Status of double virtual NNLO QCD corrections for high multiplicity processes

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Introduction

Era of precision physics at the LHC



• LHC will still be running for years, future colliders

- Many observables probed at percent level precision
- Discovery via precision: search anomalous deviations from Standard Model

Theory must reach comparable precision target!

At least NNLO QCD and NLO EW corrections generally required

 $(\oplus$ parton shower, resummation, etc.)

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[ATL-PHYS-PUB-2022-009]

Why fixed order?



• Stabilization of perturbation series



[Czakon, Mitov, Poncelet '21] [Michal's talk]

Why fixed order?



• Stabilization of perturbation series



[Michal's talk]

• Giant K factors



[Kallweit, VS, Wiesemann '20]

Why fixed order?



Clean theoretical prediction

NNLO QCD: status



Les Houches wishlist 2019 [2003.01700]

Five-point two-loop amplitudes: challenges

Scattering amplitudes for phenomenology



Transcendental part: the canonical way



Transcendental part: the canonical way



Transcendental part: the canonical way



Transcendental part: the five-point massless story



Rational coefficients: algebraic complexity











Overview of recent results

Summary two-loop five-point massless amplitudes

	Comment	Complete analytic results	Public numerical code	Cross sections
$pp \rightarrow jjj$	l.c.	[1]	[1]	[8, 9]
$pp \rightarrow \gamma \gamma j$	l.c.*	[2, 3]	[2]	[10]
$pp \to \gamma\gamma\gamma$	l.c.*	[4, 5]	[4]	[11, 12]
$pp \rightarrow \gamma \gamma j$		[6]		
$gg ightarrow \gamma \gamma g$	NLO loop induced	[7]	[7]	[13]

- [1] [Abreu, Febres Cordero, Ita, Page, <u>VS</u> '21]
- [2] [Agarwal, Buccioni, von Manteuffel, Tancredi '21]
- [3] [Chawdry, Czakon, Mitov, Poncelet '21]
- [4] [Abreu, Page, Pascual, <u>VS</u> '20]
- [5] [Chawdry, Czakon, Mitov, Poncelet '20]
- [6] [Agarwal, Buccioni, von Manteuffel, Tancredi '21]

- [8] [Czakon, Mitov, Poncelet '21]
- [9] [Chen, Gehrmann, Glover, Huss, Marcoli '21]
- [10] [Chawdry, Czakon, Mitov, Poncelet '21]
- 11] [Chawdry, Czakon, Mitov, Poncelet '19]
- [12] [Kallweit, <u>VS</u>, Wiesemann '20]
- [13] [Badger, Gehrmann, Marcoli, Moodie '21]
- [7] [Badger, Brønnum-Hansen, Chicherin, Gehrmann, Hartanto, Henn, Marcoli, Moodie, Peraro, Zoia '21]

One-mass Feynman integrals and functions

One-mass kinematics

e.g. $pp \rightarrow Vjj$



One-mass Feynman integrals and functions





DITTEXP [Moriello '19] [Hidding '20] AMFLow [Liu, Ma, Wang '17] [Liu, Ma '21] Initial values, validation, small scale sampling



One-mass two-loop five-point amplitudes





One-mass two-loop five-point amplitudes



[Badger, Hartanto, Kryś, Zoia '21]

Leading color $(N_c \to \infty)^{\star}$

Setup: Feynman diagrams, FiniteFlow (Laporta IBPs), one-mass pentagon functions

Reconstruction: functional reconstruction [Peraro '16]



sample size 1000k



One-mass pentagon functions

Basis construction

[Chicherin, <u>VS</u>, Zoia '21] (see also [Chicherin, <u>VS</u> '20] [Badger, Hartanto, Zoia '21])



Basis construction

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Planar integral families



Loop order	Master	Master	
Loop order	integrals, $\sigma_{\rm id}$	integrals, S_4	
1	13	56	
2	167	1361	

The alphabet

[Abreu, Ita, Moriello, Page, Tschernow, Zeng '20]

$$A = \sum_{i} \mathrm{d} \log W_i(\mathbf{s}) \ A_i$$

156 letters, 108 contribute up to weight 4

Rational

- 27 linear, $s_{12} p_1^2$
- 54 quadratic, $s_{12}s_{23} + p_1^2s_{45} s_{12}s_{45}$

