

# Models of Lepton Flavor From Small Groups

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# Neutrino Mixing Matrix

What we know about the mixing angles . . .

$$\begin{aligned}
 U_{\text{PMNS}} &= \begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -c_{23}s_{12} - s_{13}s_{23}c_{12}e^{i\delta} & c_{23}c_{12} - s_{13}s_{23}s_{12}e^{i\delta} & s_{23}c_{13} \\ s_{23}s_{12} - s_{13}c_{23}c_{12}e^{i\delta} & -s_{23}c_{12} - s_{13}c_{23}s_{12}e^{i\delta} & c_{23}c_{13} \end{pmatrix} \\
 &= \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & e^{-i\delta}s_{13} \\ 0 & 1 & 0 \\ -e^{i\delta}s_{13} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}
 \end{aligned}$$

T. Schwetz, M. A. Tortola, and J. W. F. Valle, "Three-flavour neutrino oscillation update," *New J. Phys.* **10** (2008) 113011, [0808.2016](#).

Angle	$1\sigma$	$2\sigma$	$3\sigma$
$\theta_{12}$	$32.46^\circ - 34.82^\circ$	$31.31^\circ - 36.27^\circ$	$30.00^\circ - 37.47^\circ$
$\theta_{23}$	$41.55^\circ - 49.02^\circ$	$38.65^\circ - 52.54^\circ$	$36.87^\circ - 54.94^\circ$
$\theta_{13}$	$0.00^\circ - 9.28^\circ$	$0.00^\circ - 11.54^\circ$	$0.00^\circ - 13.69^\circ$

# Harrison-Perkins-Scott Matrix

Presently our best guess . . .

P. F. Harrison, D. H. Perkins, and W. G. Scott, "Tri-bimaximal mixing and the neutrino oscillation data," *Phys. Lett.* **B530** (2002) 167, [hep-ph/0202074](#).

$$U_{\text{HPS}} = \begin{pmatrix} \sqrt{2/3} & 1/\sqrt{3} & 0 \\ -1/\sqrt{6} & 1/\sqrt{3} & -1/\sqrt{2} \\ -1/\sqrt{6} & 1/\sqrt{3} & 1/\sqrt{2} \end{pmatrix}$$

Indicative of an underlying symmetry . . .

Some groups that have been considered in the literature:

Review: G. Altarelli and F. Feruglio, "Discrete Flavor Symmetries and Models of Neutrino Mixing," [1002.0211](#).

$S_3$ ,  $D_4$ ,  $D_7$ ,  $A_4$ ,  $A_5$ ,  $\tilde{T}$ ,  $S_4$ ,  $(C_3 \times C_3) \rtimes_{\varphi} C_3$ ,  $C_7 \rtimes_{\varphi} C_3$ ,  $\text{PSL}_2(7)$

$\rightsquigarrow$  As a paradigm, we will consider a model with  $A_4 \times C_3$  symmetry and then generalize it to other symmetry groups

# Altarelli-Feruglio Model Revisited

G. Altarelli and F. Feruglio, "Tri-Bimaximal Neutrino Mixing,  $A_4$  and the Modular Symmetry," *Nucl. Phys.* **B741** (2006) 215–235, [hep-ph/0512103](https://arxiv.org/abs/hep-ph/0512103).

## 1 Symmetries of the model

$$SU(2)_L \times U(1)_Y \times U(1)_R \times A_4 \times C_3$$

## 2 Particle content and charges

Field	$SU(2)_L \times U(1)_Y$	$U(1)_R$	$A_4$	$C_3$	$A_4 \times C_3$
$L$	(2, -1)	1	<b>3</b>	$\omega$	$3'$
$e$	(1, 2)	1	<b>1</b>	$\omega^2$	$1'$
$\mu$	(1, 2)	1	<b>1''</b>	$\omega^2$	$1^{(8)}$
$\tau$	(1, 2)	1	<b>1'</b>	$\omega^2$	$1^{(5)}$
$h_u$	(2, 1)	0	<b>1</b>	1	1
$h_d$	(2, -1)	0	<b>1</b>	1	1
$\varphi_T$	(1, 0)	0	<b>3</b>	1	3
$\varphi_S$	(1, 0)	0	<b>3</b>	$\omega$	$3'$
$\xi$	(1, 0)	0	<b>1</b>	$\omega$	$1''$

## 3 Breaking the family symmetry

$$\varphi_T = (v_T, 0, 0), \quad \varphi_S = (v_S, v_S, v_S), \quad \xi = v_\xi,$$

# Group Information from GAP

The GAP Group, "GAP – Groups, Algorithms, and Programming, Version 4.4.12.", <http://www.gap-system.org>

“GAP is a system for computational discrete algebra, with particular emphasis on Computational Group Theory.”

```
group := SmallGroup(36,11);
Display(StructureDescription(group));
Display(GeneratorsOfGroup(group));
for e1 in GeneratorsOfGroup(group) do
  Display(Factorization(group, e1));
od;
for e1 in Elements(group) do
  Display(Factorization(group, e1));
od;
chartab := Irr(group);
Display(chartab);
SizesConjugacyClasses(CharacterTable(group));
LoadPackage("repsn");
for i in [1..Size(chartab)] do
  Display(IrreducibleAffordingRepresentation(chartab[i]));
od;
```

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➤ Specify the group that we will work with

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➤ The “human readable” name of the group

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od;
```

➤ The generators of the group



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od;
```

➤ Express all group elements in terms of the generators

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od;
```

➤ The character table

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```

➤ Dimensions of the conjugacy classes

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od;
```

