

RECENT RESULTS ON EXTRACTIONS OF QUARK TMDs

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RESULTS OBTAINED WITH CONTRIBUTIONS FROM

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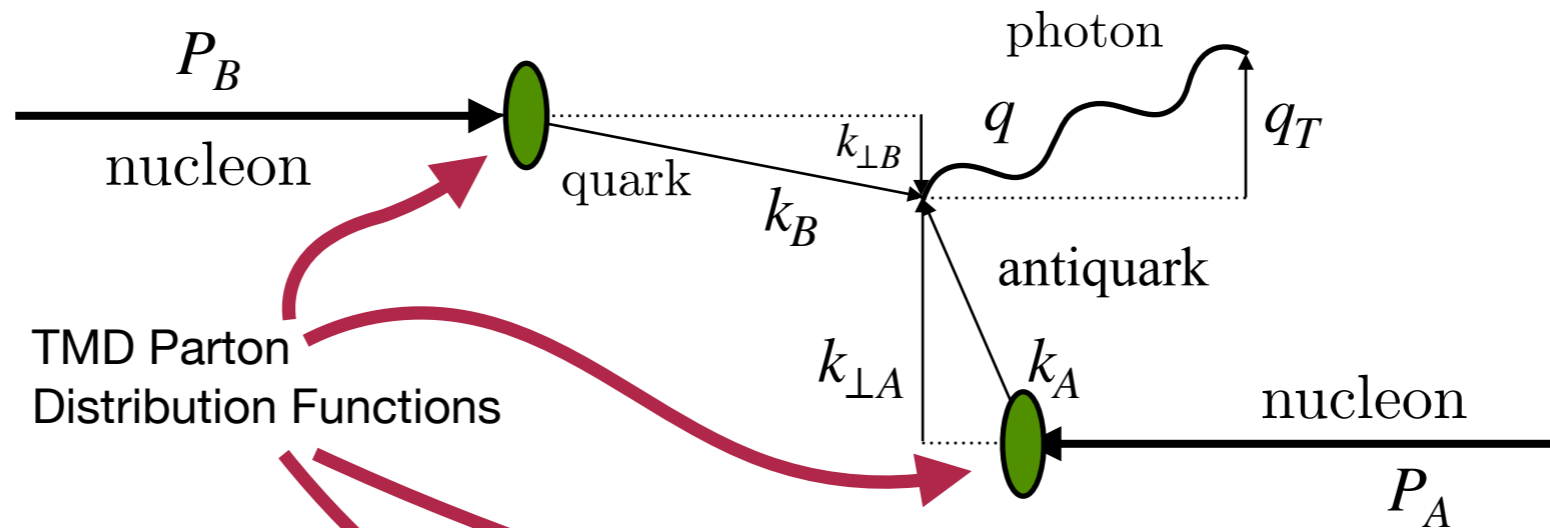


Matteo Cerutti



UNPOLARISED QUARK TMDS

TMDS IN DRELL-YAN PROCESSES



$$F_{UU}^1(x_A, x_B, \mathbf{q}_T^2, Q^2)$$

$$= \sum_a \mathcal{H}_{UU}^{1a}(Q^2, \mu^2) \int d^2\mathbf{k}_{\perp A} d^2\mathbf{k}_{\perp B} f_1^a(x_A, \mathbf{k}_{\perp A}^2; \mu^2) f_1^{\bar{a}}(x_B, \mathbf{k}_{\perp B}^2; \mu^2) \delta^{(2)}(\mathbf{k}_{\perp A} - \mathbf{q}_T + \mathbf{k}_{\perp B})$$

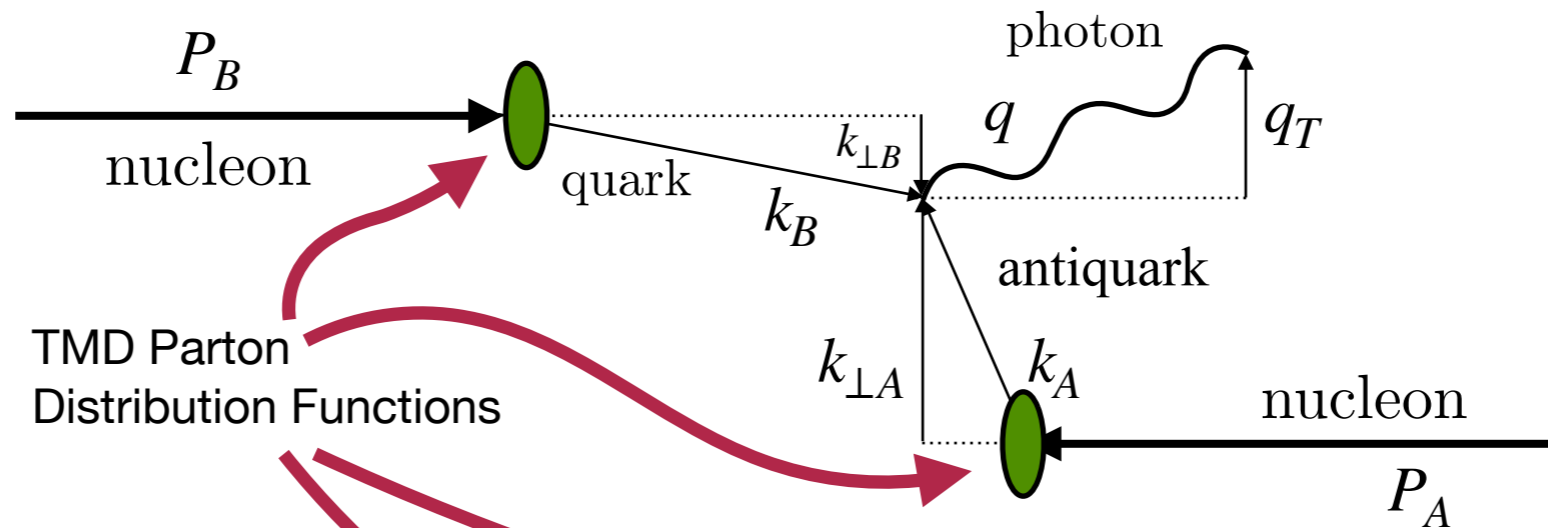
W term

$$+ Y_{UU}^1(Q^2, \mathbf{q}_T^2) + \mathcal{O}(M^2/Q^2)$$

The W term, dominates at low transverse momentum ($q_T \ll Q$)

So far, the Y term has been excluded in the Pavia analyses

TMDS IN DRELL-YAN PROCESSES



$$F_{UU}^1(x_A, x_B, \mathbf{q}_T^2, Q^2)$$

$$\approx \sum_q \mathcal{H}_{UU}^{1q}(Q^2, \mu^2) \int d^2\mathbf{k}_{\perp A} d^2\mathbf{k}_{\perp B} f_1^q(x_A, \mathbf{k}_{\perp A}^2; \mu^2) f_1^{\bar{q}}(x_B, \mathbf{k}_{\perp B}^2; \mu^2) \delta^{(2)}(\mathbf{k}_{\perp A} - \mathbf{q}_T + \mathbf{k}_{\perp B})$$

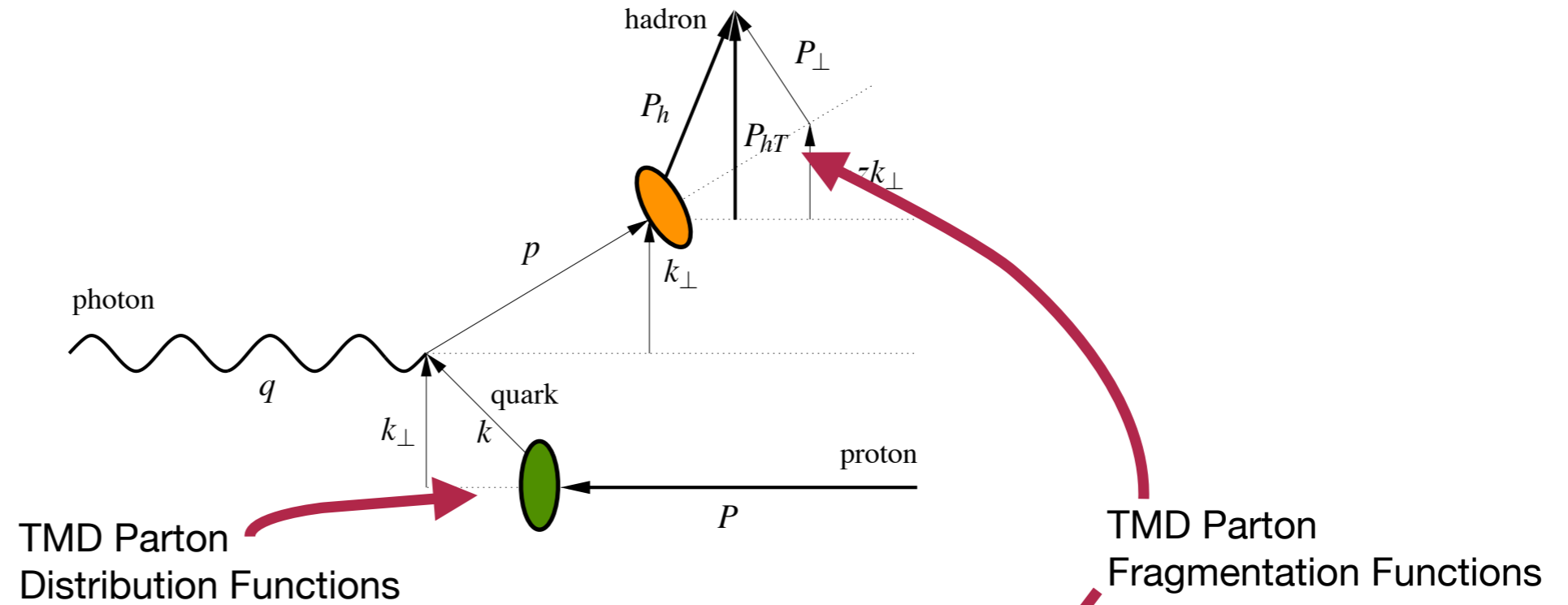
$$= \sum_q \mathcal{H}_{UU}^{1q}(Q^2, \mu^2) \int db_T b_T J_0(b_T |\mathbf{q}_T|) \hat{f}_1^q(x_A, b_T^2; \mu^2) \hat{f}_1^{\bar{q}}(x_B, b_T^2; \mu^2)$$

At small transverse momentum, the dominant part is given by TMDs.

The analysis is usually done in Fourier-transformed space

TMDs formally depend on two scales, but we set them equal.

TMDS IN SEMI-INCLUSIVE DIS



$$F_{UU,T}(x, z, \mathbf{P}_{hT}^2, Q^2)$$

$$= x \sum_q \mathcal{H}_{UU,T}^q(Q^2, \mu^2) \int d^2 \mathbf{k}_{\perp} d^2 \mathbf{P}_{\perp} f_1^a(x, \mathbf{k}_{\perp}^2; \mu^2) D_1^{a \rightarrow h}(z, \mathbf{P}_{\perp}^2; \mu^2) \delta(z \mathbf{k}_{\perp} - \mathbf{P}_{hT} + \mathbf{P}_{\perp})$$

$$= x \sum_a \mathcal{H}_{UU,T}^a(Q^2, \mu^2) \int db_T b_T J_0(b_T |\mathbf{P}_{h\perp}|) \hat{f}_1^a(x, z^2 b_{\perp}^2; \mu^2) \hat{D}_1^{a \rightarrow h}(z, b_{\perp}^2; \mu^2)$$

hard factor

TMD STRUCTURE

$$\hat{f}_1^q(x, b_T; \mu^2) = \int d^2 \mathbf{k}_\perp e^{i \mathbf{b}_T \cdot \mathbf{k}_\perp} f_1^q(x, \mathbf{k}_\perp^2; \mu^2)$$

$$\hat{f}_1^q(x, b_T; \mu^2) = \sum_i (C_{qi} \otimes f_1^i)(x, b_*; \mu_b) e^{\tilde{S}(b_*; \mu_b, \mu)} e^{g_K(b_T) \ln \frac{\mu}{\mu_0}} \hat{f}_{\text{NP}}^q(x, b_T)$$

perturbative Sudakov form factor
nonperturbative part of evolution
nonperturbative part of TMD

collinear PDF

$\mu_b = \frac{2e^{-\gamma_E}}{b_*}$

matching coefficients (perturbative)

see, e.g.,
Collins, "Foundations of Perturbative QCD" (11)

