

# Rare Decays at CMS

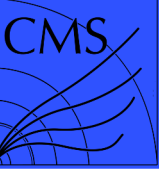
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On behalf of the CMS Collaboration

EPS-HEP 2021,  
26-30 Jul 2021, DESY and University of Hamburg  
Flavor Physics Parallel Talk  
30<sup>th</sup> July 2021

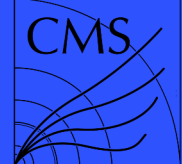


# Outline of the talk



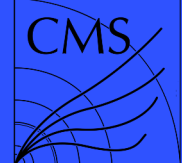
- CLFV decay of  $\tau \rightarrow \mu\mu\mu$  (DOI: [10.1007/JHEP01\(2021\)163](https://doi.org/10.1007/JHEP01(2021)163))
- Rare Decay of  $B \rightarrow \mu\mu$  (DOI: [10.1007/JHEP04\(2020\)188](https://doi.org/10.1007/JHEP04(2020)188), CMS-PAS-BPH-20-003)
- Angular Analysis  $B^+ \rightarrow K^{*+}\mu\mu$  (DOI: [10.1007/JHEP04\(2021\)124](https://doi.org/10.1007/JHEP04(2021)124))

# Why Rare Decays?

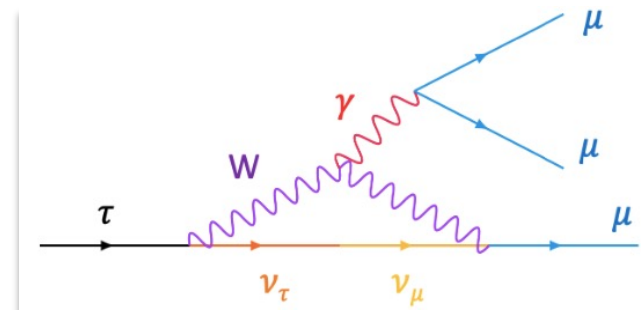


- Flavor change in charged leptons is not necessarily forbidden by any symmetry of the SM; yet no such decay is observed
- The Flavor Changing Neutral Currents (FCNCs) are suppressed in the SM; Larger branching ratios could indicate BSM physics
- Precision measurements of observables needed to check for potential discrepancies in the SM predictions

# Search for CLFV Decay of $\tau \rightarrow \mu\mu\mu$

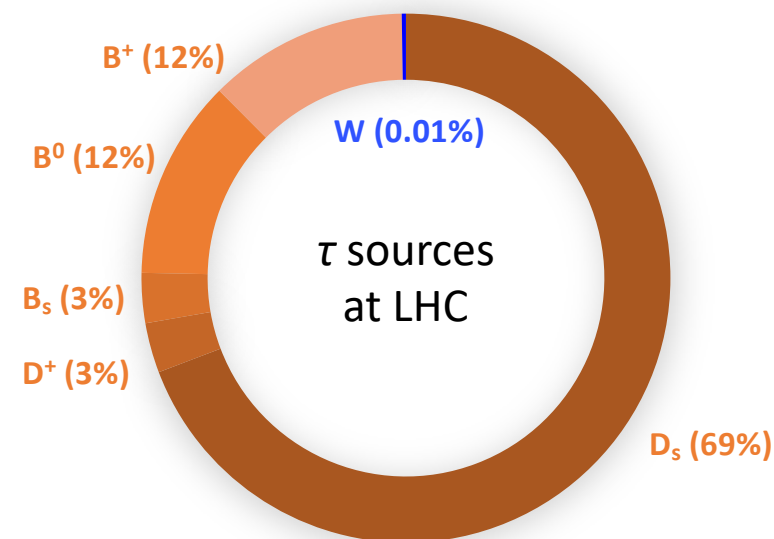


- $\tau \rightarrow \mu\mu\mu$  decay is suppressed in the SM, but several BSM theories predict higher branching ratios ( $\sim O(10^{-54})$ ) accessible by current detector technologies
- Searches have been performed by many experiments over last six decades, but no excess has been observed yet
- At LHC, searches for CLFV  $\tau$  decays are favored by their abundant production  $O(10^{11})/\text{fb}^{-1}$
- As opposed to  $e^+e^-$  colliders,  $\tau$ s produced at LHC have significantly more background and hence the analysis is more challenging



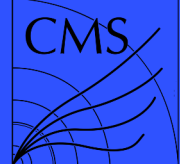
SM decay of  $\tau \rightarrow \mu\mu\mu$  occurring via neutrino oscillation with a branching ratio of  $O(10^{-54})$

Experiment	Observed (Expected) upper limit on $B(\tau \rightarrow 3\mu)$ ( $\times 10^8$ at 90% C.L.)
Belle	2.1 (-)
BaBar	3.3 (4.0)
LHCb (Run I data)	4.6 (5.0)
ATLAS (Run I data)	38 (39)

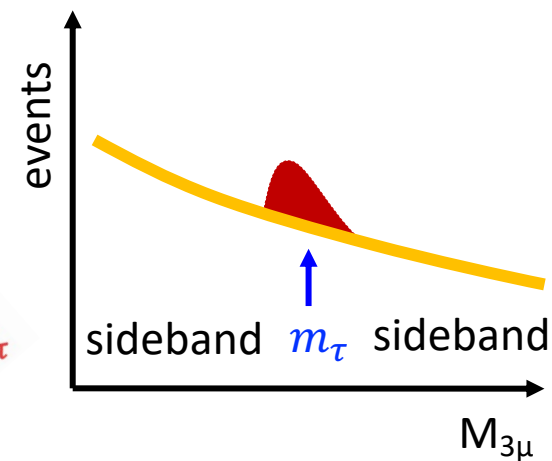
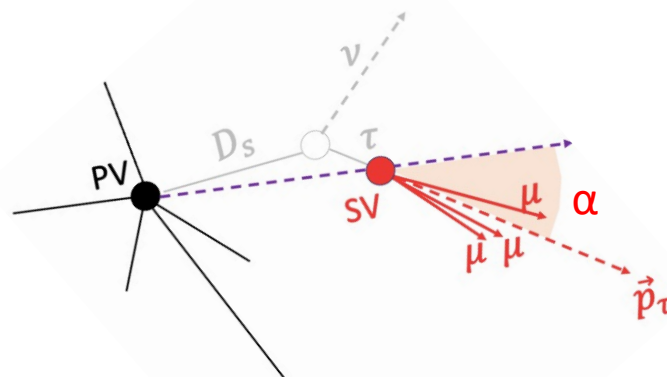
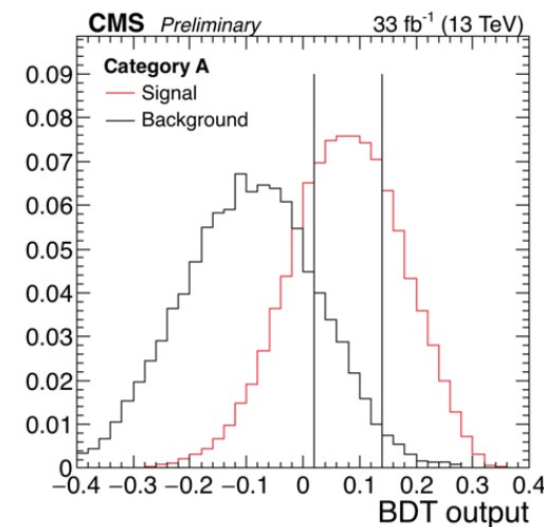
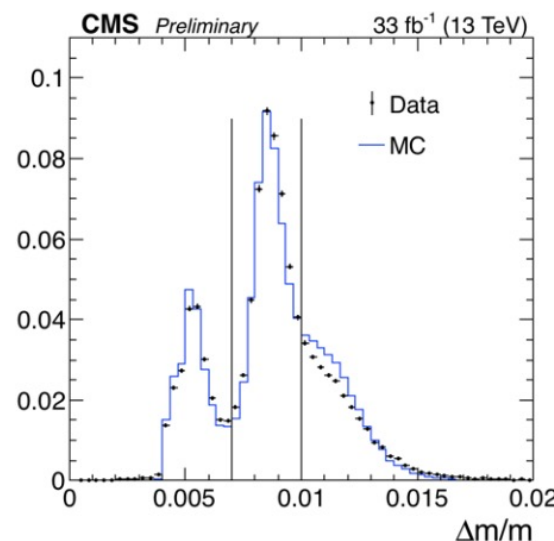




