

Axion Miniclusters and Structure Formation

Henrique Rubira

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Preamble

Let's talk about three observables

The matter power spectrum

The halo mass function

The halo profile

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Let's talk about three observables

The matter power spectrum

The halo mass function

The halo profile

Let's not talk about problems to connect those observables to galaxy surveys e.g., galaxy bias, completeness, purity, baryonic feedback

Continuity

$$\frac{\partial \delta(\mathbf{x}, \tau)}{\partial \tau} + \nabla \cdot \{ [1 + \delta(\mathbf{x}, \tau)] \, \mathbf{u}(\mathbf{x}, \tau) \} = 0,$$
$$\frac{\partial \mathbf{u}(\mathbf{x}, \tau)}{\partial \tau} + \mathcal{H}(\tau) \, \mathbf{u}(\mathbf{x}, \tau) + \mathbf{u}(\mathbf{x}, \tau) \cdot \nabla \mathbf{u}(\mathbf{x}, \tau) =$$

Euler

$$\nabla^2 \Phi(\mathbf{x}, au) = rac{3}{2} \Omega_m(au) \, \mathcal{H}^2(au) \, \delta(\mathbf{x}, au).$$

Poisson

*** During radiation domination we need to include photons

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Initial conditions (inflation + adiabatic)

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Transfer function: Behaviour of the mode by entering the horizon

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Growth rate: late time linear growth

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Please, remind those guys

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Structure formation (LCDM) T(k1.0 1 Gpc 0.8 93 Mpc . $\Phi(x)/\Phi_{in}$ 0.6 0.4 20 Mpc 1.2 0.2 RD MD 0.0 -10-15 -5 0 $\log a$

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You get the transfer function from Boltzmann solvers (CAMB/CLASS)



$$P_L(z,k) = \frac{8\pi^2}{25} A_s \left(\frac{D(z)}{H_0^2 \Omega_m}\right)^2 k T^2(k) \left(\frac{k}{k_\star}\right)^{n_s}$$



The Press-Schechter formalism



The Press-Schechter formalism

Fraction of collapsed objects

$$F_{\text{coll,PS}}(M, z) = 2 \times \frac{1}{\sqrt{2\pi}\sigma(R_L[M], z)} \int_{\delta_{\text{cr}}}^{\infty} d\delta e^{-\delta^2/2\sigma^2(R_L[M], z)}$$



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Delta is a Gaussian variable with variance sigma(R)

$$\sigma^2(R,z) = \langle \delta_R^2(\mathbf{x}) \rangle = \int_{\mathbf{k}} P_L(z,k) |W_R(k)|^2$$

Filter



The Press-Schechter formalism



Fraction of collapsed objects



The Press-Schechter formalism

Fraction of collapsed objects



overdensity $\delta_{R}^{(1)}(q)$

position q







$$\frac{\rho_{\rm NFW}(r)}{\rho_{\rm crit.}} = \frac{\delta_{\rm NFW}}{r/r_s(1+r/r_s)^2} \,,$$

(from v. d. Bosch slides, 2003.06411)





$$\frac{\rho_{\rm NFW}(r)}{\rho_{\rm crit.}} = \frac{\delta_{\rm NFW}}{r/r_s(1+r/r_s)^2} \,,$$

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Problem: observations favor core over cusp profiles (0910.3538 for a review). More later

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(from v. d. Bosch slides, 2003.06411)

I will address in this talk

How those observables change for ULA ...

... within the misalignment mechanism?

+ Mini-halos

... within the LARGE misalignment mechanism?

... within monodromies?



(1510.07633)

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Initial condition for ϕ close to the minimum of its potential



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Two regimes:

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Easy to calculate ULA abundance

$$\rho_a(a) \approx \rho_a(a_{\rm osc})(a_{\rm osc}/a)^3$$

$$\Omega_a \approx \begin{cases} \frac{1}{6} (9\Omega_r)^{3/4} \left(\frac{m_a}{H_0}\right)^{1/2} \left\langle \left(\frac{\phi_i}{M_{pl}}\right)^2 \right\rangle & \text{if } a_{\text{osc}} < a_{\text{eq}} ,\\ \frac{9}{6} \Omega_m \left\langle \left(\frac{\phi_i}{M_{pl}}\right)^2 \right\rangle & \text{if } a_{\text{eq}} < a_{\text{osc}} \lesssim 1 ,\\ \frac{1}{6} \left(\frac{m_a}{H_0}\right)^2 \left\langle \left(\frac{\phi_i}{M_{pl}}\right)^2 \right\rangle & \text{if } a_{\text{osc}} \gtrsim 1 , \end{cases}$$

