On the sources of CP violation in the Lepton Sector



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Current experimental situation



Where is New Physics (in neutrino oscillations) ?

$$P \sim |A^{SM} + \epsilon A^{NP}|^2 \sim P^{SM} + 2 \epsilon \Re (A^{SM} A^{NP})$$

in the standard 3-v paradigm



- in the absence of correlation between NP and standard parameters, strong constraints

- if correlation is strong, thus bounds can be (partially) relaxed



complementary studies compared to sterile and NSI searches at NT

Future Experimental Alternatives (some of them)



- New Physics

- Super-Kamiokande-Gd
- IceCube-Gen2, KM3NeT, ARA
- PTOLEMY

CP violation

Our concern here: study the existence of <u>new sources</u> of CP violation

Let us set our strategy: look for leptonic CP asymmetries



- we DO NOT want to study the sensitivity to some New Physics phase

- we want to assess whether all asymmetries are described in terms of the single leptonic Jarlskog invariant as predicted in the absence of New Physics effects
- to achieve this task:
 - * compute the leptonic asymmetries in the Standard Model
 - * evaluate the experimental uncertainty on them
 - * re-compute the asymmetries including the effects of New Physics
 - * check whether the new results are sufficiently away from the Standard model predictions

A. Giarnetti and DM New Sources of Leptonic CP Violation at the DUNE Neutrino Experiment Universe 2021, 7(7), 240

CP Asymmetries in the Standard Model: vacuum

$$A_{\alpha\beta} = \frac{P_{\alpha\beta} - P_{\bar{\alpha}\bar{\beta}}}{P_{\alpha\beta} + P_{\bar{\alpha}\bar{\beta}}}$$

Use perturbation theory to evaluate them (given here the vacuum case):

$$\Delta_{21} = \Delta m_{21}^2 L/4E_{\nu} \ll 1. \qquad s_{13} = \frac{r}{\sqrt{2}}, \quad s_{12} = \frac{1}{\sqrt{3}}(1+s), \quad s_{23} = \frac{1}{\sqrt{2}}(1+a) \qquad \text{r, s, a = } \lambda \sim O(10\%)$$

$$A_{\mu e} = \frac{-12 r \Delta_{21} \sin \delta \sin^2 \Delta_{31}}{4 \Delta_{21}^2 + 9 r^2 \sin^2 \Delta_{31} + 6 r \Delta_{21} \cos \delta \sin 2 \Delta_{31}} \sim O(1)$$

$$A_{\mu\tau} = \frac{4}{3} r \Delta_{21} \sin \delta \sim O(\lambda^2)$$

 $(A_{\mu\mu}=0)$

- $A_{\mu e}$ and $~A_{\mu \tau}$ asymmetries in vacuum are suppressed by the small quantities Δ_{21} and θ_{13}
- since the denominators $A_{\mu e}$ is also suppressed, a partial cancellation is at work and, in particular, one generically expects $A_{\mu e} > A_{\mu \tau}$

CP Asymmetries in the Standard Model: the case of small matter effects

$$A_{\alpha\beta} = A_{\alpha\beta}^{SM_{1}} + V_{CC} A_{\alpha\beta}^{SM_{1}}$$

$$a = \frac{\Delta_{31}}{\Delta_{31}} \quad V_{CC} \sim O(\lambda) \quad f_{1} \sim O(r^{2})$$

$$A_{\mu e}^{SM_{1}} = -\frac{6}{f_{1}} r(\Delta_{31} \cos \Delta_{31} - \sin \Delta_{31}) [2\alpha \Delta_{31} \cos \delta \cos \Delta_{31} + 3r \sin \Delta_{31} + 2r \alpha^{2} \sin^{2} \delta \Delta_{31}^{2} \sin^{3} \Delta_{31}], \longrightarrow O(\lambda)$$

$$A_{\mu \tau}^{SM_{1}} = -2r^{2}(1 - \Delta_{31} \cot \Delta_{31}) + \frac{8}{27}\alpha^{2} \Delta_{31}^{3} \cot \Delta_{31}, \longrightarrow O(\lambda^{2})$$

$$A_{\mu\mu}^{SM_{1}} = \frac{4}{3}r\alpha \Delta_{31} \cos \delta(\Delta_{31} - \tan \Delta_{31}) - \frac{8}{27}\alpha^{2} \Delta_{31}^{3} \tan \Delta_{31}$$

$$O(\lambda^{2})$$

$$A_{\mu\mu}^{SM_{1}} = \frac{4}{3}r\alpha \Delta_{31} \cos \delta(\Delta_{31} - \tan \Delta_{31}) - \frac{8}{27}\alpha^{2} \Delta_{31}^{3} \tan \Delta_{31}$$