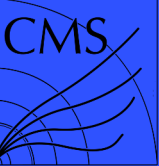
# Search Strategy (Heavy Flavor Channel)



- Characterized by
  - Displaced Secondary Vertex
  - low momenta muons boosted in the forward region
  - High QCD background
- Dedicated online triggers for online event selection
  - Two muons and a track forming a displaced vertex (HF)
- Event categorization based on mass resolution
  - Mass resolution of the tau candidate can vary upto 3 times depending on the  $\eta$
- Multivariate Analysis (BDT) for background rejection
  - vertex  $\chi^2$
  - vertex displacement significance
  - $\alpha$  angle
  - muon quality observables
- Signal extraction in multiple categories



# Signal Normalization (Heavy Flavor channel)



- Number of signal events:

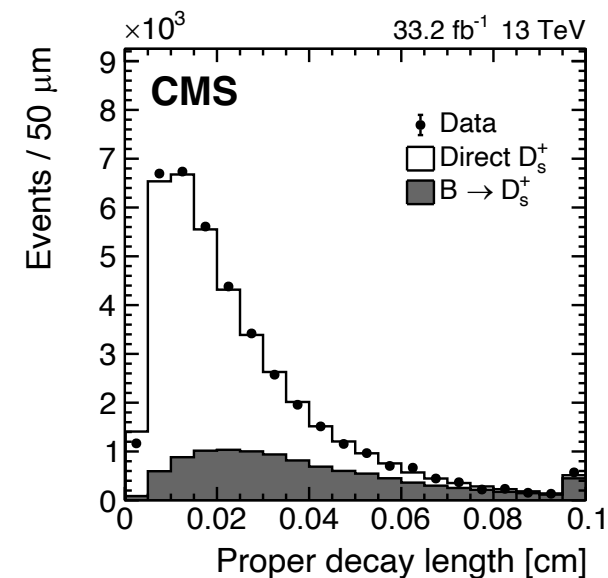
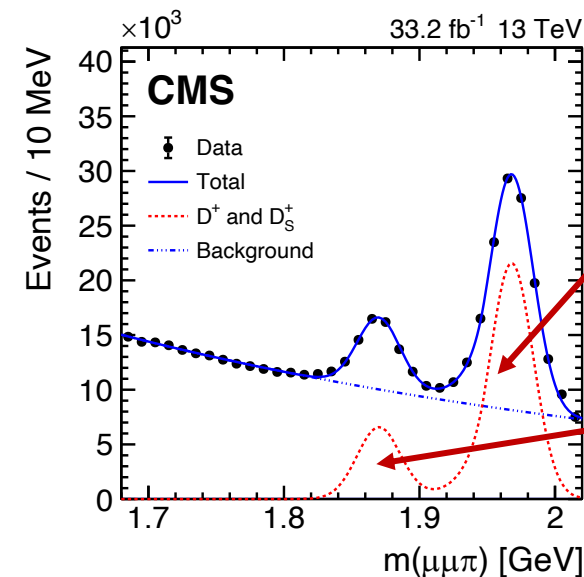
$$N_{\text{sig}(D_s)} = \mathcal{L} \sigma(pp \rightarrow D_s) \mathcal{B}(D_s \rightarrow \tau) \mathcal{B}(\tau \rightarrow 3\mu) A_{3\mu(D_s)} \epsilon_{\text{reco}}^{3\mu} \epsilon_{\text{trig}(\text{sig})}^{3\mu}$$

- $D_s$  Yield from  $D_s \rightarrow \Phi (\mu\mu)\pi$  decay:

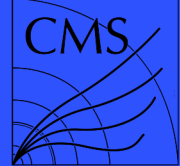
$$N = \mathcal{L} \sigma(pp \rightarrow D_s) \mathcal{B}(D_s \rightarrow \varphi\pi \rightarrow \mu\mu\pi) A_{\mu\mu\pi(D_s)} \epsilon_{\text{reco}}^{2\mu\pi} \epsilon_{\text{trig}(\mu\mu\pi)}^{2\mu}$$

$$\text{➤ } N_{\text{sig}(D_s)} = N \frac{\mathcal{B}(D_s \rightarrow \tau)}{\mathcal{B}(D_s \rightarrow \Phi\pi \rightarrow \mu\mu\pi)} \frac{A_{3\mu(D_s)}}{A_{\mu\mu\pi(D_s)}} \frac{\epsilon_{\text{reco}}^{3\mu}}{\epsilon_{\text{reco}}^{2\mu\pi}} \frac{\epsilon_{\text{trig}(\text{sig})}^{3\mu}}{\epsilon_{\text{trig}(\mu\mu\pi)}^{2\mu}} \mathcal{B}(\tau \rightarrow 3\mu)$$

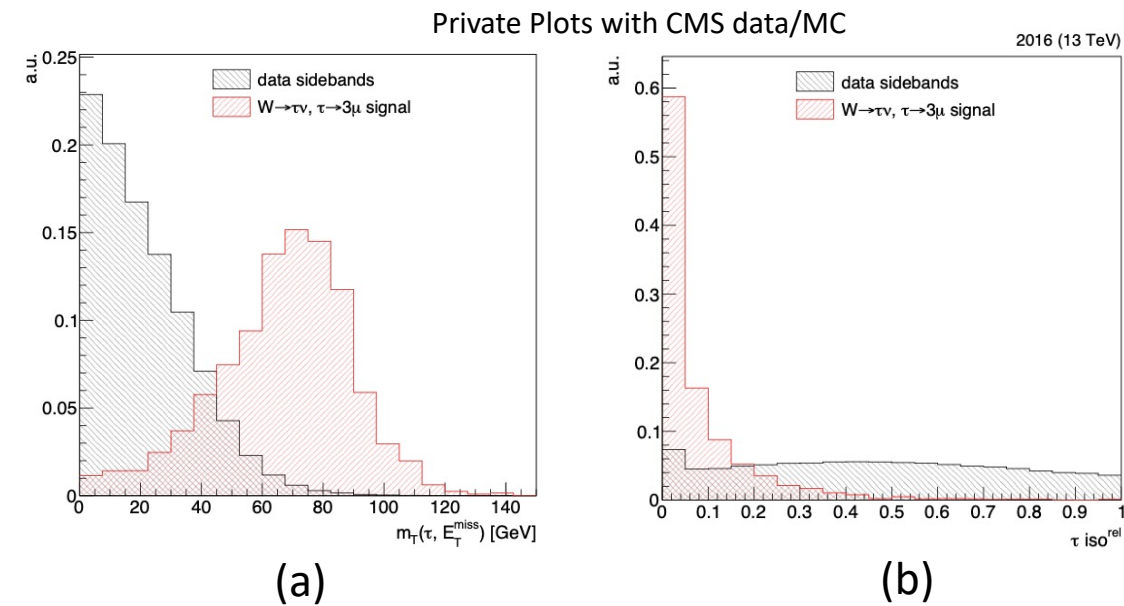
- The signal yield is independent of the luminosity and cross section measurements
- The B contribution to the signal is assessed through  $D_s \rightarrow \varphi(\mu\mu)\pi$  channel
  - Fit two templates of  $D_s$  decay length corresponding to prompt and non-prompt decays, derived from MC, to background subtracted data and extract the relative fraction



