



Analytic Structure of Banana Amplitudes

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Joint work with Kilian Bönisch, Fabian Fischbach, Albrecht Klemm & Reza Safari

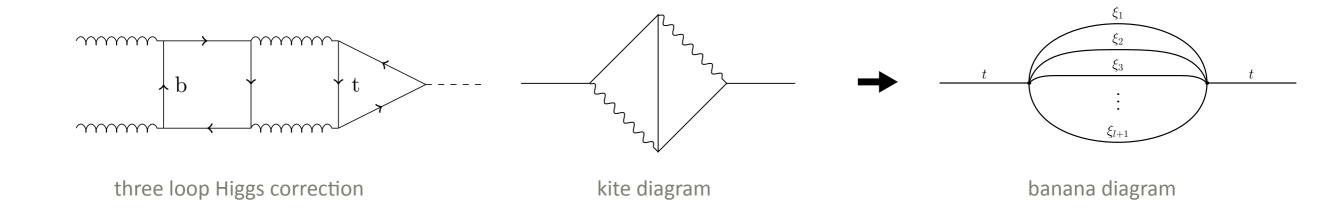
"The l-loop Banana Amplitude from GKZ Systems and relative Calabi-Yau Periods" [1912.06201] "Analytic Structure of all Loop Banana Amplitudes" [2008.10574]

Motivation

Physical interest:

Scattering processes are calculated perturbatively through Feynman diagrams

Precision measurements require high loop calculations



Mathematical interest:

Appearance of special functions (elliptic Polylogarithms, iterated integrals, Bessel functions, ...)

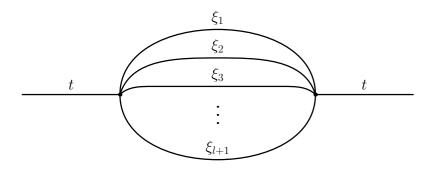
Function space of Feynman integrals?

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Algebraic and geometric structure behind Feynman integrals

Projective varieties, motives, period integrals, $\widehat{\Gamma}$ -class conjecture, ...

Representations of Feynman Integrals



E: # propagators \rightarrow l+1

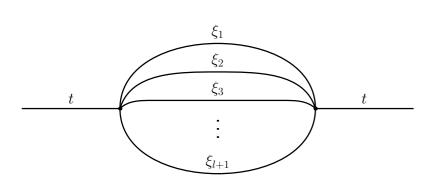
v: # vertices \longrightarrow 2

l: # loops → *l*

 u_k : # powers \rightarrow 1 or higher propagators $\nu = \sum \nu_k$

Representations of Feynman Integrals

Feynman Representation

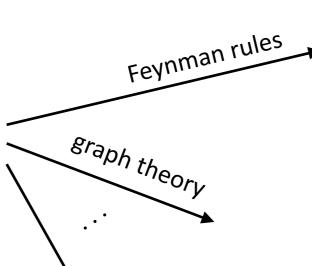


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Feynman rules
$$\mathcal{I} \sim \int_{(\mathbb{R}^{1,D-1})^l} \frac{\prod_{k=1}^l \mathrm{d}^D l_k}{\prod_{k=1}^l (q_k^2 - \xi_k^2 + i\epsilon)^{\nu_k}}$$

Symanzik Representation

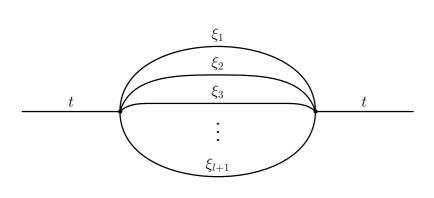
$$\mathcal{I} = \int_{x_k \ge 0} \prod_{k=1}^{E} x_k^{\nu_k - 1} \frac{\mathcal{U}^{\nu - (l+1)D/2}}{\mathcal{F}^{\nu - lD/2}} \mu_l$$

$$\mu_l = \sum_{k=1}^{l+1} (-1)^{k+1} x_k dx_1 \wedge \ldots \wedge \widehat{dx_j} \wedge \ldots \wedge dx_{l+1}$$



Representations of Feynman Integrals

Feynman Representation

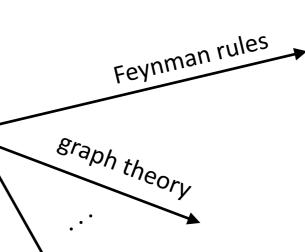


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Graph polynomials

ullet First Symanzik polynomial ${\cal U}$

No kinematic dependence

ullet Second Symanzik polynomial ${\mathcal F}$

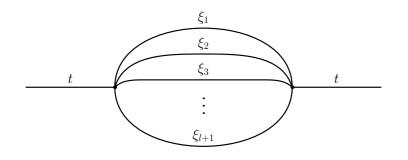
homogeneous polynomial of degree l+1 in \boldsymbol{x}_i dependence on masses and momenta

homogeneous polynomial of degree l in x_i



Geometric Realization

In D=2 the banana diagram is given by:

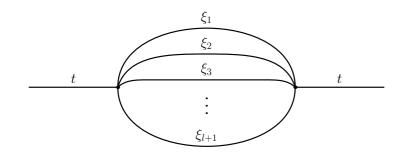


$$\mathcal{I}_l(t,\xi_i) = \int_{\sigma_l} \frac{\mathcal{U}^0}{\mathcal{F}^1} \mu_l = \int_{\sigma_l} \frac{\mu_l}{P_{\Delta_l}(t,\xi_i;x_i) \prod_{k=1}^{l+1} x_k}$$

$$P_{\Delta_l}(t, \xi_i; x_i) \coloneqq \left(t - \left(\sum_{k=1}^{l+1} \xi_k^2 x_k\right) \left(\sum_{k=1}^{l+1} \frac{1}{x_k}\right)\right)$$
$$\sigma_l \coloneqq \left\{ [x_1 : \dots : x_{l+1}] \in \mathbb{RP}^l | x_i \ge 0, \forall i \right\}$$

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1) Hypersurface in toric ambient space

$$M_{l-1} = \{ P_{\Delta_l}(x) = 0 | x \in \mathbb{P}_{\hat{\Delta}_l} \}$$

- Clear connection from differential to geometry
- Generic hypersurface constraint hast far too many parameters
 - Subslice problem
- Differential equations for periods "for free"

2) Complete intersection model

$$W_{l-1} = \begin{pmatrix} \mathbb{P}_1^1 & 1 & 1 \\ \vdots & \vdots & \vdots \\ \mathbb{P}_{l+1}^1 & 1 & 1 \end{pmatrix} \subset \begin{pmatrix} \mathbb{P}_1^1 & 1 & 1 \\ \vdots & \vdots & \vdots \\ \mathbb{P}_{l+1}^1 & 1 & 1 \end{pmatrix} = F_l$$

- Hidden connection through maximal cut period
- Number of parameters fits to physical ones
- Differential equations for periods "for free"



Equal Mass Case: Operators

We set all masses to unity, i.e. $\xi_i=1$ for $i=1,\ldots,l+1$

Maximal cut integral

$$\mathcal{T}_{T^{l}}(s) = \int_{T^{l}} \frac{\mu_{l}}{\left(1/s - \left(\sum_{k=1}^{l+1} \xi_{k}^{2} x_{k}\right) \left(\sum_{k=1}^{l+1} \frac{1}{x_{k}}\right)\right)}$$

$$\sim \sum_{n=0}^{\infty} s^{n+1} \sum_{|k|=n} \binom{n}{k_{1}, \dots, k_{l+1}}^{2} =: \varpi_{0}(s)$$

Find differential operator annihilating maximal cut integral



Set of solutions (almost) describe Feynman integral as linear combination

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Observation: $\varpi_0(s)$ is double Borel sum of the (l+1)th symmetric power of

$$\sum_{k=0}^{\infty} \frac{1}{(k!)^2} z^k = I_0(2\sqrt{z})$$

 \longrightarrow

Get easily PF equations from this

$$\mathcal{L}_l \mathcal{T}_{T^l} = 0$$

E.g. for
$$l=4$$
 we find:

$$\mathcal{L}_4 = 1 - 5s + (-4 + 28s)\theta + (6 - 63s + 26s^2 - 225s^3)\theta^2 + (-4 + 70s - 450s^3)\theta^3 - (-1 + s)(-1 + 9s)(-1 + 25s)\theta^4$$

$$\theta = s\partial_s$$



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For full Feynman integral we get inhomogeneity $\mathcal{L}_{l}\mathcal{F}_{\sigma_{l}}=S_{l}=-(l+1)!$ s

$$\mathcal{L}_l \mathcal{F}_{\sigma_l} = S_l = -(l+1)!$$



Equal Mass Case: Local Solutions

Around the MUM point s=0 we have a local Frobenius basis of the form:

$$\varpi_k = \sum_{j=0}^k {k \choose j} \log(s)^j \, \Sigma_{k-j} \qquad \text{for } k = 1, \dots, l-1$$

$$\varpi_l = (-1)^{l+1} (l+1) \sum_{j=0}^l {l \choose j} \log(s)^j \, \Sigma_{l-j}$$