One square root

• 3+12 letters with
$$\sqrt{\Delta_3^{(i)}}$$
, $\frac{p_1^2 - s_{23} + s_{45} + \sqrt{\Delta_3^{(1)}}}{p_1^2 - s_{23} + s_{45} - \sqrt{\Delta_3^{(1)}}}$

• 1+8 letters with $\sqrt{\Delta_5}$

Square roots:

$$\begin{split} \Delta_5 &= 16\,G(p_1,p_2,p_3,p_4)\\ \Delta_3^{(1)} &= -4\,G\,(p_1,p_2+p_3)\\ \Delta_3^{(2)} &= -4\,G\,(p_1,p_2+p_4)\\ \Delta_3^{(3)} &= -4\,G\,(p_1,p_3+p_4) \end{split}$$

Two square roots

3 letters with
$$\sqrt{\Delta_5}$$
 and one of $\sqrt{\Delta_3^{(i)}}$,
 $\frac{\Omega^- - \Omega^{++}}{\Omega^{++} \Omega^{++-}}$,
 $\Omega^{\pm\pm} = s_{12}s_{15} - s_{12}s_{23} - s_{15}s_{45} \pm s_{34}\sqrt{\Delta_3^{(1)}} \pm \sqrt{\Delta_5}$

Basis features



Weights 1 and 2

Well-defined combinations of $\log,\,\mathrm{Li}_2$ functions



Weights 3 and 4

- Numerical one-fold integration [Caron-Huot, Henn '14] of analytic integrands

 exponential convergence [Takahasi, Mori '73]
- No crossing of physical thresholds \implies no analytic continuation needed
- Dedicated series expansions around spurious singularities

Evaluation of all functions: any one mass planar five-point amplitude in all "crossings"

Sample over physical phase space

Precision	Digits	Timing (s)
double	12	0.19
quadruple	28	159
octuple	60	1695



Available as a C++ library PentagonFunctions++ https://gitlab.com/pentagon-functions/PentagonFunctions-cpp (also Mathematica interface)

Beyond planar

Hexa-box alphabet [Abreu, Ita, Page, Tschernow '21]

204 letters, all appear at weight 4 (+96 from planar)

6 new square roots:

$$\Sigma_5^{(1)} = (s_{12}s_{15} - s_{12}s_{23} - s_{15}s_{45} + s_{34}s_{45} + s_{23}s_{34})^2 - 4s_{23}s_{34}s_{45}(s_{34} - s_{12} - s_{15}) + 5 \text{ permutations}$$

$$\begin{array}{ll} 6+24 \mbox{ letters with } \sqrt{\Sigma_5^{(i)}}, & 6 \mbox{ letters with } \sqrt{\Delta_5} \mbox{ and one of } \sqrt{\Sigma_5^{(i)}}, \\ \frac{s_{12}s_{15}-s_{45}s_{15}-s_{12}s_{23}-s_{23}s_{34}+s_{34}s_{45}-\sqrt{\Sigma_5^{(1)}}}{s_{12}s_{15}-s_{45}s_{15}-s_{12}s_{23}-s_{23}s_{34}+s_{34}s_{45}+\sqrt{\Sigma_5^{(1)}}} & \tilde{\Omega}^{\pm\pm} = p_1^2s_{34} \pm \sqrt{\Delta_5} \pm \sqrt{\Sigma_5^{(1)}} \end{array}$$

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- Nightmare for rationalization-based methods
- A minor inconvenience for our approach: Σ_5 can vanish inside physical phase space \implies new spurious divergences

Conclusions & Outlook

Conclusions & Outlook

Two-loop amplitudes remain major bottleneck in $2 \rightarrow 3$ NNLO QCD computations

- · Leading color for jet and photon production known, full color still in progress
- First results beyond massless scattering
 - · ramp up in complexity, but techniques not stretched to limits
 - planar function basis available, reliable evaluation
 - analytic two-loop amplitudes for Wjj, $Wj\gamma$, $Hb\bar{b}$

Outlook

Need NNLO corrections for Vjj, $Vb\bar{b}$, Hjj, $t\bar{t}j$, $t\bar{t}V$, $t\bar{t}H$, etc.

- Full two-loop five-point one-mass function space
- Top quarks in loops: function bases with non-logarithmic kernels?
- N³LO applications

Looking forward to more exciting NNLO phenomenology!