➤ The matrices for the representations

The Character Table of  $A_4 \times C_3$ 

	$K_1$	$K_2$	$K_3$	$K_4$	$K_5$	$K_6$	$K_7$	$K_8$	$K_9$	$K_{10}$	$K_{11}$	$K_{12}$
<b>1</b>	1	1	1	1	1	1	1	1	1	1	1	1
<b>1'</b>	1	1	$\omega^2$	1	1	$\omega^2$	$\omega$	$\omega^2$	$\omega^2$	$\omega$	$\omega$	$\omega$
<b>1''</b>	1	1	$\omega$	1	1	$\omega$	$\omega^2$	$\omega$	$\omega$	$\omega^2$	$\omega^2$	$\omega^2$
<b>1'''</b>	1	$\omega^2$	1	1	$\omega$	$\omega^2$	1	1	$\omega$	$\omega^2$	1	$\omega$
<b>1<sup>(4)</sup></b>	1	$\omega$	1	1	$\omega^2$	$\omega$	1	1	$\omega^2$	$\omega$	1	$\omega^2$
<b>1<sup>(5)</sup></b>	1	$\omega^2$	$\omega^2$	1	$\omega$	$\omega$	$\omega$	$\omega^2$	1	1	$\omega$	$\omega^2$
<b>1<sup>(6)</sup></b>	1	$\omega$	$\omega$	1	$\omega^2$	$\omega^2$	$\omega^2$	$\omega$	1	1	$\omega^2$	$\omega$
<b>1<sup>(7)</sup></b>	1	$\omega^2$	$\omega$	1	$\omega$	1	$\omega^2$	$\omega$	$\omega^2$	$\omega$	$\omega^2$	1
<b>1<sup>(8)</sup></b>	1	$\omega$	$\omega^2$	1	$\omega^2$	1	$\omega$	$\omega^2$	$\omega$	$\omega^2$	$\omega$	1
<b>3</b>	3	0	3	-1	0	0	3	-1	0	0	-1	0
<b>3'</b>	3	0	$3\omega$	-1	0	0	$3\omega^2$	$\omega$	0	0	$1 + \omega$	0
<b>3''</b>	3	0	$3\omega^2$	-1	0	0	$3\omega$	$1 + \omega$	0	0	$\omega$	0

$\omega = e^{2\pi i/3}$  is the primitive third root of unity

## Match Irreps to Altarelli-Feruglio's Notation

Compare matrix representations of irreps for  $A_3$ ,  $C_3$  and  $A_3 \times C_3$ :

Field	$SU(2)_L \times U(1)_Y$	$U(1)_R$	$A_4$	$C_3$	$A_4 \times C_3$
$L$	$(\mathbf{2}, -1)$	1	$\mathbf{3}$	$\omega$	$\mathbf{3}'$
$e$	$(\mathbf{1}, 2)$	1	$\mathbf{1}$	$\omega^2$	$\mathbf{1}'$
$\mu$	$(\mathbf{1}, 2)$	1	$\mathbf{1}''$	$\omega^2$	$\mathbf{1}^{(8)}$
$\tau$	$(\mathbf{1}, 2)$	1	$\mathbf{1}'$	$\omega^2$	$\mathbf{1}^{(5)}$
$h_u$	$(\mathbf{2}, 1)$	0	$\mathbf{1}$	$\mathbf{1}$	$\mathbf{1}$
$h_d$	$(\mathbf{2}, -1)$	0	$\mathbf{1}$	$\mathbf{1}$	$\mathbf{1}$
$\varphi_T$	$(\mathbf{1}, 0)$	0	$\mathbf{3}$	$\mathbf{1}$	$\mathbf{3}$
$\varphi_S$	$(\mathbf{1}, 0)$	0	$\mathbf{3}$	$\omega$	$\mathbf{3}'$
$\xi$	$(\mathbf{1}, 0)$	0	$\mathbf{1}$	$\omega$	$\mathbf{1}''$

## Decomposition of Tensor Products

From the character table and the dimensions of the conjugacy classes:

$$\begin{array}{lllll}
 1 \otimes 1 = 1 & 1 \otimes 1' = 1' & 1 \otimes 1'' = 1'' & 1 \otimes 1''' = 1''' & 1 \otimes 1^{(4)} = 1^{(4)} \\
 1 \otimes 1^{(5)} = 1^{(5)} & 1 \otimes 1^{(6)} = 1^{(6)} & 1 \otimes 1^{(7)} = 1^{(7)} & 1 \otimes 1^{(8)} = 1^{(8)} & 1 \otimes 3 = 3 \\
 1 \otimes 3' = 3' & 1 \otimes 3'' = 3'' & 1' \otimes 1' = 1'' & 1' \otimes 1'' = 1 & 1' \otimes 1''' = 1^{(5)} \\
 1' \otimes 1^{(4)} = 1^{(8)} & 1' \otimes 1^{(5)} = 1^{(7)} & 1' \otimes 1^{(6)} = 1^{(4)} & 1' \otimes 1^{(7)} = 1''' & 1' \otimes 1^{(8)} = 1^{(6)} \\
 1' \otimes 3 = 3'' & 1' \otimes 3' = 3 & 1' \otimes 3'' = 3' & 1'' \otimes 1'' = 1' & 1'' \otimes 1''' = 1^{(7)} \\
 1'' \otimes 1^{(4)} = 1^{(6)} & 1'' \otimes 1^{(5)} = 1''' & 1'' \otimes 1^{(6)} = 1^{(8)} & 1'' \otimes 1^{(7)} = 1^{(5)} & 1'' \otimes 1^{(8)} = 1^{(4)} \\
 1'' \otimes 3 = 3' & 1'' \otimes 3' = 3'' & 1'' \otimes 3'' = 3 & 1''' \otimes 1''' = 1^{(4)} & 1''' \otimes 1^{(4)} = 1 \\
 1''' \otimes 1^{(5)} = 1^{(8)} & 1''' \otimes 1^{(6)} = 1'' & 1''' \otimes 1^{(7)} = 1^{(6)} & 1''' \otimes 1^{(8)} = 1' & 1''' \otimes 3 = 3 \\
 1''' \otimes 3' = 3' & 1''' \otimes 3'' = 3'' & 1^{(4)} \otimes 1^{(4)} = 1''' & 1^{(4)} \otimes 1^{(5)} = 1' & 1^{(4)} \otimes 1^{(6)} = 1^{(7)} \\
 1^{(4)} \otimes 1^{(7)} = 1'' & 1^{(4)} \otimes 1^{(8)} = 1^{(5)} & 1^{(4)} \otimes 3 = 3 & 1^{(4)} \otimes 3' = 3' & 1^{(4)} \otimes 3'' = 3'' \\
 1^{(5)} \otimes 1^{(5)} = 1^{(6)} & 1^{(5)} \otimes 1^{(6)} = 1 & 1^{(5)} \otimes 1^{(7)} = 1^{(4)} & 1^{(5)} \otimes 1^{(8)} = 1'' & 1^{(5)} \otimes 3 = 3'' \\
 1^{(5)} \otimes 3' = 3 & 1^{(5)} \otimes 3'' = 3' & 1^{(6)} \otimes 1^{(6)} = 1^{(5)} & 1^{(6)} \otimes 1^{(7)} = 1' & 1^{(6)} \otimes 1^{(8)} = 1''' \\
 1^{(6)} \otimes 3 = 3' & 1^{(6)} \otimes 3' = 3'' & 1^{(6)} \otimes 3'' = 3 & 1^{(7)} \otimes 1^{(7)} = 1^{(8)} & 1^{(7)} \otimes 1^{(8)} = 1 \\
 1^{(7)} \otimes 3 = 3' & 1^{(7)} \otimes 3' = 3'' & 1^{(7)} \otimes 3'' = 3 & 1^{(8)} \otimes 1^{(8)} = 1^{(7)} & 1^{(8)} \otimes 3 = 3'' \\
 1^{(8)} \otimes 3' = 3 & 1^{(8)} \otimes 3'' = 3' & & & 
 \end{array}$$

