

# Quenching features of quark and gluon initiated parton cascades in expanding media

Based on : JHEP 07 (2020) 150;  
arXiv: 2106.02592 (2021)

Souvik Priyam Adhya

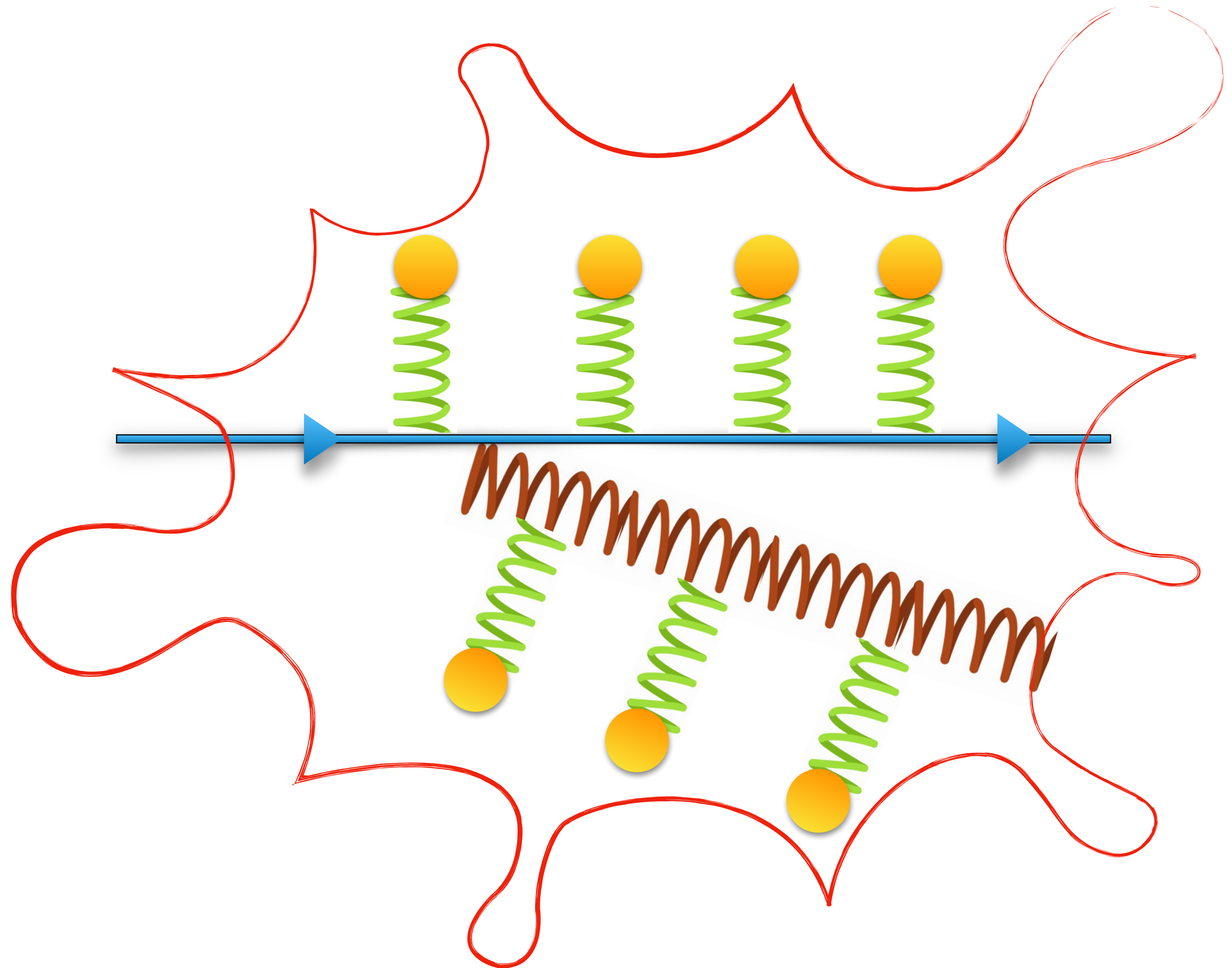
[in collaboration with K. Tywoniuk<sup>1</sup>, M. Spousta<sup>2</sup> and C. Salgado<sup>3</sup>]

Institute of Nuclear Physics, Polish Academy of Sciences

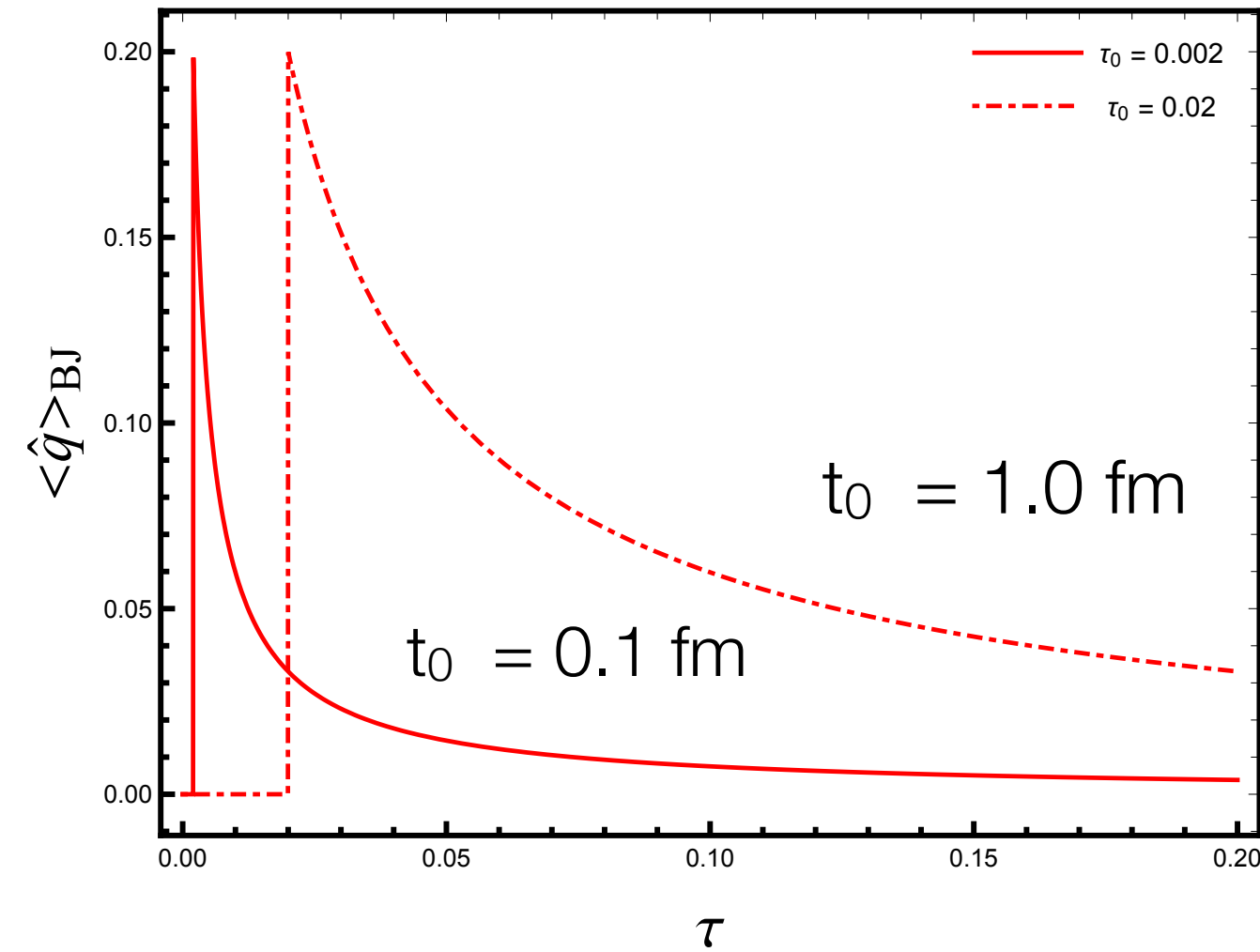
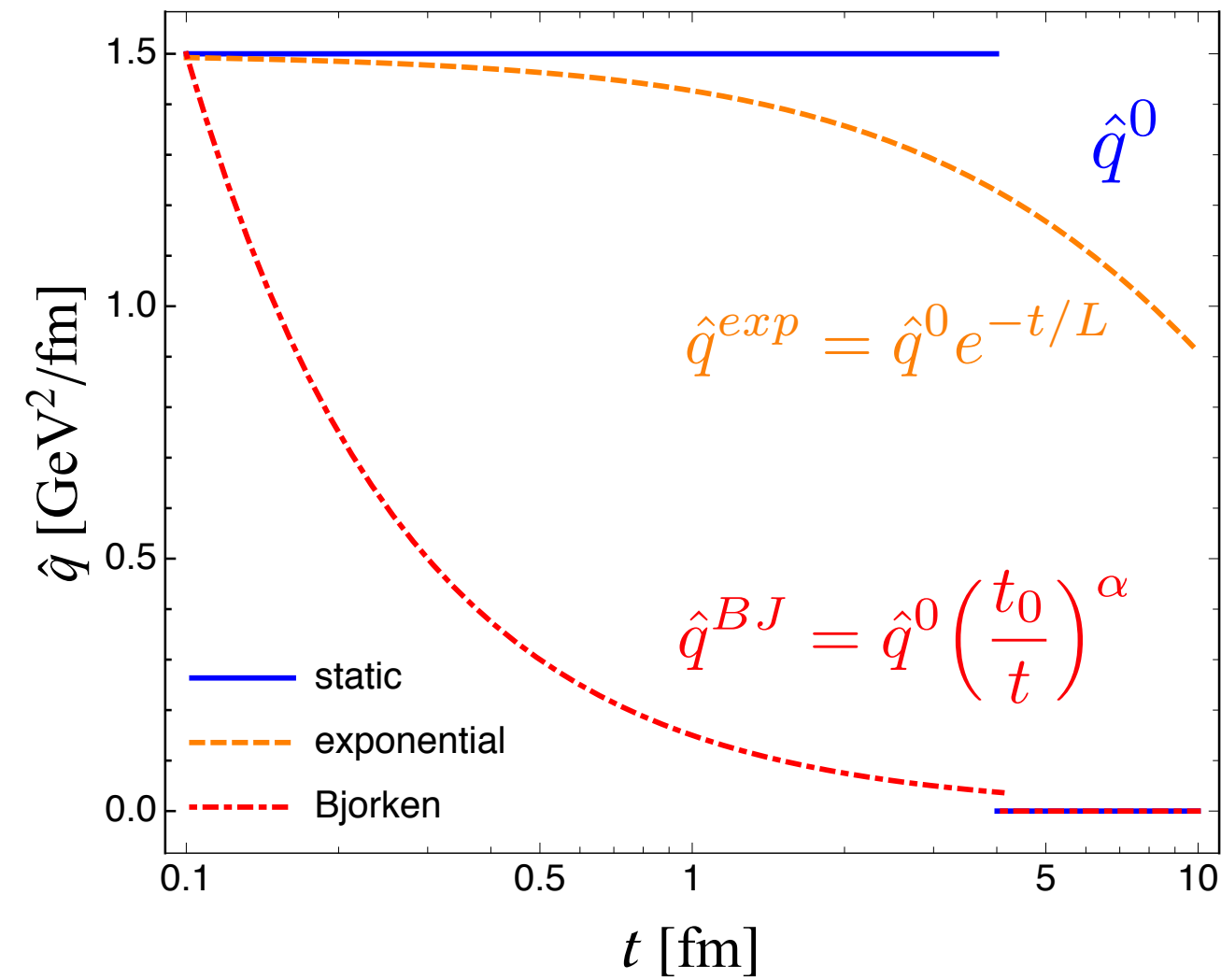
1. University of Bergen, Bergen.
2. Charles University, Prague.
3. IGFAE, Santiago de Compostela.

