INSTITUTO DE FÍSICA DE PARTÍCULAS Y DEL COSMOS



## EPS-HEP CONFERENCE 2021



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### 1. Introduction

- Higgs couplings to gauge bosons and top quark are still compatible with the SM with deviations of O (10%). For other fermions (e.g bottom) and the triple-Higgs coupling larger deviations are not excluded .[1]
- These deviations may come **from strongly interacting new physics**, where the Higgs boson and the Goldstone Bosons are composite states.
- We will focus on heavy fermion loop corrections to VBS (imaginary part) with top quark because of its large mass, 175 GeV. Fermion corrections are often neglected because the bosons ones dominate at high energy. (~ 3 TeV)

#### But how important are fermion loops?

The imaginary parts enter in the NLO counting

Is it possible to find values for the modified couplings that lead to a significant contribution?

[1] Handbook of LHC Higgs Cross Sections: 4. - LHC Higgs Cross SectionWorking Group

#### 2.Electroweak Chiral Lagrangian (EFT)

• Electroweak Chiral Lagrangian : EW GB transform non-linearly and a Higgslike field which transforms linearly under  $SU(2)_L xSU(2)_R$  which breaks to the Custodial Symmetry  $SU(2)_{L+R}$ .

 $SU(2)_L \times SU(2)_R \xrightarrow{SSB} SU(2)_{L+R}$ 

• Systematic expansion in chiral power counting (different to the SMEFT canonical expansion). Renormalizable order by order.

$$\mathscr{L}_{EChL} = \mathscr{L}_2 + \mathscr{L}_4 + \dots$$

• It is often used the Equivalence Theorem [2], where we relate the gauge bosons with the would-be-Goldstones at high energies.

$$\mathcal{A}(W_L^a W_L^b o W_L^c W_L^d) = \mathcal{A}(\omega^a \omega^b o \omega^c \omega^d) + O\left(\frac{M_W}{\sqrt{s}}\right)$$

• Because of exact cancellations of some amplitudes we need go beyond the ET.

[2] P.B. Pal, What is the equivalence theorem really? (1994)

The lagrangian at lowest order (chiral dimension 2)

$$\mathscr{L}_{2} = \frac{v^{2}}{4} \mathscr{F}(h) \operatorname{Tr}\left[\left(D_{\mu}U\right)^{\dagger} D^{\mu}U\right] + \frac{1}{2} \partial_{\mu}h \partial^{\mu}h - V(h) + i \bar{Q} \partial Q - v\mathscr{G}(h) \left[\bar{Q}'_{L}UH_{Q}Q'_{R} + \text{h.c.}\right]$$

GB + h + Yukawa sector

Just the top for this case

Spherical parametrization

$$U = \sqrt{1 - \frac{\omega^2}{v^2}} + i\frac{\bar{\omega}}{v}$$

$$Q^{(\prime)} = \begin{pmatrix} \mathscr{U}^{(\prime)} \\ \mathscr{D}^{(\prime)} \end{pmatrix} - \begin{bmatrix} \mathscr{U}^{\prime} = (u, c, t)' \\ \mathscr{D}^{\prime} = (d, s, b)' \\ \mathbb{Q}^{\prime} = (d, s, b$$

Analytic functions of powers of the Higgs field. Inspired by most of low energy HEFT models.

$$V(h) = v^{4} \sum_{n=3}^{\infty} V_{n} \left(\frac{h}{v}\right)^{n} \text{ for } V_{2} = V_{3} = \frac{M_{h}^{2}}{2v^{2}}, \quad V_{4} = \frac{M_{h}^{2}}{8v^{4}}, \quad V_{n>4} = 0 \quad \textcircled{Recover the SM}$$

$$\mathscr{F}(h) = 1 + 2a\frac{h}{v} + b\frac{h^{2}}{v^{2}} + \dots \quad \mathscr{G}(h) = 1 + c_{1}\frac{h}{v} + c_{2}\frac{h^{2}}{v^{2}} + \dots \quad a = b = 1$$

$$c_{1} = 1$$

$$c_{2} = c_{3} = \dots c_{n} = 0$$

### 3. Loops

We have calculated the contribution of top quark loops to VBS via the generating functional, obtaining the scattering for gauge bosons. Renormalized the relevant couplings and fields and compared to the existing literature [3].

