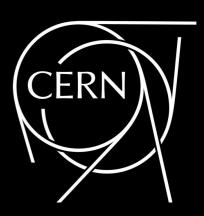
Measurements of diffractive physics and soft QCD at ATLAS

Lydia Beresford on behalf of the ATLAS Collaboration EPS-HEP Conference, 26th July 2021



Soft QCD laboratory



Run: 267638 Event: 242090708 2015-06-14 01:01:14 CEST $Z \rightarrow \mu \mu$

T

Motivation

 Key area of SM where knowledge of fundamental processes is limited

Theoretically:

- Beyond pQCD regime
- Employ phenomenological models with tunable parameters
- Measurements are vital
- Crucial input for other LHC searches + measurements & beyond!

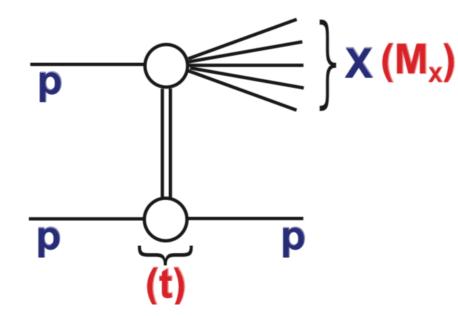
Soft QCD in the sky!

cosmic ray air shower

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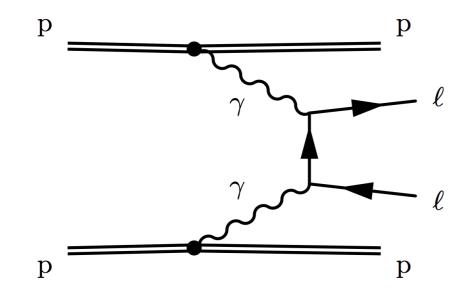
Focus: Proton tagging to probe QCD

Measurement of single diffraction using ALFA



$$\frac{\text{JHEP 02 (2020) 042}}{\sqrt{s}} = 8 \text{ TeV, very low } \langle \mu \rangle$$

Observation & measurement of $(\gamma\gamma \rightarrow \ell \ell) + p$ using AFP



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PRL 125 (2020) 261801

$$\sqrt{s}$$
 = 13 TeV, standard high $\langle \mu \rangle$

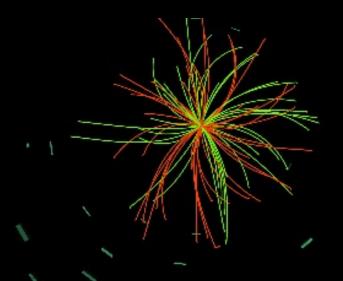
 $\langle \mu \rangle$ = mean number of interactions per bunch crossing 4

Single diffraction



Pb+Pb, 5.02 TeV Run: 365681 Event: 1064766274 2018-11-11 22:00:07 CEST

Rapidity gap example in PbPb



 $\Sigma E_{T}^{FCal} = 71 \text{ GeV} (left), 0.9 \text{ GeV} (right)$ 71 tracks, $p_{\rm T} > 0.4$ GeV

Single Diffraction

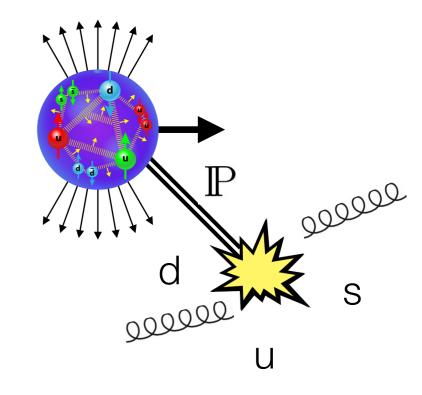
Colour singlet 'pomeron' exchange

Diffractive system X as result of interaction between proton and pomeron

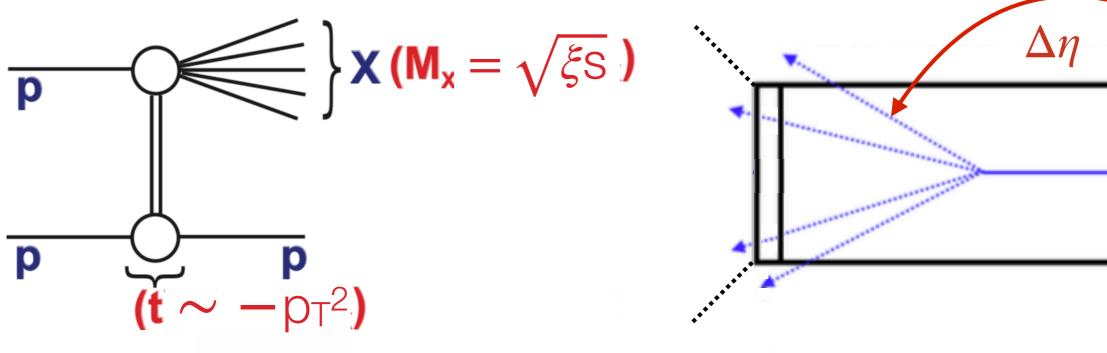
Large cross-section: ~10% of total LHC cross-section but poorly understood

Input for MC generators, improve:

- Pile-up modelling
- Modelling of cosmic ray air showers



Goals & motivation Hereboxic Proton fractional energy loss Measure cross-section differentially in $\Delta \eta$, [t] & ξ



Measured wrt edge of inner detector on side of intact proton

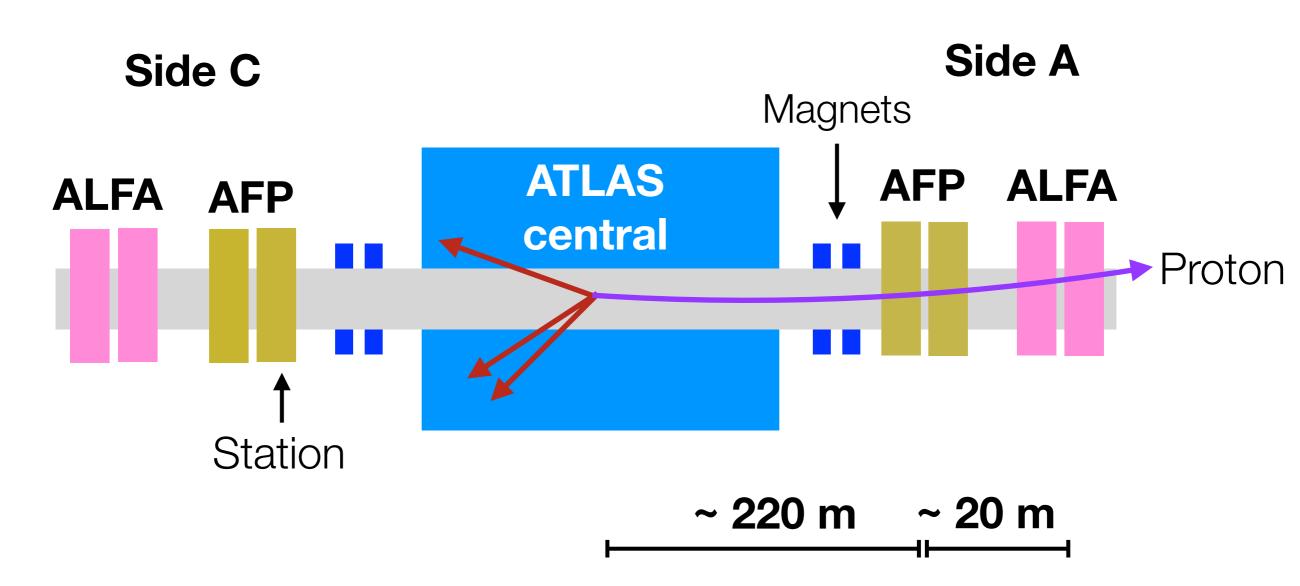
Measure intact proton:

