Compton spectra reconstruction

Borysova Maryna (KINR) 14/10/20 LUXE weekly technical meeting



Photon spectra reconstruction using Bethe-Heitler pair spectrum

The classical Bethe-Heitler formula (H.Bethe, W.Heitler, Proc.Roy.Soc.A146 (34)83)

$$\Phi (\mathbf{E}_0) d\mathbf{E}_0 = \frac{Z^2}{137} \left(\frac{e^2}{mc^2}\right)^2 4 \frac{\mathbf{E}_{0+}^2 \mathbf{E}_{+}^2 + \frac{2}{3} \mathbf{E}_0 \mathbf{E}_{+}}{(h\nu)^3} d\mathbf{E}_0 \left(\log \frac{2\mathbf{E}_0 \mathbf{E}_{+}}{h\nu mc^2} - \frac{1}{2}\right)$$

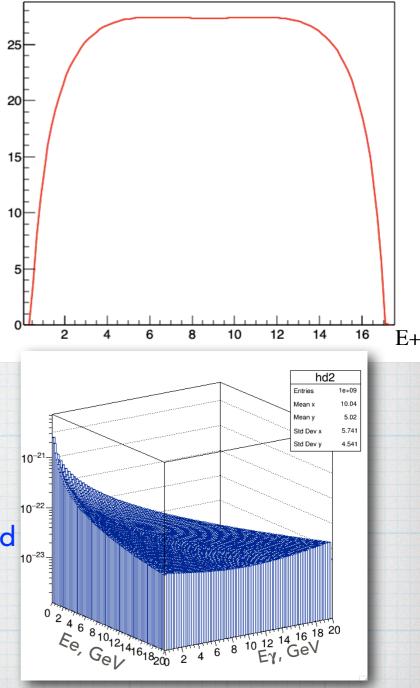
energies involved are large compared with mc²

$\Phi(E+, E_0=E\gamma)$

 $\sigma(E\gamma, Ee) = \Phi(E\gamma, Ee)^*N_a \qquad N_a - Number of atoms$ Photon spectra g(E γ) can be reconstructed by fitting $N(Ee) = \int \sigma(E\gamma, Ee)g(E\gamma)dE\gamma$ Where N(Ee) positron/electron spectra measured in detector after the conversion.

Since σ(Eγ, Ee) depends on number of scatters N_a defined by the thickness of the target the approach can be tested by using the thickness as fit parameter

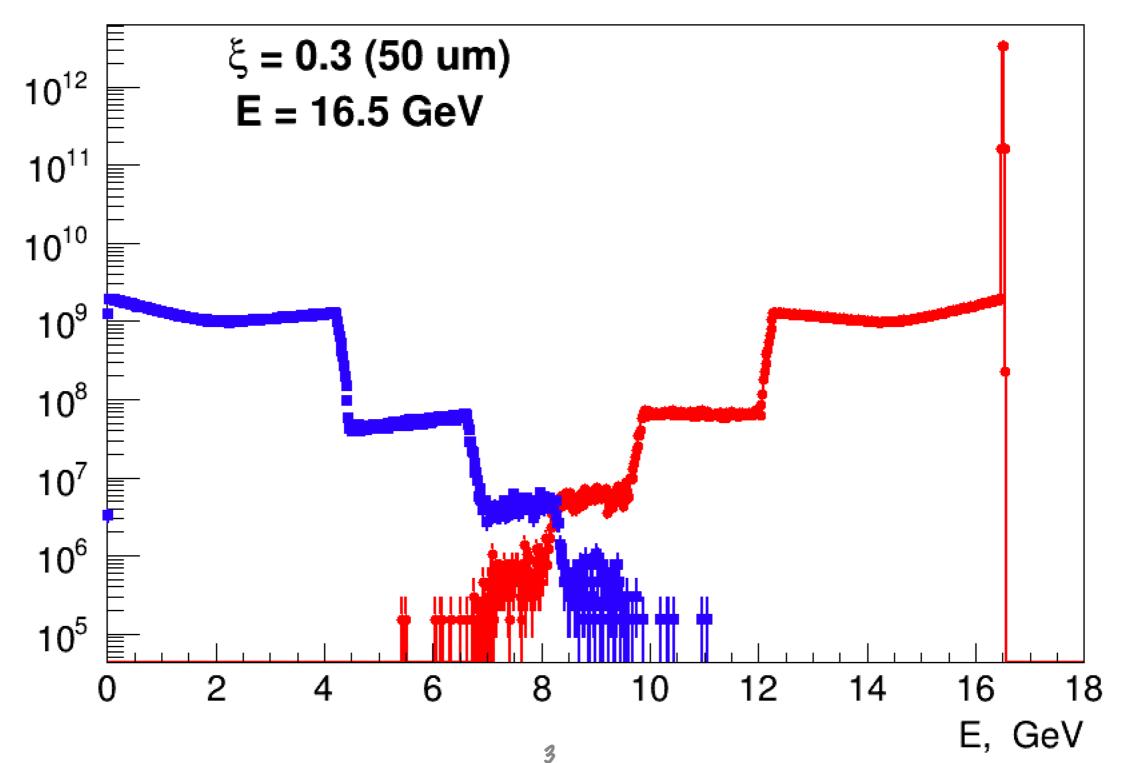
Used Bethe-Heitler class from Geant4, with corrections and extended for various effects (the screening, the pair creation in the field of atomic electrons, correction to the Born approximation, the LPM suppression mechanism, etc.) to calculate differential cross-section



True electron/photon spectra

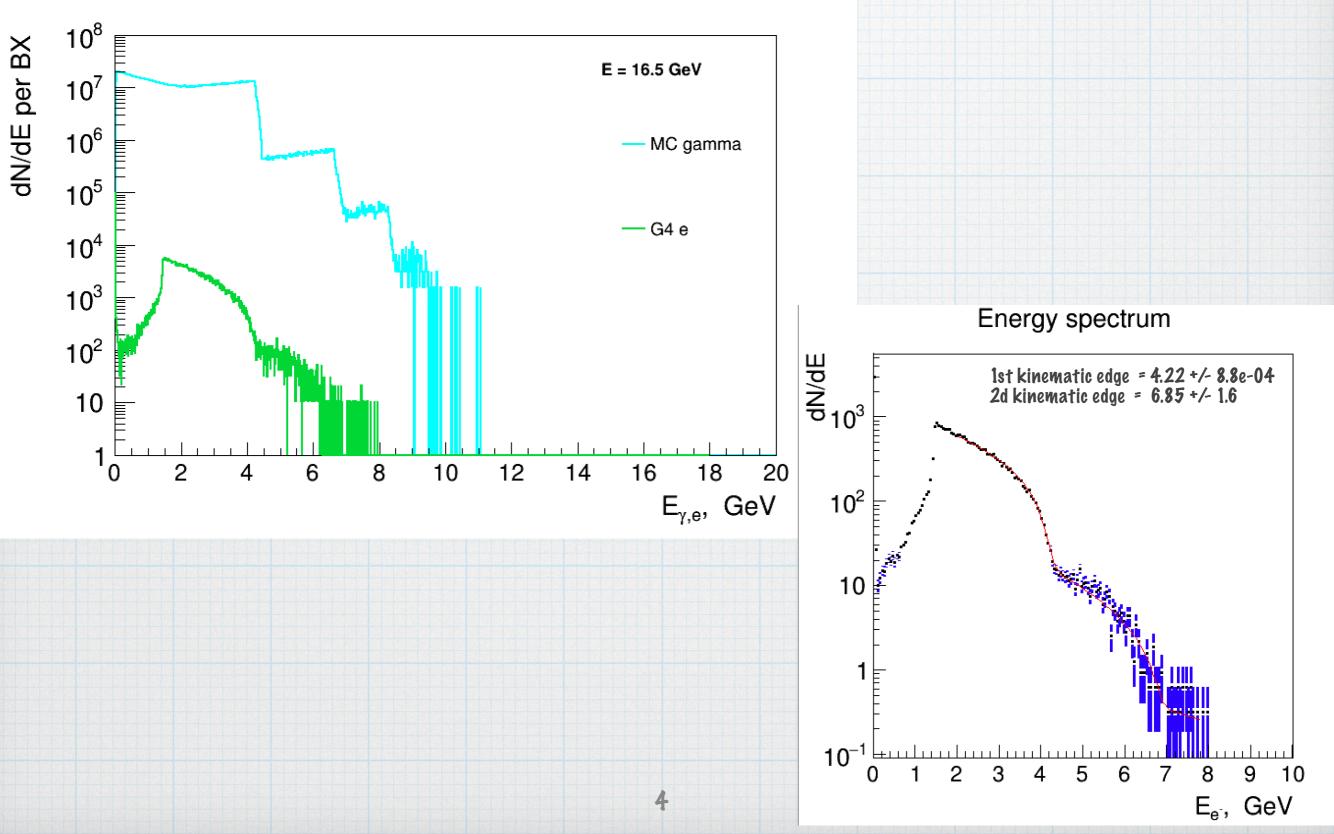
4764BX out of 5000 BX at the laser intensity xi = 0.3 for 16.5 GeV electron beam

(5% of files have NaN so they are ignored)