$$\begin{aligned}
 3 \otimes 3 &= 1 + 1''' + 1^{(4)} + 2 \otimes 3 \\
 3 \otimes 3'' &= 1' + 1^{(5)} + 1^{(8)} + 2 \otimes 3'' \\
 3' \otimes 3'' &= 1 + 1''' + 1^{(4)} + 2 \otimes 3
 \end{aligned}$$

$$\begin{aligned}
 3 \otimes 3' &= 1'' + 1^{(6)} + 1^{(7)} + 2 \otimes 3' \\
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 1 \otimes 1^{(5)} = 1^{(5)} & 1 \otimes 1^{(6)} = 1^{(6)} & 1 \otimes 1^{(7)} = 1^{(7)} & 1 \otimes 1^{(8)} = 1^{(8)} & 1 \otimes 3 = 3 \\
 1 \otimes 3' = 3' & 1 \otimes 3'' = 3'' & 1' \otimes 1' = 1'' & 1' \otimes 1'' = 1 & 1' \otimes 1''' = 1^{(5)} \\
 1' \otimes 1^{(4)} = 1^{(8)} & 1' \otimes 1^{(5)} = 1^{(7)} & 1' \otimes 1^{(6)} = 1^{(4)} & 1' \otimes 1^{(7)} = 1''' & 1' \otimes 1^{(8)} = 1^{(6)} \\
 1' \otimes 3 = 3'' & \mathbf{1' \otimes 3' = 3} & 1' \otimes 3'' = 3' & 1'' \otimes 1'' = 1' & 1'' \otimes 1''' = 1^{(7)} \\
 1'' \otimes 1^{(4)} = 1^{(6)} & 1'' \otimes 1^{(5)} = 1''' & 1'' \otimes 1^{(6)} = 1^{(8)} & 1'' \otimes 1^{(7)} = 1^{(5)} & 1'' \otimes 1^{(8)} = 1^{(4)} \\
 1'' \otimes 3 = 3' & 1'' \otimes 3' = 3'' & 1'' \otimes 3'' = 3 & 1''' \otimes 1''' = 1^{(4)} & 1''' \otimes 1^{(4)} = 1 \\
 1''' \otimes 1^{(5)} = 1^{(8)} & 1''' \otimes 1^{(6)} = 1'' & 1''' \otimes 1^{(7)} = 1^{(6)} & 1''' \otimes 1^{(8)} = 1' & 1''' \otimes 3 = 3 \\
 1''' \otimes 3' = 3' & 1''' \otimes 3'' = 3'' & 1^{(4)} \otimes 1^{(4)} = 1''' & 1^{(4)} \otimes 1^{(5)} = 1' & 1^{(4)} \otimes 1^{(6)} = 1^{(7)} \\
 1^{(4)} \otimes 1^{(7)} = 1'' & 1^{(4)} \otimes 1^{(8)} = 1^{(5)} & 1^{(4)} \otimes 3 = 3 & 1^{(4)} \otimes 3' = 3' & 1^{(4)} \otimes 3'' = 3'' \\
 1^{(5)} \otimes 1^{(5)} = 1^{(6)} & 1^{(5)} \otimes 1^{(6)} = 1 & 1^{(5)} \otimes 1^{(7)} = 1^{(4)} & 1^{(5)} \otimes 1^{(8)} = 1'' & 1^{(5)} \otimes 3 = 3'' \\
 \mathbf{1^{(5)} \otimes 3' = 3} & 1^{(5)} \otimes 3'' = 3' & 1^{(6)} \otimes 1^{(6)} = 1^{(5)} & 1^{(6)} \otimes 1^{(7)} = 1' & 1^{(6)} \otimes 1^{(8)} = 1''' \\
 1^{(6)} \otimes 3 = 3' & 1^{(6)} \otimes 3' = 3'' & 1^{(6)} \otimes 3'' = 3 & 1^{(7)} \otimes 1^{(7)} = 1^{(8)} & 1^{(7)} \otimes 1^{(8)} = 1 \\
 1^{(7)} \otimes 3 = 3' & 1^{(7)} \otimes 3' = 3'' & 1^{(7)} \otimes 3'' = 3 & 1^{(8)} \otimes 1^{(8)} = 1^{(7)} & 1^{(8)} \otimes 3 = 3'' \\
 \mathbf{1^{(8)} \otimes 3' = 3} & 1^{(8)} \otimes 3'' = 3' & & & 
 \end{array}$$

$$\begin{array}{l}
 3 \otimes 3 = 1 + 1''' + 1^{(4)} + 2 \otimes 3 \\
 3 \otimes 3'' = 1' + 1^{(5)} + 1^{(8)} + 2 \otimes 3'' \\
 \mathbf{3' \otimes 3'' = 1 + 1''' + 1^{(4)} + 2 \otimes 3}
 \end{array}$$

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 \mathbf{3' \otimes 3' = 1' + 1^{(5)} + 1^{(8)} + 2 \otimes 3''} \\
 \mathbf{3'' \otimes 3'' = 1'' + 1^{(6)} + 1^{(7)} + 2 \otimes 3'}
 \end{array}$$



# Invariant Lagrangian

- Terms that are invariant, have 2 leptons and mass dimension  $\leq 5$ :

$$LL h_u h_u \varphi_S + LL h_u h_u \xi + L e h_d \varphi_T + L \mu h_d \varphi_T + L \tau h_d \varphi_T$$

# Invariant Lagrangian

➤ Terms that are invariant, have 2 leptons and mass dimension  $\leq 5$ :

$$LL h_u h_u \varphi_S + LL h_u h_u \xi + L e h_d \varphi_T + L \mu h_d \varphi_T + L \tau h_d \varphi_T$$

$$3' \otimes 3' \otimes 1 \otimes 1 \otimes 3' = (1' + 1^{(5)} + 1^{(8)} + 2 \times 3'') \otimes 3' = 2 \times 1 + 2 \times 1''' + 2 \times 1^{(4)} + 7 \times 3$$