PAVIA 2019 b_* PRESCRIPTION

$$\hat{f}_1^q(x, b_T; \mu^2) = \sum_i (C_{qi} \otimes f_1^i)(x, b_*; \mu_b) e^{\tilde{S}(b_*; \mu_b, \mu)} e^{g_K(b_T) \ln \frac{\mu}{\mu_0}} \hat{f}_{\text{NP}}^q(x, b_T)$$

$$\mu_0 = 1 \text{ GeV}$$

$$b_* \equiv \frac{b_T}{\sqrt{1 + b_T^2/b_{\text{max}}^2}}$$

Collins, Soper, Sterman, NPB250 (85)

$$\mu_b = 2e^{-\gamma_E}/b_* \quad \bar{b}_* \equiv b_{\text{max}} \left(\frac{1 - e^{-b_T^4/b_{\text{max}}^4}}{1 - e^{-b_T^4/b_{\text{min}}^4}} \right)^{1/4} \quad b_{\text{max}} = 2e^{-\gamma_E}$$

$$b_{\text{min}} = \frac{2e^{-\gamma_E}}{Q}$$

These are all choices that should be at some point checked/challenged

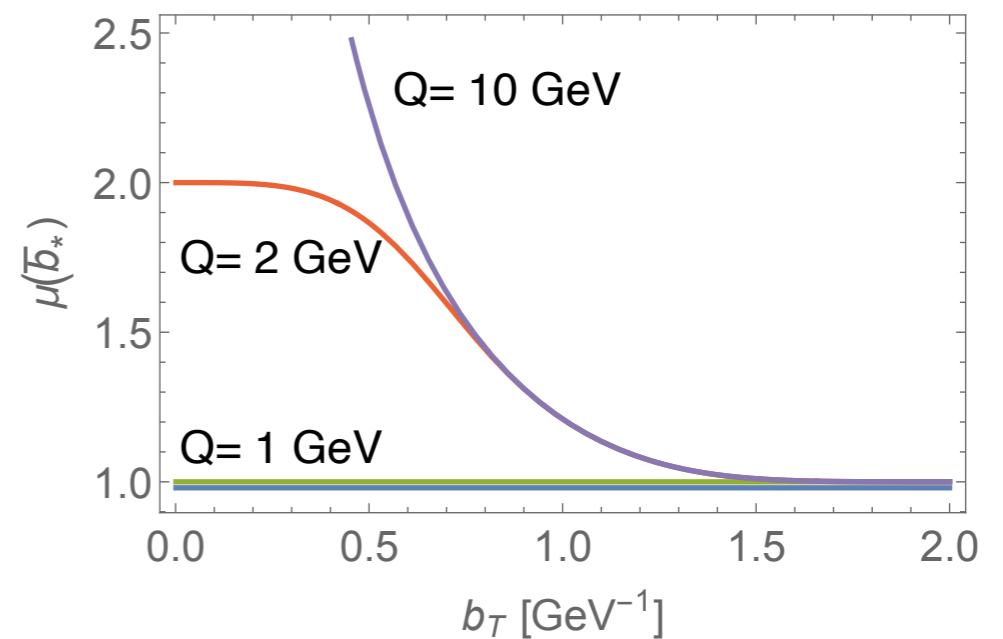
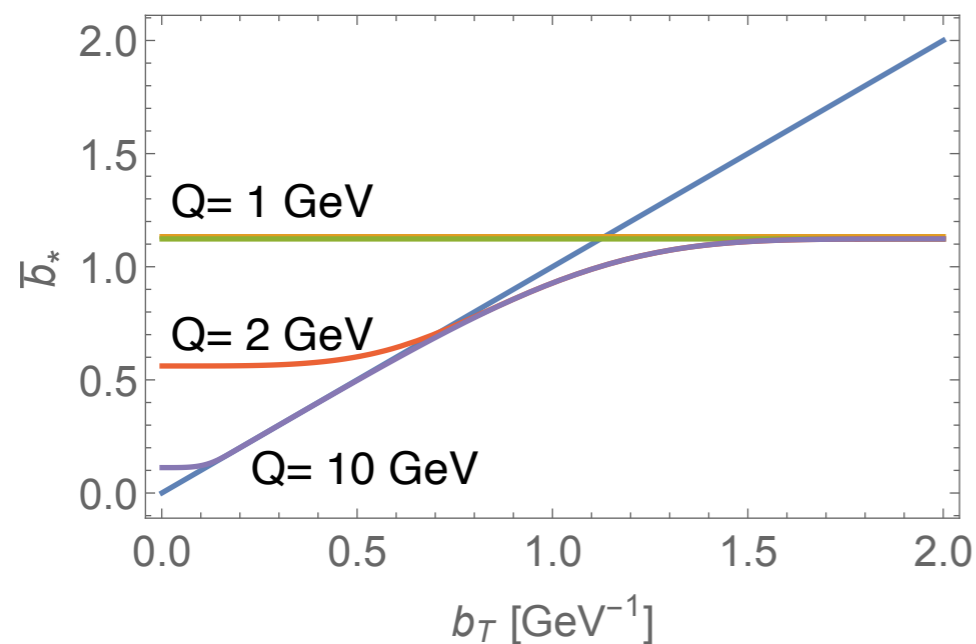
EFFECTS OF b_* PRESCRIPTION

$$\mu_b = 2e^{-\gamma_E} / b_*$$

$$\bar{b}_* \equiv b_{\max} \left(\frac{1 - e^{-b_T^4/b_{\max}^4}}{1 - e^{-b_T^4/b_{\min}^4}} \right)^{1/4}$$

$$b_{\max} = 2e^{-\gamma_E}$$

$$b_{\min} = \frac{2e^{-\gamma_E}}{Q}$$



No significant effect at high Q , but large effect at low Q
(inhibits perturbative contribution)

EFFECTS OF b_* PRESCRIPTION

- At $b_T > b_{\max}$, the perturbative contributions are inhibited and replaced by nonperturbative contributions (roughly speaking, at small transverse momentum)
- At $b_T < b_{\min}$ the TMD is modified, but it should influence the region outside the applicability of TMD factorization (roughly speaking, at high transverse momentum)
- At $Q=1$, only, the integral of the TMD is by construction equal to the collinear PDF (see also talk by Z. Sun)

LOW- b_T MODIFICATIONS

$$\log(Q^2 b_T^2) \rightarrow \log(Q^2 b_T^2 + 1)$$

see, e.g., Bozzi, Catani, De Florian, Grazzini
[hep-ph/0302104](https://arxiv.org/abs/hep-ph/0302104)

$$b_*(b_c(b_T)) = \sqrt{\frac{b_T^2 + b_0^2/(C_5^2 Q^2)}{1 + b_T^2/b_{\max}^2 + b_0^2/(C_5^2 Q^2 b_{\max}^2)}}$$

$$b_{\min} \equiv b_*(b_c(0)) = \frac{b_0}{C_5 Q} \sqrt{\frac{1}{1 + b_0^2/(C_5^2 Q^2 b_{\max}^2)}}$$

Collins et al.
[arXiv:1605.00671](https://arxiv.org/abs/1605.00671)

- The justification is to recover the integrated result (“unitarity constraint”)
- Modification at low b_T is allowed because resummed calculation is anyway unreliable there

RECENT TMD FITS OF UNPOLARIZED DATA

	Framework	HERMES	COMPASS	DY	Z production	N of points	χ^2/N_{points}
Pavia 2017 arXiv:1703.10157	NLL	✓	✓	✓	✓	8059	1.55
SV 2017 arXiv:1706.01473	NNLL'	✗	✗	✓	✓	309	1.23
BSV 2019 arXiv:1902.08474	NNLL'	✗	✗	✓	✓	457	1.17
SV 2019 arXiv:1912.06532	NNLL'	✓	✓	✓	✓	1039	1.06
Pavia 2019 arXiv:1912.07550	N ³ LL	✗	✗	✓	✓	353	1.02

What have we been doing in the last two years



PV17

PV19



MAP21

work still in progress...

NANGA PARBAT: A PUBLIC PLATFORM FOR TMD STUDIES

<https://github.com/MapCollaboration/NangaParbat>



Nanga Parbat: a TMD fitting framework

Nanga Parbat is a fitting framework aimed at the determination of the non-perturbative component of TMD distributions.

Download

You can obtain NangaParbat directly from the github repository:

<https://github.com/vbertone/NangaParbat/releases>

For the last development branch you can clone the master code:

```
git clone git@github.com:vbertone/NangaParbat.git
```

PERTURBATIVE ORDER OF EACH INGREDIENT

Order in powers of α_s

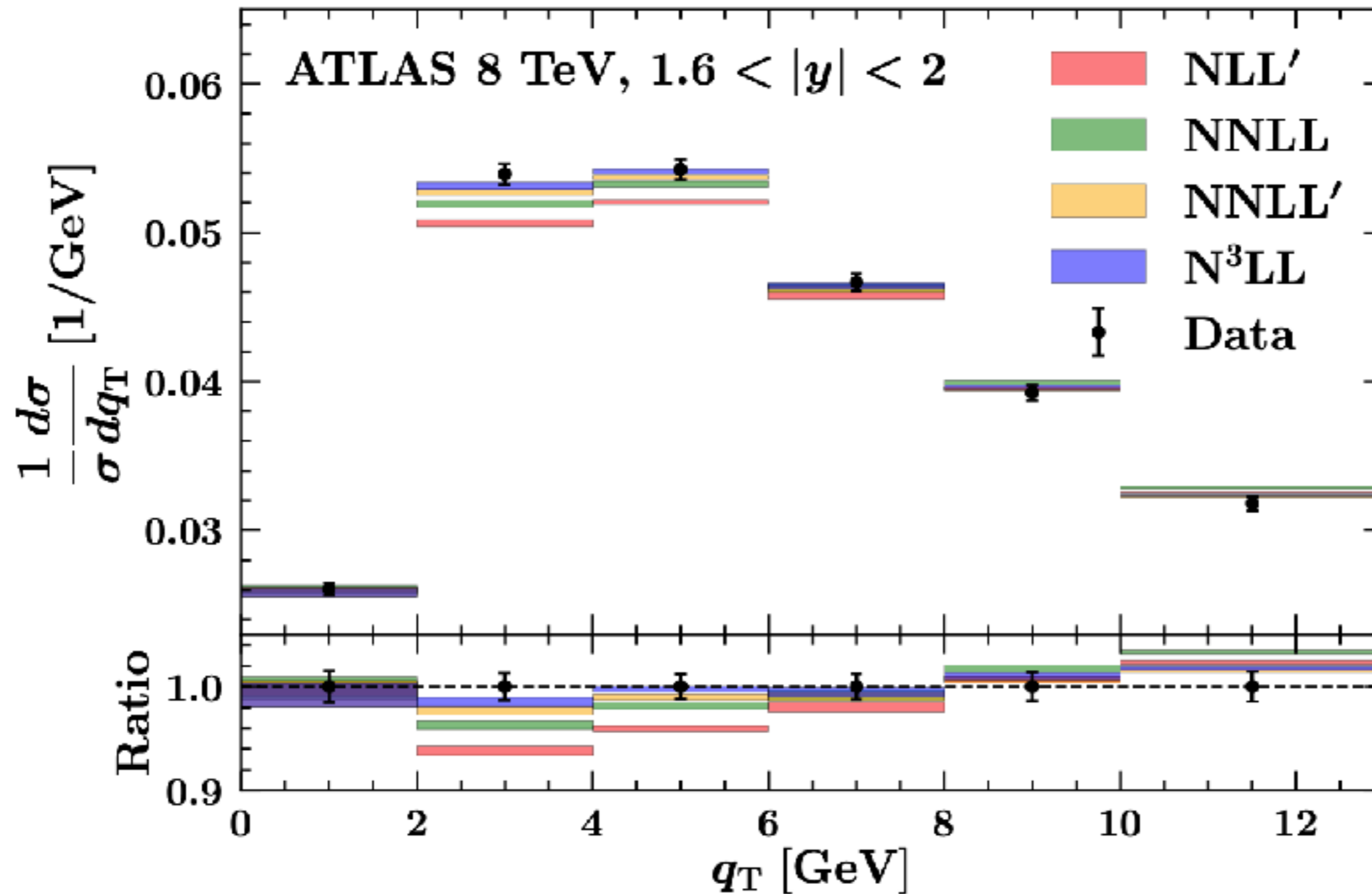
hard factor and matching coefficients

ingredients in perturbative Sudakov form factor

Accuracy	H and C	K and γ_F	γ_K	PDF and α_s evol.
LL	0	-	1	-
NLL	0	1	2	LO
NLL'	1	1	2	NLO
NNLL	1	2	3	NLO
NNLL'	2	2	3	NNLO
N ³ LL-	2	3	4	NLO
N ³ LL	2	3	4	NNLO
N ³ LL'	3	3	4	N ³ LO

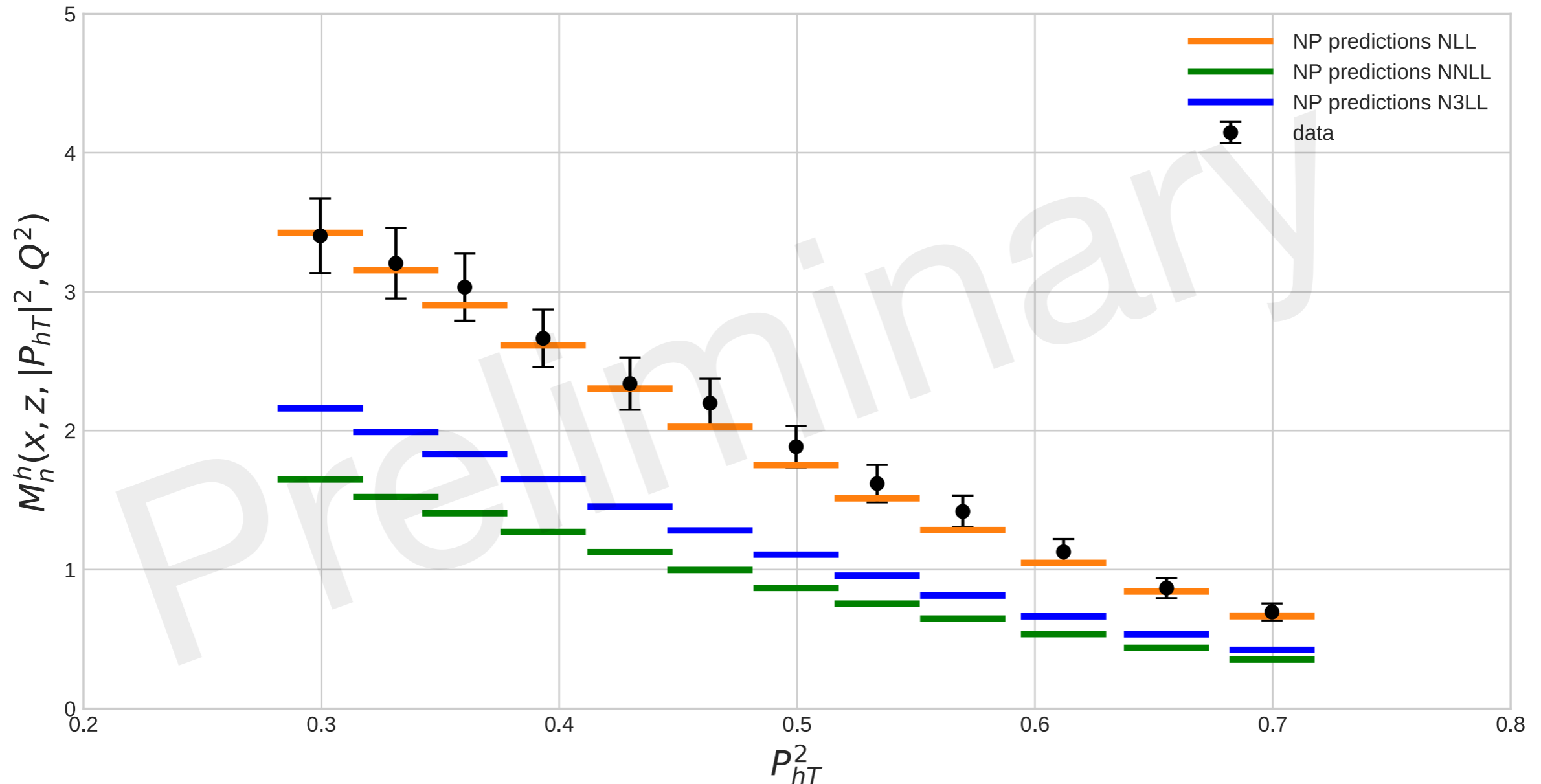
Collinear fragmentation functions not available beyond NLO