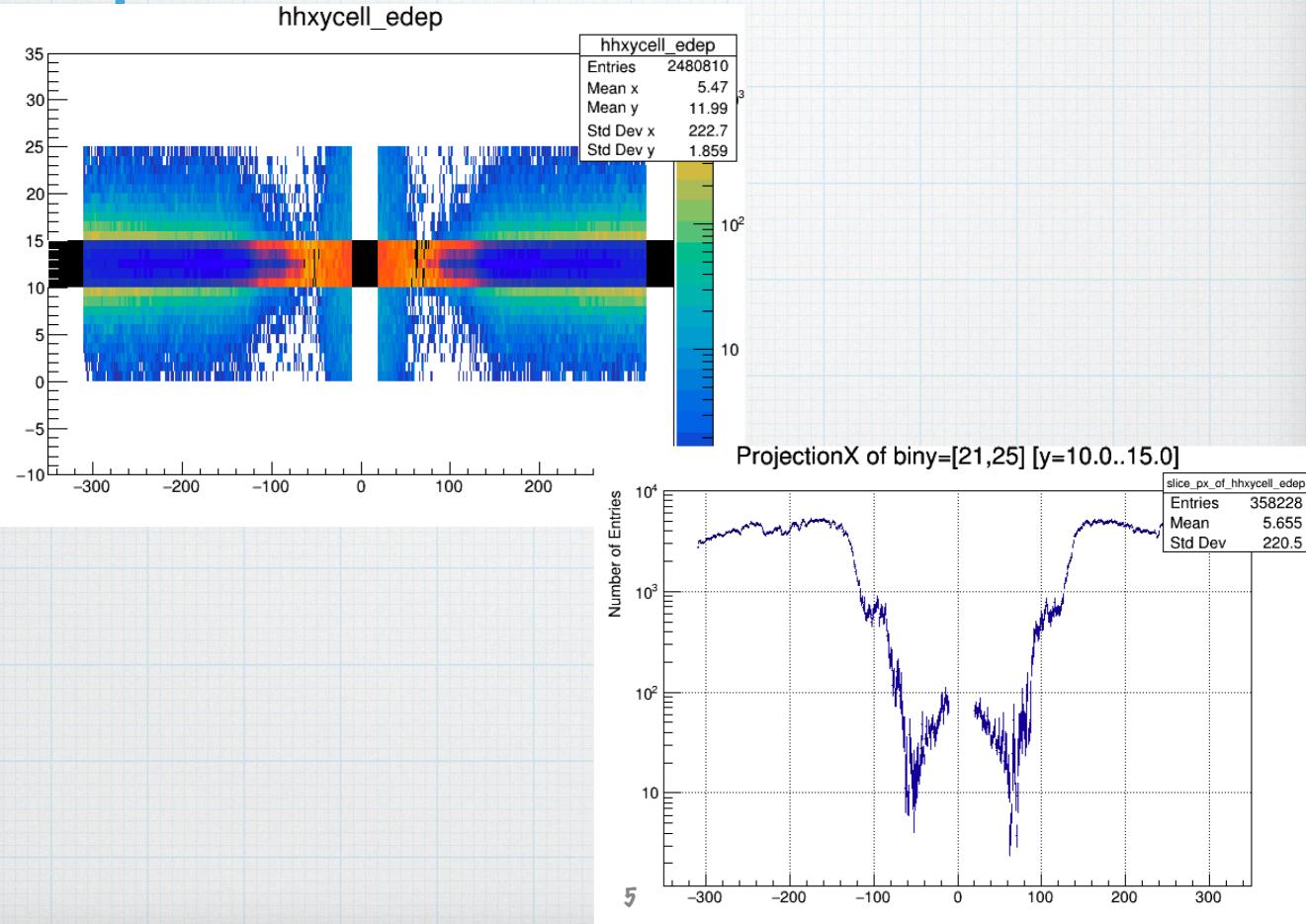


Spectra MC vs G4

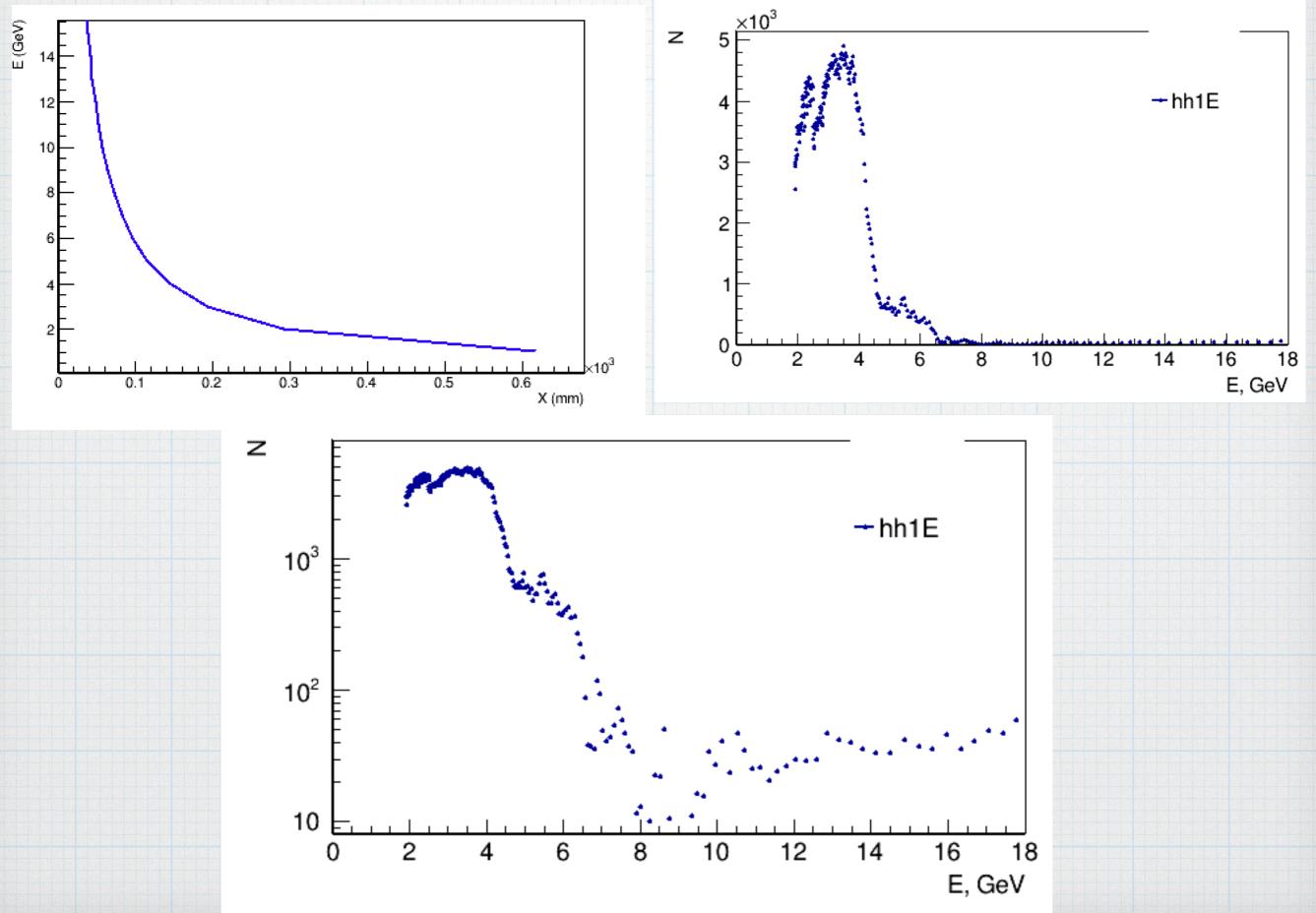
~4700BX out of 5000 BX at the laser intensity xi = 0.3 for 16.5 GeV electron beam



