

# Vacuum distribution and tuning for F-type SUSY breaking

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# Outline

Motivations

The effective field theory method and vacuum distributions

Including mass scales, results and implications

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- ▶ Including different mass scales  $M_S, M_K, M_P$ . As their ratio changes, how does the distribution behave?
- ▶ What is the limiting behaviour as some ratio of scales goes to 0? e.g. gravity decoupling, nearly canonical Kahler.

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## Comparison to the SUSY vacuum distribution

- ▶  $N \sim F_0^6 \Lambda_0$  favors the highest  $F$ . But SUSY vacua live on other branches. How is the total number of SUSY-breaking vacua compared to SUSY ones?



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- ▶ R-symmetries are widely used to get SUSY breaking.
- ▶ Alternatively, one can tune parameters to get metastable SUSY breaking vacua. How much tuning in general do we need?

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- ▶ More than one light field is possible, but need much more tuning (or symmetries).
- ▶ The model is

$$W = \sum_n a_n z^n, \quad K = \sum_{n,m} c_{nm} \bar{z}^n z^m,$$

$$\text{for SUSY: } V = \frac{1}{\bar{\partial} \partial K} \bar{\partial} \bar{W} \partial W,$$

$$\text{for SUGRA: } V = e^K \left( \frac{1}{\bar{\partial} \partial K} \bar{D} \bar{W} D W - 3 \bar{W} W \right).$$

# The general method (2)

## Constraints on parameters

- ▶ Conditions to get SUSY breaking vacua of interest, Each condition gives some constraint on  $a_n$ 's:
  1. Small SUSY breaking scale:  $\Theta(F < F_0)$ ,
  2. Stationary:  $\delta(V')$ ,
  3. Metastable:  $\Theta(V'' > 0)$ ,
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- ▶ The total number of vacua is

$$\begin{aligned} N(F < F_0, 0 < V < \Lambda_0) \\ = \int d\mu(a_n) \Theta(F < F_0) \delta(V') \Theta(V'' > 0) \Theta(0 < V < \Lambda_0) . \end{aligned}$$

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- ▶ In small region  $d\mu(a_n) \sim d^2 a_0 d^2 a_1 \dots$

# Distributions (1)

## General model calculation

- ▶ Take the effective one-field model  $W = \sum_n a_n z^n$ ,  
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- ▶ Make  $c_{11} = 1$ ,  $c_{0n} = 0$ , we have

$$V = a_1^* a_1 - 3a_0^* a_0 ,$$

$$\partial V = 2a_1^* a_2 - 2c_{12} a_1^* a_1 - 2a_0^* a_1 ,$$

$$\begin{aligned} \partial^2 V &= 6a_1^* a_3 - 8c_{12} a_1^* a_2 + (8c_{12}^2 - 6c_{13}) a_1^* a_1 + \\ &\quad - 2a_0^* a_2 - 2c_{12} a_0^* a_1 , \end{aligned}$$

$$\begin{aligned} \bar{\partial} \partial V &= 4a_2^* a_2 - 4c_{12}^* a_1^* a_2 - 4c_{12} a_1 a_2^* + \\ &\quad + (8c_{12}^* c_{12} - 4c_{22}) a_1^* a_1 - 2a_0^* a_0 . \end{aligned}$$

## Distributions (2)

### Counting vacua

- ▶ Constraints on  $a_n$ 's:

$$a_1 \sim a_2 \sim F, \quad a_3 \lesssim F, \quad a_0 \sim F - \frac{\Lambda}{F}.$$

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- ▶ The result is

$$N(F < F_0, 0 < V < \Lambda_0) \sim F_0^2 \cdot F_0^2 \cdot F_0^2 \cdot \Lambda_0 \sim F_0^6 \Lambda_0.$$

# Counting with mass scales (1)

## General model with mass scales

- ▶ Assume 3 mass scales:
  - ▶  $M_S$ : scale of SUSY dynamics,
  - ▶  $M_K$ : scale of non-minimal corrections to Kähler,
  - ▶  $M_P$ : Planck scale.



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$$DW = \partial W + \frac{1}{M_P^2} W\partial K.$$

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$$a_1 \sim F ,$$

$$a_2 \sim FM_S \left( \frac{1}{M_K} + \frac{1}{M_P} \right) ,$$

$$a_3 \lesssim FM_S^2 \left( \frac{1}{M_K^2} + \frac{1}{M_P^2} \right) ,$$

$$a_0 \sim \frac{M_P}{M_S} \left( F - \frac{\Lambda}{FM_S^4} \right) .$$

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- ▶  $\delta(V') = \delta^2\left(\frac{1}{M_S}(z - z_0)\right) \sim F^2 M_S^4\left(\frac{1}{M_K^4} + \frac{1}{M_P^4}\right)\delta^2(a_2 - a_{2(0)})$ .

# Counting with mass scales (3)

## The result

- ▶ For SUSY:

$$N(F < F_0) \sim F_0^6 \frac{M_S^8}{M_K^8} .$$

$N \rightarrow 0$  as  $M_K \rightarrow \infty$ .

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$M_K \rightarrow \infty$ ,  $N$  is still finite.

# Comparison to SUSY vacua

## Conditions for SUSY vacua

- ▶  $F = |a_1| = 0$ ,  $a_0 \sim \sqrt{\Lambda} \frac{M_P}{M_S^3}$ , no constraint on other  $a_n$ 's.

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- ▶ For SUGRA:

$$N(F = 0, 0 < V < \Lambda_0) \sim \frac{\Lambda_0 M_P^2}{M_S^6} \text{ (non-SUSY: } F_0^6 \frac{\Lambda_0 M_S^2 M_P^2}{M_K^8} \text{)} .$$

# SUSY breaking by parameter tuning

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- ▶ Low energy SUSY breaking is rare even for metastable vacua.

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