

Local Unitarity: A new approach to perturbative computations in Quantum Field Theories

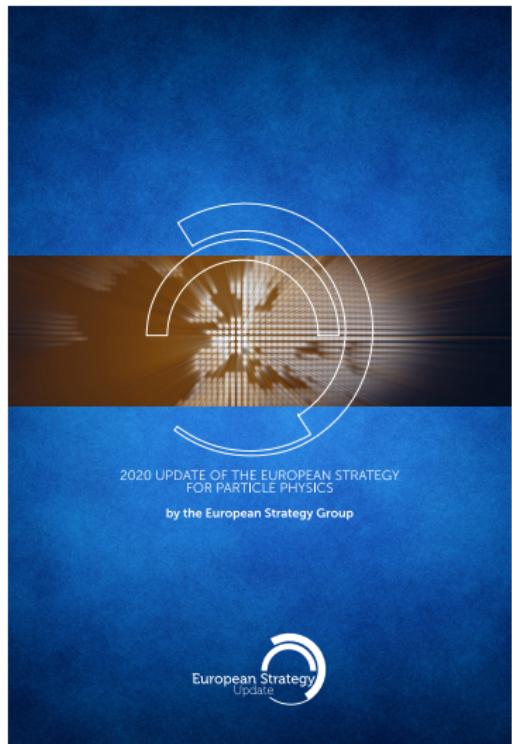
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July 29, 2021

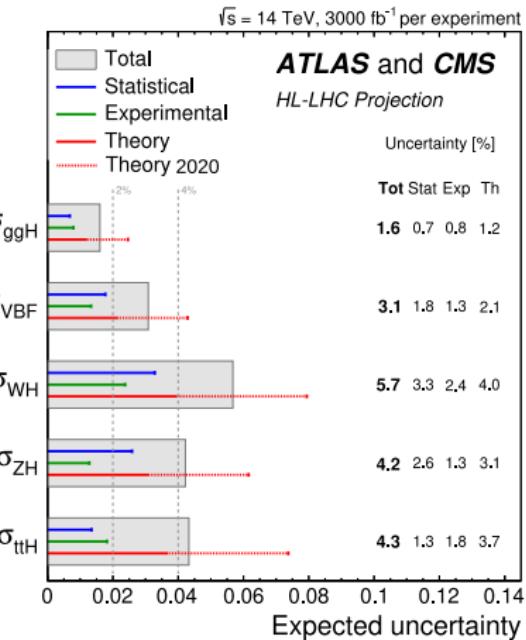
The European Strategy for Particle Physics

“ The vision is to prepare a Higgs factory, followed by a future hadron collider with sensitivity to energy scales an order of magnitude higher than those of the LHC ”



Keeping up with experiment

- High Luminosity LHC is expected to drastically reduce statistical uncertainty
- Theory uncertainty will be the **dominating contribution**
- NLO cross sections are available with the push of a button
- Why no automation at NNLO?



[CERN-LPCC-2018-04]

Cross sections

A cross-section can be represented as a sum of interference diagrams

$$\sigma = -\text{---} \circ \text{---} + -\text{---} \circ \text{---} \text{---} \circ \text{---} + -\text{---} \circ \text{---} \text{---} \circ \text{---} \text{---} \circ \text{---} + \dots$$

Cross sections

A cross-section can be represented as a sum of interference diagrams

$$\sigma = \text{---} \circlearrowleft + \text{---} \circlearrowleft \text{---} + \text{---} \circlearrowleft \text{---} + \text{---} \circlearrowright + \text{---} \circlearrowright \text{---} + \dots$$

Each interference can be represented as integrals of amplitudes:

$$\text{---} \circlearrowleft = \int d^3 \vec{p} \delta(p^0 - |\vec{p}'| - |\vec{p} - \vec{p}'|) \cdot \text{---} \quad \cdot \int d^4 k \text{---} \circlearrowright$$

Both the phase space integral and loop integral are divergent!

Challenges

A Feynman diagram showing a loop with four external lines. The top-left line is labeled p , the top-right p' , the bottom-left $p - k$, and the bottom-right $p - p'$. The loop itself is labeled k . Arrows on the lines indicate the direction of flow.

$$= \int d^4k \frac{1}{k^2(k-p')^2(p-k)^2}$$

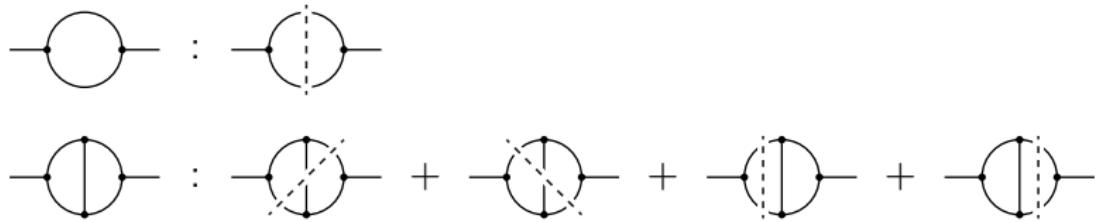
- This integral diverges when $p'^2 = 0$
- Workarounds to integrate require changing the space-time dimension or introducing a counterterm:

$$\int \underbrace{I - I_{CT}}_{\text{finite}} + \underbrace{\int I_{CT}}_{\text{easy}}$$

- Some integrals are elliptic functions that cannot be evaluated

Real-virtual cancellations

- Do not consider the amplitude as the fundamental component anymore, but elements of the cross section with the same *supergraph*: [Capatti, Hirschi, Kermanschah, Pelloni, Ruijl, JHEP 2020]



- Each of these collections is finite (KLN theorem)
- If we can make it locally finite, we can Monte Carlo integrate them

Local Unitarity

| Wishes | Local Unitarity |
|---|-----------------|
| <ul style="list-style-type: none">Differential cross sectionsNo restrictions on mass scalesNo reduction to master integralsNo complicated function evaluationsNo IR counterterms for loop degrees of freedomNo IR counterterms for real degrees of freedomFully automated renormalizationIntegration in lower dimensionsMethod generic for any process to any order | |

