

Gravitational Wave Echo of Relaxion Trapping

Abhishek Banerjee, E.M., Gilad Perez, Wolfram Ratzinger and Pedro Schwaller
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Eric Madge

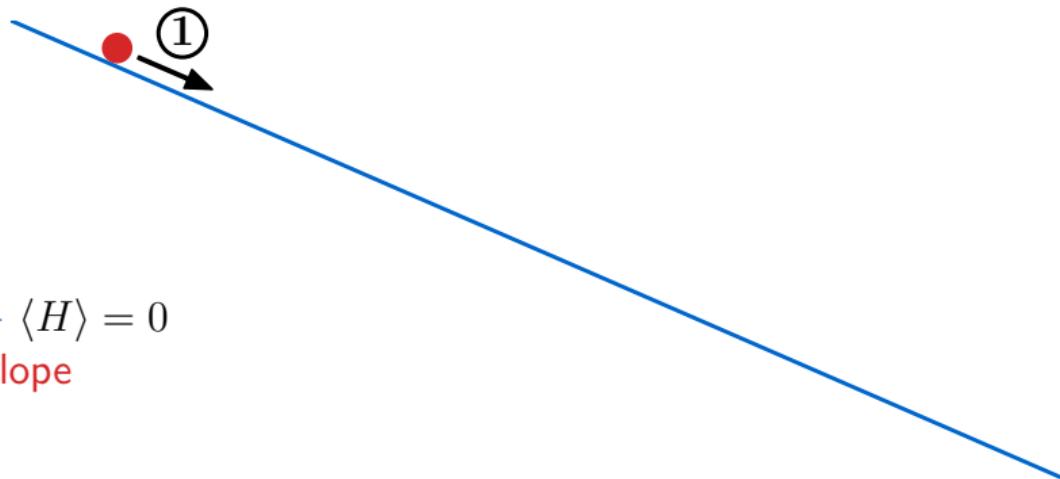
Weizmann Institute of Science

DESY Theory Workshop 2021 – September 21, 2021

Relaxion (during inflation)

[Graham, Kaplan, Rajendran (2015)]

$$V(H, \phi) = \underbrace{(\Lambda^2 - g\Lambda\phi)}_{\mu_H^2(\phi)} |H|^2 + \lambda |H|^4 - cg\Lambda^3\phi - \Lambda_{\text{br}}^4 \frac{|H|^2}{v_H^2} \cos \frac{\phi}{f_\phi}$$