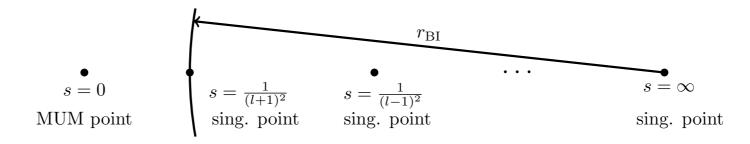
Again for l=4 we have:

$$\varpi_0 = s + 5s^2 + 45s^3 + 545s^4 + 7885s^5 + \cdots
\Sigma_1 = 8s^2 + 100s^3 + \frac{4148}{3}s^4 + \frac{64198}{3}s^5 + \cdots
\Sigma_2 = 2s^2 + \frac{197}{2}s^3 + \frac{33637}{18}s^4 + \frac{2402477}{72}s^5 + \cdots
\Sigma_3 = -12s^2 - \frac{267}{2}s^3 - \frac{19295}{18}s^4 - \frac{933155}{144}s^5 + \cdots
\Sigma_4 = 1830s^3 + \frac{112720}{3}s^4 + \frac{47200115}{72}s^5 + \cdots$$

Singularity structure of PF equation determines radius of convergence

Discriminant:
$$\Delta(\mathcal{L}_l) = s \prod_{j=0}^{\lfloor \frac{l+1}{2} \rfloor} \left(1 - s(l+1-2j)^2\right)$$

Moduli space: $\mathbb{P}^1 \setminus \left(\bigcup_{j=0}^{\lfloor \frac{l+1}{2} \rfloor} \left\{ \frac{1}{(l+1-2j)^2} \right\} \cup \{0\} \cup \{\infty\} \right)$



Equal Mass Case: λ -Coefficients and $\widehat{\Gamma}$ -Conjecture

Actual Feynman integral as linear combination of Frobenius basis:

$$\mathcal{F}_{\sigma_l} = \sum_{k} \lambda_k^{(l), loc} \varpi_k^{loc} \quad \text{with } \lambda_k^{(l), loc} \in \mathbb{C}$$

Numerical computation of \mathcal{F}_{σ_l} yields λ -coefficients (Bessel function representation, analytic continuation)

We could guess their analytic form ($l \le 20$, ~ 300 digits), e.g. for l = 4 around the MUM point s = 0:

$$\lambda_0^{(4)} = -450\zeta(4) - i\pi \cdot 80\zeta(3) \qquad \qquad \lambda_1^{(4)} = 80\zeta(3) - i\pi \cdot 120\zeta(2)$$

$$\lambda_2^{(4)} = 180\zeta(2) \qquad \qquad \lambda_3^{(4)} = i\pi \cdot 20$$

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Actually, we find a generating functional for it:

$$\sum_{l=0}^{\infty} \lambda_0^{(l)} \frac{x^l}{(l+1)!} = -\frac{\Gamma(1-x)}{\Gamma(1+x)} \mathrm{e}^{-2\gamma x + i\pi x} \qquad \text{and} \qquad \lambda_k^{(l)} = (-1)^k \binom{l+1}{k} \lambda_0^{(l-k)}$$

Geometric argument/interpretation of it? —— Yes! They follow from a (modified) $\widehat{\Gamma}$ -conjecture

$$\operatorname{Im}(\mathcal{F}_{\sigma_l}) = \int_{W_{l-1}} e^{\omega \cdot \mathfrak{t}} \widehat{\Gamma}(TW_{l-1}) + \mathcal{O}(e^{\mathfrak{t}}) \quad \text{and} \quad \operatorname{Re}(\mathcal{F}_{\sigma_l}) = \int_{F_l} e^{\omega \cdot \mathfrak{t}} \widehat{\Gamma}(1 - c_1)^2 \frac{\sin(2\pi c_1)}{2\pi c_1} + \mathcal{O}(e^{\mathfrak{t}})$$

Mirror map: $\mathfrak{t} = \frac{1}{2\pi i} \frac{\overline{\omega}_1}{\overline{\omega}_0}$



Non Equal Mass Case

Similar structure as in equal mass case:

(inhomogeneous) PF equations computed from GKZ method

$$\mathcal{L}_l, S_l$$
 \longrightarrow \mathcal{D} -modul $\{\mathcal{D}_l^{(k)}\}$, inhomogeneities $\{S_l^{(k)}\}$

Larger Frobenius basis (logarithmic solutions split)

$$\varpi_k$$
 for $s=1,\ldots$ (primitive vertical Hodge numbers)

symmetric splitting of λ -coefficients

$$\lambda_k^{(l)}$$
 (" $\sim \lambda_k^{(l)}/\mathrm{Hodge/komb}$ ")

Computed explicitly the non equal mass banana Feynman integral up to $l \leq 4$



Conclusions

Found full analytic structure of l-loop banana Feynman integrals

- ullet Equal mass case explicitly for $l \leq 20$ and general results
- ullet Non equal mass case explicitly for $l \leq 4$ and understanding of splitting

 $\widehat{\Gamma}$ -conjecture allows to proof mathematically our results (Iritani)

Guiding principle: Search for associated CY motive of Feynman graph

whole machinery from algebraic geometry, number theory

Possible extension to other Feynman graphs

- Extension to other CY period integrals (traintracks, ice cream cone, kite, ...)
- General non CY graphs? What structure survives? Underlying motive?



Thank you for your attention