- Suppress backgrounds
- Can measure t dependence (& alternative ξ measurement)

±2.5

ALFA & AFP detectors

Detectors reach inside beam pipe down to 2 mm of incoming proton beam



Reconstruct tracks in station \rightarrow **Proton**

Complementary acceptance in AFP & ALFA

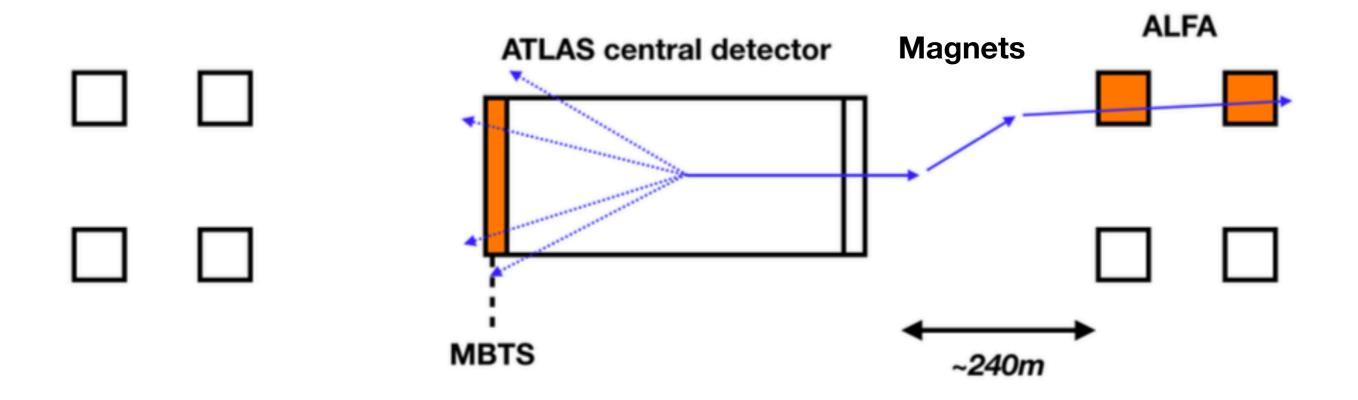
ALFA active during special low pile-up, high β^* runs

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How?

Track $p_T > 200 \text{ MeV}$

Special run from 2012: \sqrt{s} = 8 TeV, $\langle\mu\rangle$ < 0.08 , high β^*

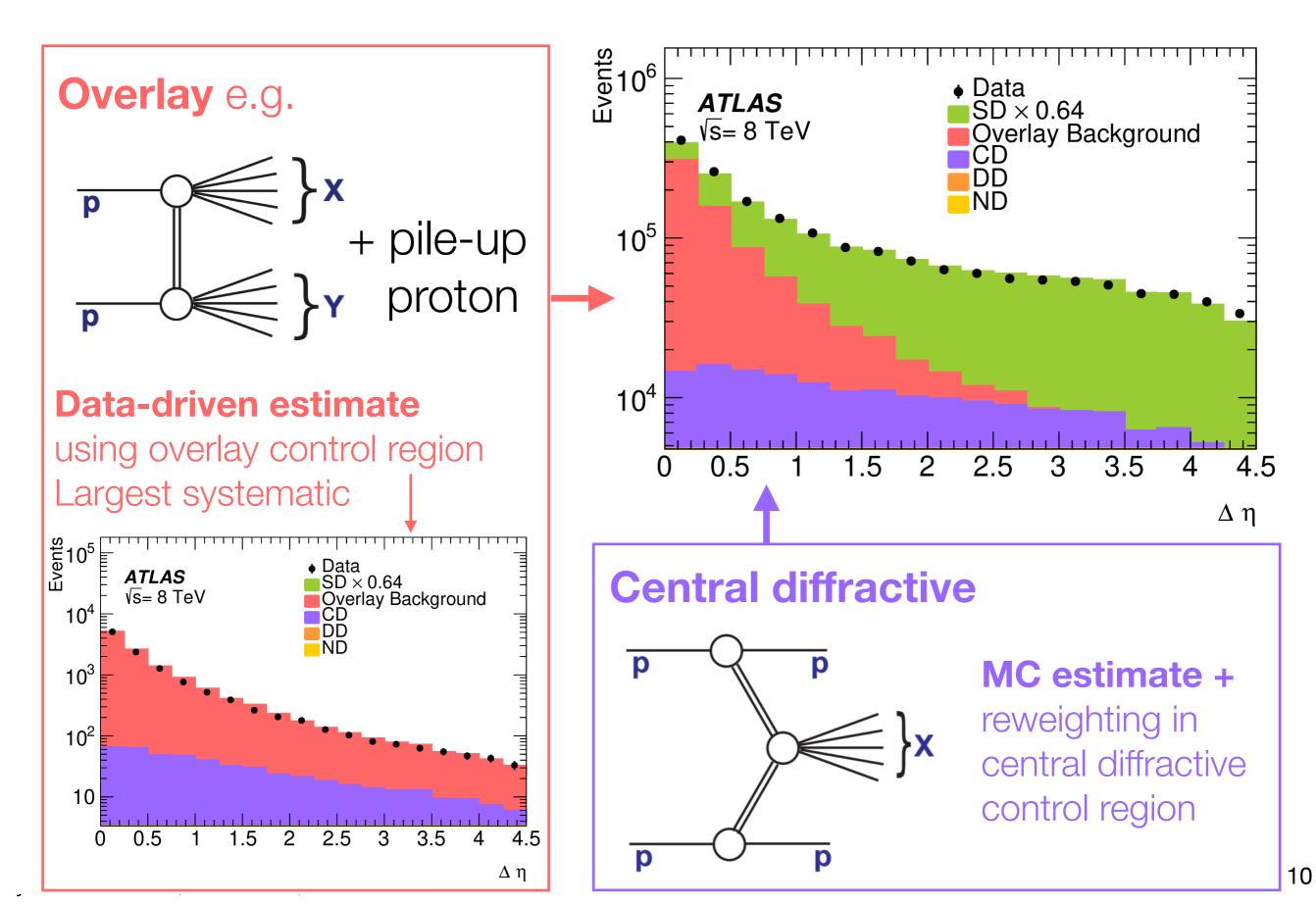


Trigger: ALFA signal & Min Bias Trigger Scintillator

Single proton in ALFA & one good vertex

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How? Backgrounds estimated & subtracted