# Invariant Lagrangian

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- Contract family indices:

$$\begin{aligned} \frac{1}{\sqrt{3}} L_2 L_3 h_u h_u \varphi_{S,1} + \frac{1}{\sqrt{3}} L_3 L_1 h_u h_u \varphi_{S,2} + \frac{1}{\sqrt{3}} L_1 L_2 h_u h_u \varphi_{S,3} + \frac{1}{\sqrt{3}} L_1 L_1 h_u h_u \xi \\ + \frac{1}{\sqrt{3}} L_2 L_2 h_u h_u \xi + \frac{1}{\sqrt{3}} L_3 L_3 h_u h_u \xi \end{aligned}$$

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$$L L h_u h_u \varphi_S + L L h_u h_u \xi + L e h_d \varphi_T + L \mu h_d \varphi_T + L \tau h_d \varphi_T$$

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- Contract SU(2) indices and substitute vevs  $\langle \varphi_S \rangle = (v_S, 0, 0)$ , etc:

$$\frac{1}{\sqrt{3}} L_2^{(1)} L_3^{(1)} v_u v_u v_S + \frac{1}{\sqrt{3}} L_1^{(1)} L_1^{(1)} v_u v_u \xi + \frac{1}{\sqrt{3}} L_2^{(1)} L_2^{(1)} v_u v_u \xi + \frac{1}{\sqrt{3}} L_3^{(1)} L_3^{(1)} v_u v_u \xi$$

# Finally, the Mixing Angles

➤ Mass matrices

$$M_{\ell^+} = \begin{matrix} & e & \mu & \tau \\ \begin{matrix} L_1^{(2)} \\ L_2^{(2)} \\ L_3^{(2)} \end{matrix} & \begin{pmatrix} -\frac{1}{\sqrt{3}} & -\frac{1}{\sqrt{3}} & -\frac{1}{\sqrt{3}} \\ -\frac{1}{\sqrt{3}} & \frac{1}{2\sqrt{3}} & \frac{1}{2\sqrt{3}} \\ -\frac{1}{\sqrt{3}} & \frac{1}{2\sqrt{3}} & \frac{1}{2\sqrt{3}} \end{pmatrix}, & & \end{matrix}$$

$$M_\nu = \begin{matrix} & L_1^{(1)} & L_2^{(1)} & L_3^{(1)} \\ \begin{matrix} L_1^{(1)} \\ L_2^{(1)} \\ L_3^{(1)} \end{matrix} & \begin{pmatrix} \frac{1}{\sqrt{3}} & 0 & 0 \\ 0 & \frac{1}{\sqrt{3}} & \frac{1}{2\sqrt{3}} \\ 0 & \frac{1}{2\sqrt{3}} & \frac{1}{\sqrt{3}} \end{pmatrix} & & \end{matrix}$$

# Finally, the Mixing Angles

➤ Mass matrices

$$M_{\ell^+} = \begin{matrix} & e & \mu & \tau \\ L_1^{(2)} & \left( -\frac{1}{\sqrt{3}} & -\frac{1}{\sqrt{3}} & -\frac{1}{\sqrt{3}} \right) \\ L_2^{(2)} & \left( -\frac{1}{\sqrt{3}} & \frac{1}{2\sqrt{3}} & \frac{1}{2\sqrt{3}} \right) \\ L_3^{(2)} & \left( -\frac{1}{\sqrt{3}} & \frac{1}{2\sqrt{3}} & \frac{1}{2\sqrt{3}} \right) \end{matrix}, \quad M_\nu = \begin{matrix} & L_1^{(1)} & L_2^{(1)} & L_3^{(1)} \\ L_1^{(1)} & \left( \frac{1}{\sqrt{3}} & 0 & 0 \right) \\ L_2^{(1)} & \left( 0 & \frac{1}{\sqrt{3}} & \frac{1}{2\sqrt{3}} \right) \\ L_3^{(1)} & \left( 0 & \frac{1}{2\sqrt{3}} & \frac{1}{\sqrt{3}} \right) \end{matrix}$$

➤ Singular value decomposition:  $\hat{M}_{\ell^+} = D_L M_{\ell^+} D_R^\dagger$ ,  $\hat{M}_\nu = U_L M_\nu U_R^\dagger$

$$D_L = \begin{pmatrix} -0.5774+i0.0000 & -0.5774+i0.0000 & -0.5774+i0.0000 \\ 0.5738-i0.0636 & -0.2319+i0.5287 & -0.3420-i0.4652 \\ 0.5731-i0.0702 & -0.3474-i0.4612 & -0.2257+i0.5314 \end{pmatrix}, \quad U_L = \begin{pmatrix} 0.0000 & -0.7071 & -0.7071 \\ -1.0000 & 0.0000 & 0.0000 \\ 0.0000 & 0.7071 & -0.7071 \end{pmatrix}$$

# Finally, the Mixing Angles

➤ Mass matrices

$$M_{\ell^+} = \begin{matrix} & e & \mu & \tau \\ L_1^{(2)} & \left( -\frac{1}{\sqrt{3}} & -\frac{1}{\sqrt{3}} & -\frac{1}{\sqrt{3}} \right) \\ L_2^{(2)} & \left( -\frac{1}{\sqrt{3}} & \frac{1}{2\sqrt{3}} & \frac{1}{2\sqrt{3}} \right) \\ L_3^{(2)} & \left( -\frac{1}{\sqrt{3}} & \frac{1}{2\sqrt{3}} & \frac{1}{2\sqrt{3}} \right) \end{matrix}, \quad M_\nu = \begin{matrix} & L_1^{(1)} & L_2^{(1)} & L_3^{(1)} \\ L_1^{(1)} & \left( \frac{1}{\sqrt{3}} & 0 & 0 \right) \\ L_2^{(1)} & \left( 0 & \frac{1}{\sqrt{3}} & \frac{1}{2\sqrt{3}} \right) \\ L_3^{(1)} & \left( 0 & \frac{1}{2\sqrt{3}} & \frac{1}{\sqrt{3}} \right) \end{matrix}$$