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$$A_{\mu\mu}^{SM_{1}} = \frac{4}{3}r\alpha \Delta_{31} \cos \delta(\Delta_{31} - \tan \Delta_{31}) - \frac{8}{27}\alpha^{2} \Delta_{31}^{3} \tan \Delta_{31}$$

$$O(\lambda^{2})$$

$$A_{\mu\mu}^{SM_{1}} = 0 + \lambda^{2}$$

NSI Oscillation Probabilities

Oscillations with Neutral Current NSI: •

$$\begin{aligned} |\nu_{\alpha}^{s}\rangle &= |\nu_{\alpha}\rangle + \sum_{\beta=e,\mu,\tau} \varepsilon_{\alpha\beta}^{s} |\nu_{\beta}\rangle \\ \langle\nu_{\beta}^{d}| &= \langle\nu_{\beta}| + \sum_{\alpha=e,\mu,\tau} \varepsilon_{\alpha\beta}^{d} \langle\nu_{\alpha}|. \end{aligned} P\left(\nu_{\alpha}^{s} \rightarrow \nu_{\beta}^{d}\right) = \left|\langle\nu_{\beta}^{d}|e^{-i(H+V_{NSI})L}|\nu_{\alpha}^{s}\rangle\right|^{2} \end{aligned} V_{NSI} = \sqrt{2}G_{F}N_{e}\begin{pmatrix}\varepsilon_{ee}^{m} & \varepsilon_{e\mu}^{m} & \varepsilon_{e\tau}^{m} \\ \varepsilon_{e\mu}^{m*} & \varepsilon_{\mu\mu}^{m} & \varepsilon_{\mu\tau}^{m} \\ \varepsilon_{e\tau}^{m*} & \varepsilon_{\mu\tau}^{m*} & \varepsilon_{\tau\tau}^{m}\end{pmatrix} \end{aligned}$$

$$P(\nu_{\alpha}^{s} \rightarrow \nu_{\beta}^{d}) = \left| \left[(1 + \epsilon^{d})^{T} e^{-i(H + V_{NSI})L} (1 + \epsilon^{s})^{T} \right]_{\beta \alpha} \right|^{2}$$

three more possible sources of CP violation: $\delta_{e \mu}$, $\delta_{e \tau}$, $\delta_{\mu \tau}$

Existing bounds •

Blennow, Choubey, Ohlsson, Pramanik and Raut, JHEP08 (2016), 090 Biggio, Blennow, and Fernandez-Martinez, JHEP08, 090 (2009), 0907.0097

from $G_{_{\rm F}}$, pion decay, unitarity of CKM, oscillation experiments

$$|\varepsilon_{\alpha\beta}^{s/d}| < \begin{bmatrix} 0.041 & 0.025 & 0.041 \\ 0.026 & 0.078 & 0.013 \\ 0.12 & 0.018 & 0.13 \end{bmatrix}$$

mainly from neutrino-electron scattering and neutrino oscillations

$$\varepsilon_{\alpha\beta}^{m}| < \begin{bmatrix} 4.2 & 0.3 & 3.0\\ 0.3 & - & 0.04\\ 3.0 & 0.04 & 0.15 \end{bmatrix}$$

CP asymmetries - Adding the NSI contributions

by construction, the new terms are all proportional to V_{CC}

$$A_{\alpha\beta} = A_{\alpha\beta}^{SM_0} + V_{CC} A_{\alpha\beta}^{SM_1} + V_{CC} A_{\alpha\beta}^{NSI}$$

first order dependence on NSI parameters but suppressed by matter effects \rightarrow globally is an O(λ) correction

$$A_{\mu\tau}^{NSI} = 8 \epsilon_{\mu\tau} \cos \delta_{\mu\tau} \Delta_{31} \cot \Delta_{31} + G[\epsilon_{e\mu}, \epsilon_{e\tau}, \epsilon_{\mu\tau}, \epsilon_{\tau\tau}]$$
first order correction on NSI parameters
$$A_{\mu\mu}^{NSI} = -8 \epsilon_{\mu\tau} \cos \delta_{\mu\tau} \Delta_{31} \tan \Delta_{31} + H[\epsilon_{e\mu}, \epsilon_{e\tau}, \epsilon_{\mu\tau}, \epsilon_{\tau\tau}]$$
SM
$$SM + NSI$$