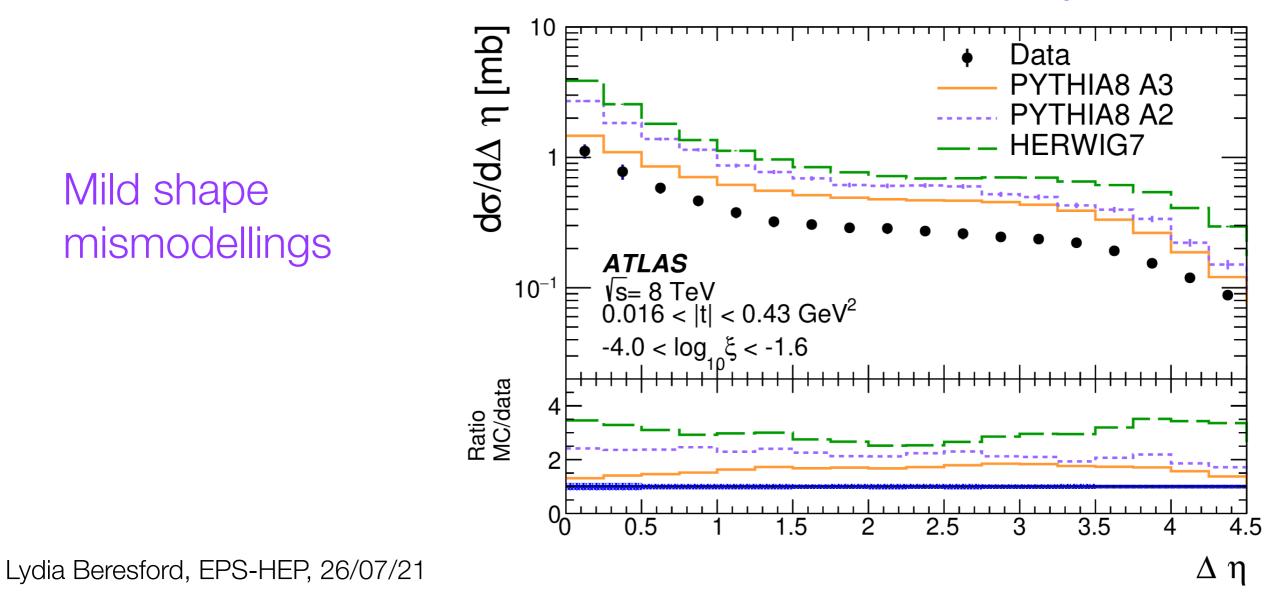
$0.016 < |t| < 0.43 \text{ GeV}^2$

$-4.0 < \log_{10}(\xi) < -1.6$ (80 < M_x < 1270 GeV)

	Distribution	$\sigma_{\mathrm{SD}}^{\mathrm{fiducial}(\xi,t)}$ [mb]	$\sigma_{\rm SD}^{t-{\rm extrap}}$ [mb]
MC	Data	1.59 ± 0.13	1.88 ± 0.15
over-prediction	Рутніа8 A2 (Schuler–Sjöstrand)	3.69	4.35
over-prediction	PYTHIA8 A3 (Donnachie–Landshoff)	2.52	2.98
	Herwig7	4.96	6.11

Cross-section as function of rapidity gap size $\Delta\eta$

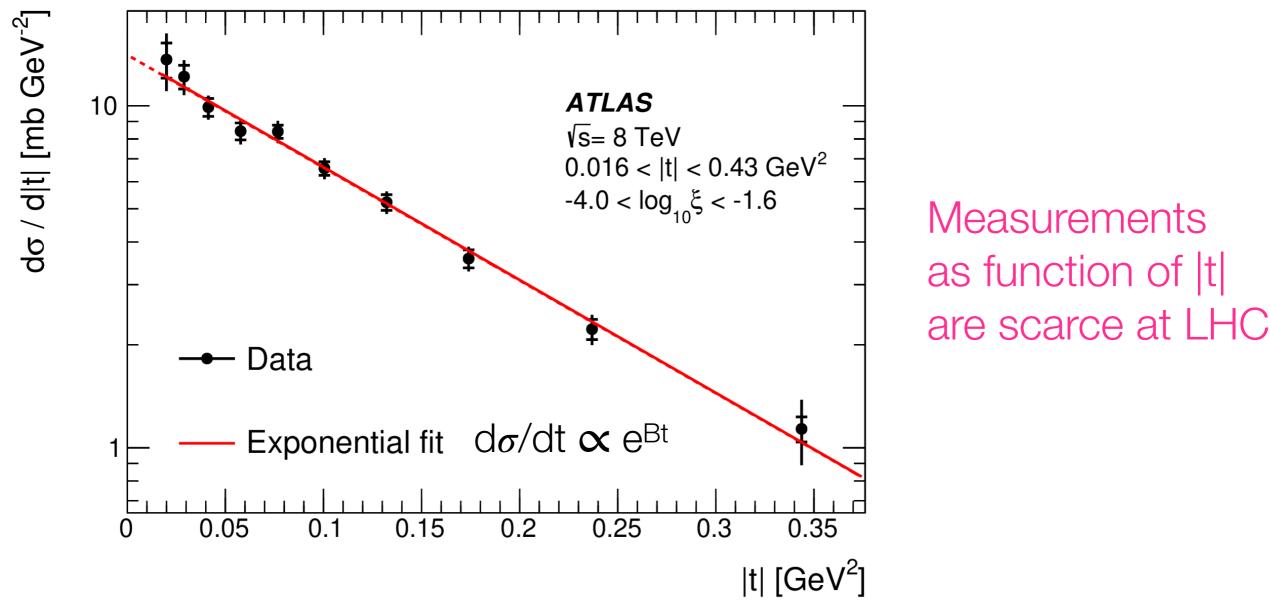
Results



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Results

Cross-section as function of momentum transfer |t|



Fit to data

 $B = 7.65 \pm 0.26$ (stat.) ± 0.22 (syst.) GeV⁻²

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MC Predictions

Pythia8 A2: B = 7.82 GeV^{-2} Pythia8 A3: B = 7.10 GeV^{-2}

Photon collision

$\gamma\gamma \to \mu\mu$



Run 183081, Event 94526500 Time 2011-06-05, 16:37 CEST

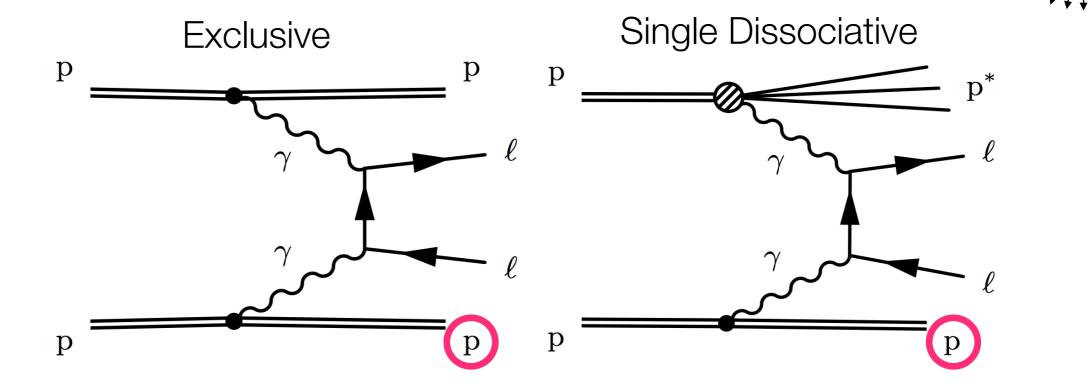
LHC is world's highest energy photon collider

μ

Goals & motivation

Observe & measure $(\gamma \gamma \rightarrow \ell \ell) + p$

First high- $\langle \mu \rangle$ physics publication using **AFP detectors** \checkmark



Direct access to protons provides truly new info to exploit!

- Powerful background rejection
- Novel constraints on proton soft survival probability

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How?