Spectra reconstruction for the LYSO case

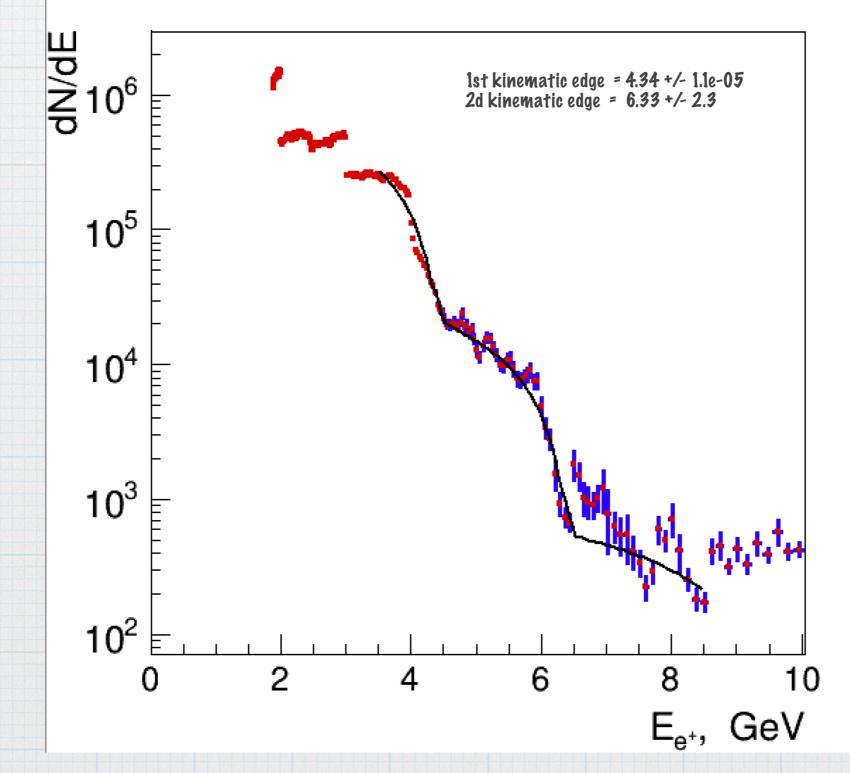


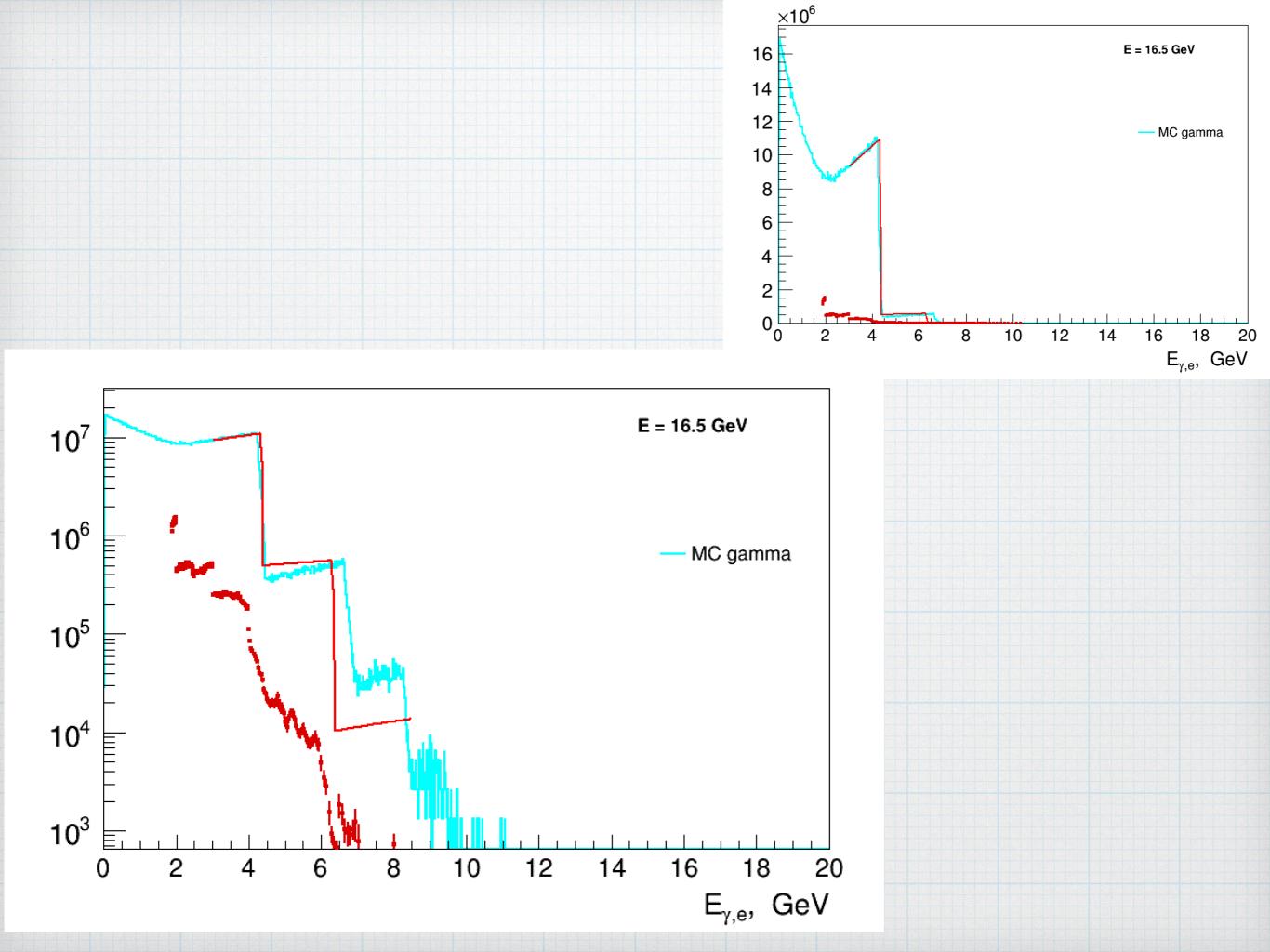
Spectra reconstruction for the LYSO case



Kinematic edges reconstruction

spectrum_positron





High Intensity Compton Scattering

$$e^- + n\omega \to e^- + \gamma$$

The rate of High Intensity Compton Scattering (HICS) is proportional to:

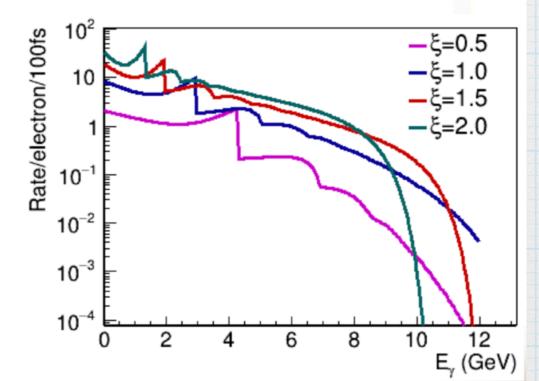
 $\left|\sum_{\mathsf{n}} \delta^{(4)} \left[p_{\mathsf{i}} + k \frac{\xi^3}{2\chi_{\mathsf{i}}} + nk - p_{\mathsf{f}} - k \frac{\xi^3}{2\chi_{\mathsf{f}}} - k_{\mathsf{f}} \right] \right|$

Momentum conservation is a sum over external field photon contributions, *nk*

Even for *n*=0 there is an irreducible contribution:

$$p_{\rm i} + k \frac{\xi^3}{2\chi_{\rm i}} \to p_{\rm i}^2 = m^2(1+\xi^2)$$

- Strong field leads to increase in electron rest mass.
- Observation of Compton edge shift.