# Jet formation

- High energy partons, resulting from an initial hard scattering, will create a high energy collimated spray of particles → **JETS**
- Partons traveling through a dense color medium are expected to lose energy via **medium induced gluon radiation**, "jet quenching". We have adopted the *BDMPS-Z* (Baier, Dokshitzer, Mueller, Peigné, Schiff; Zakharov) formalism (*multiple soft scattering in medium*).

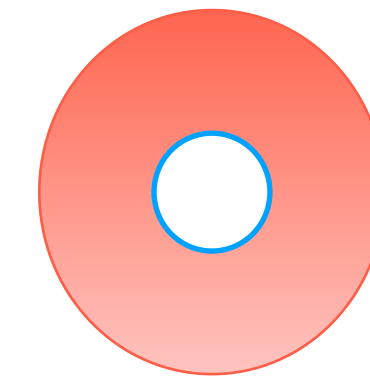


# Medium profiles and calculation workflow :

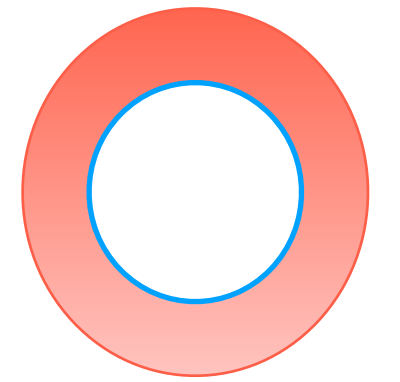


Early quenching

Late quenching



$t_0 = 0.1$  fm



$t_0 = 1.0$  fm

Bjorken initial conditions

Evolution equations =>

$$\star \frac{\partial}{\partial \tau} D_g(x, \tau) = \int_0^1 dz K_{gg} \left[ \sqrt{\frac{z}{x}} D_g\left(\frac{x}{z}\right) - \frac{z}{\sqrt{x}} D_g(x) \right] - \int_0^1 z K_{qg}(z) \frac{z}{\sqrt{x}} D_g(x) + \int_0^1 z K_{gq}(z) \sqrt{\frac{z}{x}} D_S\left(\frac{x}{z}\right)$$

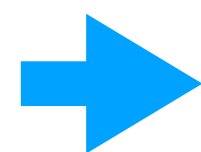
Calculation flowchart

$$\star \frac{\partial}{\partial \tau} D_S(x, \tau) = \int_0^1 dz K_{qq}(z) \left[ \sqrt{\frac{z}{x}} D_S\left(\frac{x}{z}\right) - \frac{1}{\sqrt{x}} D_S(x) \right] + \int_0^1 dz K_{qg}(z) \sqrt{\frac{z}{x}} D_g\left(\frac{x}{z}\right)$$

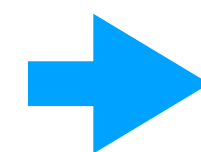
$D_S$  = q singlet spectra  
 $D_g$  = gluon spectra  
 $K$  = splitting rate  
 $\tau$  = evolution variable

S. S and Y. M-T ; JHEP 09 (2018) 144.

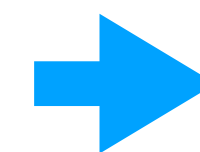
Single parton emission spectra (D) in BDMPS-Z formalism for static, exponential and Bjorken expanding media



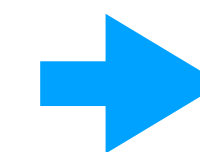
Splitting rates in static, exponential and Bjorken expanding media



Kinematic rate equation taking into account all the possible splittings for quark & gluon initiated jets



Optimisation in the Quenching factor for jets with combined q and g fractions through modified power law



Study of rapidity dependence and estimation of elliptic flow



# Scaling behaviour of the spectrum

The single gluon emission spectra are given as :

$$\frac{dI^{static,soft}}{dz} \simeq \frac{\alpha_s P(z)}{\pi} \sqrt{\frac{\omega_c}{2\omega}}$$

$$\frac{dI^{static}}{dz} = \frac{\alpha_s}{\pi} P(z) \operatorname{Re} \ln[\cos(\Omega_0 L)]$$