COMPARISON OF DIFFERENT ORDERS IN DY



COMPARISON OF DIFFERENT ORDER SIDIS

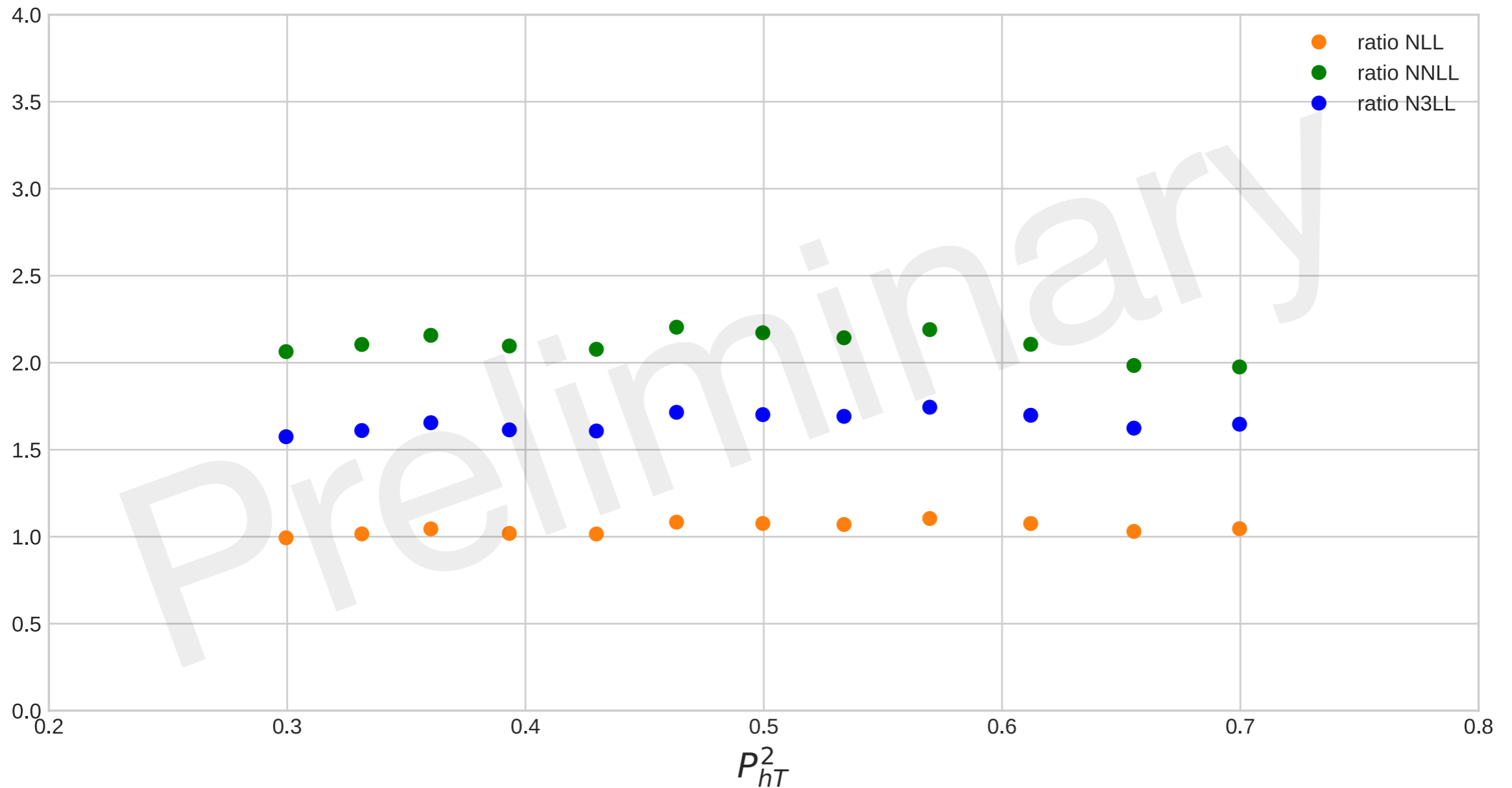
COMPASS multiplicities (one of many bins)



The description considerably worsens at higher orders

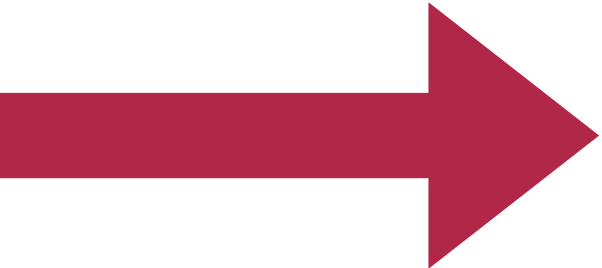
RATIO DATA/PREDICTIONS

COMPASS multiplicities (one of many bins)



The discrepancy amounts to a constant factor

OUR TENTATIVE SOLUTION



PREFACTOR $= \frac{\left. \frac{d\sigma}{dx dz dQ^2} \right|_{\text{nonmix.}}}{\int W d^2 q_T}$

$$\left. \frac{d\sigma^h}{dx dQ^2 dz} \right|_{O(\alpha_S)} = \sigma_0 \sum_{f,f'} \frac{e_f^2}{z^2} (\delta_{f'f} + \delta_{f'g}) \frac{\alpha_S}{\pi} \left\{ [D_1^{h/f'} \otimes C_1^{f'f} \otimes f_1^{f/N}](x, z, Q) \right\} \Big|_{\text{nonmix.}}$$

$$\int W \Big|_{O(\alpha_S)} = \sigma_0 \sum_{f,f'} \frac{e_f^2}{z^2} (\delta_{f'f} + \delta_{f'g}) \frac{\alpha_S}{\pi} [D_1^{h/f'} \otimes C_{\text{TMD}}^{f'f} \otimes f_1^{f/N}](x, z, Q)$$

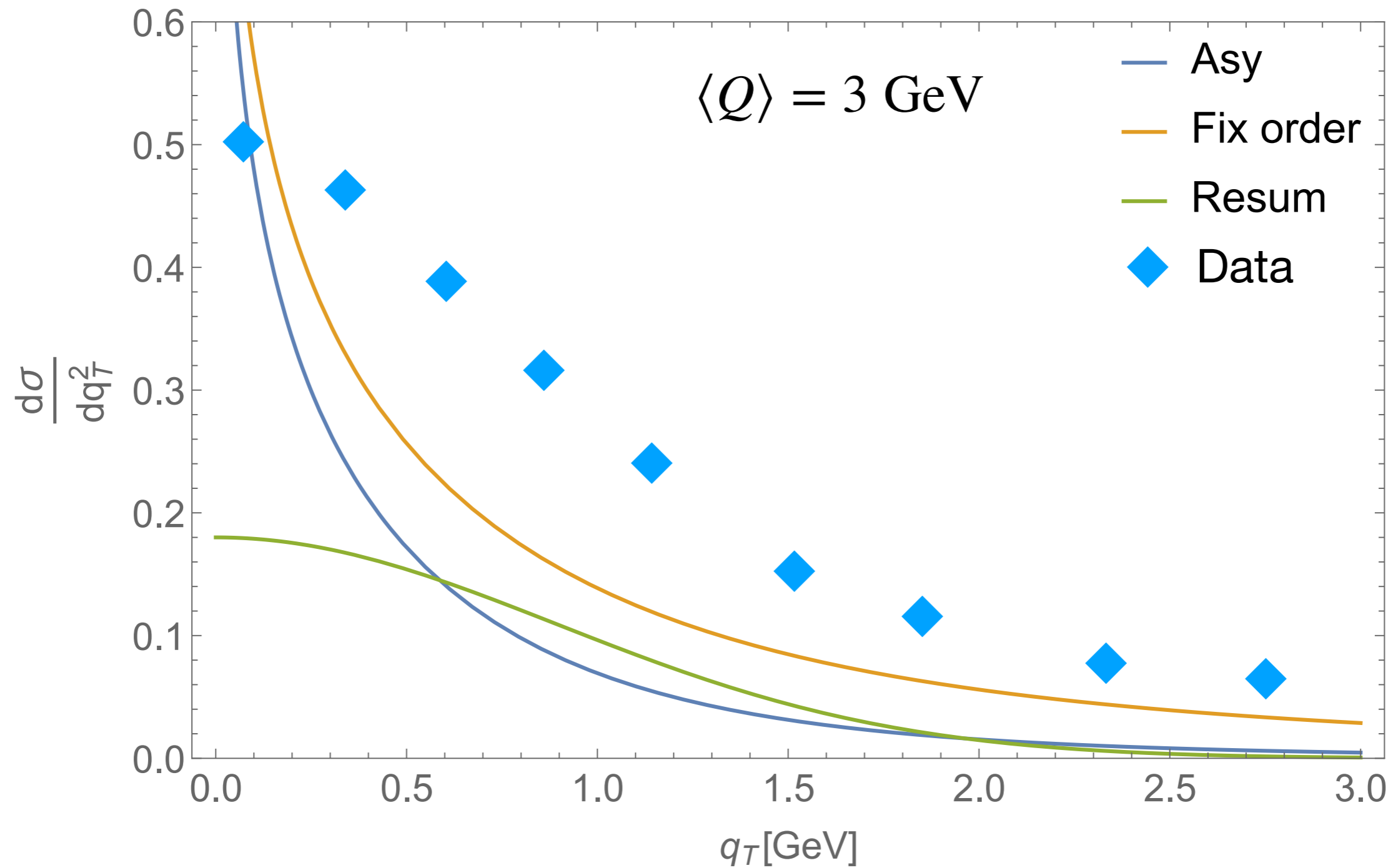
The prefactor is independent of the fitting parameters

NONMIXED TERMS IN COLLINEAR SIDIS CROSS SECTION

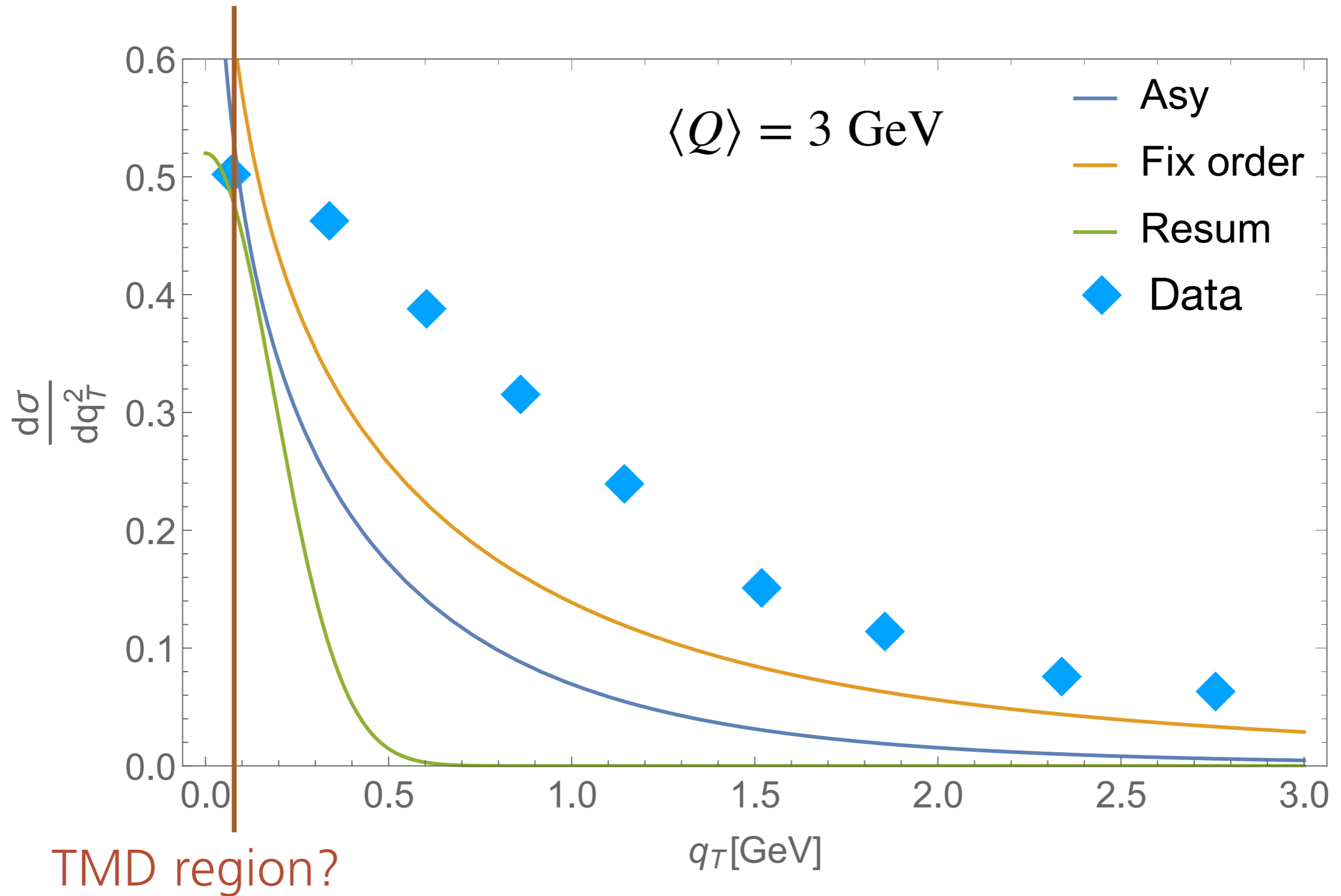
$$\frac{d\sigma^h}{dx dQ^2 dz} \Big|_{O(\alpha_s^1)} = \sigma_0 \sum_{ff'} \frac{e_f^2}{z^2} (\delta_{f'f} + \delta_{f'g}) \frac{\alpha_s}{\pi} \left\{ \left[D_1^{h/f'} \otimes C_1^{f'f} \otimes f_1^{f/N} \right] (x, z, Q) \right. \\ \left. + \frac{1-y}{1+(1-y)^2} \left[D_1^{h/f'} \otimes C_L^{f'f} \otimes f_1^{f/N} \right] (x, z, Q) \right\},$$

$$C_1^{qq} = \frac{C_F}{2} \left\{ -8\delta(1-x)\delta(1-z) \right. \\ \left. + \delta(1-x) \left[P_{qq}(z) \ln \frac{Q^2}{\mu_F^2} + L_1(z) + L_2(z) + (1-z) \right] \right. \\ \left. + \delta(1-z) \left[P_{qq}(x) \ln \frac{Q^2}{\mu^2} + L_1(x) - L_2(x) + (1-x) \right] \right. \\ \left. + 2 \frac{1}{(1-x)_+} \frac{1}{(1-z)_+} - \frac{1+z}{(1-x)_+} \frac{1+x}{(1-z)_+} + 2(1+xz) \right\},$$