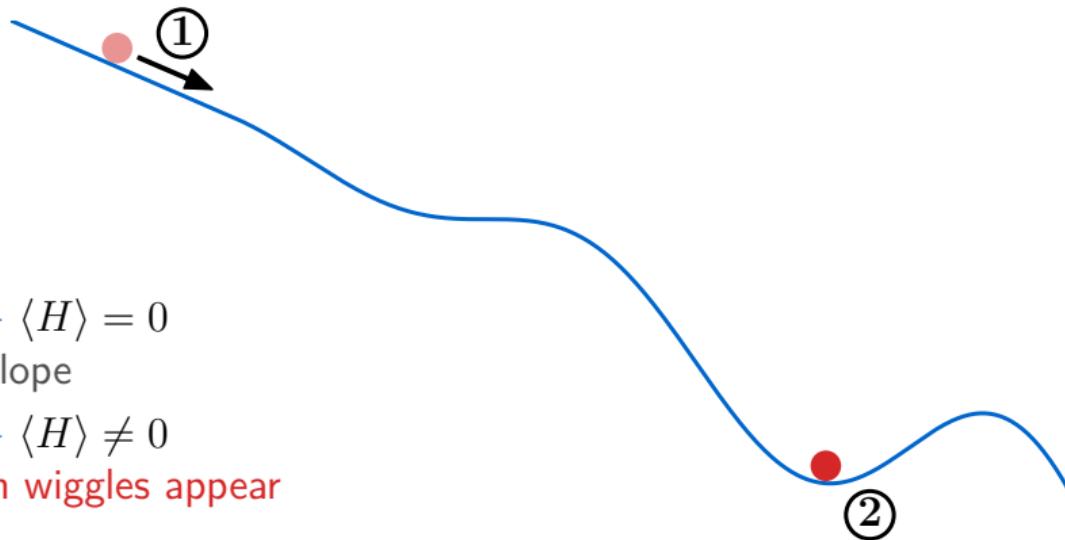


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only linear slope

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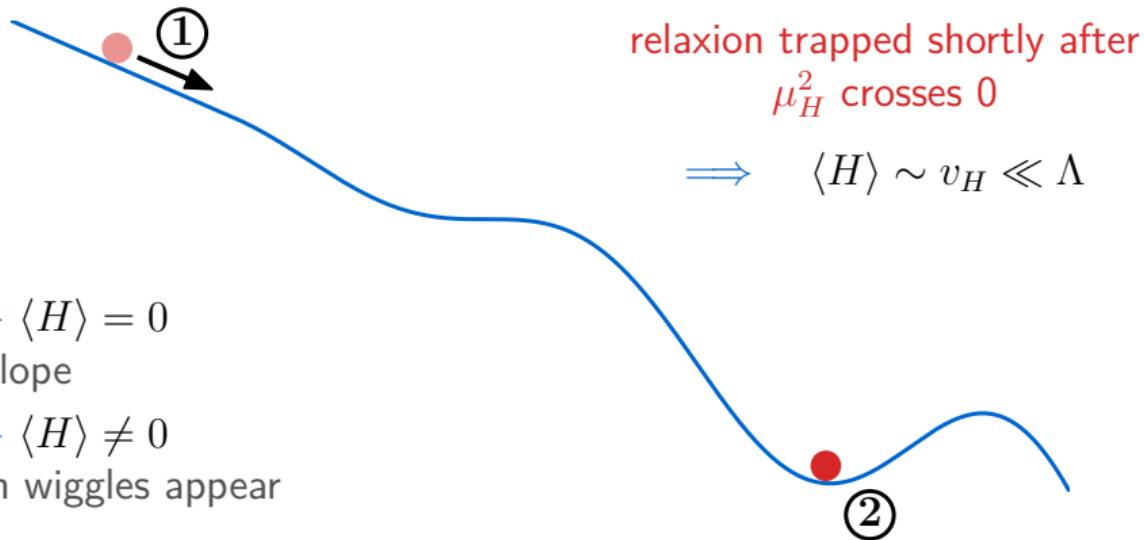


1. $\mu_H^2 > 0 \implies \langle H \rangle = 0$
only linear slope
2. $\mu_H^2 < 0 \implies \langle H \rangle \neq 0$
backreaction wiggles appear