Local Unitarity

| Wishes | Local Unitarity |
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| Differential cross sections | ✓ |
| No restrictions on mass scales | ✓ |
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| No complicated function evaluations | ✓ |
| No IR counterterms for loop degrees of freedom | ✓ |
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| Fully automated renormalization | ✓ |
| Integration in lower dimensions | ✓ |
| Method generic for any process to any order | ✓ |

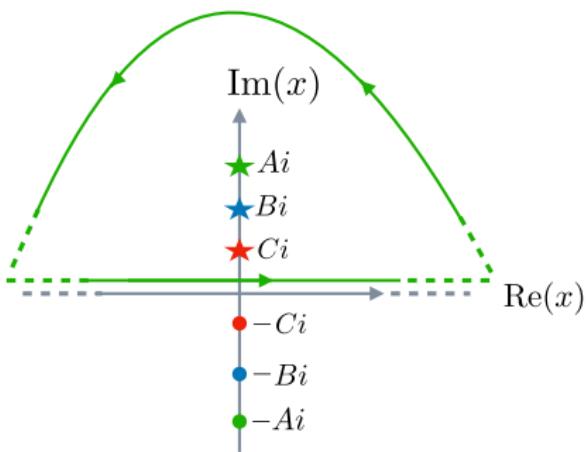
Amplitudes

- First we study the singularities of amplitudes
- Our treatment of amplitudes will guide the extension to cross sections
- Loop-Tree Duality (LTD): integrate out the loop energies analytically

Residue theorem

$$I = \int dx F(x)$$

$$F(x) = \frac{1}{(x - Ai)(x + Ai)} + \frac{1}{(x - Bi)(x + Bi)} + \frac{1}{(x - Ci)(x + Ci)}$$



Cauchy: $R(x^*) \equiv \text{Res}(F, x = x^*)$

$$I = (-2\pi i) [R(Ai) + R(Bi) + R(Ci)]$$

Inverse propagator notation

$$(k + p_i)^2 - m_i^2 + i\epsilon = (k^0 + p_i^0)^2 - (\vec{k} + \vec{p}_i)^2 - m_i^2 + i\epsilon$$

Inverse propagator notation

$$(k + p_i)^2 - m_i^2 + i\epsilon = (k^0 + p_i^0)^2 - (\vec{k} + \vec{p}_i)^2 - m_i^2 + i\epsilon$$

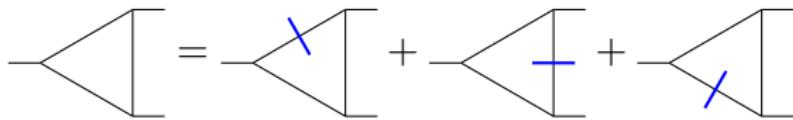
Factorizing:

$$\left(k^0 + p_i^0 + \sqrt{(\vec{k} + \vec{p}_i)^2 + m_i^2 - i\epsilon} \right) \left(k^0 + p_i^0 - \sqrt{(\vec{k} + \vec{p}_i)^2 + m_i^2 - i\epsilon} \right)$$

We call the spatial and mass part Δ :

$$(k^0 + p_i^0 + \Delta_i) (k^0 + p_i^0 - \Delta_i)$$

Residue theorem for propagators



We take the positive energy solution:

$$\delta(k^0 + p_i^0 - \Delta_i) \rightarrow k^0 = \Delta_i - p_i^0$$

Effect of δ on propagator i :

$$(k^0 + p_i^0 + \Delta_i) \rightarrow 2\Delta_i$$

Effect of δ on other propagators j :

$$(\Delta_i - p_i^0 + p_j^0 + \Delta_j) (\Delta_i - p_i^0 + p_j^0 - \Delta_j) = E_{ij} H_{ij}$$

Loop Tree Duality (LTD)

- LTD yields a sum of cut diagrams without loops [Catani, Gleisberg, Krauss, Rodrigo, Winter, JHEP 2009]
- Expression at one-loop:

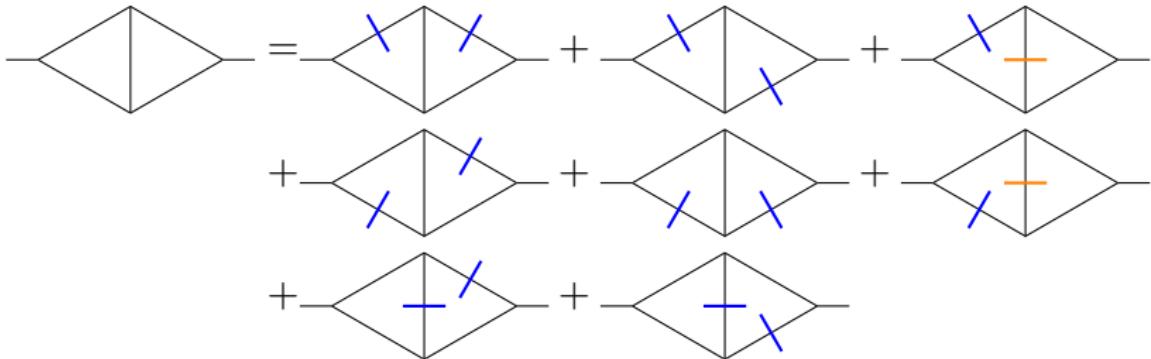
$$I = - \int d^4 k \sum_i^N \tilde{\delta}(q_i^2) \prod_{\substack{j=1 \\ j \neq i}}^N \frac{1}{q_j^2 - i0\eta(q_j - q_i)}$$

- Alternatively:

$$I = - \int d^3 k \sum_i^N \frac{1}{2\Delta_i} \prod_{\substack{j=1 \\ j \neq i}}^N \frac{1}{(\Delta_i - p_i^0 + p_j^0 + \Delta_j)(\Delta_i - p_i^0 + p_j^0 - \Delta_j)}$$