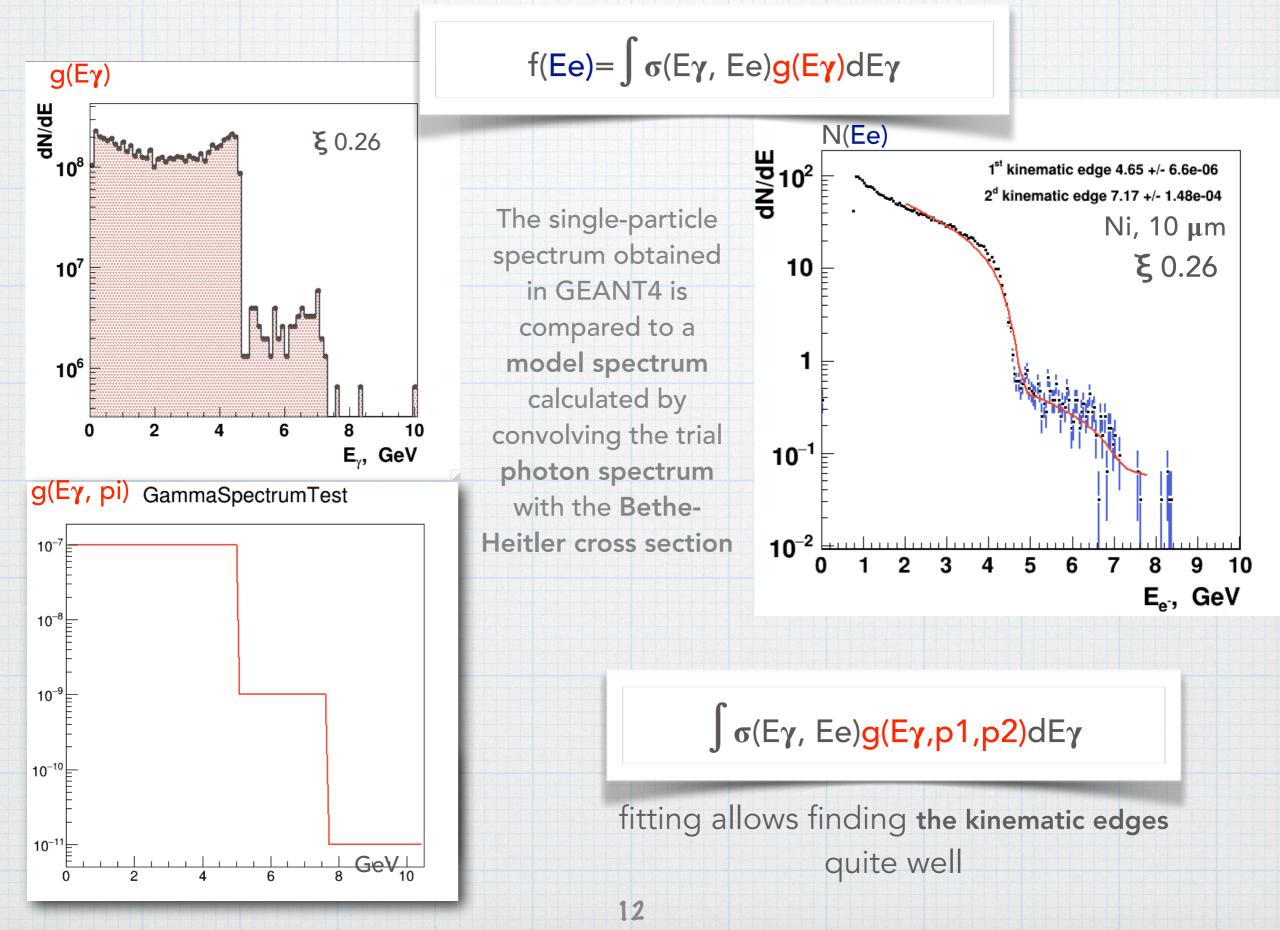


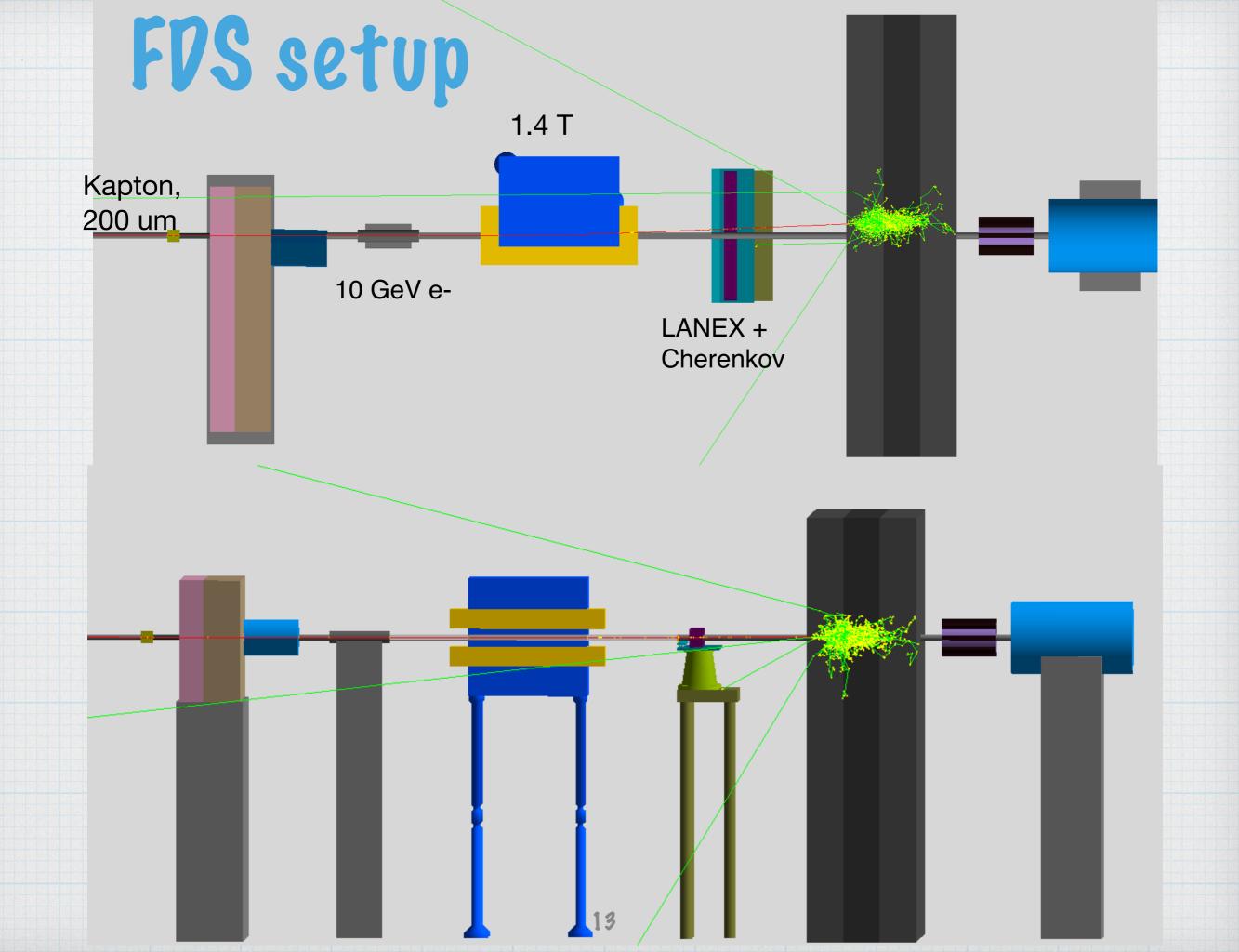


- * The performance of FDS setup with beam pipe from the target to Compton detectors was studied
- * First glimpse at the reconstruction in LYSO



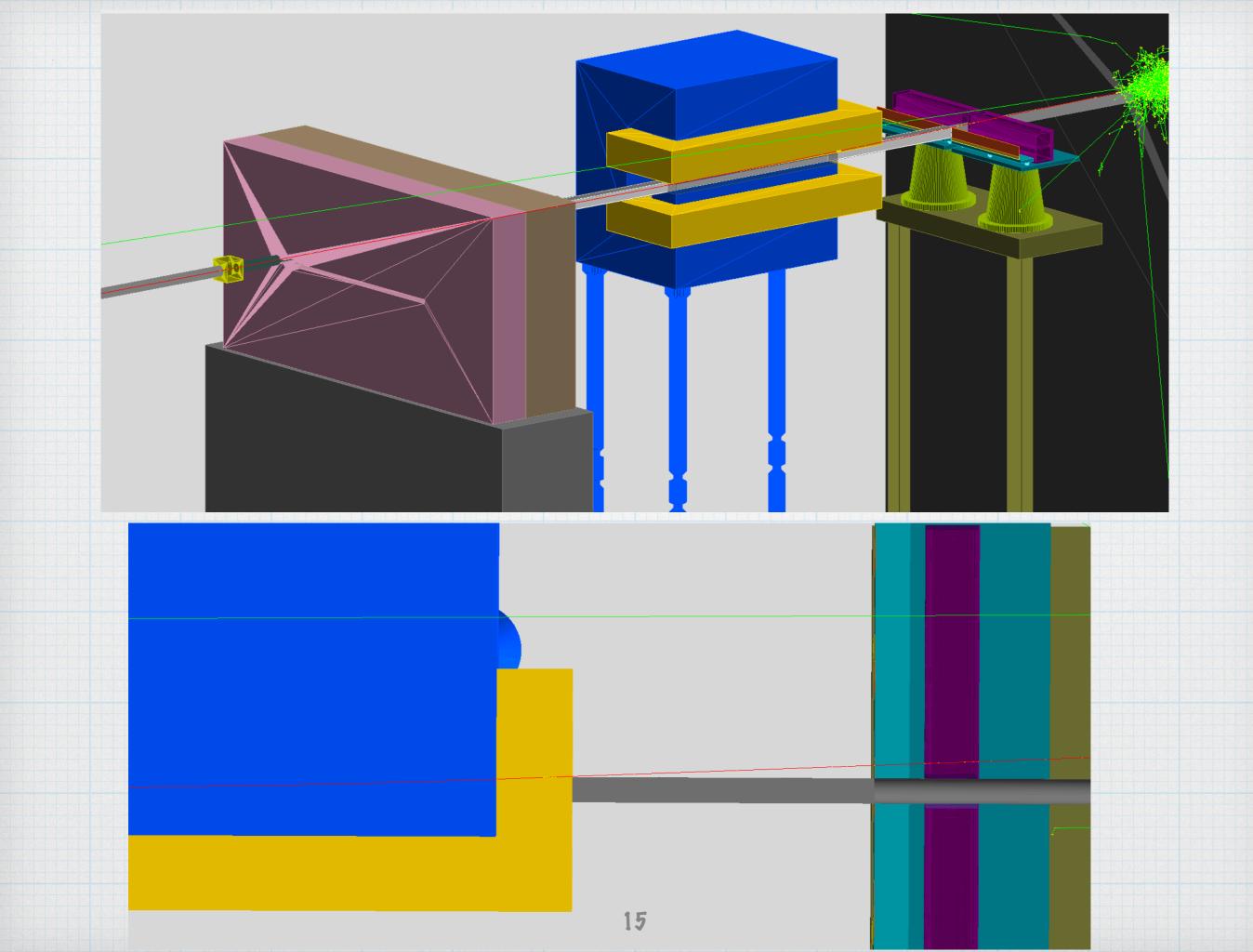
Kinematic edges with accurate pair spectrum



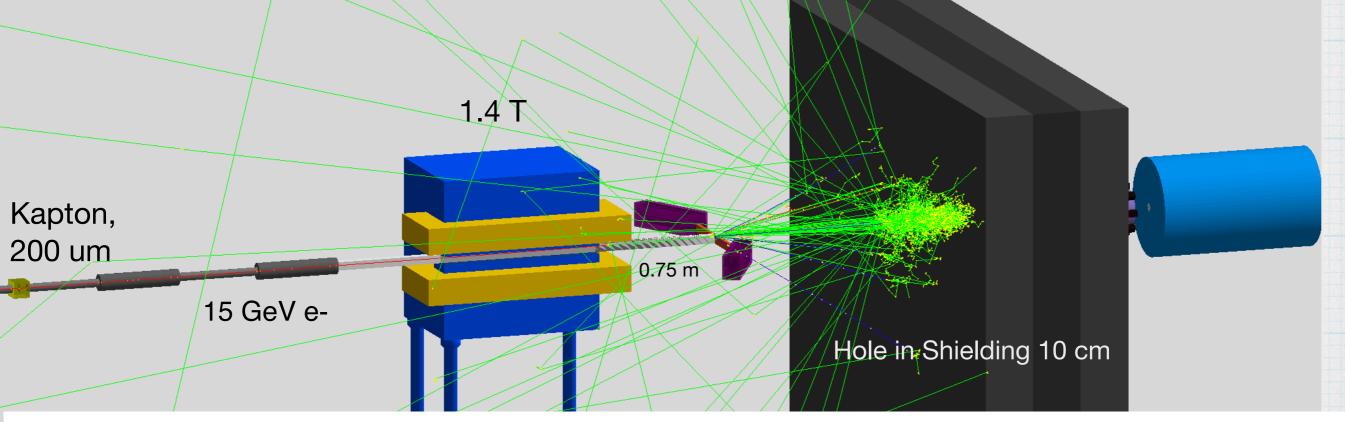


Beam pipe through the electron shielding

14



FDS with LYSO calorimeters



Aug 2020 Data Runs, bunch/pulse crossings completed

Experiment Config	$w_0 = 3\mu m$	$w_0 = 3.5 \mu \text{m}$	$w = 0, 4.0 \mu \text{m}$	$w_0 = 4.5 \mu \text{m}$	$w_0 = 5.0 \mu \text{m}$	$w_0 = 8.0 \mu \text{m}$	$w_0 = 20.0 \mu m$	$w_0 = 50.0 \mu m$	$w_0 = 100.0 \mu \text{m}$
peak SQED ξ	5.12	4.44	3.88	3.45	3.1	1.94	0.78	0.31	0.15
peak SQED χ (16.5 GeV)	0.9	0.79	0.69	0.61	0.55	0.34	0.138	0.055	0.028
JETI40 e-laser 16.5 GeV	10000	1000	1000	1000	1000	1000	500	5000	500

