

BEYOND THE MINIMAL
COMPOSITE HIGGS MODEL
2 HIGGSES AS COMPOSITE PSEUDO-NGB'S

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work (to appear) in collaboration with: J.Mrazek, A.Pomarol, R.Rattazzi, M.Redi, A.Wulzer

WHY A COMPOSITE-NGB-HIGGS ?

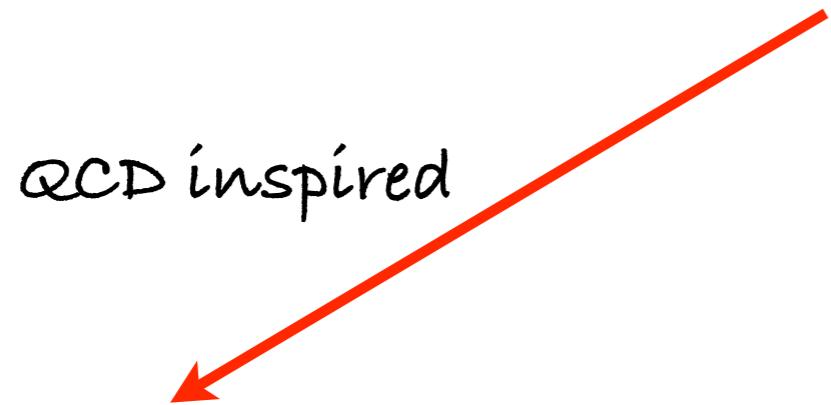
WHY A COMPOSITE-NGB-HIGGS ?

QCD inspired

STRONG
SECTOR

↓ dynamically

EW SCALE



**NO Naturalness or Hierarchy
problems**

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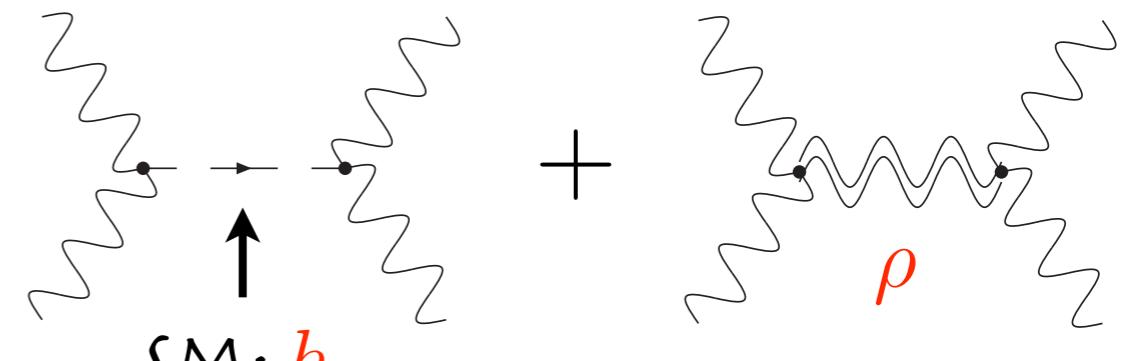
**NO Naturalness or Hierarchy
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dynamically
EW SCALE

consistency with EWPT (S , T)

