Far-from-equilibrium dynamics of axion-like particles with broken shift symmetry

Aleksandr Chatrchyan

In collaboration with Jürgen Berges and Joerg Jaeckel







ALPs with a broken shift symmetry

• ALPs can arise as pseudo Nambu-Goldstone bosons (e.g. QCD axion) or string theory axions

Peccei, Quinn, PRL 38, 1440.

Weinberg, PRL 40, 223. Svrcek, Witten,

Arvanitaki et al.,

Cicoli et al., JHEP 06, 051 (2006). PRD 81, 123530 (2010). JHEP 10, 146 (2012).

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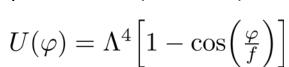
Svrcek, Witten, JHEP 06, 051 (2006). PRD 81, 123530 (2010). JHEP 10, 146 (2012).

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ALPs enjoy an approximate shift symmetry, which can be broken by

nonperturbative (instanton) effects

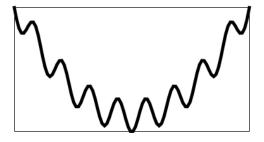


Callan Curtis et al., PRD 17, 2717 (1978).



$$U(\varphi) = \mu^{4-m} \varphi^m$$

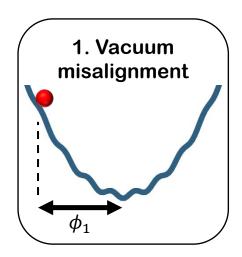
• Combination of both effects: wiggly potentials



Misalignment production of ALP dark matter

- Considered potential: $U(\varphi) = \frac{1}{2}m^2\varphi^2 + \Lambda^4\Big(1-\cos\frac{\varphi}{f}\Big)$
- Homogenous initial conditions, $\langle \hat{\varphi}(\mathbf{x}) \rangle = \phi_1$.

Preskill, et al.,
PLB 120, 127–132 (1983).
Arias et al.,
JCAP 1206, 013 (2012).

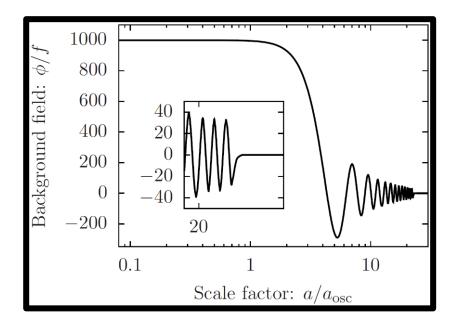


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- Homogenous initial conditions, $\langle \hat{\varphi}(\mathbf{x}) \rangle = \phi_1$.
- Coherent oscillations around the minimum.

1. Vacuum misalignment $H \sim \frac{m_a}{3}$ 2. Coherent oscillations Growth of fluctuations

Jaeckel et al., JCAP 1701, 036 (2017). Preskill, et al.,
PLB 120, 127–132 (1983).
Arias et al.,
JCAP 1206, 013 (2012).



Berges, AC, Jaeckel, JCAP 1908 (2019) 020

Fonseca et al., 1911.08473

Fonseca et al., JHEP 04, 010 (2020).

Gravitational wave production

• The frequency is determined by the mass, $\nu_{\star} \propto \sqrt{m_a}$

see also

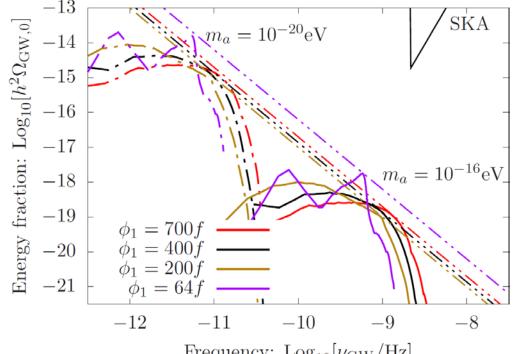
Machado et. al, 1912.01007

Machado et al., JHEP 01, 053 (2019).

Kitajima et al., JCAP 10 (2018) 008.

The strength of the signal

$$h_{\star} \propto \rho_{\varphi}(a_{\mathrm{emit}}) = \rho_{\varphi}(a_{\mathrm{today}}) \left(\frac{a_{\mathrm{today}}}{a_{\mathrm{emit}}}\right)^{3} \frac{1}{\mathcal{Z}}$$
 Fixed
$$\mathcal{Z} = \exp\left(-3 \int_{a_{\mathrm{emit}}}^{a_{\mathrm{today}}} w_{\varphi}(\widetilde{a}) d\ln \widetilde{a}\right)$$