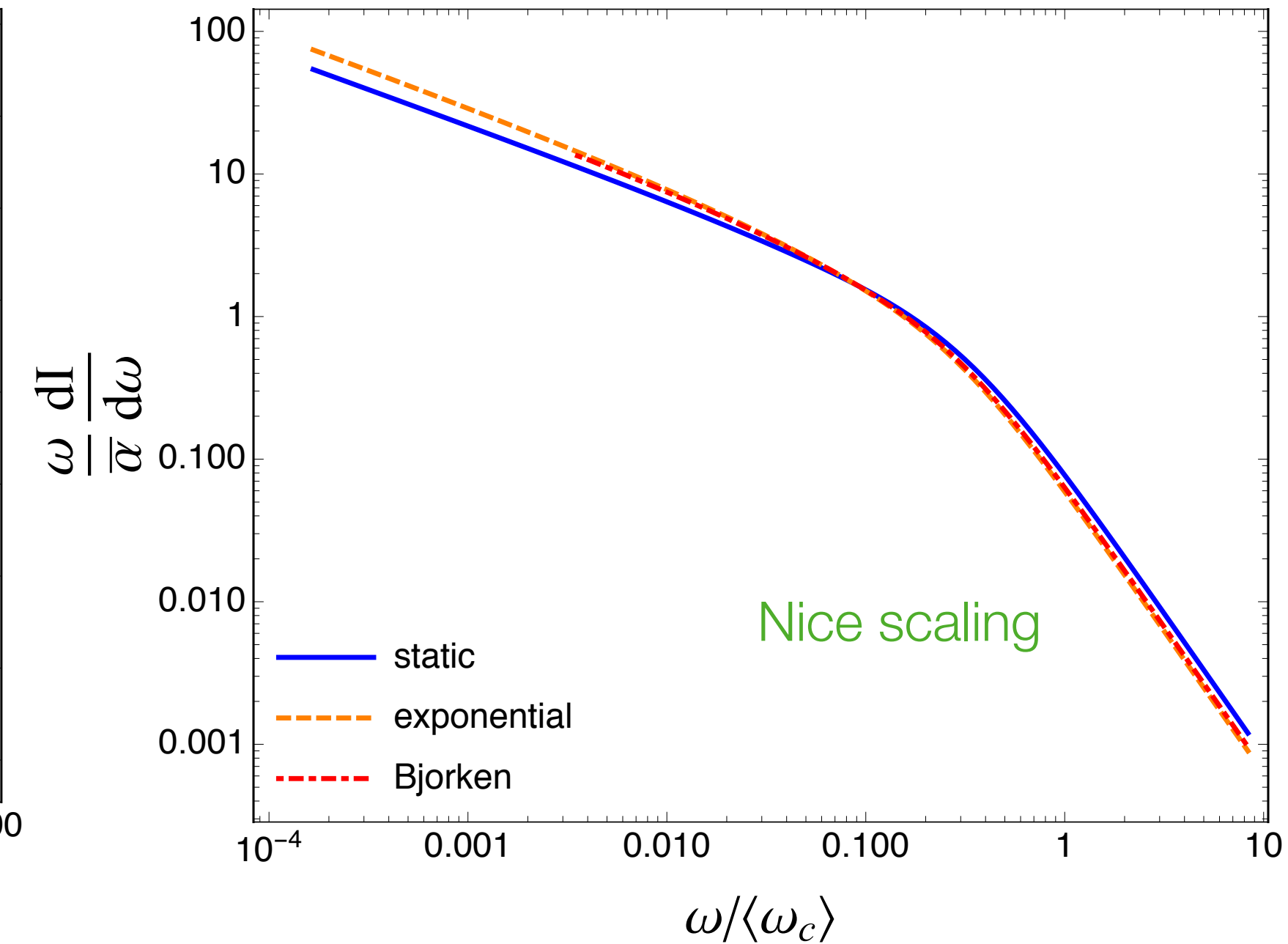
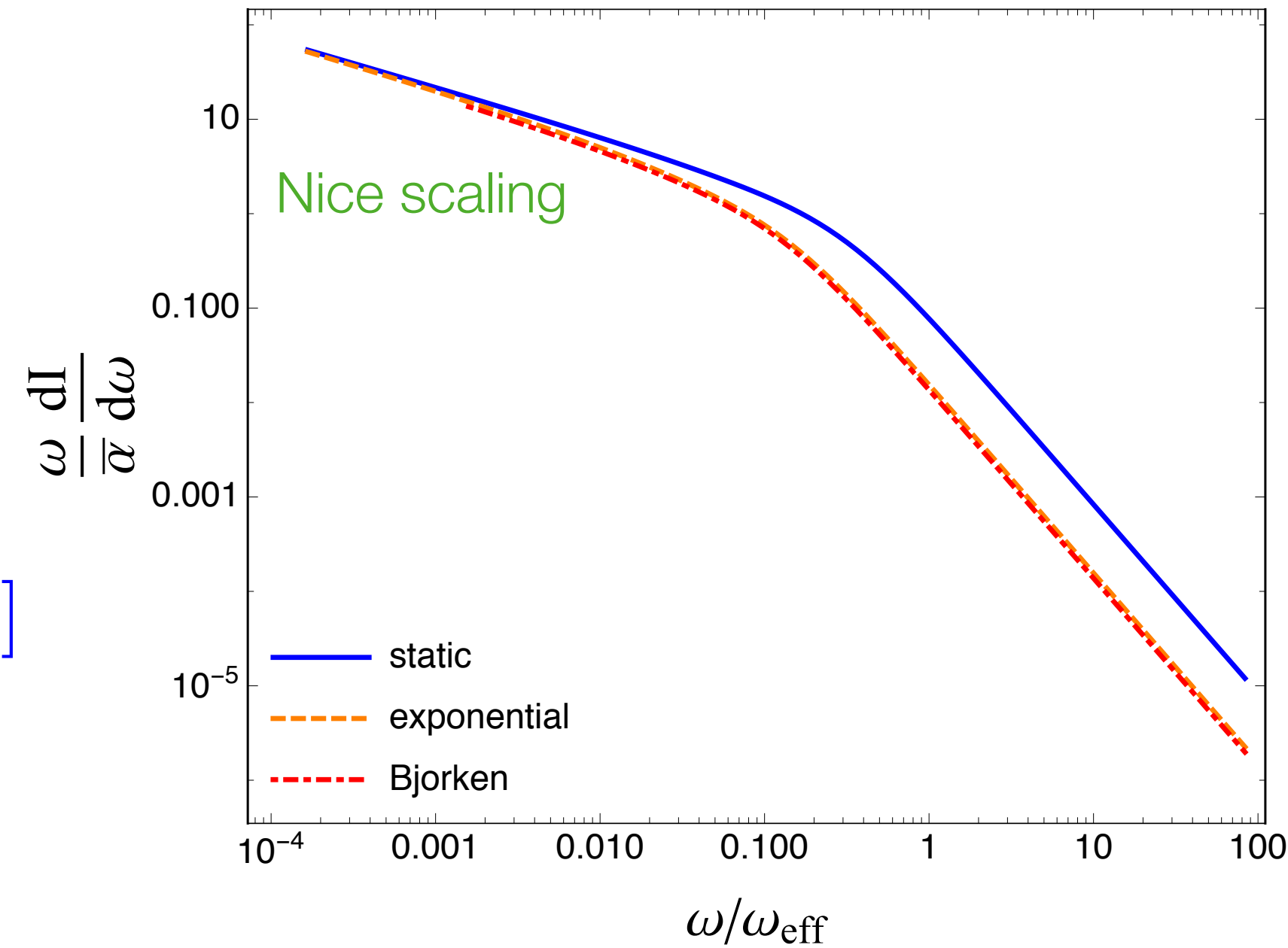
$$\frac{dI^{expo}}{dz} = \frac{\alpha_s}{\pi} P(z) \operatorname{Re} \ln J_0(2\Omega_0 L)$$

$$\frac{dI^{BJ}}{dz} = \frac{\alpha_s}{\pi} P(z) \operatorname{Re} \ln \left[ \left( \frac{t_0}{L+t_0} \right)^{1/2} \frac{J_1(z_0)Y_0(z_L) - Y_1(z_0)J_0(z_L)}{J_1(z_L)Y_0(z_L) - Y_1(z_L)J_0(z_L)} \right]$$

$$P_{gg} = 2C_A \frac{(1-z(1-z))^2}{z(1-z)} \quad \tau \equiv \sqrt{\frac{\hat{q}_0}{p}} L \quad z_0 \equiv (1-i)\kappa(z)\tau_0$$

$$z_L \equiv (1-i)\kappa(z)\sqrt{\tau_0(\tau+\tau_0)},$$

Can we interpret the scalings in different kinematical limits ?



Effective parameter

$$\frac{dI^{static,sing}}{dz} \simeq \frac{dI^{expo,sing}}{dz} \simeq \frac{dI^{BJ,sing}}{dz}$$

The singular spectra can be re-scaled

$$\omega_{\text{eff}} = \begin{cases} \frac{1}{2}\hat{q}_0 L^2 & \text{static medium} \\ 2\hat{q}_0 L^2 & \text{exponentially expansion} \\ 2\hat{q}_0 t_0 L & \text{Bjorken expansion} \end{cases}$$