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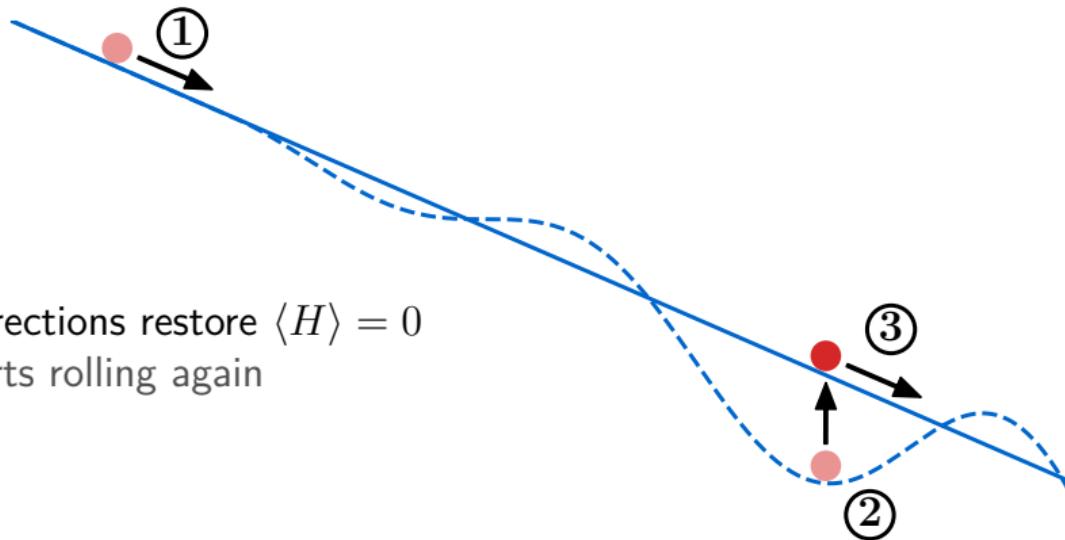
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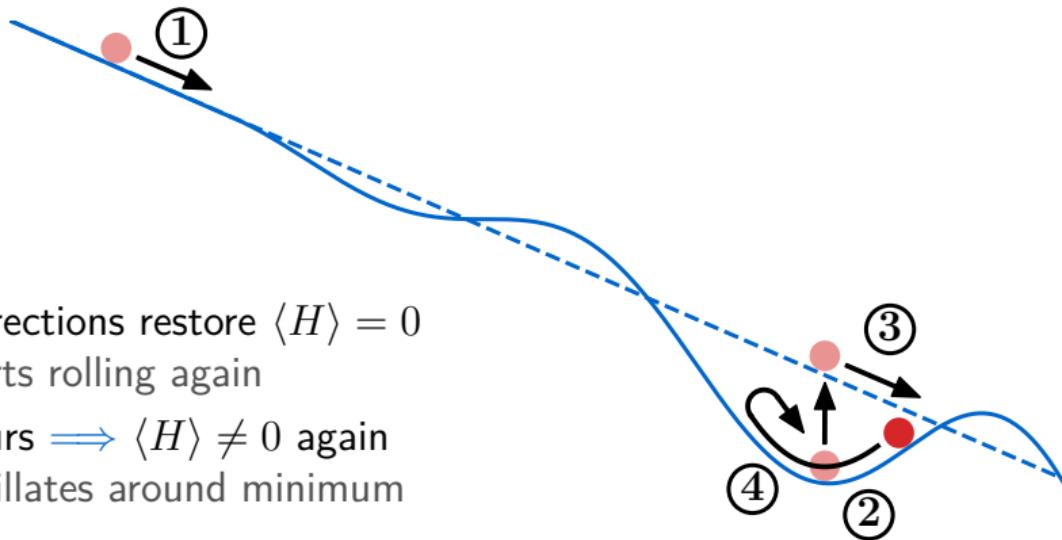


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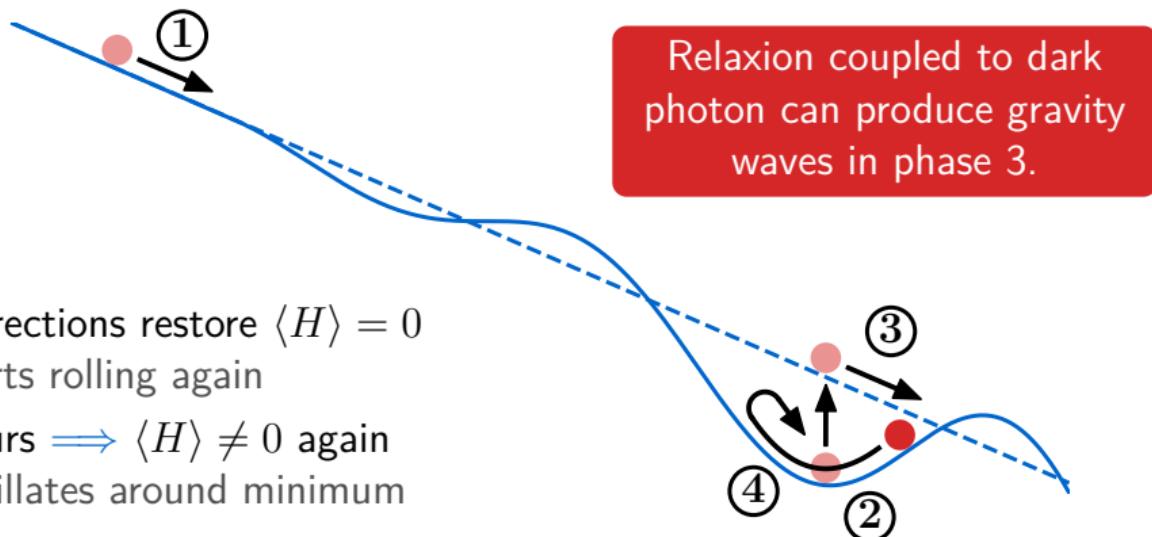


