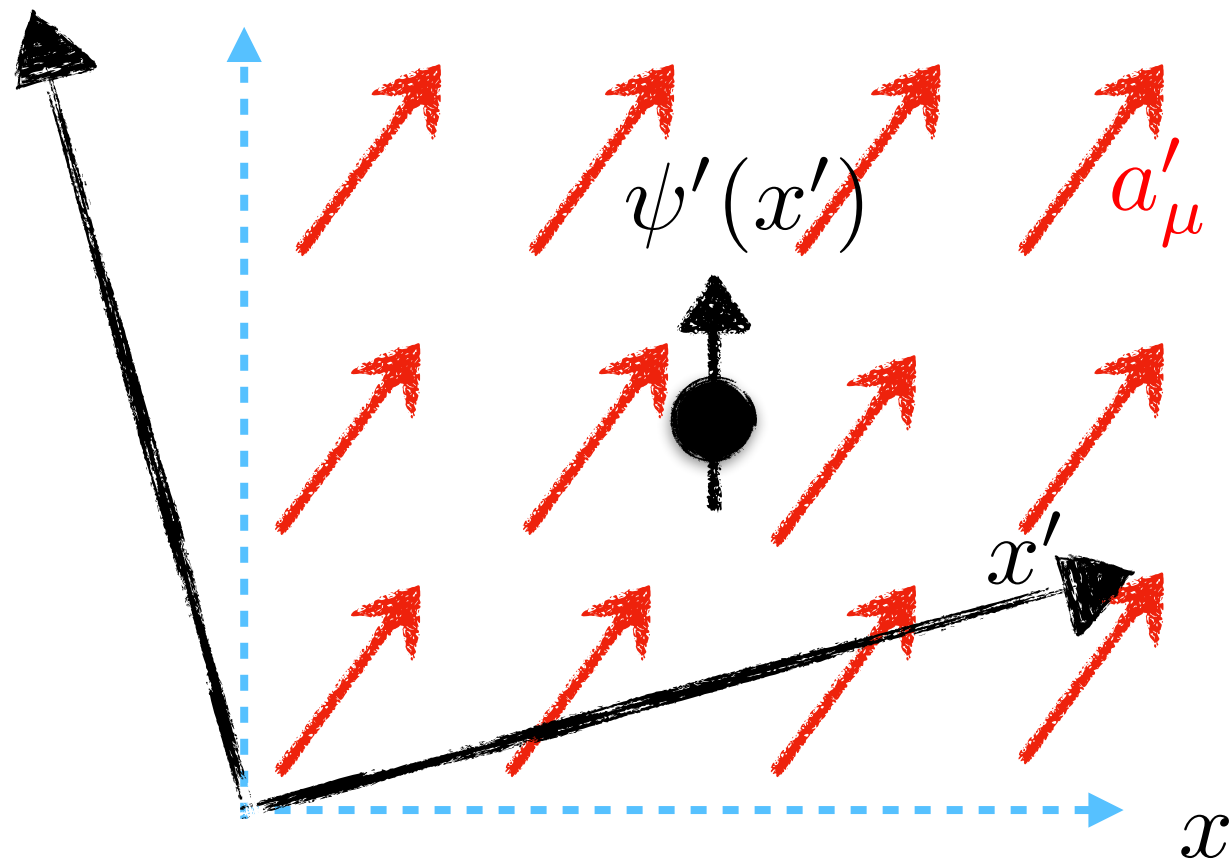
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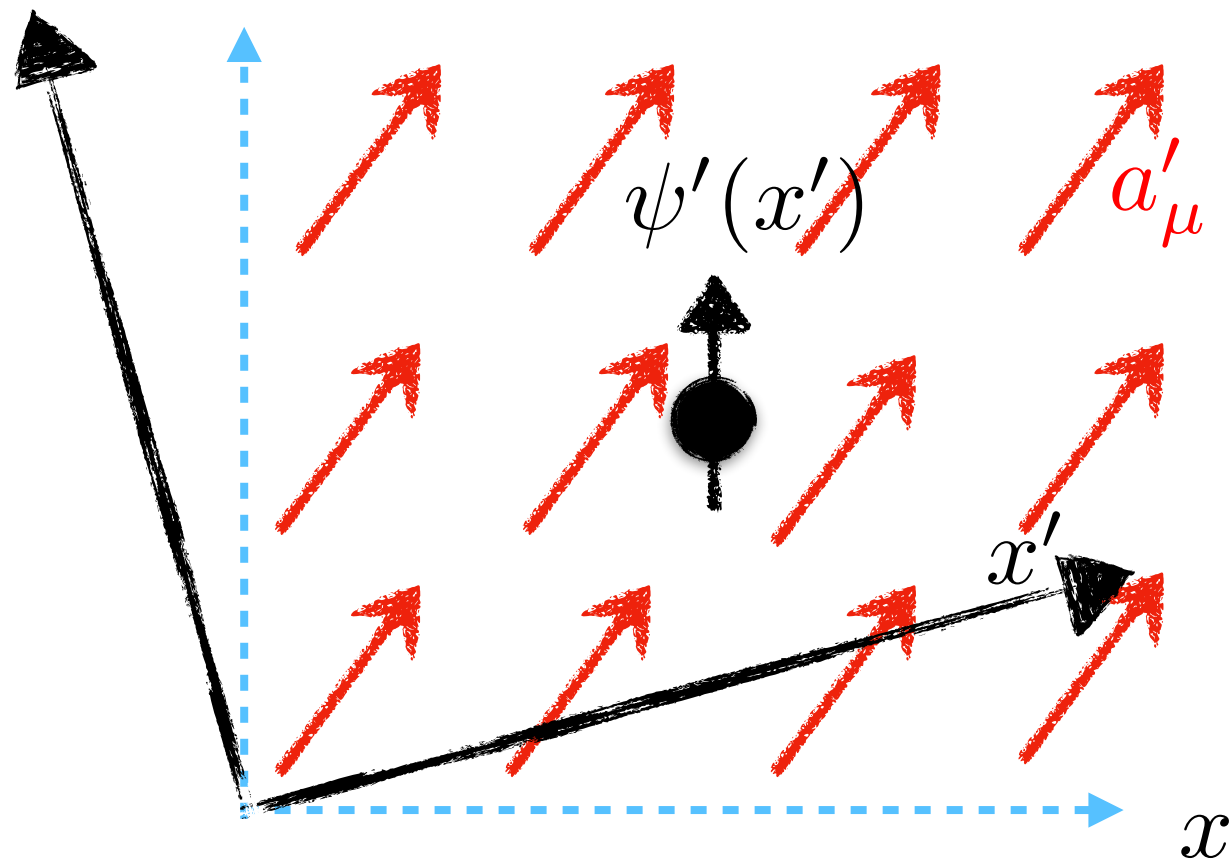
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Under an OLT the background a_μ transforms like an ordinary four vector

Hence, there is no change in the physics; the background cannot be seen by performing observer transformations (changing coordinates)

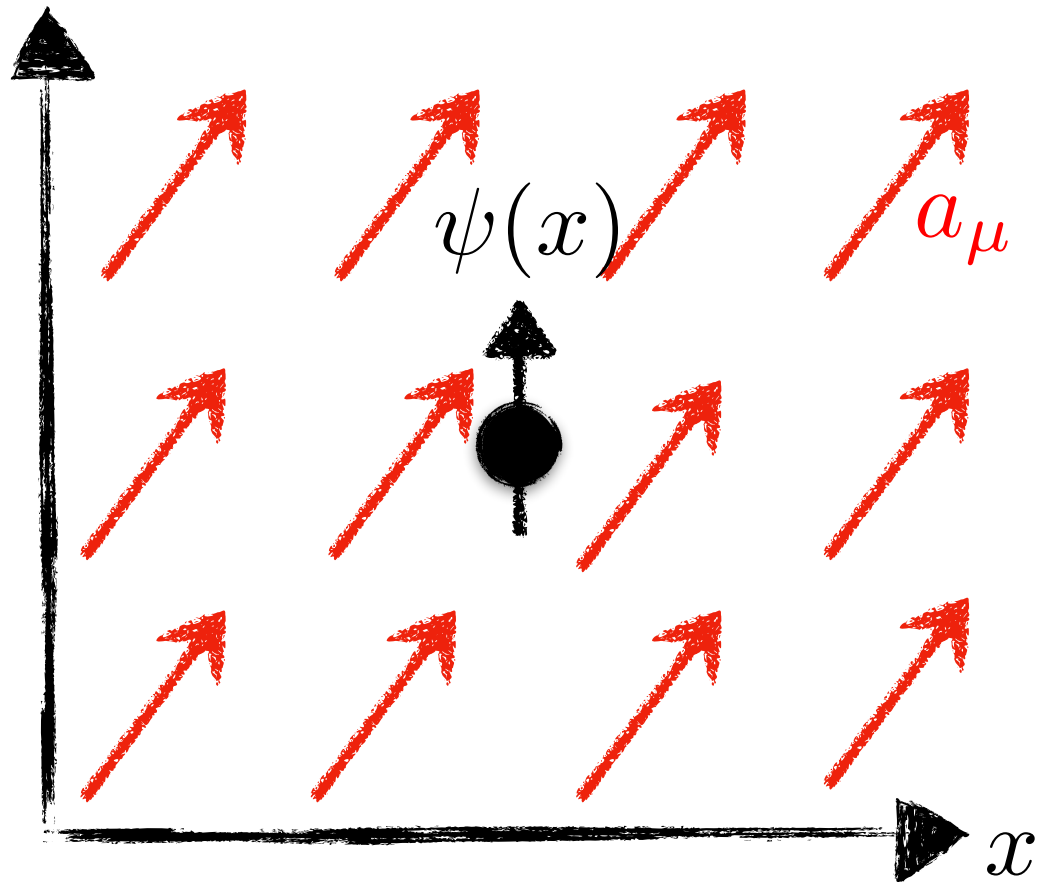
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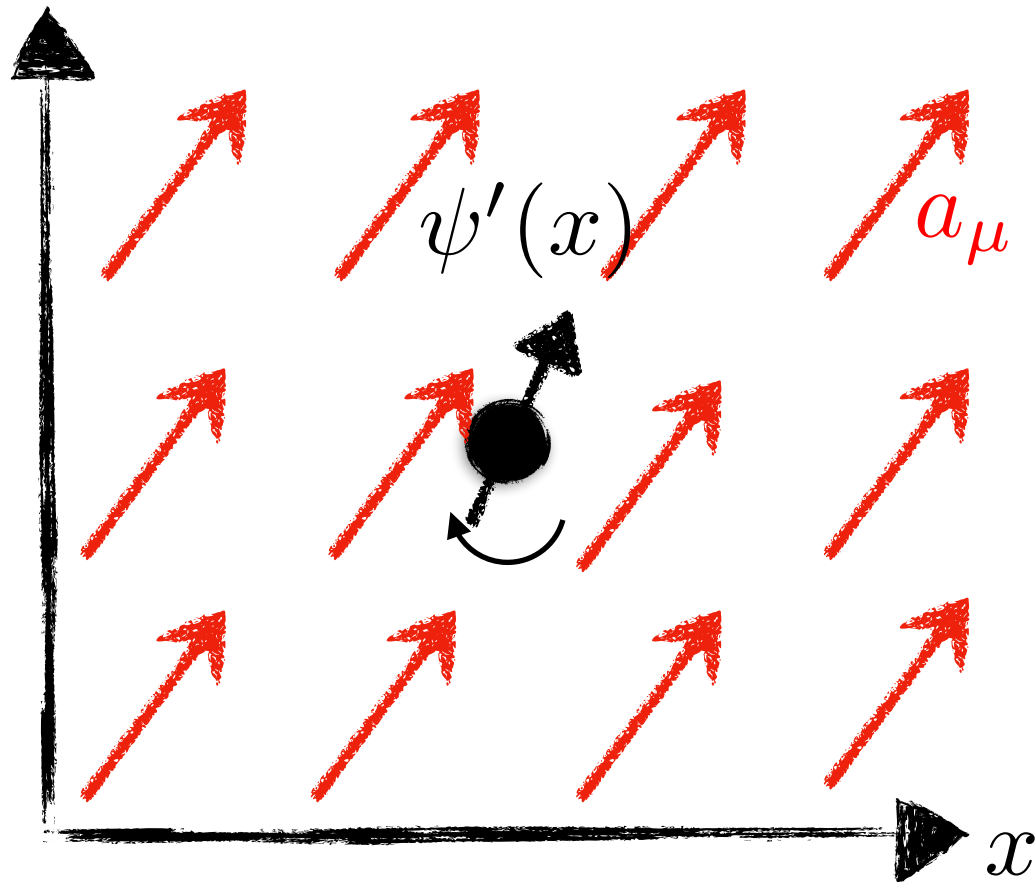
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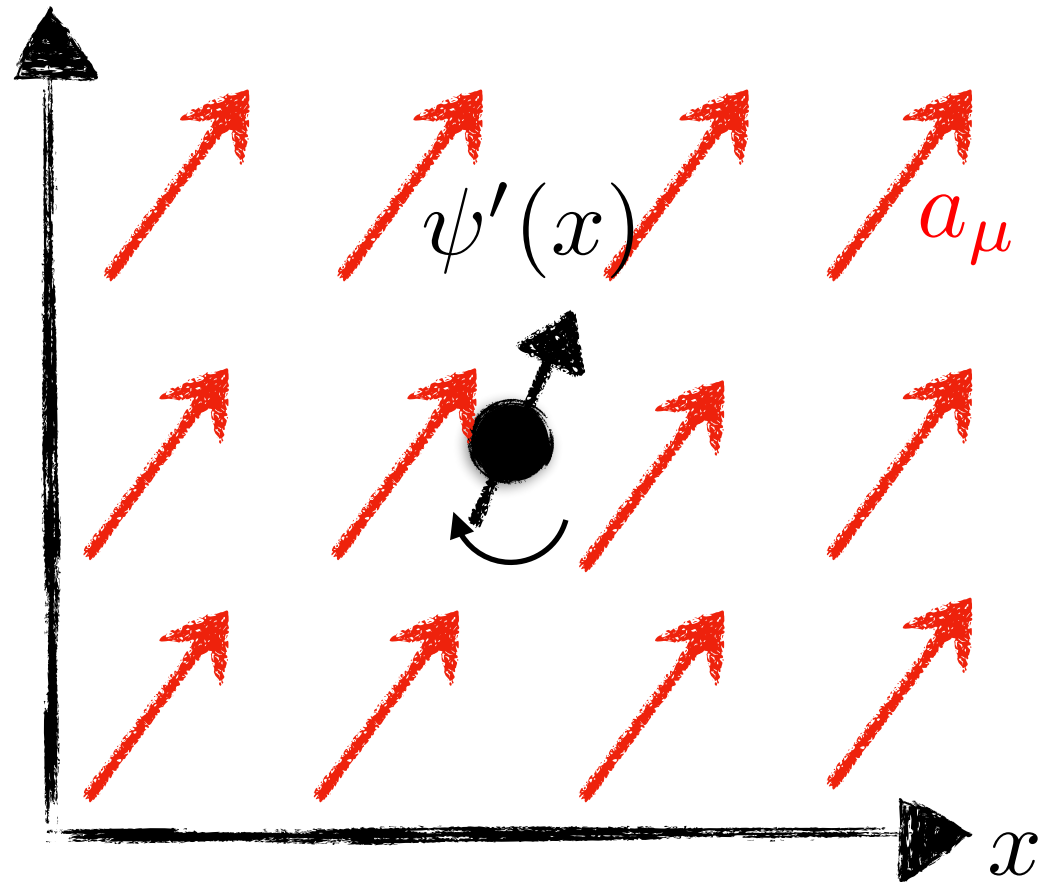
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$$a_\mu \rightarrow a_\mu$$