Use 2017 standard high $\langle \mu \rangle$ dataset \sqrt{s} = 13 TeV

Signal

AFP AFP

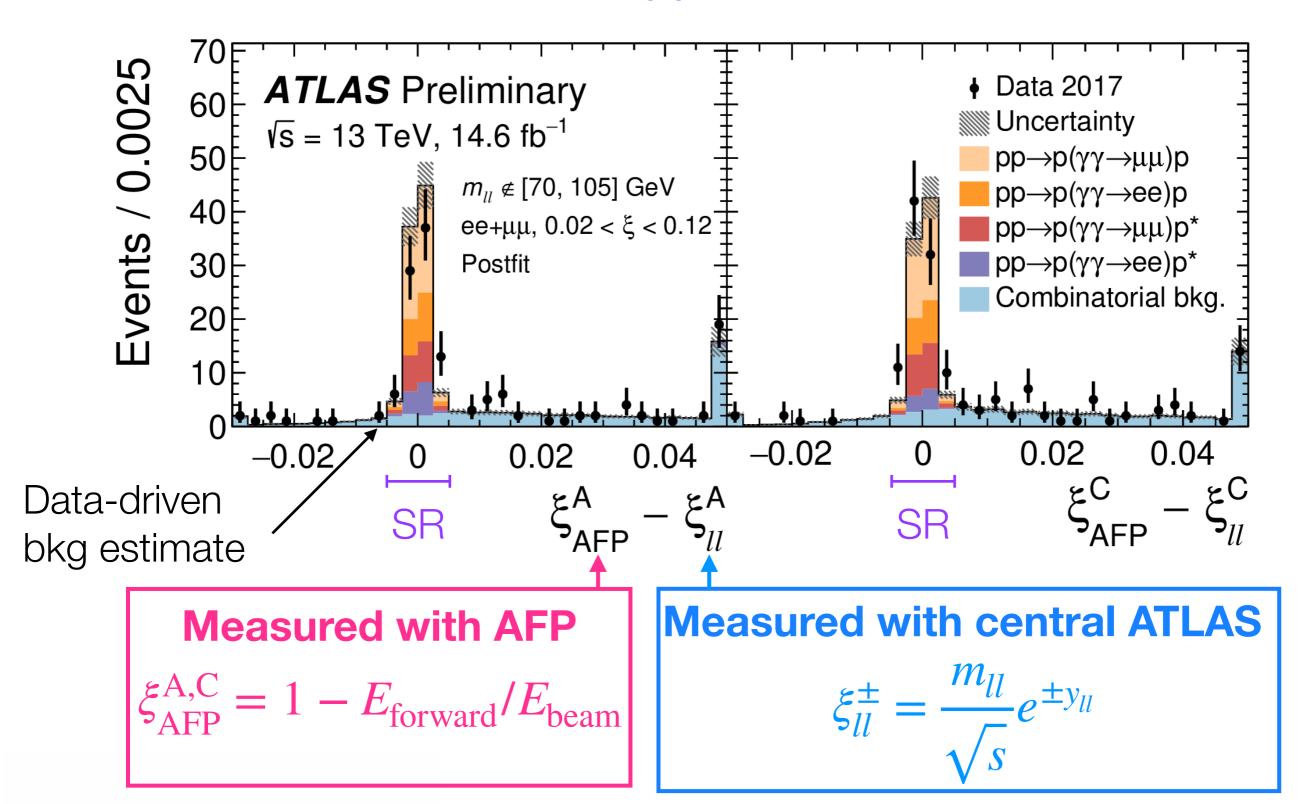
Background

e.g. $Z \rightarrow \ell \ell$ + pileup proton

How?

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Observe ee+p and $\mu\mu$ +**p > 9** σ in each channel



Results

1st LHC cross-section measurements for this process with a tagged proton

$\sigma_{\mathrm{HERWIG+LPAIR}} imes S_{\mathrm{surv}}$	$\sigma_{ee+p}^{ ext{fid.}}$ (fb)	$\sigma^{ ext{fid.}}_{\mu\mu+p}$ (fb)
$S_{\rm surv} = 1$	15.5 ± 1.2	13.5 ± 1.1
S_{surv} using Refs. [33,34]	10.9 ± 0.8	9.4 ± 0.7
SUPERCHIC 4 [97]	12.2 ± 0.9	10.4 ± 0.7
Measurement	11.0 ± 2.9	7.2 ± 1.8

Fiducial cross-sections $\xi \in [0.035, 0.08]$ [33] Harland-Lang et al EPJC 76 (2016) 9Compared to proton soft survival models[34] Dyndal & Schoffel PLB 741 (2015) 66[97] Harland-Lang et al EPJC 80, 925 (2020)

Outlook:

- Statistically limited \rightarrow Improvements with larger dataset in LHC Run 3
- Inclusive measurement \rightarrow Differential measurements e.g. versus m_{II}, y_{II}
- Single proton tag \rightarrow Measure both single and double tag events

Summary

Diverse array of QCD phenomena at the LHC, including:

- Single diffraction
- Proton break up

Target using array of techniques:

- Rapidity gaps
- Forward proton detectors

Advance our understanding of nature

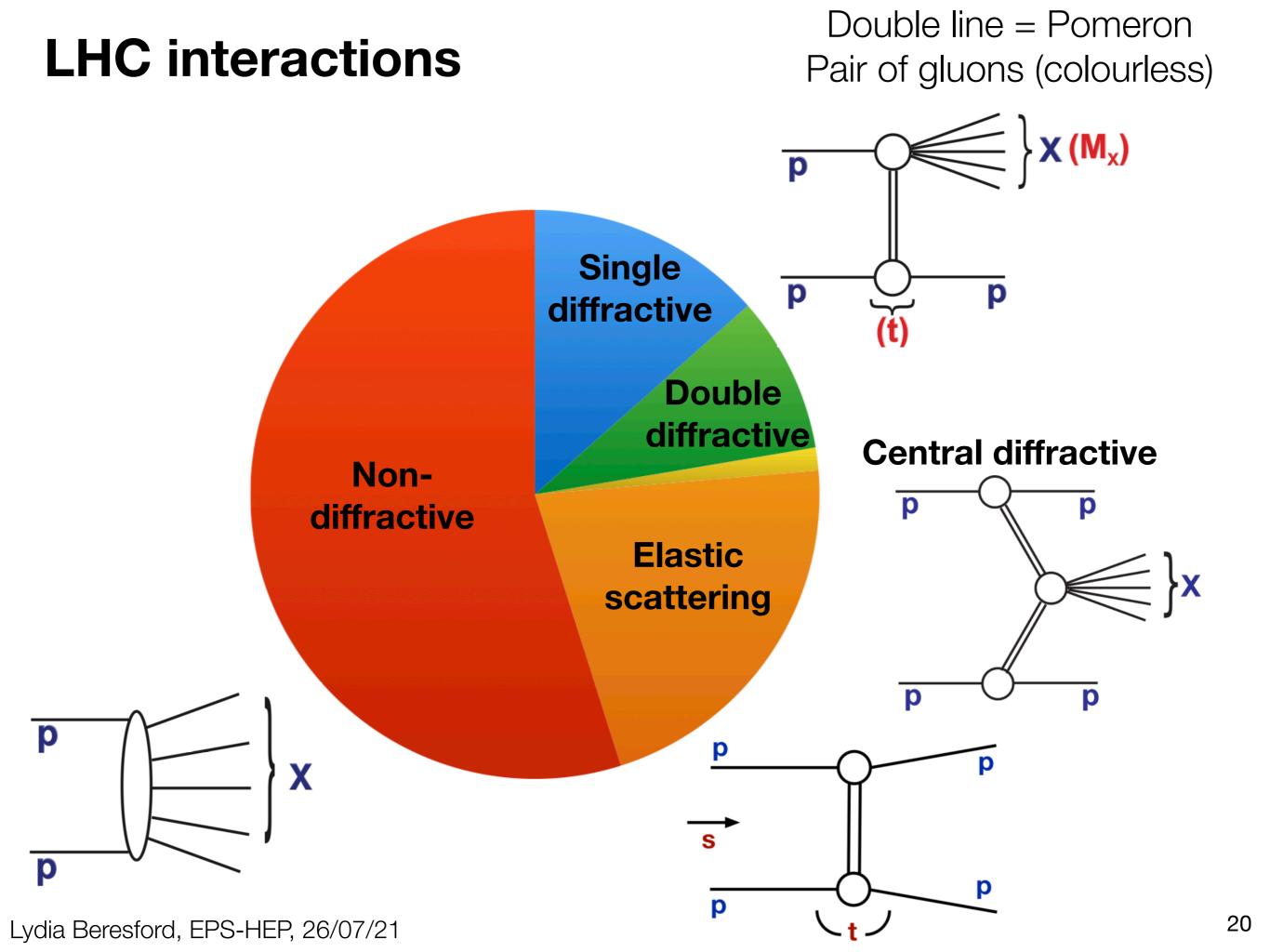
Just as nonperturbative QED contains very interesting phenomena, nonperturbative QCD is a most interesting portion of that theory. To me, it is *the* most interesting and most important portion of QCD to address, despite the evident difficulty in doing so. *Bjorken 1996*