➤ Singular value decomposition:  $\hat{M}_{\ell^+} = D_L M_{\ell^+} D_R^\dagger$ ,  $\hat{M}_\nu = U_L M_\nu U_R^\dagger$

$$D_L = \begin{pmatrix} -0.5774+i0.0000 & -0.5774+i0.0000 & -0.5774+i0.0000 \\ 0.5738-i0.0636 & -0.2319+i0.5287 & -0.3420-i0.4652 \\ 0.5731-i0.0702 & -0.3474-i0.4612 & -0.2257+i0.5314 \end{pmatrix}, \quad U_L = \begin{pmatrix} 0.0000 & -0.7071 & -0.7071 \\ -1.0000 & 0.0000 & 0.0000 \\ 0.0000 & 0.7071 & -0.7071 \end{pmatrix}$$

➤ Neutrino mixing matrix:  $U_{\text{PMNS}} = D_L U_L^\dagger$  (needs rephasing)

$$U_{\text{PMNS}} = \begin{pmatrix} 0.8165 + i0.0000 & 0.5774 + i0.0000 & 0.0000 + i0.0000 \\ 0.4058 - i0.0449 & -0.5738 + i0.0636 & 0.0778 + i0.7028 \\ 0.4052 - i0.0497 & -0.5731 + i0.0702 & -0.0860 - i0.7019 \end{pmatrix}$$

➤ Mixing angles:  $\theta_{12} = 35.26$ ,  $\theta_{23} = 45.00$ ,  $\theta_{13} = 0.00$  Tribimaximal ✓

# Finally, the Mixing Angles

➤ Mass matrices

$$M_{\ell^+} = \begin{matrix} & e & \mu & \tau \\ L_1^{(2)} & \left( -\frac{1}{\sqrt{3}} & -\frac{1}{\sqrt{3}} & -\frac{1}{\sqrt{3}} \right) \\ L_2^{(2)} & \left( -\frac{1}{\sqrt{3}} & \frac{1}{2\sqrt{3}} & \frac{1}{2\sqrt{3}} \right) \\ L_3^{(2)} & \left( -\frac{1}{\sqrt{3}} & \frac{1}{2\sqrt{3}} & \frac{1}{2\sqrt{3}} \right) \end{matrix}, \quad M_\nu = \begin{matrix} & L_1^{(1)} & L_2^{(1)} & L_3^{(1)} \\ L_1^{(1)} & \left( \frac{1}{\sqrt{3}} & 0 & 0 \right) \\ L_2^{(1)} & \left( 0 & \frac{1}{\sqrt{3}} & \frac{1}{2\sqrt{3}} \right) \\ L_3^{(1)} & \left( 0 & \frac{1}{2\sqrt{3}} & \frac{1}{\sqrt{3}} \right) \end{matrix}$$

➤ Singular value decomposition:  $\hat{M}_{\ell^+} = D_L M_{\ell^+} D_R^\dagger$ ,  $\hat{M}_\nu = U_L M_\nu U_R^\dagger$

$$D_L = \begin{pmatrix} -0.5774+i0.0000 & -0.5774+i0.0000 & -0.5774+i0.0000 \\ 0.5738-i0.0636 & -0.2319+i0.5287 & -0.3420-i0.4652 \\ 0.5731-i0.0702 & -0.3474-i0.4612 & -0.2257+i0.5314 \end{pmatrix}, \quad U_L = \begin{pmatrix} 0.0000 & -0.7071 & -0.7071 \\ -1.0000 & 0.0000 & 0.0000 \\ 0.0000 & 0.7071 & -0.7071 \end{pmatrix}$$

➤ Neutrino mixing matrix:  $U_{\text{PMNS}} = D_L U_L^\dagger$  (needs rephasing)

$$|U_{\text{PMNS}}| = \begin{pmatrix} 0.8165 & 0.5774 & 0.0000 \\ 0.4082 & 0.5774 & 0.7071 \\ 0.4082 & 0.5774 & 0.7071 \end{pmatrix}$$

➤ Mixing angles:  $\theta_{12} = 35.26$ ,  $\theta_{23} = 45.00$ ,  $\theta_{13} = 0.00$  **Tribimaximal** ✓



# Finally, the Mixing Angles

➤ Mass matrices

$$M_{\ell^+} = \begin{matrix} & e & \mu & \tau \\ L_1^{(2)} & \left( -\frac{1}{\sqrt{3}} & -\frac{1}{\sqrt{3}} & -\frac{1}{\sqrt{3}} \right) \\ L_2^{(2)} & \left( -\frac{1}{\sqrt{3}} & \frac{1}{2\sqrt{3}} & \frac{1}{2\sqrt{3}} \right) \\ L_3^{(2)} & \left( -\frac{1}{\sqrt{3}} & \frac{1}{2\sqrt{3}} & \frac{1}{2\sqrt{3}} \right) \end{matrix}, \quad M_\nu = \begin{matrix} & L_1^{(1)} & L_2^{(1)} & L_3^{(1)} \\ L_1^{(1)} & \left( \frac{1}{\sqrt{3}} & 0 & 0 \right) \\ L_2^{(1)} & \left( 0 & \frac{1}{\sqrt{3}} & \frac{1}{2\sqrt{3}} \right) \\ L_3^{(1)} & \left( 0 & \frac{1}{2\sqrt{3}} & \frac{1}{\sqrt{3}} \right) \end{matrix}$$

➤ Singular value decomposition:  $\hat{M}_{\ell^+} = D_L M_{\ell^+} D_R^\dagger$ ,  $\hat{M}_\nu = U_L M_\nu U_R^\dagger$

$$D_L = \begin{pmatrix} -0.5774+i0.0000 & -0.5774+i0.0000 & -0.5774+i0.0000 \\ 0.5738-i0.0636 & -0.2319+i0.5287 & -0.3420-i0.4652 \\ 0.5731-i0.0702 & -0.3474-i0.4612 & -0.2257+i0.5314 \end{pmatrix}, \quad U_L = \begin{pmatrix} 0.0000 & -0.7071 & -0.7071 \\ -1.0000 & 0.0000 & 0.0000 \\ 0.0000 & 0.7071 & -0.7071 \end{pmatrix}$$

➤ Neutrino mixing matrix:  $U_{\text{PMNS}} = D_L U_L^\dagger$  (needs rephasing)

$$|U_{\text{PMNS}}| = \begin{pmatrix} \sqrt{2/3} & 1/\sqrt{3} & 0 \\ 1/\sqrt{6} & 1/\sqrt{3} & 1/\sqrt{2} \\ 1/\sqrt{6} & 1/\sqrt{3} & 1/\sqrt{2} \end{pmatrix}$$

➤ Mixing angles:  $\theta_{12} = 35.26$ ,  $\theta_{23} = 45.00$ ,  $\theta_{13} = 0.00$  Tribimaximal ✓