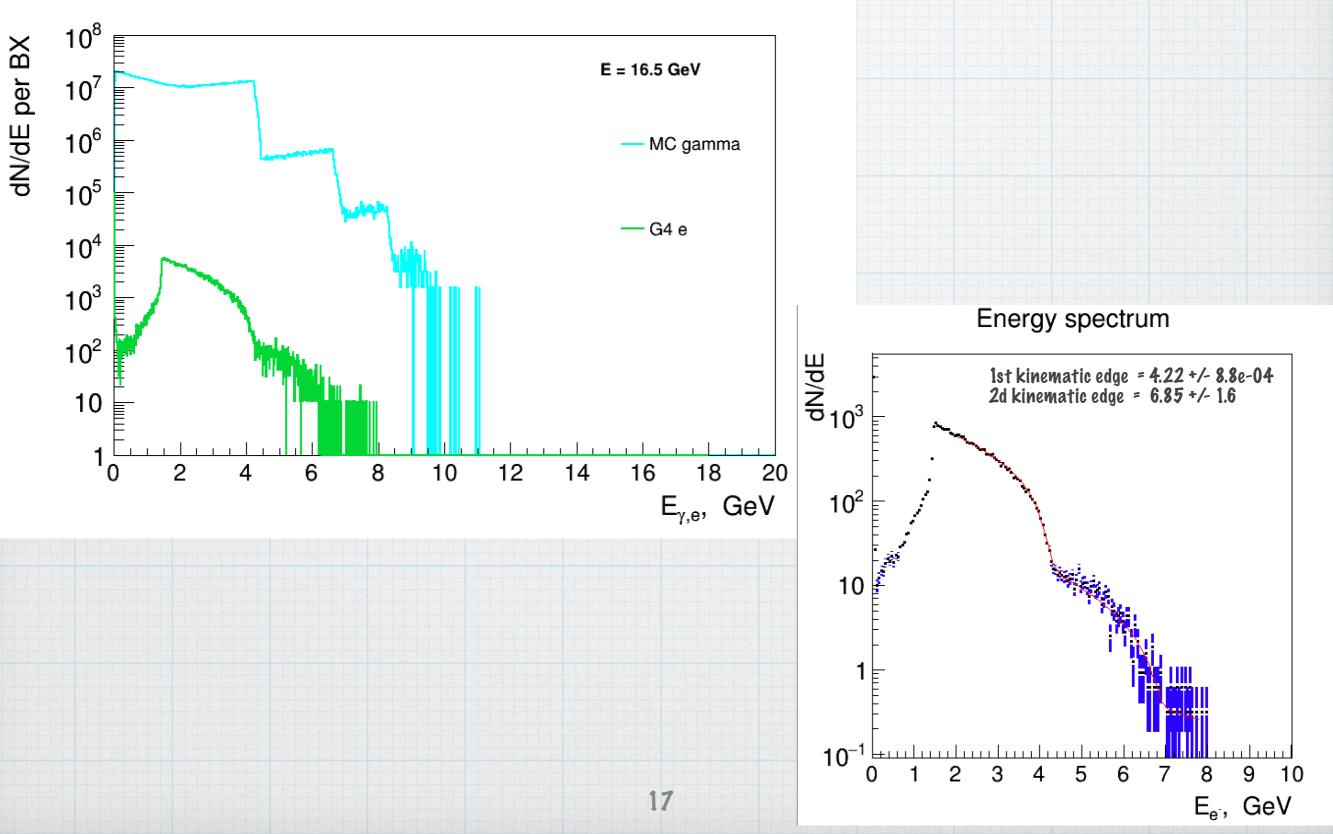
- * The scintillators are modelled as a 15x5x2 cm (x:y:z) layer of lyso material
- * The crystal (bin) size of the scintillators are 2 x 1 mm (finer segmentation in x; the deflection direction) giving 25 x 300 bins.

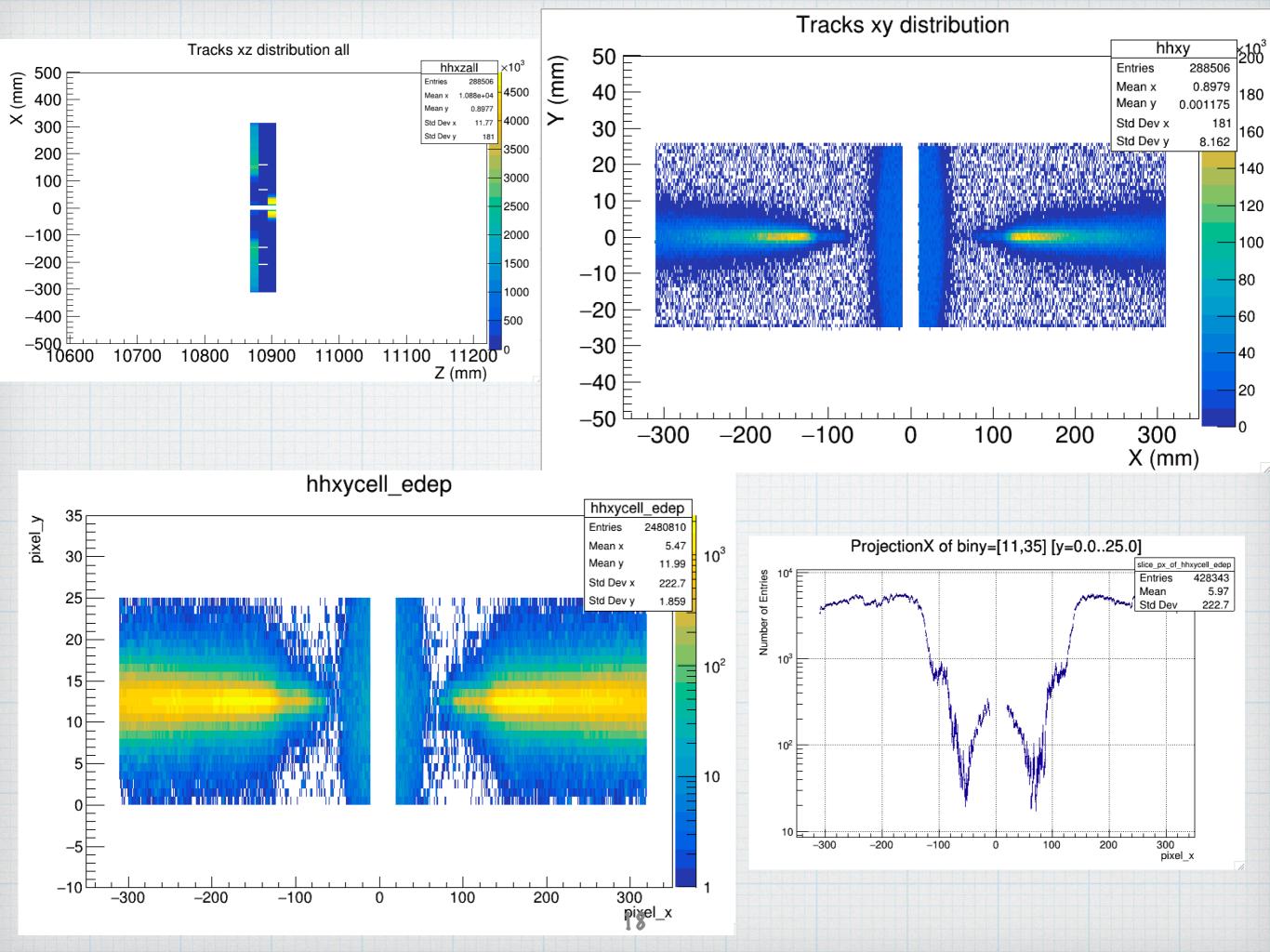
All studies were performed with 5000 BX at the laser intensity xi = 0.3 for 16.5 GeV electron beam

LYSO $(Lu_{1.8}Y_{0.2}SiO_5)$

Spectra MC vs G4

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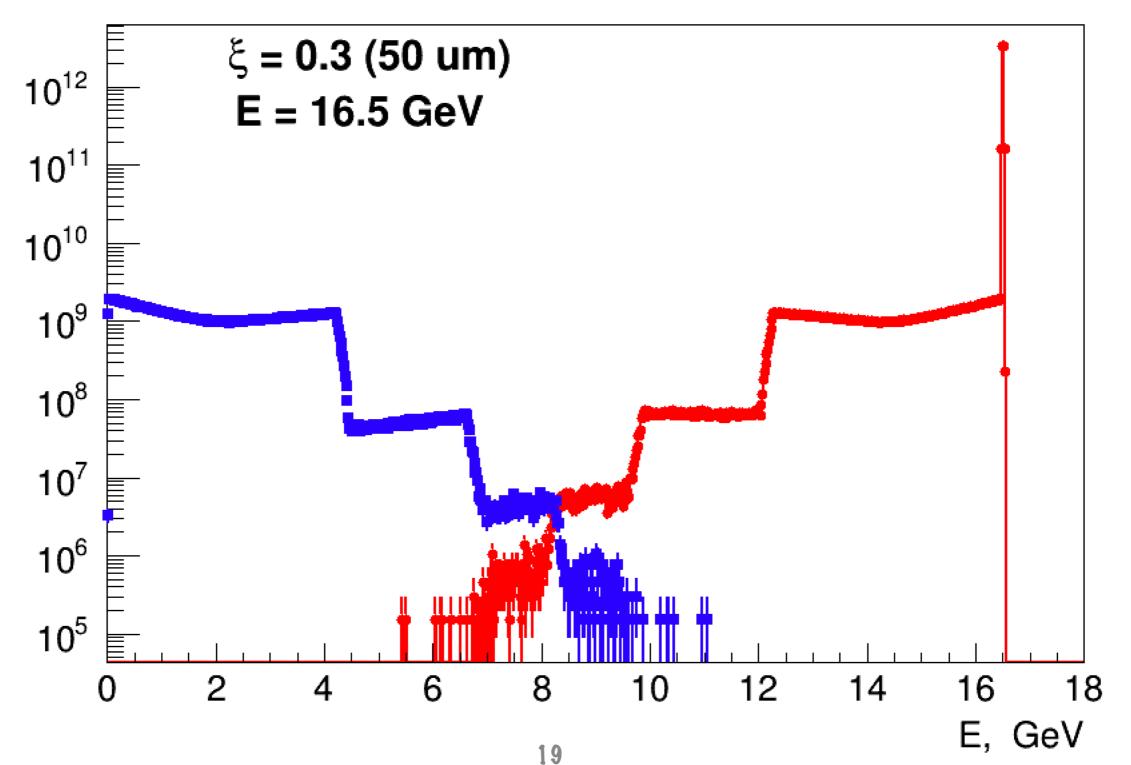




