

Soft Gluon Resummation for the Associated Single Top and Higgs Production at the LHC

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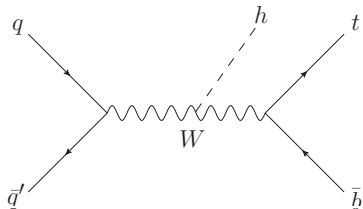
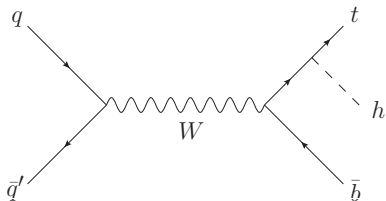


- ▶ Beyond-the-Standard Model (BSM) physics search
 - ▶ Precision calculations \Rightarrow find deviations between theoretical predictions and experimental data
 - ▶ Higgs and top quark processes are of particular interest: direct access to Yukawa coupling
- ▶ $pp \rightarrow Htj$: very small cross section, $g_{hWW}/g_{ht\bar{t}}$, very **sensitive** to new physics, currently being measured together with $t\bar{t}H$
[ATLAS Collaboration PRL.125.061802 (2015)] [CMS Collaboration PRD.99.092005 (2019)]

GOAL: extend the precision of theoretical predictions beyond the known NLO accuracy by resumming soft-gluon corrections to all-orders for the Htj production

FIRST STEP: s-channel (ca. 30% NLO QCD corrections)

- Tree-level Feynman diagrams in s-channel $pp \rightarrow Htj$



$$\sigma_{pp \rightarrow Htj}(\{m^2\}) = \sum_{i,k} \int dx_1 \int dx_2 \underbrace{f_i(x_1, \mu_F^2) f_k(x_2, \mu_F^2)}_{\text{parton distribution functions (PDFs)}} \times \underbrace{\hat{\sigma}_{ik \rightarrow Htj}(x_1, x_2, \alpha_s(\mu_R^2), \{m^2\}, \mu_F^2)}_{\text{partonic cross section}}$$

► Higher order contributions to σ

- Next-to-leading order (NLO) known [Demartin, Maltoni, Mawatari, Zaro (2015)]
- NLO QCD+EW corrections for all combined channels [Pagani, Tsinikos, Vryonidou (2020)]
- Full NNLO for Htj not available
 - aNNLO t-channel recently calculated [Forslund, Kidonakis (2020, 2021)]
- Soft gluon corrections: well-defined class of higher-order corrections, can be resummed to all orders (achieved for $t\bar{t}H$) [Kulesza, Motyka, Schwartländer, Stebel, Theeuwes (2020)] [Broggio, Ferrogli, Frederix, Pagani, Pecjak, Tsinikos (2019)]

- ▶ Improving theoretical precision: adding contributions from soft gluon emission
 - ▶ Gluons emitted from coloured particles
 - ▶ Close to production threshold, only soft gluons
 - ▶ Logarithmic form with logarithms becoming large as threshold is approached
 - ▶ Depending on the observable, various measures of softness w (\sim distance to threshold) are used

$$w = 1 - z = 1 - Q^2/s \quad (\text{Invariant Mass Threshold}) \quad \rightarrow \quad \left[\frac{\ln(1 - z)}{1 - z} \right]_+$$

▶ Soft Gluon Resummation

- ▶ Large logs can invalidate predictive power of perturbative series in α_s
- ▶ Systematic treatment to all orders: resummation
- ▶ Relies on: $|\mathcal{M}|^2$ and Phase Space factorization \rightarrow Mellin space

$$\int_0^1 dz z^{N-1} f(z) = \{\mathcal{M}f\}(N) \equiv f_N \equiv \tilde{f}_N$$

▶ Technically complicated

- ▶ Resummed Logarithms: $\alpha_s^n \left[\frac{\ln^k(1-z)}{1-z} \right]_+$, $k \leq 2n - 1$