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Initial condition for ϕ close to the minimum of its potential $H > m_a$ Overdamped 1) Two regimes: Underdamped $H \, < \, m_a$ 2) ~ DE ~ DM Equation of State w0 10^{0} 10^{1} 10^{2} (1510.07633)Scale Factor a/a_i

Easy to calculate ULA abundance $\rho_a(a) \approx \rho_a(a_{\rm osc})(a_{\rm osc}/a)^3$ $\left(\frac{1}{6} (9\Omega_r)^{3/4} \left(\frac{m_a}{H_0} \right)^{1/2} \left\langle \left(\frac{\phi_i}{M_{pl}} \right)^2 \right\rangle \text{ if } a_{\text{osc}} < a_{\text{eq}} ,$ $\Omega_a \approx \begin{cases} \frac{9}{6} \Omega_m \left\langle \left(\frac{\phi_i}{M_{pl}}\right)^2 \right\rangle \text{ if } a_{\text{eq}} < a_{\text{osc}} \lesssim 1 ,\\ \frac{1}{6} \left(\frac{m_a}{H_0}\right)^2 \left\langle \left(\frac{\phi_i}{M_{pl}}\right)^2 \right\rangle \text{ if } a_{\text{osc}} \gtrsim 1 , \end{cases}$ 18 Excluded: Dake Do. 12 + 17 $\log_{10}(\phi_i/\text{GeV})$ $\log_{10}(\Omega_a h^2)$ -3 -5 14 -9 -24 -22 -20 -18 -16 -14 -12 $\log_{10}(m_a/\text{eV})$





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I hope you remind those guys from 15 min ago

$$\delta = \frac{2}{5} \frac{k^2}{H_0^2 \Omega_m} T(k) D(z) \mathcal{R}(k)$$



Initial conditions (inflation)



Transfer function: Behaviour of the mode by entering the horizon



Growth rate: late time linear growth

Misalignment - Initial conditions

(1510.07633)
Remember what **adiabatic** means:



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Feel free to play with **isocurvature** modes (PQ breaking during inflation generates those)

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Feel free to play with **isocurvature** modes (PQ breaking during inflation generates those)

$$S_{ij} = 3(\zeta_i - \zeta_j) \qquad \qquad \zeta_i = -\Psi - H \frac{\delta \rho_i}{\dot{\rho}_i}.$$

** They will be relevant for miniclusters, as we will see

(Maggiore 2, 1510.07633, 1410.2896)

$\begin{aligned} & \mathsf{CDM} \\ \delta_{\mathbf{c}}' + iku_{\mathbf{c}} = -3\Phi', \\ & u_{\mathbf{c}}' + \frac{a'}{a}u_{\mathbf{c}} = -ik\Psi, \end{aligned}$

ULA

$$\begin{split} \dot{\delta}_{a} &= -ku_{a} - (1+w_{a}) \dot{\beta}/2 - 3\mathcal{H} (1-w_{a}) \delta_{a} \\ &-9\mathcal{H}^{2} \left(1 - c_{ad}^{2}\right) u_{a}/k, \\ \dot{u}_{a} &= 2\mathcal{H} u_{a} + k\delta_{a} + 3\mathcal{H} \left(w_{a} - c_{ad}^{2}\right) u_{a}, \end{split}$$

ULA



ULA





ULA





ULA





Misalignment - Growth function

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 $\ddot{\delta}_a + 2H\dot{\delta}_a + (k^2 c_{s,\text{eff}}^2 / a^2) + 4\pi G\rho_a)\delta_a = 0.$



Misalignment - Growth function



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$$\begin{split} & \underset{\begin{subarray}{c} \text{Misalignment - Growth function} \\ & \underset{\begin{subarray}{c} \delta_a \\ \hline \delta_a + 2H\dot{\delta}_a + (k^2c_{s,\text{eff}}^2/a^2) + 4\pi G\rho_a)\delta_a = 0 \,. \end{split}$$

After equating both terms

$$k_J = (16\pi G a \rho_{a,0})^{1/4} m_a^{1/2} = 66.5 a^{1/4} \left(\frac{\Omega_a h^2}{0.12}\right)^{1/4} \left(\frac{m_a}{10^{-22} \text{ eV}}\right)^{1/2} \text{ Mpc}^{-1}.$$

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and for the growth function

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Scale dependent growth

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Misalignment - LSS + CMB Constraints



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(1410.2896)

Misalignment - LSS + CMB Constraints



Factor out time part that oscillates fast (WKB)

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 $|\ddot{\psi}| \ll m |\dot{\psi}| = -2$

 $i \partial_t \psi = -\frac{\nabla^2}{2m} \psi + m \Phi \psi \,.$

$$-\Box\phi + m^2\phi = 0\,,$$

Non-relativistic

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9

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= $m|\psi|^2$ $\vec{v} = \frac{1}{m} \vec{\nabla} \theta$

