

Positive moments for scattering amplitudes

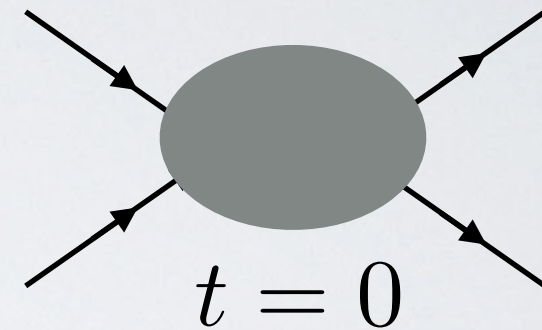
from

Bellazzini, Elias Miro, Rattazzi, Riembau, Riva
2011.00037 [hep-th]

(see also Tolley, Wang, Zhou, 2011.02400 [hep-th],
and probably others)

Positivity bounds : sign constraints on amplitudes

[Adams, Arkani-Hamed, Dubovsky, Nicolis, Rattazzi, 06]



$$\mathcal{A}(s)$$

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Ingredients :

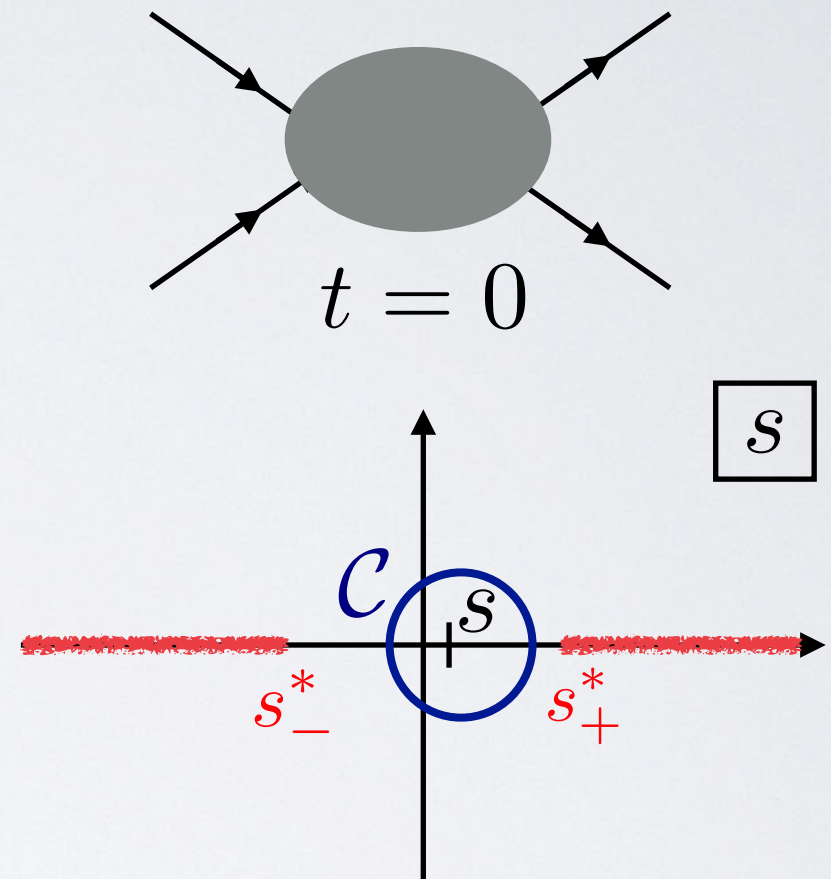
- **analyticity**

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$$\mathcal{A}(s) = \int_{\mathcal{C}} \frac{\mathcal{A}(s')}{2\pi i (s' - s)}$$