# Search Strategy (W channel)



- Even after having a small contribution to the  $\tau$  production at LHC ( $\sim 0.01\%$ ), W channel has significant advantages
  - Fewer background sources
  - Isolated high momentum muons
  - Higher trigger and offline selection efficiency
  - Presence of large MET in the final state
- Online trigger selection
  - Dedicated three muon triggers
- MVA analysis using BDT for background rejection
  - Transverse mass W
  - Tau Isolation
  - Muon Id variable
- Event categorization based on the mass resolution (barrel, endcap)

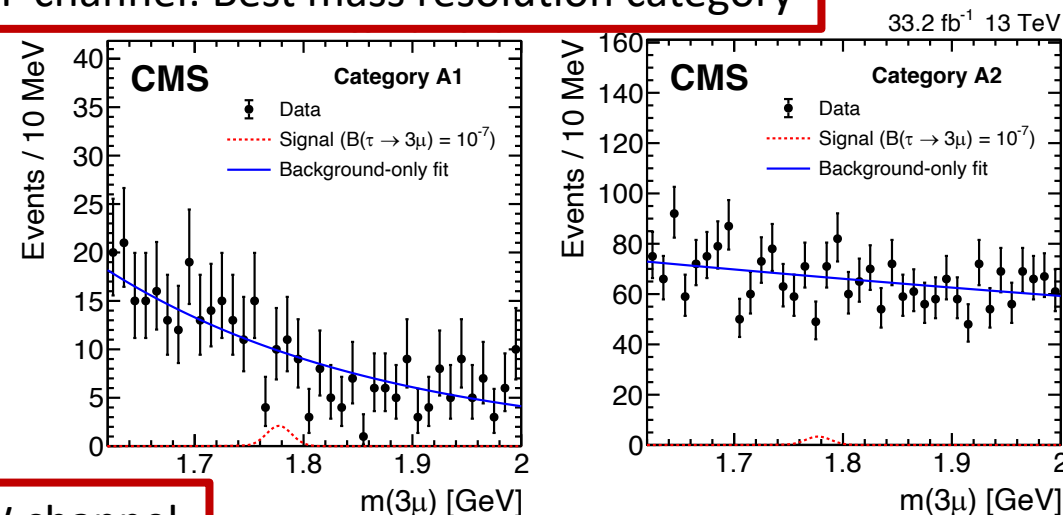


Signal Selection Efficiency		
	HF channel	W channel
Muon acceptance*	1.5%	64%
Trigger efficiency	2%	23.5%

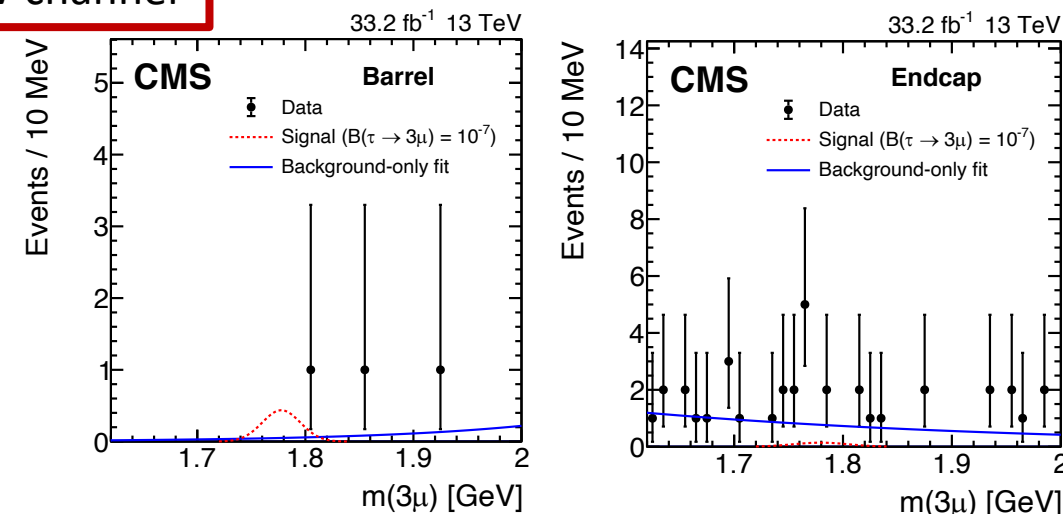
\*Muon acceptance: all 3 muons have  $p > 2.5$  and  $\eta < 2.4$

- 2 categories from W channel are combined with 6 categories from the HF channel
- Observed (Expected) limit is  $8.0$  ( $6.9$ )  $\times 10^{-8}$  at 90% C.L.
- Fitting W boson and HF events separately returns an observed(expected) upper limits of  $20 \times 10^{-8}$  ( $13 \times 10^{-8}$ ) and  $9.2 \times 10^{-8}$  ( $10.0 \times 10^{-8}$ ) at 90% C.L.
- Full Run 2 analysis is currently ongoing and the results are going to be released soon
  - 4 times more integrated luminosity
  - new pixel detector since 2017

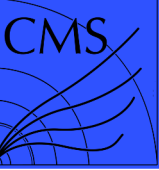
## HF channel: Best mass resolution category



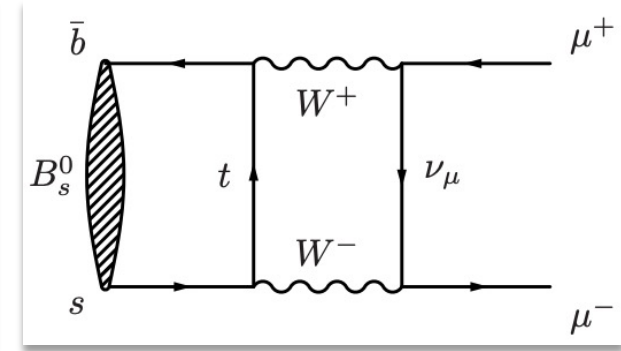
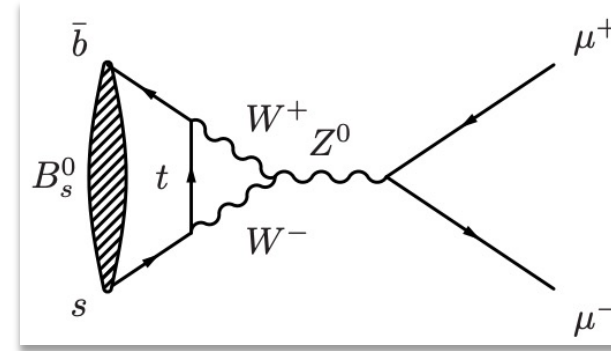
## W channel