3. thermal corrections restore $\langle H \rangle = 0$
relaxion starts rolling again
4. EWPT occurs $\Rightarrow \langle H \rangle \neq 0$ again
relaxion oscillates around minimum

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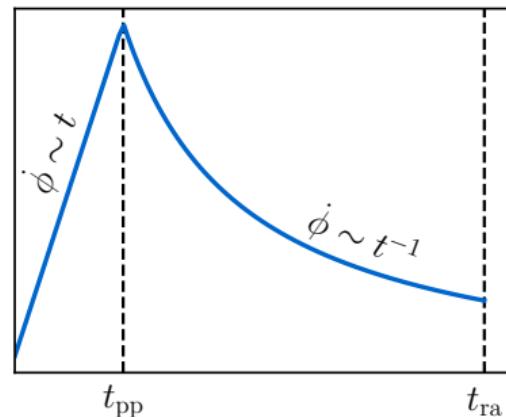
Relaxion and dark photon evolution

- relaxion coupled to dark photon X : $\mathcal{L} \supset -\frac{r_X}{4} \frac{\phi}{f_\phi} X_{\mu\nu} \tilde{X}^{\mu\nu}$

$$\implies \ddot{\phi} + 3H\dot{\phi} - \frac{\Lambda_{\text{br}}^4}{f_\phi} + \frac{r_X}{f_\phi} \frac{\langle \tilde{X}_{\mu\nu} X^{\mu\nu} \rangle}{4a^4} = 0$$

- relaxion reaches terminal velocity when $\frac{r_X}{4a^4} \langle \tilde{X} X \rangle \sim \Lambda_{\text{br}}^4$:

$$\frac{\dot{\phi}}{f_\phi} = \frac{\xi H}{r_X} \quad \xi \sim \mathcal{O}(10 - 100)$$



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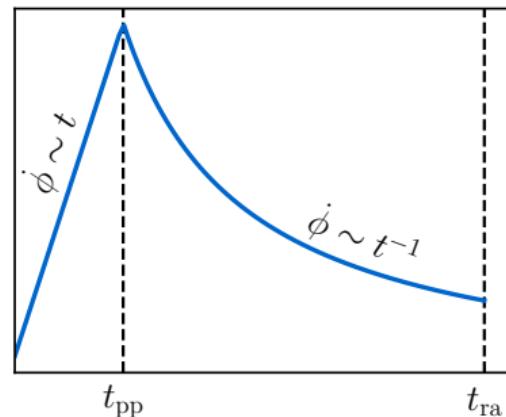
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$$X''_\lambda(\tau, k) + \left(k^2 - \lambda k \frac{r_X \phi'(\tau)}{f_\phi} \right) X_\lambda(\tau, k) = 0$$



