

Inert Neutralino Dark Matter in the E_6 SSM

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Motivations for the E_6 SSM

S. F. King, S. Moretti and R. Nevzorov [[arXiv:hep-ph/0510419](https://arxiv.org/abs/hep-ph/0510419)]

- The E_6 SSM is a string theory inspired model where the GUT group E_6 is broken below a large energy scale
- The model has automatic anomaly cancellation
- A type-I see-saw mechanism explains light neutrino masses
- Being a supersymmetric model, softly broken at the TeV scale, this model does not suffer from the hierarchy problem
- It also does not suffer from the μ problem of the MSSM
- In the E_6 SSM, like in other singlet extended models, the μ term is generated radiatively when a SM-singlet acquires a VEV, naturally related to the soft breaking scale

The E_6 SSM

S. F. King, S. Moretti and R. Nevzorov [arXiv:hep-ph/0510419]

- E_6 is broken below a large energy scale

$$\begin{aligned} E_6 &\rightarrow SO(10) \times U(1)_\psi \rightarrow SU(5) \times U(1)_\psi \times U(1)_\chi \\ &\rightarrow SU(3) \times SU(2) \times U(1)_Y \times U(1)_N \end{aligned}$$

- Anomaly cancellation is ensured by having complete 27 representations of E_6 (three generations i)

$$\begin{aligned} 27_i &\rightarrow (10, 1)_i + (\bar{5}, 2)_i \\ &\quad + (\bar{5}, -3)_i + (5, -2)_i \\ &\quad + (1, 5)_i + (1, 0)_i \end{aligned}$$

$$\mathcal{W} = \lambda_{ijk} S_i H_{dj} H_{uk} + \kappa_{ijk} S_i D_j \bar{D}_k + \dots$$

$$\mu = \lambda_{333} \langle S_3 \rangle = \lambda s / \sqrt{2}$$

- A generalised version of matter parity is imposed, but which allows exotics to decay
- There are two options under which exotics are interpreted as either di-quarks or lepto-quarks
- However, an approximate symmetry Z_2^H is also imposed in order to prevent FCNCs
- Under Z_2^H only S_3 , H_{d3} and H_{h3} are even
- The first and second generations of Higgs doublets and “Higgs singlets” then have suppressed couplings to ordinary matter
- This then explains why the first and second generations are inert

The E_6 SSM Neutralinos and Charginos

S. F. King, S. Moretti and R. Nevzorov [arXiv:hep-ph/0510419]

- In the USSM neutralino interaction basis $\tilde{\chi}_{\text{int}}^0 =$

$$\left(\tilde{B} \quad \tilde{W}^3 \quad \tilde{H}_{d3}^0 \quad \tilde{H}_{u3}^0 \mid \tilde{S}_3 \quad \tilde{B}' \right)^T$$

the USSM neutralino mass matrix $M_{\text{USSM}}^n =$

$$\left(\begin{array}{cccc|cc} M_1 & 0 & -m_Z s_W c_\beta & m_Z s_W s_\beta & 0 & 0 \\ 0 & M_2 & m_Z c_W c_\beta & -m_Z c_W s_\beta & 0 & 0 \\ -m_Z s_W c_\beta & m_Z c_W c_\beta & 0 & -\mu & -\mu_s s_\beta & g'_1 v c_\beta Q_d^N \\ m_Z s_W s_\beta & -m_Z c_W s_\beta & -\mu & 0 & -\mu_s c_\beta & g'_1 v s_\beta Q_u^N \\ \hline 0 & 0 & -\mu_s s_\beta & -\mu_s c_\beta & 0 & g'_1 s Q_s^N \\ 0 & 0 & g'_1 v c_\beta Q_d^N & g'_1 v s_\beta Q_u^N & g'_1 s Q_s^N & M'_1 \end{array} \right)$$

where

- $\mu_s = \lambda v / \sqrt{2}$,
- $\langle H_{d3} \rangle = v \cos(\beta) / \sqrt{2}$, and
- $\langle H_{u3} \rangle = v \sin(\beta) / \sqrt{2}$