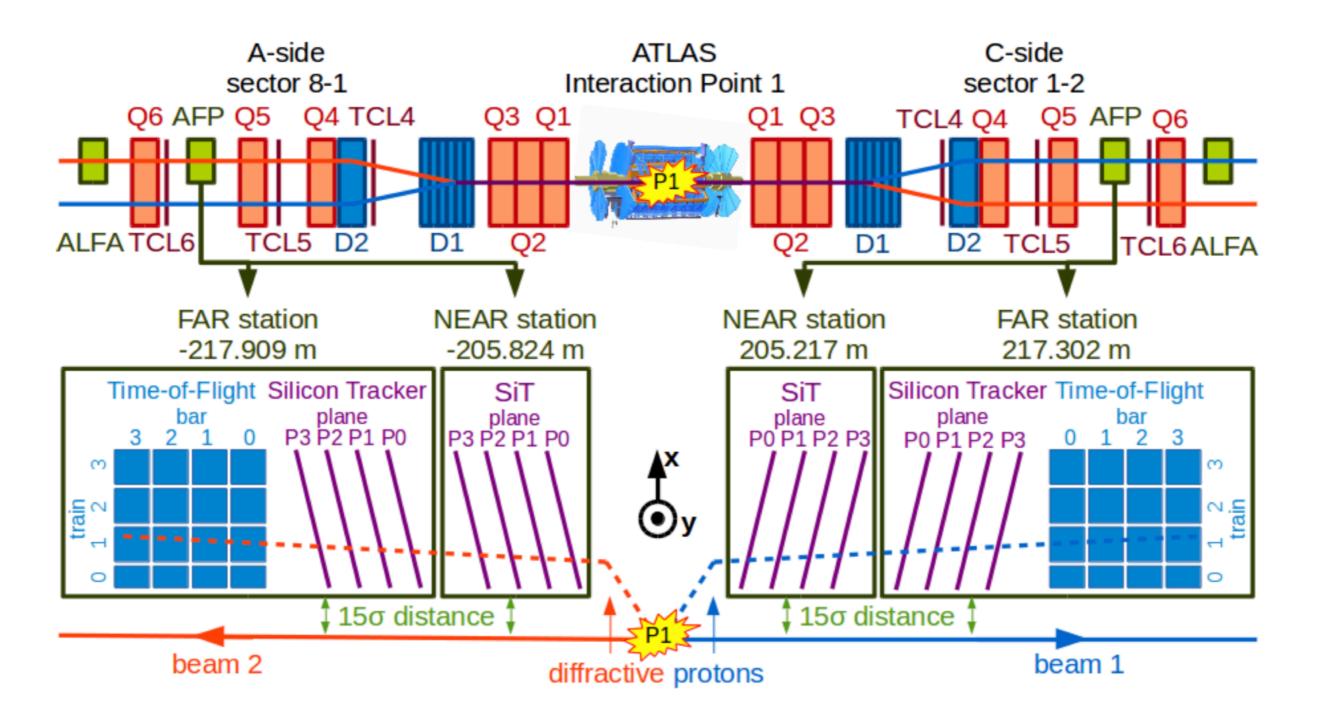
cosmic ray air shower 1

Backup

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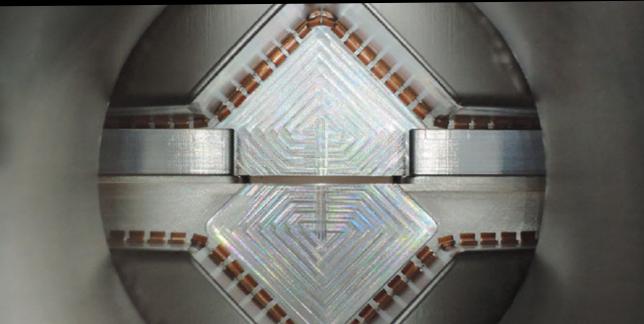
ALFA & AFP



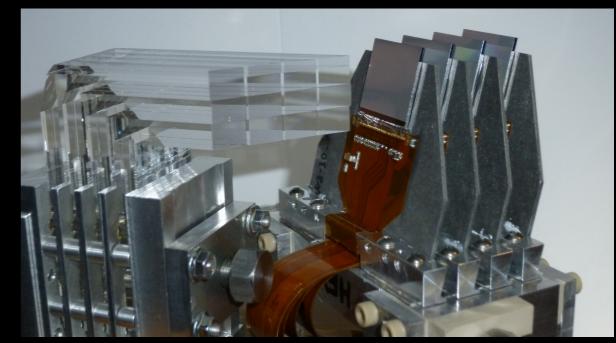
ALFA detector

AFP detector









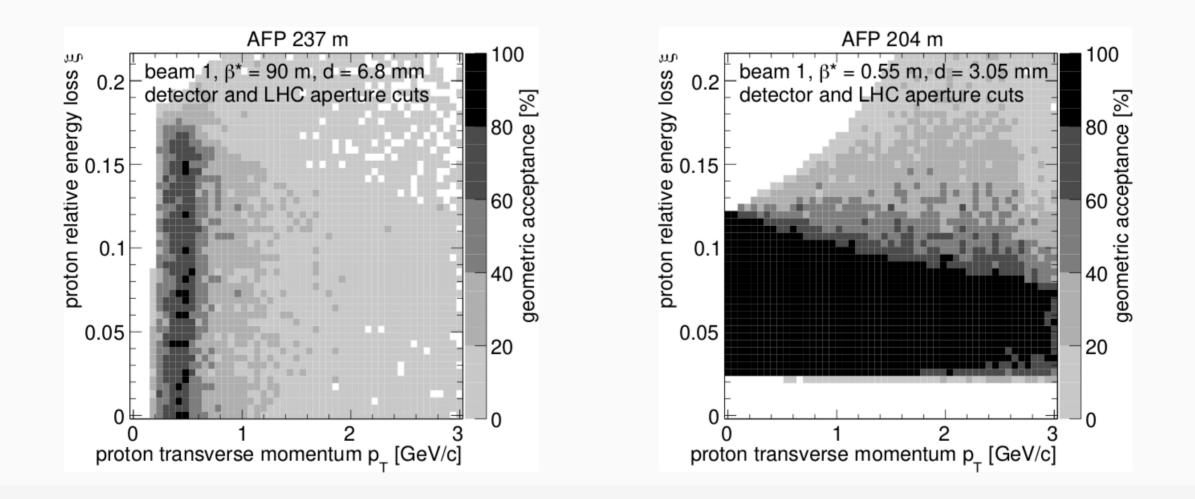
Complementarity of ALFA and AFP

ALFA

- vertical roman pots
- special optics
- (almost) full acceptance in ξ
- good resolution in t