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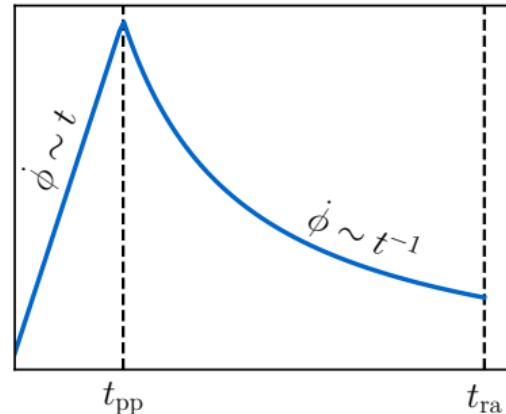
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- anisotropic stress in dark photon energy-momentum tensor sources GWs
 \implies stochastic GW background



Gravitational wave spectra

- peak frequency:

$$f_{\text{peak}} \sim 2 \frac{k_{\text{peak}}^X}{a} \sim \frac{a_{\text{ra}}}{a_0} \xi H_{\text{ra}} \sim 1 \mu\text{Hz} \left(\frac{\xi}{25} \right) \left(\frac{T_{\text{ra}}}{1 \text{ GeV}} \right)$$

- peak amplitude:

$$\Omega_{\text{GW}}^{\text{peak}} \sim \frac{\left(\rho_X^{\text{ra}} / f_{\text{peak}}^{\text{ra}} \right)^2}{\rho_c M_{\text{Pl}}^2} \left(\frac{a_{\text{ra}}}{a_0} \right)^4 \sim 10^{-10} \left(\frac{25}{\xi} \right)^2 \left(\frac{m_\phi f_\phi}{T_{\text{ra}}^2} \right)^4$$