- ▶ Resummed partonic cross section for $pp \rightarrow Htj$

$$\begin{aligned}\tilde{\sigma}_{ij \rightarrow tHk}^{(\text{NLL})} &= \text{Tr} [\mathbf{H}_{ij \rightarrow tHk} \mathbf{S}_{ij \rightarrow tHk}] \Delta_i \Delta_j \mathcal{J}_k \\ &= \text{Tr} \left[\mathbf{H}_{ij \rightarrow tHk} \bar{\mathbf{U}}_{ij \rightarrow tHk} \tilde{\mathbf{S}}_{ij \rightarrow tHk} \mathbf{U}_{ij \rightarrow tHk} \right] \Delta_i \Delta_j \mathcal{J}_k \\ \mathbf{U}_{ij \rightarrow tHk}(N, Q^2, \mu_F^2, \mu_R^2) &= \text{P exp} \left[\int_{\mu_F}^{Q/\tilde{N}} \frac{dq}{q} \Gamma_{ij \rightarrow tHk}(\alpha_s(q^2)) \right]\end{aligned}$$

- ▶ One-loop Soft Anomalous Dimension $\Gamma_{ij \rightarrow tHk}^{(1)}$ checked with [Forslund, Kidonakis (2020)]

$$\Gamma_{ij \rightarrow tHk}(\alpha_s) = \left(\frac{\alpha_s}{\pi} \right) \Gamma_{ij \rightarrow tHk}^{(1)} + \left(\frac{\alpha_s}{\pi} \right)^2 \Gamma_{ij \rightarrow tHk}^{(2)} + \dots$$

$$\Gamma_{11, ij \rightarrow tHk}^{(1)} = C_F \left(\log \frac{s_{34} - m_t^2}{m_t \sqrt{s}} - \frac{1}{2} \right)$$

$$\Gamma_{12, ij \rightarrow tHk}^{(1)} = \frac{C_F}{2N_c} \log \frac{\tilde{t}_{13} \tilde{t}_{24}}{\tilde{t}_{23} \tilde{t}_{14}} \quad \Gamma_{21, ij \rightarrow tHk}^{(1)} = \log \frac{\tilde{t}_{13} \tilde{t}_{24}}{\tilde{t}_{23} \tilde{t}_{14}} \quad \text{with } \tilde{t}_{ik} = (p_i - p_k)^2 - m_k^2$$

$$\Gamma_{22, ij \rightarrow tHk}^{(1)} = C_F \left(\log \frac{s_{34} - m_t^2}{m_t \sqrt{s}} - \frac{1}{2} \right) + \frac{4C_F - N_c}{2} \log \frac{\tilde{t}_{13} \tilde{t}_{24}}{\tilde{t}_{23} \tilde{t}_{14}} + \frac{N_c}{2} \log \frac{\tilde{t}_{14} \tilde{t}_{23}}{s(s_{34} - m_t^2)}$$

- ▶ We take exponential functions at NLL accuracy
- ▶ \mathbf{H} and $\tilde{\mathbf{S}}$ up to $\mathcal{O}(\alpha_s)$

$$\mathbf{H}(\mu_R, \mu_F) = \mathbf{H}^{(0)} + \frac{\alpha_s(\mu_R)}{\pi} \mathbf{H}^{(1)}(\mu_R, \mu_F)$$

$$\tilde{\mathbf{S}}(1) = \tilde{\mathbf{S}}^{(0)} \left(1 + \frac{\alpha_s(\mu_R)}{\pi} \tilde{\mathbf{S}}^{(1)}(1) \right)$$

- ▶ Resulting in NLLwC accuracy (NLL' in SCET notation)
- ▶ Invariant mass threshold:

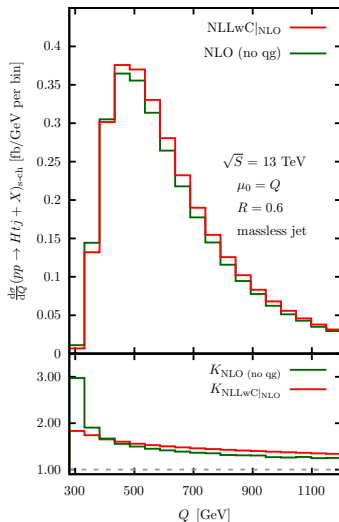
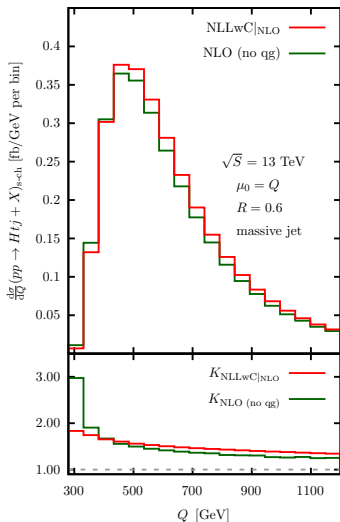
- ▶ Massless jet at threshold ($p_j^2 = 0$): $1 - z = 1 - \frac{(p_H + p_t)^2 + 2p_j \cdot (p_H + p_t)}{\hat{s}}$
- ▶ Massive jet at threshold: $1 - z = 1 - \frac{Q^2}{\hat{s}}$, with $Q^2 = (p_H + p_t + p_j)^2$