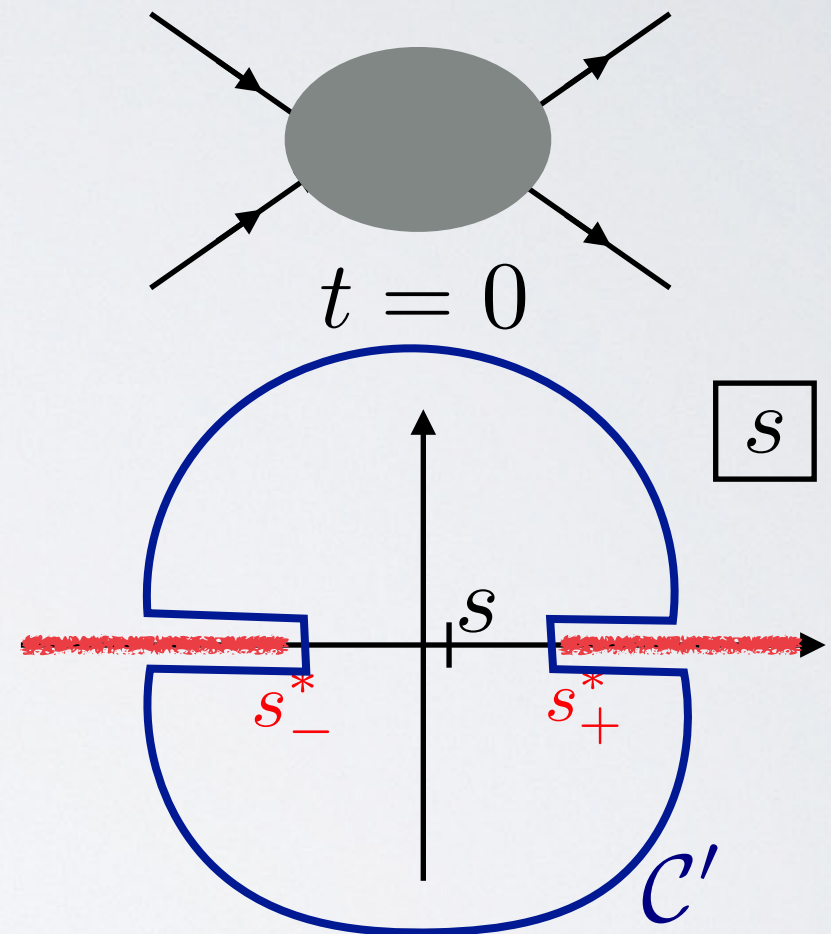


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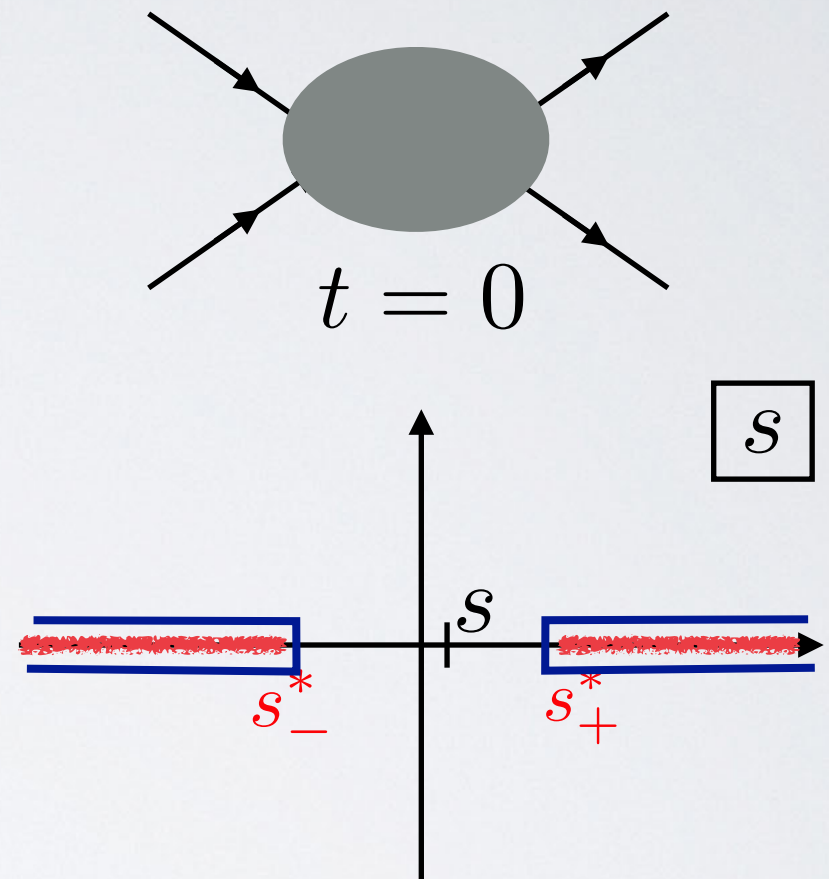
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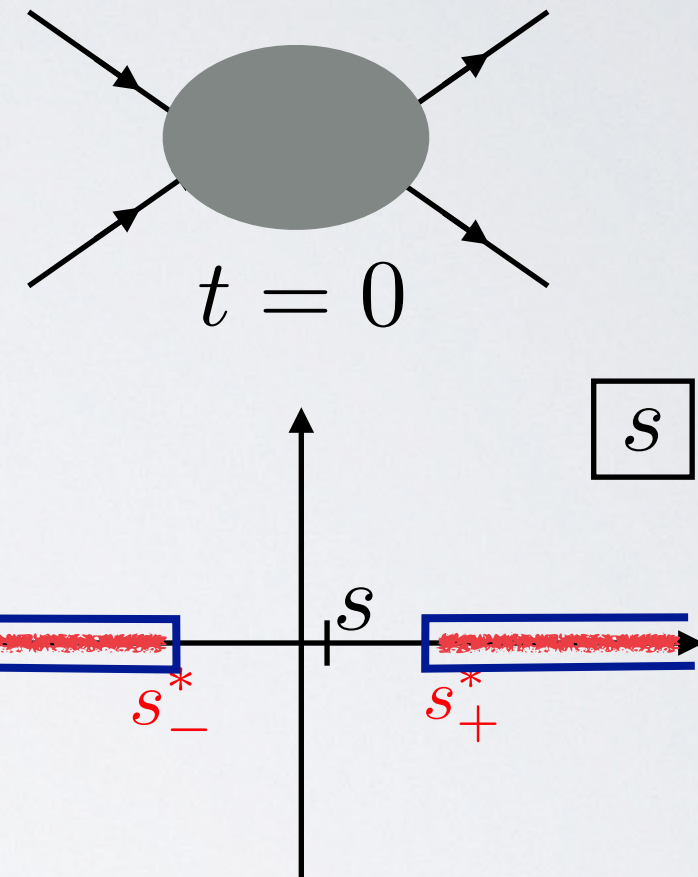
$$\mathcal{A}(s) = \int_{-\infty}^{s_-^*} \frac{\text{Im}\mathcal{A}(s')}{2\pi i(s' - s)} + \int_{s_+^*}^{+\infty} \frac{\text{Im}\mathcal{A}(s')}{2\pi i(s' - s)}$$