- In the full E_6 SSM neutralino interaction basis $\tilde{\chi}_{\text{int}}^0 =$

$$\left(\tilde{B} \quad \tilde{W}^3 \quad \tilde{H}_{d3}^0 \quad \tilde{H}_{u3}^0 \mid \tilde{S} \quad \tilde{B}' \mid \tilde{H}_{d2}^0 \quad \tilde{H}_{u2}^0 \quad \tilde{S}_2 \mid \tilde{H}_{d1}^0 \quad \tilde{H}_{u1}^0 \quad \tilde{S}_1 \right)^T$$

the E_6 SSM neutralino mass matrix $M^n =$

$$\begin{pmatrix} M_{\text{USSM}}^n & B_2 & B_1 \\ B_2^T & A_{22} & A_{21} \\ B_1^T & A_{21}^T & A_{11} \end{pmatrix}$$

where

$$A_{\alpha\beta} = -\frac{1}{\sqrt{2}} \begin{pmatrix} 0 & \lambda_{\alpha\beta} s & f_{\beta\alpha}^u v \sin(\beta) \\ \lambda_{\beta\alpha} s & 0 & f_{\beta\alpha}^d v \cos(\beta) \\ f_{\alpha\beta}^u v \sin(\beta) & f_{\alpha\beta}^d v \cos(\beta) & 0 \end{pmatrix}$$

where $\lambda_{\alpha\beta} = \lambda_{3\alpha\beta}$, $f_{\beta\alpha}^d = \lambda_{\alpha\beta 3}$, and $f_{\beta\alpha}^u = \lambda_{\alpha\beta 3}$

The E_6 SSM Neutralinos and Charginos continued

J. P. Hall and S. F. King [arXiv:0905.2696 [hep-ph]]

- In the E_6 SSM chargino interaction basis $\tilde{\chi}_{\text{int}}^{\pm} =$

$$\left(\tilde{W}^+ \quad \tilde{H}_{u3}^+ \quad \tilde{H}_{u2}^+ \quad \tilde{H}_{u1}^+ \mid \tilde{W}^- \quad \tilde{H}_{d3}^- \quad \tilde{H}_{d2}^- \quad \tilde{H}_{d1}^- \right)^T$$

the E_6 SSM chargino mass matrix $M^c =$

$$\begin{pmatrix} & C^T \\ C & \end{pmatrix}$$

where

$$C = \begin{pmatrix} M_2 & \sqrt{2}m_W \sin(\beta) & 0 & 0 \\ \sqrt{2}m_W \cos(\beta) & \mu & \frac{1}{\sqrt{2}}x_2^d s & \frac{1}{\sqrt{2}}x_1^d s \\ 0 & \frac{1}{\sqrt{2}}x_2^u s & \frac{1}{\sqrt{2}}\lambda_{22} s & \frac{1}{\sqrt{2}}\lambda_{21} s \\ 0 & \frac{1}{\sqrt{2}}x_1^u s & \frac{1}{\sqrt{2}}\lambda_{12} s & \frac{1}{\sqrt{2}}\lambda_{11} s \end{pmatrix}$$

- As an example, demonstrating the relevant physics for determining the amount of dark matter, let us briefly consider the following parametrisation:

$$\lambda' = \lambda_{22} = \lambda_{11} \quad f = f_{22}^d = f_{22}^u = f_{11}^d = f_{11}^u$$

varying λ' , f and $\tan(\beta)$

- The A matrices become

$$A_{22} = A_{11} = -\frac{1}{\sqrt{2}} \begin{pmatrix} 0 & \lambda' s & f v \sin(\beta) \\ \lambda' s & 0 & f v \cos(\beta) \\ f v \sin(\beta) & f v \cos(\beta) & 0 \end{pmatrix}$$