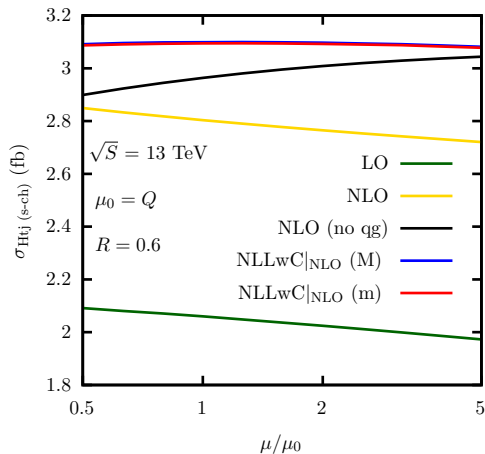
[Kidonakis, Oderda, Sterman (1998)]

► Matching to NLO cross section

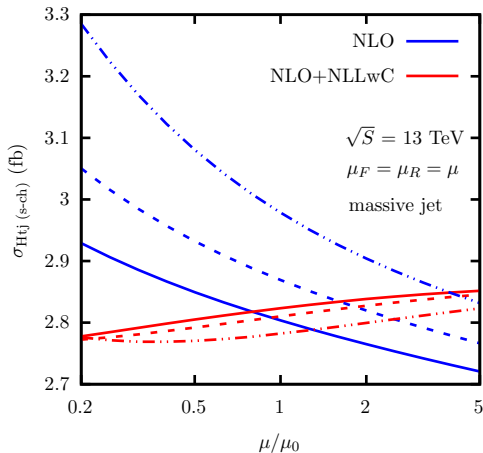
$$\begin{aligned} \sigma_{pp \rightarrow Htj}^{NLO+NLLwC} &= \sigma_{pp \rightarrow Htj}^{NLO} \\ &+ \sum_{i,k} \int \frac{dN}{2\pi i} \rho_h^{-N} \tilde{f}_i(N+1, \mu_F^2) \tilde{f}_k(N+1, \mu_F^2) \\ &\times \left[\frac{\tilde{\sigma}_{ik \rightarrow Htj}^{NLLwC}(N) - \tilde{\sigma}_{ik \rightarrow Htj}^{NLO}(N)}{\text{avoid double counting with NLO!}} \right] \end{aligned}$$

- $q\bar{q}'$ channel at NLO+NLLwC
- qg channel at NLO

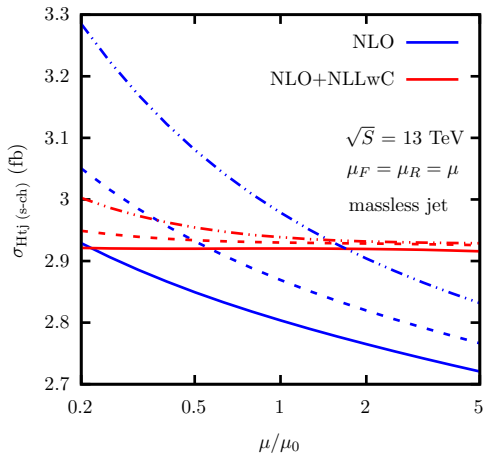




Different central scale choices: Q : —, $H_T/2$: ---, $H_T/6$:-·-·-



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- ▶ Associated top and Higgs production very important for BSM searches: coupling $t\bar{t}H$
- ▶ Extending precision of theoretical calculations beyond NLO
- ▶ Effects of soft gluon emission at NLLwC
- ▶ Different jet treatment: massive and massless at threshold
- ▶ Reduced scale dependence NLO+NLLwC

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