$$\mathcal{M}(W_L W_L \rightarrow W_L W_L)$$



MSSM: h, H, A, H^\pm

$$m_\rho \sim 2.5 \text{ TeV}$$

resonances pushed up in energy



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QCD inspired

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**NO Naturalness or Hierarchy
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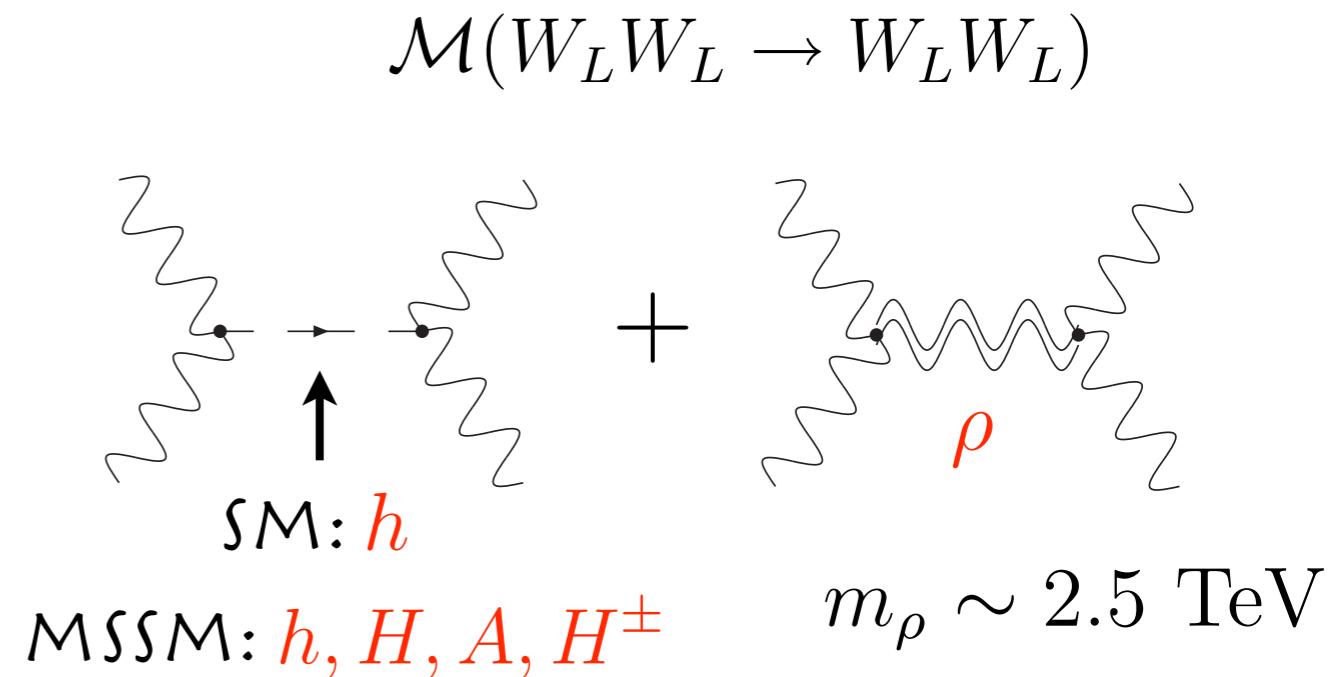


EW SCALE
dynamically

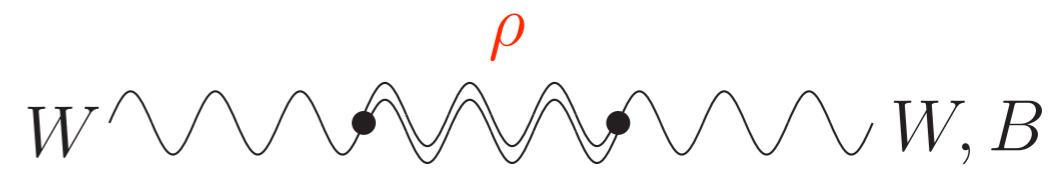
like QCD pions

Higgs lighter than other resonances

$$m_h \ll m_\rho$$



resonances pushed up in energy



WHAT DO WE NEED ?

Spontaneous Symmetry Breaking:

$$G \rightarrow H$$

$$\Sigma(h/f)$$

= G/H sigma model scale

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$$\Sigma(h/f)$$

$$i) \quad G \supset \mathrm{SU}(2)_L \times \mathrm{U}(1)_Y$$

= G/H sigma model scale

gauged

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$$\Sigma(h/f)$$

- i) $G \supset \text{SU}(2)_L \times \text{U}(1)_Y$
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= G/H sigma model scale

- ii) $H \supset \text{O}(4) \cong \text{SU}(2)_A \times \text{SU}(2)_B \times P_{AB}$

$$\delta g_{Zbb} = 0$$

$$\langle h \rangle = \begin{pmatrix} 0 \\ 0 \\ 0 \\ v \end{pmatrix} \quad \text{SO}(3) \cong \text{SU}(2)_{A+B}$$

Custodial symmetry

$$\delta \rho = 0$$

WHAT DO WE NEED ?

Spontaneous Symmetry Breaking:

$$G \rightarrow H$$

- i) $G \supset \text{SU}(2)_L \times \text{U}(1)_Y$ $\Sigma(h/f)$
gauged $= G/H$ sigma model scale

- ii) $H \supset \text{O}(4) \cong \text{SU}(2)_A \times \text{SU}(2)_B \times P_{AB}$ $\delta g_{Zbb} = 0$

- iii) $G/H \supset 4 = (2, 2)$ $\langle h \rangle = \begin{pmatrix} 0 \\ 0 \\ 0 \\ v \end{pmatrix}$ $\text{SO}(3) \cong \text{SU}(2)_{A+B}$
Custodial symmetry