Positivity bounds : sign constraints on amplitudes

[Adams, Arkani-Hamed, Dubovsky, Nicolis, Rattazzi, 06]

Ingredients :

- **analyticity**
- **polynomial boundedness**
(Froissart bound)
-
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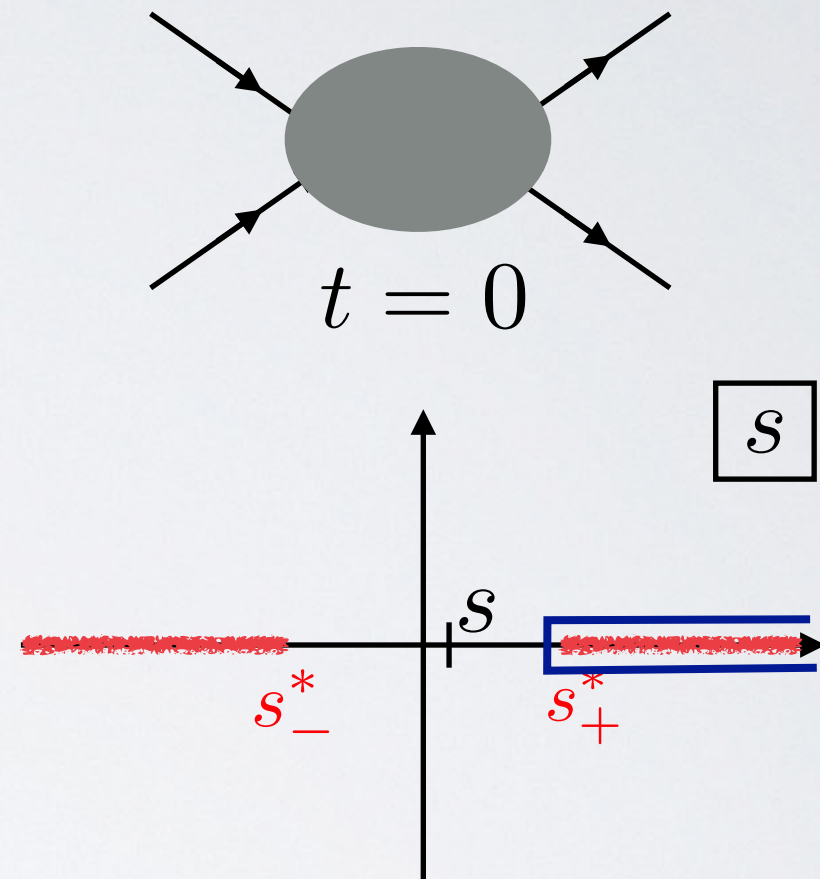
$$\tilde{\mathcal{A}}(s) = \int_{-\infty}^{s^*_-} \frac{\text{Im}\mathcal{A}(s')}{2\pi i(s' - s)^3} + \int_{s^*_+}^{+\infty} \frac{\text{Im}\mathcal{A}(s')}{2\pi i(s' - s)^3}$$

Positivity bounds : sign constraints on amplitudes

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Ingredients :

- **analyticity**
- **polynomial boundedness**
(Froissart bound)
- **crossing symmetry**
-