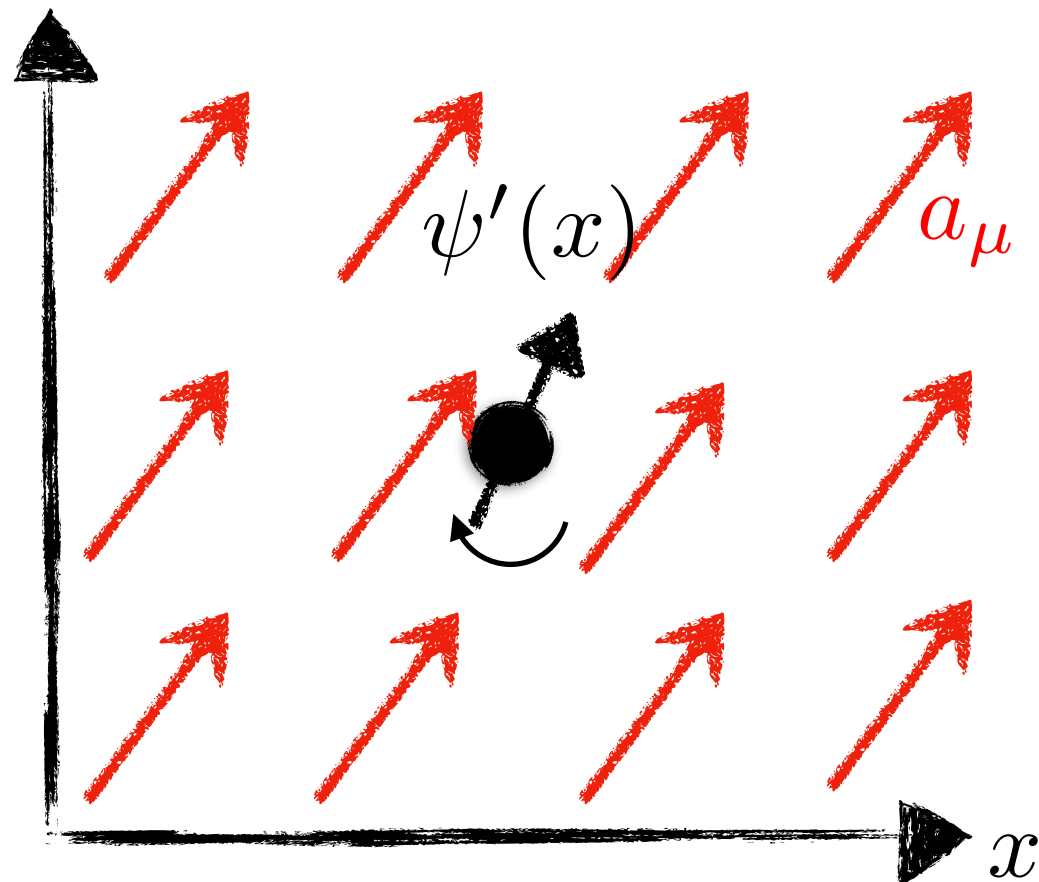
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Net physical effect

$$\begin{aligned} -a_\mu \bar{\psi} \gamma^\mu \psi &\rightarrow -(\Lambda^{-1})_{\mu\nu} a^\nu \bar{\psi} \gamma^\mu \psi \\ &\neq -a_\mu \bar{\psi} \gamma^\mu \psi \end{aligned}$$

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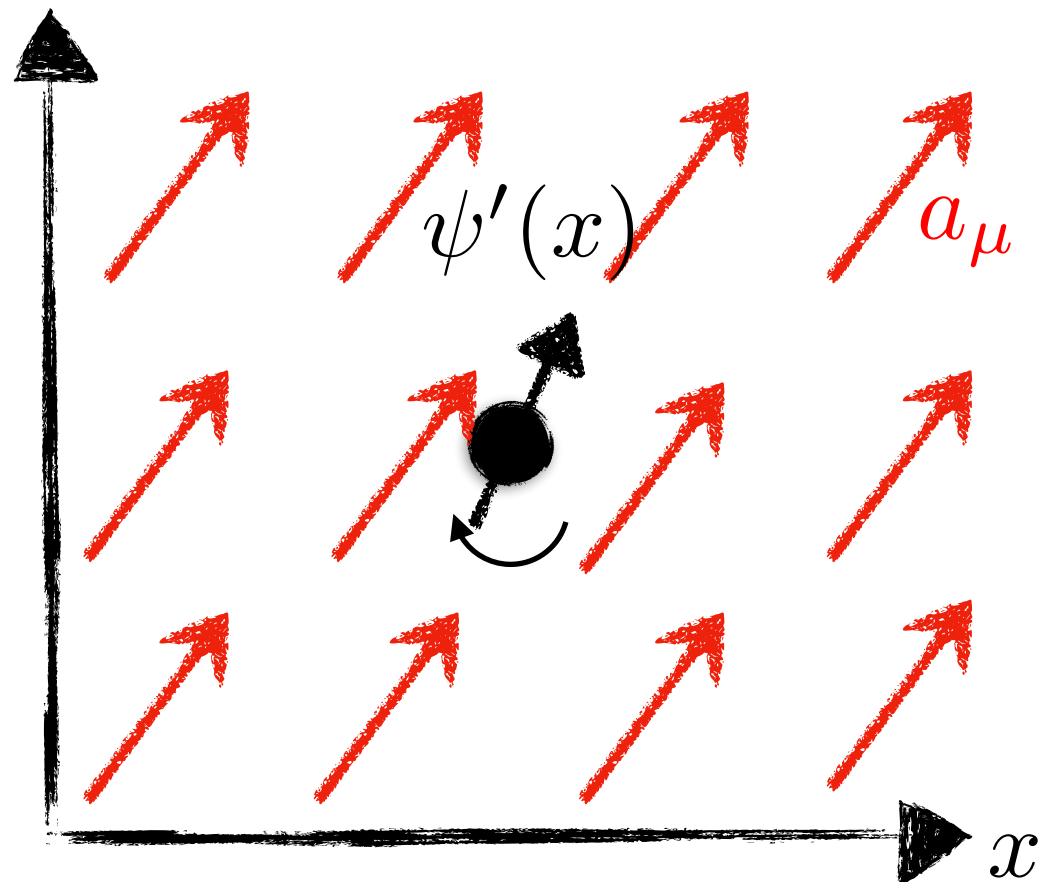
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Unlike OLTs, PLTs can produce physical effects as a result of the background

The rotated system obeys a different
physical law than the same system with
rotated coordinates

\Rightarrow Lorentz violation!

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We use a model-independent, effective field theory framework: the Standard-Model Extension (SME)*

*D. Colladay, V. A. Kostelecký, PRD 55, 6760 (1997); PRD 58, 1166002 (1998)

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- “Coefficients for Lorentz violation”
- Observer Lorentz tensors
- Coupling constants
- Necessarily small (perturbative)
- Experimentally accessible!

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Data Tables for Lorentz and CPT Violation

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January 2020 update of *Reviews of Modern Physics* **83**, 11 (2011) [arXiv:0801.0287]

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Table D19. Nonminimal photon sector, $d = 7$

Combination	Result	System	Ref.
$ \sum_{jm} Y_{jm}(110.47^\circ, 71.34^\circ) k_{(V)jm}^{(7)} $	$< 2 \times 10^{-6} \text{ GeV}^{-3}$	Spectropolarimetry	[170]
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Much of the QCD sector is yet to be explored!

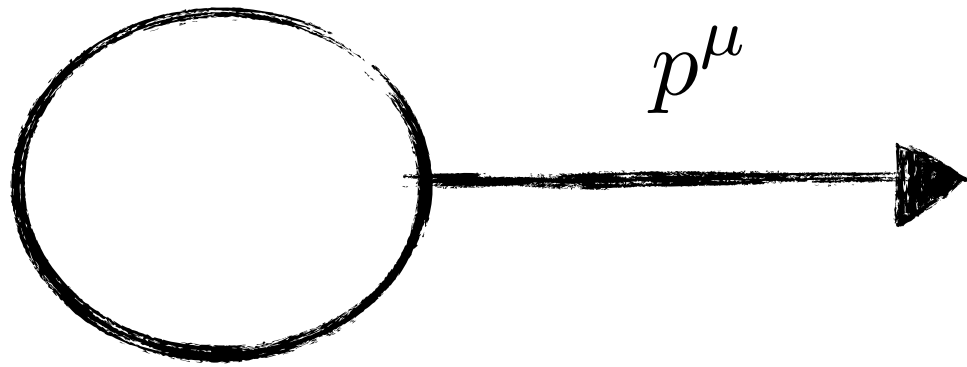
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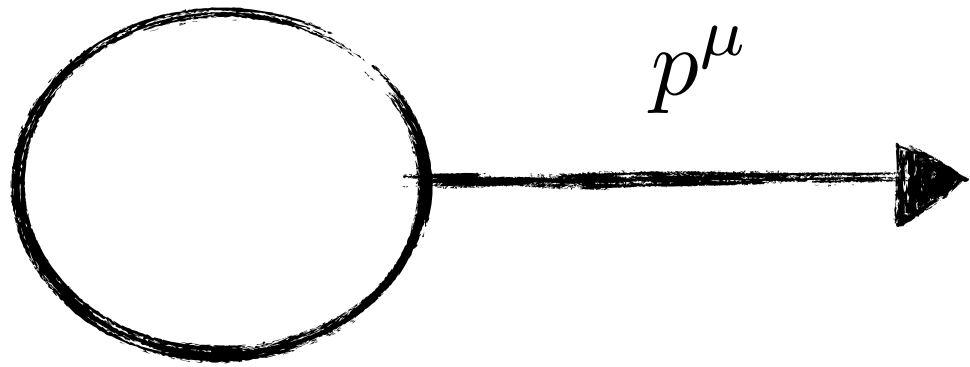
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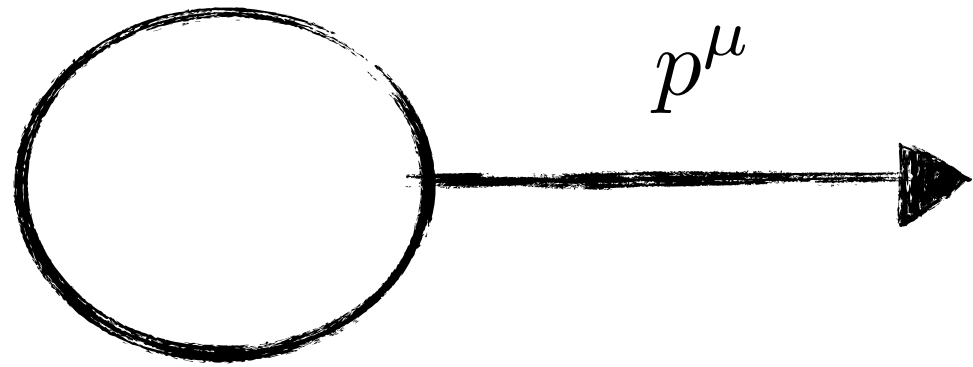
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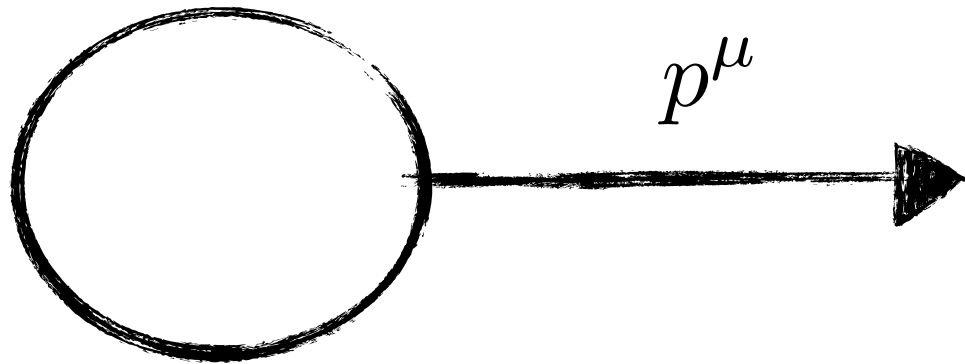
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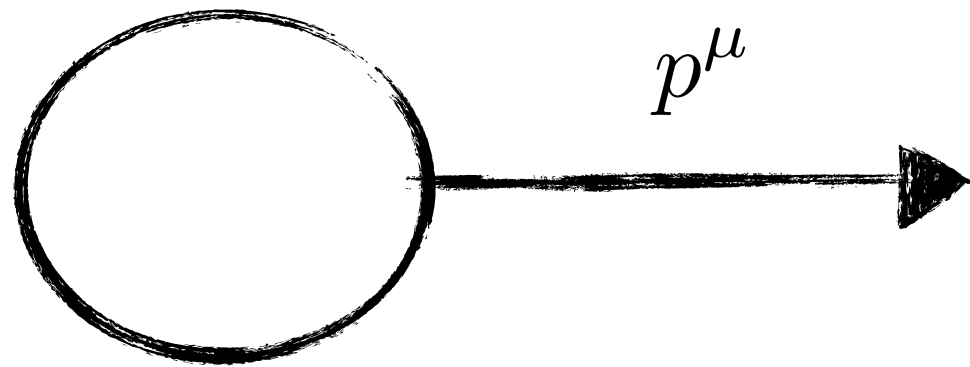
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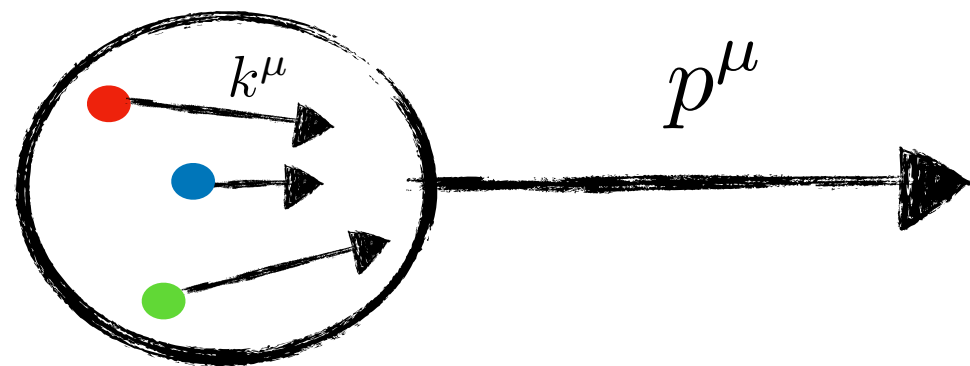
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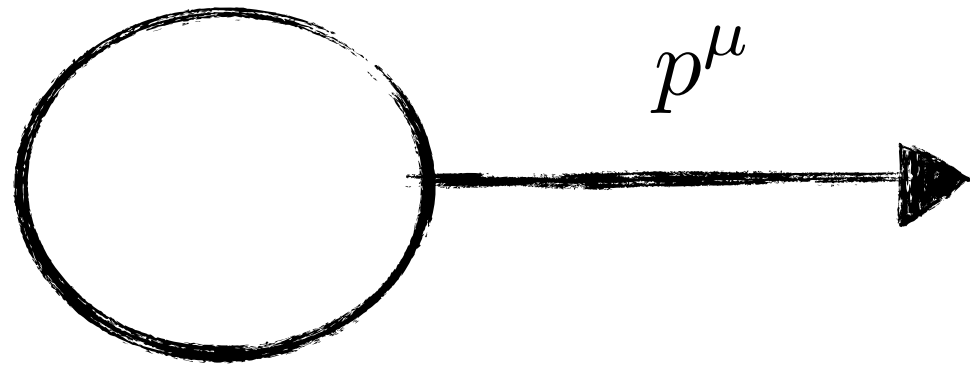