h , Higgs complex doublet

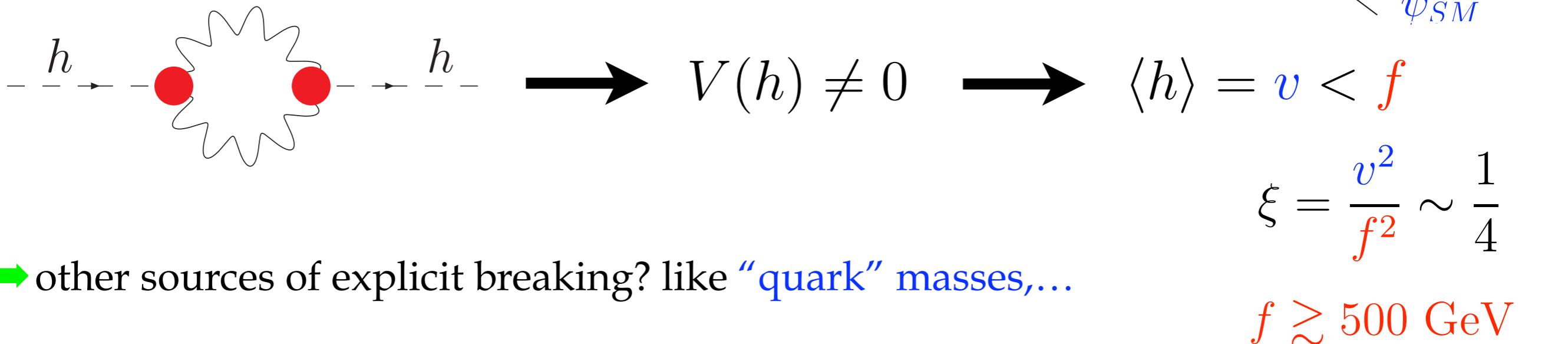
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WHAT DO WE NEED ?

Explicit Symmetry Breaking: from the Strong Sector alone: $V(h) = 0$

from interactions with SM fields:

- SM gauge interactions (subgroup K of G is gauged)
- SM fermion interactions (proto-Yukawas, mostly the top)

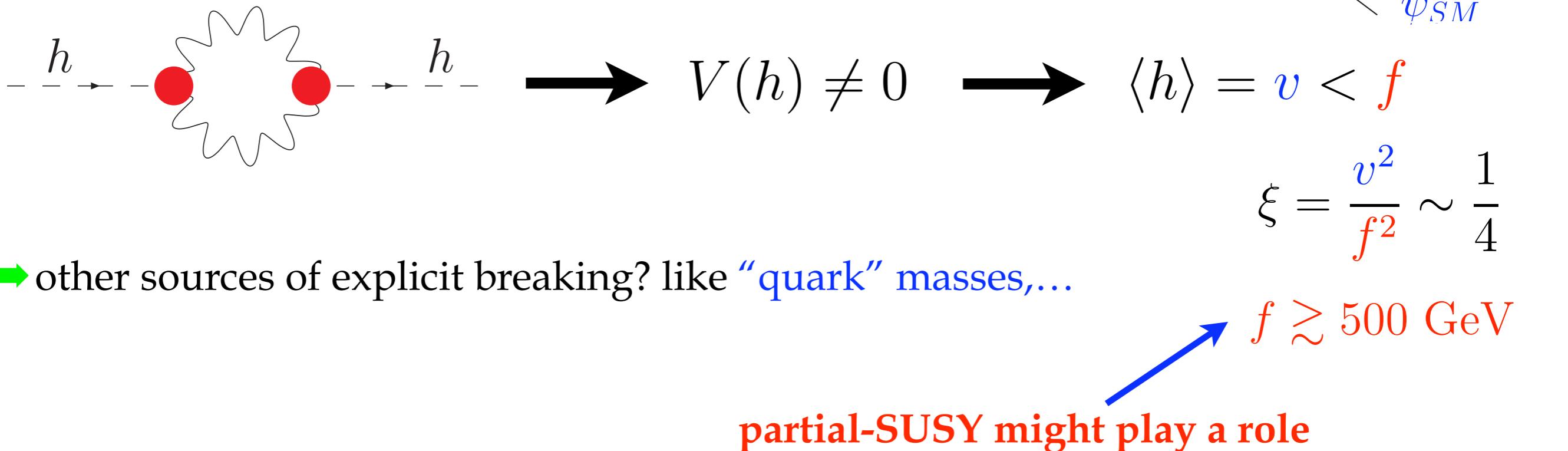


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R.Sundrum
B.Gripaios, M.Redi

