



Axion Gegenschein

Dark Countersources of Bright Radio Objects

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» Photons from axion decay

Interaction term involving axions and photons

$$\mathcal{L}_{a\gamma} = -\frac{1}{4}g_{a\gamma}aF_{\mu\nu}\tilde{F}_{\mu\nu} \quad (1)$$

Axion decay $a \rightarrow \gamma\gamma$ leads to two photons each of frequency ν s.t.

$$h\nu_d = m_a c^2 / 2$$

with an estimated decay lifetime

$$\begin{aligned} \tau_{a\gamma} &= \frac{64\pi\hbar}{m_a^3 c^6 g_{a\gamma}^2} = 1.7 \times 10^{43} \left(\frac{5.7 \times 10^{-6} \text{eV}}{m_a c^2} \right)^3 \\ &\times \left(\frac{C_{a\gamma} 1.16 \times 10^{-15} \text{GeV}^{-1}}{g_{a\gamma}} \right)^2 \text{yr} \end{aligned} \quad (2)$$

» Stimulated decay

Evolution of axion number density

$$\dot{n}_a \simeq -n_a \Gamma_a (1 + 2N_\gamma) \quad (3)$$

with decay rate $\Gamma_a = \tau_a^{-1}$

Photon occupation number in the ambient radiation related to its energy density

$$N_\gamma = \frac{\pi^2 \rho_\gamma}{E_\gamma^3} \quad (4)$$

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Stimulated decay of axions: decay rate enhanced by quantum occupation number of photons of frequency $\sim m_a c^2 / \hbar$

↓

One of the two final state photons in a momentum state opposite to that of incoming radiation

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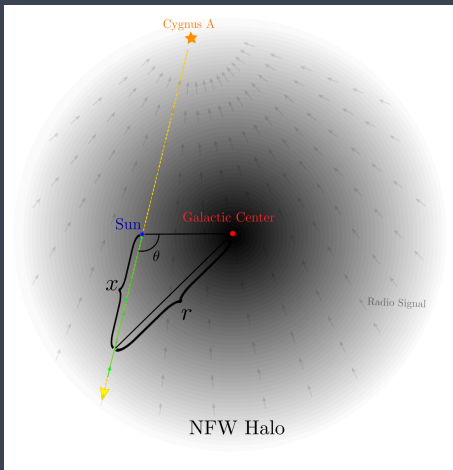
Creation of a “countersource” made of axions



Analogous to zodiacal light, the decay photons traveling backwards \leadsto axion gegenschein

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Backscattering of astrophysical radio pulses in the MW Halo along line of sight \rightarrow dark countersources for every radio-bright object

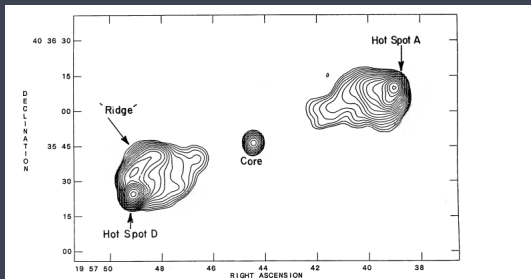


Axion gegenschein in the Milky Way's dark matter halo

» Source selection

Cygnus A ($\ell = 76.20^\circ$, $b = 5.75^\circ$)

Brightest extragalactic radio source: Type II Seyfert galaxy hosting AGN



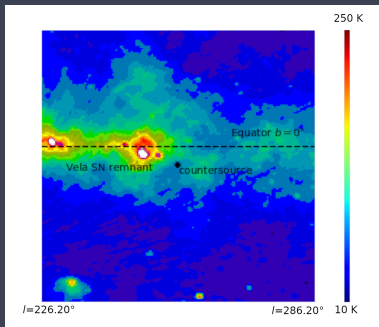
Spectral flux:

$$\log S_{A\nu_d}(\nu_d) = a + b \log \nu_d + c \log^2 \nu_d \quad (5)$$

$\theta = 103.74^\circ$, angular separation between Cygnus A & Sun wrt GC

» Countersource of Cygnus A

Location: 180° apart ($\ell = 256.20^\circ, b = -5.75^\circ$)



Countersource in radio sky ($60^\circ \times 60^\circ$ cutout of 408-MHz sky)

Characteristics:

- * Size comparable to source size with spatial features preserved
- * Smoothing owing to dark matter velocity dispersion \perp l.o.s.
- * Distant sources advantageous owing to longer round-trip time

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Intensity due to stimulated decay (I_g)

\approx wave intensity from spontaneous decay (I_{sp}) \times
photon occupation number in incoming wave (N_γ)