$$A_{\mu e} \approx 1 + \lambda$$

$$A_{\mu e} \approx 1 + \lambda$$

$$A_{\mu r} \approx \lambda^{2} + \lambda^{2}$$

$$A_{\mu r} \approx \lambda^{2} + \lambda^{2}$$

$$A_{\mu \mu} \approx 0 + \lambda^{2}$$

$$A_{\mu \mu} \approx 0 + \lambda^{2}$$

CP asymmetries - Adding the NSI contributions

Integrated asymmetries: working with number of events in DUNE



- scatter plot: asymmetries obtained varying NSI parameters
- square = error bars from statistics + systematic errors



3+1 sterile neutrino model

• Three more rotation in the definition of the PMNS mixing matrix

 $U = R(\theta_{34})R(\theta_{24})R(\theta_{23},\delta_2)R(\theta_{14})R(\theta_{13},\delta_3)R(\theta_{12},\delta_1)$

two more possible sources of CP violation: δ_1 and δ_2

three small new mixing angles: θ_{14} , θ_{24} , θ_{34}

Hamiltonian

matter potential from NC interactions

$$\begin{bmatrix} \frac{1}{2E_{\nu}}U \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & \Delta m_{21}^2 & 0 & 0 \\ 0 & 0 & \Delta m_{31}^2 & 0 \\ 0 & 0 & 0 & \Delta m_{41}^2 \end{pmatrix} U^{\dagger} + \begin{pmatrix} A_{CC} + A_{NC} & 0 & 0 & 0 \\ 0 & A_{NC} & 0 & 0 \\ 0 & 0 & 0 & A_{NC} & 0 \\ 0 & 0 & 0 & 0 & 0 \end{pmatrix}$$

CP asymmetries in the 3+1 neutrino model

$$A_{\alpha\beta} = A^{SM}_{\alpha\beta} + A^{3+1}_{\alpha\beta}$$



No dependence on the 34 rotation

$$A_{\mu\mu}^{3+1} = O(\lambda^2) = 4 s_{24} s_{34} V_{NC} \tan \Delta_{31} \Delta_{31} \cos \delta_3$$

$$A_{\mu\tau}^{3+1} = O(\lambda^2) = 2s_{24}s_{34}\cot\Delta_{31}(\sin\delta_3 - 2V_{NC}\Delta_{31}\cos\delta_3)$$

significantly large contributions from sterile mixing 34, barely constrained



 $A_{\mu e}$ and $A_{\mu \tau}$ show the best sensitivity to sterile neutrinos

Conclusions

- On-going and planned neutrino experiments will probe the PMNS with huge precision
- Good chance to investigate <u>New Physics effects</u> in Neutrino oscillations:

several "Beyond the Standard Model" scenarios, including sterile neutrinos and Non-Standard Interactions

• the μ -e CP asymmetry shows the best sensitivity to both NP scenarios (with $\mu \mu$ also good for NSI and $\mu \tau$ for 3+1)

Asymmetry	\mathbf{SM}	NSI	$^{3+1}$
$A_{\mu e}$	1	$\lambda^{}$	λ
$A_{\mu\mu}$	λ^3	λ^2	λ^2
$A_{\mu\tau}$	λ^2	λ^2	λ^2
$A_{\mu s}$	-	-	1



Non-standard Neutrino Interactions (NSI)

- A matter NSI operator is induced in fermionic seesaw models once the heavy fermions (singletsor triplets) are integrated out leading to a d= 6 operator that modifies the neutrino kinetic energy.
- After a transformation to obtain canonical kinetic terms, modified couplings of the leptons to the gauge bosons, characterized by deviations from unitarity of the leptonic mixing matrix, are induced.
- Upon integrating out the gauge bosons with their modified couplings, NSI operators are therefore obtained.

SU(2) formulation

- Large NSI could be generated by some other new physics at an energy above the electroweak scale. As a consequence, an SU(2) gauge invariant formulation of NSI is mandatory
- However, in that case, strong bounds stemming from four-charged fermion processes would apply
- In order to avoid these constraints, cancellations among different higher-dimensional operators are required

Non-standard Neutrino Interactions (NSI)

• Many new-physics parameters, huge parameter space:



arbitrary complex matrices



hermitean complex matrices

there exists arguments to reduce the parameter space

- > for the non-standard matter effects, only coupling to electrons, up quarks, and down quarks is important
- non-standard couplings involving τ leptons are irrelevant in reactor and beam sources since τ-production is impossible
- > for I_{α} = e, all corresponding ε 's are vanishing in superbeams because of no-e production
- > in Superbeam source and detector: f=u, f'=d.