$$\hat{q}_{\text{eff}}^{expo} = 4\hat{q}_0$$

$$\hat{q}_{\text{eff}}^{BJ} = 4\hat{q}_0 t_0 / L$$

# Medium evolved gluon spectra

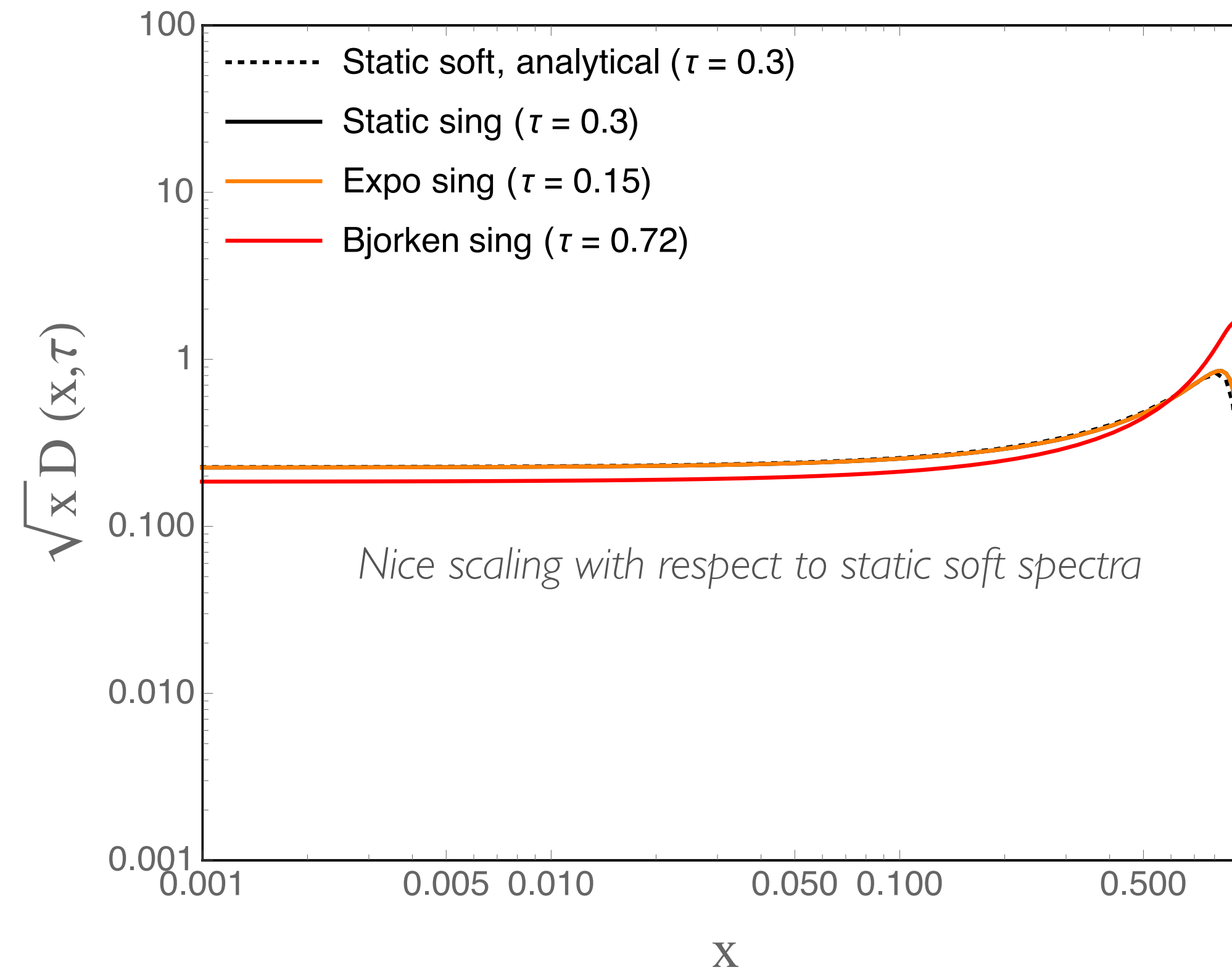
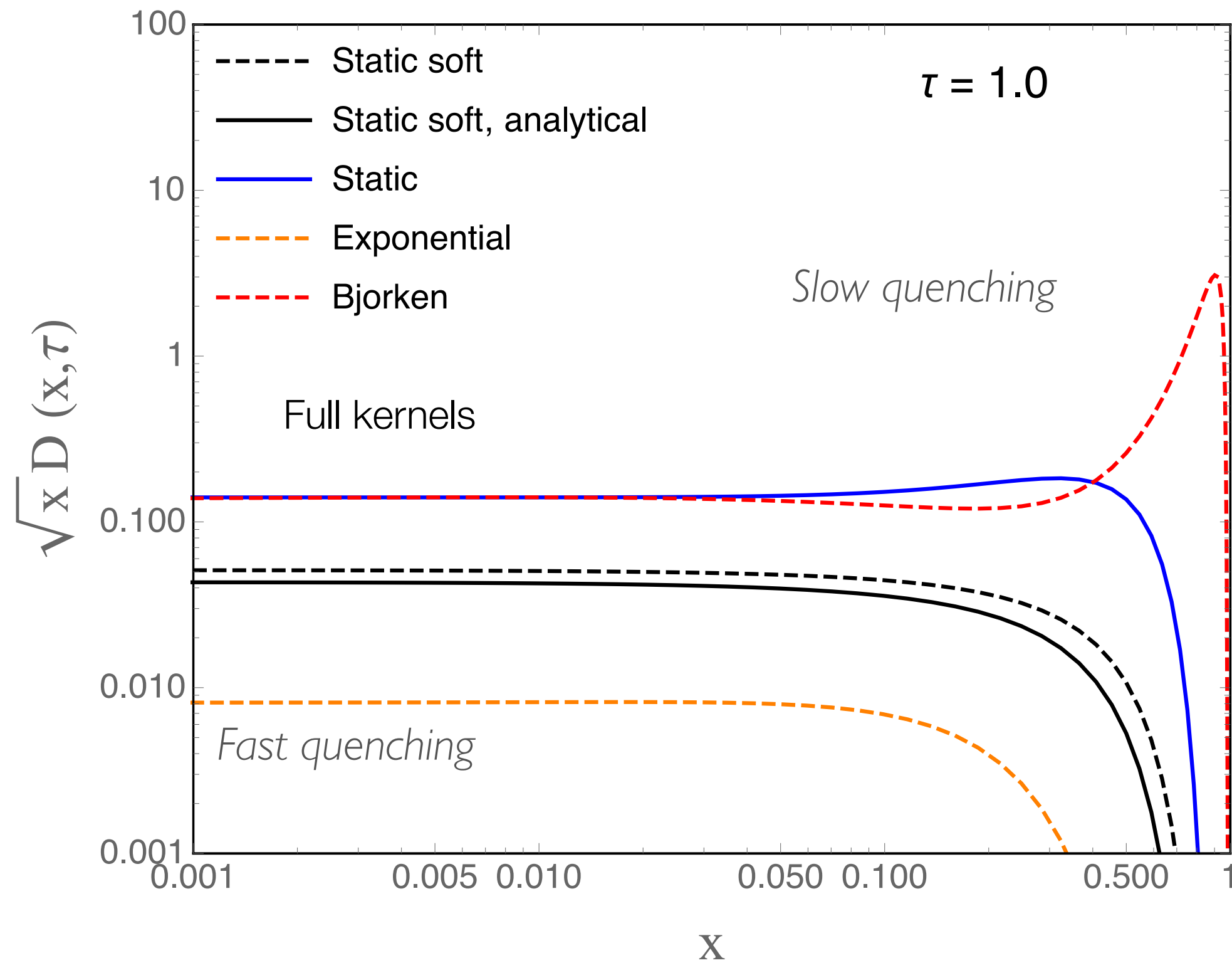
J.-P. Blaizot, E. Iancu, and Y. Mehtar-Tani, Phys.Rev.Lett. 111 (2013) 052001.

- The kinematic evolution equation (GAIN + LOSS terms) in terms of gluon spectra :

$$\frac{\partial D(x, t)}{\partial \tau} = \int dz \mathcal{K}(z, \tau | p) \left[ \sqrt{\frac{z}{x}} D\left(\frac{x}{z}, \tau\right) - \frac{z}{\sqrt{x}} D(x, \tau) \right]$$

