

# Statistical Analysis of the $N_{\text{DW}} = 1$ QCD Axion Mass Window from Topological Defects

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[arxiv.org/abs/2108.09563](https://arxiv.org/abs/2108.09563)

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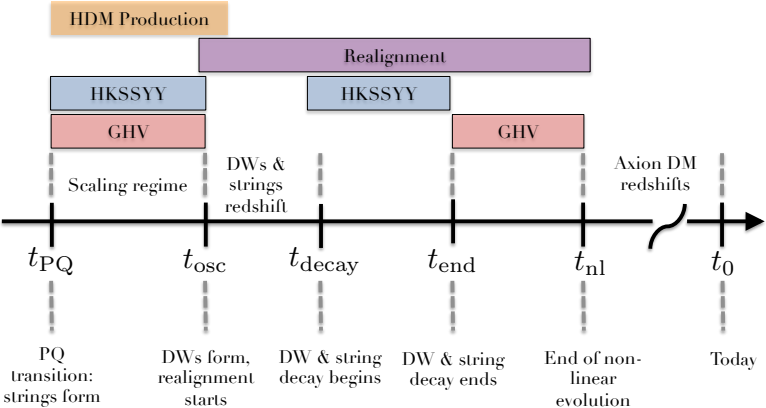
Georg-August Universität Göttingen

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# Motivation

- ▶ Axions are a promising dark matter candidates.
- ▶ Symmetry breaking after inflation  $\Rightarrow$  axion misalignment angle varies over one Hubble patch, in principle no open parameter.
- ▶ Contribution from **topological defects** makes computation more difficult, **Issue for the last decades**.
- ▶ What QCD axion mass do we expect if it is the dark matter?
- ▶ **Goal: Provide framework for comparing different results including statistical errors.**

# Different Contributions



# Different Contributions

**Focus on string scaling in this talk.**

# What are Topological Strings?

- ▶ Axion = angular degree of freedom of the PQ field:  $\phi = f_a \arg \Phi$
- ▶ PQ field: “wine bottle potential”

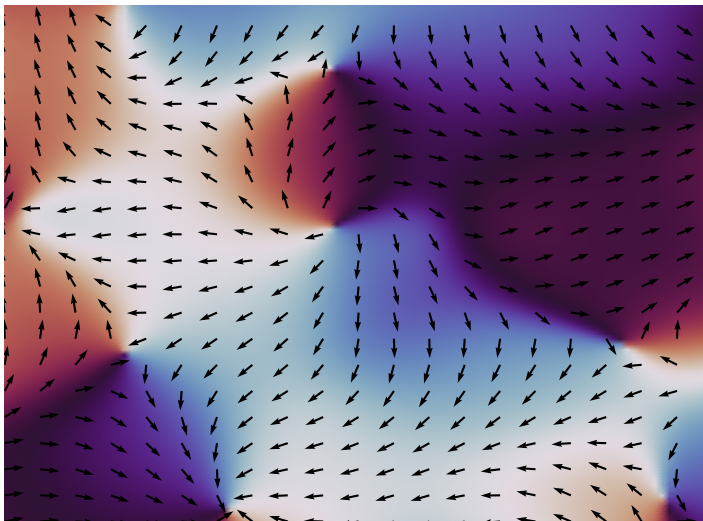
$$V(\Phi) = \lambda_r \left( |\Phi|^2 - \frac{f_a^2}{2} \right)^2. \quad (1)$$