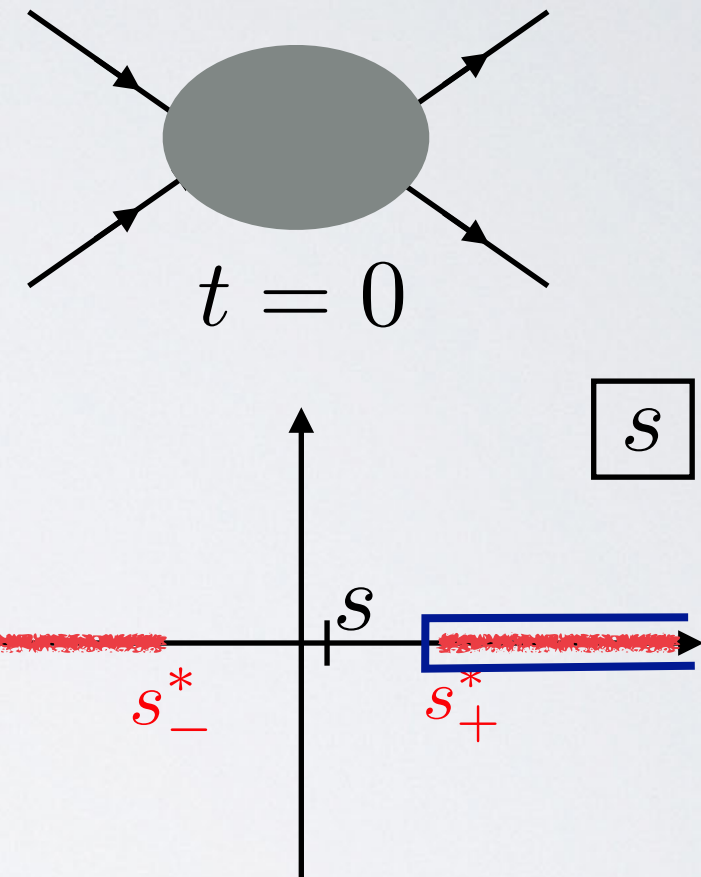
$$\tilde{\mathcal{A}}(s) = \int_{4m^2 - s_-^*}^{+\infty} \frac{\text{Im}\mathcal{A}(s')}{2\pi i(s' - 4m^2 + s)^3} + \int_{s_+^*}^{+\infty} \frac{\text{Im}\mathcal{A}(s')}{2\pi i(s' - s)^3}$$

Positivity bounds : sign constraints on amplitudes

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Ingredients :

- **analyticity**
- **polynomial boundedness**
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- **unitarity**



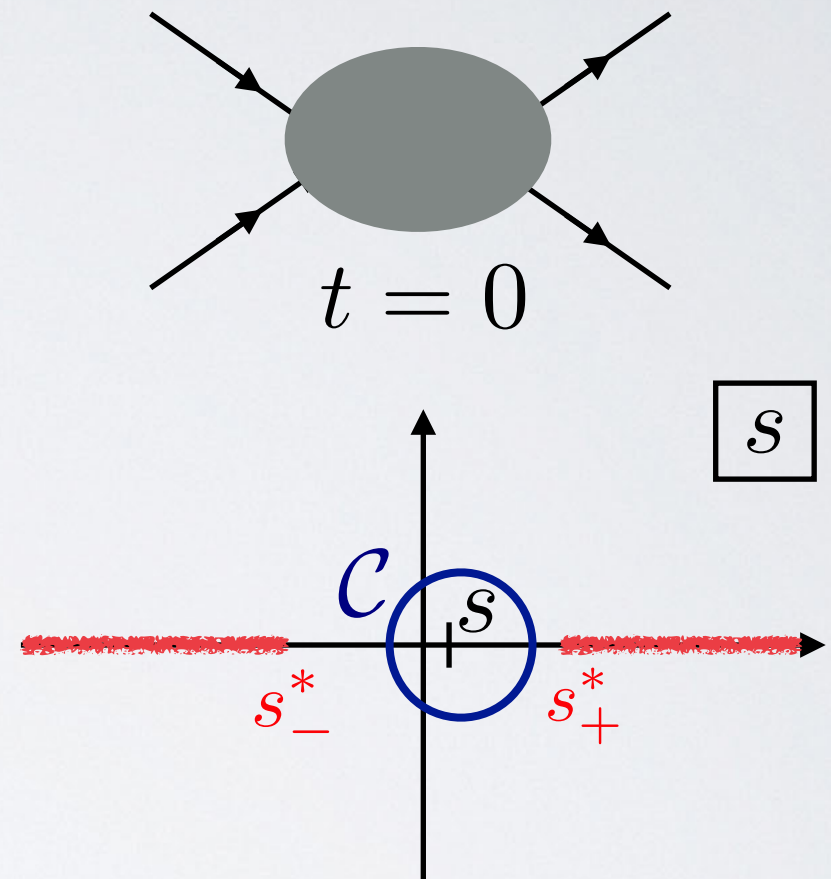
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Positivity bounds : sign constraints on amplitudes

[Adams, Arkani-Hamed, Dubovsky, Nicolis, Rattazzi, 06]

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(Froissart bound)
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$$\tilde{\mathcal{A}}(s) = \int_C \frac{\mathcal{A}(s')}{2\pi i (s' - s)^3} > 0$$

At small s : computed in the EFT,
bound on EFT coefficients

[see papers of the authors, Remmen/Rodd, de Rham/Tolley/Zhou, Zhang/Wang...]