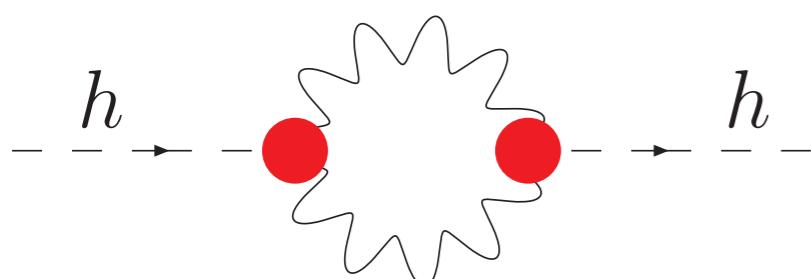
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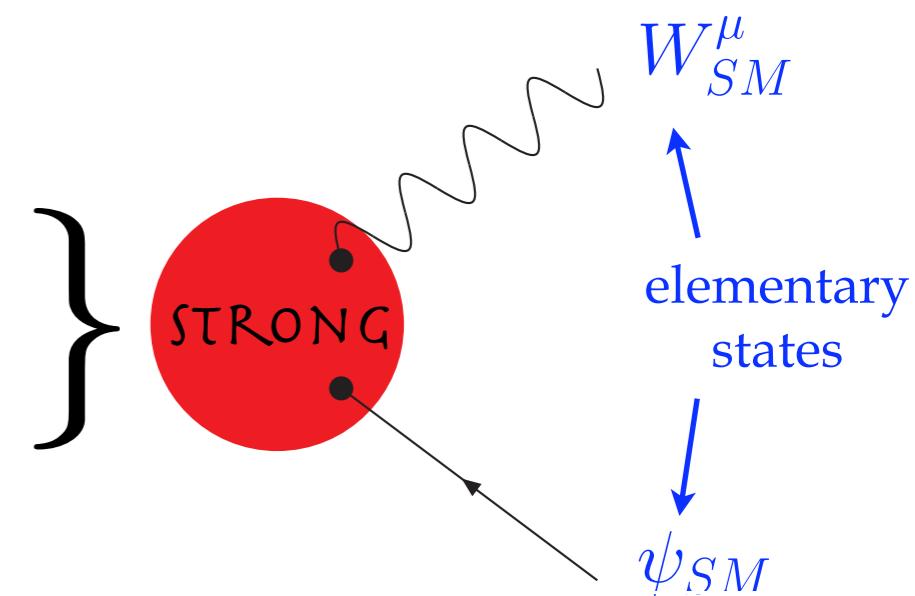
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$$V(h) \neq 0$$



$$\langle h \rangle = v < f$$

$$\xi = \frac{v^2}{f^2} \sim \frac{1}{4}$$

$$f \gtrsim 500 \text{ GeV}$$

- other sources of explicit breaking? like “quark” masses,...

Calculability: Effective low-energy “chiral” Lagrangian

+

Holography and Warped extra-d

WHAT DO WE GET ?

Minimal Composite Higgs Model:

$$G = \text{SO}(5) \rightarrow H = \text{O}(4) \quad G/H = \textcolor{blue}{4} = (2,2)$$

Pheno: new resonances + modifications of Higgs couplings

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Example of *minimal* cosets from constituent “quarks”:

$$\text{SU}(4) \rightarrow \text{Sp}(4)$$

$$G/H = 5 = \mathbf{4} + 1 = (2,2) + (1,1)$$

$$\text{SU}(5) \rightarrow \text{SO}(5)$$

$$G/H = 14$$

extra singlet
B.Gripaios, A.Pomarol, F.Riva, JS
J.Galloway, J.Evans, M.Luty, R.Tacchi

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$$\text{SO}(6) \rightarrow \text{SO}(4) \times \text{U}(1)$$

$$G/H = 8 = \mathbf{4}_{+2} + \mathbf{4}_{-2} = 2 \times (2,2)$$

2 Higgs doublets! and not MSSM

2 COMPOSITE-NGB-HIGGS: IS IT VIABLE?

$$h_i = (2, 2), \quad i = 1, 2$$

→ Large contributions to $\delta\rho$:

$$\langle h_1 \rangle = \begin{pmatrix} 0 \\ 0 \\ 0 \\ v_1 \end{pmatrix} \quad \langle h_2 \rangle = \begin{pmatrix} 0 \\ 0 \\ v_2 \\ 0 \end{pmatrix} \quad \left\{ \begin{array}{l} \text{SO}(2) \cong \text{U}(1)_Q \\ \text{custodial SO(3) breaking} \end{array} \right.$$

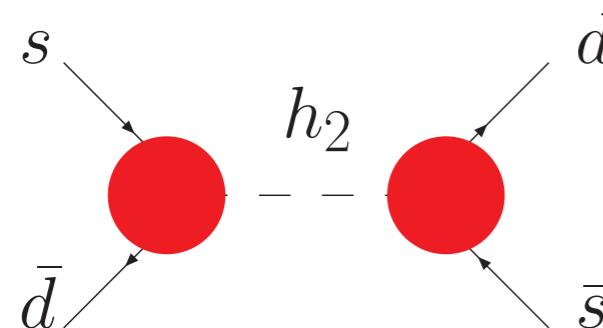
misaligned VEV's

$$\frac{c_T}{f^2} (h_1^T D_\mu h_2)^2 \rightarrow \delta\rho \sim c_T \frac{v^2}{f^2}$$

$$v^2/f^2 \lesssim 1 \times 10^{-3}$$

→ Dangerous Flavor transitions:

$$\bar{q}_L (y_1^u \tilde{h}_1 + y_2^u \tilde{h}_2) u_R + \bar{q}_L (y_1^d \tilde{h}_1 + y_2^d \tilde{h}_2) d_R \quad \text{misaligned Yukawa's}$$



$$\sim \frac{m_s^2 V_{us}^2}{2 m_{h_2}^2} (\bar{s}_L d_R)^2 \rightarrow m_{h_2} \gtrsim 2 \text{ TeV}$$