(2101.11735, 1911.04505)

Continuity + Euler (with extra pressure term)

$$\partial_t \rho + \vec{\nabla} \cdot (\rho \vec{v}) = 0.$$
$$\partial_t \vec{v} + (\vec{v} \cdot \vec{\nabla}) \vec{v} = -\vec{\nabla} \Phi + \frac{1}{2m^2} \vec{\nabla} \left(\frac{\nabla^2 \sqrt{\rho}}{\sqrt{\rho}} \right)$$

Misalignment - Halo Profile



Misalignment - Halo Profile

From numerical simulations





(1407.7762)

Misalignment - Halo Mass function


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 $\delta_{\rm crit}(M,z) = 1.686\mathcal{G}(M,z)$



Misalignment - Halo Mass function

Critical density now depends on scale (halo mass)

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Jeans scale introduces threshold dependence



(1409.3544)

Misalignment: Mini-Halos

(1707.03310)

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$$\begin{split} M_0(m_a,n=0) \approx & 2.3 \times 10^{-7} M_{\odot} \left(\frac{m_a}{10^{-10} \text{ eV}}\right)^{-3/2} \\ & \left(\frac{\Omega_c h^2}{0.12}\right) \left(\frac{\Omega_m}{0.32}\right)^{-3/4} \left(\frac{1+z_{\text{eq}}}{3403}\right)^{3/4} \\ & \left(\frac{1+z_{\text{eq}}}{3403}\right)^{-3/4} \\ & \left(\frac{1+z_{\text{eq}}}{3403}\right)^{-3/4} \\ & \left(\frac{1+z_{\text{eq}}}{3403}\right)^{-1/2} \\ & \left(\frac{1+z_{\text{eq}$$

 $m_a \; [eV]$

(1707.03310)

 10^{-3}

Misalignment: Mini-Halos

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(1707.03310)

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Different Hubble patches with different θ_{PQ}

Topological defects that decay generating axions (complicated stuff)

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$$P_{\theta}^{\rm TH}(k) = \frac{2\pi^4}{K^3} \Theta(k - k) \,. \label{eq:PH}$$

K the typical comoving Hubble at ~ $1~GeV_{\rm c}$

Misalignment: MiniHalos, Miniclusters for QCD axion

As pointed out by Hogan&Rees88, Kolb&Tkachev, there is a mechanism to generate **minihalos** for QCD axion.



We can write down the energy density as

$$\rho(\vec{x}) = \frac{f_{\rm PQ}^2}{2} \left[\dot{\theta}^2 - \frac{1}{a^2} (\vec{\nabla}\theta)^2 + m^2 \left(T\right) \theta^2 \right]$$

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Redshift

K the typical comoving Hubble at ~ $1~GeV_{\rm c}$

Misalignment: MiniHalos, Miniclusters for QCD axion

As pointed out by Hogan&Rees88, Kolb&Tkachev, there is a mechanism to generate **minihalos** for QCD axion.



PQ broken after inflation (f < H_l):

Different Hubble patches with different θ_{PQ}

Topological defects that decay generating axions (complicated stuff)



$$P_{\theta}^{\rm TH}(k) = \frac{2\pi^4}{K^3} \Theta (k - k) \,. \label{eq:P_theta}$$

K the typical comoving Hubble at ~ $1~GeV_{\rm c}$

We can write down the energy density as

$$\rho(\vec{x}) = \frac{f_{\rm PQ}^2}{2} \left[\dot{\theta}^2 - \frac{1}{a^2} (\vec{\nabla}\theta)^2 + m^2 (T) \theta^2 \right]$$

After some developments (Wick's theo, etc.)



Large misalignment

(mostly from 1909.11665)

Large misalignment

 $V = m^2 f^2 \left[1 - \cos\left(\frac{\phi}{f}\right) \right]$

Attractive self-interaction

$$\ddot{\phi} + 3H\dot{\phi} + m^2\phi - \lambda\phi^3 + \dots = 0$$

(mostly from 1909.11665)

Large misalignmentAttractive
self-interaction $V = m^2 f^2 \left[1 - \cos\left(\frac{\phi}{f}\right) \right]$ $\ddot{\phi} + 3H\dot{\phi} + m^2\phi \left(-\lambda\phi^3 + \cdots = 0\right)$ Start to be relevant if

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Attractive

solf_interaction

Pheno consequences:

- delay of oscillations;
- amplification of fluctuations;

Start to be relevant if

$$|\phi^2| \gtrsim m^2/\lambda,$$

(mostly from 1909.11665)

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colf interaction

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$$\frac{\phi(t,\mathbf{x})}{f} = \Theta(t) + \sum_{\mathbf{k}} \theta_{\mathbf{k}}(t) e^{-i\mathbf{k}\cdot\mathbf{x}}$$