How to go further (in simple scalar models)?

$$\tilde{A}(s) = \int_{4m^2 - s_-^*}^{+\infty} \frac{\text{Im}\mathcal{A}(s')}{2\pi i(s' - 4m^2 + s)^3} + \int_{s_+^*}^{+\infty} \frac{\text{Im}\mathcal{A}(s')}{2\pi i(s' - s)^3} > 0$$

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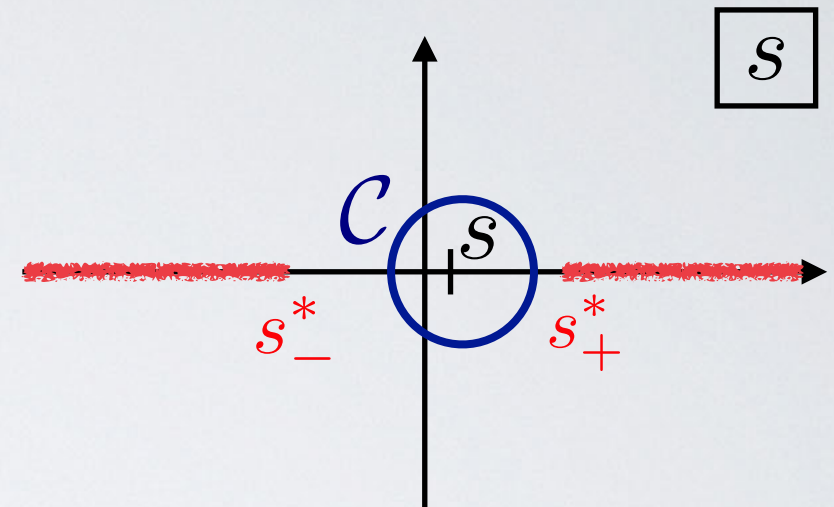
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Integrands are positive :

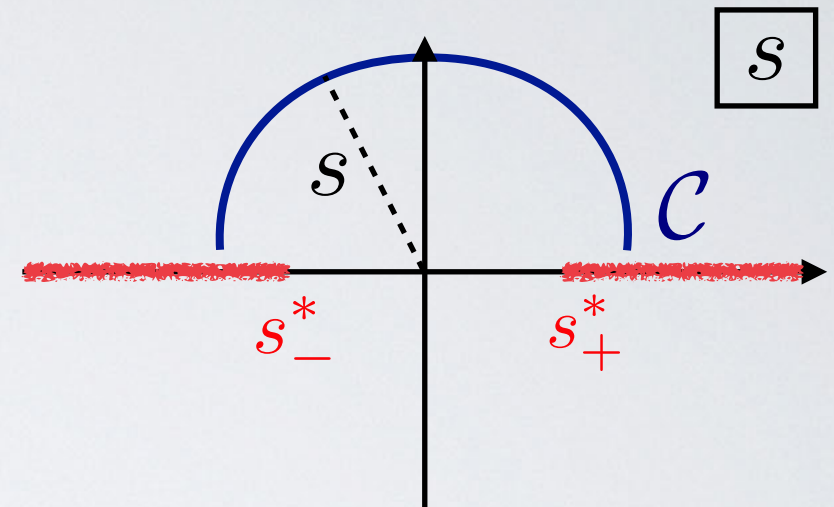
$$\int_{s_+^*}^{+\infty} \frac{\text{Im}\mathcal{A}(s')}{2\pi i(s' - s)^3} f(s') > 0$$