Multi-Loop-Tree Duality

- Integrate the energies from all loop variables one by one [Capatti, Hirschi, Kermanschah, Ruijl PRL 2019; Bierenbaum, Catani, Draggiotis, Rodrigo JHEP 2010]
- Effect: cut all propagator combinations that leave no loops
- Due to the iterative procedure, sometimes the negative energy solution has to be taken (orange)



Singular structure

When is the inverse propagator 0?

$$E_{ij} \equiv \Delta_i - p_i^0 + p_j^0 + \Delta_j = 0$$

$$H_{ij} \equiv \Delta_i - p_i^0 + p_j^0 - \Delta_j = 0$$

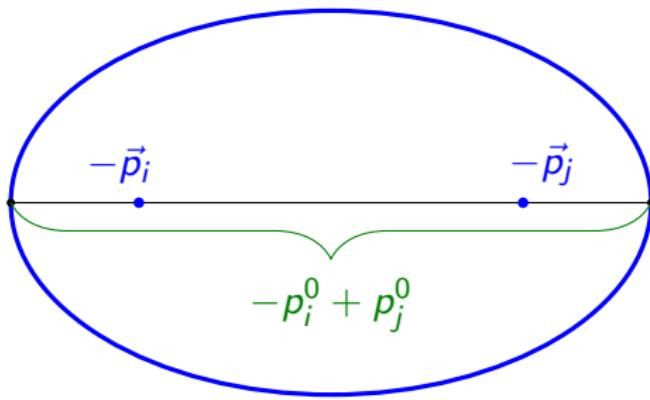
- Δ_i is always ≥ 0
- E_{ij} is an ellipsoid
- H_{ij} is a hyperboloid and cancels with H_{ji}
- All singularities are ellipsoids and are thus **bounded!**

Ellipsoids

$$E_{ij} = \sqrt{(\vec{k} + \vec{p}_i)^2 + m_i^2 - i\epsilon} + \sqrt{(\vec{k} + \vec{p}_j)^2 + m_j^2 - i\epsilon} - p_i^0 + p_j^0 = 0$$

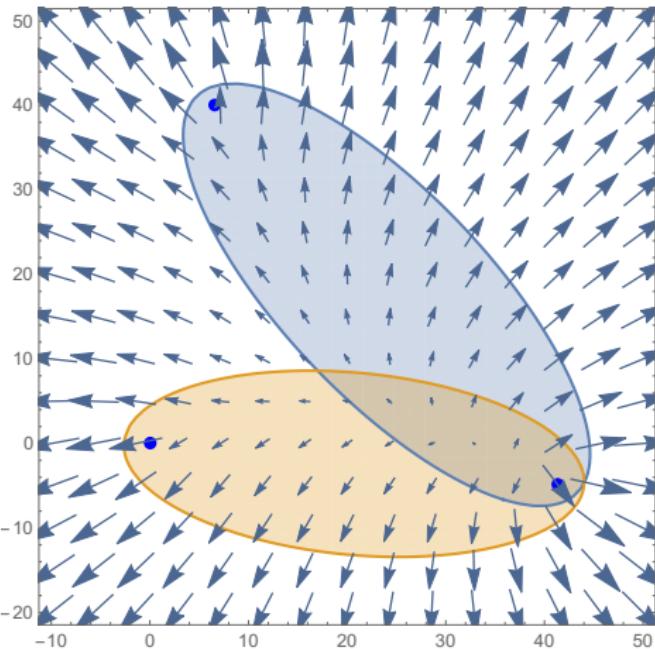
Massless case (and ignoring $i\epsilon$):

$$|\vec{k} + \vec{p}_i| + |\vec{k} + \vec{p}_j| - p_i^0 + p_j^0 = 0$$



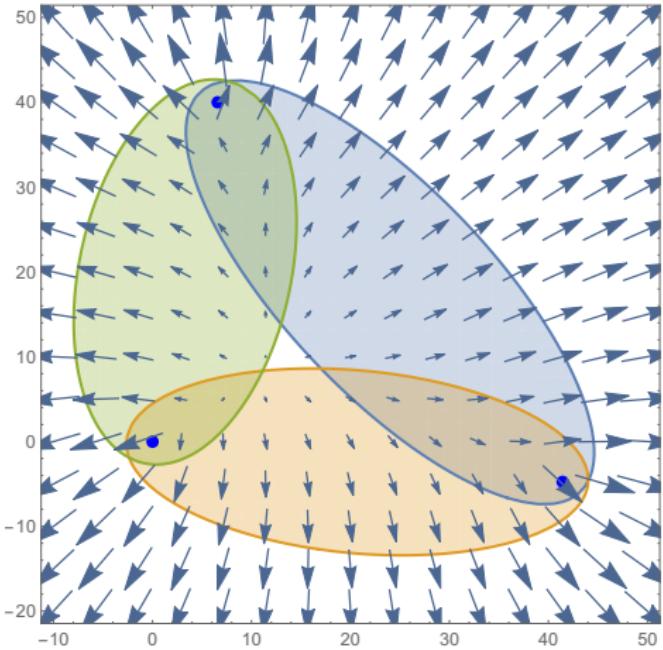
Deformation

- A deformation $k \rightarrow k - \kappa(k)$ needs to have a positive projection on the surface normal
- The sum of normals works for intersections of two ellipsoids [Buchta, Chachamis, Draggiotis, Rodrigo, JHEP 2017]



