

Role of parton fragmentation for associated J/ψ production at high energies

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Outline

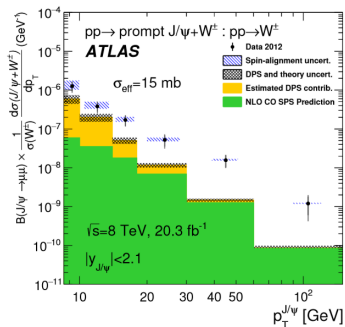
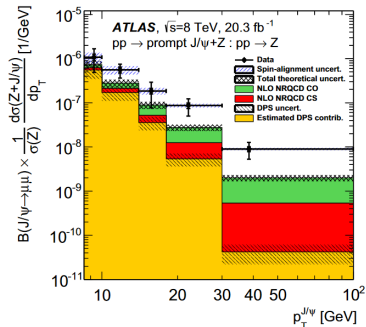
- ▶ Introduction
- ▶ Motivation and goals
- ▶ New contributions to the $J/\psi + Z/W$
- ▶ Fragmentation to the charmonium
- ▶ Modelling events
- ▶ Comparison with experimental data and summary

- ▶ **Non-relativistic QCD** (NRQCD):

$$\sigma(pp \rightarrow J/\psi + X) = \sum_n \hat{\sigma}(pp \rightarrow c\bar{c}(^{2S+1}L_J^{[a]} + X) \langle \mathcal{O}^{J/\psi}[n] \rangle$$

- ▶ $\hat{\sigma}(pp \rightarrow c\bar{c}(^{2S+1}L_J^{[n]} + X)$ is the cross section of production unbound $c\bar{c}$ pair at the Fock state $n = ^{2S+1}L_J^{(a)}$ with definite spin S , orbital angular momentum L , total angular momentum J and color representation a (color singlet (CS) [1] and color octet (CO) [8]) - can be calculated in the framework of pQCD
- ▶ **LDME** (long distance matrix element) $\langle \mathcal{O}^{J/\psi}[n] \rangle$ corresponds to transition from unbound state to the physical J/ψ meson - nonperturbative part.
- ▶ Progress in NRQCD evaluation of prompt $J/\psi + Z/W^\pm$ production: complete NLO calculations [Phys. Rev. D66, 114002 (2002)], [Phys. Rev. D83, 014001 (2011)], [JHEP02, 071 (2011)]; differential cross sections at the LO are significantly enhanced by the NLO corrections.

Introduction



Complete NLO NRQCD predictions with the double parton scattering (DPS) underestimate the latest ATLAS [Eur.Phys.J.C. 75, 229 (2015)] and ATLAS [J.High.Energ.Phys. 2020,95 (2020)] data by the factor 2 - 10 (depending on the J/ψ transverse momentum).

Motivation and goals

- ▶ We consider the new contributions to the prompt $J/\psi + Z/W^\pm$ production: **flavor excitation subprocesses** (charm for Z boson and strange for W) followed by the subsequent **charm fragmentation**, $c \rightarrow J/\psi + c$. **Our goal is to estimate such contributions**
- ▶ Recently we found that contribution of multiple gluon radiation to the cross section of double J/ψ production are very important [Eur. Phys. J. C80,1046 (2020)]. The multiple gluon radiation can be taken into account using the CCFM evolution equation. One can expect a sizeable contribution from **multiple gluon radiation** for $J/\psi + Z/W^\pm$ processes.
- ▶ **Our goal is to investigate a role of multiple gluon fragmentation to the prompt $J/\psi + Z/W^\pm$ production**

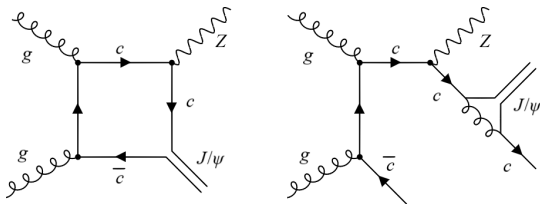
k_T -factorization approach

- ▶ We use the k_T -factorization approach with CCFM-evolved (Catani, Ciafaloni, Fiorani, Marchesini) Transverse Momentum Dependent (TMD) gluon densities
- ▶ Cross section in k_T -factorization approach:

$$d\sigma(pp \rightarrow J/\psi + Z/W) = \int dx_1 dx_2 \sum_{i,j} d^2\vec{k}_{\perp 1} d^2\vec{k}_{\perp 2} f_i(x_1, \vec{k}_{\perp 1}^2 \mu^2) f_j(x_2, \vec{k}_{\perp 2}^2 \mu^2) \cdot d\hat{\sigma}(i^* + j^* \rightarrow J/\psi + Z/W)$$