Static, soft (analytical) gluon spectra

$$D(x, \tau) = \frac{\tau}{\sqrt{x}(1-x)^{3/2}} e^{-\pi \frac{\tau^2}{1-x}}$$

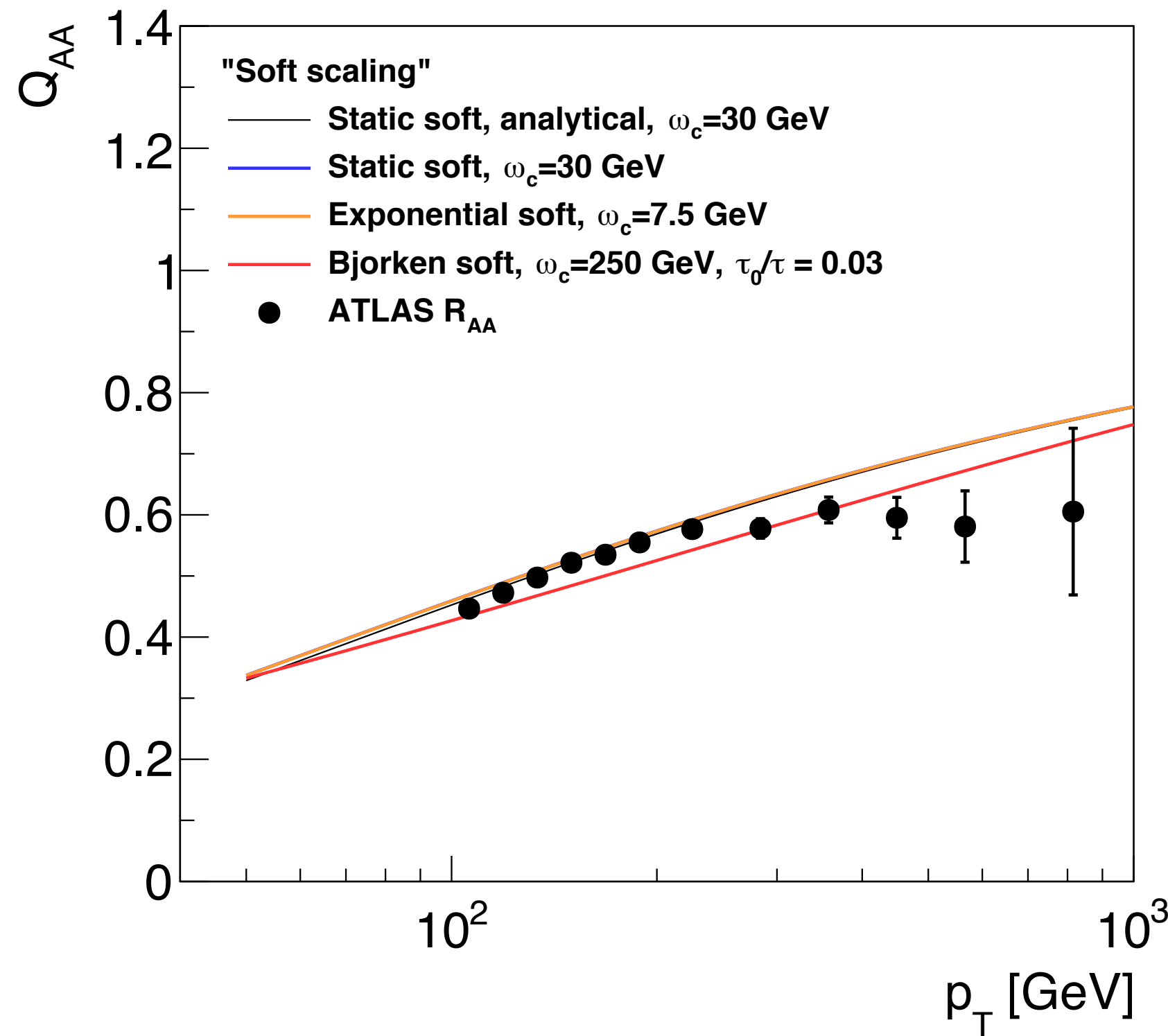


- Gluon spectra:
  - Singular rates ==> Nice scaling in  $\tau_{\text{eff}}$ .
  - Full rates ==> No scaling in  $\tau_{\text{eff}}$ .

See  $k_T$  broadening of cascades : M. Rohmoser talk

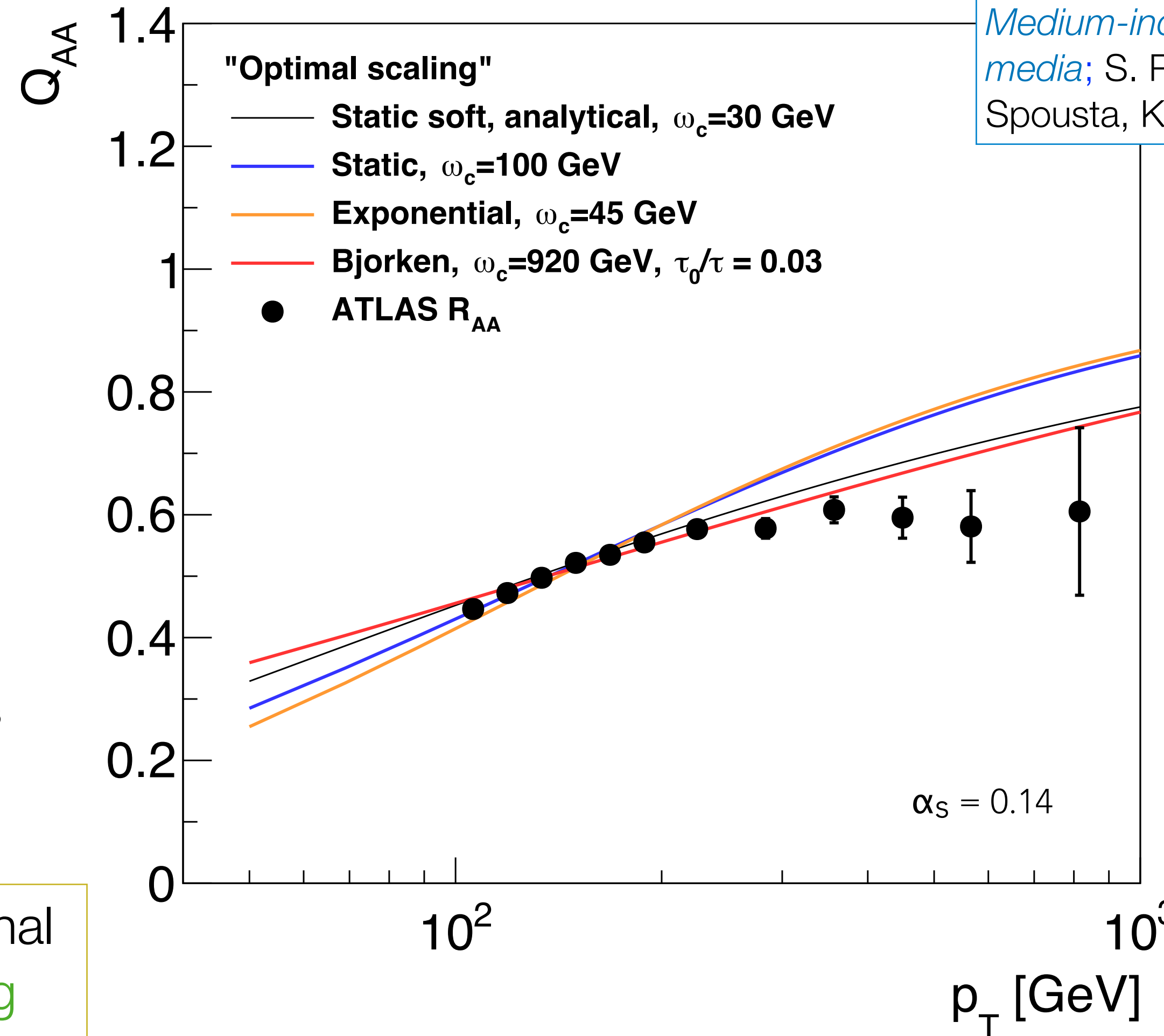
- At low  $x$ , we see a  $1/(\sqrt{x})$  behaviour of all the profiles >> recovered from the similar gluon splitting at low  $x$ .

# Is it possible to re-scale $Q_{AA}$ for different medium ?



The Bjorken profile depends on additional choice of  $(\tau_0/\tau)$  : **No universal scaling**

- Good, but not perfect scaling is achieved by optimisation.
- Scaling for expo medium ~ average scaling.



Medium-induced cascade in expanding media; S. P. Adhya, C. A. Salgado, M. Spousta, K. Tywoniuk; JHEP 07 (2020) 150.

- Significant differences in values of the quenching parameter for different types of medium and kinematical ranges point to the importance of precise modelling of the jet quenching phenomenon !

But we did for **GLUONIC** cascades only. Next, we go for **multi-partonic cascades** ...