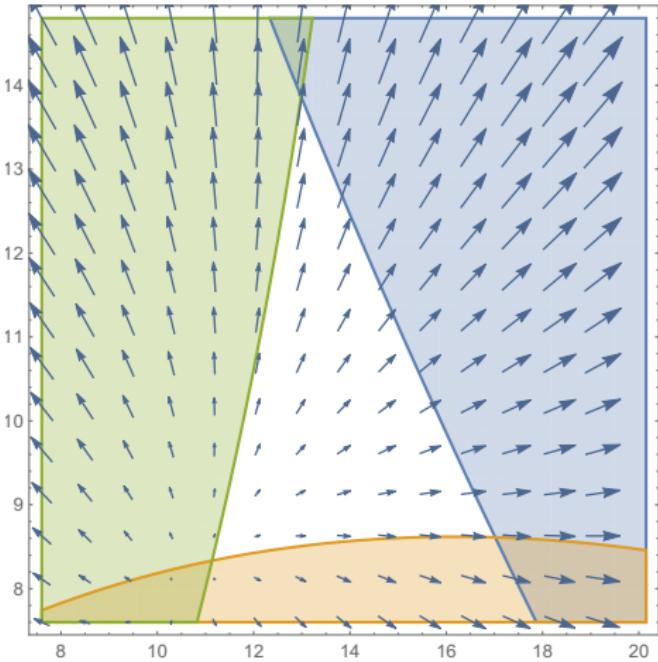
Normal vectors I

- In general the sum will not work
- Exponential dampening around the ellipsoids does not work without extensive finetuning

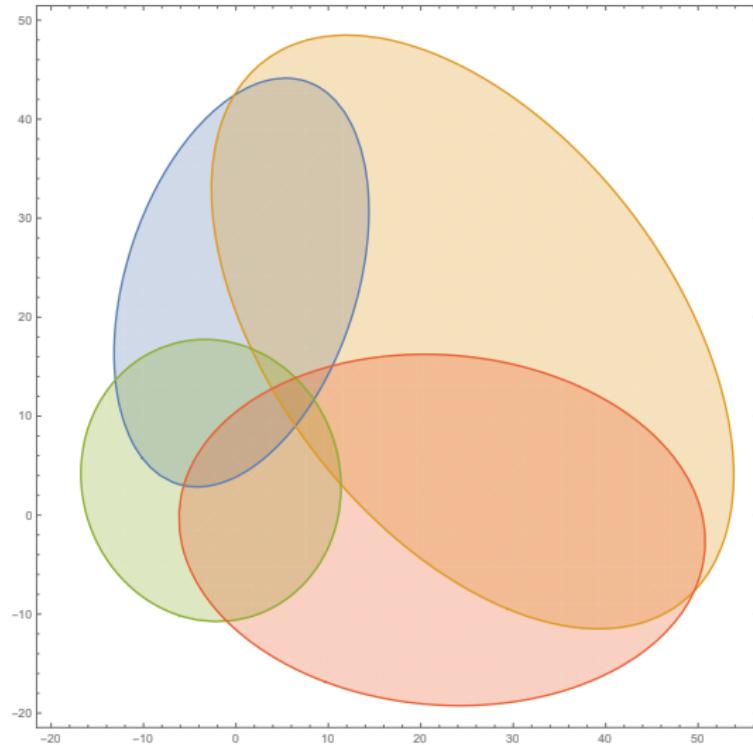


Normal vectors II

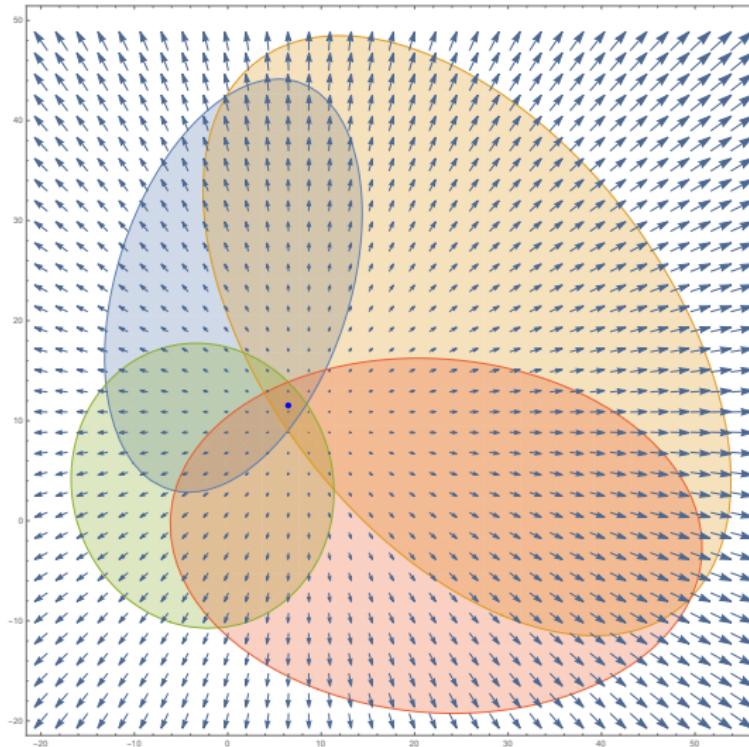
- In general the sum will not work
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How to deform here?