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Introducing DUNE

"Deep Underground Neutrino Experiment"

- 1300 km baseline
- Large (70 kt) LArTPC far detector
- 1.5 km underground
- Near Detector (ND) w/LAr component

"Physics goals"

- v and v oscillations (δ_{CP} , θ_{13} , θ_{23} , ordering of nu masses)
- Supernova burst neutrinos
- Beyond Standard Model processes



DUNE events

- <u>neutrino signal channels:</u>
- ν_{e} appearance and ν_{μ} disappearance channels (2% and 5% systematic normalization errors)

T. Alionet al[DUNE Collaboration], arXiv:1606.09550 [physics.ins-det]

	Background	Normalization Uncertainty	Correlations		
	For $\nu_e/\bar{\nu}_e$ appe	earance:			
inels	Beam ν_e	5%	Uncorrelated in $ u_e$ and $ u_e$ samples		
	NC	5%	Correlated in $ u_e$ and $ u_e$ samples		
	ν_{μ} CC	5%	Correlated to NC		
	ν_{τ} CC	20%	Correlated in $ u_e$ and $\bar{ u}_e$ samples		
	For $\nu_{\mu}/\bar{\nu}_{\mu}$ disa	ppearance:			
et]	NC	5%	Uncorrelated to $ u_e/ar{ u}_e$ NC background		
	$\nu_{ au}$	20%	Correlated to $ u_e/ar u_e u_ au$ background		
electron mode - 6% overall det - signal-to-backo - signal systema		- 6% overall det - signal-to-backg - signal systema	ection efficiency for the signal ground ratio of 2.45 tic uncertainty of 20%		
hadro	nic mode	- we take into ac τ-s are detecte	- we take into account that only 30% of the τ-s are detected		

- 0.5% of the NC events as a background
- overall 90% signal detection efficiency
- systematic uncertainty at 10%
- backgrounds come from the mis-identification of CC events (mainly a conservative 10% of the v_{μ} and $v_{e}^{\ CC}$ events)
- <u>neutral current events</u> (hadronic shower with a certain visible energy)

ν<u>appearance channel</u>

Flux options in DUNE

M. Bishai and M. Dolce, *Optimization of the LBNF/DUNE beamline for tau neutrinos*, in Document Database (DocDB) for DUNE and LBNF [http://docs.dunescience.org/cgi-bin/RetrieveFile?docid=2013&filename=DOLCE_M_report.pdf&version=1].



standard

ν mode]	$\bar{\nu}$ mode	
ν_{τ} Signal	277		ν_{τ} Signal	68
$\bar{\nu}_{\tau}$ Signal	26		$\bar{\nu}_{\tau}$ Signal	85
Total Signal	303	1	Total Signal	153
$\nu_e + \bar{\nu}_e$ CC Bkg (beam)	333 + 38		$\nu_e + \bar{\nu}_e$ CC Bkg (beam)	117 + 104
$\nu_e + \bar{\nu}_e$ CC Bkg (oscillation)	1753 + 12		$\nu_e + \bar{\nu}_e$ CC Bkg (oscillation)	90 + 188

A factor of ~10 more tau events

optimized

u mode	
ν_{τ} Signal	2673
$\bar{ u}_{ au}$ Signal	34
Total Signal	2707
$\nu_e + \bar{\nu}_e$ CC Bkg (beam)	688 + 63
$\nu_e + \bar{\nu}_e$ CC Bkg (oscillation)	1958 + 11

$\bar{\nu}$ mode	
ν_{τ} Signal	98
$\bar{\nu}_{\tau}$ Signal	983
Total Signal	1081
$\nu_e + \bar{\nu}_e$ CC Bkg (beam)	176 + 177
ν_e CC Bkg (oscillation)	76 + 324

Effects of NSI on the event spectra



Dotted blue lines: SM events

Red bands: range of events while varying matter NSI parameters in their allowed ranges



Effects of a sterile state on the event spectra





Red bands: range of events while varying 3+1 sterile parameters in their allowed ranges



 $\nu_{\mu} \rightarrow \nu_{e} ~~and ~\nu_{\mu} \rightarrow \nu_{\tau}$ are the most sensitive channels