# Does the media behave differently as a function of rapidity ?

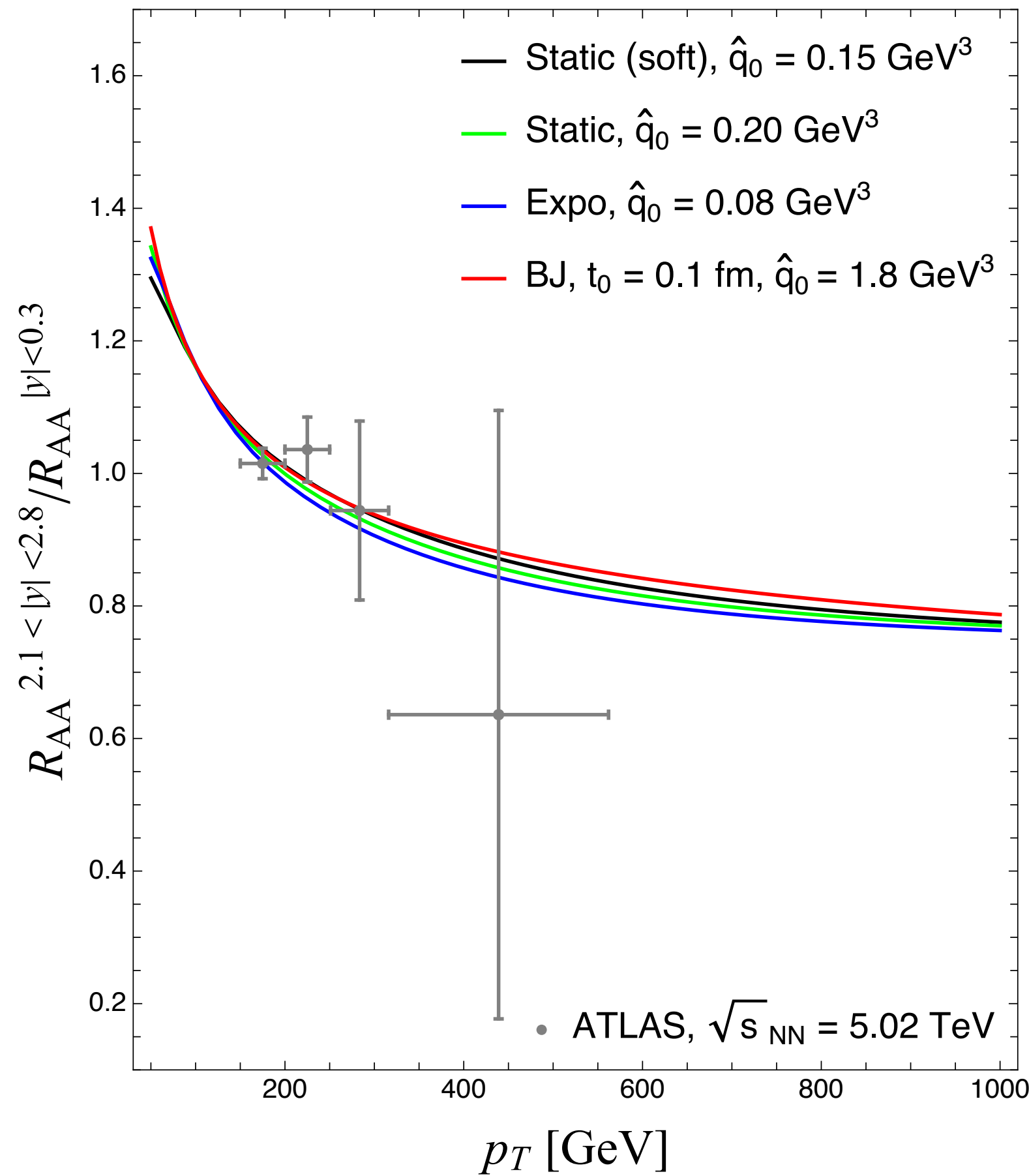
S. P. Adhya., C. A. S., M. S., K. T.  
[arXiv:2106.02592](https://arxiv.org/abs/2106.02592)  
 (2021)