A diagram showing a large circle representing a hadron. Inside the circle are three smaller colored circles (red, blue, and green) representing partons. Each parton has a horizontal arrow pointing to the right, labeled k^μ . A larger horizontal arrow points to the right from the center of the hadron circle, labeled p^μ .

$$k^\mu \sim \left(p^+, \frac{M^2}{2p^+}, M \right) + \mathcal{O}(M/p^+)$$

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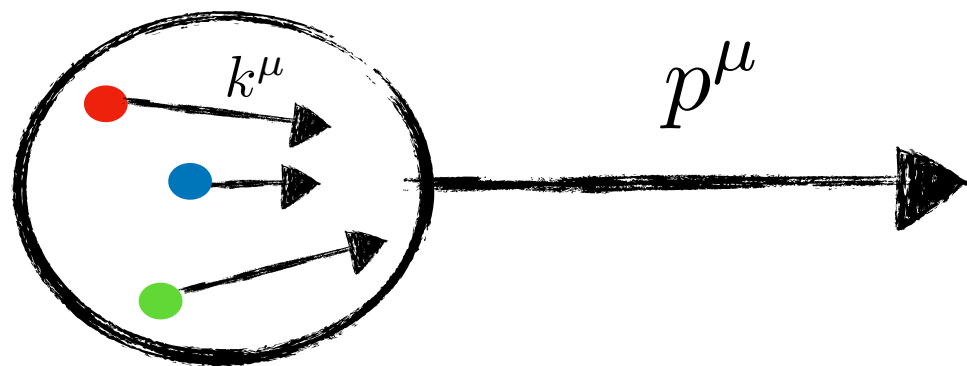
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$$k^\mu \sim \left(p^+, \frac{M^2}{2p^+}, M \right) + \mathcal{O}(M/p^+)$$

Fraction of plus momenta is boost invariant, leading to familiar parameterization for high-energy, massless, on-shell partons within hadrons

$$\xi \equiv k^+ / p^+$$

$$k^\mu = \xi p^\mu$$

Covariant expression; can be used in any frame

Quark-sector Lorentz-violating effects

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Massless quarks modified by Lorentz-violating effects

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We consider the following (spin-independent, flavor-diagonal) effects

$$\begin{aligned} \mathcal{L} = \sum_{f=u,d} & \frac{1}{2} \bar{\psi}_f \gamma^\mu i D_\mu \psi_f + \frac{1}{2} \left(c_f^{(4)} \right)^{\mu\nu} \bar{\psi}_f \gamma_\mu i D_\nu \psi_f \\ & - \left(a_f^{(5)} \right)^{\mu\alpha\beta} \bar{\psi}_f \gamma_\mu i D_{(\alpha} i D_{\beta)} \psi_f + \text{h.c.} \end{aligned}$$

Quark-sector Lorentz-violating effects

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Modified Dirac equation

$$[(\eta^{\mu\nu} + c_f^{\mu\nu})\gamma_\mu i\partial_\mu - a_f^{(5)\mu\alpha\beta}\gamma_\mu i\partial_\alpha i\partial_\beta]\psi_f = 0$$

Quark-sector Lorentz-violating effects

Modified Dirac equation

$$[(\eta^{\mu\nu} + c_f^{\mu\nu})\gamma_\mu i\partial_\mu - a_f^{(5)\mu\alpha\beta}\gamma_\mu i\partial_\alpha i\partial_\beta]\psi_f = 0$$

Dispersion relation

$$\tilde{k}^2 = k^2 + \mathcal{O}(\text{coefficients}) = 0$$

$$E^2 = |\vec{k}|^2 + \mathcal{O}(\text{coefficients})$$

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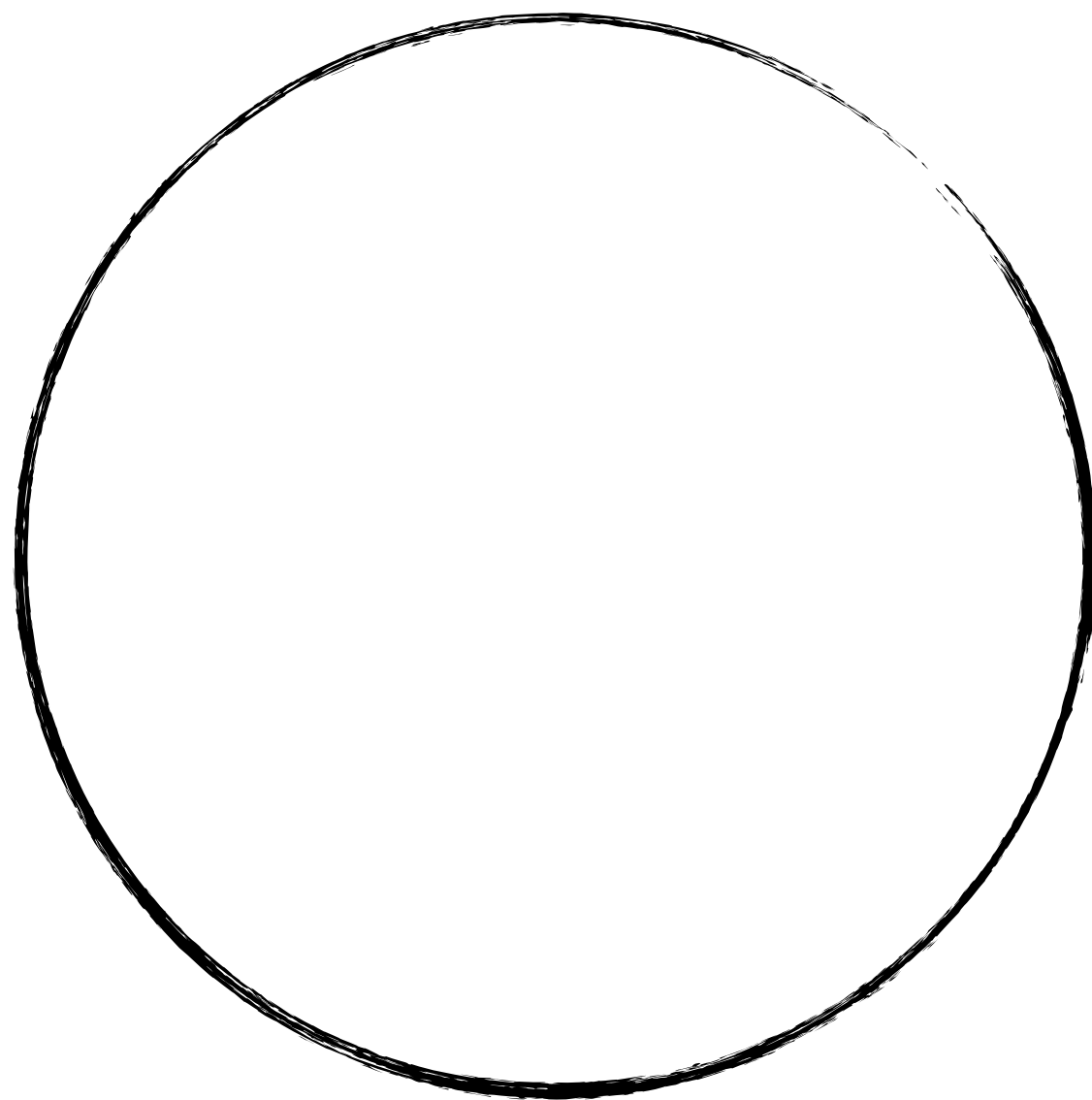
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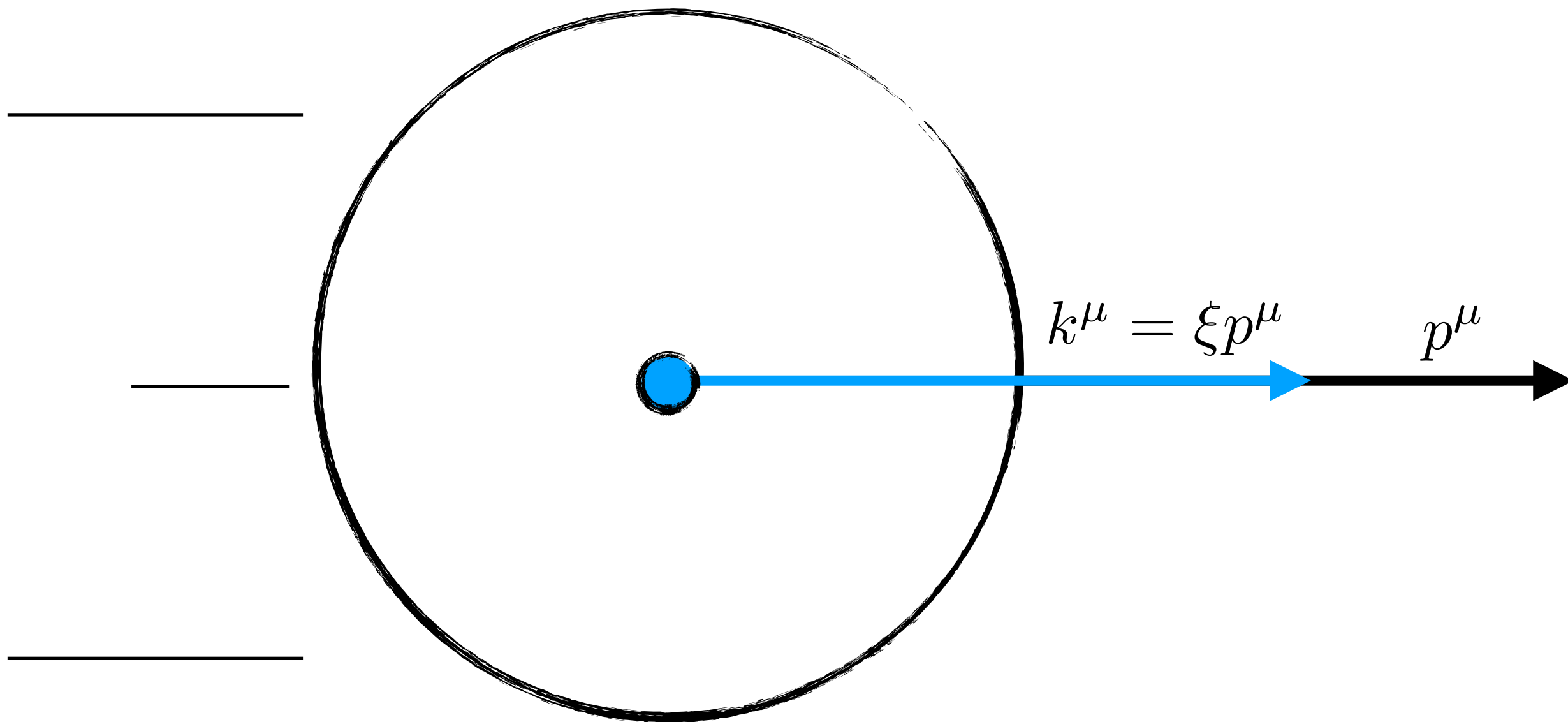
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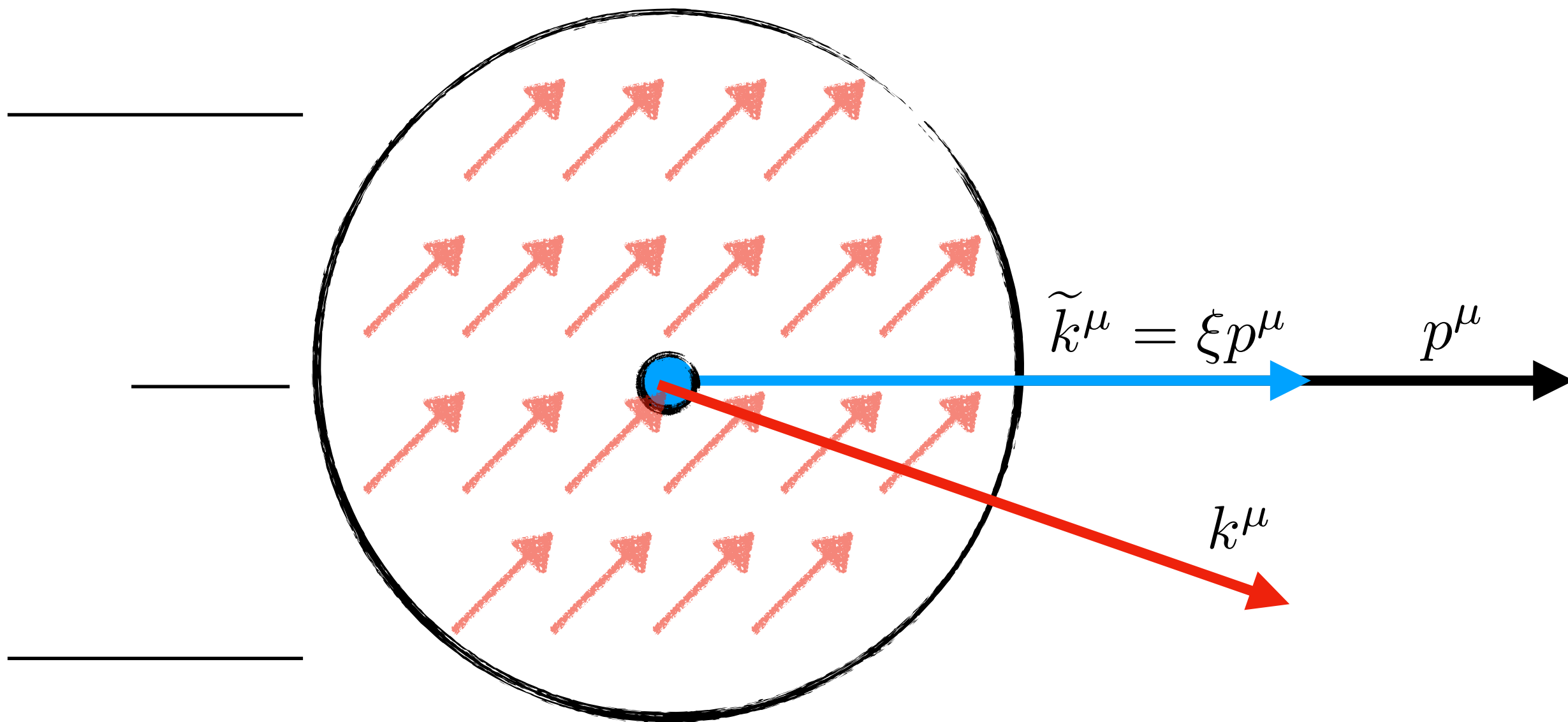
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Instead, for a covariant definition to be retained $\tilde{k}^\mu = \xi p^\mu$