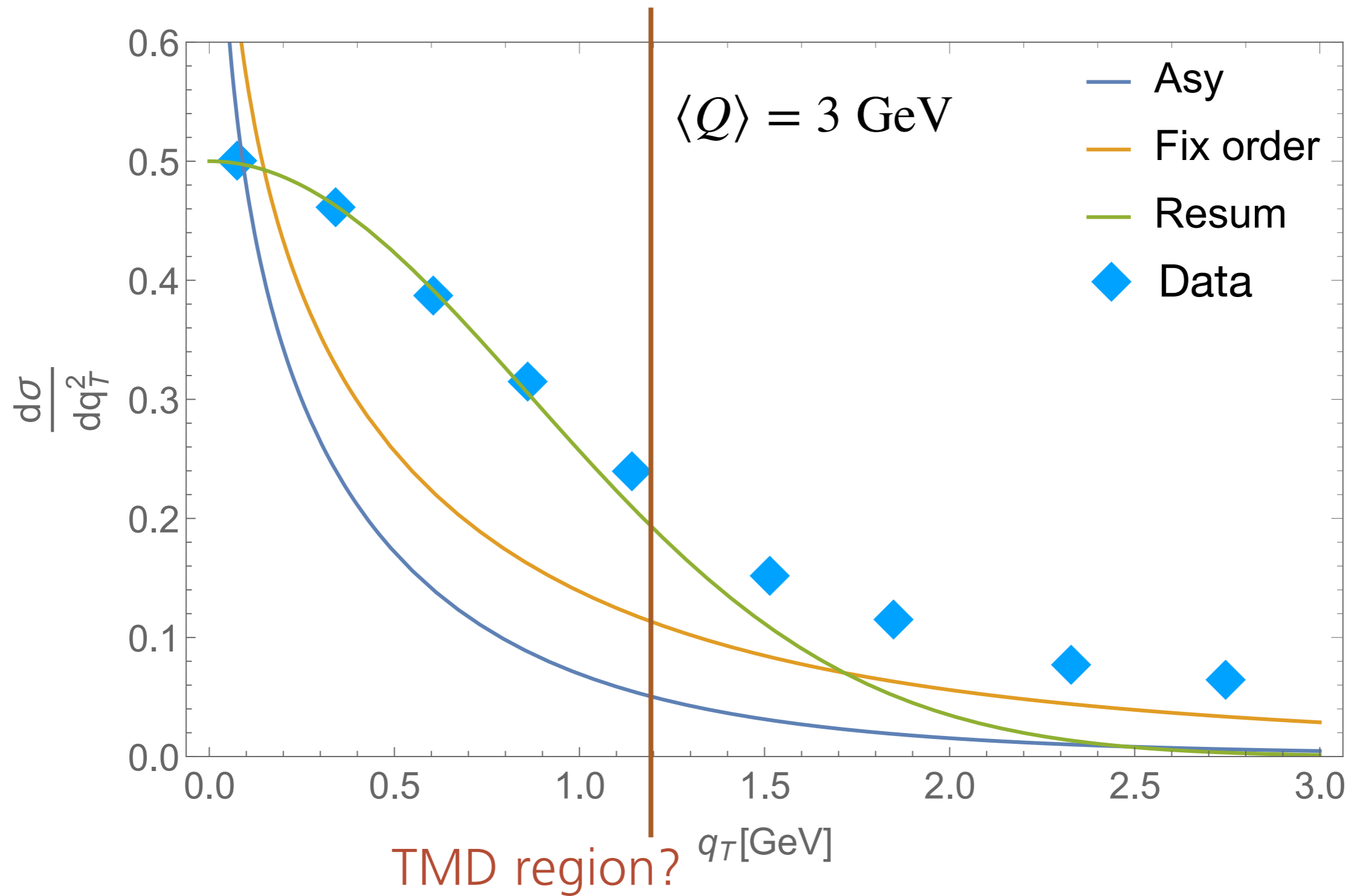
SOME JUSTIFICATION: INITIAL SITUATION



SOLUTION 1: RESTRICT TMD REGION



SOLUTION 2: ENHANCE TMD CONTRIBUTIONS

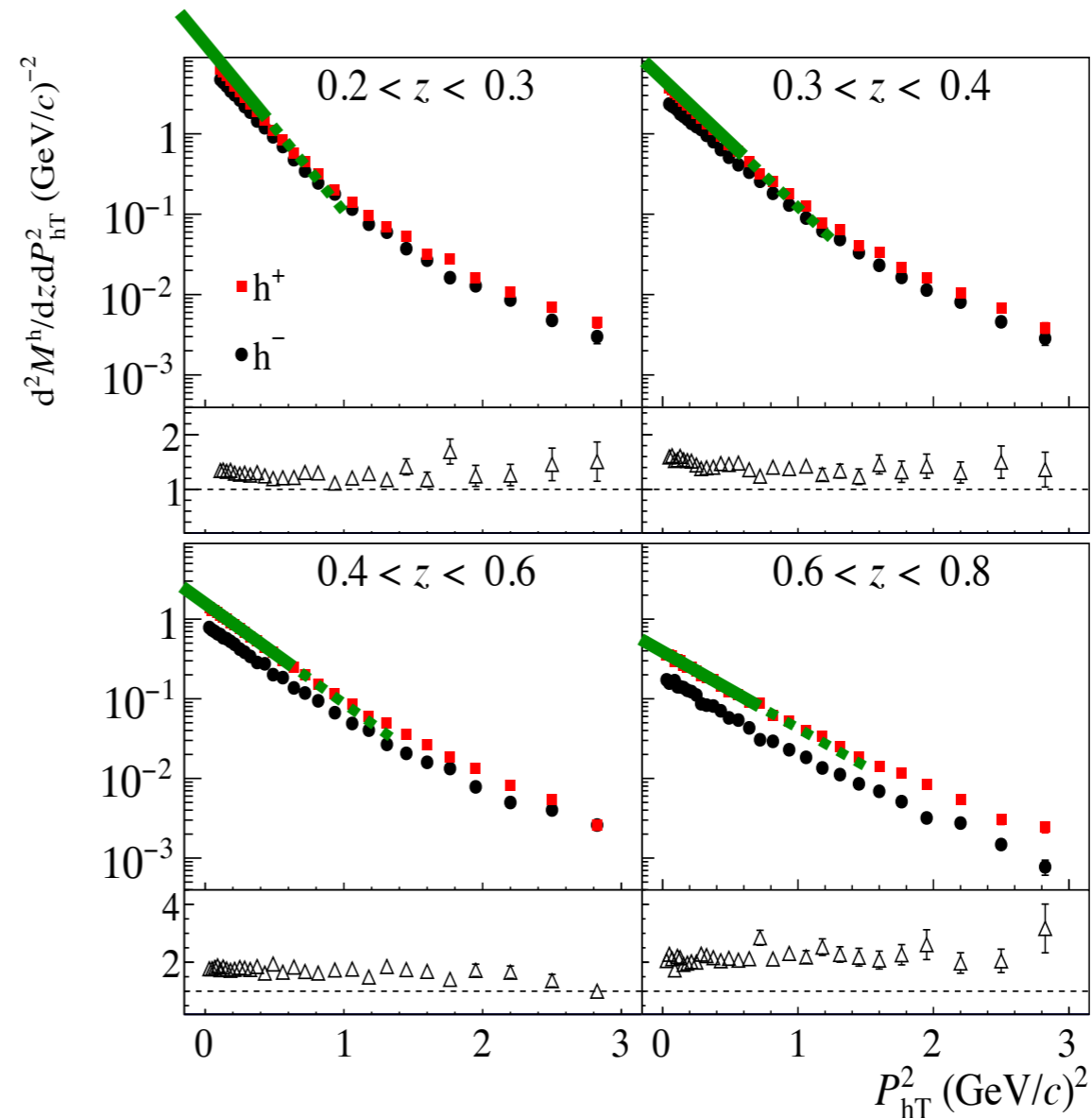


TMD REGIONS AND DATA

COMPASS Collab., arXiv:1709.07374



$$\langle Q \rangle = 3.1 \text{ GeV}$$



It seems that the same “physics” is dominating at least for
 $0 \leq P_{hT} \leq 0.7 \text{ GeV}$,
which means $0 \leq q_T \leq 2.8 \text{ GeV}$ in the lowest- z bin

MAP21 TMD FIT CHOICES (PRELIMINARY)

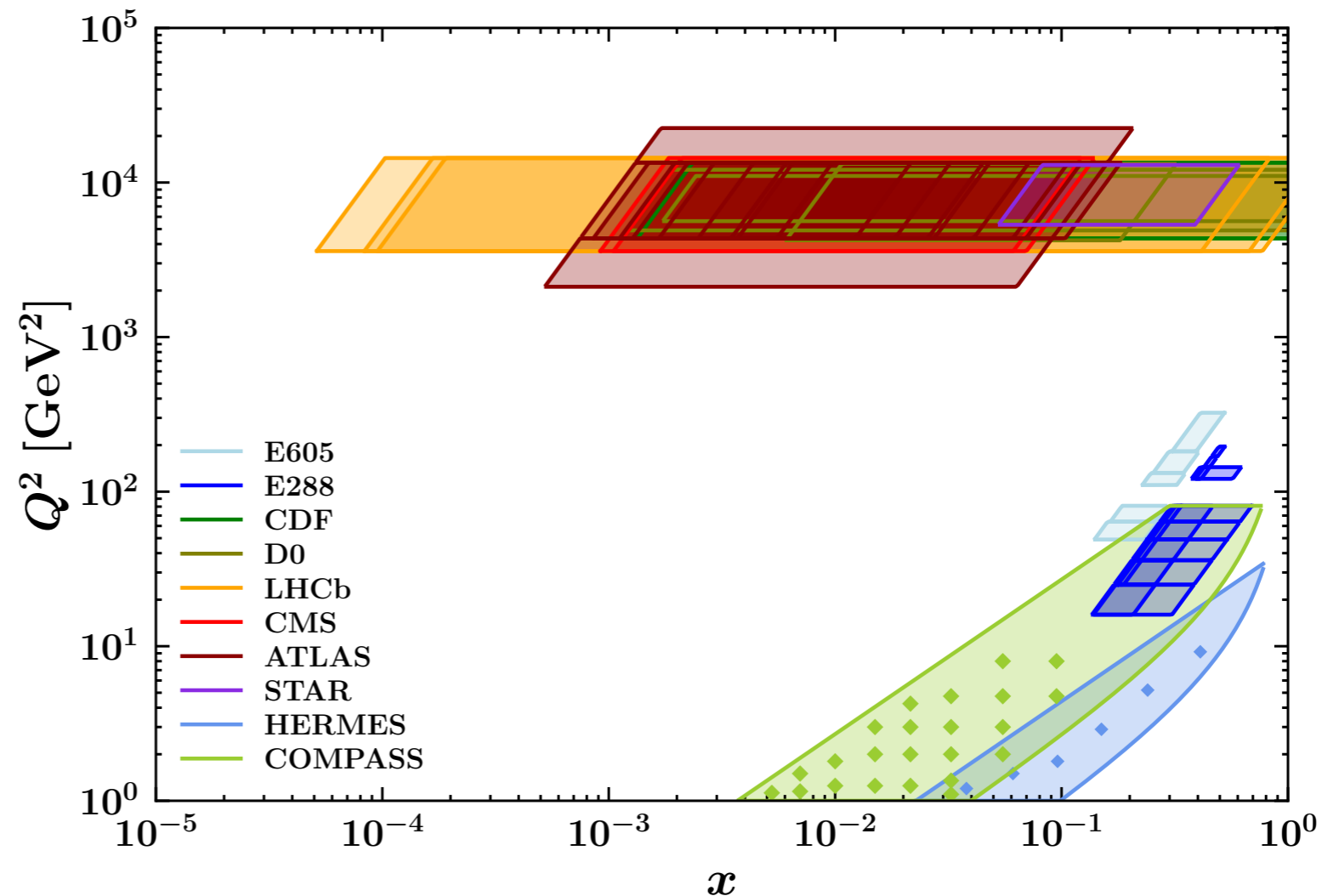
$$Q^2 > 1.4 \text{ GeV}^2$$

$$0.2 < z < 0.6$$

$$q_T < 0.2 Q \quad (\text{DY})$$

$$P_{hT} < \min [\min [0.2 Q, 0.5zQ] + 0.3 \text{ GeV}, zQ] \quad (\text{SIDIS})$$

Number of points: 1520



FUNCTIONAL FORM

$$f_{1\text{NP}}(x, b_T^2) \propto \text{F.T. of } \left(e^{-\frac{k_\perp^2}{g_{1A}}} + \lambda_B k_\perp^2 e^{-\frac{k_\perp^2}{g_{1B}}} + \lambda_C e^{-\frac{k_\perp^2}{g_{1C}}} \right)$$

