# Constructing the 6 Loop $\mathcal{N}=4$ sYM Integrand 

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## Motivations

Goal of program: UV behavior of 7 loop $\mathcal{N}=8$ SUGRA Why?

- SUSY arguments predict $L=7$ counterterm in $D_{c}=4$ (Bossard, Howe, Stelle; Green, Russo, Vanhove; Beisert, Elvang, Freedman, Kiermaier, Morales, Stieberger; many more)
- Similar counterterms proven absent for $\mathcal{N}=4,5$ at $L=\mathcal{N}-1$ (Bossard, Howe, Stelle, Vanhove; Bern, Davies, Dennen; Bern, Davies, Dennen, Huang)
- Improved behavior observed in $D=4$ kinematics (AE, Hermann, Parra-Martinez, Trnka)


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History of direct calculations:
- $1 \& 2$ loops '80 -'90s (Green, Schwarr, Brink; Bern, Dixon, Dunbar, Perelstein, Rozowsky)
- 3 loops '07-'10 (Bern, Carrasco, Dixon, Johansson, Kosower, Roiban)
- 4 loops '09-'12 (Bern, Carrasco, Dixon, Johansson, Roiban)
- 5 loops 2018 (Bern, Carrasco, Chen, AE, Johansson, Parra-Martinez, Roiban, Zeng)


## Computational Challenges

| Obstacle | Solution |
| :---: | :---: |
| Many Feynman diagrams, cancellations between diagrams | Cuts contain minimal needed data |
| Cuts from state sums: 256 states per cut propagator | Double copy: $\mathrm{GR}=\mathrm{YM}^{2}$ |

sYM state sums still hard, Color-kinematics + tricks for special cuts rapidly exploding Dirac traces

Problems with CK at 5L - still
New recursive tools for cuts \& integrands need to solve previous problems

## X Identity

Four-point ordered YM tree amplitudes only have $s$ and $t$ channel poles.
What if we try to "sit on the $u$ pole" anyway, via $p_{3} \rightarrow p_{1}, p_{4} \rightarrow p_{2}$ ?


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Dimensionless, must respect all symmetries: Can only get identity insertions! The diagram disconnects!


Same for all other supersymmetric states

## H Identity

Can we find a similar identity that maintains planarity?

- Consider $\left.s_{13} A(1,3,2,4)\right|_{p_{3}=p_{1}}$ : cut with zero momentum exchange - Apply $s_{13} A(1,3,2,4)=s_{14} A(1,2,3,4)$ with X ID:


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e.g.: $\quad s_{13} A(1,3,2,4)=\frac{t_{8} F^{4}}{s_{23}} \xrightarrow{p_{3}=p_{1}}$

$s_{14}\left(\epsilon_{1} \cdot \epsilon_{3}\right)\left(\epsilon_{2} \cdot \epsilon_{4}\right)$

Same for all other supersymmetric states. Extends to (super)gravity. N.B.: Physical identity (soft factorization), not heuristic rule.

## Method of Maximal Cuts

Systematic integrand construction from cuts (1) Enumerate diagram basis, striate by cut depth ( $k$ )
(1) Build all cubic vacuums
(2) Attach four external legs
(3) Collapse internal legs
(4) Cubic $=$ max; one quartic $=$ next-to-max; $\ldots$
 from lower- $k$ and constructing new numerators

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Trying to determine $\quad \mathcal{P}_{\gamma, \text { ans }}^{(k)}=\mathcal{C}_{\gamma}^{(k)}-\mathcal{R}_{\gamma, \mathrm{MMC}}^{(k)} \underbrace{}_{\text {Known higher cuts: }} \sum \frac{n}{p^{2}}$

## Integrand Search Space



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## Applications of X ID in MMC

X Identity: Evaluate $\mathcal{C}^{(k)}$ directly from limit of higher-loop cut


Challenges:
(1) Edge crossing is NP-Hard ${ }^{1}$
(2) Quickly outpace known planar cuts ( $11+$ loops)
(3) Only works for color-ordered cuts

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${ }^{1} \mathrm{ex}$ : There're better crossing schemes than in the diagram. Can you find one?

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H Identity: Evaluate $\mathcal{P}$ via constraining limits, on which $\mathcal{C}_{\gamma}^{(k)} \rightarrow \mathcal{C}_{\gamma_{L-1}}^{(k)}$


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(1) Need to merge conditions
(2) Many evaluations of lower-loop cuts

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## 6L Numerator Construction



Resolving the conditions
Fix a basis of momentum invariants
$\lim _{\ell \rightarrow 0}$ induces linear relations between invariants: $\pi_{\ell_{m}}$ $\ell_{m} \rightarrow 0$

```
Brute force: use an ansatz
- }\mp@subsup{\mathcal{P}}{\mathrm{ ans }}{}\mathrm{ as literal polynomial ansatz in basis
- H ID gives linear equations between ansatz parameters
More clever: intersection of polynomial ideals
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$$
\begin{aligned}
& \pi_{\ell_{m}} \mathcal{P}_{\gamma, \text { ans }}^{(k)}=h_{\gamma, \ell_{m}} \mathcal{C}_{\gamma \ell_{m}}^{(k)}-\pi_{\ell_{m}} \mathcal{R}_{\gamma, M M C}^{(k)} \\
& \Rightarrow \mathcal{P}_{\gamma, a n s}^{(k)}=\pi_{\ell_{m}} \mathcal{P}_{\gamma, \text { ans }}^{(k)}+\operatorname{ker} \pi_{\ell_{m}} \subset\left\langle\pi_{\ell_{m}} \mathcal{P}_{\gamma, a n s}^{(k)}, \pi_{\ell_{m}}\right\rangle \\
& \mathcal{P}_{\gamma, a n s}^{(k)} \sim \bigcap_{\ell_{m} \in \gamma}\left\langle\pi_{\ell_{m}} \mathcal{P}_{\gamma, \text { ans }}^{(k)}, \pi_{\ell_{m}}\right\rangle
\end{aligned}
$$

## The 6L Integrand

| $N^{k} M$ | 0 | 1 | 2 | 3 | 4 | 5 | $\sum$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| cuts | 5548 | 41649 | 131907 | 216961 | 202271 | 107945 | 706,281 |
| non-zero <br> contacts | 4420 | 16776 | 37373 | 53472 | 32465 | 0 | 144,506 |

Fits on a CD!
Contact terms carry ladder color factors

- Longest numerator: 62,511 terms
- Shortest numerator: 1 term (ladder diagrams)
- Average numerator: 90 terms


## Looking Forward

- UV Integration
- Prepping tools for SUGRA: KLT, IBPs
- Improve efficiency for 7 loops
- Intersection of ideals is senstive to many superficial choices
- Minimize number of limts to evaluate
- Ansatz requires efficient inversion/row reduction
- Application to other theories: QCD, open string eff.
- Cubic representation: generalized double-copy, color-kinematics duality?


## Thanks!

## Questions?