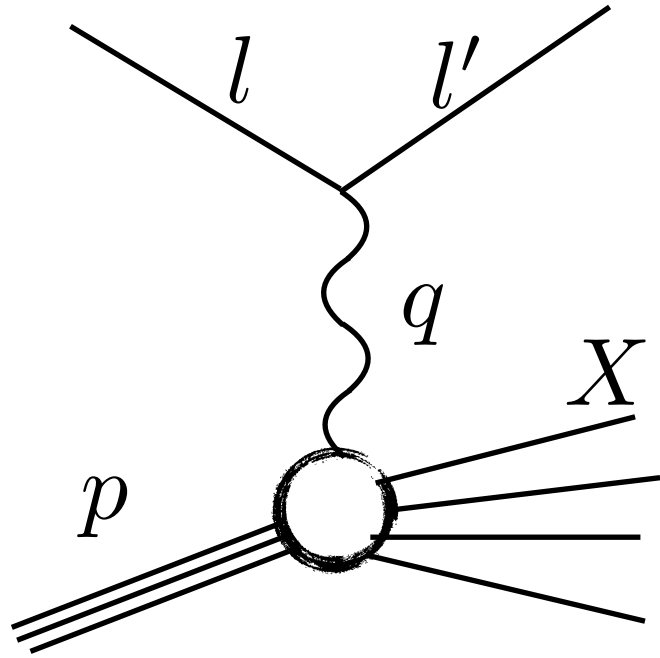




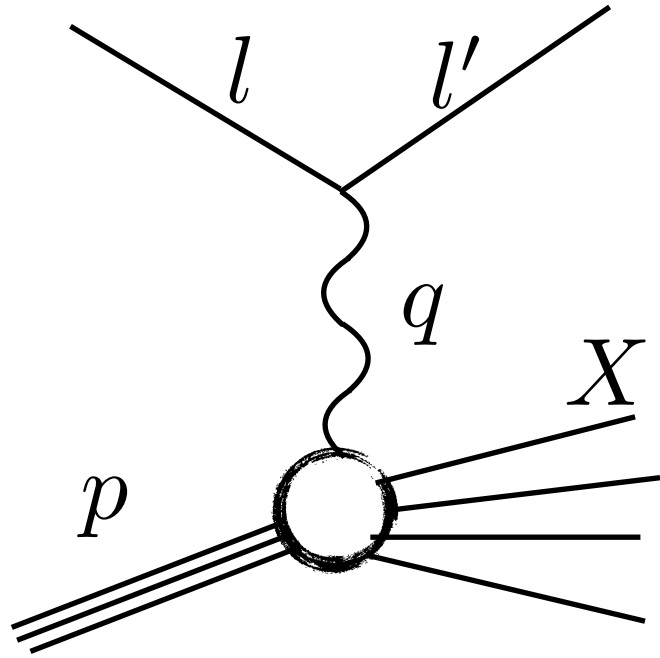


Application I: deep inelastic scattering (DIS)

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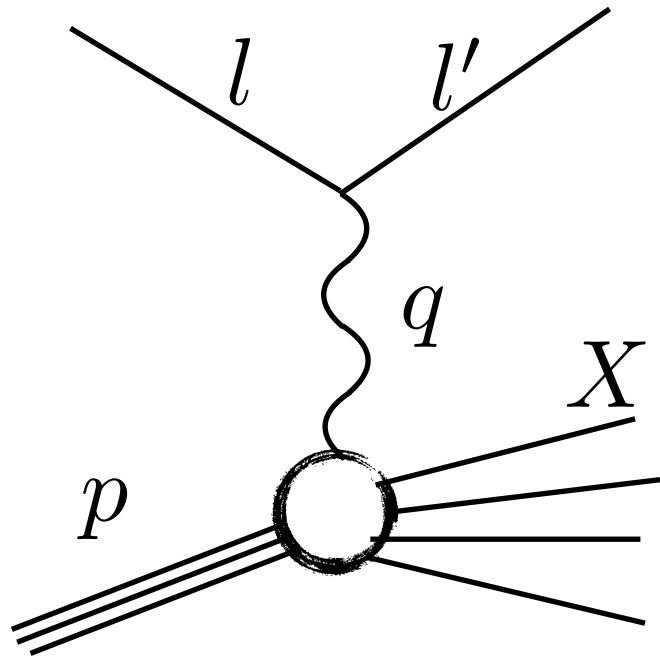
Application I: deep inelastic scattering (DIS)



$$T_{\mu\nu} = i \int d^4w e^{iq \cdot w} \langle p, s | T j_\mu^\dagger(w) j_\nu(0) | p, s \rangle_c$$

$$\sigma \propto L_{\mu\nu} \text{Im} T^{\mu\nu}$$

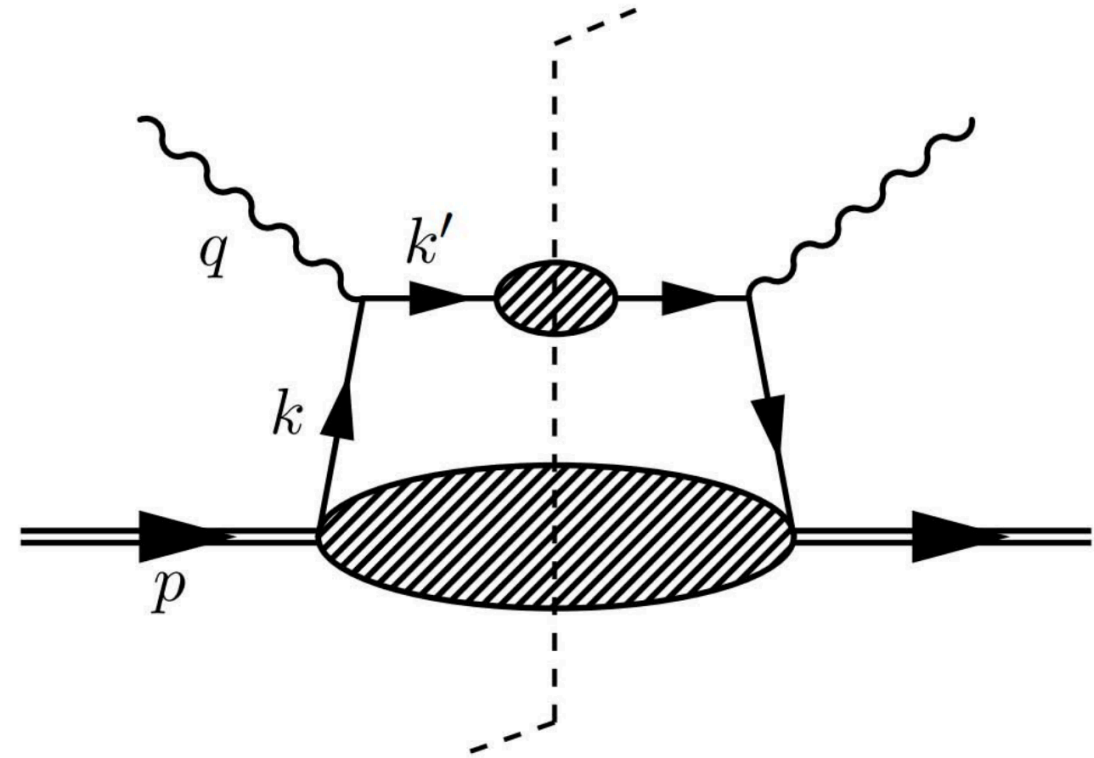
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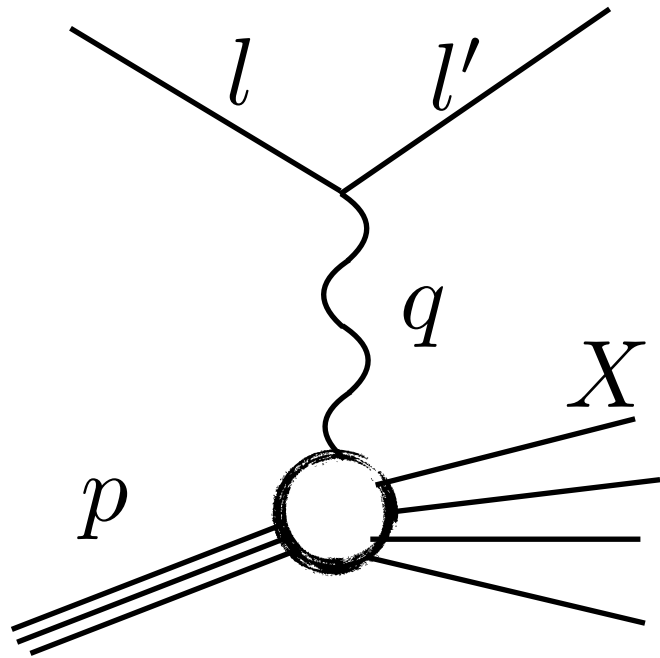
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Factorization in DIS limit most simply shown
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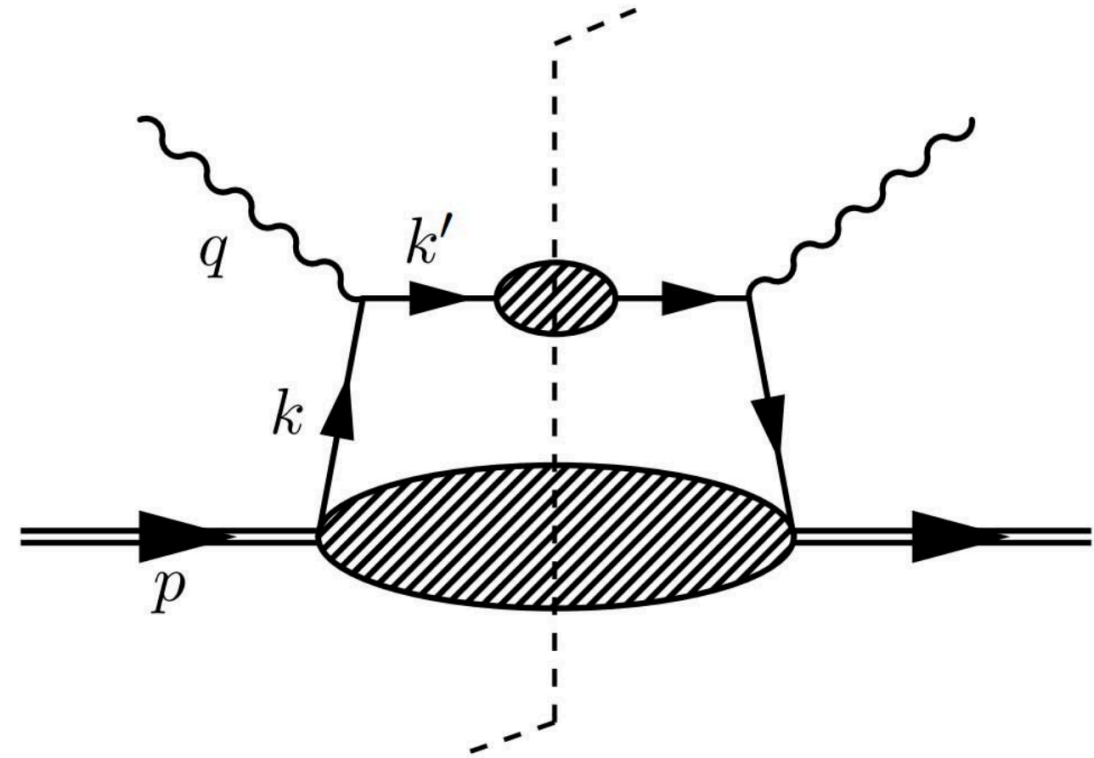
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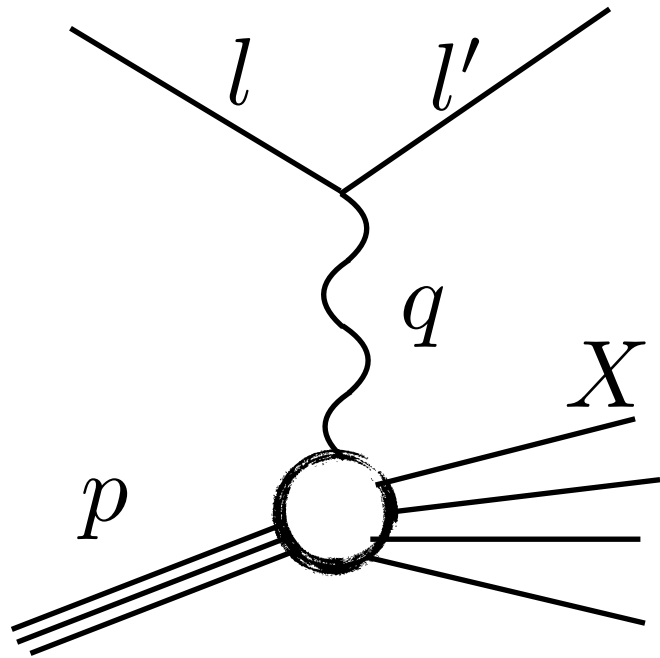
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$$\sigma \sim \int d\xi \sigma_{\text{parton}}(\xi) f(\xi) + \text{small corrections}$$