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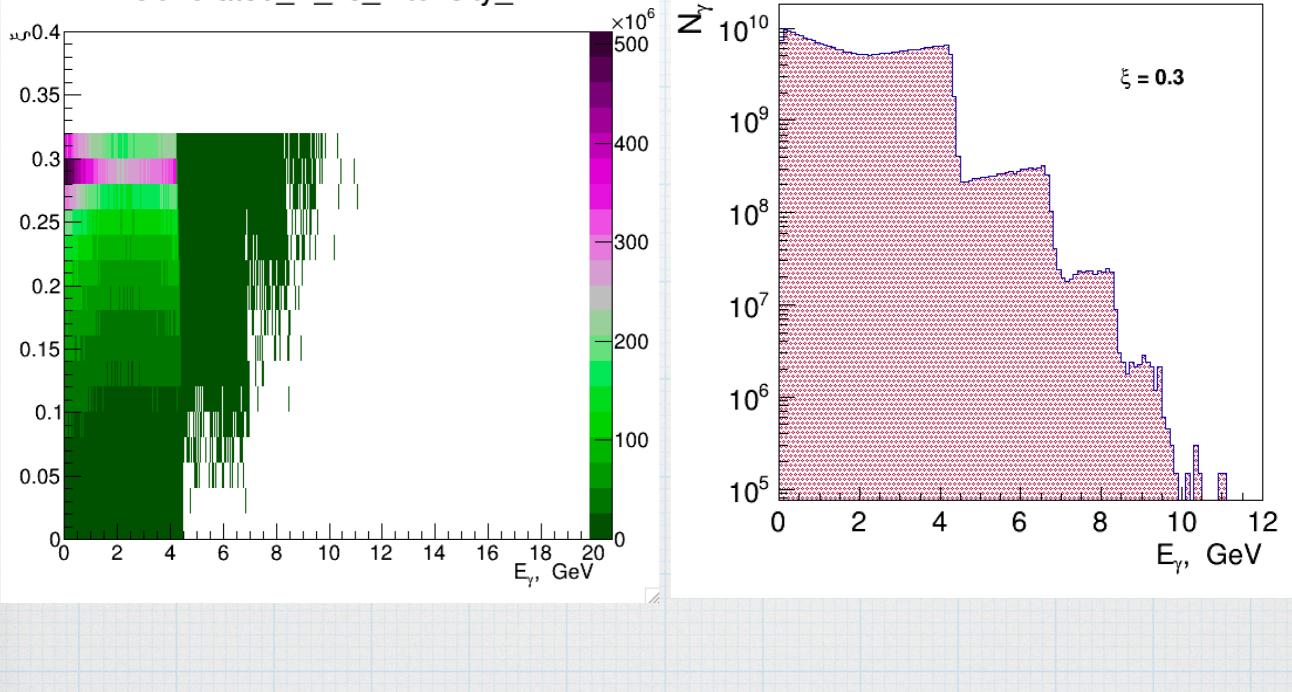
(5% of files have NaN so they are ignored)



& vs photon energy in MC

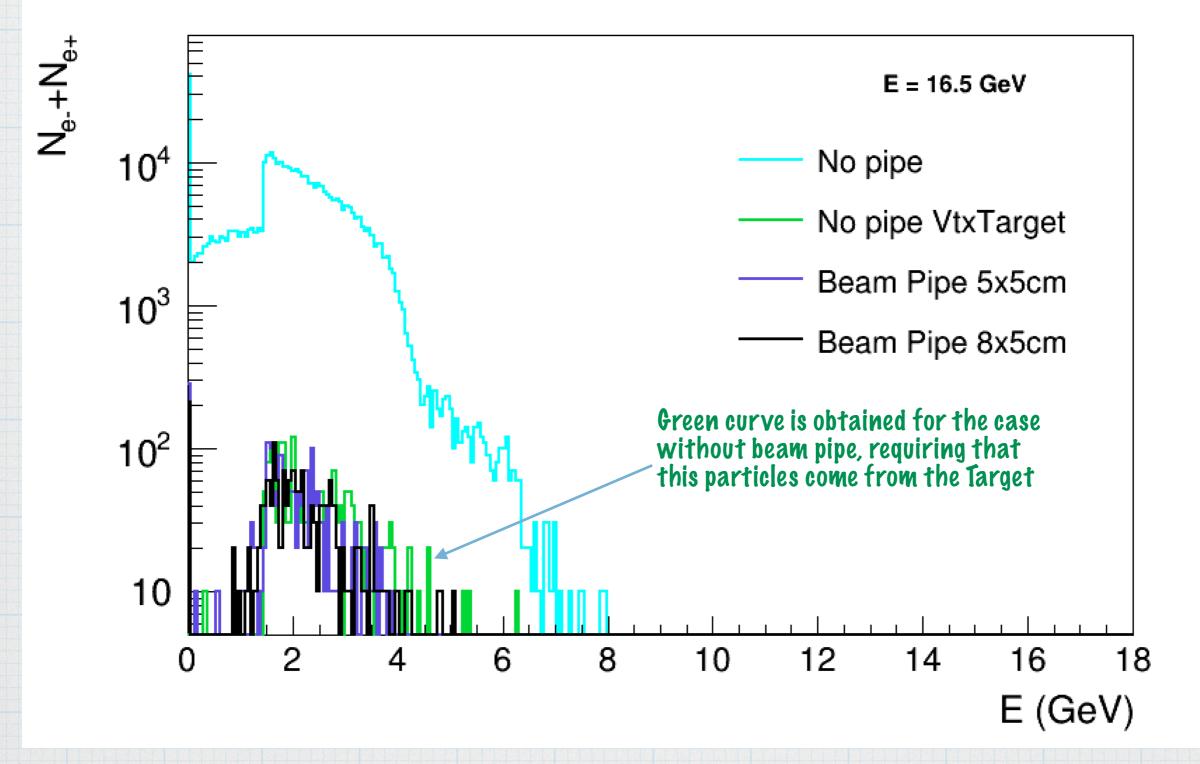
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Generated_E_vs_Intensity_2



20





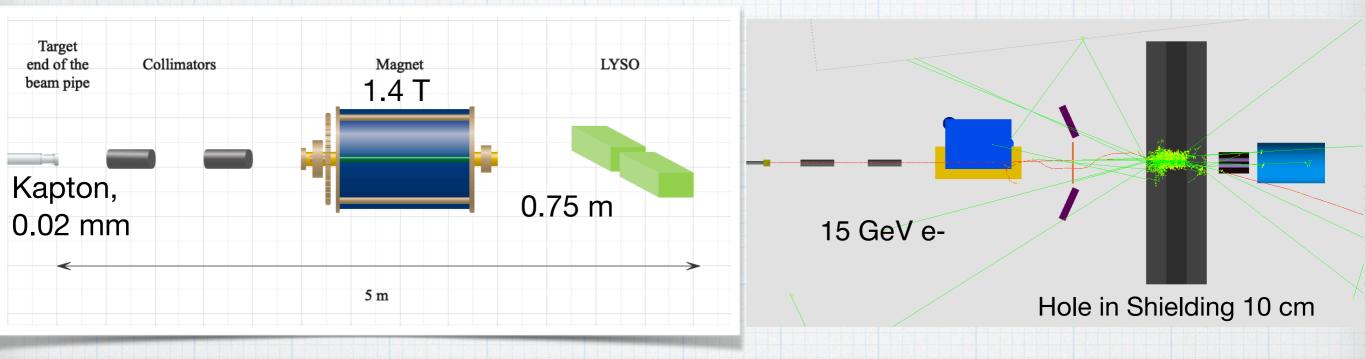
Without beam pipe we measure in Compton detectors a lot e-/e+ pairs that were created in the air. Only 4% e-/e+ are generated in the Target

Setup with the beam pipe

Added round beam pipe of 5cm diameter between the collimators

And beam pipe w/ square Xsection of 5x5cm (8x5cm) from collimators to the LYSO detectors