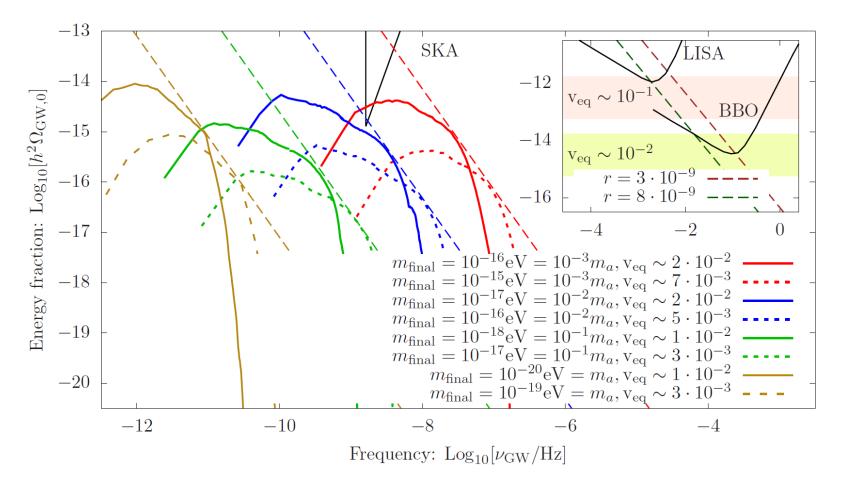


Frequency: $Log_{10}[\nu_{GW}/Hz]$

GWs from ALP dark matter

• Small final mass: extended relativistic phase after fragmentation, $\,\mathcal{Z}\ll 1.\,$

AC, Jaeckel 2004.07844



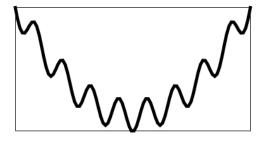
ALPs driving inflation: Preheating

• $m \sim 10^{-5} M_{\rm Pl}$, GW production during preheating at frequencies $\nu \sim (10^7 - 10^9) {\rm Hz}$.

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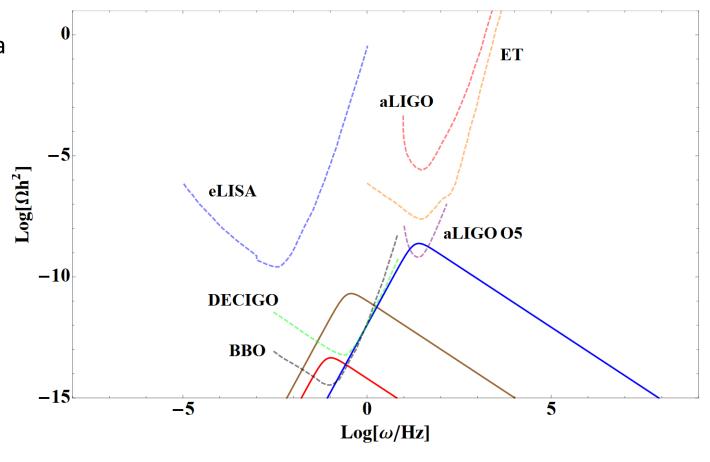
Phase transitions between minima



• GW production from PTs

Figure taken from

Hebecker et al. JCAP 11 (2016) 003



Bubble nucleation out of equilibrium

- Instanton (static) picture by Coleman, Callan, de Luccia and Linde.
 - Valid when the field is initially in vacuum or thermal state.
 - The nucleation rate is determined by the Euclidean action of the O(4) or O(3) bounce solution:

$$\Gamma_4 = A_4 \exp(-S_4), \qquad \Gamma_3 = A_3 T \exp(-S_3/T).$$

- Not applicable in our case
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- Stochastic approach to tunneling, proposed by Linde
 - Based on classical-statistical approximation
 - (Classical) fluctuations randomly add up in position space to form a bubble

A. Linde Nucl.Phys.B 372 (1992)

Braden et al., PRL 123 (2019)

Hertzberg, Yamada PRD 100 (2019)

Stochastic (real-time) approach

• Analytical estimation:

$$\Gamma \sim \exp(-rac{arphi_b^2}{2\langlearphi^2
angle_{\mathrm{p}<1/\mathrm{r}}}), \qquad \mathrm{where} \ \langlearphi^2
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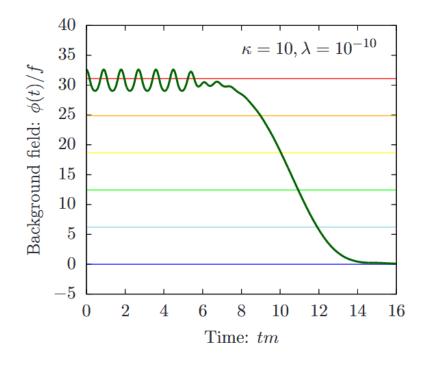
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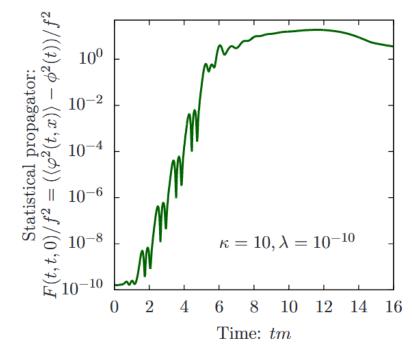
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• Numerical simulations:

More details: in preparation





Thanks for your attention!