We have obtained the real and imaginary part of the Partial Wave Amplitudes (PWA) or pseudo-PWA's.

#### But how important are fermion loops?

The imaginary parts enter in the NLO counting. In general the bosons dominate at high energy. ( $\sqrt{s} \sim 3$  TeV)

$$Im[Bosons] = Im[a_{J}]\Big|_{\gamma\gamma,\gamma Z,\gamma H,W^{+}W^{-},ZZ,ZH,HH}$$
$$Im[Fermions] = Im[a_{J}]\Big|_{t\bar{t},b\bar{b}}$$
$$R_{J} = \frac{Im[Fermions]}{Im[Boson] + Im[Fermions]}$$
$$R \sim 1 \rightarrow \text{Fermions dominate}$$
$$R \sim 0 \rightarrow \text{Bosons dominate}$$

We will inspect this ratio for the PWA of the process  $W^+W^- \rightarrow W^+W^-$ 

[3] G. Buchalla et al. LMU-ASC 13/20 [4] D. Espriu and J. Matías Phys. Rev. D 52, 6530

### Im[Bosons] depend on *a*, *b* and and *d*<sub>3</sub>

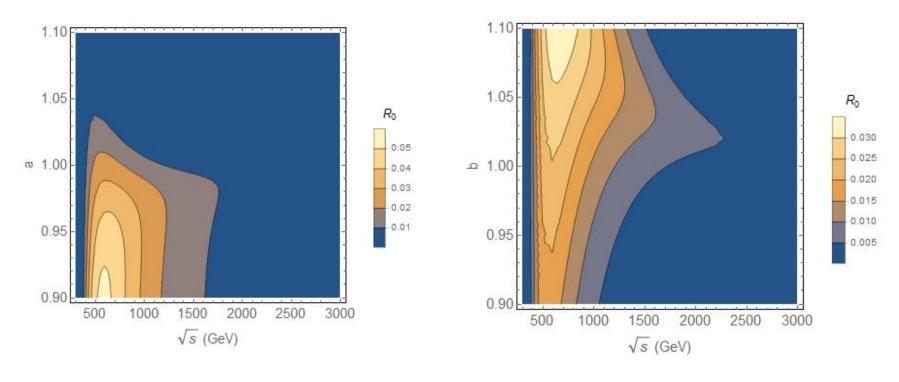
### Im[Fermions] depend on *a* and *c*<sub>1</sub>

# We will alow a 10% deviation from 1

[1] Handbook of LHC Higgs Cross Sections: 4. - LHC Higgs Cross SectionWorking Group