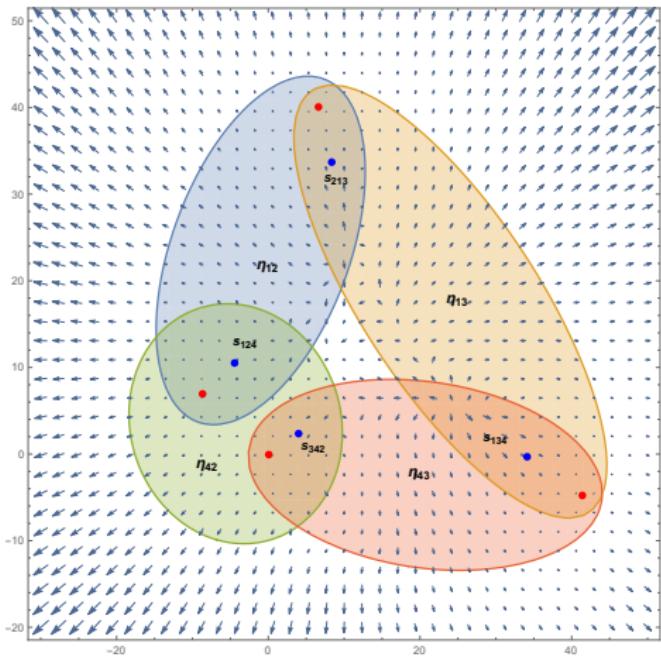
Deformation sources



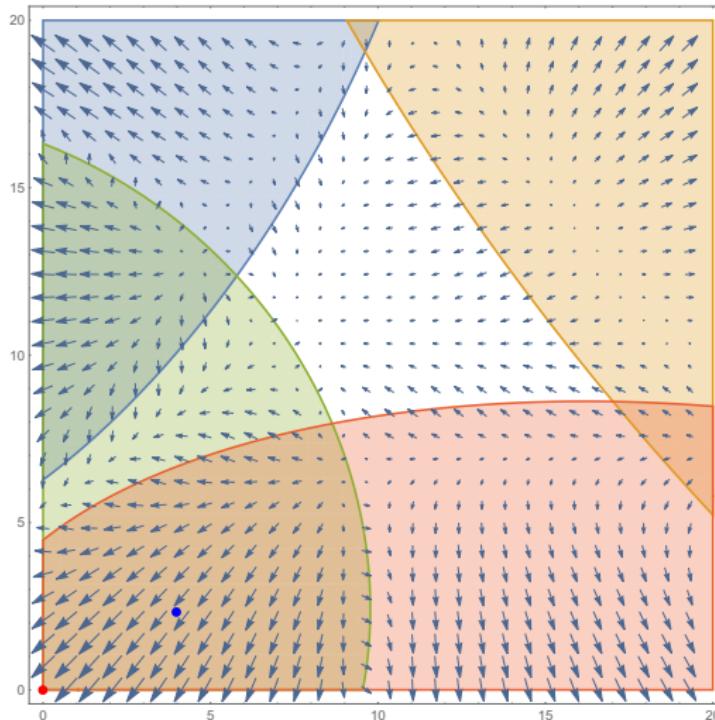
Deformation exclusion

- Deformation field is sum of radial fields multiplied by function $T(E_{ij})$ that is 0 on a surface E_{ij} and goes to 1 away from it
- Exclusion guarantees correct deformation

$$\vec{\kappa} = (\vec{k} - \vec{s}_{124}) T(\textcolor{blue}{E}_{13}) T(\textcolor{red}{E}_{43}) + (\vec{k} - \vec{s}_{213}) T(\textcolor{green}{E}_{42}) T(\textcolor{red}{E}_{43}) + (\vec{k} - \vec{s}_{134}) T(\textcolor{green}{E}_{42}) T(\textcolor{blue}{E}_{12}) + (\vec{k} - \vec{s}_{342}) T(\textcolor{blue}{E}_{12}) T(\textcolor{blue}{E}_{13})$$

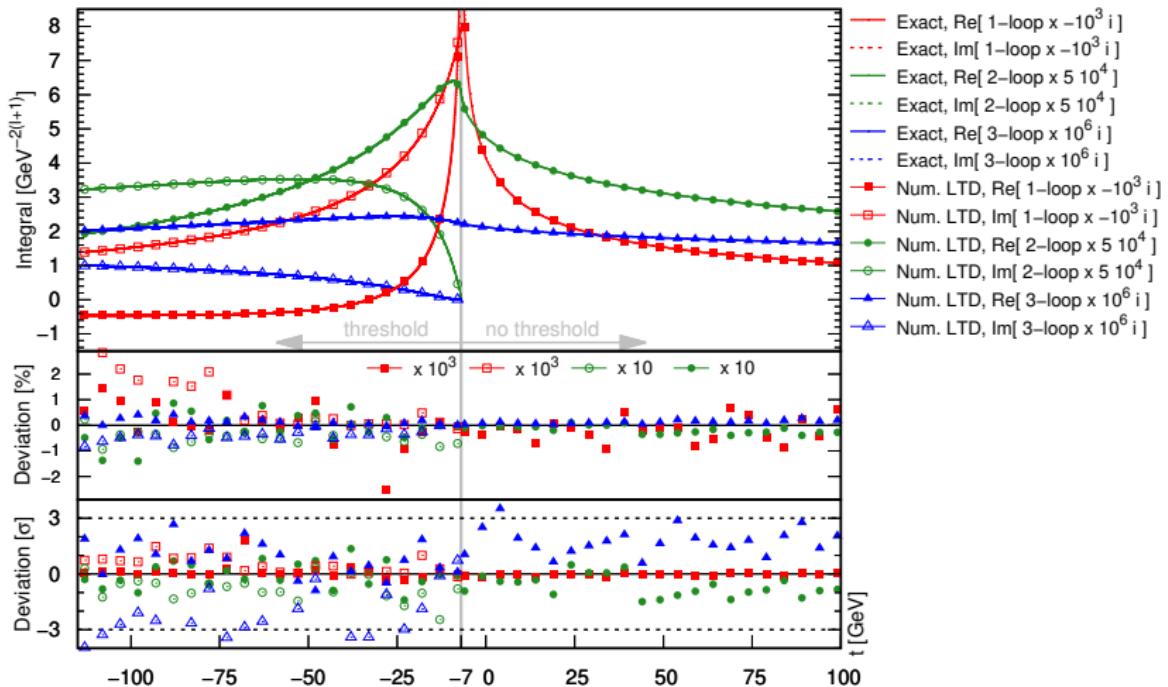


A general and valid deformation



[Capatti, Hirschi, Kermanschah, Pelloni, Ruijl, JHEP 2019]