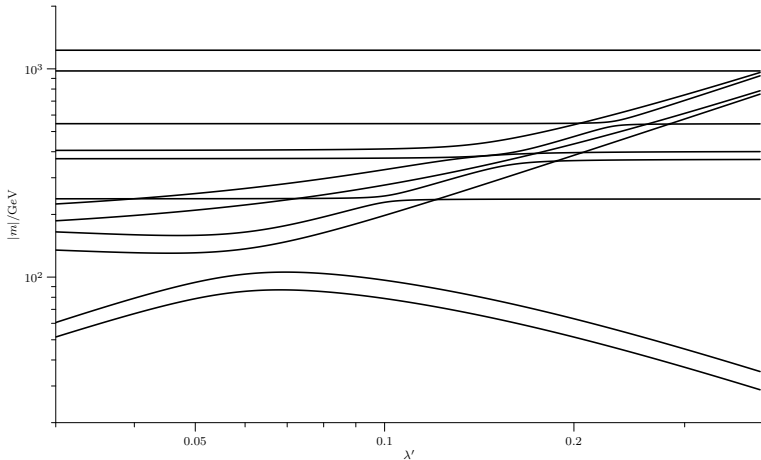
$$A_{21} = \epsilon A_{22}$$

- The inert scalars and exotics are assumed heavy

The Inert Neutralino LSP Neutralino spectrum

J. P. Hall and S. F. King [arXiv:0905.2696 [hep-ph]]

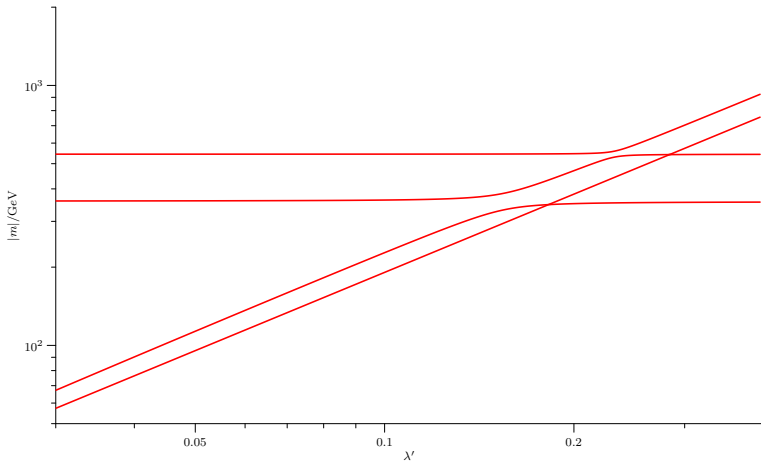
● $s = 3000 \text{ GeV}$, $\tan(\beta) = 1.5$, $f = 1$, $\epsilon = 0.1$