- ▶ radial mode mass  $m_r$ :

$$\lambda_r \equiv m_r^2 / 2f_a^2 \quad (2)$$

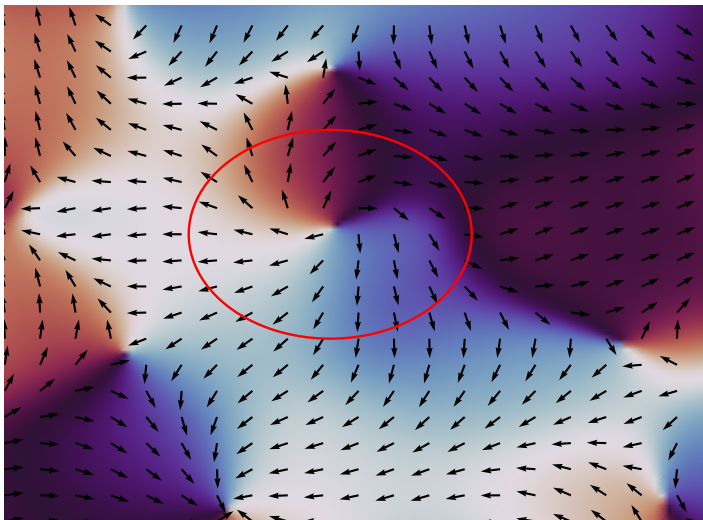
# What are Topological Strings?

- ▶ In our scenario it takes random values across the Hubble patch:



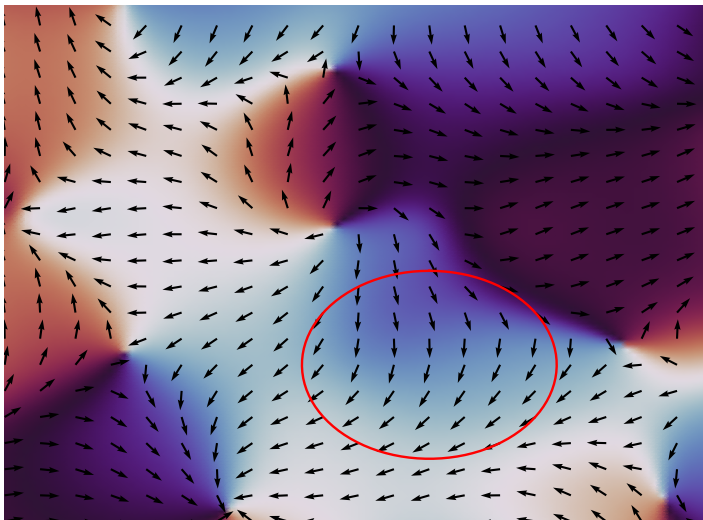
# What are Topological Strings?

- ▶ Closed loops integrate to a winding number.
- ▶ This loop **does** contain a string.



# What are Topological Strings?

- ▶ Closed loops integrate to a winding number.
- ▶ This loop **does not** contain a string.





# What are Topological Strings?

- ▶ In 3D: a one-dimensional structure (strings)
- ▶ Change in the string energy density is converted into axions.

# Difficulties in Simulations

Two scales:

▶ Size of the string core  $\sim 1/m_r$

▶ Size of the Hubble patch  $\sim 1/H$

$\Rightarrow l \equiv \log(m_r/H) \sim \mathcal{O}(70)$ ,  $m_r/H \sim 10^{30}$

Hence:

▶ Not feasible to simulate in the physical regime!!!

▶ Hopeless? No!

# The Attractor Solution

The string network is thought to approach an attractor solution. Characterized by the parameters:

- ▶ The co-moving string length per Hubble volume:

$$\xi(t) \equiv \lim_{V \rightarrow \infty} \left[ \frac{L_{\text{tot}}(V) t^2}{V} \right], \quad (3)$$

- ▶ The emission spectrum into axions:

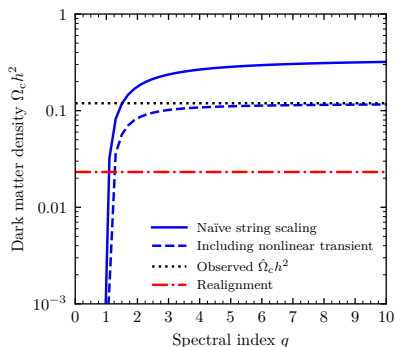
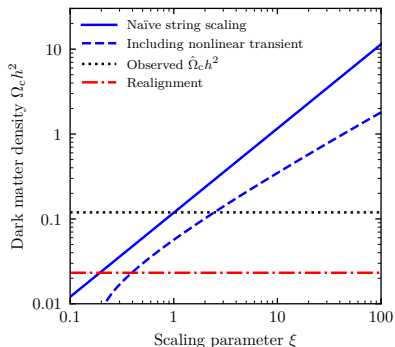
$$\dot{\rho}_s + 2H\rho_s = -\Gamma \approx -\Gamma_a \quad (4)$$