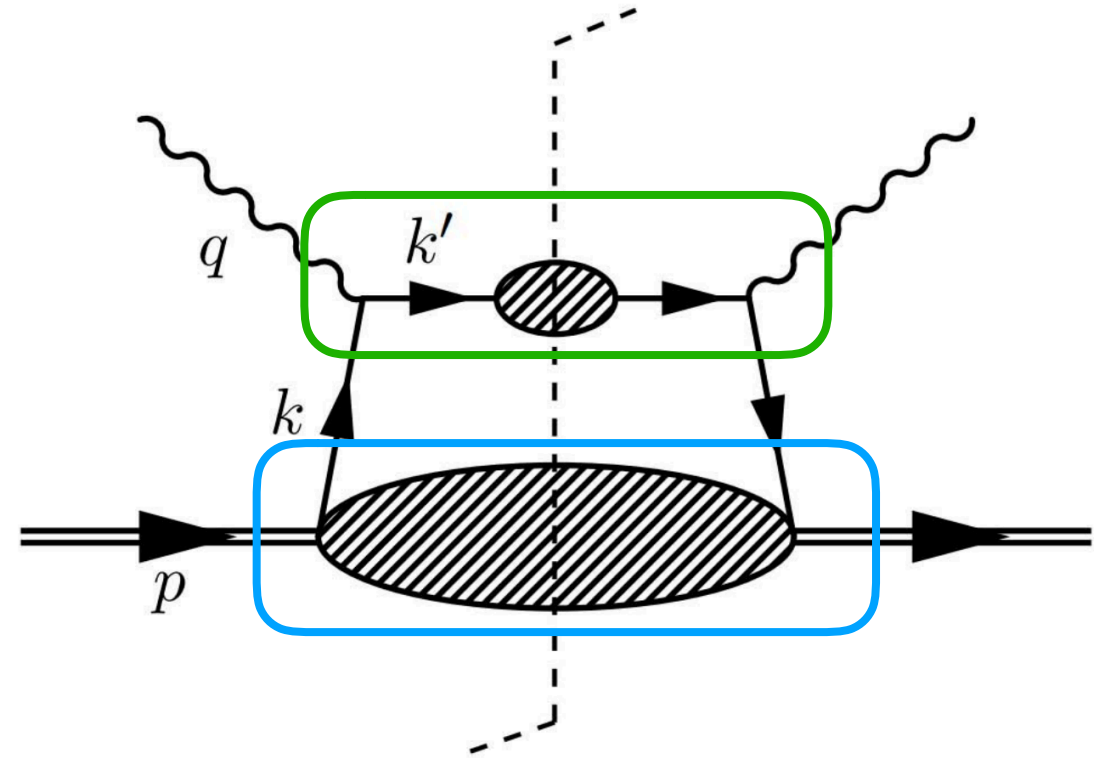
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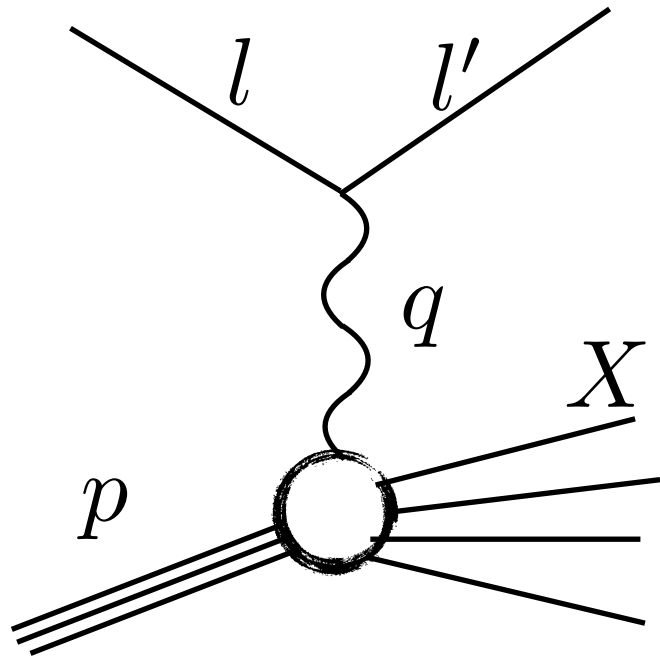
$$\sigma \sim \int d\xi \sigma_{\text{parton}}(\xi) f(\xi) + \text{small corrections}$$

$$\sim \left| \text{Feynman diagram} \right|^2 + \dots$$

$$\sim \langle \text{hadron} | \Gamma^+ | \text{hadron} \rangle$$

- kinematical corrections
- radiative effects

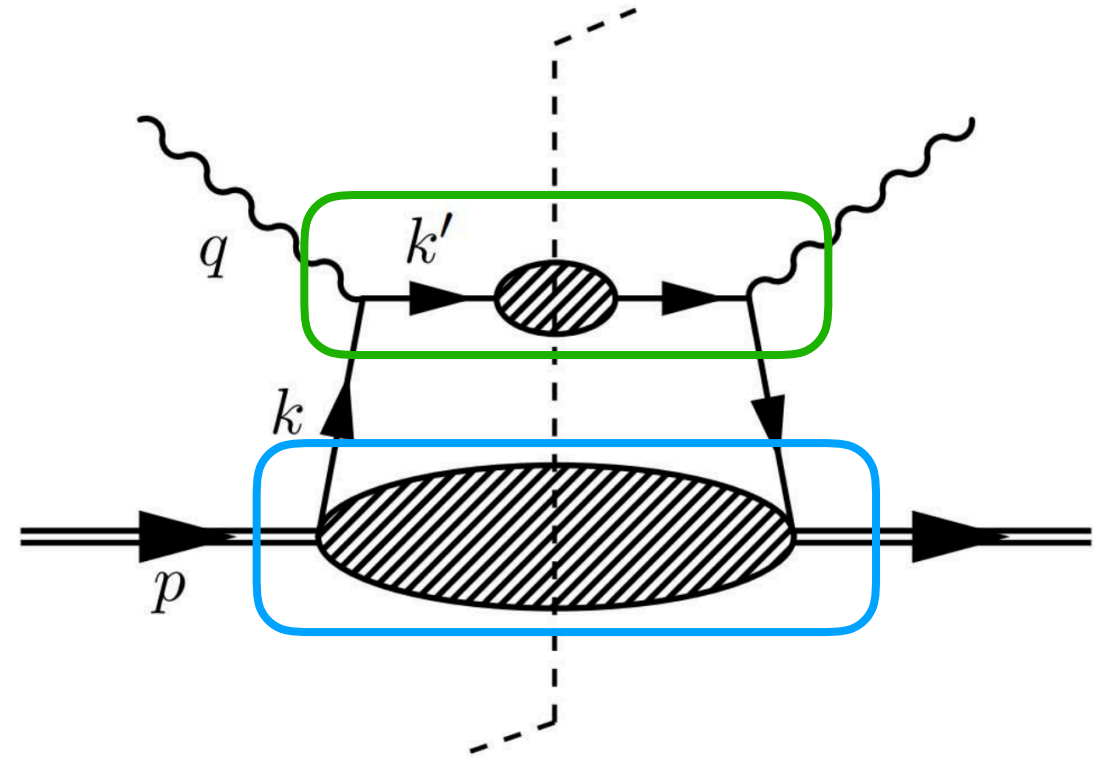
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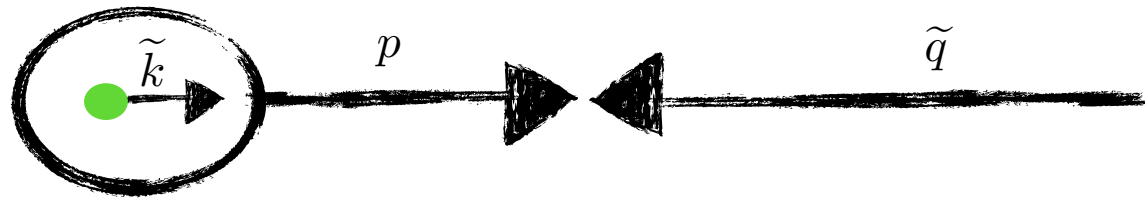
What happens when Lorentz violation is present?

Application I: deep inelastic scattering (DIS)

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In the presence of Lorentz violation, factorization occurs in a modified Breit frame

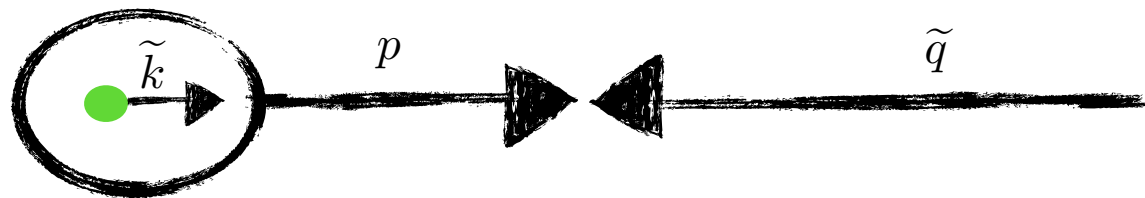
$$\vec{p} + \vec{\tilde{q}} = \vec{0} \qquad \vec{\tilde{q}} \equiv \widetilde{\vec{k} + \vec{q}} - \vec{\tilde{k}}$$



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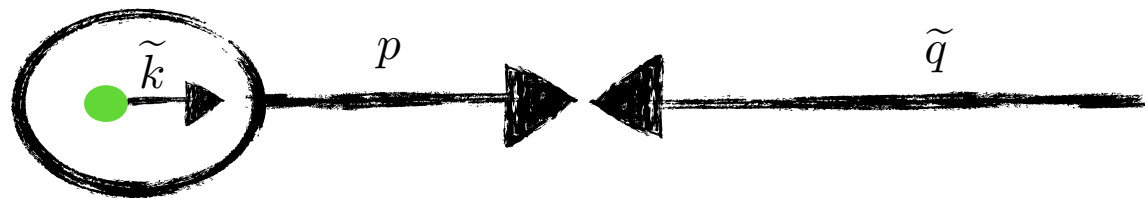
Calculate the imaginary part of internal propagator

$$\text{Im} \frac{1}{\widetilde{\vec{k}}^2 + i\epsilon} = -\pi \left[\delta(\widetilde{\vec{k}}^2) \theta(k^0) + \delta(-\widetilde{\vec{k}}^2) \theta(-k^0) \right]$$

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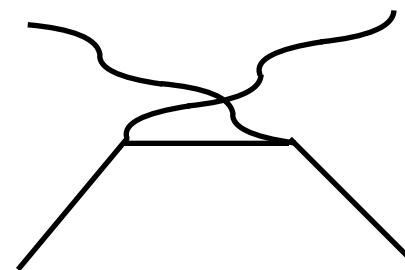
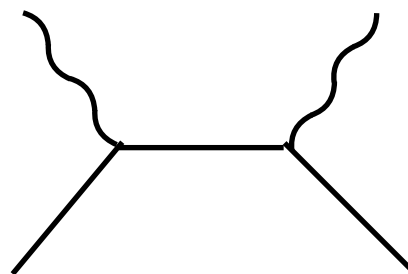
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Quark initiated

Antiquark initiated

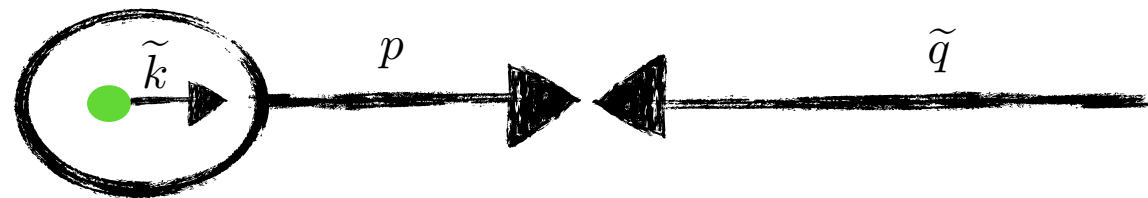
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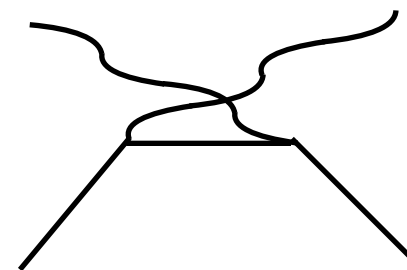
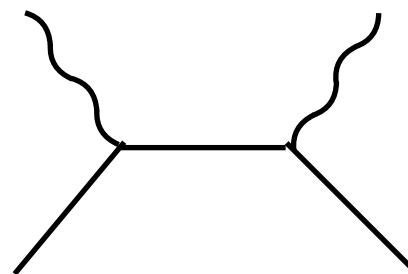
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Will focus on quark contribution