Bkg + fluctuations

(mostly from 1909.11665)

-

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~ ` '

 $\phi(t, \mathbf{x})$ $heta_{\mathbf{k}}$ (, $i\mathbf{k} \cdot \mathbf{x}$

 $+\sin(\Theta) = 0$ Dimensionless time (~mt) Θ''

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Bkg + fluctuations

$$\Theta'' + \underbrace{\frac{3}{2t_m}} \Theta' + \sin(\Theta) = 0$$

Dimensionless time (~mt)

Start to be relevant if

Lets analyse the fluctuations

 $|\phi^2| \gtrsim m^2/\lambda,$

(mostly from 1909.11665)

-

$$\theta_{\mathbf{k}}^{\prime\prime} + \frac{3}{2t_m}\theta_{\mathbf{k}}^{\prime} + \left[\cos(\Theta) + \frac{\tilde{k}^2}{t_m}\right]\theta_{\mathbf{k}} = S\left(\tilde{k}, t_m\right)$$

Sourced by curvature perturbations (even under null IC and isocurvature modes)

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Three regimes:

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Three regimes:

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For the energy density

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Transfer Function $|\delta_k/\Phi_0|^2$, $\pi - |\Theta_0| = 10^{-10}$ 100 10² 104 106 1010 108 3 X 0.3 0.1 0.03 -10² 101 10^{3} (1909.11665, 1702.07065) tm

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Mass function

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Which we can solve together to other Boltzmann eq.





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Oscillons

Oscillons Metastable and compact, bound by self-interaction (different than solitons that are bound by gravity)

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Numerical simulations

Initial condition (gaussian)

$$\mathbf{n} \quad \theta(t_{m,0}, \mathbf{x}_m) = \theta_0 \left[1 + \frac{1}{2} \delta_0 \exp\left(-\frac{r_m^2}{2R_{m,0}^2}\right) \right],$$

Initial time can be mapped into different initial mis. angle
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r_m

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Decay: emission of scalar + tensor waves



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GW emission

$$\frac{\Omega_{\rm GW}}{\Omega_{\rm DM}} \simeq 10^{-10} \eta^2 \alpha^5 \left[\frac{10^{-22} \,\mathrm{eV}}{m} \right] \frac{\sqrt{t_{m,\mathrm{coll}}}}{\tilde{k}^2} \left[\frac{\rho_{\pi/2}}{\rho} \right]^{1/3}$$

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GW emission Relation between GW freq and radius of oscillons (~3 by sim)

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Eccentricity (~1 by sim)

r_m

3

Numerical simulations

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Eccentricity (~1 by sim)

$$f_{\rm GW} \simeq 6 \times 10^{-15} \,\mathrm{Hz} \,\alpha \left[\frac{m}{10^{-22} \,\mathrm{eV}}\right]^{1/2} \tilde{k} \sqrt{t_{m,\mathrm{coll}}} \left[\frac{\rho_{\pi/2}}{\rho}\right]^{1/3}$$

(1909.11665)

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(1909.11665)





$$m_a(T=0) = 5.70 \,\mu \text{eV}\left(\frac{10^{12}\,\text{GeV}}{f_a}\right)$$

Temperature corrections from light quarks

$$V(\phi) = -m_{\pi}^2 f_{\pi}^2 \sqrt{1 - \frac{4m_u m_d}{(m_u + m_d)^2} \sin^2\left(\frac{\phi}{2}\right)}$$



 10^{-4} 10^{-3} 10^{-2} 10^{-1} 10^{0}

 $\pi - |\Theta_0|$

10-8 10-7 10-6 10-5

QCD axion

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(1909.11665)



IC fine tuning



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To have a large effect we need theta~pi



People have cook up inflation-based/anthropic-based mechanisms to get theta~pi (1812.11192, 0810.0703)

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*** For more about astrophysical effects (star formation, lensing), see 1909.11665

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$$U(\phi) = \frac{1}{2}m^2\phi^2 + \Lambda^4 \left(1 - \cos\frac{\phi}{f}\right).$$

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Explicitly break shift sym Realized in string theory, see McAllister, Silverstein and Westphal 08

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Aleksandr Chatrchyan +, 1903.03116

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Aleksandr Chatrchyan +, 1903.03116



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Aleksandr Chatrchyan +, 1903.03116



Using Press-Schechter formalism (linear evolution in time), they show it didn't form miniclusters yet Too small, pressure is too high



Depending on axion potential and initial conditions, we have a whole zoo of structures that can form in a broad range of masses

Rich pheno and observational prospects (GW, lensing, star formation, tidal disruption, ...)

Non-linear evolution? SP solver

How baryons affect this picture?



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Such a broad area with many directions to follow. I would be glad to chat more ;)