- ▶ $f_{i,j}(x_{1,2}, \vec{k}_{\perp 1,2}, \mu^2)$ - TMD parton distribution functions (TMD PDF) in a proton obeying the BFKL or CCFM evolution equation
- ▶ $d\hat{\sigma}(i^* + j^* \rightarrow J/\psi + Z/W)$ - off-shell partonic cross section

Flavor excitation



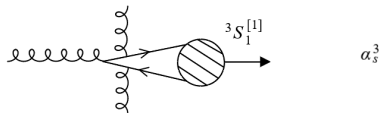
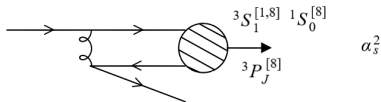
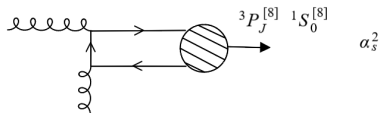
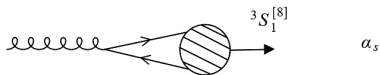
- ▶ Examples of Feynman diagram taken into account in the NRQCD calculations (left panel) and diagram of **charm excitation followed by the c-quark fragmentation to J/ψ** (right panel)
- ▶ Since the charm quark contribution can be obtained via gluon splitting ($g \rightarrow q_s \bar{q}_s$) for CCFM evolved gluon densities, the processes of flavor excitation: $g + c \rightarrow Z + c$, $g + s \rightarrow W^- + c$ turn to gluon-gluon fusion $g + g \rightarrow Z + c + \bar{c}$, $g + g \rightarrow W^- + c + \bar{s}$
- ▶ Gluon-gluon fusion followed by the fragmentation $c \rightarrow J/\psi + c$

Fragmentation to the charmonium \mathcal{H}

- ▶ We consider not only the direct production of J/ψ but also the feeddown contribution from radiative decay of $\psi' \rightarrow J/\psi X$ and $\chi_{cJ} \rightarrow J/\psi \gamma$
- ▶ Fragmentation function in NRQCD formalism at the starting scale $\mu_0^2 = m_{\mathcal{H}}^2$:

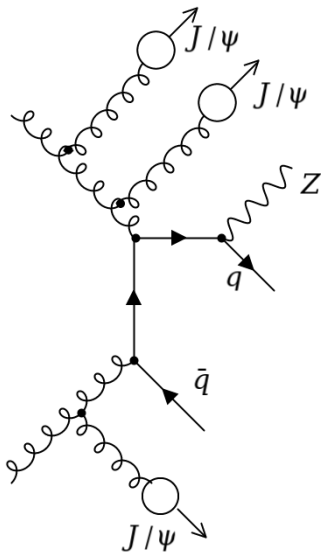
$$D_a^{\mathcal{H}}(z, \mu_0^2) = \sum_n d_a^n(z, \mu_0^2) \langle \mathcal{O}^{\mathcal{H}}[n] \rangle$$

- ▶ Typical diagrams of gluons and charm quarks fragmentation into charmonium

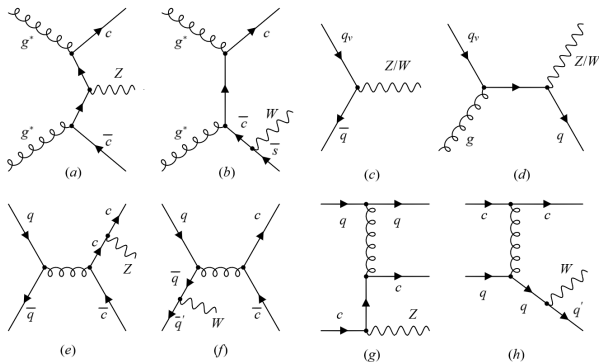


Multiple gluon radiation

- ▶ Additional contribution comes from **multiple initial gluon radiation** that accompanies the Z/W production
- ▶ Initial gluon cascade can be described by the CCFM evolution equation
- ▶ Subprocesses $g + g \rightarrow Z + q + \bar{q}$,
 $g + g \rightarrow W + q + \bar{q}'$ give additional contribution via fragmentation
 $g \rightarrow cc(^3S_1^{[8]}) \rightarrow J/\psi$
- ▶ Circles on the plot denote the possible channels of partons fragmentation into J/ψ mesons