$$\frac{\partial \Gamma_a(t, k)}{\partial k} = \frac{\Gamma_a(t)}{H(t)} F \left[ \frac{k}{H(t)}, \frac{m_r}{H(t)} \right] \quad (5)$$

$$F[x, y] \propto \frac{\left(\frac{x}{x_0}\right)^{q'} \left[ 1 + \Theta(x - x'') \left( \left(\frac{x''}{x}\right)^{q''-q} - 1 \right) \right]}{\left(\frac{x}{x_0}\right)^{q'+1} + 1}, \quad (6)$$

# Relic Density from Strings

Integrating  $F[x, y]$  appropriately over momentum  $k$  and time  $t$  yields the axion number density and hence the relic density.

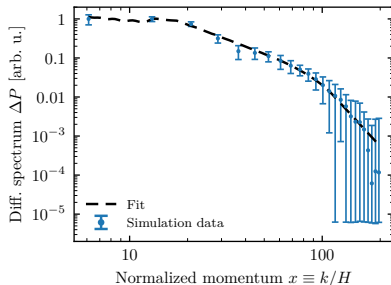
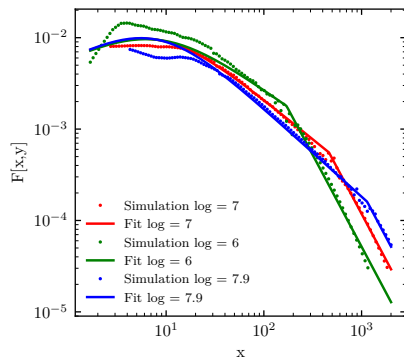


# Where are the disagreements?

- ▶ What values of  $\xi$  and  $q$  does the system approach?
- ▶  $q > 1 \Rightarrow \Omega_{\text{str}} > \Omega_{\text{align}}$
- ▶  $q < 1 \Rightarrow \Omega_{\text{str}} \ll \Omega_{\text{align}}$
- ▶ Constant or scaling-violation (they change/increase with  $l$ )?
- ▶ Our work: Consider the results of two groups:
  - ▶ Takashi Hiramatsu, Masahiro Kawasaki, Toyokazu Sekiguchi, Ken'ichi Saikawa, Masahide Yamaguchi, Jun'ichi Yokoyama (HKSSYY)
  - ▶ Marco Gorghetto, Edward Hardy, Giovanni Villadoro (GHV)

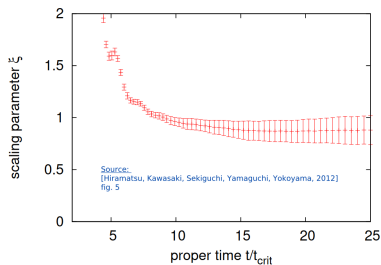
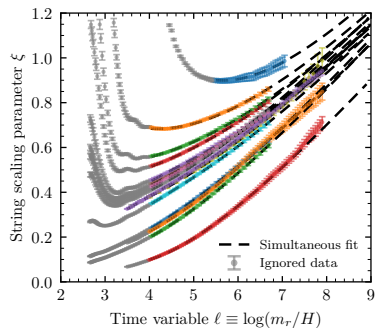
[Hiramatsu, Kawasaki, Sekiguchi, Yamaguchi, Yokoyama, 2010] [Hiramatsu, Kawasaki, Saikawa, Sekiguchi, 2012] [Kawasaki, Saikawa, Sekiguchi, 2015] [Gorghetto, Hardy, Villadoro, 2018]  
[Gorghetto, Hardy, Villadoro, 2020]