### 4. Results for $W^+W^- \rightarrow W^+W^-$

#### 4.1 Partial wave $a_0$ (J=0)



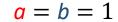
 $b = c_1 = 1$ 

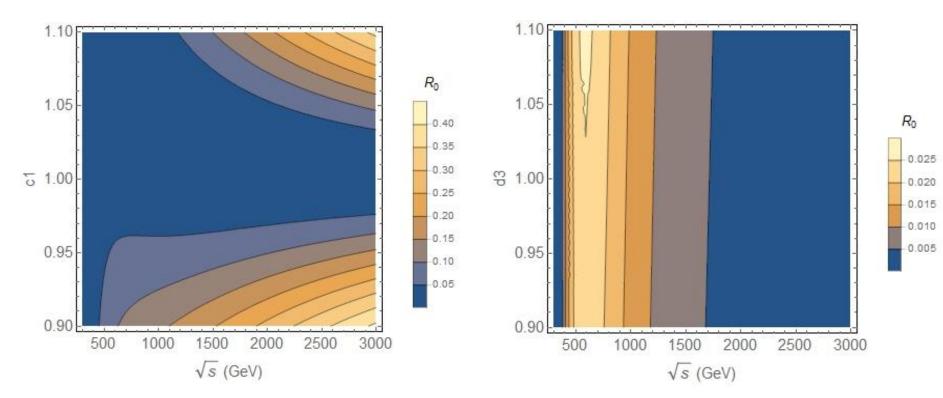
 $a = c_1 = 1$ 

5 % corretions at 500 GeV máximum for *a* around 0.9

3 % correction at 500 GeV for **b** around 1.1

Bosons completely dominate over 1 TeV for a and b

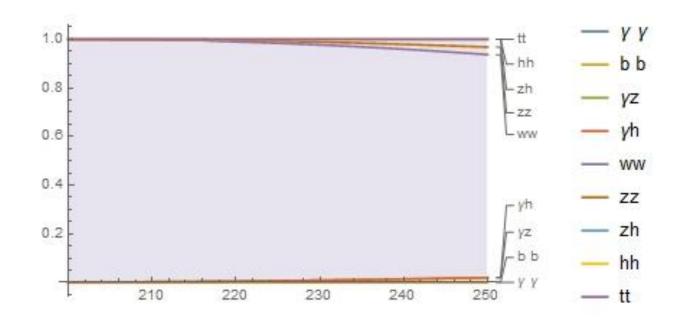




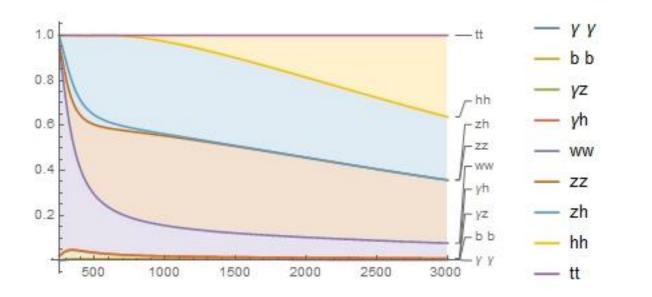
We find corrections of 40% at high energies around  $c_1$ =0.90 and  $c_1$ =1.1

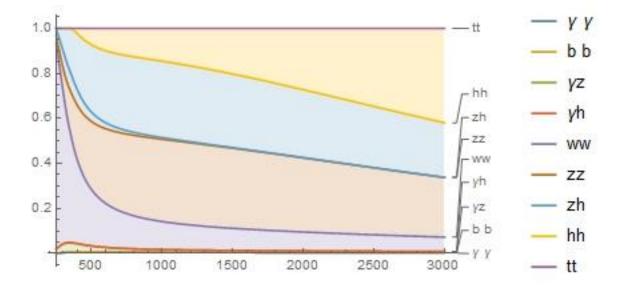
Again 2% corrections. Neglibible

# At the SM



Fermion contributions are negilible





#### Fermion contributions are indeed important

C1=0.9

C1=1.1

#### Parameter scan for $a_0$

We inspect *a*, *b*,  $c_1$  and  $d_3 \in [0.90, 1.10]$  [1]

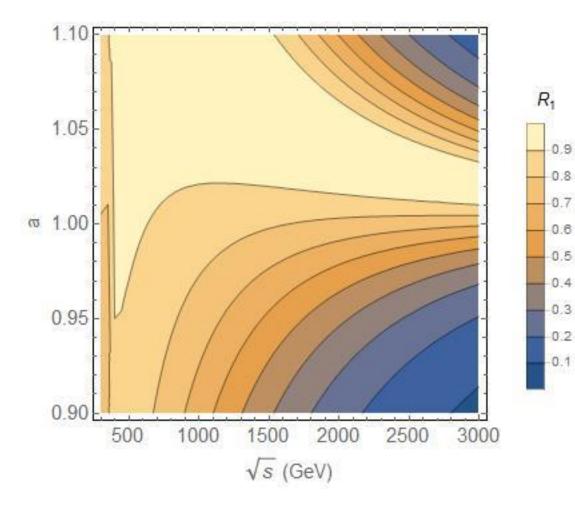
$\sqrt{s} (\text{TeV})$	a	b	$c_1$	$d_3$	$R_0$
1.50	1.00	1.00	0.90	1.10	0.42
3.00	1.00	1.05	0.90	0.90	0.30