INERT COMPOSITE-HIGGS MODEL

Defined by extra \mathbf{Z}_2 :

$$h_1 \rightarrow + h_1$$

$$h_2 \rightarrow - h_2 \quad + \quad \langle h_2 \rangle = 0$$

$$\text{SM} \rightarrow + \text{SM}$$

\mathbf{Z}_2 exactly conserved

NO $(h_2)^{2n+1}$ couplings

Immediate consequences:

- ➡ Automatic solution of $\delta\rho$ and Flavor problems.
- ➡ Stability of lightest h_2 component (DM candidate).
- ➡ Double h_2 production @ colliders.

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$$G = \mathbf{O}(6) \longrightarrow H = \mathbf{O}(4) \times \mathbf{O}(2)$$

\mathbf{Z}_2 = outer automorphism of $\mathbf{SO}(6)$

INERT HIGGS: $O(6)/O(4) \times O(2)$

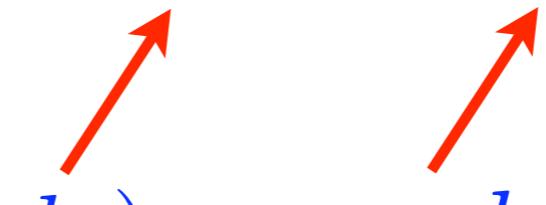
Breakings:

$(SU(2)_L, SU(2)_R)_{Z_2}$

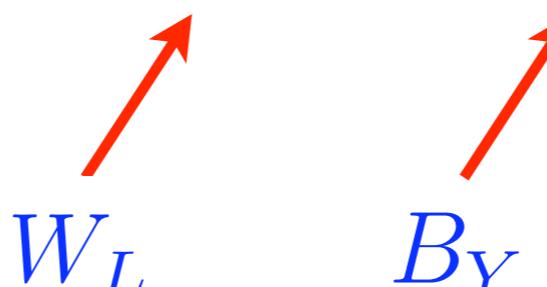
spontaneous, $\Sigma(h_1, h_2) = SO(6)/SO(4) \times SO(2) \in \mathbf{20'}$

$$= (2, 2)_+ + (2, 2)_-$$

explicit, $\psi_{SM} \in \mathbf{6} \quad \mathbf{6} = (2, 2)_+ + (1, 1)_+ + (1, 1)_-$