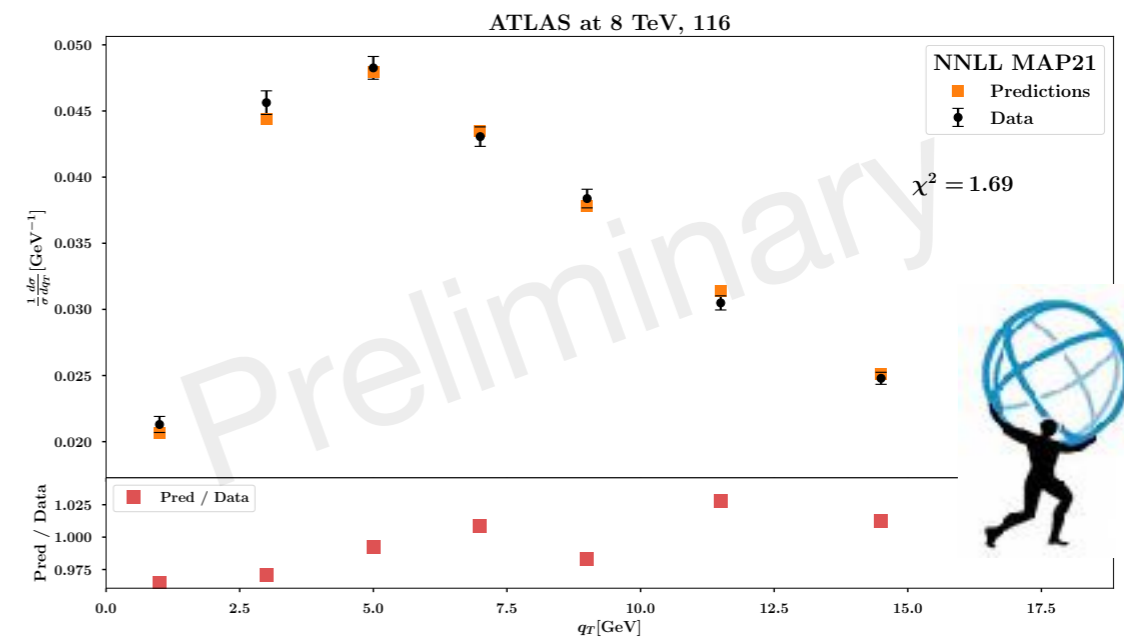
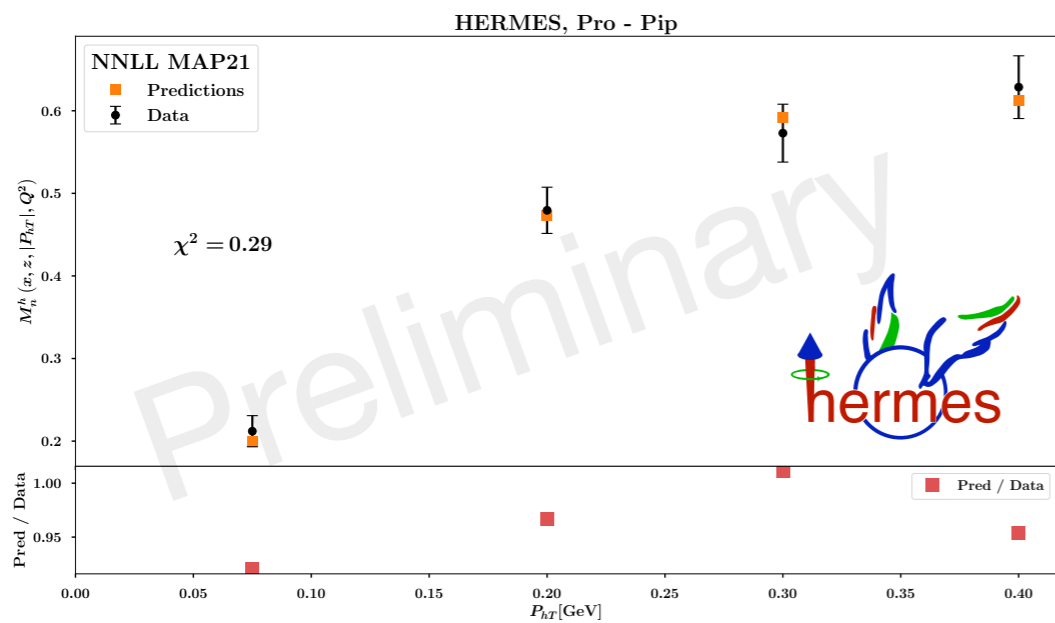
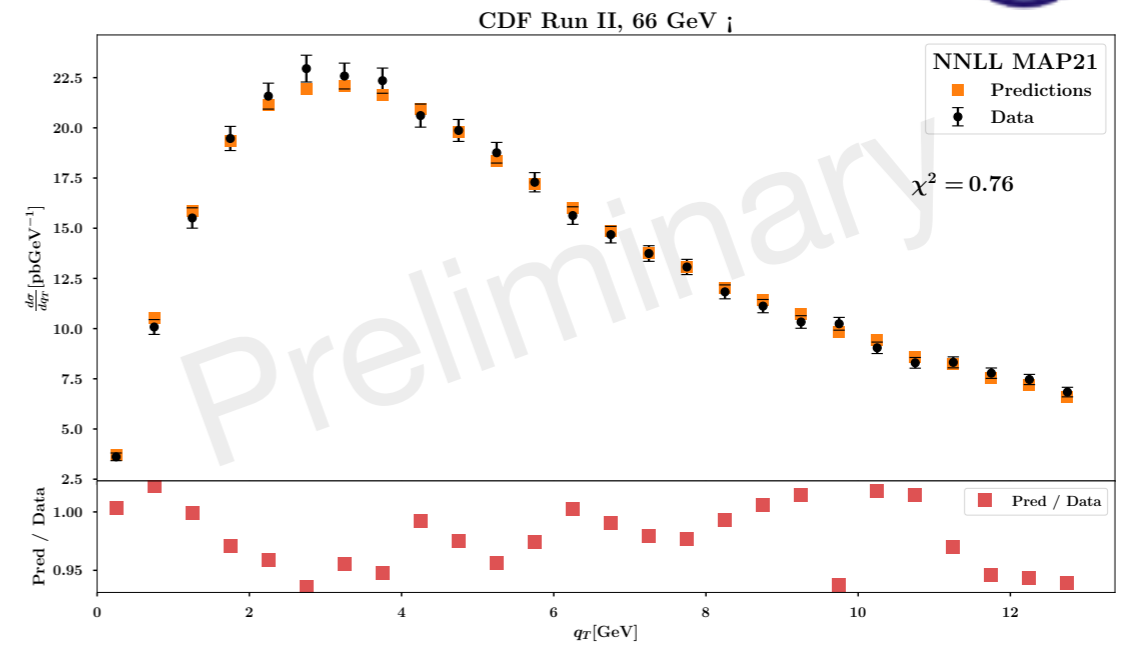
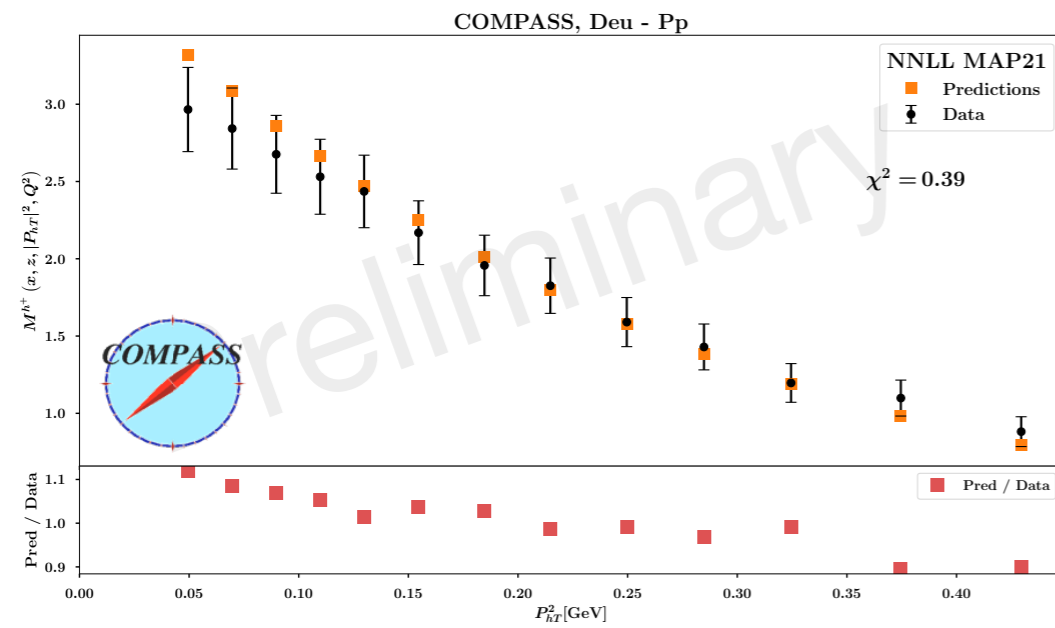
Still working on the final form

$$g_1(x) = N_1 \frac{(1-x)^\alpha x^\sigma}{(1-\hat{x})^\alpha \hat{x}^\sigma}$$

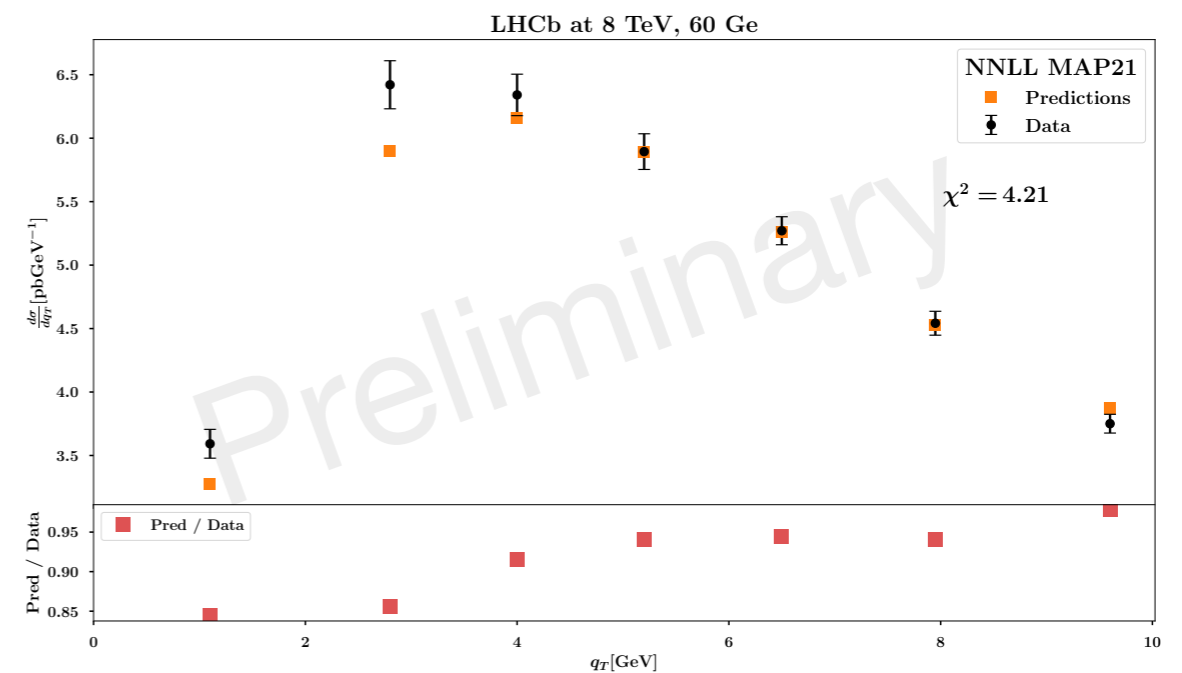
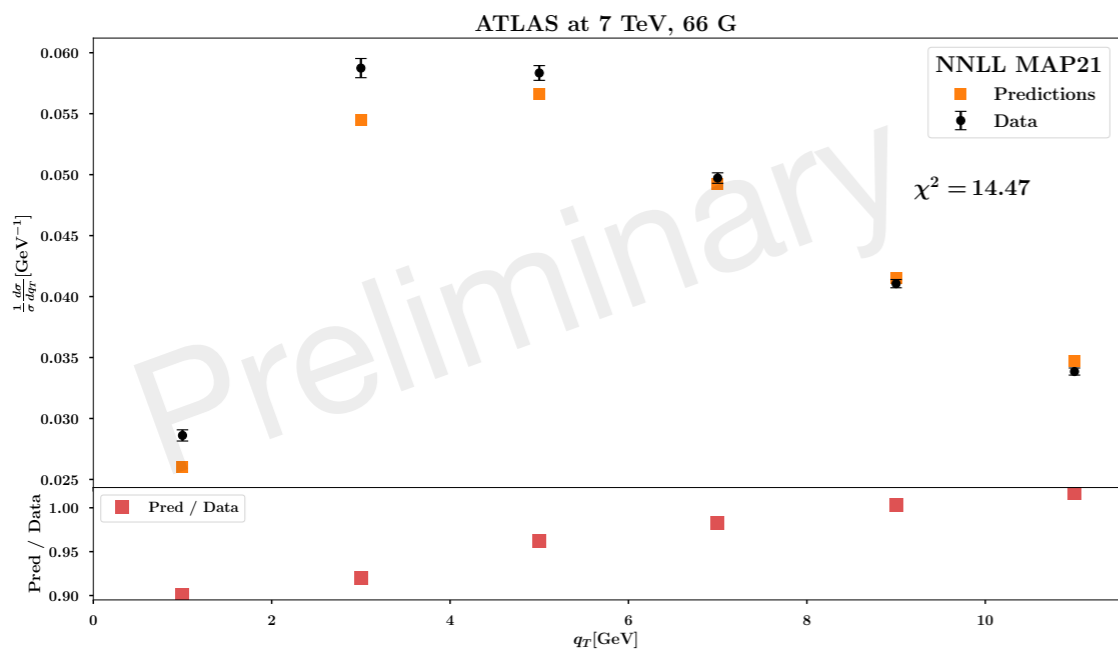
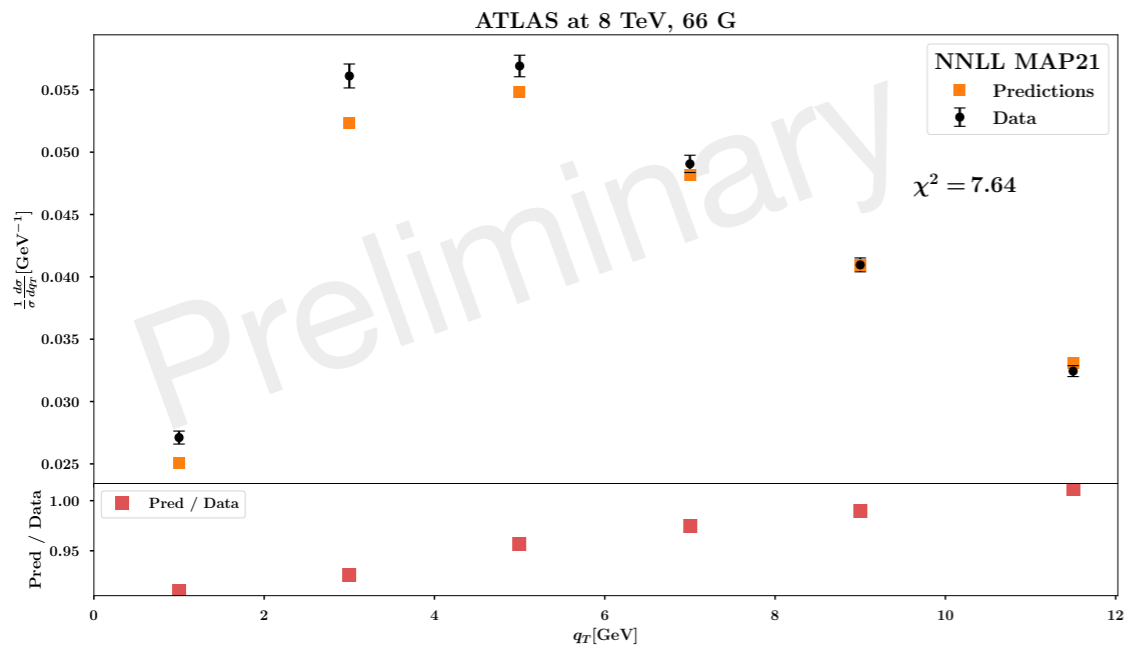
$$g_K(b_T^2) = -\frac{g_{2A}}{2} b_T^2 - \frac{g_{2B}}{2} b_T^4$$