# Fitting the Spectra:



- ▶ GHV: scaling violation:  $q = (0.154 \pm 0.057) + (0.1030 \pm 0.0071) /$
- ▶ HKSSYY: non-scaling violation:  $q = 1.44 \pm 0.16$

## Fitting the scaling parameter:



- ▶ GHV: scaling violation:  $\xi = (-1.618 \pm 0.038) + (0.2428 \pm 0.0025)/$
- ▶ HKSSYY: non-scaling violation:  $\xi = 0.87 \pm 0.14$

# Other Effects included

## ▶ Realignment

- ▶ Including QCD axion potential at finite  $T$  and  $g_*(T)$  dependence.
- ▶ Average over oscillations.
- ▶ Numerical averaging over initial field values.  $[-\pi, \pi)$ .

## ▶ Domain Wall Decay

- ▶ Simulated only by the HKSSYY group.
- ▶ Parameters taken from their paper.
- ▶ Only for  $N_{\text{DW}} = 1$ .

## ▶ Non-linear transit

- ▶ Pointed out and simulated by GHV.
- ▶ Large gradients delay the point where the axion mass becomes effective.
- ▶ Suppresses relic density.

## ▶ Thermal Production

- ▶ Coupling to gluons and pions.
- ▶ Upper-bound from  $\Delta N_{\text{eff}}$



# Uncertainties

## Included:

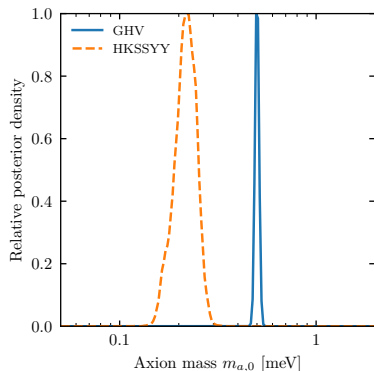
- ▶ *Statistical* errors of the fits of the string parameters  $\xi$ ,  $q$  and  $x_0$  (with bootstrapping).
- ▶ Uncertainties of the domain wall parameters  $\sigma_{\text{DW}}$  and  $\varepsilon$ .
- ▶ Uncertainties of the QCD axion mass inc. temperature dependence.
- ▶ Uncertainties of standard model parameters.

## Not included:

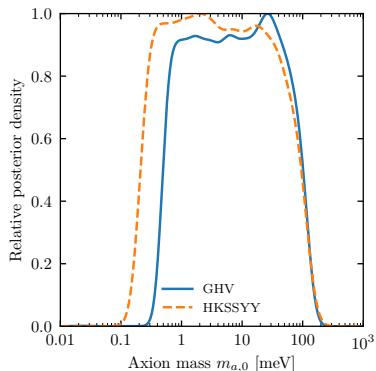
- ▶ Errors on  $g_*(T)$

# Results from the MCMC Scans

$$\Omega_a = \Omega_{\text{DM}}$$



$$\Omega_a < \Omega_{\text{DM}}$$



# Summary and Conclusion

## Work done:

- ▶ Statistical analysis of the axion relic density in the post-inflationary symmetry breaking scenario.
- ▶ Using common framework to analyse different simulations.
- ▶ Results:
  - ▶  $0.49 \text{ meV} < m_a < 84 \text{ meV}$  for GHV
  - ▶  $0.23 \text{ meV} < m_a < 82 \text{ meV}$  for HKSSYY

## Work required:

- ▶ Do these comparisons directly with simulation data.
- ▶ Investigate systematic errors by comparing simulations.
- ▶ Requires work by / collaboration with simulation groups.

# Thank you for your Attention!

## Questions?

I would like to thank:

- ▶ The organizers of this workshop.
- ▶ My supervisors/co-workers Sebastian Hoof and David J. E. Marsh.
- ▶ Laura Covi for encouraging me to submit an abstract to this workshop.
- ▶ Kim W. and Melina A. for their personal support.