# Rare Decay of $B \rightarrow \mu\mu$



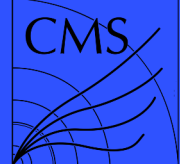
- SM decay suppressed by CKM and helicity
  - $B(B_s^0 \rightarrow \mu\mu) = (3.57 \pm 0.17) \times 10^{-9}$
  - $B(B^0 \rightarrow \mu\mu) = (1.06 \pm 0.09) \times 10^{-10}$
- Theoretical uncertainties in the branching ratios have been reduced in the recent years
- Contributions from BSM processes could either enhance or further suppress the branching ratios
- $B_s^0$  can be expressed as a linear combination of the two mass eigenstates- heavy ( $B_{sH}^0$ ) and light( $B_{sL}^0$ )
- The SM predicts that the lifetime of the  $B_s^0 \rightarrow \mu\mu$  receives contribution from the heavy state
- Contributions from new physics could change the contributions from heavy and light mass states ( $R_H^{\mu\mu}$ ,  $R_L^{\mu\mu}$ )
- Searches have been performed in ATLAS, CMS and LHCb



Experiment	$B(B_s^0 \rightarrow \mu\mu) \times 10^9$	$B(B^0 \rightarrow \mu\mu) \times 10^{10}$
CMS (Run I)	$3.0^{+1.0}_{-0.9}$	$<11$ @ 95% C.L.
CMS+LHCb (Run I)	$2.8^{+0.7}_{-0.6}$	$3.9^{+1.6}_{-1.4}$
ATLAS (Run I+2015+2016)	$2.8^{+0.8}_{-0.7}$	$<2.1$ @ 95% C.L.
LHCb (Run I+ Run II)	$3.0 \pm 0.6^{+0.3}_{-0.2}$	$<3.4$ @ 95% C.L.



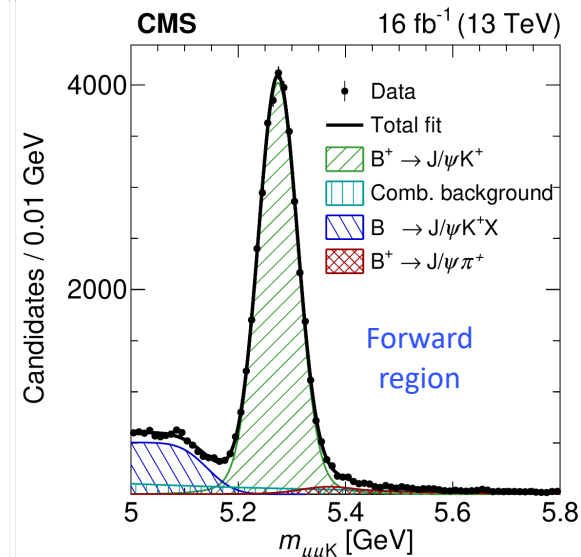
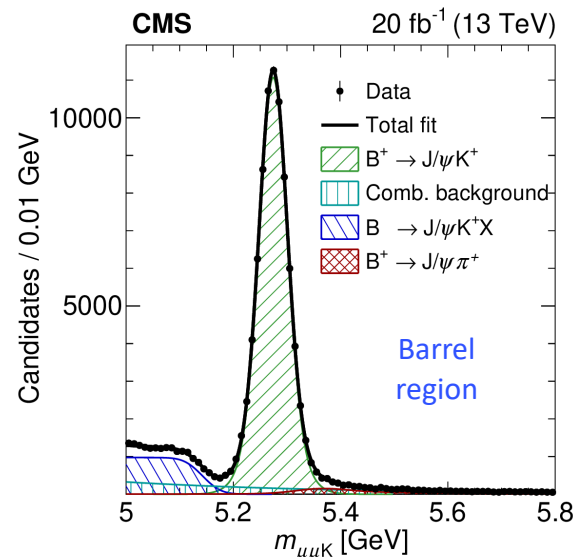
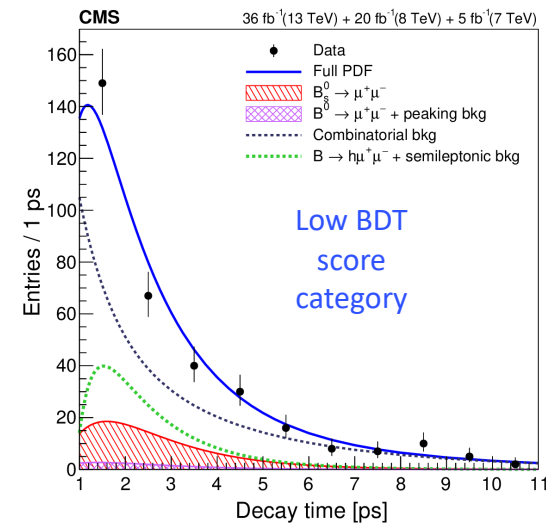
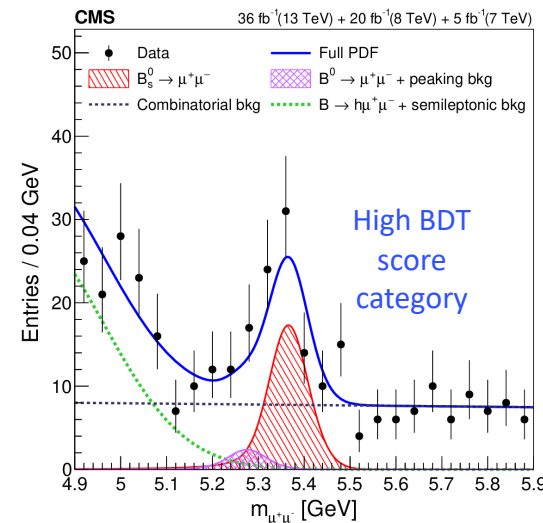
# Rare Decay of $B \rightarrow \mu\mu$



- HLT selection using opposite sign dimuons displaced from the vertex
- Background contributions from semi-leptonic and hadronic combinatorial decays, rare decays of  $B_s \rightarrow K\mu\nu$ ,  $\Lambda_b \rightarrow p\mu\nu$ ,  $B_s \rightarrow KK$  with the hadrons misidentified as a muon
- MVA analysis used to perform signal-background discrimination
- Signal Normalization performed using the  $B \rightarrow J/\psi(\mu\mu)K^+$

$$\mathcal{B}(B_s^0 \rightarrow \mu\mu) = \frac{n_{B_s^0}^{obs}}{N(B \rightarrow J/\psi K) f_d} \frac{f_s}{A_{B_s^0}} \frac{A_{B^+}}{\varepsilon_{B_s^0}} \frac{\varepsilon_{B^+}}{\varepsilon_{B_s^0}} \mathcal{B}(B \rightarrow J/\psi(\mu\mu)K)$$

- Signal yields are extracted in the bins of the output of the MVA
- The ratio of fragmentation fraction ( $f_s/f_d$ ) is known with limited precision and is the dominant source of systematics



- Binned two-dimensional profile likelihoods obtained by each experiment from their fit to the dimuon invariant mass distributions

- The branching ratio of  $B_s^0 \rightarrow \mu\mu$  has been measured to be

$$\mathcal{B}(B_s^0 \rightarrow \mu^+ \mu^-) = (2.69^{+0.37}_{-0.35}) \times 10^{-9}$$

- An upper limit has been set on the branching ratios of  $B^0 \rightarrow \mu\mu$ :