$$q_L = (u_L, d_L) \quad u_R, d_R$$


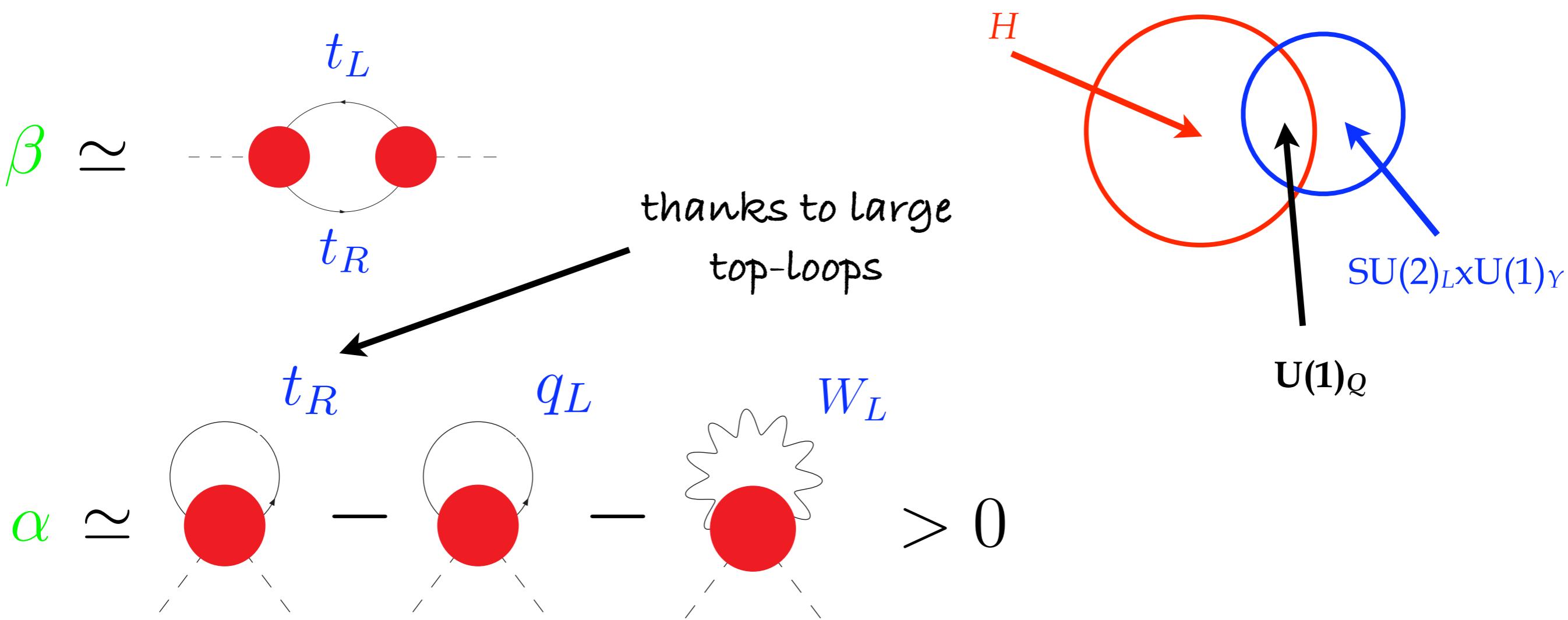
$$W_{SM}^\mu \in \mathbf{15}$$

$$\mathbf{15} = (1, 1)_+ + (3, 1)_+ + (1, 3)_+ + (2, 2)_+ + (2, 2)_-$$
$$W_L \quad B_Y$$


INERT HIGGS: $O(6)/O(4)\times O(2)$

EW Symmetry Breaking: $\langle h_2 \rangle = 0$

$$V(h_1) = \alpha \cos(h_1/f) - \beta \sin^2(h_1/f) \rightarrow \cos(h_1/f) = -\alpha/2\beta$$

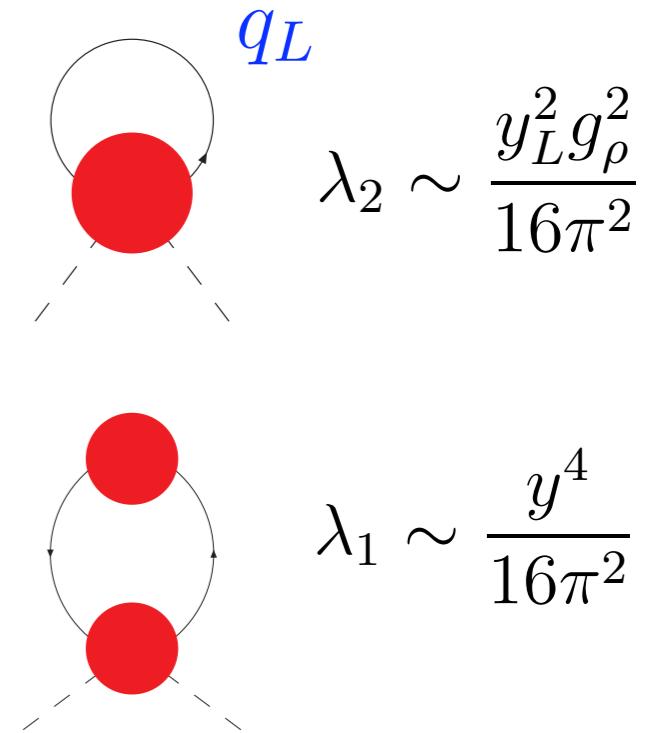


INERT HIGGS: $O(6)/O(4)\times O(2)$

Spectrum:

$$\left. \begin{array}{c} \equiv \\ \equiv \\ \equiv \end{array} \right\} \sim m_{h_2} \sim \sqrt{\lambda_2} \ f \sim 450 \text{ GeV}$$

$$m_{h_1} \sim \sqrt{\lambda_1} \ v \sim 150 \text{ GeV}$$



$$\lambda_1 \sim \frac{y^4}{16\pi^2}$$

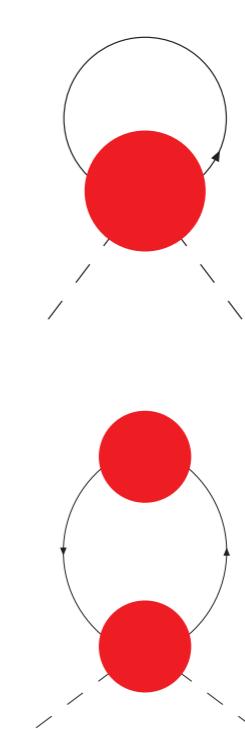
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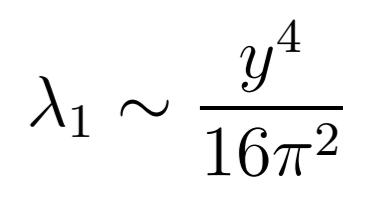
$$\left. \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \right\} \sim m_{h_2} \sim \sqrt{\lambda_2} f \sim 450 \text{ GeV}$$