Highest R

c<sub>1</sub> is the most important parameter for **J=0** 

[1] Handbook of LHC Higgs Cross Sections: 4. - LHC Higgs Cross Section Working Group

#### **<u>4.2 Partial wave a\_1 (J=1)</u>**



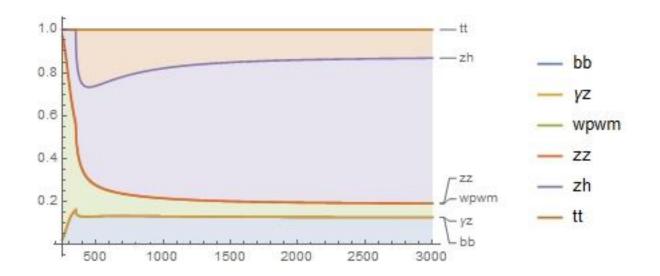
Im[Bosons] = f(a) 
$$\approx \left[\frac{(1-a^2)^2 s}{96 \pi v^2}\right]^2$$

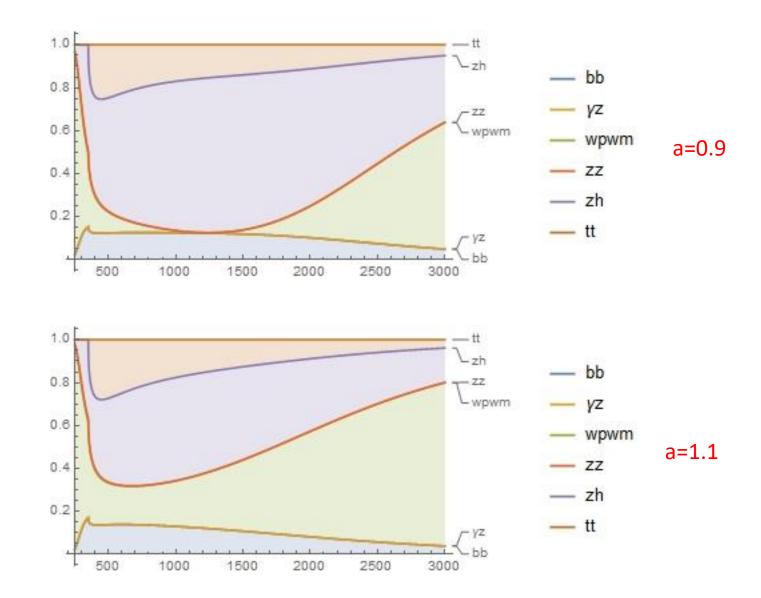
Im[Fermions]=*Im*[*Fermions*]<sub>*SM*</sub>

Does not depend on  $b, c_1$  or  $d_3$ , just a

High corrections for a close to 1

# At the SM





#### Fermion contribution are always important. Even close to the SM

#### 4. Specific Scenarios: Minimal Composite Higgs Model 0.07 0.06 0.05 $\xi = v^2/f^2$ f=0.56 TeV (a=0.9) 0.04 Å f=0.78 TeV (a=0.95) 0.03 $b^* = 1 - 2\xi$ f=1.7 TeV (a=0.99) 0.02 f=5.5 TeV (a=0.999) 0.01 $a^* = c_1^* = \sqrt{1-\xi}$ SM (a=1) 0.00 500 2500 1000 1500 2000 3000 √s (GeV) R1 is significally larger than RO 0.40 0.35 a1 more sensitive to 0.30 f=0.56 TeV (a=0.9) fermion corrections f=0.78 TeV (a=0.95) 0.20 f=1.7 TeV (a=0.99) f=5.5 TeV (a=0.999) 0.15 SM (a=1) 0.10 40% corrections at Low energy and for 2000 500 1000 1500 2500 3000 values close to SM s (GeV)

### 5. Conclusions

- We estimate fermion corrections to WW scattering: negligible in most of the parameter space in some cases but not always.
- For instance, the PWA's in the range considered:

- The MCHM shows R1 than R0 hence its more sensitive to the fermion corrections.
- Future work: considering the whole amplitude (real and imaginary) and unitarizing.

Thank you.