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Example: $\mathcal{L}_c \supset \frac{1}{2} c_f^{\mu\nu} \bar{\psi}_f(x) i \gamma_\mu \overleftrightarrow{\partial}_\nu \psi_f(x)$

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With on-shell parameterization $\tilde{k} = \xi p$

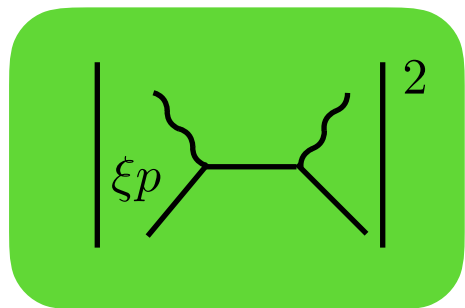
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$$\sim \text{Tr} \left[(\gamma^\mu + c_f^{\alpha\mu} \gamma_\alpha) \frac{1}{(\xi p^\alpha + q^\alpha + c_f^{\alpha\beta} q_\beta) \gamma_\alpha + i\epsilon} (\gamma^\nu + c_f^{\alpha\nu} \gamma_\alpha) \gamma_\beta \xi p^\beta \right]$$

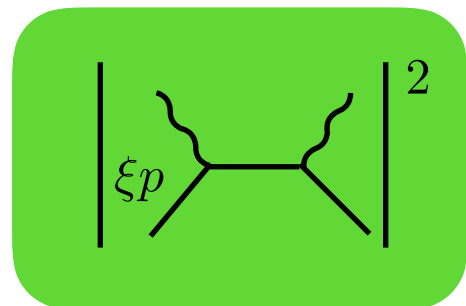
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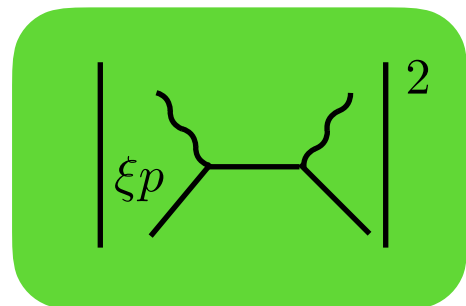
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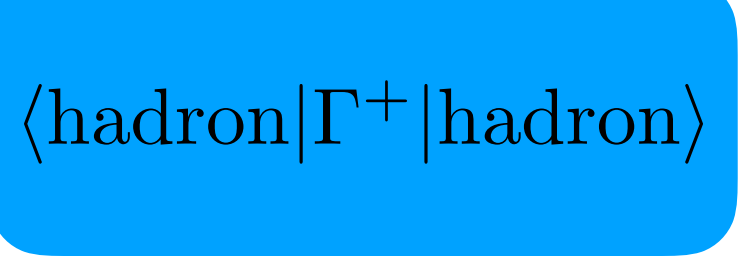
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$$\langle \text{hadron} | \Gamma^+ | \text{hadron} \rangle \sim f_f(\xi, \dots) = \int \frac{d\lambda}{2\pi} e^{-i\xi p \cdot n \lambda} \langle p | \bar{\psi}(\lambda \tilde{n}_f) \frac{\gamma_\mu n^\mu}{2} \psi(0) | p \rangle$$

\searrow
 $n^\mu + c_f^{\mu\alpha} n_\alpha$

“Shifted” conventional scenario

Estimating sensitivities at colliders

Using data from HERA, the LHC, and the future electron-ion collider (EIC) we obtain *estimates* on the sensitivity to the coefficients of interest