Spectral intensity being $I_\nu(\nu) = (\nu/c)^2 h\nu N_\gamma$, spontaneous intensity owing to axion number density $\rho_a(x)/m_a$ over time dt is

$$dI_{sp} = \rho_a / (m_a \tau_a) h\nu_d (cdt/4\pi) \quad (6)$$

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Thus, the gegenschein intensity is

$$I_g = \int_0^\infty dx \frac{\rho_a(x)}{m_a \tau_a} \frac{c^2}{4\pi\nu_d^2} I_{s\nu}(\nu_d, x) \quad (7)$$

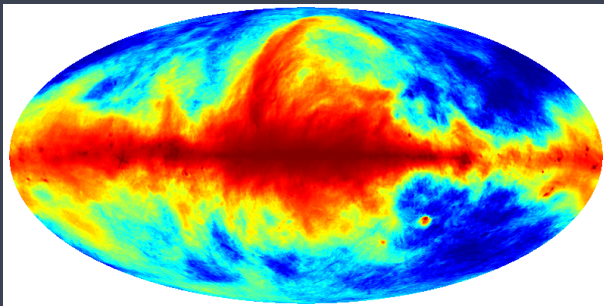
Replacing decay lifetime τ_a

$$I_g = \frac{\hbar c^4}{16} g_{a\gamma}^2 \int_0^\infty dx I_{s\nu}(\nu_d, x) \rho_a(x) \quad (8)$$

Note: the incoming radio intensity (from the astrophysical source) $I_{s\nu}$ is conserved along the line of sight

» Radio sensitivity

Galactic and extragalactic background in synchrotron radio emission



408-MHz Haslam all-sky map

» Radio sensitivity

Considering the contribution of atmospheric radio noise and CMB (blackbody)

$$T_s = T_a + T_{\text{CMB}} + T_{\text{bg}} + T_r \quad (9)$$

Synchrotron background estimated as power law $T_{\text{bg}}(\nu) \propto \nu^\beta$ with $\beta = -2.55$.

Thus

$$T_{\text{bg}} = 60 \left(\frac{300\text{MHz}}{\nu} \right)^{2.55} \text{ K} \quad (10)$$

$T_{\text{bg}}=27 \text{ K}$ at the sky coordinate of the countersource at 408 MHz.

» Radio sensitivity



SKA-low (50-350 MHz)

Dipole array with number of elements $\sim 131,000$

Collection area (A_{coll}) $\sim 419,000 \text{ m}^2$

$$T_r = 40K$$

SKA-mid: (350 MHz-15.3 GHz)

$$N_{tele} \sim 5600$$

Diameter $D=15 \text{ m}$ for each dish

$$T_r = 20K$$

Efficiency of SKA: $\eta = 0.8$

» Radio sensitivity

Radio power of background

$$P_{\text{noise}} = 2k_B T_s \sqrt{\frac{\Delta\nu}{t_{\text{obs}}}} \quad (11)$$

Radio power of gegenschein signal

$$P_{\text{signal}} = \eta AS \quad (12)$$

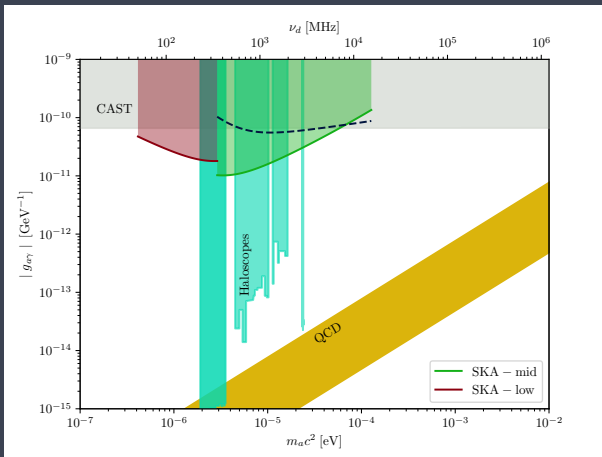
For an $n - \sigma$ detection, $P_{\text{signal}} / P_{\text{noise}} = n$

The spectral flux owing to gegenschein emission of Cygnus A is

$$S_{Ag} = \frac{\hbar c^4}{16} g_{a\gamma}^2 S_{A\nu}(\nu_d) \int_0^\infty dx \rho_a[r(x)] \quad (13)$$