New quantities : **arcs**



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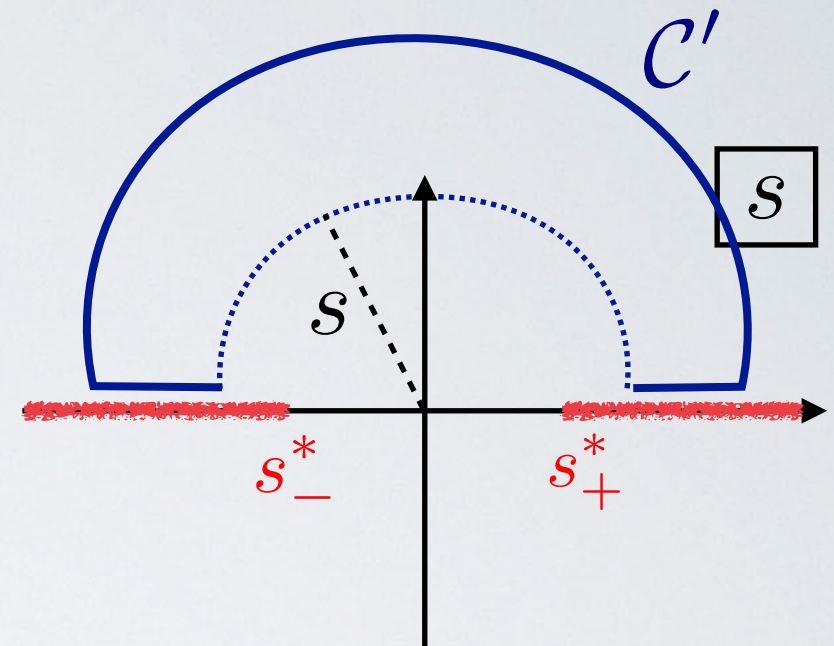
$$a_n(s) \equiv \frac{1}{\pi i} \int_{\mathcal{C}(s)} ds' \frac{\mathcal{A}(s')}{s'^{2n+3}}$$



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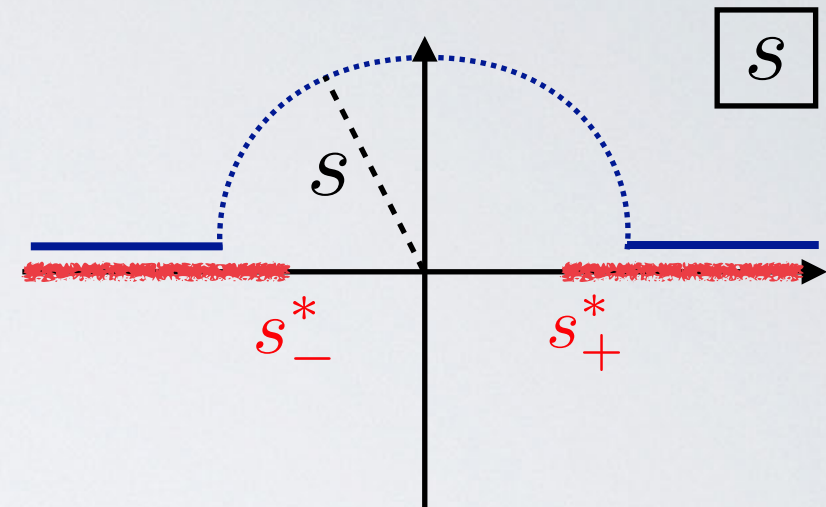
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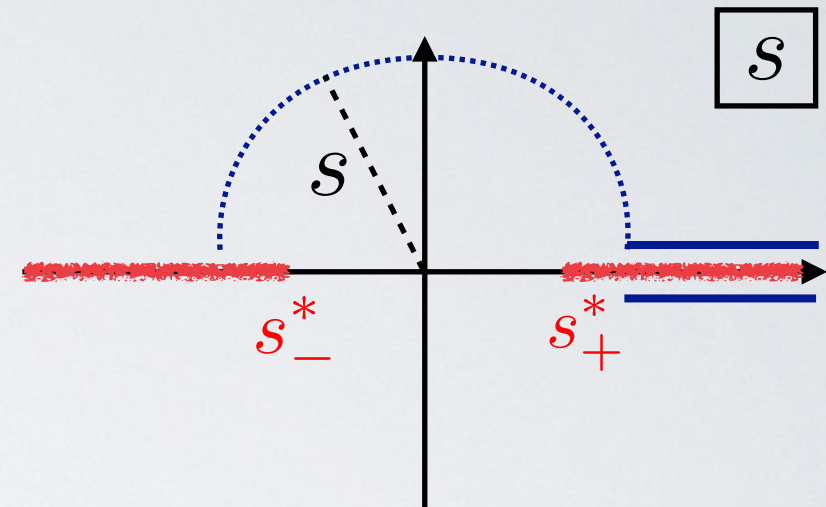