List of considered subprocesses



- ▶ Gluon-gluon fusion (a)-(b) are calculated in k_T -factorization approach QCD. The initial multiple gluon radiation can be taken into account using the CCFM evolved gluon densities
- ▶ Quark-involved subprocesses (c)-(h) are calculated in collinear QCD. The initial multiple gluon radiation are reconstructed with PYTHIA routine. Subprocesses (c)-(d) involve only valence quarks (sea quark effectively included in gluon-gluon fusion)

J/ψ production via fragmentation

- ▶ We took only **LO contributions to the FFs**: $D_g^{\mathcal{H}}({}^3S_1^{[8]})$ and $D_c^{\mathcal{H}}({}^3S_1^{[1]})$ for $J/\psi, \psi'$; $D_g^{\mathcal{H}}({}^3S_1^{[8]})$ and $D_c^{\mathcal{H}}({}^3P_J^{[1]})$ for χ_{cJ} . Charm fragmentation into octet color states suppressed due to color factor.
- ▶ **LO DGLAP evolution equation** \Rightarrow FFs $D_c^{\mathcal{H}}(z, \mu^2)$ and $D_g^{\mathcal{H}}(z, \mu^2)$ at the any scale μ^2

$$\frac{d}{d \log \mu^2} \begin{pmatrix} D_c \\ D_g \end{pmatrix} = \frac{\alpha_s(\mu^2)}{2\pi} \begin{pmatrix} P_{cc} & P_{gc} \\ P_{cg} & P_{gg} \end{pmatrix} \otimes \begin{pmatrix} D_c \\ D_g \end{pmatrix}$$

where P_{ab} standard LO DGLAP splitting function

- ▶ Cross section of $J/\psi + Z/W$ production via charm fragmentation can be written:

$$\frac{d\sigma(pp \rightarrow J/\psi + Z/W)}{dp_T} = \int dz \frac{d\hat{\sigma}(pp \rightarrow c + Z/W)}{dp_T^c} D_c^{J/\psi}(z, \mu^2) \delta\left(z - \frac{p}{p^c}\right)$$

Modelling events

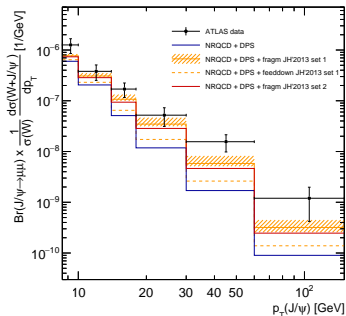
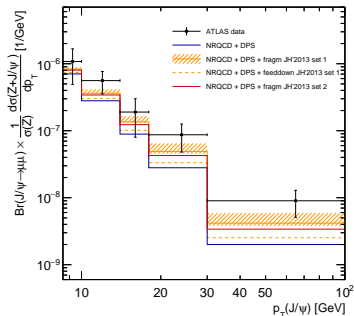
We used:

- ▶ **JH'2013 set1 and set2 TMD gluon densities** for numerical calculations of gluon-gluon fusion subprocesses in k_T -factorization approach; Monte Carlo event generator CASCADE for reconstruction of CCFM initial gluon emissions
- ▶ **MMHT2014LO PDF** for numerical calculation of quark-involved subprocesses in collinear QCD; PYTHIA routine for initial gluon emissions
- ▶ **numerical solution of DGLAP evolution of FFs** with appropriate LDME's
 $\langle \mathcal{O}^{\mathcal{H}}[n] \rangle$ (list of used LDME: $\langle \mathcal{O}^{J/\psi} [{}^3S_1^{(1)}] \rangle = 1.16 \text{ GeV}^3$, $\langle \mathcal{O}^{\psi'} [{}^3S_1^{(1)}] \rangle = 0.7038 \text{ GeV}^3$, $\langle \mathcal{O}^{\chi_{c1}} [{}^3P_1^{(1)}] \rangle = 0.2 \text{ GeV}^5$, $\langle \mathcal{O}^{\chi_{c2}} [{}^3P_2^{(1)}] \rangle = 0.0496 \text{ GeV}^5$, $\langle \mathcal{O}^{J/\psi, \psi'} [{}^3S_1^{(8)}] \rangle = 0.0012 \text{ GeV}^3$, $\langle \mathcal{O}^{\chi_{c0}} [{}^3S_1^{(8)}] \rangle = 0.0004 \text{ GeV}^3$)