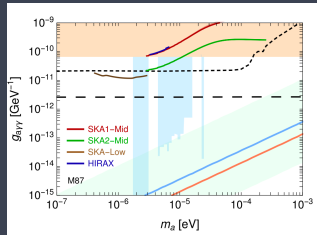
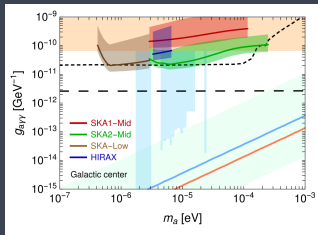
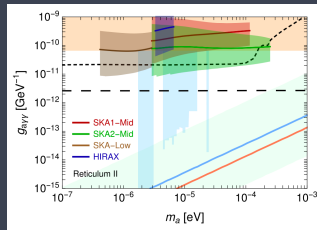
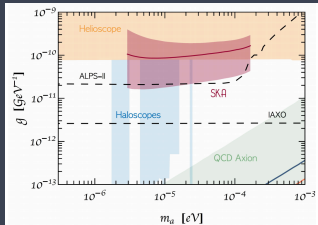
» Radio sensitivity

For a conservative $S/N = 1$ and standard NFW halo model,



Solid lines: Axionogenesis in MW, *this work*
Dashed line: Stimulated decay of axions in dSph

» Other proposed probes



Caputo et al. [1805.08780](#), [1811.08436](#)

» Other proposed probes

For a top-hat frequency profile of bandwidth, the optimal frequency width to achieve maximum S/N

$$\Delta\nu = 2.17\nu_d \frac{\sigma_d}{c} \quad (14)$$

Spectral flux S_{Ag} is reduced by

$$f_{\Delta} = 0.721 \quad (15)$$

A decrease in number of photons photon collected within each frequency bin

Radio sensitivity must be adjusted by a correction factor \Rightarrow goes down by $\sim \frac{1}{0.49}$ in radio power, thus goes up by ~ 0.7 in $g_{a\gamma}$!

» Key takeaways

- * Astrophysical radio pulses from galactic and extragalactic sources can induce stimulated emission in the Milky Way's halo, emission line feature with a spread characterized by DM velocity dispersion
- * Detectable at future-generation radio telescopes even with a conservative $S/N \sim 1$
- * Axion gegenschein provides **100-fold increase in radio sensitivity compared to proposed frameworks**
- * Probing dark matter overdensities at large distances
- * Axion gegenschein opens up a new (and relatively clean) indirect search method for dark matter: stimulated emission from radio-bright sources

Axion Gegenschein, *arXiv: 2008.02729*

OG, J. Salvado, J. Miralda-Escudé



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» The End

Thank you for your attention!

Questions?

» Backup slides

Single dish mode

Angular resolution

$$\theta_{FWHM} \simeq 1.22 \frac{\lambda}{D} \simeq 0.7^\circ \left(\frac{1\text{GHz}}{\nu} \right) \left(\frac{15\text{m}}{D} \right) \quad (16)$$

Noise temperature of the instrument

$$T_{\text{ant}} = \frac{A_{\text{eff}} \langle S \rangle}{2k_b} \quad (17)$$

For each telescope

$$\left(\frac{S}{N} \right)_{sd, \text{single}} = \frac{T_{\text{ant}}^{\text{pb}}}{T_{\text{min}}} \quad (18)$$

For an array of single-dish telescopes

$$\left(\frac{S}{N} \right)_{sd, \text{array}} = \sqrt{N_{\text{tele}} n_{\text{pol}}} \left(\frac{S}{N} \right)_{\text{single}} = \sqrt{N_{\text{tele}} n_{\text{pole}}} \frac{T_{\text{ant}}^{\text{pb}}}{T_{\text{min}}} \quad (19)$$

» Backup slides

Interferometric mode

Angular resolution of primary beam

$$\theta_{pb} = 12.5' \left(\frac{1\text{GHz}}{\nu} \right) \left(\frac{100\text{m}}{D} \right) \quad (20)$$

Minimum detectable flux density

$$S_{\min} = SNR \frac{\text{SEFD}}{\sqrt{n_{\text{pol}} \Delta B t_{\text{obs}}}} \quad (21)$$

with

$$\text{SEFD} = \frac{T_{\text{sys}}(\nu)}{G} \quad (22)$$

» Backup slides

Observation in the interferometric mode

For each synthesized beam in the interferometric mode, angular resolution

$$\theta_{\text{synth}} \approx 50'' \left(\frac{1\text{GHz}}{f} \right) \left(\frac{1\text{km}}{B_{\text{max}}} \right) \quad (23)$$

SNR is expressed as $\sqrt{\delta\chi^2}$

$$\delta\chi^2 = n_{\text{pol}} t_{\text{obs}} G_{\text{array}}^2 \sum_{i=1}^{N_{\text{synth}}} \frac{F_i^2}{B_i T_i^2} \quad (24)$$

Scaling relation for antenna gain $G_{\text{array}} \sim N_{\text{tele}} (N_{\text{tele}} - 1) G$

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