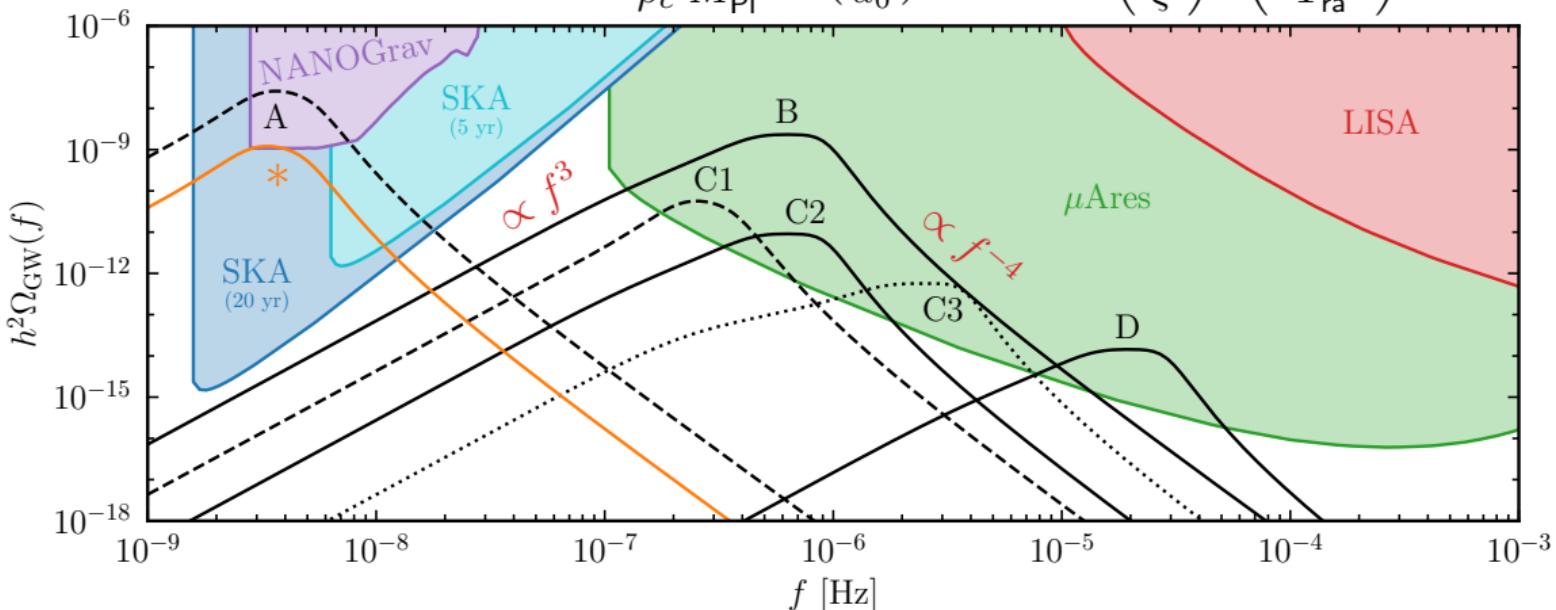
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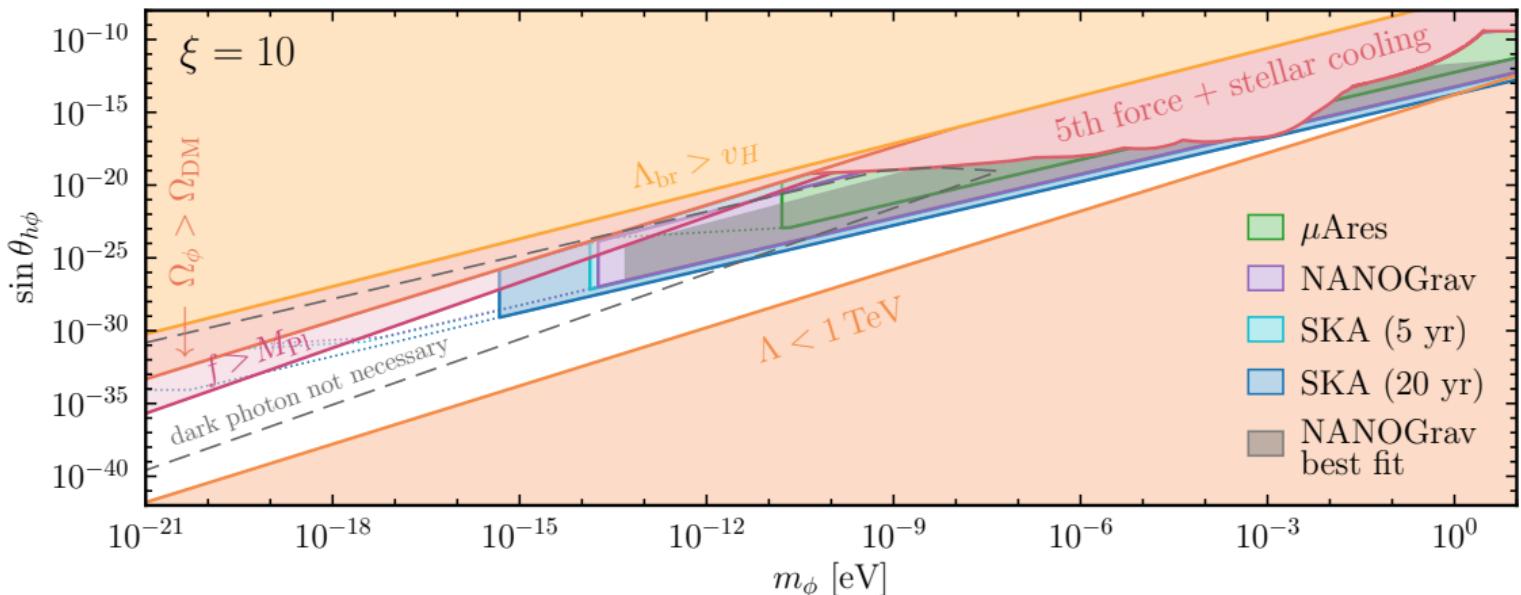
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$$f_{\text{peak}} \propto \xi T_{\text{ra}} \quad h^2 \Omega_{\text{GW}}^{\text{peak}} \propto \xi^{-2} m_{\phi}^{-12} (\sin \theta_{h\phi})^{12} \Lambda^{-8} T_{\text{ra}}^{-8}$$



$$T_{\text{BBN}} \sim 10 \text{ MeV} \lesssim T_{\text{ra}} \lesssim v_H \sim 170 \text{ GeV}$$