A scan for multi-loop box topologies



Pinched ellipsoids

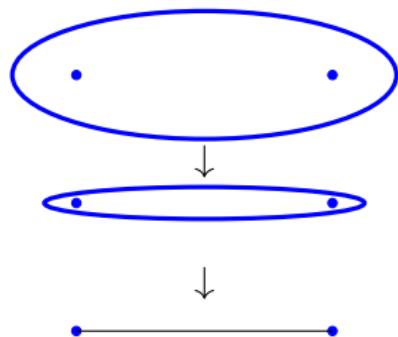
- An ellipsoid is *pinched* when it has no interior
- No deformation possible
- All IR singularities are pinched ellipsoids
- Pinching not possible if there are masses
- Occurs when external momenta are on-shell:

$$|\vec{k} + \vec{p}| + |\vec{k}| - p^0 = 0 \rightarrow$$

$$|\vec{k} + \vec{p}| + |\vec{k}| - |\vec{p}| = 0$$

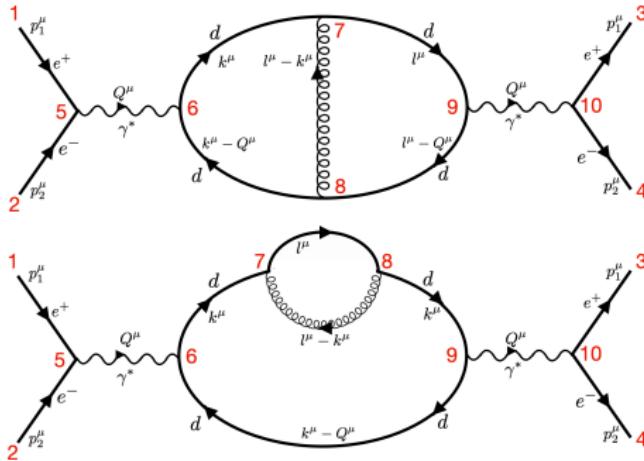
when $\vec{k} = x\vec{p}$ with $0 \leq x \leq 1$

- Collinear singularity with soft singularities on the focal points

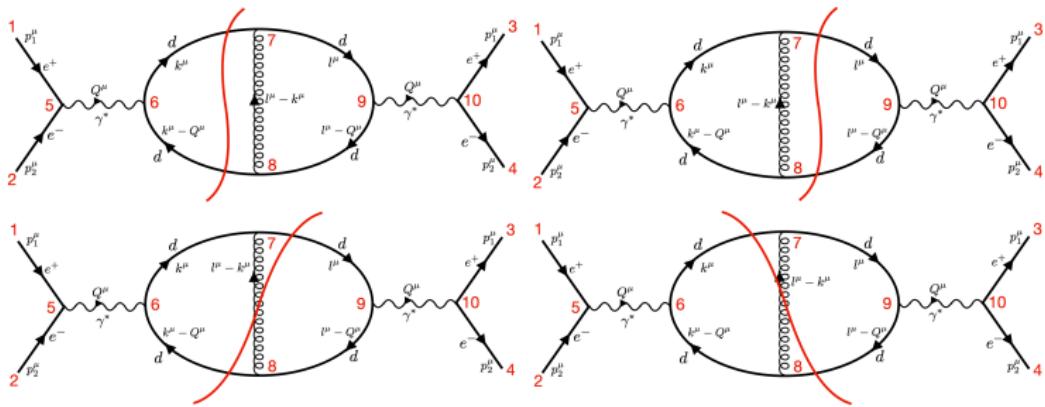


Local Unitarity

- The cross section is given by the sum of all Cutkosky cuts over all *supergraphs*
- Each cut supergraph is IR finite
- For NLO $e^+e^- \rightarrow d\bar{d}$ we have 2 supergraphs:

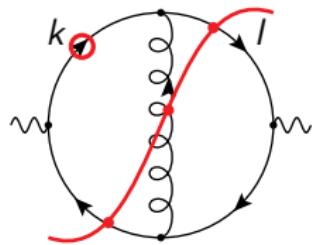


Double-triangle cuts



Cancelling singularities I

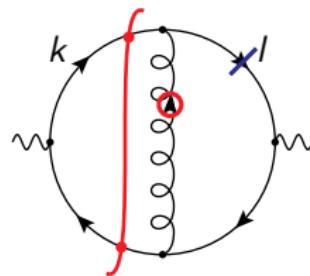
- Consider Cutkosky cuts with additional on-shell propagators:



- $k^0 \pm |\vec{k}| \rightarrow (|\vec{I}| + |\vec{I} - \vec{k}| + |\vec{k}|) (|\vec{I}| + |\vec{I} - \vec{k}| - |\vec{k}|) \xrightarrow{\lim} 2|k|$

Cancelling singularities I

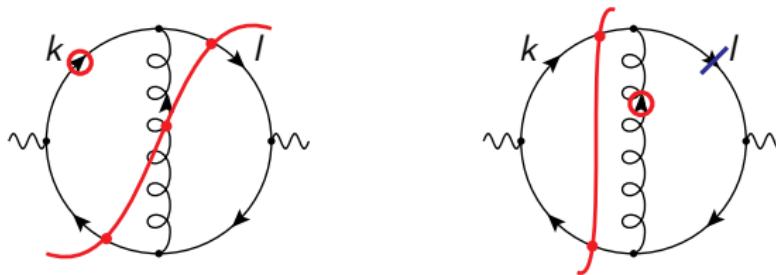
- Consider Cutkosky cuts with additional on-shell propagators:



- $I^0 - k^0 \pm |\vec{I} - \vec{k}| \rightarrow$
 $(|\vec{I}| - |\vec{k}| + |\vec{I} - \vec{k}|) (|\vec{I}| - |\vec{k}| - |\vec{I} - \vec{k}|) \xrightarrow{\lim} -2|\vec{I} - \vec{k}|$

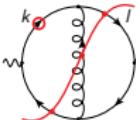
Cancelling singularities I

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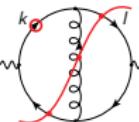
- $k^0 \pm |\vec{k}| \rightarrow (|\vec{I}| + |\vec{I} - \vec{k}| + |\vec{k}|) (|\vec{I}| + |\vec{I} - \vec{k}| - |\vec{k}|) \xrightarrow{\lim} 2|\vec{k}|$
- $I^0 - k^0 \pm |\vec{I} - \vec{k}| \rightarrow (|\vec{I}| - |\vec{k}| + |\vec{I} - \vec{k}|) (|\vec{I}| - |\vec{k}| - |\vec{I} - \vec{k}|) \xrightarrow{\lim} -2|\vec{I} - \vec{k}|$
- Exactly the same on-shell propagators (and 2Δ s), but with a **relative sign**