**11 parameters for TMD PDF
+ 2 for NP evolution +14 for FF
= 27 free parameters**

N²LL: EXAMPLE OF GOOD BINS



N²LL: EXAMPLE OF BAD BINS

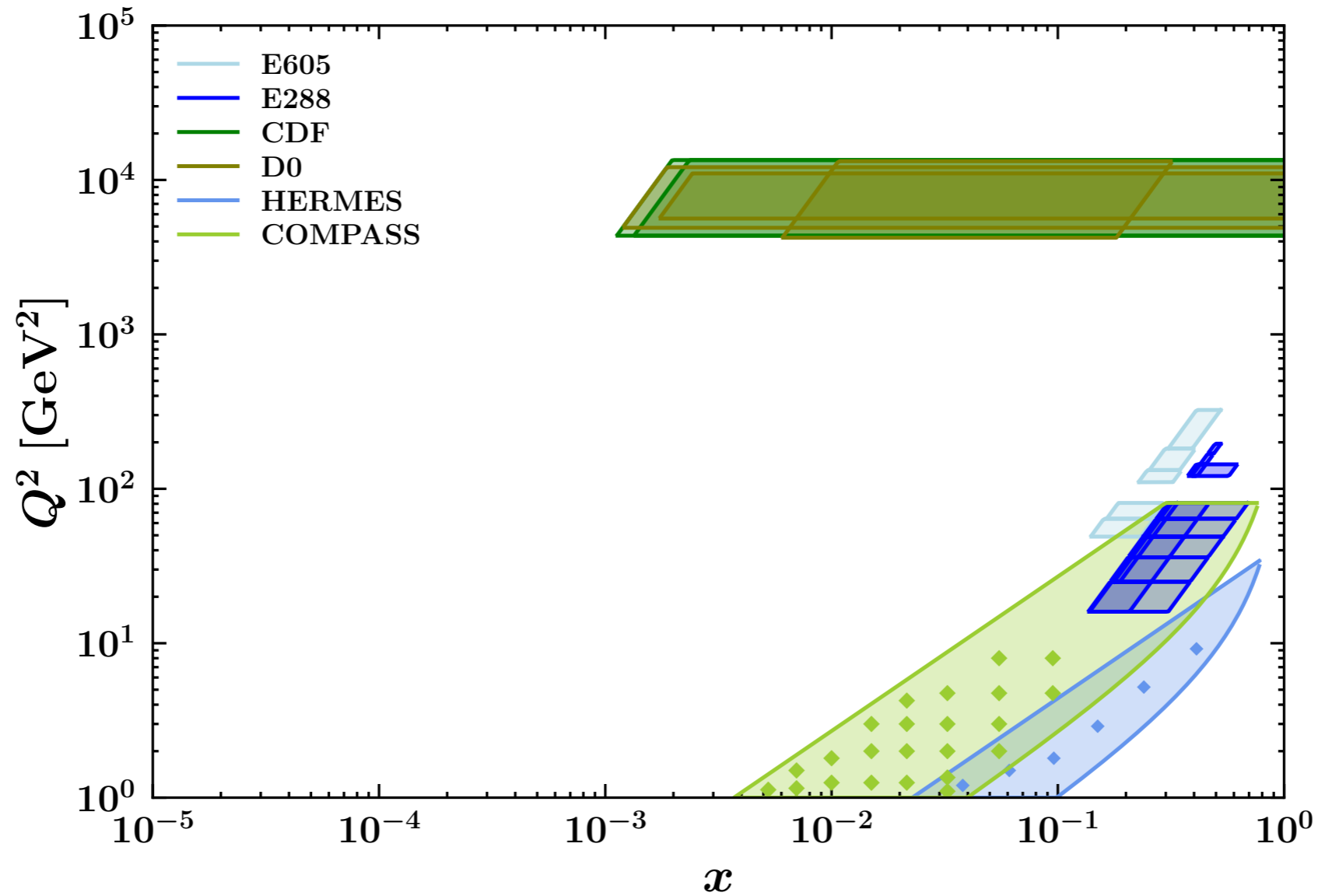


GENERAL CONSIDERATIONS ABOUT UNPOLARIZED TMDs

- SIDIS data can be described very well at NLL, but require normalisation factors at NLL' or higher
- The identification of the region of applicability of the TMD formalism is still an open issue
- Good global χ^2 can be reached at N²LL, but some LHC data remain hard to describe

BACKUP SLIDES

x - Q^2 COVERAGE PV17



Bacchetta, Delcarro, Pisano, Radici, Signori, arXiv:1703.10157

DATA SELECTION IN PAVIA 2017

$$Q^2 > 1.4 \text{ GeV}^2$$

$$0.2 < z < 0.7$$

$$P_{hT}, q_T < \text{Min}[0.2 Q, 0.7 Qz] + 0.5 \text{ GeV}$$

Total number of data points: 8059

Total $\chi^2/\text{dof} = 1.55$

DATA SELECTION IN PAVIA 2017

$$Q^2 > 1.4 \text{ GeV}^2$$

$$0.2 < z < 0.7$$

$$P_{hT}, q_T < \text{Min}[0.2 Q, 0.7 Qz] + 0.5 \text{ GeV}$$

Total number of data points: 8059

Total $\chi^2/\text{dof} = 1.55$

We checked also

$$P_{hT} < \text{Min}[0.2 Q, 0.5 Qz] + 0.3 \text{ GeV}$$

$$P_{hT} < 0.2 Qz$$

Total number of data points: 3380

Total number of data points: 477

Total $\chi^2/\text{dof} = 0.96$

Total $\chi^2/\text{dof} = 1.02$

PV17 – RESULTING TMDS

Bacchetta, Delcarro, Pisano, Radici, Signori, arXiv:1703.10157

expression in b_T space

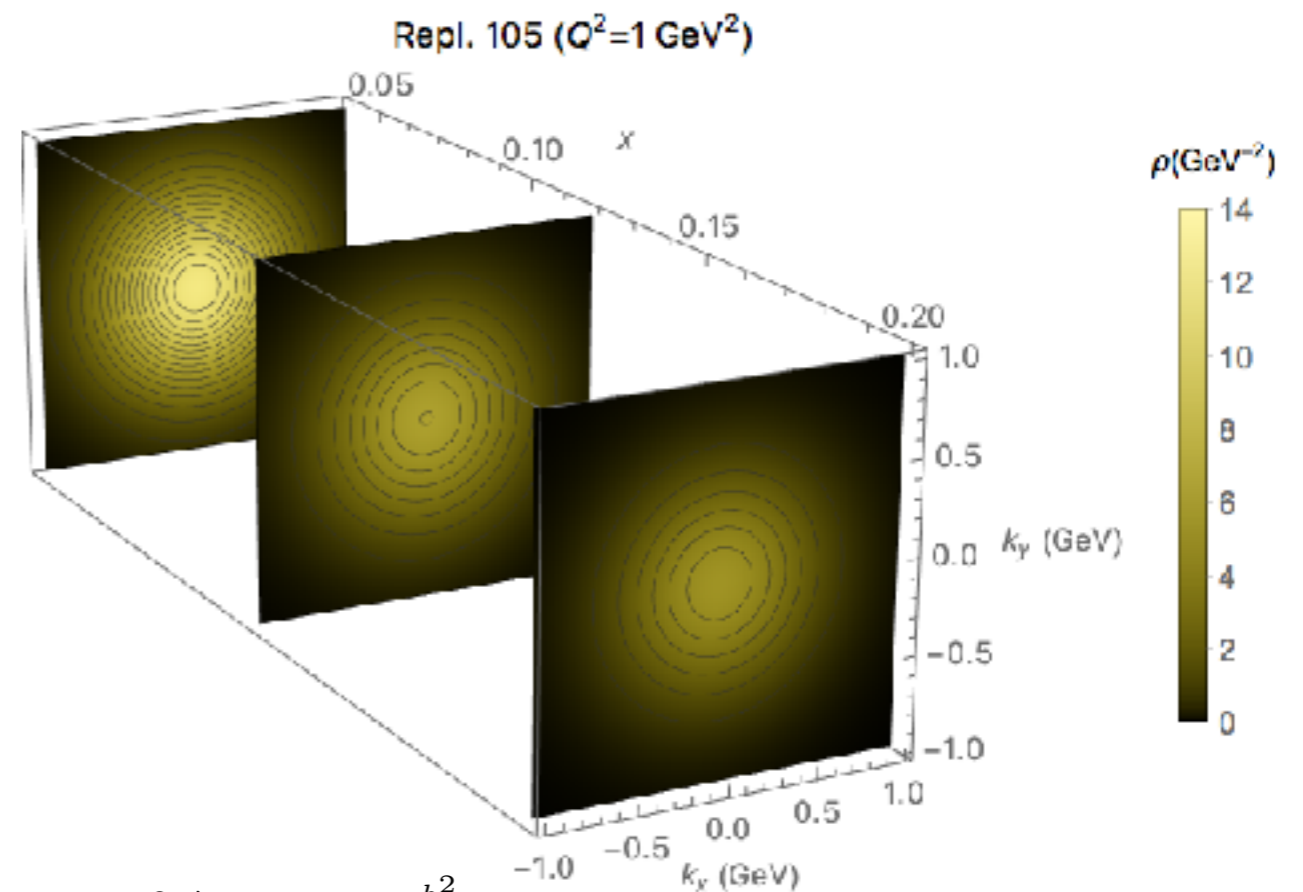
$$\hat{f}_{\text{NP}}(x, b_T) = e^{-g_1(x) \frac{b_T^2}{4}} \left(1 - \frac{\lambda g_1^2(x)}{1 + \lambda g_1(x)} \frac{b_T^2}{4} \right)$$

- Gaussian + weighted Gaussian
- nontrivial x dependence
- no flavor dependence

$$g_K(b_T) = -\frac{g_2}{2} b_T^2 \quad \text{Gaussian}$$

$$\hat{D}_{\text{NP}}(z, b_T) = \frac{g_3(z) e^{-g_3(z) \frac{b_T^2}{4z^2}} + (\lambda_F/z^2) g_4^2(z) \left(1 - g_4(z) \frac{b_T^2}{4z^2} \right) e^{-g_4^2(z) \frac{b_T^2}{4z^2}}}{z^2 \left(g_3(z) + (\lambda_F/z^2) g_4^2(z) \right)}$$