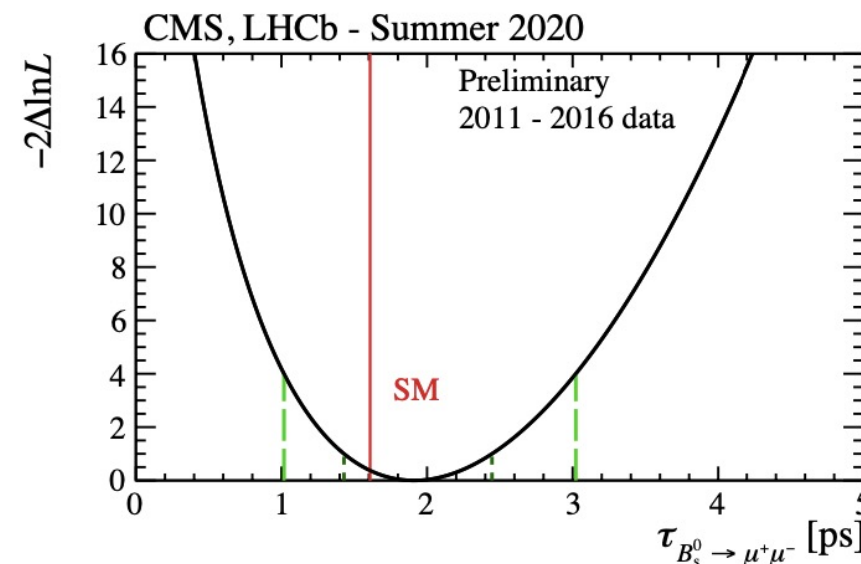
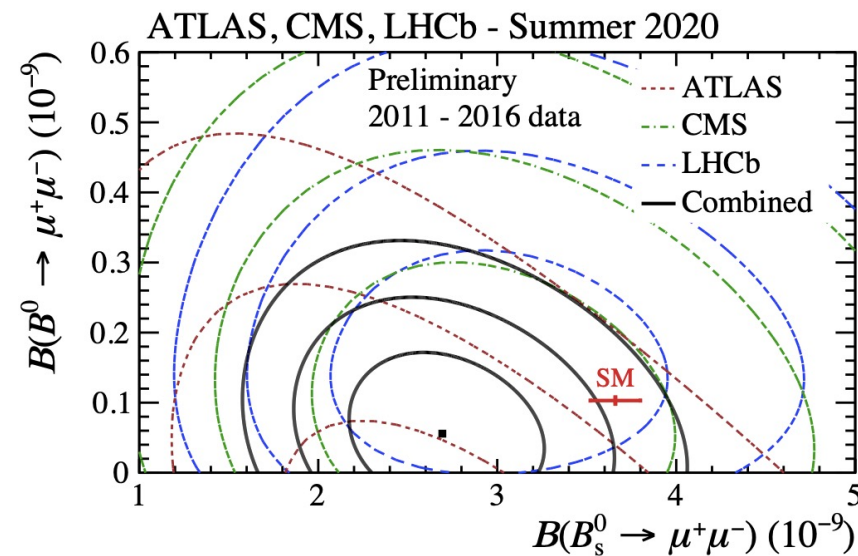
$$\mathcal{B}(B^0 \rightarrow \mu^+ \mu^-) < 1.6 \times 10^{-10} \text{ at 90\% CL}$$

$$\mathcal{B}(B^0 \rightarrow \mu^+ \mu^-) < 1.9 \times 10^{-10} \text{ at 95\% CL}$$

- The results are compatible with the SM predictions within 2.1 standard deviations with the inclusion of theoretical uncertainties

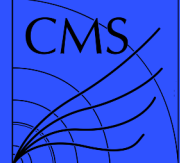
- The effective lifetime of  $B_s^0 \rightarrow \mu\mu$  has been measured to be

$$\tau_{B_s^0 \rightarrow \mu^+ \mu^-} = 1.91^{+0.37}_{-0.35} \text{ ps}$$



\*<https://cms-results.web.cern.ch/cms-results/public-results/preliminary-results/BPH-20-003/index.html>

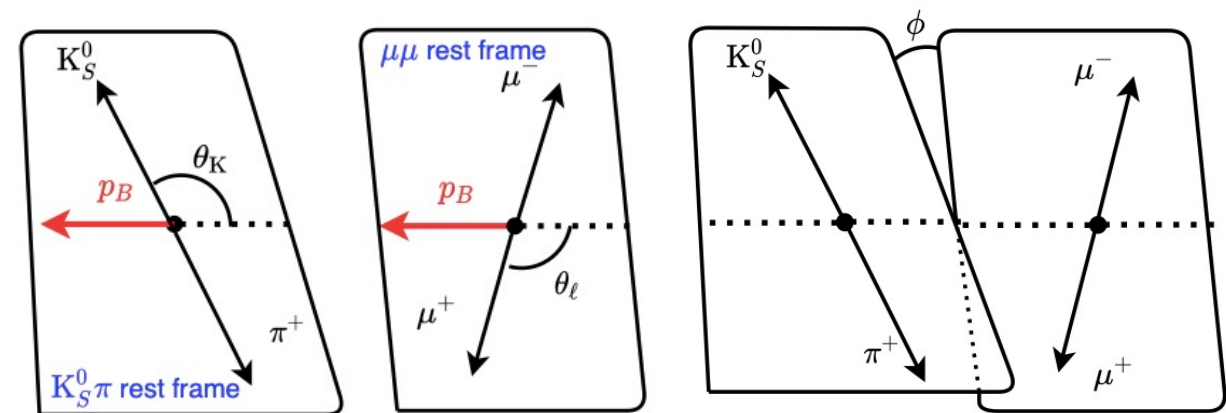
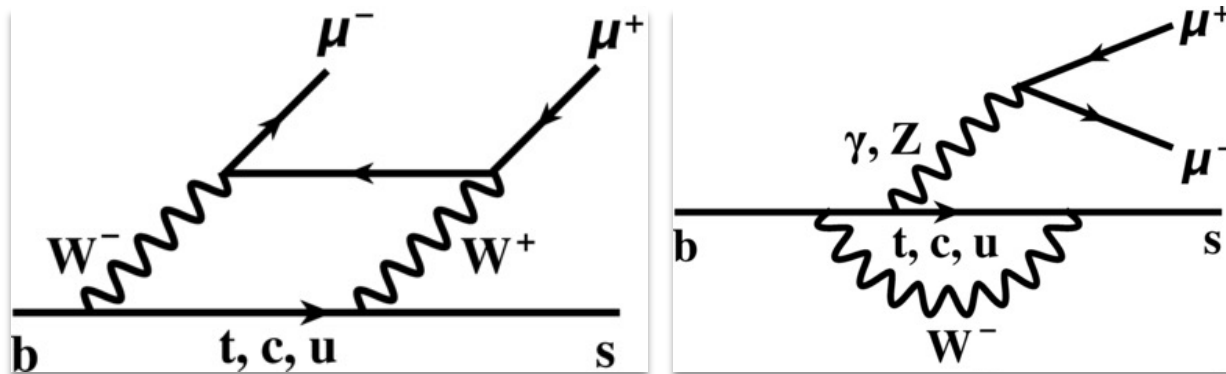
# Angular Analysis $B^+ \rightarrow K^{*+} \mu \mu$



- $b \rightarrow s ll$  FCNC process which is forbidden at tree level in SM, but only allowed through penguin and box diagrams
- Presence of BSM particles mediating the decay could alter the differential decay rate

$$\frac{1}{\Gamma_l} \frac{d^3\Gamma_l}{d\cos\theta_l d\cos\theta_k} = \frac{9}{16} \left\{ \frac{2}{3} [F_S + 2A_S \cos\theta_k](1 + \cos^2\theta_l) + (1 - F_S)[2F_L(\cos^2\theta_k)(1 + \cos^2\theta_l) + \frac{1}{2}(1 - F_L)(1 + \cos^2\theta_k)(1 + \cos^2\theta_l) + \frac{3}{4}A_{FB}(1 - \cos^2\theta_k)\cos\theta_l] \right\}$$