AFP

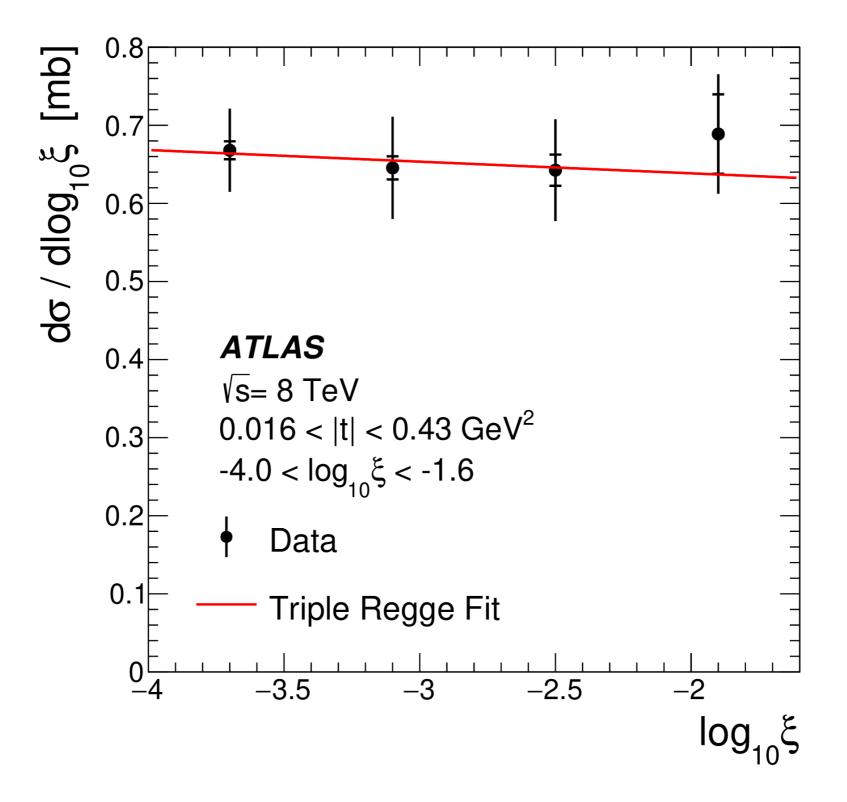
- horizontal roman pots
- standard optics
- \cdot (almost) full acceptance in t
- good resolution in ξ



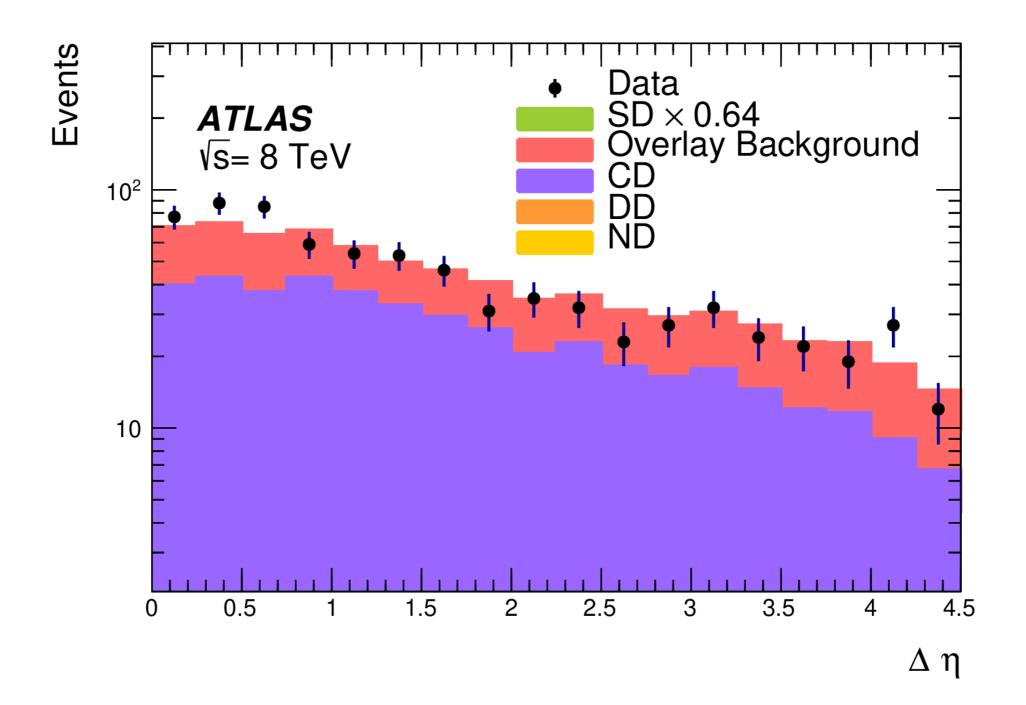
ALFA analysis MC

Monte Carlo (MC) simulations are used for the modelling of background contributions, unfolding of instrumental effects, and comparisons of models with the hadron-level cross-section measurements. The PYTHIA8 [21] generator was used to produce the main SD, ND and DD samples and also that for the 'central diffractive' (CD, $pp \rightarrow pXp$, Figure 1(c)) process. The SD, DD and CD models in PYTHIA8 are based on the exchange of a pomeron with trajectory $\alpha(t) = \alpha(0) + \alpha' t$, assuming 'triple Regge' [22] formalism (see Section 10). The models [23] are tuned using previous ATLAS data, including the total inelastic cross section [11] and rapidity gap spectra [14]. By default, the 'A3' tune [24] was used, which adopts the 'Donnachie–Landshoff' [25] choice for the pomeron flux factor to describe the ξ and t dependences in the diffractive channels with pomeron intercept $\alpha(0) = 1.07$. An alternative SD sample was produced using the A2 tune [26] and the Schuler–Sjöstrand model for the pomeron flux factor [23], which has $\alpha(0) = 1$ and therefore differs from Donnachie–Landshoff mainly in its ξ dependence. Both tunes use the H1 2006 Fit B diffractive parton densities [27] as an input to model the hadronisation in the diffractive channels. For the non-diffractive channel, the A3 tune uses the NNPDF23LO [28] proton parton densities. Generated central particles were propagated through the GEANT4 based simulation of ATLAS [29, 30] to produce the simulated signals in the central detector components. The generated protons in diffractive processes are transported from the interaction point to the ALFA detectors by representing each element of the LHC optical lattice (quadrupole and dipole magnets) as a simple matrix under the thin-lens approximation, giving the total transfer matrix once multiplied together.