# What is the Bottom Line?

- We used the Altarelli-Feruglio model only as a paradigm
- The analysis is **completely independent** of the family symmetry
- GAP gives us all the relevant information about the group
- Complexity of group is hereby irrelevant
- We use Python to interact w/GAP and do symbolic manipulations
- From symmetry to lagrangian to mixing angles takes less than 5 ms/model

# Where Do We Go From Here?

- ① Generalize family symmetry:

$$A_4 \times C_3 \longrightarrow 1048 \text{ groups of order } \leq 100$$

- ② Particle content is the same:

$$L, e, \mu, \tau, h_u, h_d, \varphi_T, \varphi_S, \xi$$

- ③ Generalize family charge assignments:

$$\begin{aligned} (L, e, \mu, \tau, h_u, h_d, \varphi_T, \varphi_S, \xi) &\rightarrow (\mathbf{3}', \mathbf{1}', \mathbf{1}^{(8)}, \mathbf{1}^{(5)}, \mathbf{1}, \mathbf{1}, \mathbf{3}, \mathbf{3}', \mathbf{1}'') \\ &\rightarrow (*, *, *, *, *, *, *, *, *) \end{aligned}$$

- ④ Generalize symmetry breaking patterns:

$$\langle \varphi_T \rangle = (*, *, *), \quad \langle \varphi_S \rangle = (*, *, *), \quad \langle \xi \rangle = *$$

# Small Groups

The first few of the 1048 groups of order  $\leq 100$

✓ =  $U(n)$  and ✓ =  $SU(n)$  for  $n = 2, 3$

GAP ID	Group	<b>3</b>	U(3)	U(2)	U(2)×U(1)	$A_4$
[1, 1]	1	✗	✗	✗	✗	✗
[2, 1]	$C_2$	✗	✗	✗	✗	✗
[3, 1]	$C_3$	✗	✗	✗	✗	✗
[4, 1]	$C_4$	✗	✗	✗	✗	✗
[4, 2]	$C_2 \times C_2$	✗	✗	✗	✗	✗
[5, 1]	$C_5$	✗	✗	✗	✗	✗
[6, 1]	$S_3$	✗	✓	✓	✓	✗
[6, 2]	$C_6$	✗	✗	✗	✗	✗
[7, 1]	$C_7$	✗	✗	✗	✗	✗
[8, 1]	$C_8$	✗	✗	✗	✗	✗

## Small Groups

The first few of the 1048 groups of order  $\leq 100$

✓ =  $U(n)$  and ✓ =  $SU(n)$  for  $n = 2, 3$

GAP ID	Group	3	U(3)	U(2)	U(2)×U(1)	$A_4$
[8, 2]	$C_4 \times C_2$	✗	✗	✗	✗	✗
[8, 3]	$D_4$	✗	✓	✓	✓	✗
[8, 4]	$Q_8$	✗	✓	✓	✓	✗
[8, 5]	$C_2 \times C_2 \times C_2$	✗	✗	✗	✗	✗
[9, 1]	$C_9$	✗	✗	✗	✗	✗
[9, 2]	$C_3 \times C_3$	✗	✗	✗	✗	✗
[10, 1]	$D_5$	✗	✓	✓	✓	✗
[10, 2]	$C_{10}$	✗	✗	✗	✗	✗
[11, 1]	$C_{11}$	✗	✗	✗	✗	✗
[12, 1]	$C_3 \rtimes_{\varphi} C_4$	✗	✓	✓	✓	✗

## Small Groups

The first few of the 1048 groups of order  $\leq 100$

✓ =  $U(n)$  and ✓ =  $SU(n)$  for  $n = 2, 3$

GAP ID	Group	3	U(3)	U(2)	U(2)×U(1)	$A_4$
[12, 2]	$C_{12}$	X	X	X	X	X
[12, 3]	$A_4$	✓	✓	X	X	✓
[12, 4]	$D_6$	X	✓	✓	✓	X
[12, 5]	$C_6 \times C_2$	X	X	X	X	X
[13, 1]	$C_{13}$	X	X	X	X	X
[14, 1]	$D_7$	X	✓	✓	✓	X
[14, 2]	$C_{14}$	X	X	X	X	X
[15, 1]	$C_{15}$	X	X	X	X	X
[16, 1]	$C_{16}$	X	X	X	X	X
[16, 2]	$C_4 \times C_4$	X	X	X	X	X

## Small Groups

The first few of the 1048 groups of order  $\leq 100$

✓ = U(n) and ✓ = SU(n) for  $n = 2, 3$

GAP ID	Group	3	U(3)	U(2)	U(2)×U(1)	A <sub>4</sub>
[16, 3]	$(C_4 \times C_2) \rtimes_{\varphi} C_2$	X	✓	X	✓	X
[16, 4]	$C_4 \rtimes_{\varphi} C_4$	X	✓	X	✓	X
[16, 5]	$C_8 \times C_2$	X	X	X	X	X
[16, 6]	$C_8 \rtimes_{\varphi} C_2$	X	✓	✓	✓	X
[16, 7]	$D_8$	X	✓	✓	✓	X
[16, 8]	$QD_8$	X	✓	✓	✓	X
[16, 9]	$Q_{16}$	X	✓	✓	✓	X
[16, 10]	$C_4 \times C_2 \times C_2$	X	X	X	X	X
[16, 11]	$C_2 \times D_4$	X	✓	X	✓	X
[16, 12]	$C_2 \times Q_8$	X	✓	X	✓	X

## Small Groups

The first few of the 1048 groups of order  $\leq 100$

✓ =  $U(n)$  and ✓ =  $SU(n)$  for  $n = 2, 3$

GAP ID	Group	3	U(3)	U(2)	U(2)×U(1)	$A_4$
[16, 13]	$(C_4 \times C_2) \rtimes_{\varphi} C_2$	✗	✓	✓	✓	✗
[16, 14]	$C_2 \times C_2 \times C_2 \times C_2$	✗	✗	✗	✗	✗
[17, 1]	$C_{17}$	✗	✗	✗	✗	✗
[18, 1]	$D_9$	✗	✓	✓	✓	✗
[18, 2]	$C_{18}$	✗	✗	✗	✗	✗
[18, 3]	$C_3 \times S_3$	✗	✓	✓	✓	✗
[18, 4]	$(C_3 \times C_3) \rtimes_{\varphi} C_2$	✗	✓	✗	✓	✗
[18, 5]	$C_6 \times C_3$	✗	✗	✗	✗	✗
[19, 1]	$C_{19}$	✗	✗	✗	✗	✗
[20, 1]	$C_5 \rtimes_{\varphi} C_4$	✗	✓	✓	✓	✗