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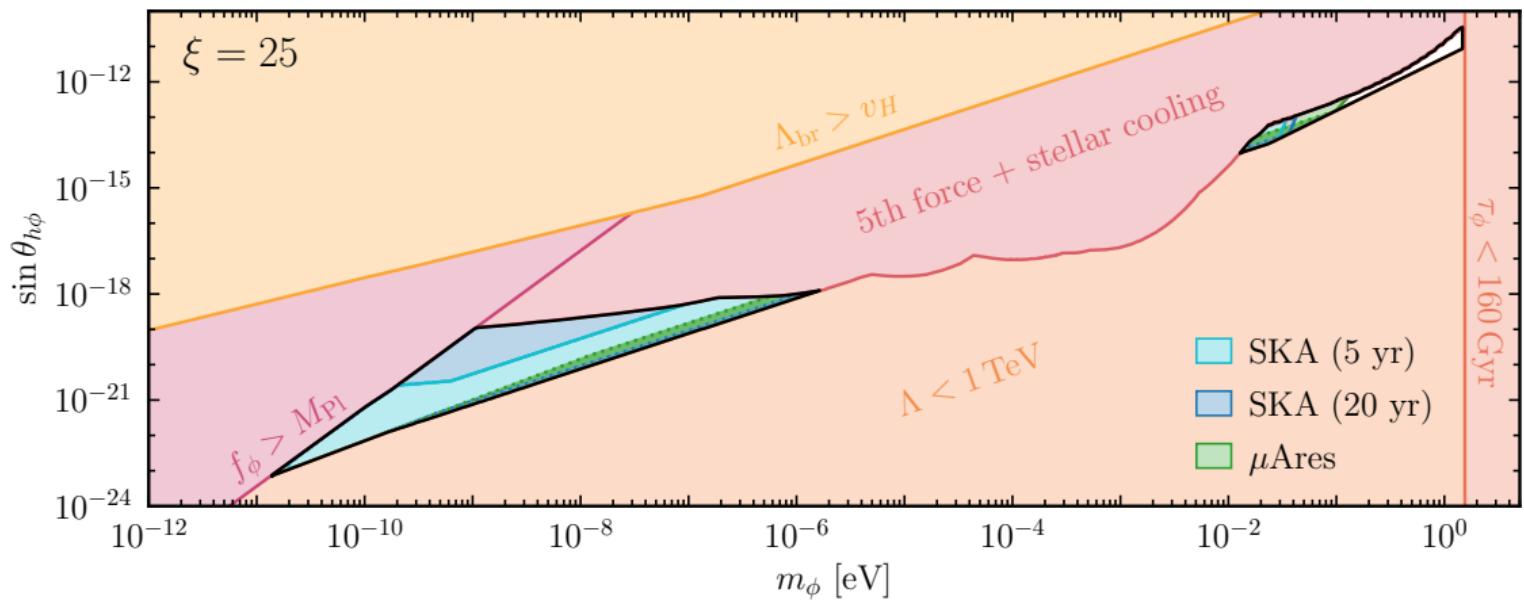
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Relaxion dark matter

displaced relaxion oscillates \implies ultra-light dark matter

[Banerjee, Kim, Perez (2018)]



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animation_slide_6.mp4

Conclusion

- Relaxion coupled to dark photon can produce gravitational waves:

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- SKA will be able to probe a part of the relaxion dark matter parameter space
- large portion of the parameter space will be accessible to μHz observatories such as μAres
- if the relaxion is not DM, we obtain constraints from NANOGrav

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Thank you for your attention!