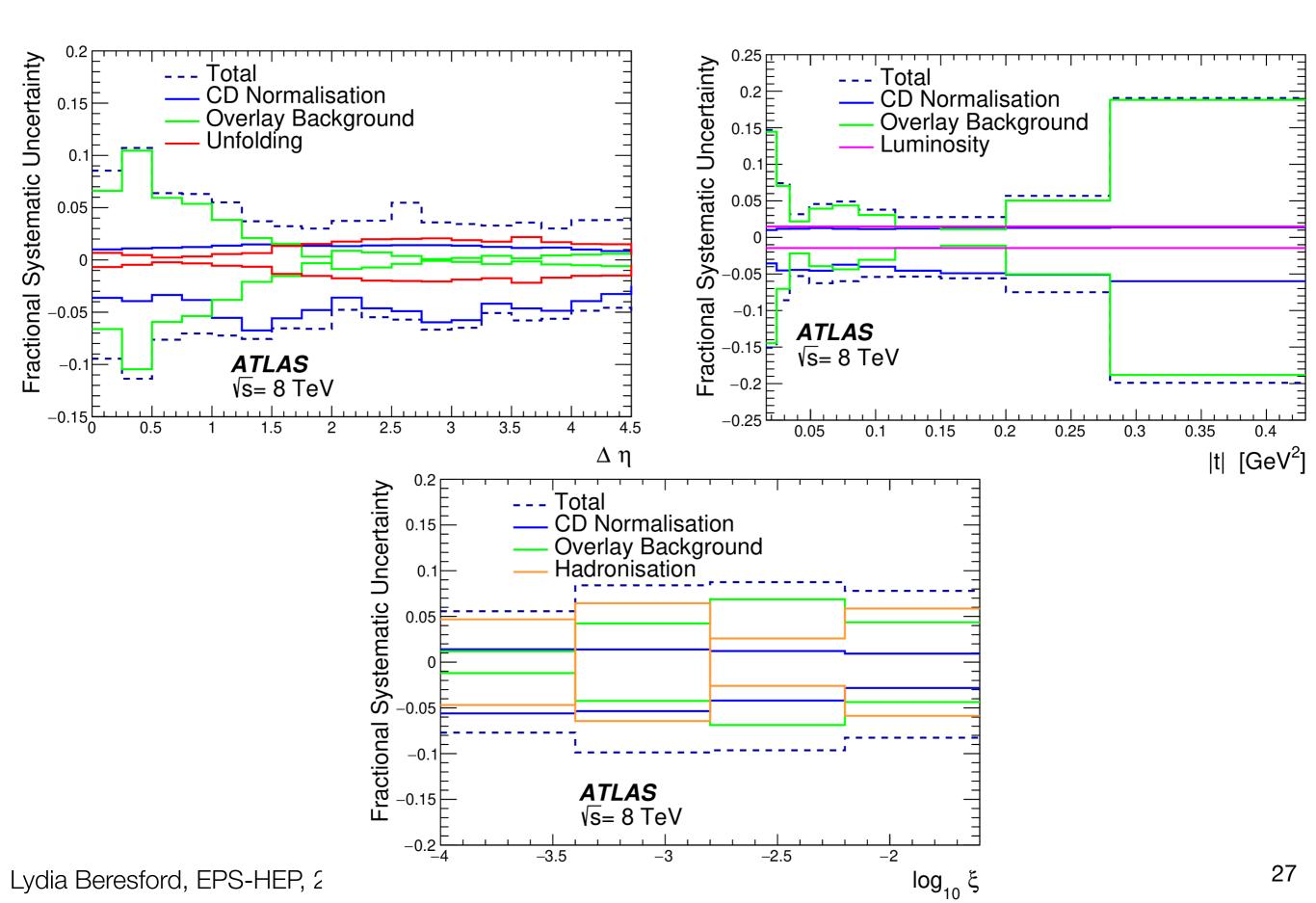
ALFA measurement



ALFA CD control region



SD analysis uncertainties



AFP MC

Simulated events of the exclusive signal $pp \rightarrow p(\gamma\gamma \rightarrow \ell^+\ell^-)p$ were produced using the HERWIG7 Monte Carlo (MC) generator [64, 65]. The single-dissociative signal $pp \rightarrow p(\gamma\gamma \rightarrow \ell^+\ell^-)p^*$ was generated using LPAIR4.0 [66], with proton dissociation modeled using the Brasse et al. [67] and Suri-Yennie [68] structure functions interfaced with JETSET7.408 [69, 70]. Simulation of these processes is detailed in Ref. [5]. To model the central-detector response, the exclusive signal sample underwent full detector simulation based on GEANT4 [71]. The single-dissociative samples employed a fast simulation [72], which uses a parametrization of the calorimeter response [73]. The response of the AFP spectrometer is modeled by a fast simulation, where a Gaussian smearing is applied to track positions based on the AFP spatial resolution. Simulated samples include the effect on the central detector of multiple pp interactions in the same and neighboring bunch crossing (pileup), as detailed in Ref. [5].

AFP selection

Reconstructed events must contain at least one interaction vertex with two or more associated inner-detector tracks that satisfy $p_{\rm T} > 500$ MeV, $|\eta| < 2.5$, and the "Loose" criterion [74, 75]. Electrons (muons) must satisfy $p_{\rm T} > 18$ (15) GeV, $|\eta| < 2.47$ (2.4), the "LooseAndBLayer" [76] ("Medium" [77]) identification criterion, and $|z_0 \sin \theta| < 0.5$ mm [78]. Electrons sharing an inner-detector track with a muon are discarded. To suppress fake and/or nonprompt lepton backgrounds, remaining electrons (muons) must satisfy transverse impact parameter significance $|d_0/\sigma_{d_0}| < 5$ (3) and isolation requirements described in Ref. [79] (Ref. [80]). Electrons must also satisfy "Medium" identification [76]. Small corrections are applied to leptons in simulated samples to match reconstruction and trigger efficiencies measured in data, as described in Refs. [76, 77].

Selected events must have exactly two same-flavor leptons with opposite electric charge $(e^+e^- \text{ or } \mu^+\mu^-)$ and be matched to the leptons that triggered the event. To suppress quarkonia and Z boson resonances, the dilepton invariant mass must satisfy $m_{\ell\ell} > 20$ and $m_{\ell\ell} \notin [70, 105]$ GeV. To select events compatible with $pp \rightarrow p(\gamma\gamma \rightarrow \ell^+\ell^-)p^{(*)}$ processes based on the simulated signals, the dilepton transverse momentum must satisfy $p_T^{\ell\ell} < 5$ GeV. This set of criteria is referred to as the preselection. Signal event candidates must additionally have small acoplanarity $A_{\phi}^{\ell\ell} = 1 - |\Delta\phi_{\ell\ell}|/\pi < 0.01$. These events must have no inner-detector tracks $(N_{\text{tracks}}^{0.5 \text{ mm}} = 0)$ that satisfy $\Delta R(\text{track}, \ell) > 0.01$ for both leptons and $|z_0^{\text{track}} - z_0^{\ell\ell}| < 0.5$ mm, where z_0^{track} is the track z_0 position and $z_0^{\ell\ell} = (z_0^{\ell_1} + z_0^{\ell_2})/2$ with $\ell_{1,2}$ denoting the two leptons. The expected proton energy loss based on lepton kinematics $\xi_{\ell\ell}$ is determined from $m_{\ell\ell}$ and the dilepton rapidity $y_{\ell\ell}$ by momentum conservation $\xi_{\ell\ell}^{\pm} = (m_{\ell\ell}/\sqrt{s})e^{\pm y_{\ell\ell}}$, where + (-) corresponds to the proton on side A(C).

AFP cutflow

	Number of events	
Requirement	$pp \rightarrow p(\gamma\gamma \rightarrow ee)p$	$pp \to p(\gamma\gamma \to \mu\mu)p$
$\sigma imes \mathcal{L}$	44790	44740
$\sigma imes \mathcal{L} imes \epsilon_{ ext{filter}}$	11570	11560
$\sigma imes \mathcal{L} imes \epsilon_{ ext{filter}} imes w_{ ext{SF}}$	11440	11190
Exactly two signal leptons	1217	3628
Trigger matched	968	2641
Opposite charge	964	2641
Same flavor	964	2641
$p_{\rm T}^{\ell\ell} < 5~GeV$	931	2594
$A_{\phi}^{\ell\ell} < 0.01$	913	2520
$N_{\rm tracks}^{0.5 \rm mm} = 0$	378	1138
$m_{\ell\ell} > 20~GeV$	378	1138
$m_{\ell\ell} \not\in [70, 105] \; GeV$	283	960
$\xi_{\ell\ell}^{A} \in [0.02, 0.12] \text{ or } \xi_{\ell\ell}^{C} \in [0.02, 0.12]$	69.8	155
$\xi_{\ell\ell}^{\rm A} \in [0.035, 0.08] \text{ or } \xi_{\ell\ell}^{\rm C} \in [0.035, 0.08]$	18.2	28.9
$ \xi_{\rm AFP} - \xi_{\ell\ell} < 0.005$	17.8	27.8