# Small Groups

The first few of the 1048 groups of order  $\leq 100$

✓ =  $U(n)$  and ✓ =  $SU(n)$  for  $n = 2, 3$

GAP ID	Group	3	U(3)	U(2)	U(2)×U(1)	$A_4$
[20, 2]	$C_{20}$	X	X	X	X	X
[20, 3]	$C_5 \rtimes_{\varphi} C_4$	X	X	X	X	X
[20, 4]	$D_{10}$	X	✓	✓	✓	X
[20, 5]	$C_{10} \times C_2$	X	X	X	X	X
[21, 1]	$C_7 \rtimes_{\varphi} C_3$	✓	✓	X	X	X
[21, 2]	$C_{21}$	X	X	X	X	X
[22, 1]	$D_{11}$	X	✓	✓	✓	X
[22, 2]	$C_{22}$	X	X	X	X	X
[23, 1]	$C_{23}$	X	X	X	X	X
[24, 1]	$C_3 \rtimes_{\varphi} C_8$	X	✓	✓	✓	X

# Small Groups

The first few of the 1048 groups of order  $\leq 100$

✓ =  $U(n)$  and ✓ =  $SU(n)$  for  $n = 2, 3$

GAP ID	Group	3	U(3)	U(2)	U(2)×U(1)	$A_4$
[24, 2]	$C_{24}$	X	X	X	X	X
[24, 3]	$SL(2, 3)$	✓	✓	✓	✓	X
[24, 4]	$C_3 \rtimes_{\varphi} Q_8$	X	✓	✓	✓	X
[24, 5]	$C_4 \times S_3$	X	✓	✓	✓	X
[24, 6]	$D_{12}$	X	✓	✓	✓	X
[24, 7]	$C_2 \times (C_3 \rtimes_{\varphi} C_4)$	X	✓	X	✓	X
[24, 8]	$(C_6 \times C_2) \rtimes_{\varphi} C_2$	X	✓	✓	✓	X
[24, 9]	$C_{12} \times C_2$	X	X	X	X	X
[24, 10]	$C_3 \times D_4$	X	✓	✓	✓	X
[24, 11]	$C_3 \times Q_8$	X	✓	✓	✓	X

# Small Groups

The first few of the 1048 groups of order  $\leq 100$

✓ =  $U(n)$  and ✓ =  $SU(n)$  for  $n = 2, 3$

GAP ID	Group	3	U(3)	U(2)	U(2)×U(1)	$A_4$
[24, 12]	$S_4$	✓	✓	✗	✓	✓
[24, 13]	$C_2 \times A_4$	✓	✓	✗	✗	✓
[24, 14]	$C_2 \times C_2 \times S_3$	✗	✓	✗	✓	✗
[24, 15]	$C_6 \times C_2 \times C_2$	✗	✗	✗	✗	✗
[25, 1]	$C_{25}$	✗	✗	✗	✗	✗
[25, 2]	$C_5 \times C_5$	✗	✗	✗	✗	✗
[26, 1]	$D_{13}$	✗	✓	✓	✓	✗
[26, 2]	$C_{26}$	✗	✗	✗	✗	✗
[27, 1]	$C_{27}$	✗	✗	✗	✗	✗
[27, 2]	$C_9 \times C_3$	✗	✗	✗	✗	✗

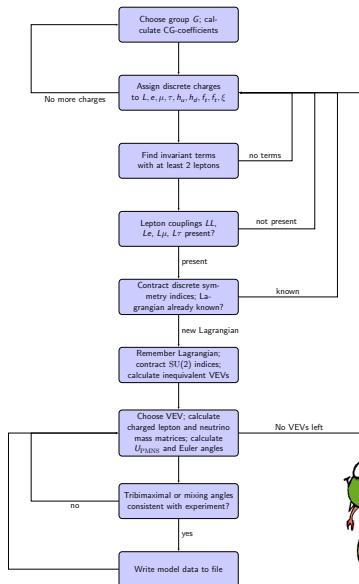
# Particles and Symmetry Breaking

## ➤ Assignment of family symmetry charges

Field	$SU(2)_L \times U(1)_Y$	$U(1)_R$	$A_4 \times C_3$	
$L$	$(\mathbf{2}, -1)$	1	$\mathbf{3}'$	→ Any 3-dim irrep
$e$	$(\mathbf{1}, 2)$	1	$\mathbf{1}'$	→ Any 1-dim irrep
$\mu$	$(\mathbf{1}, 2)$	1	$\mathbf{1}^{(8)}$	→ Any 1-dim irrep
$\tau$	$(\mathbf{1}, 2)$	1	$\mathbf{1}^{(5)}$	→ Any 1-dim irrep
$h_u$	$(\mathbf{2}, 1)$	0	$\mathbf{1}$	→ Any 1-dim irrep
$h_d$	$(\mathbf{2}, -1)$	0	$\mathbf{1}$	→ Any 1-dim irrep
$\varphi_T$	$(\mathbf{1}, 0)$	0	$\mathbf{3}$	→ Any irrep
$\varphi_S$	$(\mathbf{1}, 0)$	0	$\mathbf{3}'$	→ Any irrep
$\xi$	$(\mathbf{1}, 0)$	0	$\mathbf{1}''$	→ Any irrep

## ➤ Choose vevs that break to all subgroups

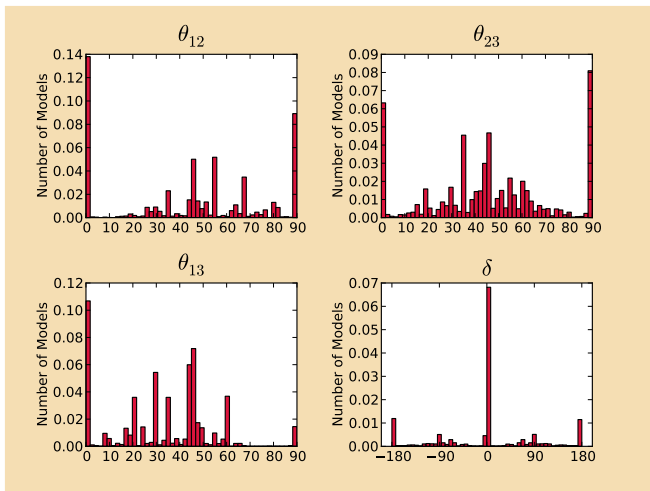
# Scanning for the Models



- Consider  $A_4 \times C_3$
- 14,594,580 different family charge assignments (particles w/same gauge/R-symmetry charges are considered identical)
- Computer (2.6 GHz, 2 GB) takes 19.35 hours  $\leadsto$  Distribute job to 10 machines and get results in 1.9 hours
- 41,086 different Lagrangians
- 8,526 cases of tribimaximal mixing, i.e. 20.75%
- Results for other groups on the way