$\text{---} \quad m_{h_1} \sim \sqrt{\lambda_1} v \sim 150 \text{ GeV}$





$$\lambda_2 \sim \frac{y_L^2 g_\rho^2}{16\pi^2}$$



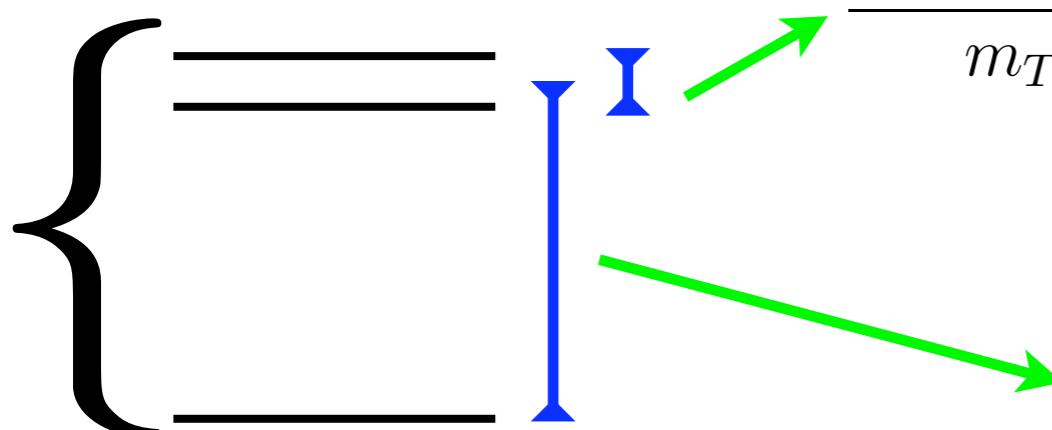
$$\lambda_1 \sim \frac{y^4}{16\pi^2}$$

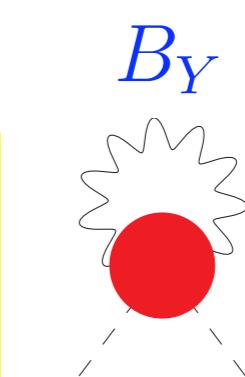
splittings, $h_2 \rightarrow T = (H^\pm, A) \oplus S$

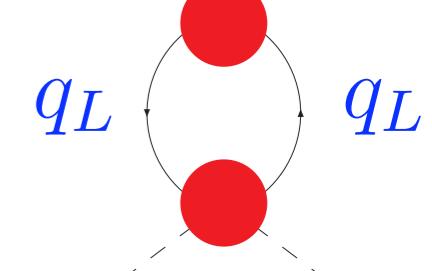
$$SO(4) \quad \mathbf{4} = \mathbf{3} + \mathbf{1} \quad SO(3)$$

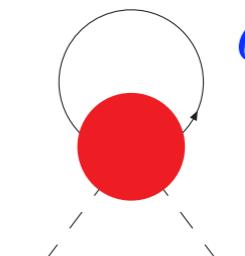
$$\left\{ \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \right\} \xrightarrow{\text{I}} \frac{|m_H^\pm - m_A|}{m_T} \sim \left(\frac{g'^2}{y_L^2} \pm \frac{y_L^2}{g_\rho^2} \right) \frac{v^2}{f^2} \sim \boxed{\frac{1}{40}}$$

$\xrightarrow{\text{II}} \frac{|m_S - m_T|}{m_T} \sim \frac{v^2}{f^2} \sim \boxed{\frac{1}{4}}$





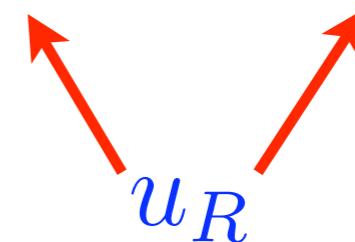




$$q_L$$

PSEUDO-INERT HIGGS: EXPLICIT Z_2 BREAKING

explicit, $\psi_{SM} \in \mathbf{6}$ $\mathbf{6} = (\mathbf{2}, \mathbf{2})_+ + (\mathbf{1}, \mathbf{1})_+ + (\mathbf{1}, \mathbf{1})_-$