FDS with LYSO calorimeters



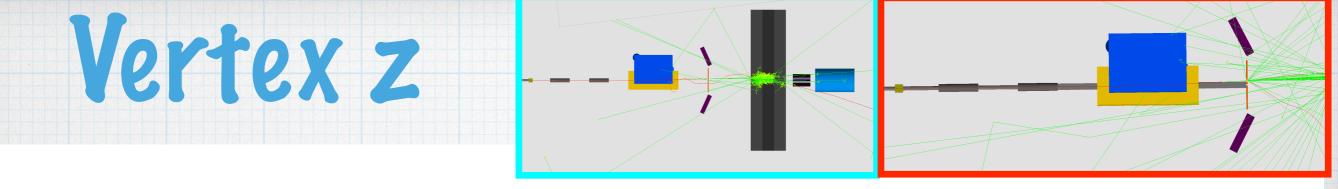
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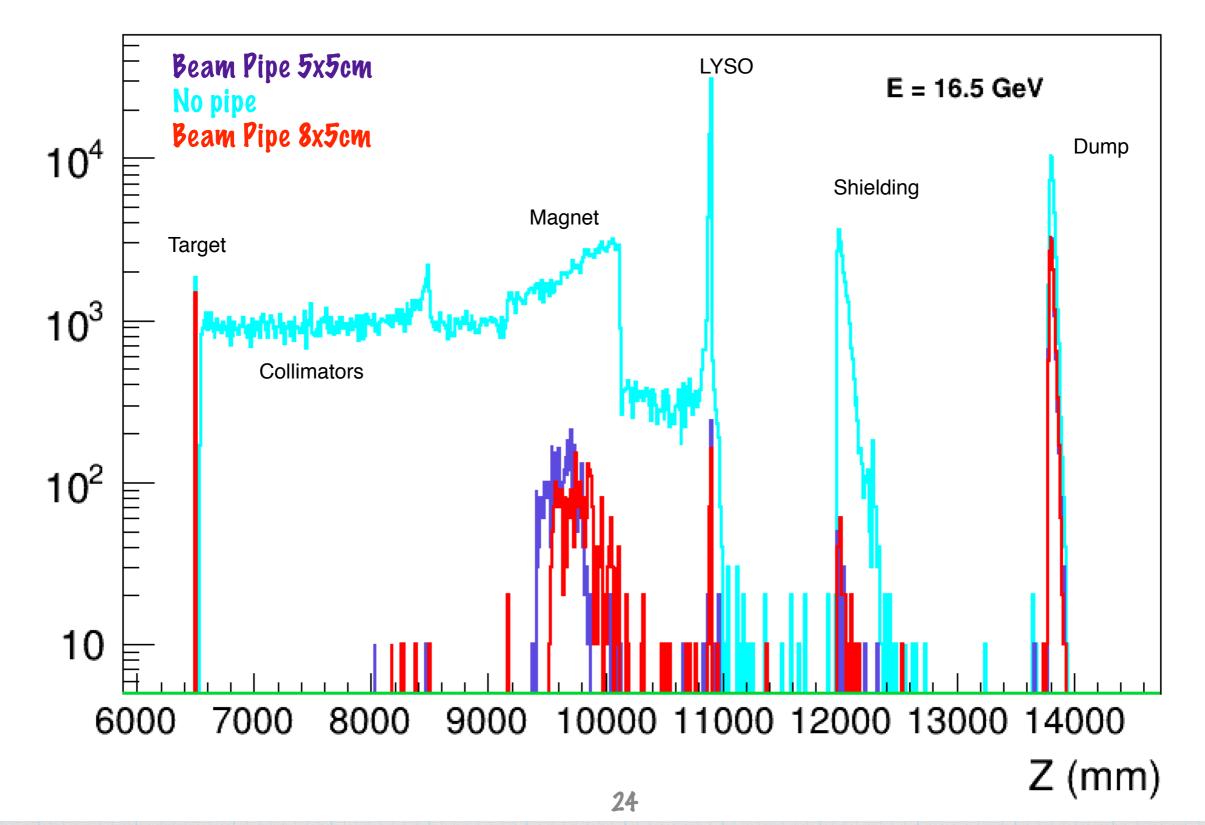
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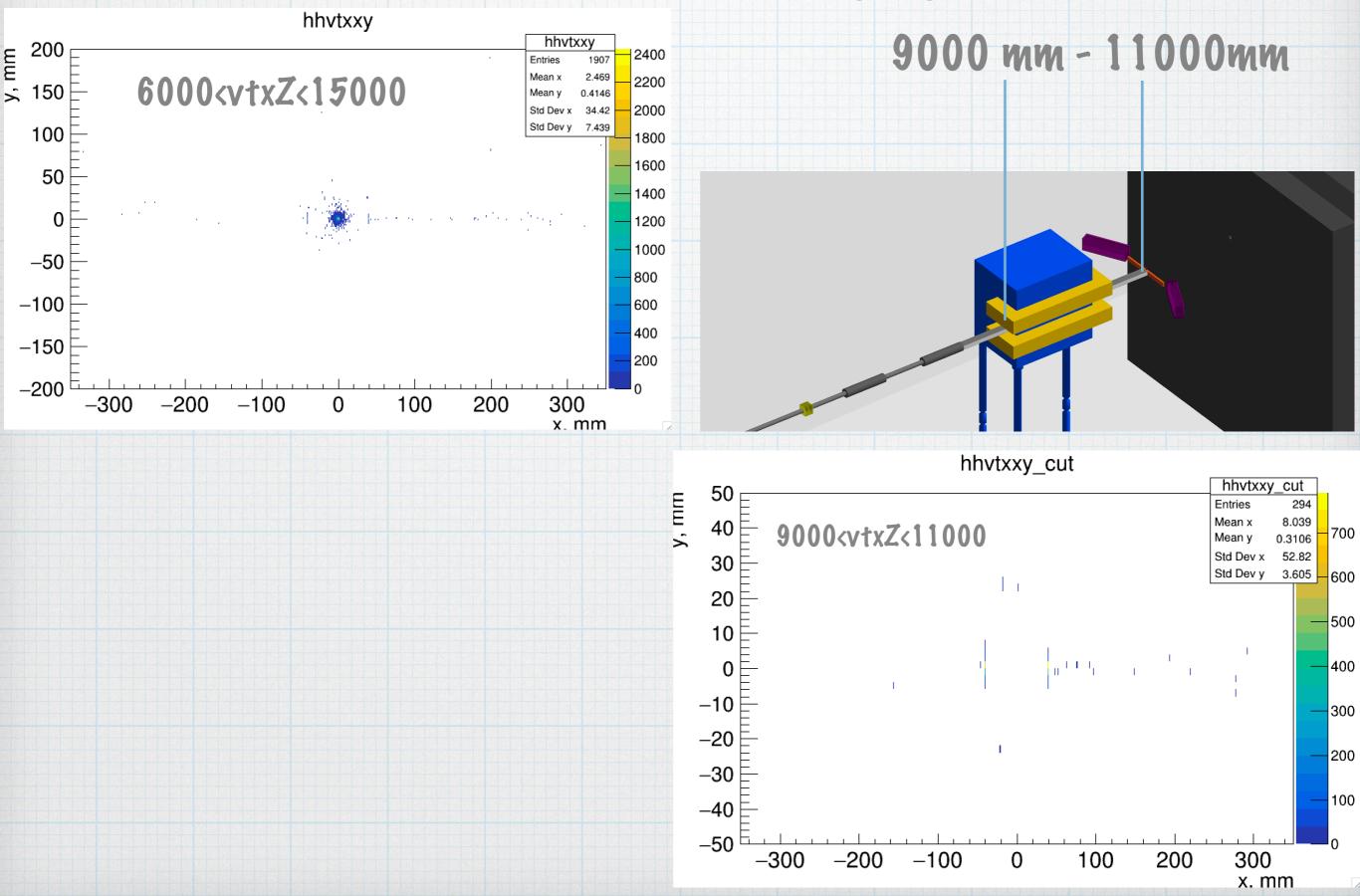
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All studies were performed with 100 BX at the laser intensity xi = 0.3 for 16.5 GeV electron beam