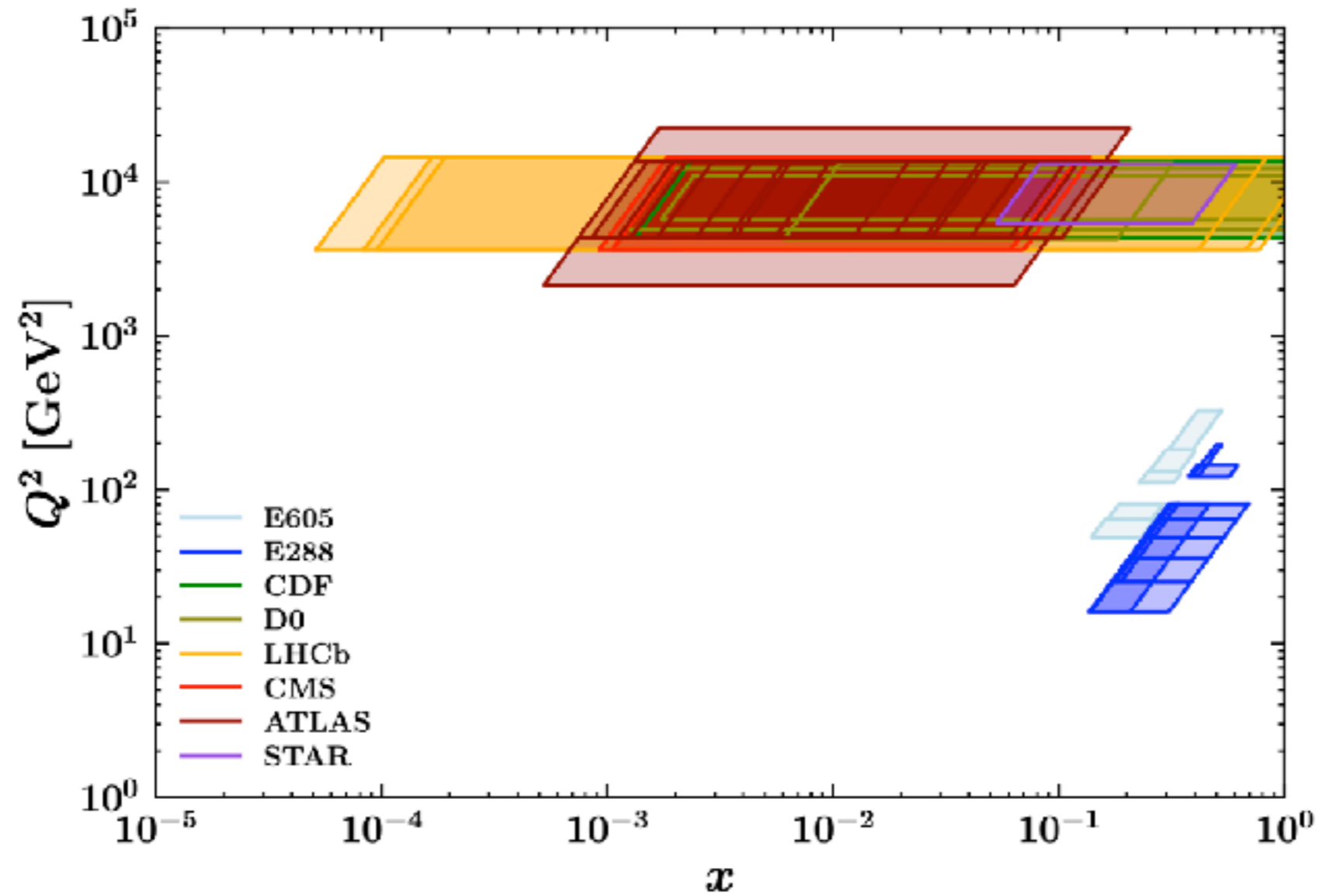
plot in k_\perp space



TMD Frag. Func.

11 free parameters

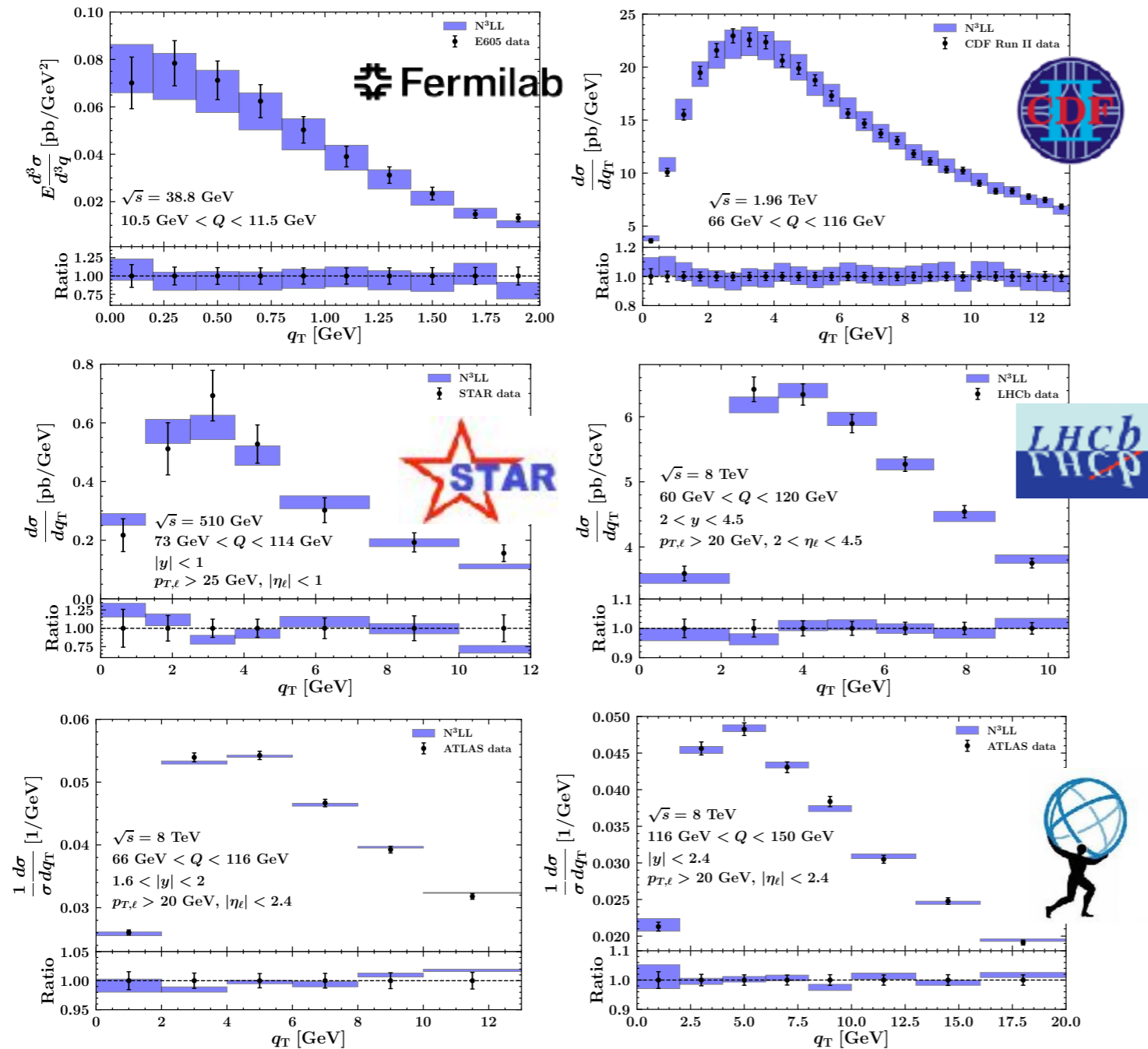
x-Q² COVERAGE PV19



Bacchetta, Bertone, Bissolotti, Bozzi, Delcarro, Piacenza, Radici, arXiv:1912.07550

THE PAVIA19 EXTRACTION

Bacchetta, Bertone, Bissolotti, Bozzi, Delcarro, Piacenza, Radici, arXiv:1912.07550

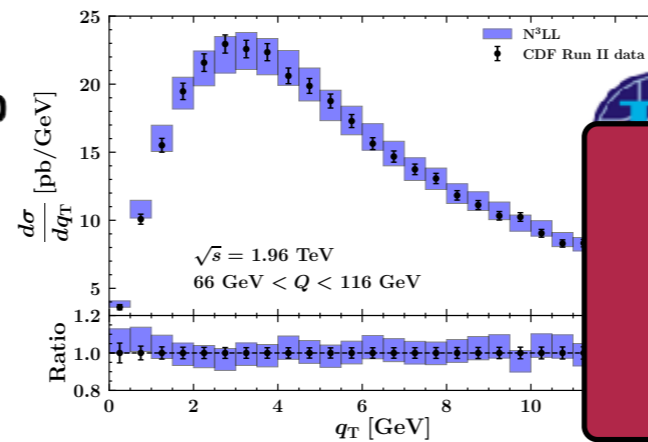
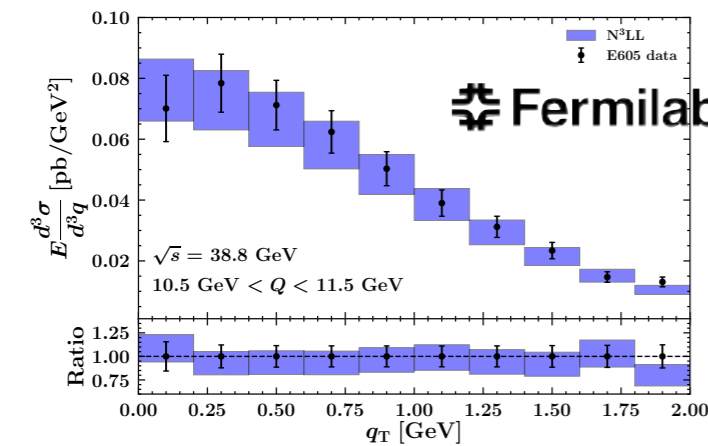


Data selection: $q_T/Q < 0.2$

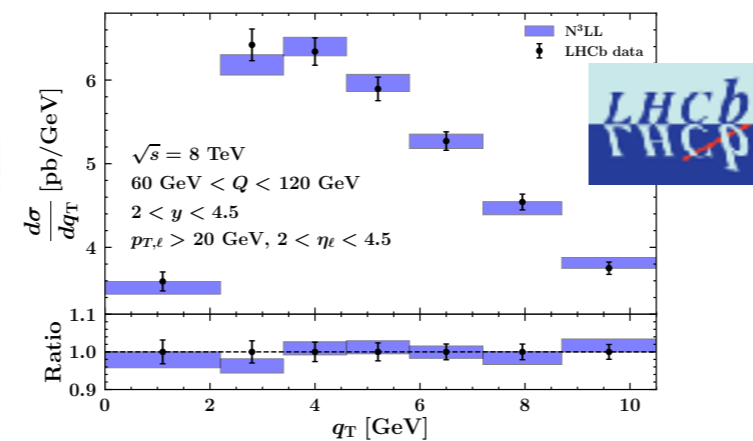
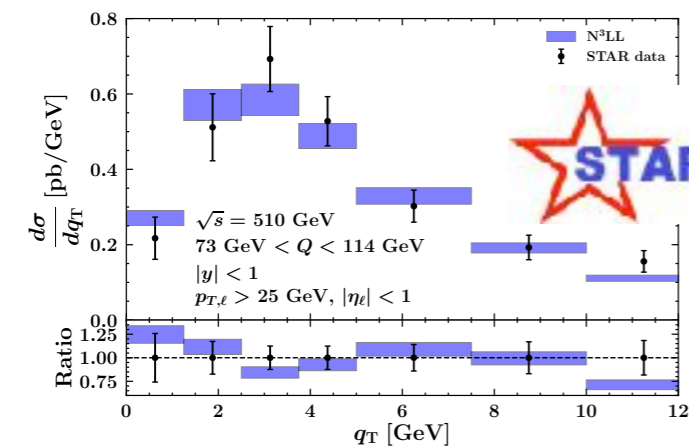
Number of data points: 353

PV19 - DATA COMPARISON

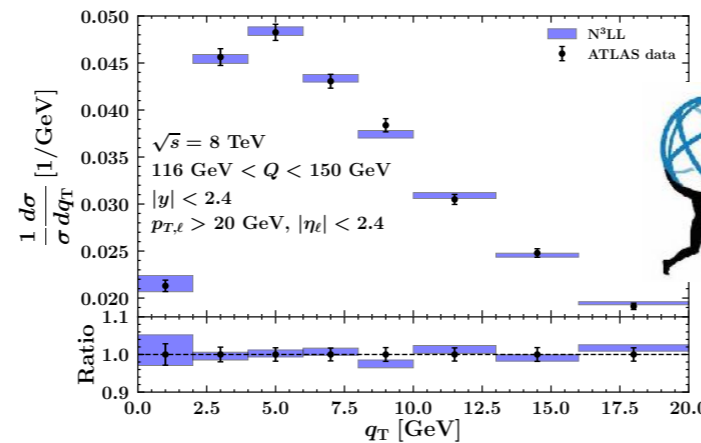
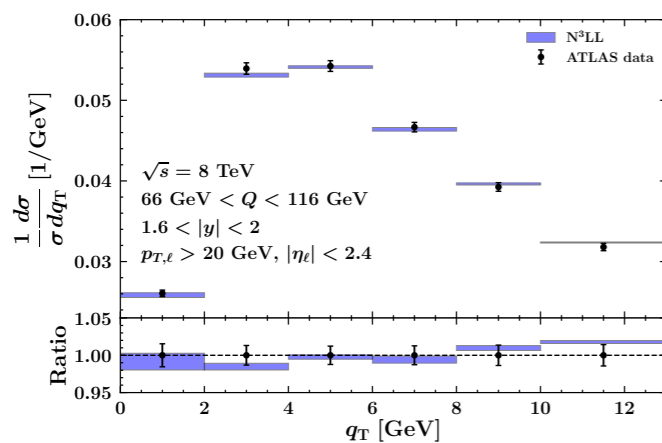
Bacchetta, Bertone, Bissolotti, Bozzi, Delcarro, Piacenza, Radici, arXiv:1912.07550



**Pavia19: first DY fit at N³LL,
 exactly reproduces normalization**



Data selection: $q_T/Q < 0.2$



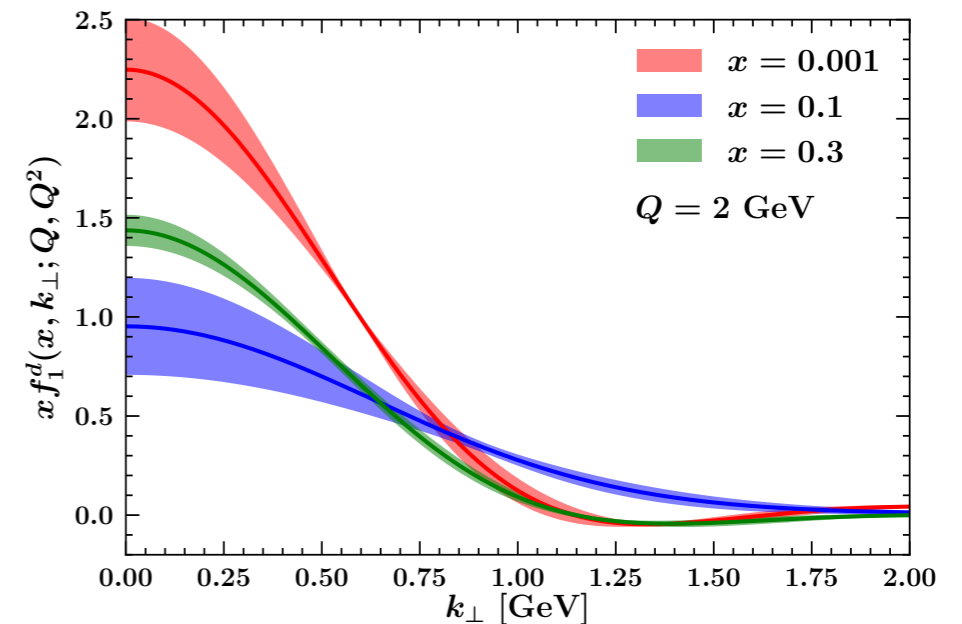
Number of data points: 353
Global $\chi^2/\text{dof} = 1.02$