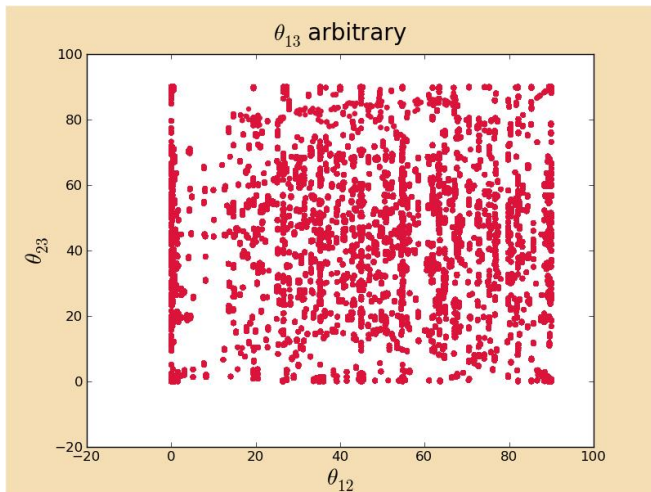
# Distribution of Angles

➤ Frequency of angles and phase



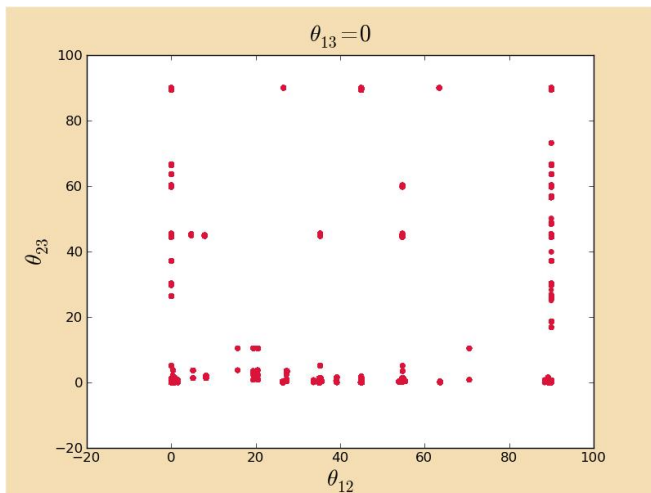
# Distribution of Angles

➤  $\theta_{12}$  vs.  $\theta_{23}$  for arbitrary  $\theta_{13}$



# Distribution of Angles

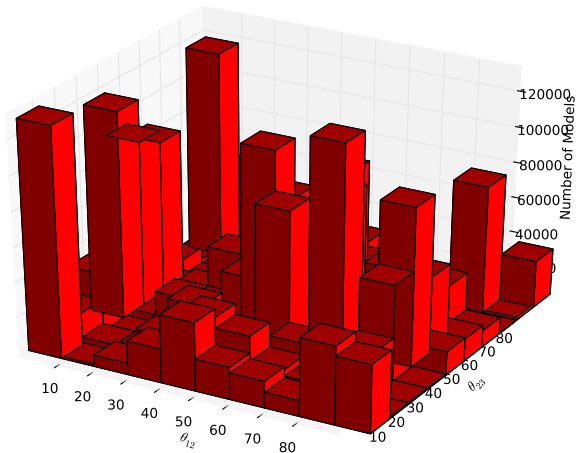
➤  $\theta_{12}$  vs.  $\theta_{23}$  for  $\theta_{13} = 0$





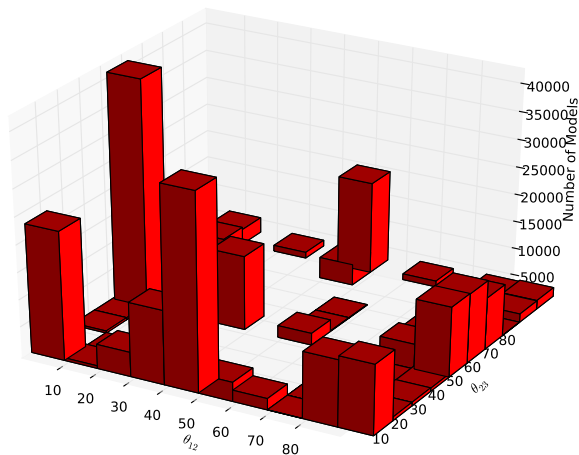
# Distribution of Angles

➤ Frequency of models:  $\theta_{12}$  vs.  $\theta_{23}$  for arbitrary  $\theta_{13}$



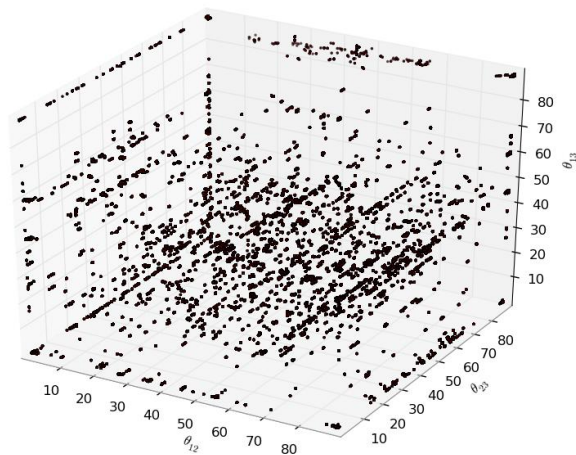
# Distribution of Angles

➤ Frequency of models:  $\theta_{12}$  vs.  $\theta_{23}$  for  $\theta_{13} = 0$



# Distribution of Angles

➤ Scatter plot for  $\theta_{12}$ ,  $\theta_{23}$ ,  $\theta_{13}$



# Conclusions

- Work in progress
- Disclaimer: I am not a landscaper
- Showing what one can do with data
- Our objective: Clarify connection between TBM and symmetry
- TBM occurs in 20.75% of all cases for  $A_4 \times C_3$ . Encouraging!
- TBM is not “privileged” or singled out, though
- Next: Extend analysis to other “small” groups