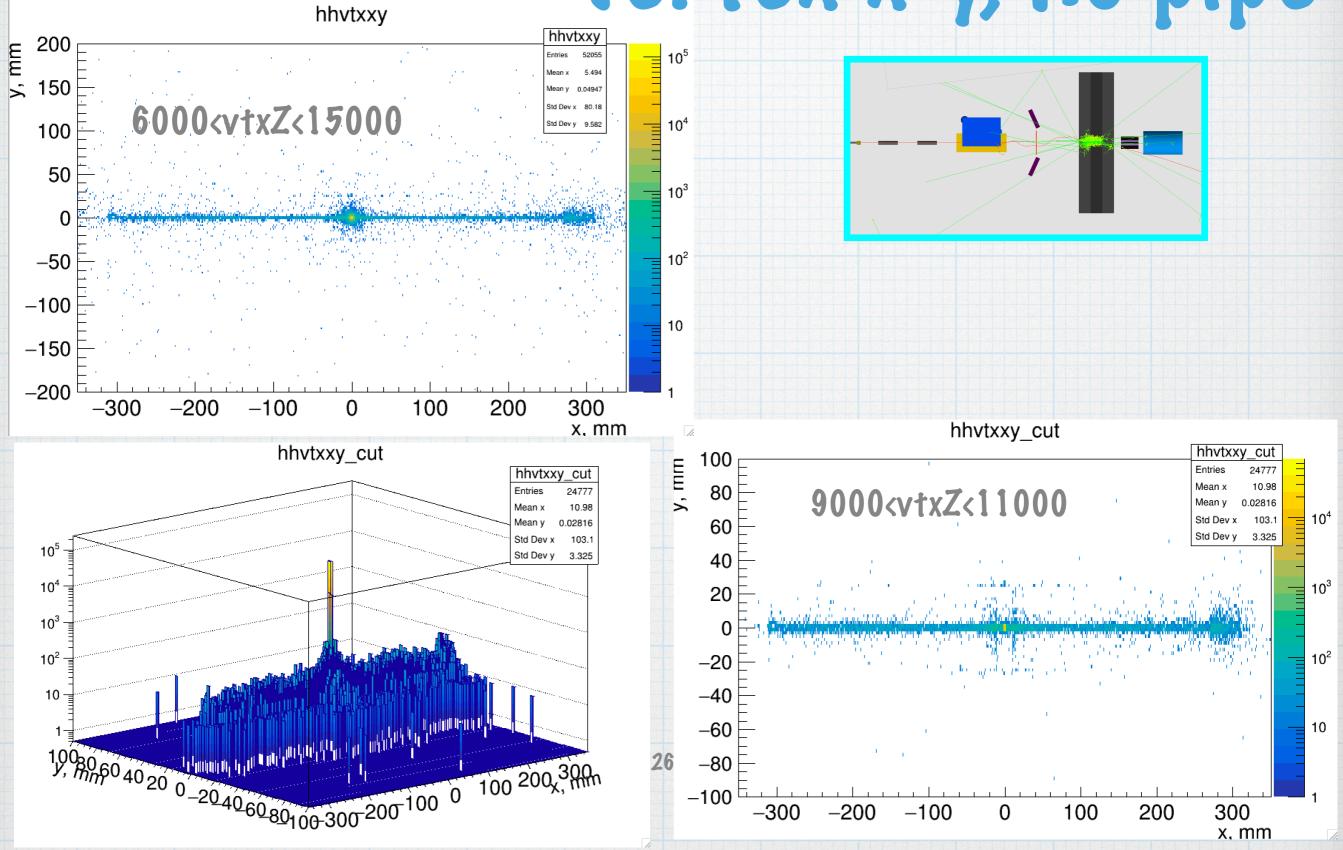




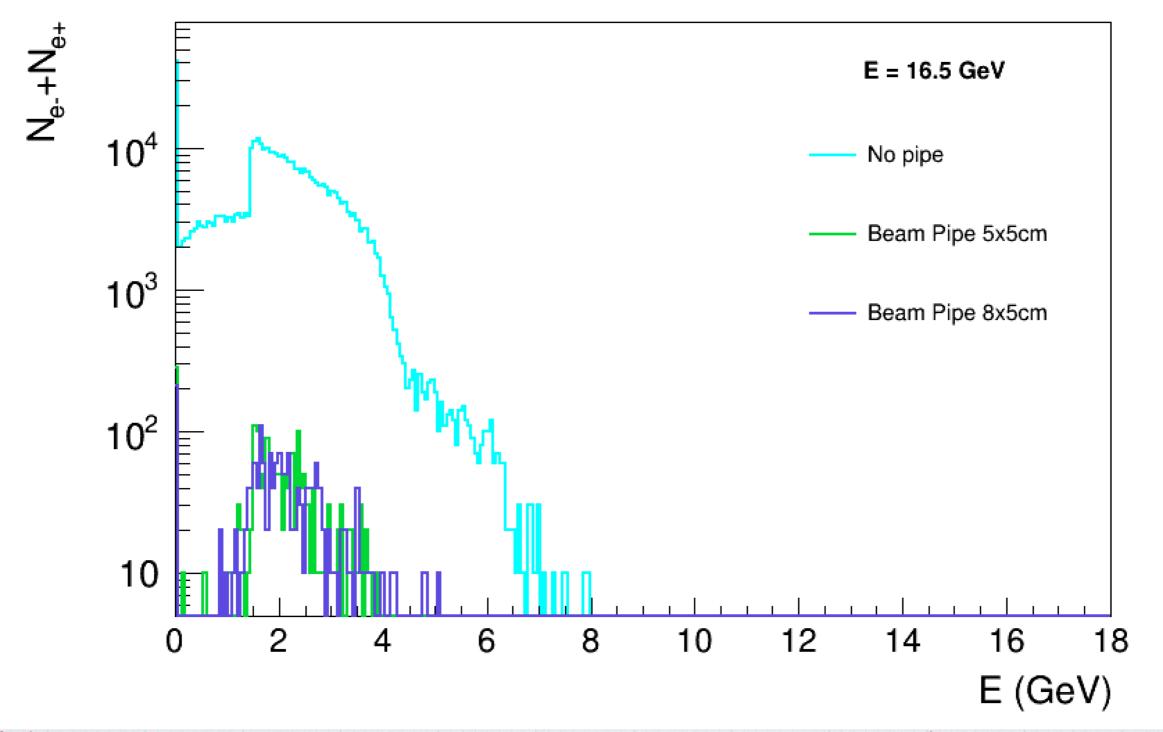
Vertex x-y, beam pipe 8x5



Vertex x-y, no pipe

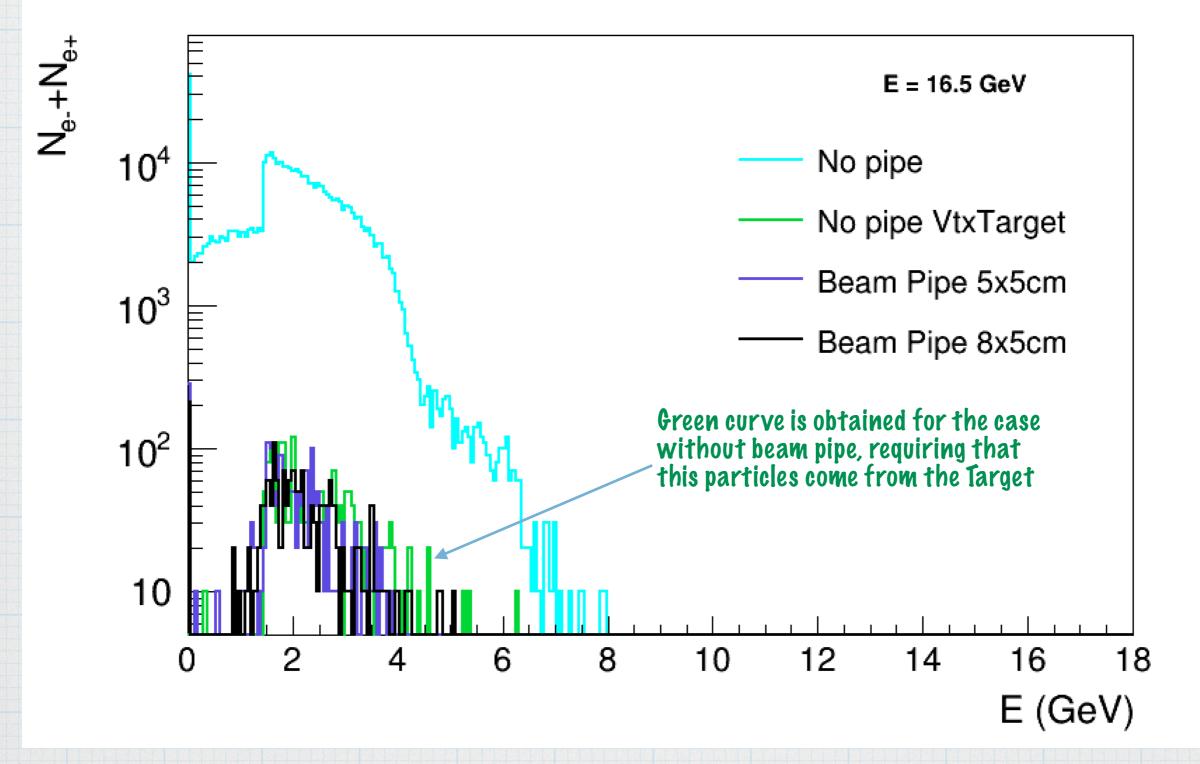






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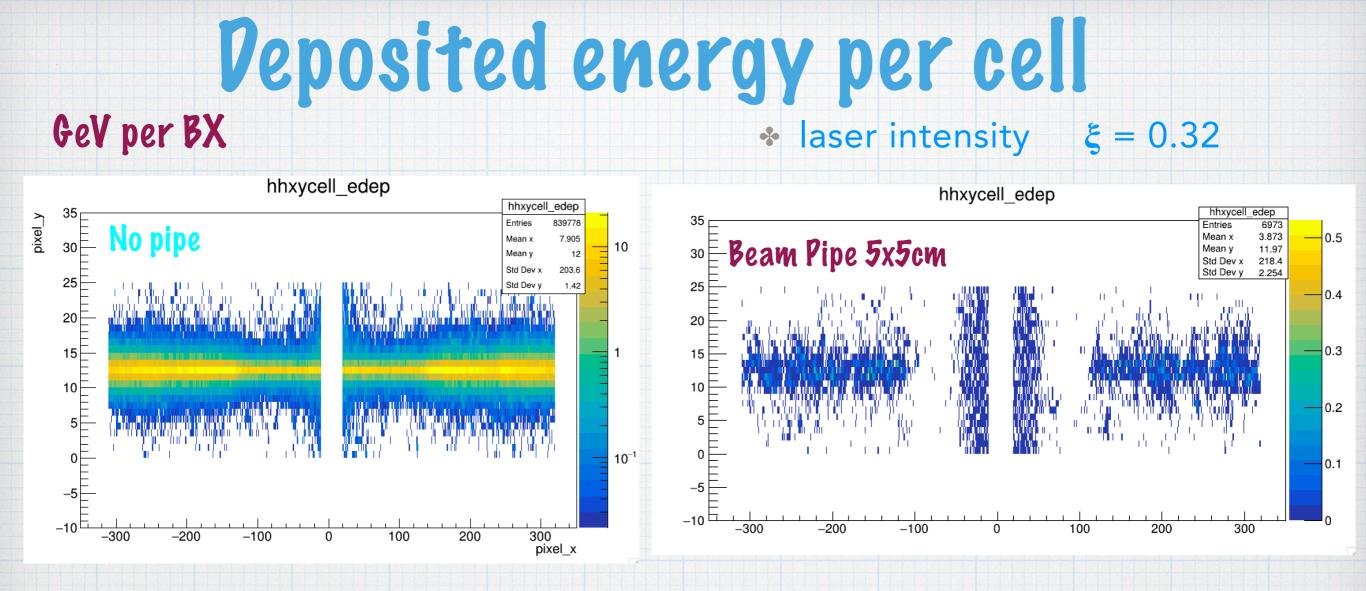




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- * The performance of FDS setup was compared with and without beam pipe from the target to Compton detectors
- * Number of particles per BX hitting LYSO detector is 25 higher without beam pipe
- * Big hole in the Shielding creates substantial background occupancy in LISO detectors.
- * All extra particles are generated in the air. Number of particles generated in the target is identical.
- In the air the vertexes are distributed almost uniformly all the way from the target to the detectors in case of no pipe.
- * As the laser intensity is low (xi =0.3), to reconstruct spectra we need more statistics. Asked Anthony to produce more; he runs now 1000BX



Compton MC2020 r for (xi=0.32), 16.5 GeV electrons. G4: Kapton foil of 20 um as a target, magnet 1.4T and 0.75m distance from magnet to LYSO .

If we take distribution of deposited energy the values around maximum are ~10 GeV.

To convert it to Gy, convert it to J: $^{-1.6e-9J}$ and then divide it to the mass of crystals in kg. Gy= J/kg

The density is 7.1 g/cm3, volume 0.1 * 0.2 * 2 = 0.04 cm3. Mass 7.1 * 0.04 = 0.284g.

Finally, 5.6e-6 Gy per BX.

Assuming 1 Hz collisions rate we get the dose of 10 kGy in LYSO crystal in about 56 years.

Vertex z