The Inert Neutralino LSP Chargino spectrum

J. P. Hall and S. F. King [arXiv:0905.2696 [hep-ph]]

• $s = 3000 \text{ GeV}$, $\tan(\beta) = 1.5$, $f = 1$, $\epsilon = 0.1$



The Inert Neutralino LSP

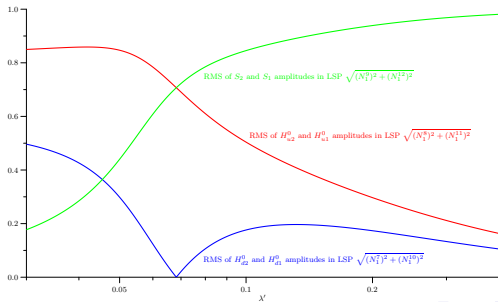
LSP composition

J. P. Hall and S. F. King [arXiv:0905.2696 [hep-ph]]

- The most important annihilation channels are



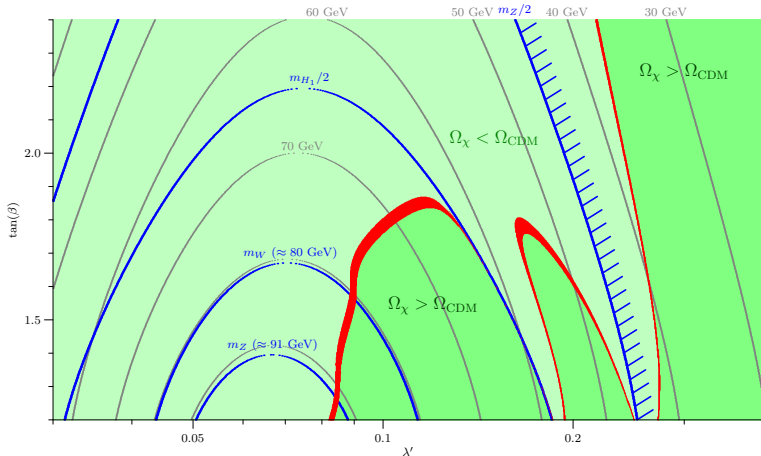
- $s = 3000$ GeV, $\tan(\beta) = 1.5$, $f = 1$, $\epsilon = 0.1$



Dark Matter Predictions MicrOMEGAs analysis

J. P. Hall and S. F. King [arXiv:0905.2696 [hep-ph]]

- $s = 3000 \text{ GeV}$, $f = 1$, $\epsilon = 0.1$

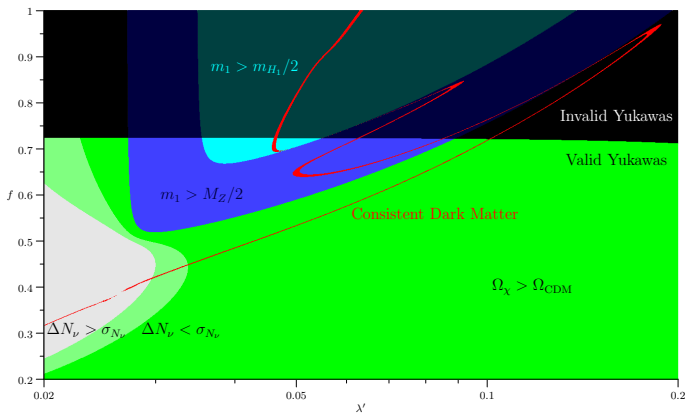


Dark Matter Predictions Additional considerations

J. P. Hall and S. F. King (in preparation)

• $s = 2400 \text{ GeV}, \tan(\beta) = 1.7$

$N_\nu = 2.9840 \pm 0.0082$



- This parametrisation is consistent with the E_6 SSM flavour theory proposed in *R. Howl and S. F. King [arXiv:0908.2067 [hep-ph]]*

Dark Matter Predictions

Implications for the lightest Higgs boson

J. P. Hall, S. F. King, R. Nevzorov, S. Pakvasa and M. Sher (in preparation)

- If the two lightest inert neutralinos are lighter than half of the Higgs mass they typically dominate its width
- This greatly reduces the chances of discovering the Higgs by its decay into photons
- As an example if $\text{Br}(H_1 \rightarrow \tilde{\chi}_2^0 \tilde{\chi}_2^0, \tilde{\chi}_2^0 \tilde{\chi}_1^0, \tilde{\chi}_1^0 \tilde{\chi}_1^0) \sim 95\%$, $\text{Br}(H_1 \rightarrow \gamma\gamma)$ is decreased by a factor of about 20
- It is possible that $H_1 \rightarrow \tilde{\chi}_2^0 \tilde{\chi}_2^0 \rightarrow \mu\mu\bar{\mu}\bar{\mu}\tilde{\chi}_1^0\tilde{\chi}_1^0$ could be detected, but the requisite mass splitting between $\tilde{\chi}_2^0$ and $\tilde{\chi}_1^0$ is not a typical feature

Conclusions

- The E_6 SSM is motivated for several reasons
- The observed dark matter relic density can be accounted for in this model
- The only hard limit on the USSM (MSSM) parameter space from dark matter considerations is that $\tan(\beta)$ should be less than about 2
- In fact the lightest USSM-like supersymmetric particle could even be charged
- We are currently looking at dark matter in a GUT constrained version of the full model, expanding on the version of the constrained E_6 SSM studied in *P. Athron, S. F. King, D. J. Miller, S. Moretti and R. Nevzorov* [[arXiv:0904.2169](https://arxiv.org/abs/0904.2169) [[hep-ph](#)]]