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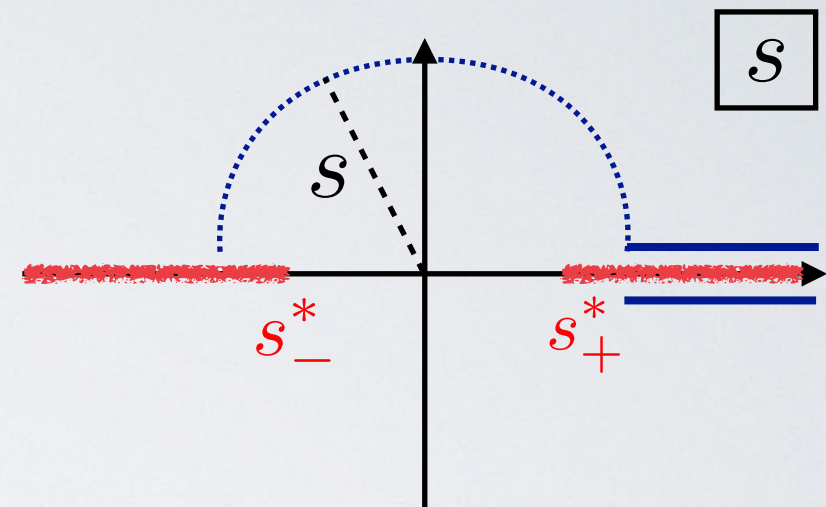


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Positive, and **computable at small s in the EFT**, for the same reasons at before

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$$d\mu(x) \equiv \frac{dx}{\pi} \text{Im}\mathcal{A}\left(\frac{s}{\sqrt{x}}\right)$$

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Largest set of inequalities on the moments

$$(-1)^k \Delta^k a_n > 0 \quad \text{with} \quad \Delta a_n \equiv s^2 a_{n-1} - a_n$$

What about the **projection on low-energy arcs?**

$$\int_0^1 d\mu(x) x^n (1 - x^k) > 0$$

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$$(H_N^l)_{ij} = a_{i+j+l}$$

$$i, j \leq \lfloor (N-l)/2 \rfloor$$

$$H_N^0 \succ 0, H_N^1 \succ 0$$

$$H_{N-1}^0 - s^2 H_N^1 \succ 0, H_{N-1}^1 - s^2 H_N^2 \succ 0$$

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Non-linear, s-dependent bounds (more and more stringent when s grows)

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At tree-level

$$a_n(s) \equiv \frac{1}{\pi i} \int_{\mathcal{C}(s)} ds' \frac{\mathcal{A}(s')}{s'^{2n+3}}$$

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$$\mathcal{A}(s) = c_0 + c_2 s^2 + c_4 s^4 \dots$$

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s-independent
bounds

$$\implies a_n = c_{2n+2}$$

s-dependent bounds :

upper bounds on the cutoff

$$\text{Ex : } s_{\text{max}}^2 < \frac{c_4}{c_6}$$

bounds on « soft » theories

$$\text{Ex : } c_2 > s^2 c_4$$

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Tree-level UV completion

$$\text{Im}\mathcal{A}(s) = \frac{\pi}{2} \sum_{i=1}^P g_k^2 M_k^2 \delta(s - M_k^2) \implies d\mu(x) \text{ is a sum of deltas at } x = \frac{s}{M_k^2}$$

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With enough arcs, one **reconstructs the full UV theory** :

- number of particles from $\det(H_{2P}^0) = 0$
- spectrum from $\det(H_{2P-2}^0 - s^2 H_{2P-1}^1) = 0$
- couplings from $a_n = \sum \frac{g_k^2}{M_k^{4n+4}}$

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Loops in the EFT

$$a_n(s) \equiv \frac{1}{\pi i} \int_{\mathcal{C}(s)} ds' \frac{\mathcal{A}(s')}{s'^{2n+3}}$$

$$a_0 = c_2 + \frac{s^2}{2} \beta_4 + \frac{s^3}{3} \beta_5 + \frac{s^4}{4} \left(\beta_6 + \frac{\beta'_6}{2} (4 \log s - 1) \right) + \dots,$$

$$a_1 = c_4(s) + s \beta_5 + \frac{s^2}{2} \left(\beta_6 + \beta'_6 (2 \log s - 1) \right) + \dots,$$

$$a_2 = -\frac{\beta_4}{2s^2} - \frac{\beta_5}{s} + c_6(s) + \dots, \quad (29)$$

$$\mathcal{A}(s) = c_2 s^2 + s^4 [c_4 + \beta_4 \log(-is)] - i\pi s^5 \beta_5 / 2$$

$$+ s^6 [c_6 + \beta_6 \log(-is) + \beta'_6 \log^2(-is)] + O(s^7)$$

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New s-dependence

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$$a_0 = c_2 + \frac{s^2}{2} \beta_4 + \frac{s^3}{3} \beta_5 + \frac{s^4}{4} \left(\beta_6 + \frac{\beta'_6}{2} (4 \log s - 1) \right) + \dots,$$

$$a_1 = c_4(s) + s \beta_5 + \frac{s^2}{2} \left(\beta_6 + \beta'_6 (2 \log s - 1) \right) + \dots,$$