Cancelling singularities: observable functions


$$= N_3 \delta(p^0 - |\vec{k} - \vec{p}| - |\vec{l} - \vec{k}| - |\vec{l}|) \mathcal{O}_3(-k + p, -l + k, l)$$

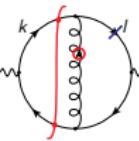
$$\xrightarrow{\lim} N \delta(p^0 - |\vec{k} - \vec{p}| - |\vec{k}|) \mathcal{O}_3(-k + p, -l + k, l) \Big|_{k^0 = |\vec{k}|}$$

Cancelling singularities: observable functions



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$$\xrightarrow{\lim} N \delta(p^0 - |\vec{k} - \vec{p}| - |\vec{k}|) \mathcal{O}_3(-k + p, -l + k, l) \Big|_{k^0 = |\vec{k}|}$$



$$= N_2 \delta(p^0 - |\vec{k} - \vec{p}| - |\vec{k}|) \mathcal{O}_2(k, -k + p)$$

$$\xrightarrow{\lim} -N \delta(p^0 - |\vec{k} - \vec{p}| - |\vec{k}|) \mathcal{O}_2(k, -k + p) \Big|_{l^0 - k^0 = -|\vec{l} - \vec{k}|}$$

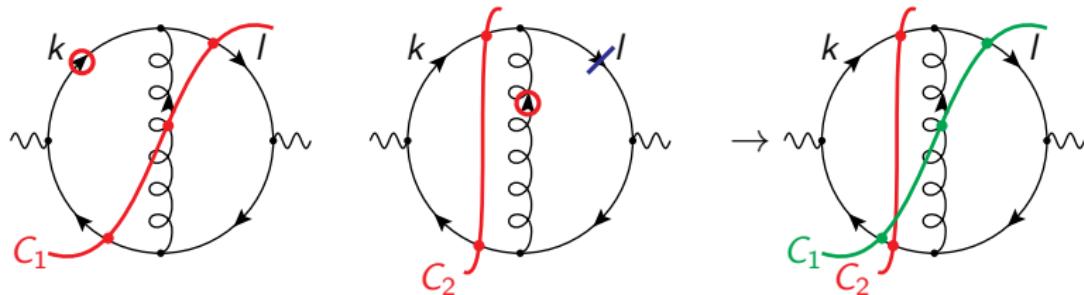
- Energy-conserving δ s are the same on pinches
- Only cancels for IR-safe observables: $\mathcal{O}_3|_{\text{on pinch}} = \mathcal{O}_2|_{\text{on pinch}}$

Cancelling singularities: delta function

$$\propto \delta(p^0 - |\vec{k} - \vec{p}| - |\vec{l} - \vec{k}| - |\vec{l}|)$$
$$\propto \delta(p^0 - |\vec{k} - \vec{p}| - |\vec{k}|)$$

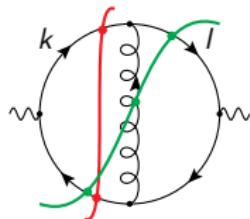
- Solve energy conserving delta function by rescaling loop momenta $\vec{k} \rightarrow t\vec{k}$
- Each interference graph has its own solution for t
- Easy Newton's method for $1 \rightarrow N$ [99 Soper]
- General solution is the *causal flow* [20, Capatti, Hirschi, Pelloni, BR]

Cancellations between Cutkosky cuts

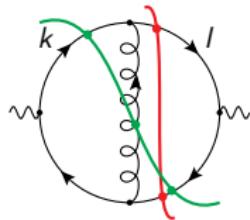
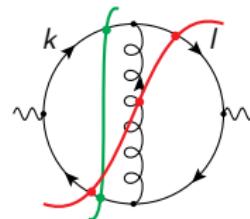


- Singularities lie on intersections of C_1 and C_2
- The intersection point of C_1 and C_2 is finite in $C_1 + C_2$
- The cancellations are pairwise (as with hyperboloids)

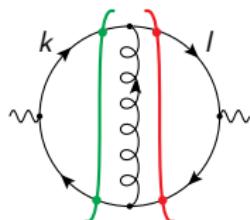
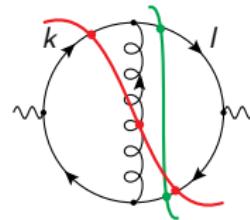
Cancellations between Cutkosky cuts



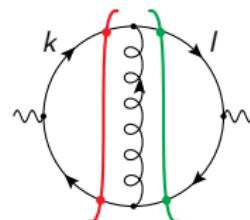
cancels



cancels



cancels



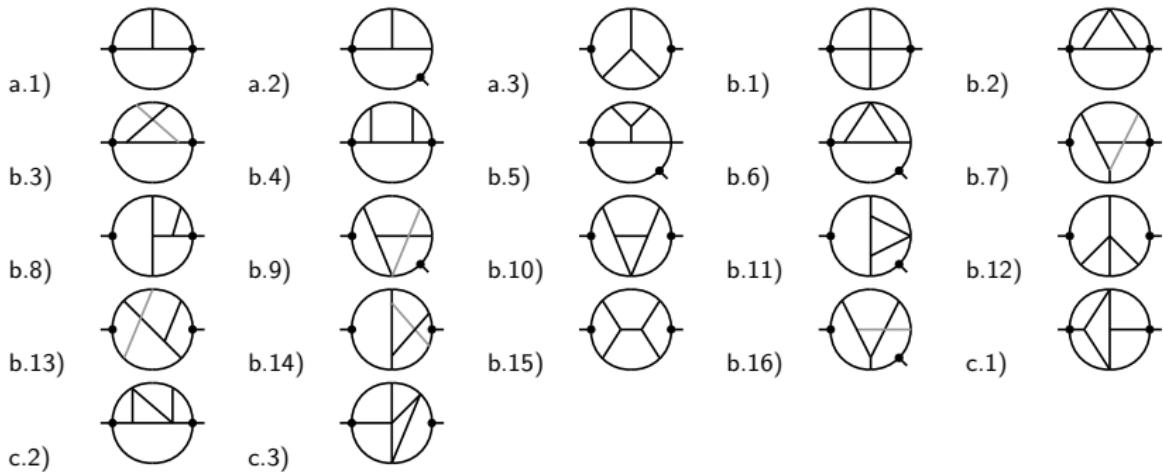
Non-pinched ellipsoid cancellation if the observable is 1!

