

Check in PYTHIA8 (reproduce results Nagy and Soper, *arXiv:2002.04125*).

No-emission probability:

$$\Pi(\mu_s^2, \mu_h^2, x) = \exp \left\{ - \int_{\mu_s^2}^{\mu_h^2} \frac{d\mu^2}{\mu^2} \int \frac{dz}{z} \frac{\alpha_s(\mu^2)}{2\pi} \hat{P}(z) \frac{f(x/z, \mu^2)}{f(x, \mu^2)} \right\} \quad (1)$$

The perturbative no-emission probability:

$$\Pi_{pert}(\mu_s^2, \mu_h^2, x) = \exp \left\{ - \int_{\mu_s^2}^{\mu_h^2} \frac{d\mu^2}{\mu^2} \int dz \frac{\alpha_s(\mu^2)}{2\pi} \hat{P}(z) \right\}, \quad (2)$$

Their relation is:

$$\Pi_{pert}(\mu_s^2, \mu_h^2, x) = \frac{f(x, \mu_s^2)}{f(x, \mu_h^2)} \Pi(\mu_s^2, \mu_h^2, x) \quad (3)$$

PDF property:

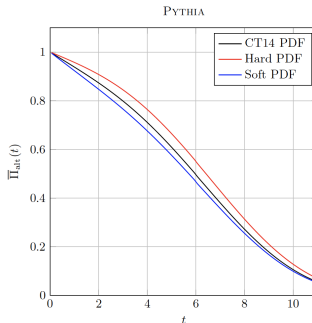
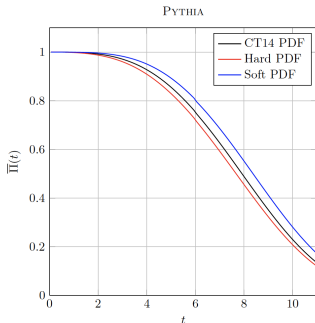
$$\Pi_{alt}(\mu_s^2, \mu_h^2, x) = \frac{f(x, \mu_h^2)}{f(x, \mu_s^2)} \Pi(\mu_s^2, \mu_h^2, x) \quad (4)$$

$$\stackrel{?}{=} \exp \left\{ - \int_{\mu_s^2}^{\mu_h^2} \frac{d\mu^2}{\mu^2} \int dz \frac{\alpha_s(\mu^2)}{2\pi} \hat{P}(z) \right\} \quad (5)$$

$$\stackrel{?}{=} \Delta_a(\mu_h^2, \mu_s^2) \quad (6)$$

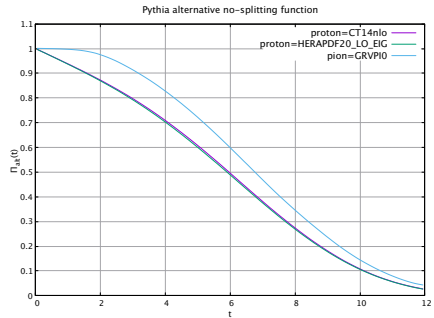
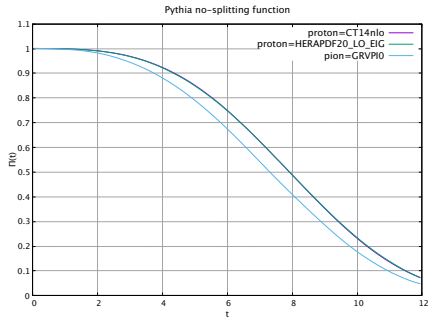
- The alternative no-emission probability should not depend on the PDFs to reproduce the correct Sudakov.
- Procedure:
 - Determine scale of 1st emission in the backward shower
 - Fill histograms for Π and Π_{alt}

Evolution of parton showers and parton distribution functions, Z. Nagy and D.E. Soper, arXiv:2002.04125 [hep-ph]

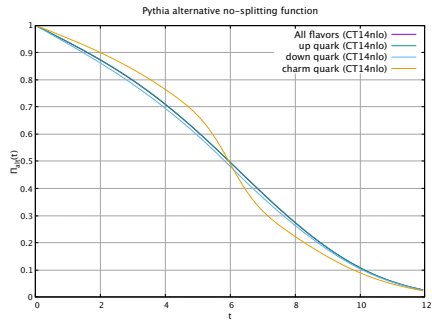
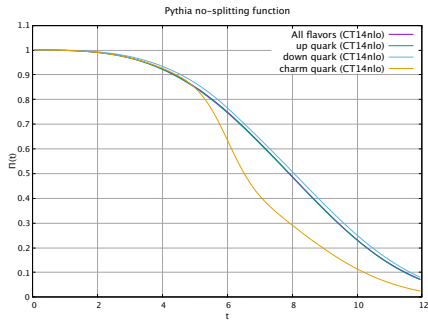


Conclusion: "Differences in the alternative no-emission probability (right) are not much smaller than the differences in the no-emission probability (left)."
⇒ PDF property does not hold.

Test with proton and pion PDFs



Test with all flavours and single flavour selection



Generate Drell-Yan events using PYTHIA8 with an e^+e^- mass range of:
 $2500 \text{ GeV} < \mu_h < 2700 \text{ GeV}$.

- Starting scale of the shower is μ_h
- Scale of the first splitting μ_s
- 'Shower time':

$$t = \ln(\mu_h^2/\mu_s^2)$$

PYTHIA:

- determine the t of the first splitting in an event
- calculate the following quantities of the shower:
 - Probability for first splitting at t : $\rho(t)$
 - No-emission probability: $\Pi(t) = \int_t^\infty d\tau \rho(\tau)$
 - Alternative probability function: $\Pi_{alt}(t) = \frac{f(x, \mu_h^2)}{f(x, \mu_s^2)} \Pi(t)$