Selection criteria:

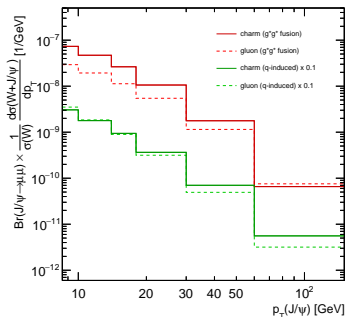
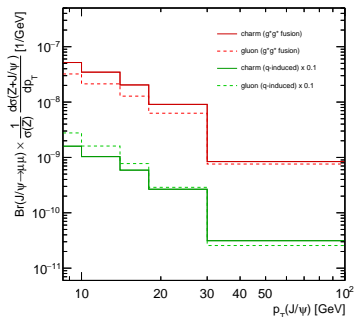
- ▶ $J/\psi + Z$: $p_T(J/\psi) > 8.5 \text{ GeV}$, $|y(J/\psi)| < 2.1$, $M(Z) = 81 \div 101 \text{ GeV}$
lead l : $p_T > 25 \text{ GeV}$, $|\eta| < 2.5$; sublead l : $p_T > 15 \text{ GeV}$, $|\eta| < 2.5$
- ▶ $J/\psi + W$: $p_T(J/\psi) > 8.5 \text{ GeV}$, $|y(J/\psi)| < 2.1$,
 $m_T(W^\pm) = \sqrt{2p_T^l p_T^\nu [1 - \cos(\phi^l - \phi^\nu)]} > 40 \text{ GeV}$
muon l : $p_T > 25 \text{ GeV}$, $|\eta| < 2.4$; neutrino ν : $p_T > 20 \text{ GeV}$, $|\eta| < 2.4$

Comparison with ATLAS data



- ▶ Contribution from considered subprocesses with their subsequent parton fragmentation into J/ψ mesons (**fragmentation contribution**) are remarkably important, especially at large transverse momenta (at $p_T^{J/\psi} \geq 20\text{-}30$ GeV it gives approximately the same contribution as NLO NRQCD + DPS)
- ▶ Feeddown contribution from radiative decay of ψ' , χ_{cJ} also play the significant role (about 30% of the estimated direct contribution at the wide $p_T^{J/\psi}$ range).
- ▶ Shaded bands represents the scale uncertainties.

Role of the multiple gluon radiation



- ▶ We consider the two qualitatively different sources of parton fragmentation into the J/ψ meson: **fragmentation of charm quarks**, originated in the hard interaction, and **fragmentation of gluons**, originated as a result of initial QCD evolution of parton densities
- ▶ In both cases of gluon-gluon fusion and quark-involved subprocesses the fragmentation of multiple gluon emission noticeably enhances the charm fragmentation and provides a sensible growth of the total and differential cross sections (especially at the region of high $p_T^{J/\psi}$)

Summary

- ▶ We investigated the role of new partonic subprocesses which yet have never been considered in the literature, namely, flavor (charm or strangeness) excitation subprocesses followed by the charm fragmentation $c \rightarrow J/\psi + c$.
- ▶ We take into account the effects of the multiple quark and gluon radiation in the initial and final states.
- ▶ Contributions from multiple gluon emissions noticeably enhance the charm fragmentation (especially at the region of high $p_T^{J/\psi}$).
- ▶ Accounting for the feeddown contribution from radiative decay of ψ', χ_{cJ} also play a significant role (about 30% of the estimated direct contribution at the wide $p_T^{J/\psi}$ range).
- ▶ Considered new contributions are remarkably important and significantly reduce the gap between the theoretical predictions and experimental results for the $J/\psi + Z/W^\pm$

Backup

Evolution of FFs

$$D_{init\ g}^{J/\psi}(z, \mu_0^2) = \frac{a_s(\mu_0^2)\pi}{8m_{J/\psi}^3} \delta(1-z)$$

$$D_{init\ c}^{J/\psi}(z, \mu_0^2) = \frac{a_s^2(\mu_0^2)}{m_{J/\psi}^3} \frac{16z(1-z)^2}{243(2-z)^6} (5z^4 - 32z^3 + 72z^2 - 32z + 16)$$

[Bernd. A.Kniehl, Gustav Kramer Phys.Rev.D. 56, 5820, (1997)]