Local Unitarity workflow

- Generate all supergraphs for a process
- For every supergraph:
 - - Fix a momentum basis and find Cutkosky cuts
 - - For every Cutkosky cut:
 - - - Determine scaling factor and scale momenta
 - - - For the two scaled amplitudes:
 - - - - Determine the deformation field
 - - - - Apply local BPHZ counterterms

This process generates a UV and IR finite integrand

Results for scalar topologies



Results for scalar topologies

- Scalar diagrams have the worst IR divergences
- Example: [Forcer+R*, Herzog, Ruijl, Ueda, Vermaseren, JHEP 2018]



optical theorem $\rightarrow \frac{5\pi}{(16\pi)^5} \frac{441}{40} \zeta_7 = 1.77832 \cdot 10^{-9}$

- LU with 1M samples: [Capatti, Hirschi, Kermanshah, Pelloni, Ruijl, JHEP 2020]

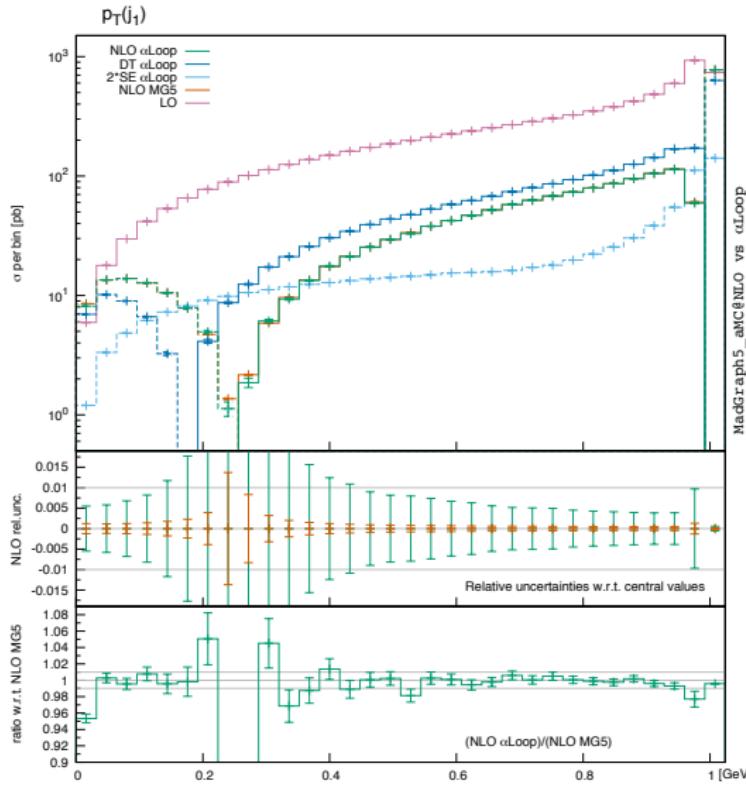
$$I = 1.7797(33) \cdot 10^{-9}$$

$$\Delta_\sigma = 0.42\sigma$$

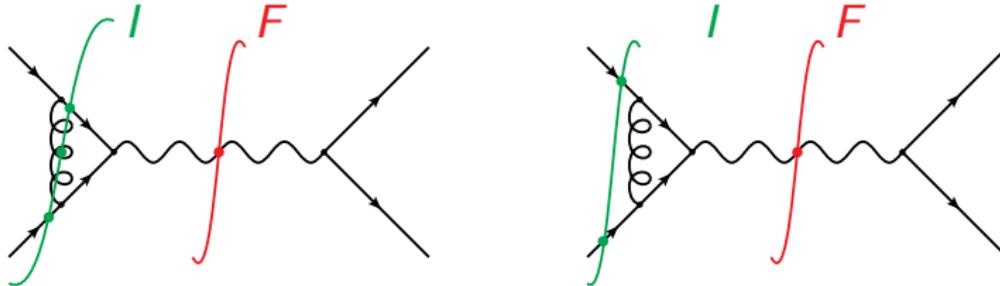
$$\Delta\% = 0.077\%$$

NLO correction to $e^+e^- \rightarrow Ht\bar{t}$

- 15 supergraphs Monte Carlo sampled over
- $\mathcal{O}(50)$ interference graphs (cuts)
- MG5_aMC@NLO: $-1.38400 \cdot 10^{-4} \pm 1.4 \cdot 10^{-7}$ [Alwall, Frederix, Frixione, Hirschi, Maltoni]
- Local Unitarity: $-1.38320 \cdot 10^{-4} \pm 5.9 \cdot 10^{-7}$

NLO correction to $e^+e^- \rightarrow d\bar{d}$ 

Initial-state radiation



- Initial-state infrared singularity cancellations to be studied
- Suggests extra particles in the initial state (as in the final state)
- Fixing the number of initial-state particles is a non-IR-safe observable
- Find connection to classical PDFs and PDF counterterms

Conclusion

The Local Unitarity approach

- locally realizes the KLN-theorem and regulates final-state radiation
- uses a contour deformation to regulate threshold singularities
- automatically renormalizes in \overline{MS} using local BPHZ
- can be extended to initial-state radiation

Status:

- Sample time around 1ms at NNLO
- Verification at NLO
- Working on automatic on-shell renormalization at NNLO
- Working on improving convergence rate

Thank you for your attention.