**Multipartonic cascades**

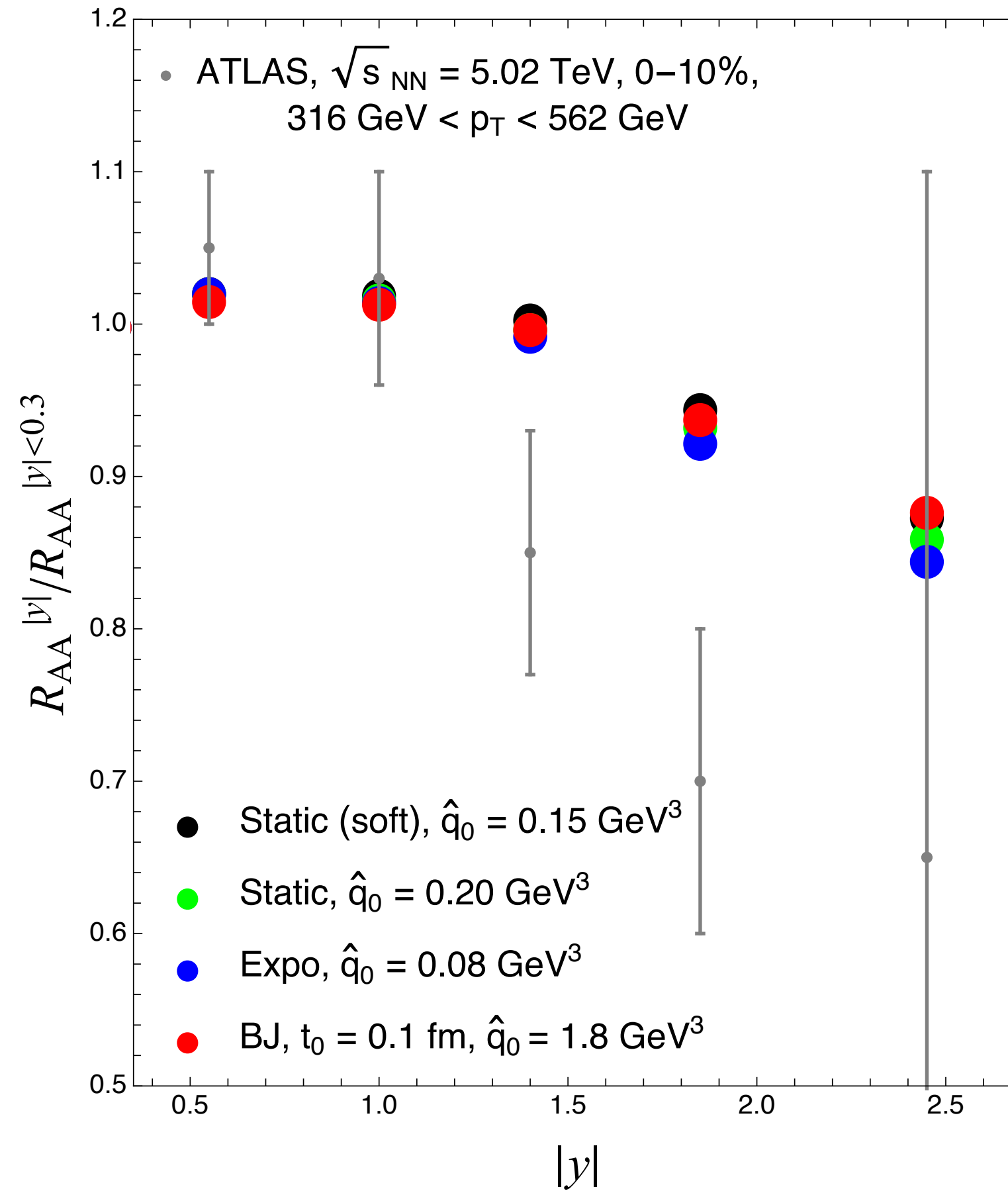
No

No

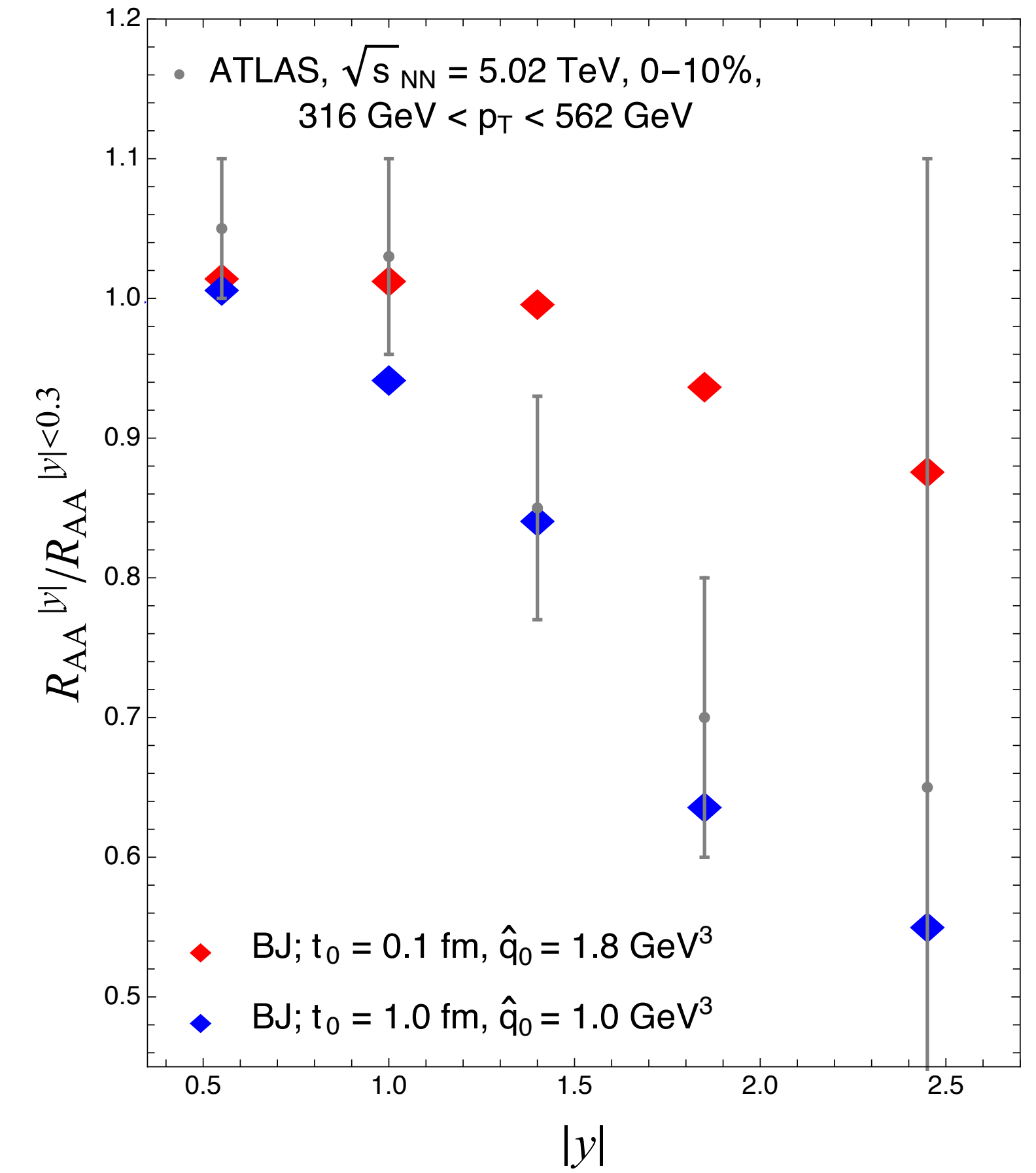
Yes



Rapidity ratio for different medium profiles wrt  $p_T$



Rapidity ratio for different medium profiles wrt  $|y|$

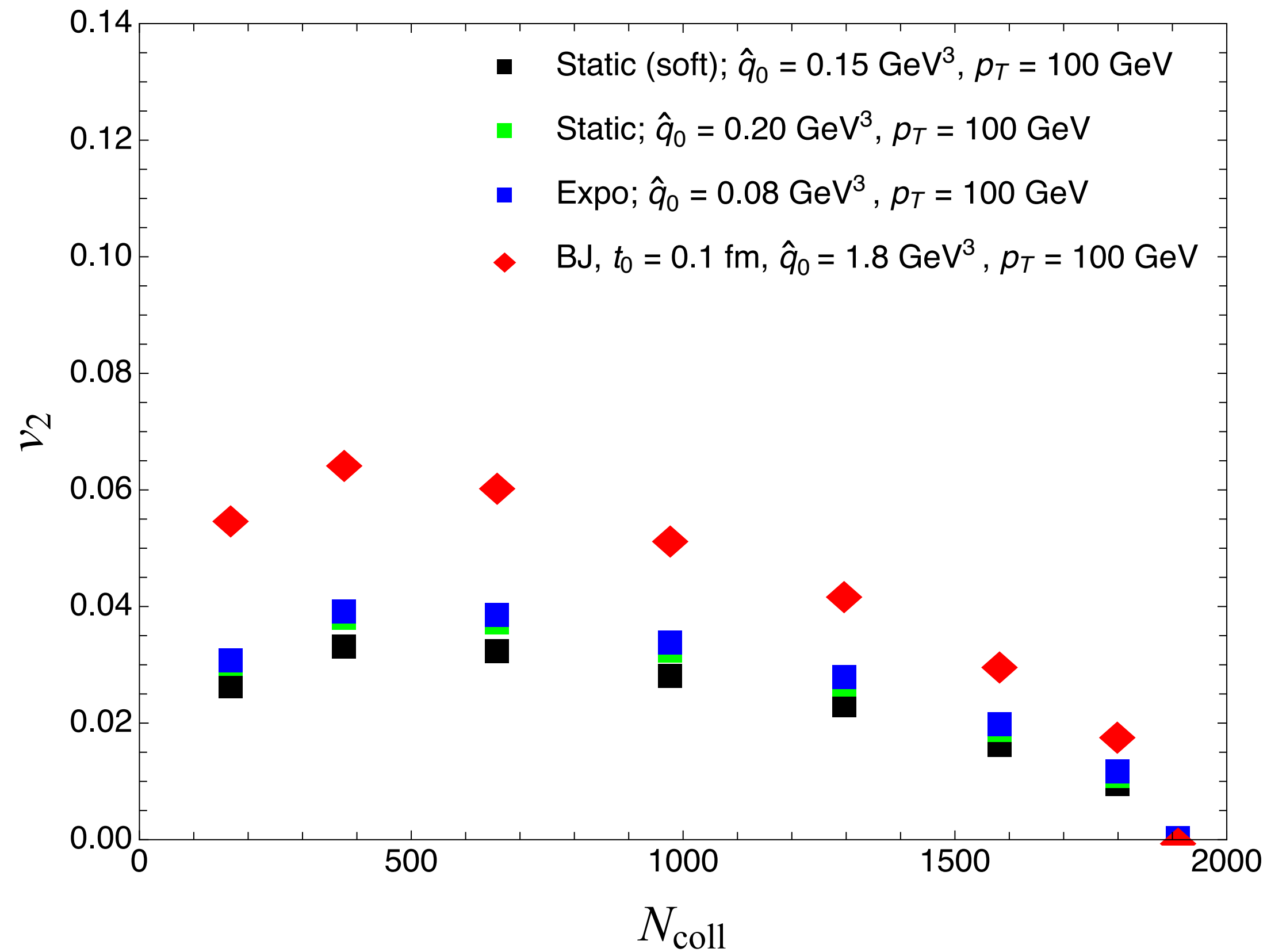


Rapidity ratio for different start of quenching time for the Bjorken

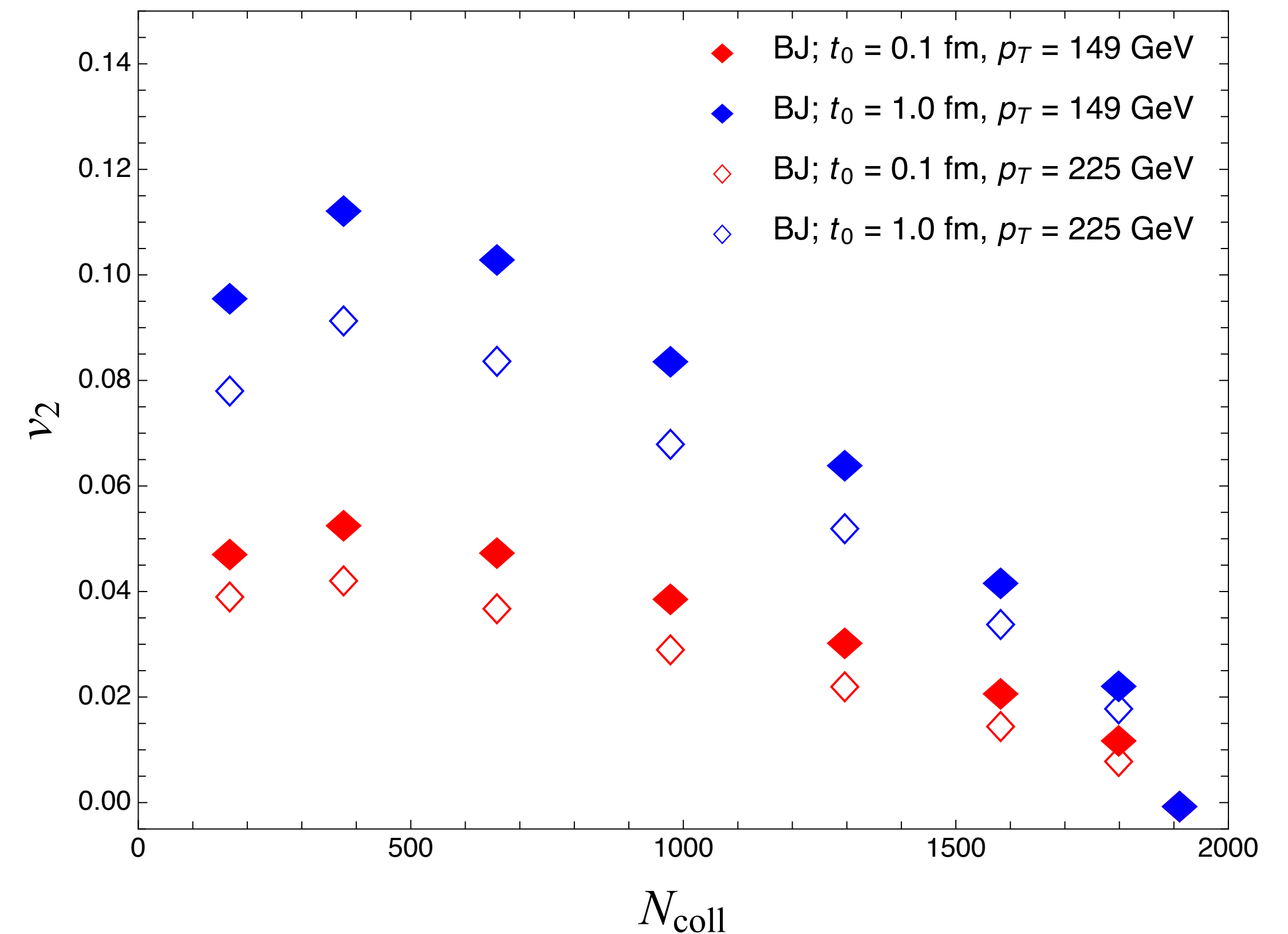


# Does the media behave differently with respect to $v_2$ ?

Yes and no



Yes



$$v_2 = \frac{1}{2} \frac{Q_{AA}(L^{\text{in}}) - Q_{AA}(L^{\text{out}})}{Q_{AA}(L^{\text{in}}) + Q_{AA}(L^{\text{out}})}$$

- The impact of the medium expansion can be largely scaled out by a **suitable choice of  $\hat{q}$**  [confirming Adhya et. al., 2020].
- The **jet  $v_2$  remains sensitive to choice of  $t_0$** .