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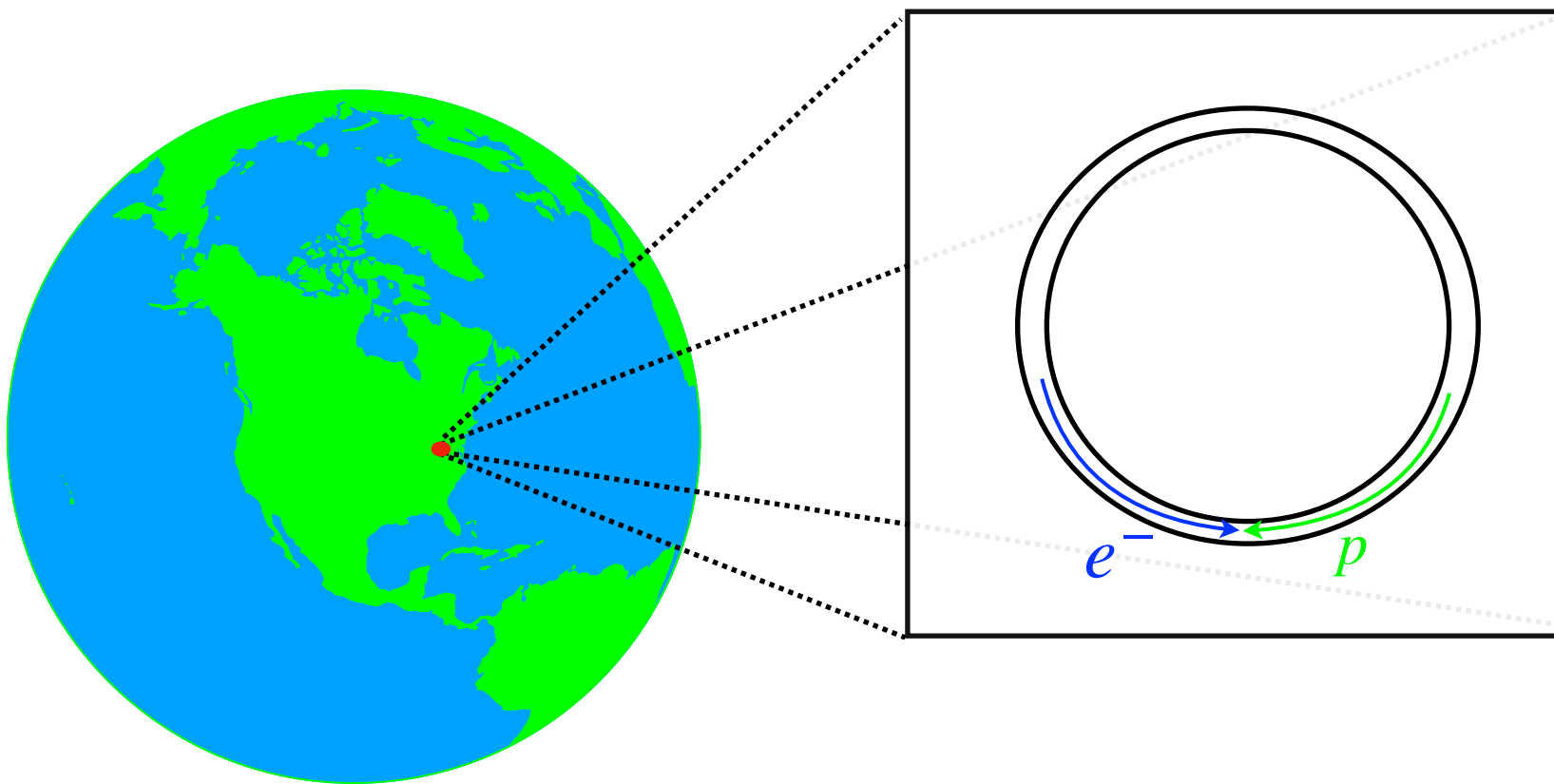
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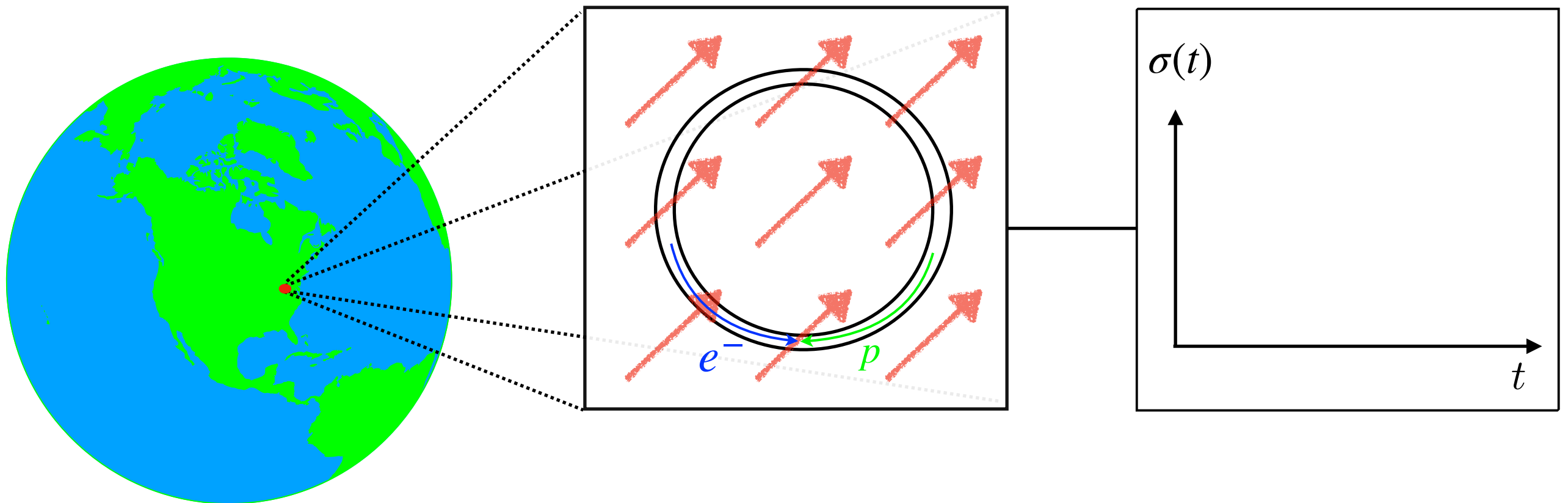
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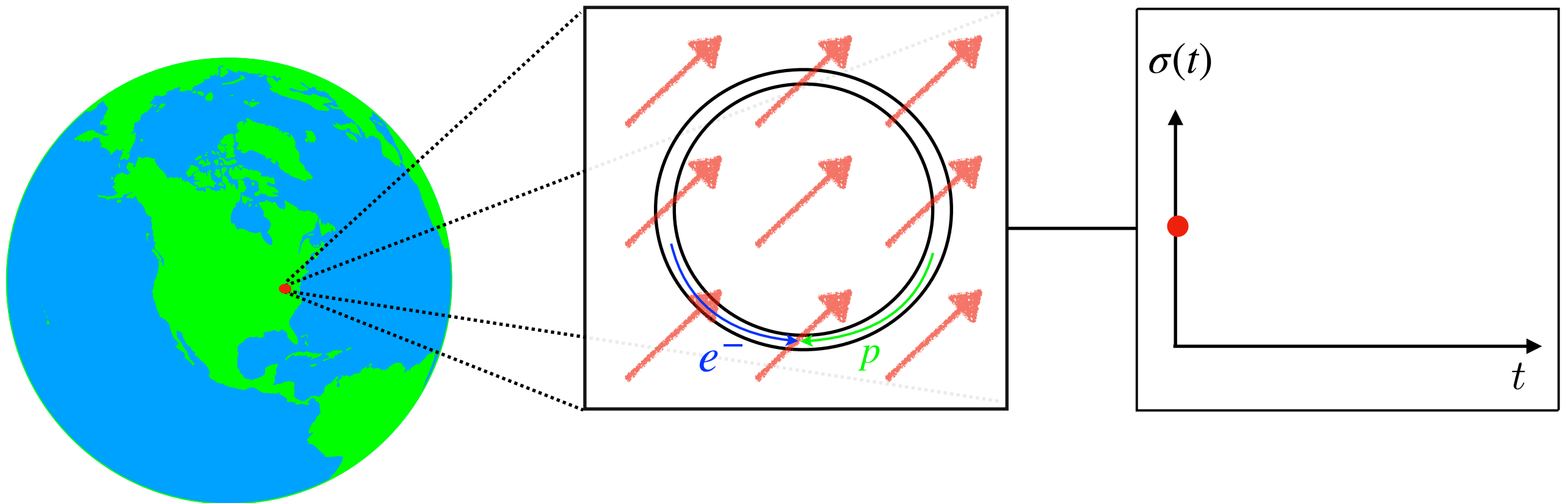
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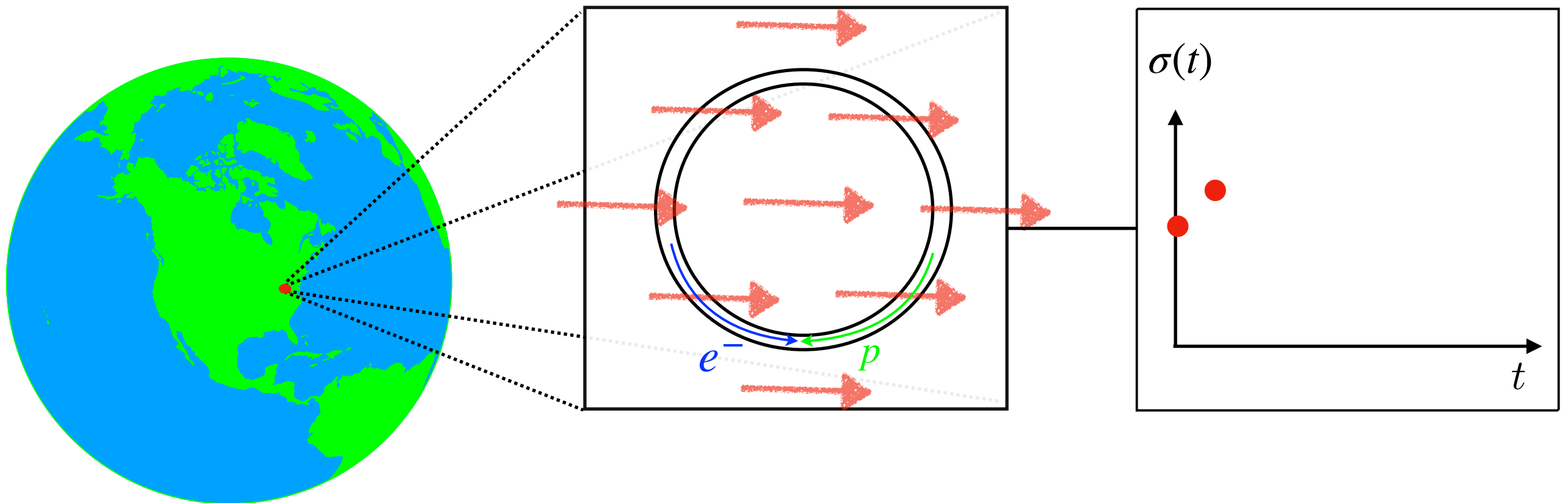
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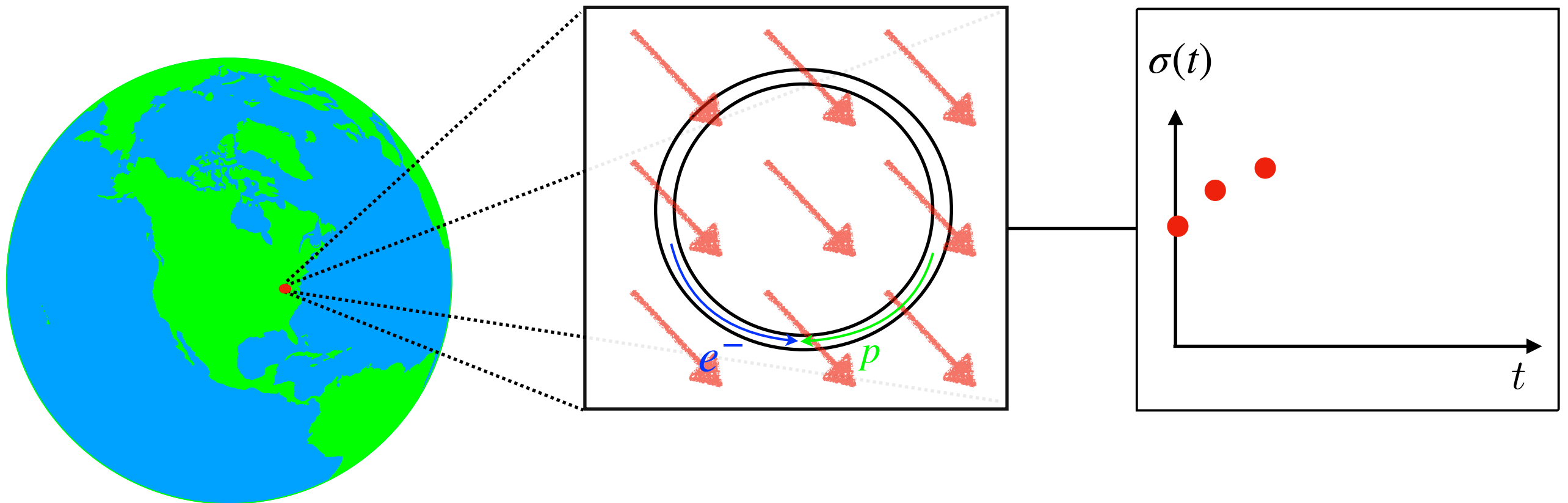
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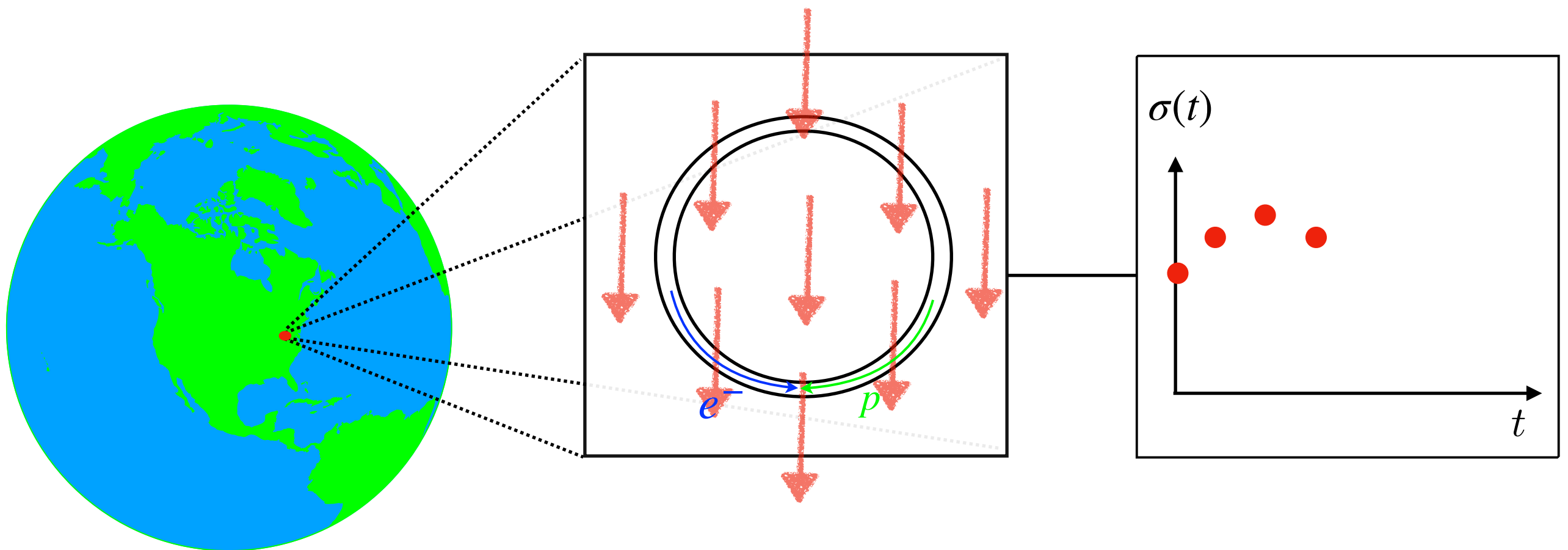
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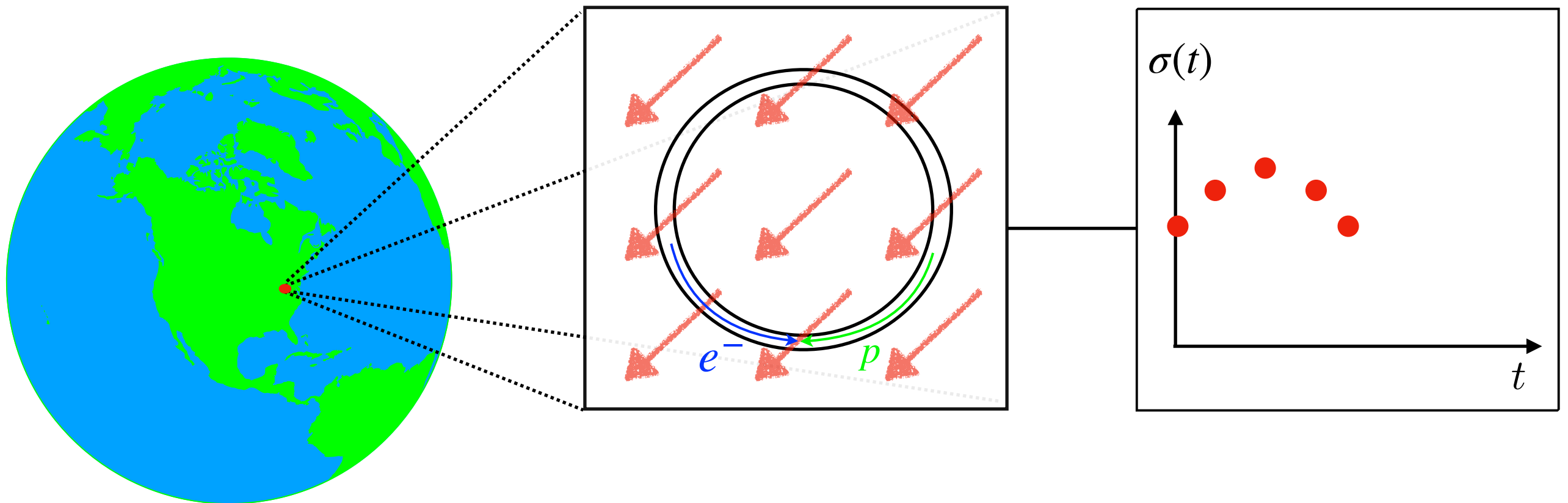
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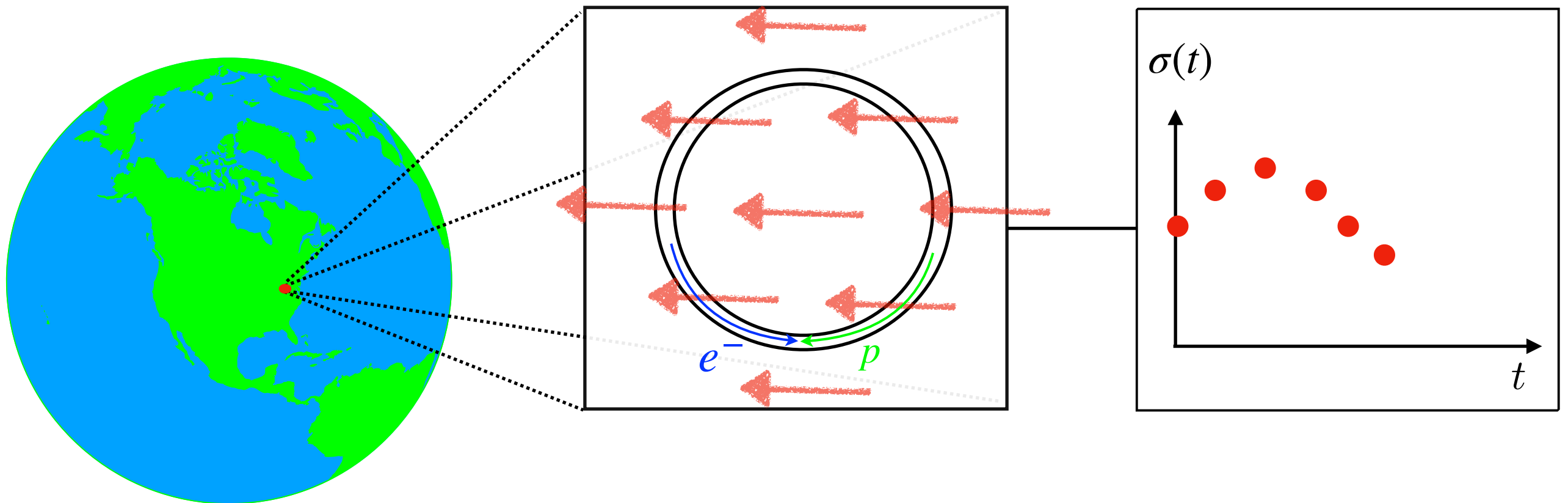
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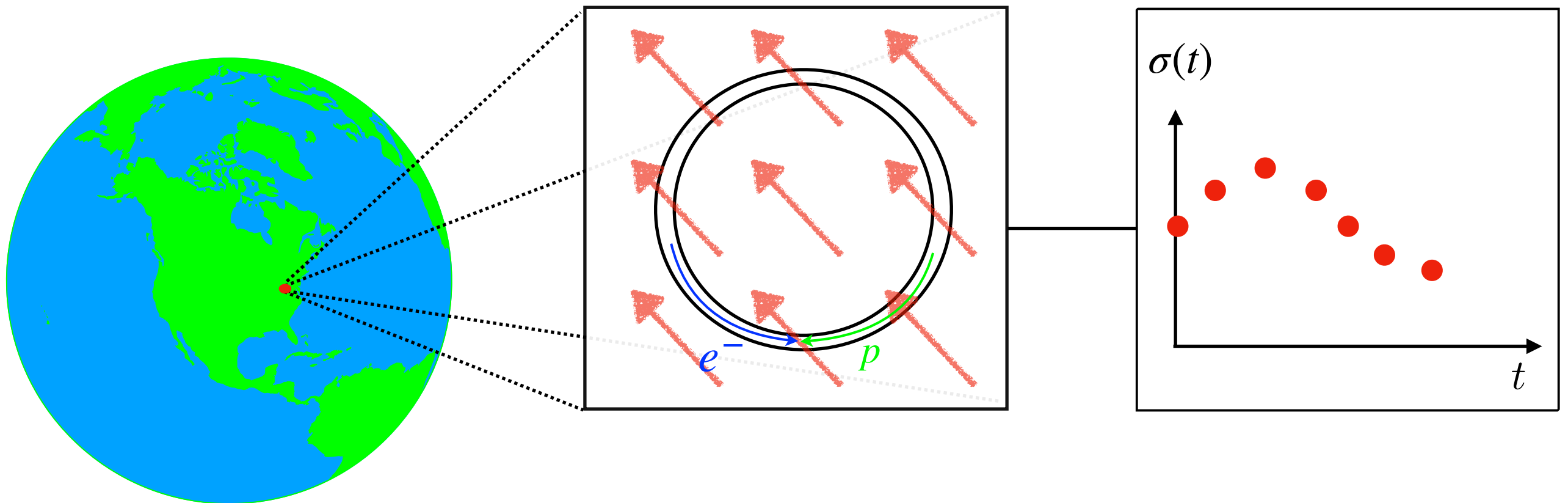
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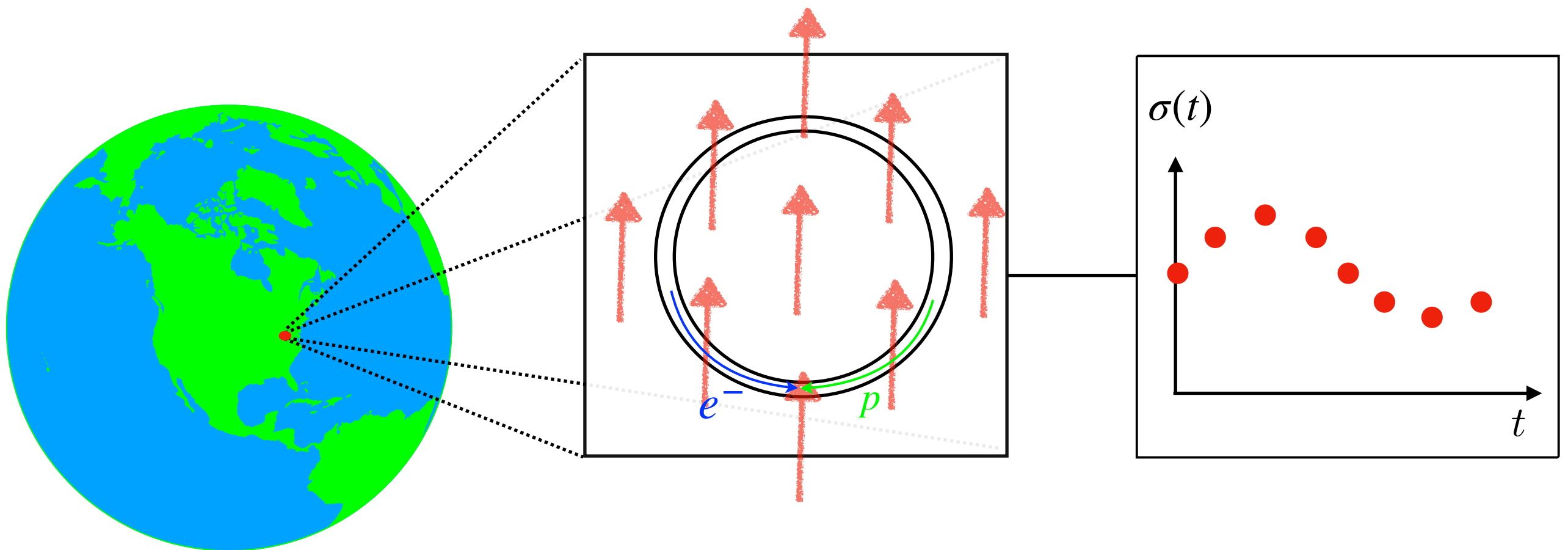
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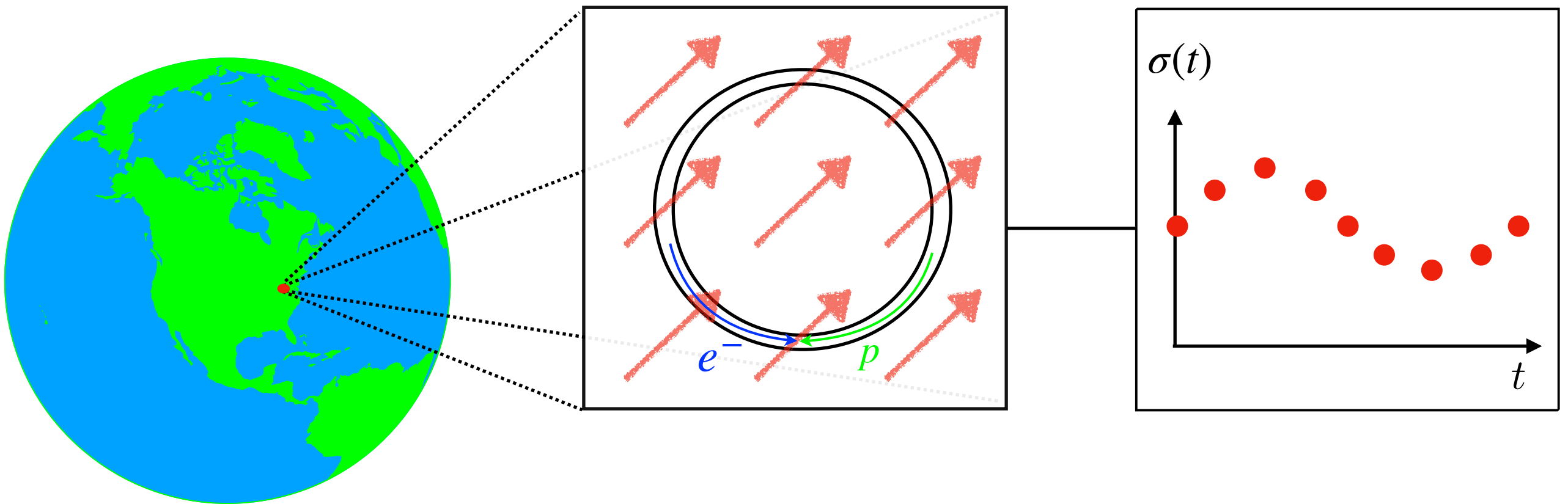
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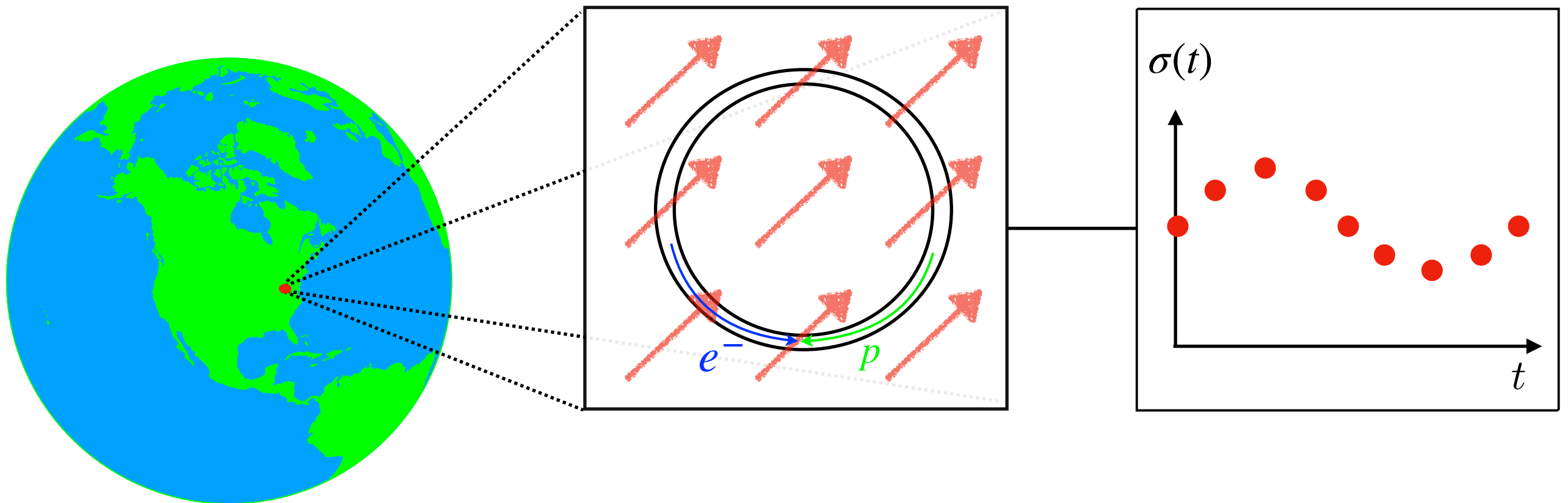
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Coefficients also depend on laboratory colatitude and beam directions!