$$a_2 = -\frac{\beta_4}{2s^2} - \frac{\beta_5}{s} + c_6(s) + \dots, \quad (29)$$

$$\mathcal{A}(s) = c_2 s^2 + s^4 [c_4 + \beta_4 \log(-is)] - i\pi s^5 \beta_5 / 2$$

$$+ s^6 [c_6 + \beta_6 \log(-is) + \beta'_6 \log^2(-is)] + O(s^7)$$

Running effects dominate the IR

New s-dependence

$$a_0 \rightarrow c_2, \quad a_1 \rightarrow \beta_4 \log s, \quad a_{n \geq 2} \rightarrow -\frac{\beta_4}{(2n-2)s^{2n-2}}$$

and enter the bounds

$$c_4(s) - c_6(s)s^2 > -\frac{\beta_4}{2}$$

$$(H_N^l)_{ij} = a_{i+j+l}$$

$$i, j \leq \lfloor (N-l)/2 \rfloor$$

$$H_N^0 \succ 0, H_N^1 \succ 0$$

$$H_{N-1}^0 - s^2 H_N^1 \succ 0, H_{N-1}^1 - s^2 H_N^2 \succ 0$$

Loops in the EFT

$$a_n(s) \equiv \frac{1}{\pi i} \int_{\mathcal{C}(s)} ds' \frac{\mathcal{A}(s')}{s'^{2n+3}}$$

$$a_0 = c_2 + \frac{s^2}{2} \beta_4 + \frac{s^3}{3} \beta_5 + \frac{s^4}{4} \left(\beta_6 + \frac{\beta'_6}{2} (4 \log s - 1) \right) + \dots,$$

$$a_1 = c_4(s) + s \beta_5 + \frac{s^2}{2} \left(\beta_6 + \beta'_6 (2 \log s - 1) \right) + \dots,$$

$$a_2 = -\frac{\beta_4}{2s^2} - \frac{\beta_5}{s} + c_6(s) + \dots, \quad (29)$$

$$\mathcal{A}(s) = c_2 s^2 + s^4 [c_4 + \beta_4 \log(-is)] - i\pi s^5 \beta_5 / 2$$

$$+ s^6 [c_6 + \beta_6 \log(-is) + \beta'_6 \log^2(-is)] + O(s^7)$$

Running effects dominate the IR

New s-dependence

Strong corrections from soft theories?
Still forbidden

$$a_0 \rightarrow c_2, \quad a_1 \rightarrow \beta_4 \log s, \quad a_{n \geq 2} \rightarrow -\frac{\beta_4}{(2n-2)s^{2n-2}}$$

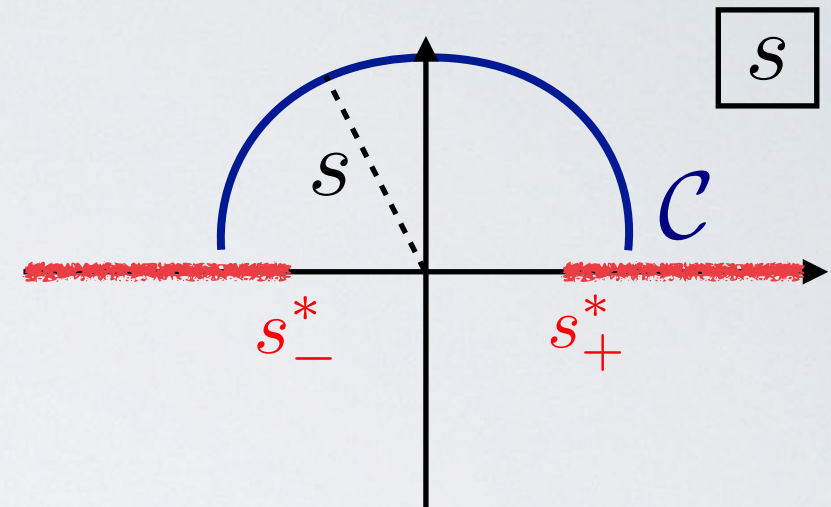
and enter the bounds

$$c_4(s) - c_6(s)s^2 > -\frac{\beta_4}{2}$$

TAKE AWAY

New quantities at finite s : **arcs**

$$a_n(s) \equiv \frac{1}{\pi i} \int_{\mathcal{C}(s)} ds' \frac{\mathcal{A}(s')}{s'^{2n+3}}$$



Most general bounds on arcs from the theory of moments :
s-dependence, non-linearity

At tree-level : bounds on the Wilson coefficients and on the EFT cutoff. UV info for weakly-coupled UV theories

At loop-level : quantum effects enter and sometimes dominate the bounds, the naive ones on Wilson coefficients can be violated. All bounds are s -dependent, they still rule out soft theories.