- $\theta_k$ : Angle between  $K_S$  and  $-\mathbf{p}(B^+)$  in the  $K^*$  rest frame
- $\theta_l$ : Angle between  $\mu^+$  and  $-\mathbf{p}(B^+)$  in the dimuon rest frame
- $F_L$ :  $K^*$  longitudinal polarization factor
- $A_{FB}$ : muon forward-backward symmetry
- $F_S$ : S-wave fraction
- $A_S$ : Interference amplitude between S and P wave



# Analysis Strategy



- Analysis performed in three different  $q^2$  (1-19 GeV) bins from 3D extended maximum-likelihood fit to angular distribution to extract the parameters  $A_{FB}$  and  $F_L$
- Dedicated HLT to trigger on low momenta dimuon events with a displaced vertex
- Offline reconstruction of displaced secondary vertex
- Veto on  $J/\psi$  and  $\psi(2S)$  masses based on  $q^2$
- Simulated shapes of the signal are used to generate PDFs and the efficiency as a function of  $\theta_k, \theta_l$

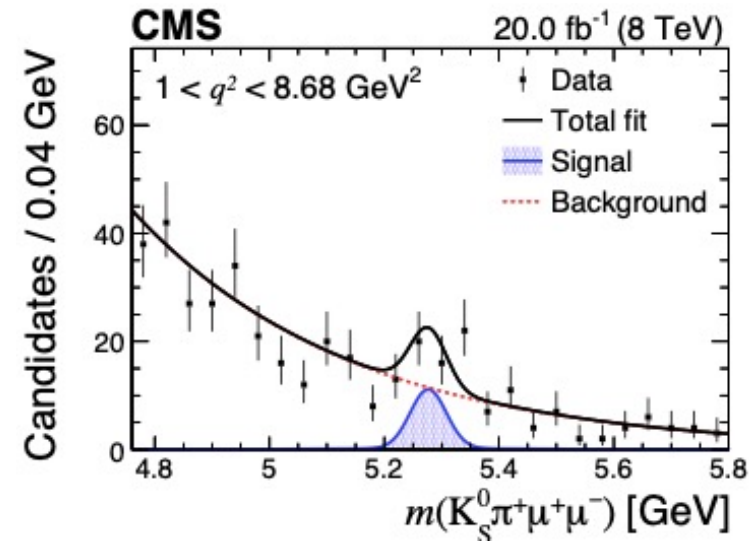
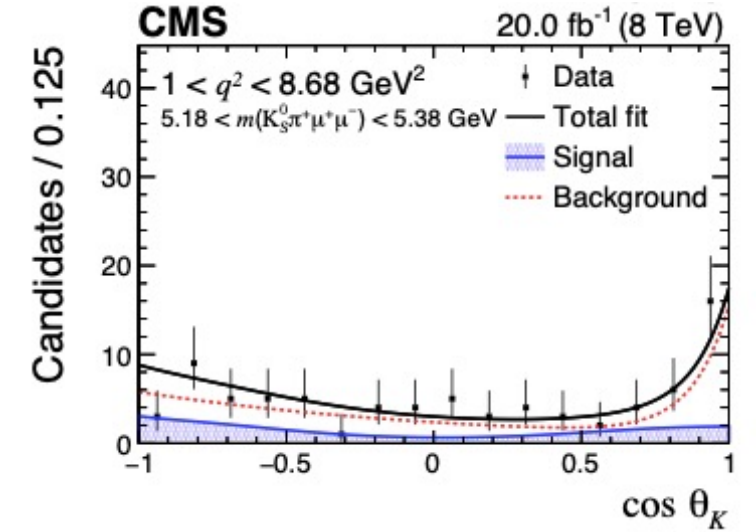
$$\text{pdf}(m, \cos \theta_k, \cos \theta_l) = Y_S S^m(m) S^a(\cos \theta_k, \cos \theta_l) \varepsilon(\cos \theta_k, \cos \theta_l) + Y_B B^m(m) B^{\theta_K}(\cos \theta_k) B^{\theta_l}(\cos \theta_l)$$

$Y_S, Y_B$ : signal and background yields (free parameters)

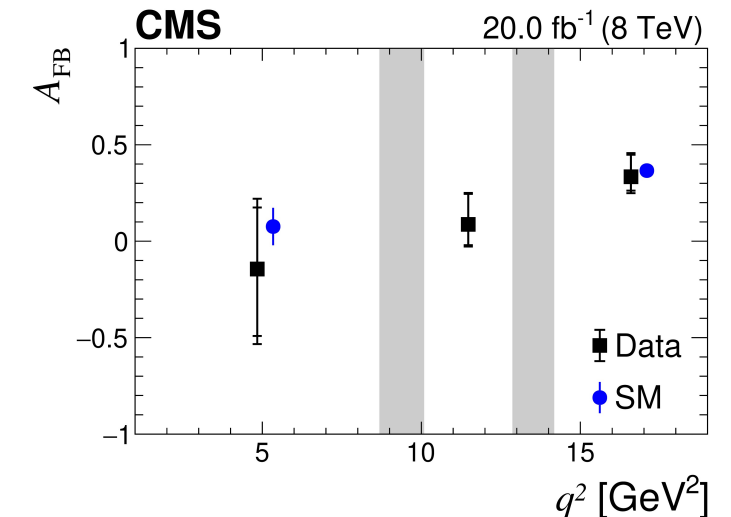
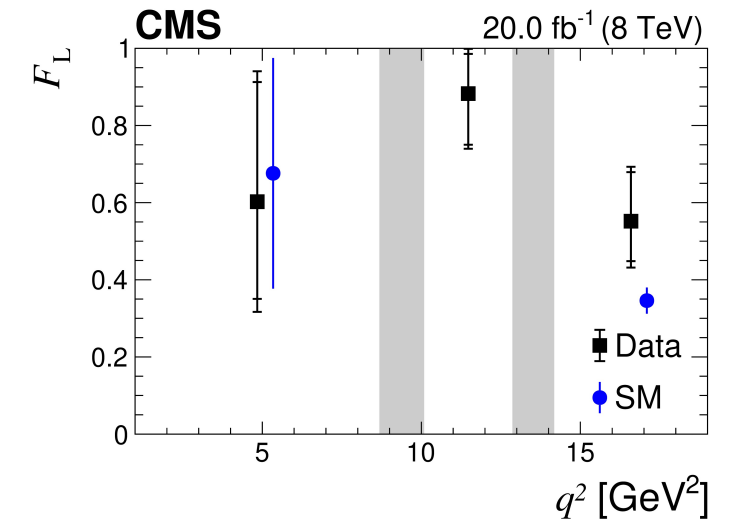
$S^m(m)$ : signal mass shapes (derived from simulations)

$B^m(m)$ : exponential function (free parameter)

$B^{\theta_K}(\cos \theta_k), B^{\theta_l}(\cos \theta_l)$ : Background shapes (sidebands)



- The results performed using Run I data indicate no deviations from the SM predictions
- The systematics were driven by the uncertainty in background shapes and extension of backgrounds to the sideband regions
- The parameters  $A_{\text{FB}}$  and  $F_L$  are extracted in each  $q^2$  bin (dimuon invariant mass)
- The results are consistent with the SM predictions
- The measurements using the full Run II dataset are currently ongoing