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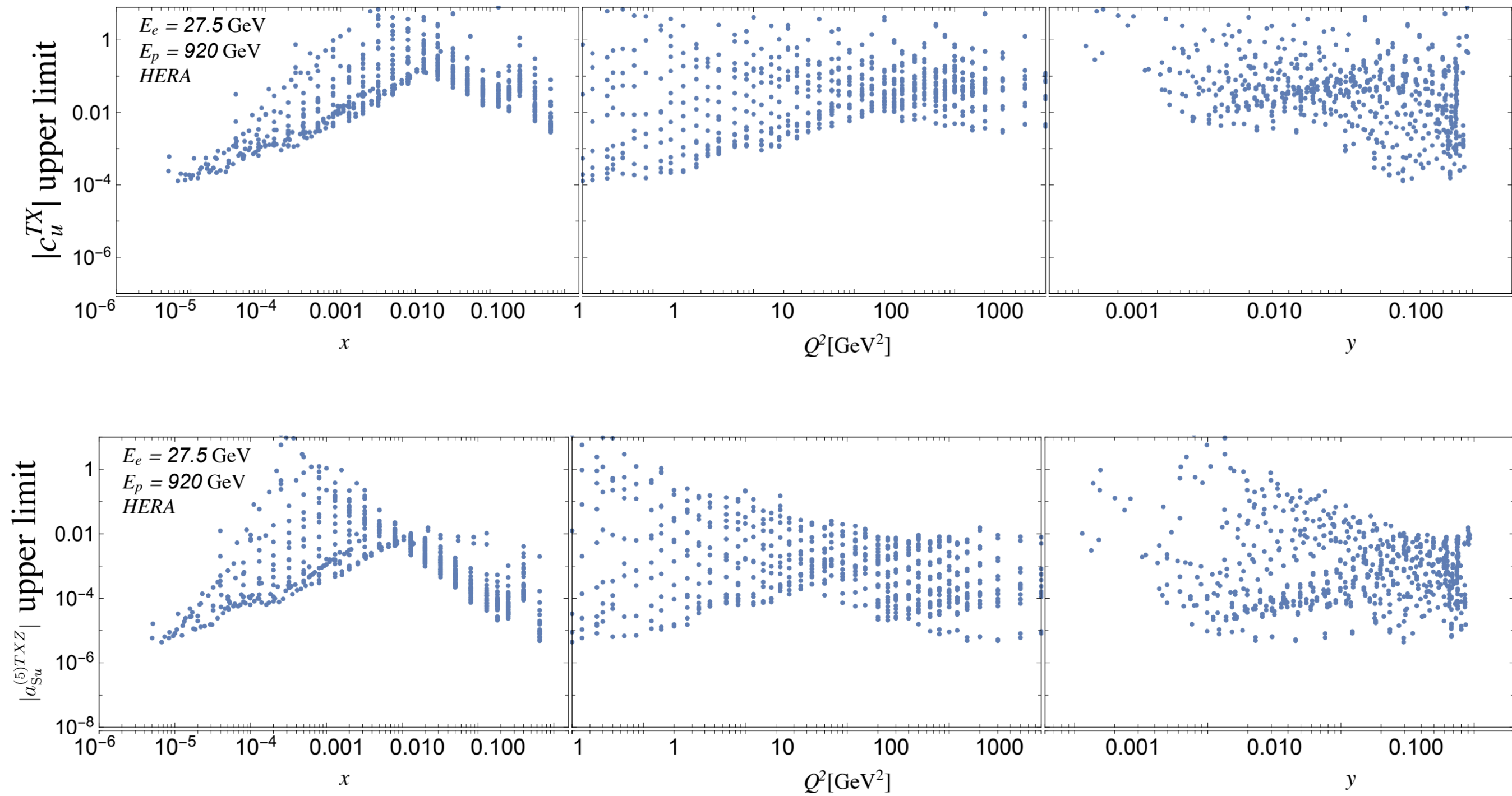
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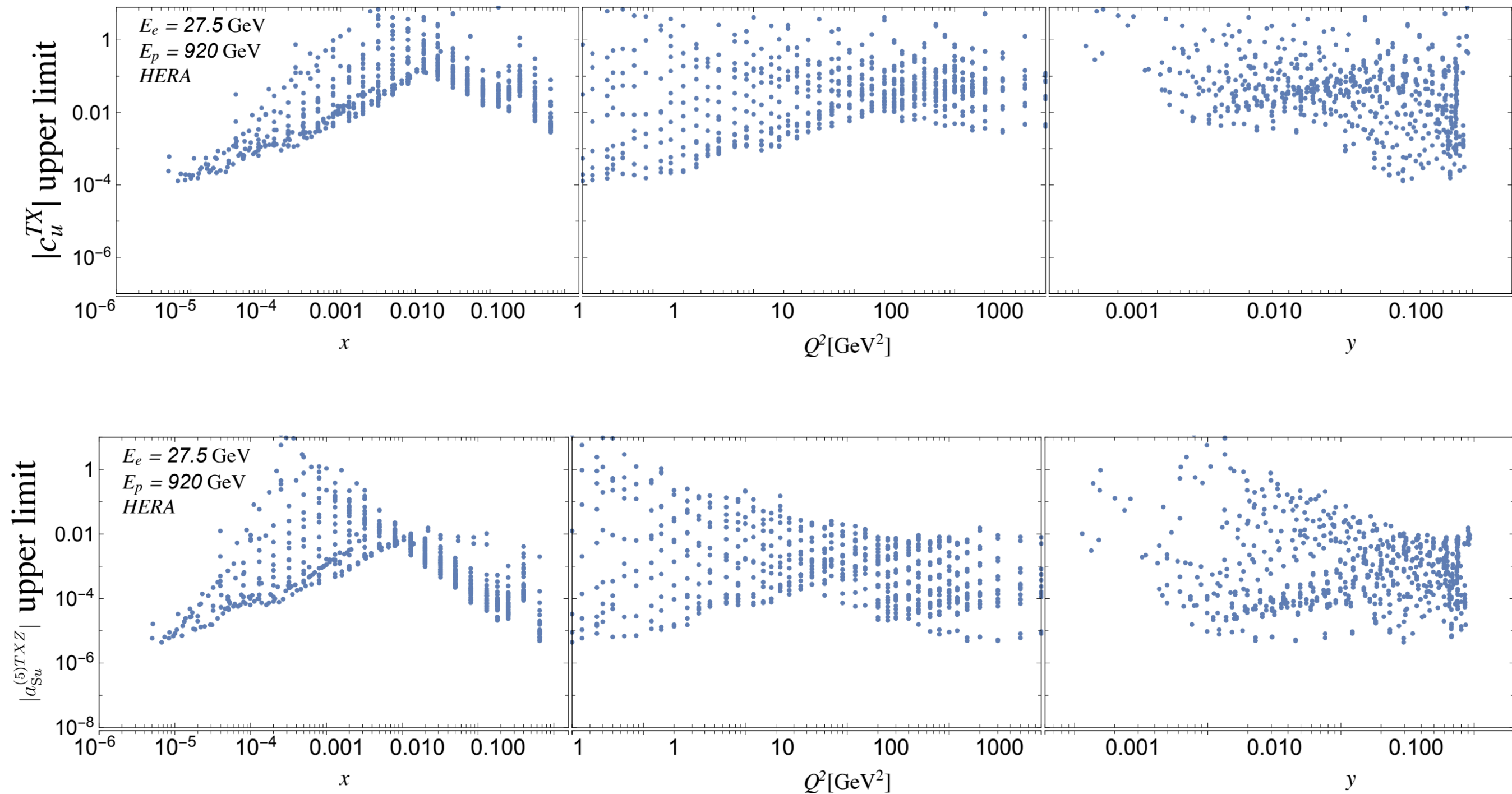
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- Minimize and extract 95% CL constraint

Estimating sensitivities at colliders



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Generally speaking, region of most sensitivity at low x , low-moderate Q , and higher collision energies

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What is needed (?):

- Cross sections as functions of (x, Q) in (4-8) bins of sidereal time
- Understanding of systematics: which uncertainties matter over the course of a few hours in a day? E.g., beam luminosity \sim constant?
- Since SME cross section is different in each bin, can construct observables that partially shield against systematics

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Thank you!