$O(2) \rightarrow Z'_2$ $u_R \in (\mathbf{1}, \mathbf{1})_+$,



we need collective breaking

$$v_2 \sim \frac{(y_R^t)^2 (y_R^b)^2}{y_L^2 g_\rho^2} \frac{1}{16\pi^2} v_1 \quad \text{large } \tan \beta$$

COSET INTERACTIONS: $O(6)/O(4)\times O(2)$

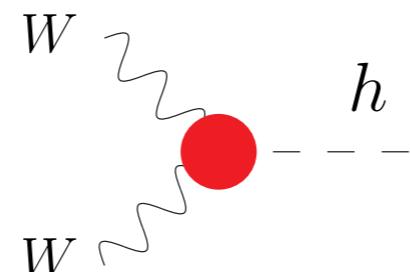
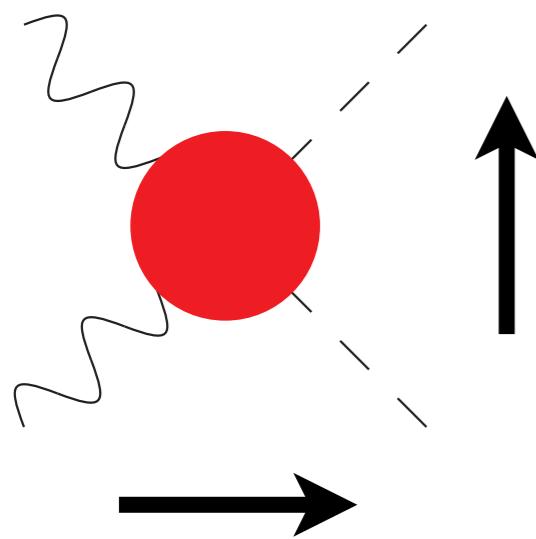
Signatures of composite-NGB:

$$\mathcal{L}_{coset}[h_1/f, h_2/f] = \frac{f^2}{2} \text{Tr}[|D_\mu \Sigma|^2]$$

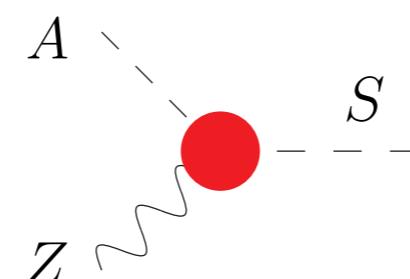
$$= \frac{1}{2}(D_\mu h_j)^2 + \frac{c_\partial}{f^2}(h_i \partial_\mu h_i)^2 + \frac{c_W}{f^2}(h_i W_\mu h_j)^2 + \frac{c_{\partial W}}{f^2}(h_i \partial_\mu h_i)(h_j W^\mu h_j)$$

(no CPV or large $\delta\rho$)

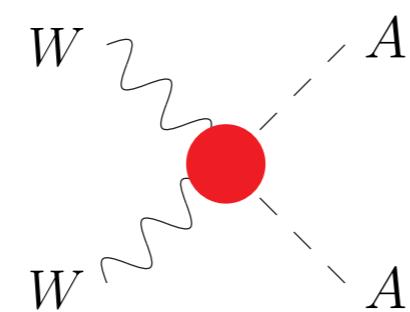
non-unitarity of amplitudes,



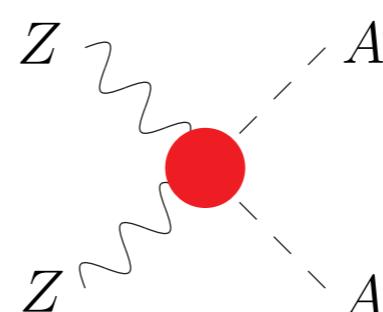
$$\sum_i g_{h_i^0 WW}^2 = g^2 m_W^2 \left(1 - \frac{1}{2} \frac{v^2}{f^2} \right)$$



$$\sum_i g_{h_i^0 AZ}^2 = \frac{g^2}{\cos^2 \theta_W} \left(1 - \frac{1}{2} \frac{v^2}{f^2} \right)$$



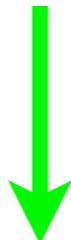
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$$= \cos^2 \theta_W \left(1 + \frac{v^2}{f^2} \right)$$

CONCLUSIONS

If the hierarchy problem is solved by **strong dynamics**, we can expect rich phenomenology of pseudo-NGB's.



motivated framework for extended Higgs sectors

example,

Inert composite-Higgs model (not near MSSM)

- completely **viable**
- predictions: large h_2 masses, $\mathbf{m_{h2} \sim 450 \text{ GeV}}$
small custodial splittings, $|\mathbf{m_{H^+} - m_A}| \sim 10 \text{ GeV}$
 $\mathbf{O(v^2/f^2) \sim 20\%}$ deviations in couplings