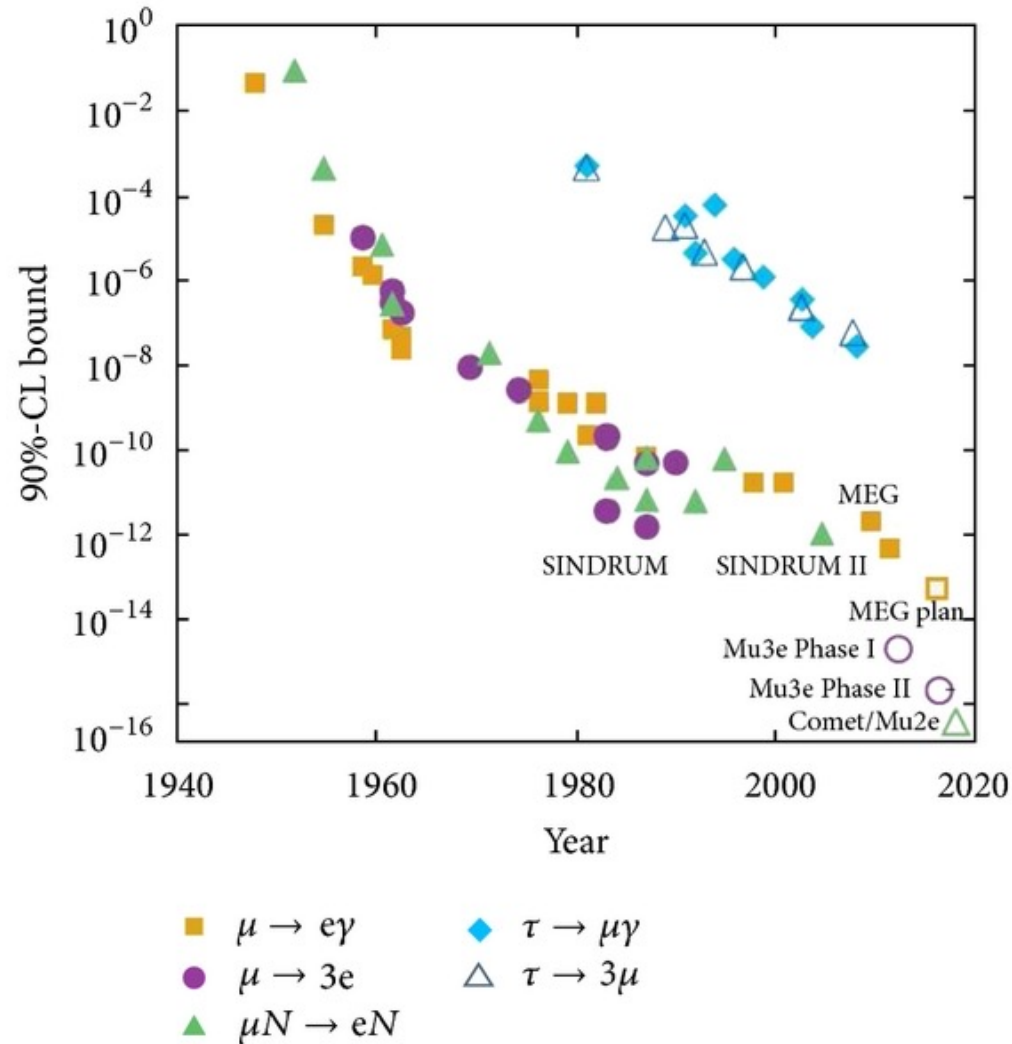
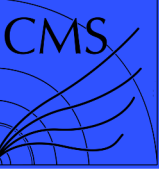




- A search for  $\tau \rightarrow \mu\mu\mu$  was performed at CMS using both Heavy Flavor (HF) and W channels; the combined limit on the branching ratio is obtained to be  $8.0 \times 10^{-8}$  at 90% C.L.; no excess is seen
- The results obtained from ATLAS, CMS and LHCb experiments on  $B_{(s)}^0 \rightarrow \mu^+\mu^-$  decays obtained from the data collected between 2011 and 2016 have been combined; The effective lifetime of  $\tau_{B_s^0 \rightarrow \mu^+\mu^-}$  is the most precise measurement to date.
- The angular analysis of  $B^+ \rightarrow K^{*+}\mu\mu$  is performed and the parameters  $F_L$  and  $A_{FB}$  are measured; no deviations from the SM are observed