S. P. A., C. A. S., M. S., K. T.  
arXiv:2106.02592 (2021)

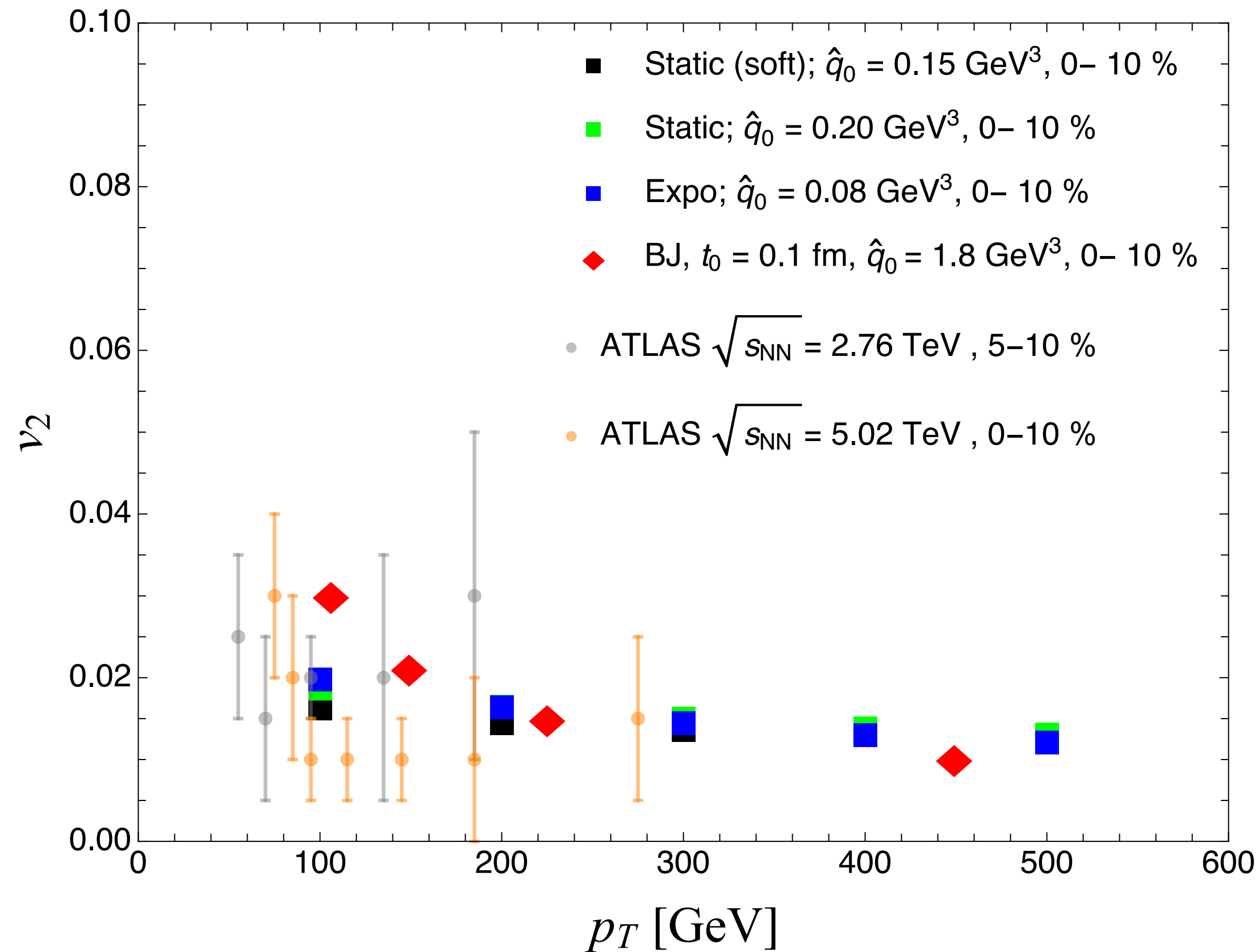


# Does the media behave differently with respect to $v_2$ ?

S. P. A., C. A. S., M. S., K. T.  
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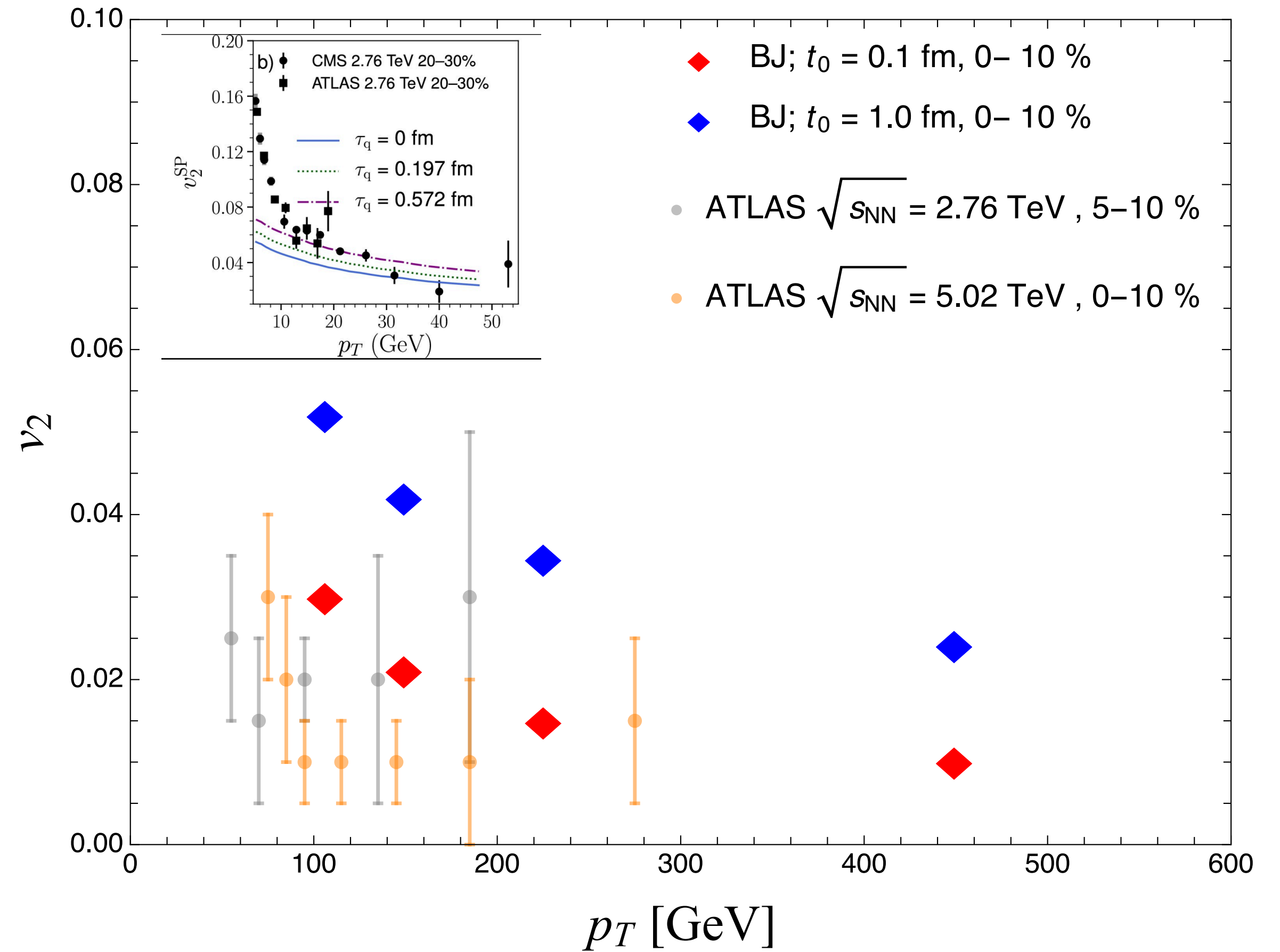
Yes and no

Elliptic flow  $v_2$  as a function of  $p_T$  for different media



Yes

Elliptic flow  $v_2$  as a function of  $p_T$  for Bjorken initial conditions



- Agreement with findings of the sensitivity of  $v_2$  on  $t_0$  [Carlota et. al., PLB, 2020 (inset)] which was done in more complex modelling of the collision geometry, but less complex modelling of the medium induced showering.

Thanks!

BACKUP

# ATLAS measurements