AFP matched event kinematics

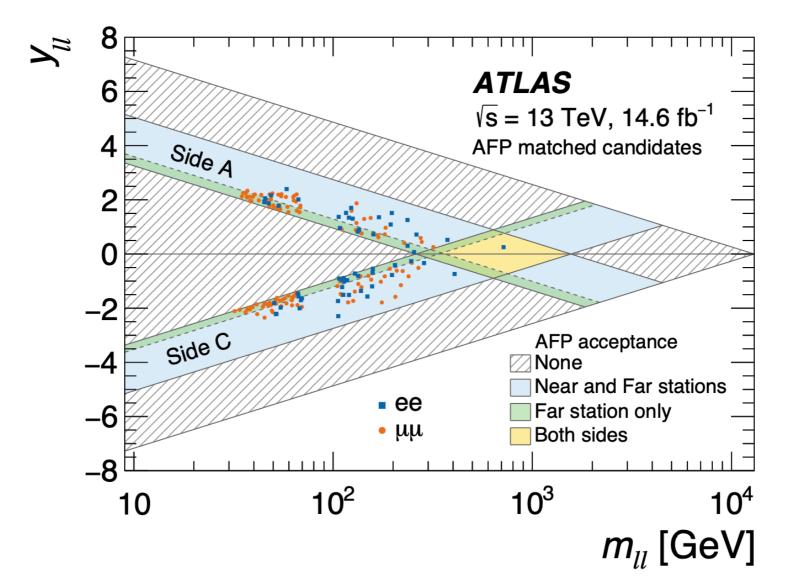


Figure 2: The 57 (123) *ee* ($\mu\mu$) data event candidates in the dilepton rapidity $y_{\ell\ell}$ vs $m_{\ell\ell}$ plane satisfying event selection and kinematic matching, $|\xi_{AFP} - \xi_{\ell\ell}| < 0.005$, on at least one side. Shaded (hatched) areas denote the acceptance (no acceptance) for the AFP stations indicated in the legend. Areas neither shaded nor hatched correspond to $\xi \notin [0, 1]$.

Probing masses up to \sim 700 GeV

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AFP matched event kinematics

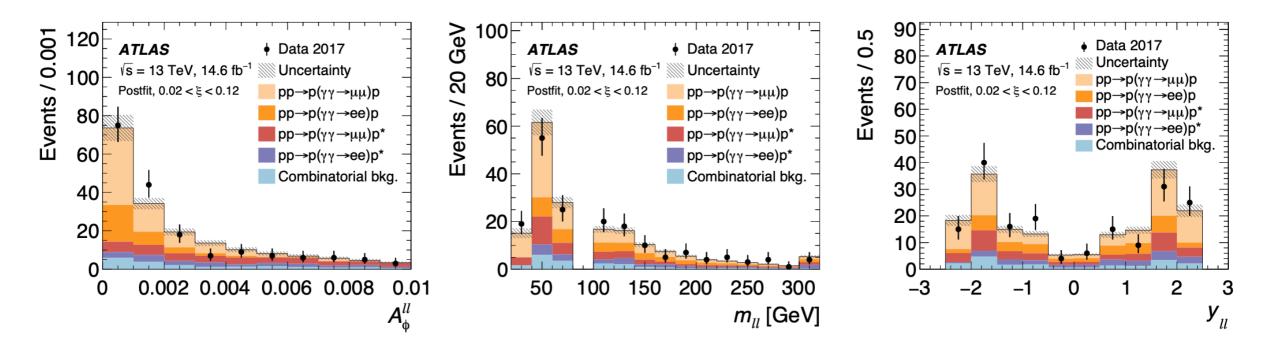


Figure 3: Distributions of dilepton acoplanarity $A_{\phi}^{\ell\ell}$ (left), invariant mass $m_{\ell\ell}$ (center), rapidity $y_{\ell\ell}$ (right) satisfying $\xi_{\ell\ell}, \xi_{AFP} \in [0.02, 0.12]$, and $|\xi_{AFP} - \xi_{\ell\ell}| < 0.005$ for at least one AFP side. Events with 70 < $m_{\ell\ell} < 105$ GeV are vetoed. The total prediction comprises the signal and combinatorial background processes, where p^* denotes a dissociated proton. The simulated predictions are normalized to data to illustrate the expected signal composition. The rightmost bin of the $m_{\ell\ell}$ distribution includes overflow. The hatched band indicates the combined statistical and systematic uncertainties of the prediction. Error bars denote statistical uncertainties of the data.

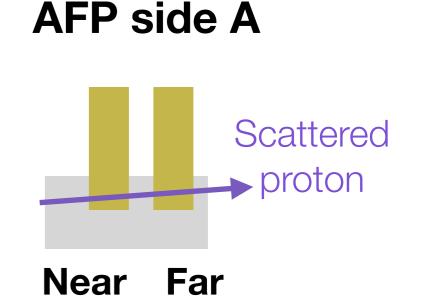
AFP Efficiency

Needed to correct for detector effects

Station tag and probe:

& vice versa

- 'Tag' track in near station (require exactly 1)
- Check if track in far (within $|x_{near}-x_{far}| < 2 \text{ mm}$)



probability [%] 102 A FAR (tag in A NEAR) **A NEAR** (tag in A FAR) — **C NEAR** (tag in C FAR) **C FAR** (tag in C NEAR) 100 98 96 94 **ATLAS**Preliminary 92 √s = 13 TeV (2017) 90 Oct Sep Aug Nov 88 4907 4907 34960 335016 335016 5222 6719 S ATLAS run number

Lower efficiency in far station due to showers

AFP cross-section break down

The measured fiducial cross sections in the *ee* and $\mu\mu$ channels are $\sigma_{ee+p}^{\text{fid.}} = 11.0 \pm 2.6 \text{ (stat)} \pm 1.2 \text{ (syst)} \pm 0.3 \text{ (lumi)}$ and $\sigma_{\mu\mu+p}^{\text{fid.}} = 7.2 \pm 1.6 \text{ (stat)} \pm 0.9 \text{ (syst)} \pm 0.2 \text{ (lumi)}$ fb, respectively. Table 1 compares these with the combined HERWIG and LPAIR predictions assuming unit soft-survival factors $S_{\text{surv}} = 1$. Soft-survival effects are included using an $m_{\ell\ell}$ -dependent reweighting of these predictions to S_{surv} calculated for exclusive processes from Ref. [34]; LPAIR predictions are additionally scaled down by 15% to account for S_{surv} being lower for single-dissociative processes [33]. SUPERCHIC 4 [97] predictions include full kinematic dependence on S_{surv} for exclusive, single-, and double-dissociative processes. The predictions for *ee* are higher than for $\mu\mu$ due to the looser $\eta(e)$ requirement [94].

Uncertainties of 7% (17%) are assigned for predictions of the exclusive (single-dissociative) processes [98].

AFP cross-section break down

	$\sigma_{ee+p}^{\text{fid.}}$ [fb]	$\sigma^{\rm fid.}_{\mu\mu+p}$ [fb]
Measurement	11.0 ± 2.9	7.2 ± 1.8
Predictions		
$S_{ m surv} = 1$		
Herwig+Lpair	15.5 ± 1.2	13.5 ± 1.1
HERWIG	9.3 ± 0.7	8.0 ± 0.6
LPAIR	6.2 ± 1.1	5.5 ± 0.9
$S_{\rm surv}$ using Refs. [31,30]		
Herwig+Lpair	10.9 ± 0.8	9.2 ± 0.7
Herwig	7.0 ± 0.5	5.9 ± 0.4
LPAIR	3.9 ± 0.7	3.4 ± 0.6
SuperChic 4 [94]		
Exclusive + single-dissociative	12.2 ± 0.9	10.4 ± 0.7
Exclusive	8.6 ± 0.6	7.3 ± 0.5
Single-dissociative	3.6 ± 0.6	3.1 ± 0.5

Uncertainties

Source of systematic uncertainty	
Forward detector	
Global alignment	6%
Beam optics	5%
Resolution and kinematic matching	3-5%
Track reconstruction efficiency	3%
Alignment rotation	1%
Clustering and track-finding procedure	< 1%
Central detector	
Track veto efficiency	5%
Pileup modeling	2-3%
Muon scale and resolution	3%
Muon trigger, isolation, reconstruction efficiencies	
Electron trigger, isolation, reconstruction efficiencies	
Electron scale and resolution	
Background modeling	
Luminosity	

Table 4: Summary of sources of systematic uncertainty and their impact on the cross-section measurement.

AFP background rejection

Powerful background suppression

Can see photo-production of di-leptons above Drell-Yan background