# Backup

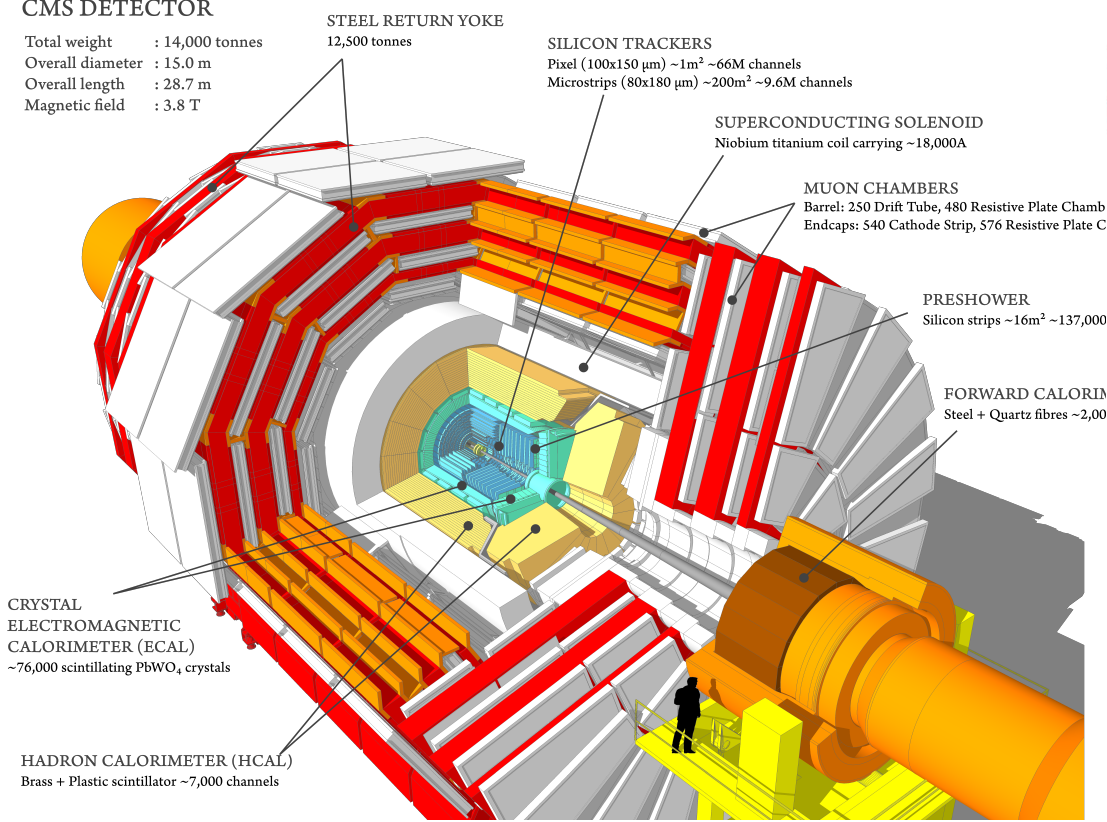
# History of CLFV decays



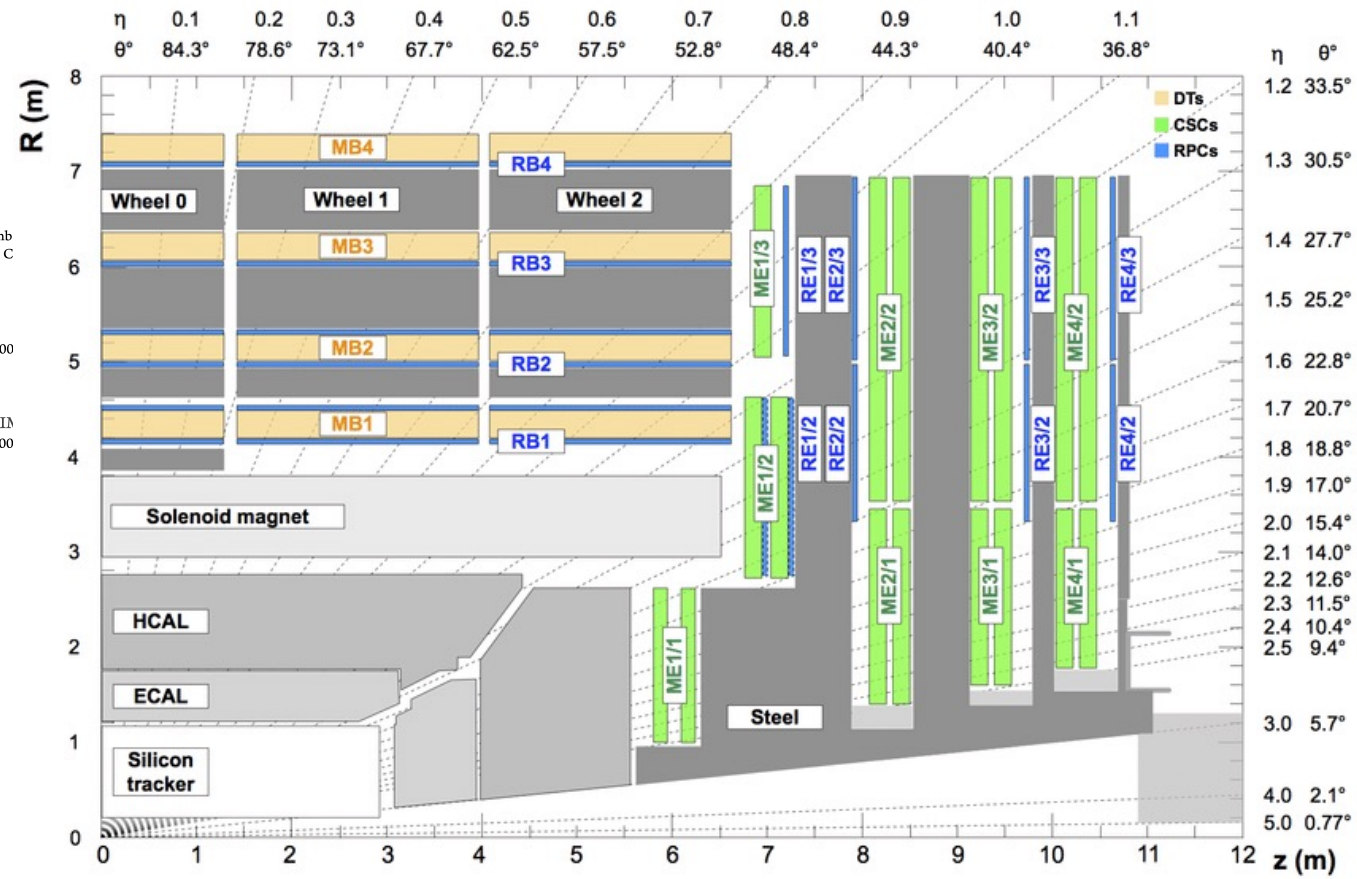
## 3D Model of the Detector Design

### CMS DETECTOR

Total weight : 14,000 tonnes  
Overall diameter : 15.0 m  
Overall length : 28.7 m  
Magnetic field : 3.8 T



## Muon System

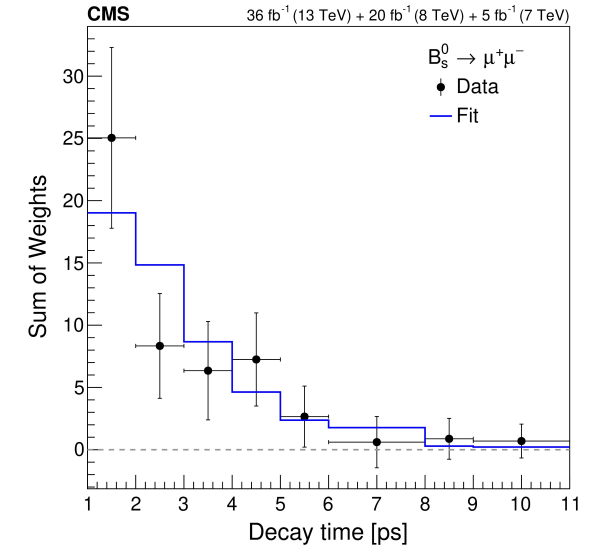
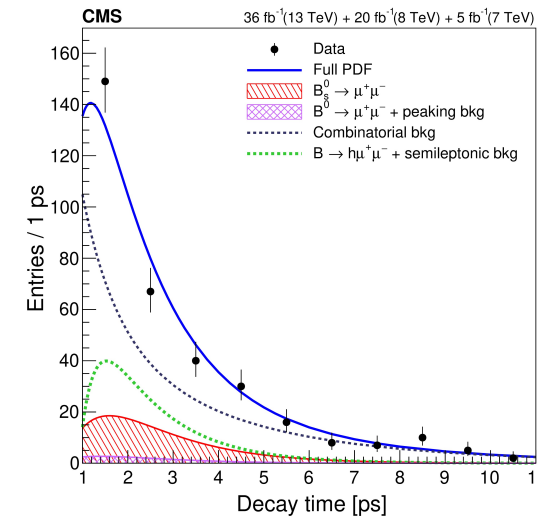


## Effective Lifetime

$$\tau_{\mu^+\mu^-} \equiv \frac{\int_0^\infty t \Gamma(B_s(t) \rightarrow \mu^+\mu^-) dt}{\int_0^\infty \Gamma(B_s(t) \rightarrow \mu^+\mu^-) dt} = \frac{\tau_{B_s^0}}{1 - y_s^2} \left( \frac{1 + 2\mathcal{A}_{\Delta\Gamma}^{\mu^+\mu^-} y_s + y_s^2}{1 + \mathcal{A}_{\Delta\Gamma}^{\mu^+\mu^-} y_s} \right)$$

$$\tau_{B_s^0 \rightarrow \mu^+\mu^-} = \frac{\tau_{B_s^0}}{1 - y_s^2} \left[ \frac{1 + 2\mathcal{A}_{\Delta\Gamma} y_s + y_s^2}{1 + \mathcal{A}_{\Delta\Gamma} y_s} \right]$$

$$y_s \equiv \frac{\Delta\Gamma_s}{2\Gamma_s}, \quad \mathcal{A}_{\Delta\Gamma} \equiv \frac{R_H^{\mu^+\mu^-} - R_L^{\mu^+\mu^-}}{R_H^{\mu^+\mu^-} + R_L^{\mu^+\mu^-}}$$



Reference: <https://cms-results.web.cern.ch/cms-results/public-results/publications/BPH-16-004/index.html>