PV19 – RESULTING TMDS

expression in b_T space

$$f_{\text{NP}}(x, b_T, \zeta) = \left[\frac{1 - \lambda}{1 + g_1(x) \frac{b_T^2}{4}} + \lambda \exp\left(-g_{1B}(x) \frac{b_T^2}{4}\right) \right] \times \exp\left[-(g_2 + g_{2B} b_T^2) \ln\left(\frac{\zeta}{Q_0^2}\right) \frac{b_T^2}{4}\right],$$

plot in k_\perp space



- q-Gaussian + Gaussian
- nontrivial x dependence
- no flavor dependence
- non-Gaussian nonperturbative TMD evolution

9 free parameters

X DEPENDENCE IN TMDS

PV17

$$g_1(x) = N_1 \frac{(1-x)^\alpha x^\sigma}{(1-\hat{x})^\alpha \hat{x}^\sigma}$$

PV19

$$g_1(x) = \frac{N_1}{x\sigma} \exp \left[-\frac{1}{2\sigma^2} \ln^2 \left(\frac{x}{\alpha} \right) \right] ,$$

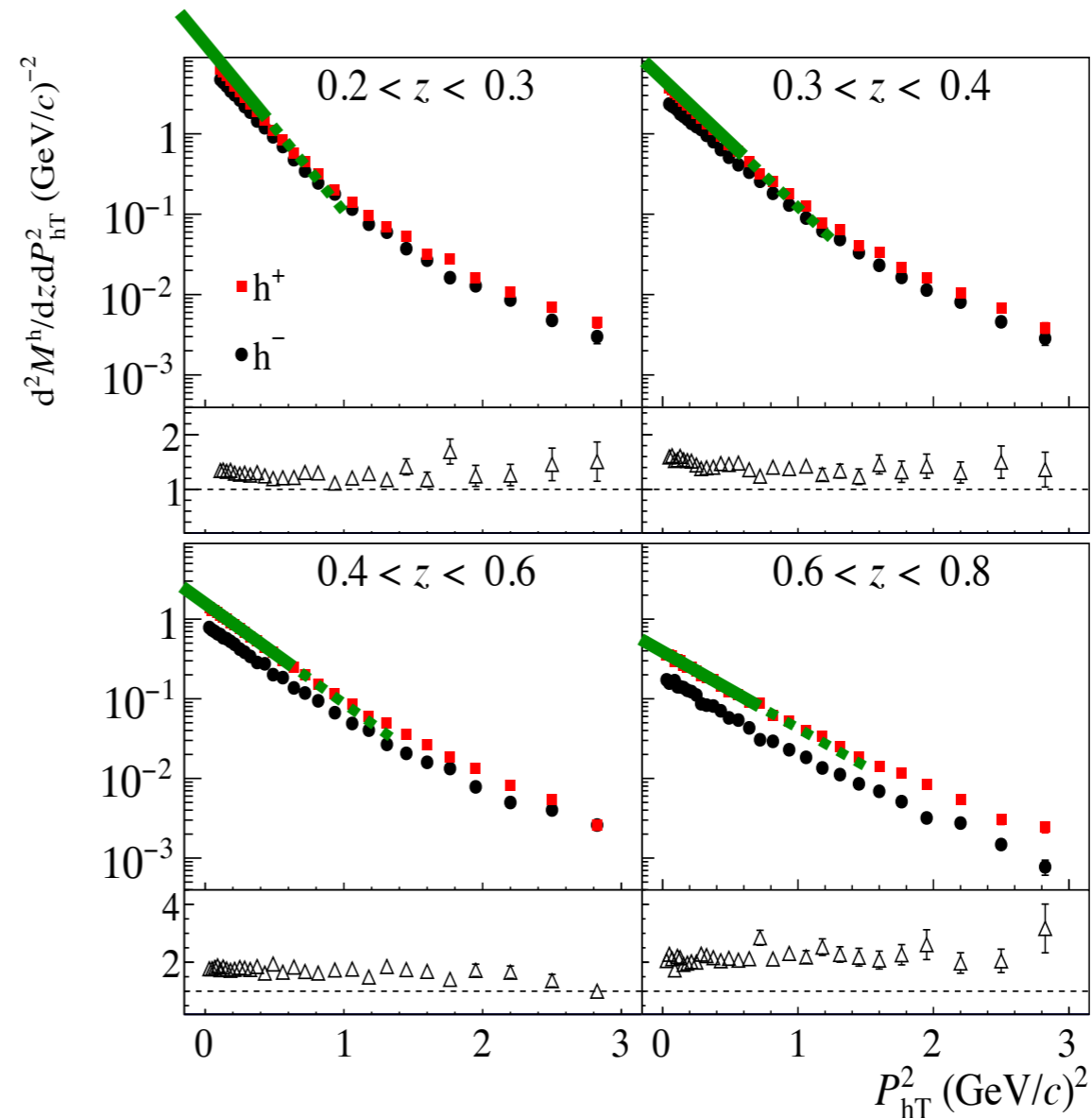
$$g_{1B}(x) = \frac{N_{1B}}{x\sigma_B} \exp \left[-\frac{1}{2\sigma_B^2} \ln^2 \left(\frac{x}{\alpha_B} \right) \right] .$$

TMD REGIONS AND DATA

COMPASS Collab., arXiv:1709.07374



$$\langle Q \rangle = 3.1 \text{ GeV}$$



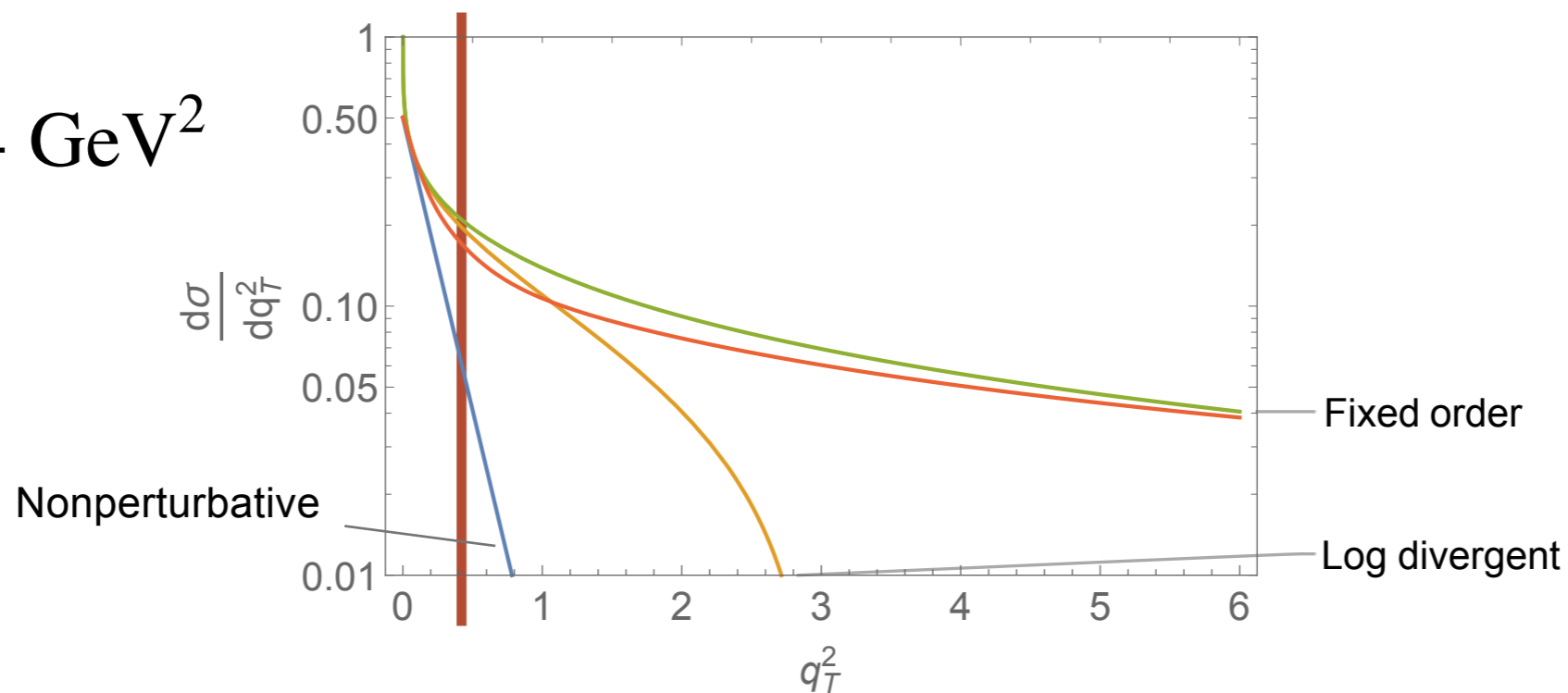
It seems that the same “physics” is dominating at least for
 $0 \leq P_{hT} \leq 0.7 \text{ GeV}$,
which means $0 \leq q_T \leq 2.8 \text{ GeV}$ in the lowest- z bin

TMD REGIONS: PERTURBATIVE VS. NONPERTURBATIVE

Perturbative approach:
TMD region = where the log
divergence of the fixed-order
calculation dominates
(resummation is required)

Nonperturbative approach:
TMD region = where either
the log divergence OR the
nonperturbative
contributions dominate

$$\langle Q^2 \rangle = 4 \text{ GeV}^2$$



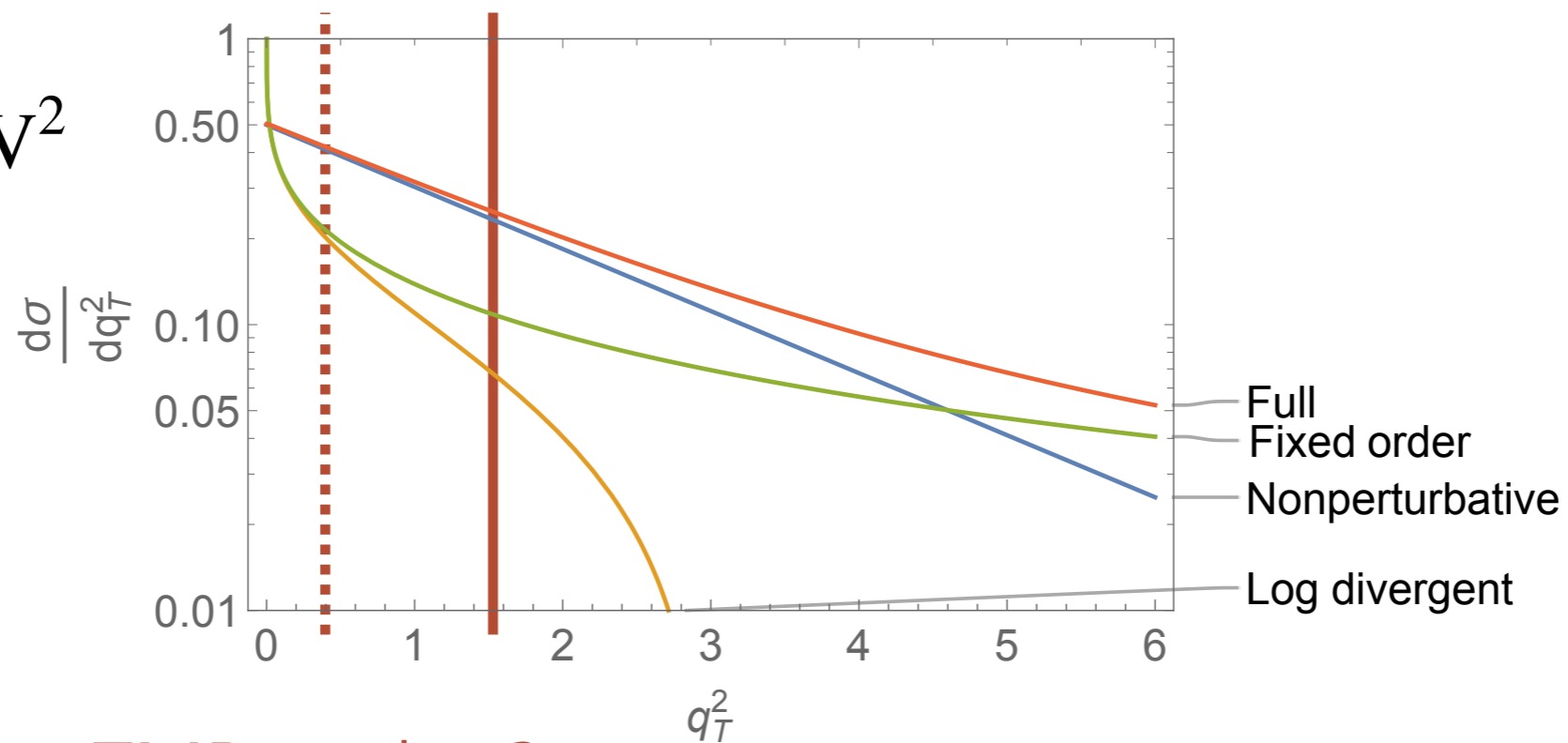
TMD region (ideal situation)

TMD REGIONS: PERTURBATIVE VS. NONPERTURBATIVE

Perturbative approach:
 TMD region = where the log divergence of the fixed-order calculation dominates
 (resummation is required)

Nonperturbative approach:
 TMD region = where either the log divergence OR the nonperturbative contributions dominate

$$\langle Q^2 \rangle = 4 \text{ GeV}^2$$



TMD region?