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Backup

Initial values

We choose an initial point $X_0 \in \mathcal{P}_{45}^+$,

$$X_0 := \left(p_1^2 = 1, s_{12} = 3, s_{23} = 2, s_{34} = -2, s_{45} = 7, s_{15} = -2 \right),$$

which satisfies the following requirements:

- 1. X_0 introduces the minimal number of distinct prime factors.
- 2. X_0 is invariant under the exchanges of momenta $2 \leftrightarrow 3$ and $4 \leftrightarrow 5$ (automorphisms of \mathcal{P}_{45}).
- 3. The four linear letters which have indefinite sign vanish at X_0 .

Algebraic relations between initial					
values required.		Line (⊕ r	ear span	Irred	ucible
Numerical evaluation of available GPL expressions	Weight	Re	Im	Re	Im
[Canko, Papadopoulos, Syrrakos '20],	1	4*	1	4*	1
[Syrrakos 20] to 3000 digits	2	12	4	5	0
	3	67	23	23	7
• Relations from $PSLQ \implies$	4	305	135	90	40
generating set					

We start with canonical DEs

$$d\tilde{g}_{\tau,\sigma}(X) = \epsilon \, d\tilde{A}_{\tau,\sigma}(X) \, \tilde{g}_{\tau,\sigma}(X) \,,$$
$$d\tilde{A}_{\tau,\sigma}(X) = \sum_{i \in \mathbb{A}} a_{\tau,\sigma}^{(i)} \, d\log W_i(X) \,.$$

for all relevant integral families τ and all permutations σ .

E.g. for planar five-point one-mass $\tau = \{1L, zmz, zzm, zzz, 1L^2\}, \sigma \in S_4$.

Step 2: solution through iterated integrals

- 1. Choose a "good" base point X_0 for all τ and σ inside a physical phase space $\mathcal{P}_{ij},$ $ij \to kl \ldots$
- 2. Determine initial values $ec{g}^{(w)}_{ au,\sigma}(X_0)$ and algebraic relations between them
- 3. Write solutions DEs along a generic path $\gamma(t)$ ($\gamma(0) = X_0$) entirely within \mathcal{P}_{ij} through iterated integrals

$$\vec{g}_{\tau,\sigma}^{(w)} = \sum_{w'=0}^{w} \sum_{i_1,\dots,i_{w'} \in \mathbb{A}} a_{\tau,\sigma}^{(i_1)} \cdot a_{\tau,\sigma}^{(i_2)} \cdots a_{\tau,\sigma'}^{(i_{w'})} \cdot \vec{g}_{\tau,\sigma}^{(w-w')}(X_0) \left[W_{i_1},\dots,W_{i_{w'}} \right]_{X_0}$$

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- $\bigcup_{\tau,\sigma} \vec{g}_{\tau,\sigma}^{(w)}$ spans a vector space at each weight ω
- No analytic continuation required by construction

Iterated integrals

[Chen '77] (see also "Iterated integrals in QFT" [Brown '11])

Let $\omega_1, \ldots, \omega_n$ be differential 1-forms on M (phase space), and path $\gamma : [0, 1] \to M$. Pull the forms back on the path $\omega_i(\mathbf{s}) \xrightarrow{\gamma^*} w_i(t) dt$. Iterated integrals $\langle \text{iints} \rangle$ are

$$M_{s_0} = \int_0^1 w_n(t_n) \, \mathrm{d}t_n \dots \int_0^{t_2} w_1(t_1) \, \mathrm{d}t_1 \tag{ii}$$

We need only logarithmic forms $\omega_i = d \log(W_i)$, use notation

$$[W_1,\ldots,W_n]_{\gamma} \coloneqq I_{\gamma}[\omega_1,\ldots,\omega_n]$$

Shuffle product

$$I_{\gamma}[\omega_1,\ldots,\omega_r] \ I_{\gamma}[\omega_{r+1},\ldots,\omega_n] = \sum_{i \in \{1,\ldots,r\} \sqcup \{r+1,n\}} I_{\gamma}[\omega_{i_1},\ldots,\omega_{i_n}] \tag{U}$$

- All functional relations manifest
- Similar to symbol, but complete information preserved

Step 3: construct basis

Proceed recursively order-by-order in ϵ (or transcendental weight w).

Denote weight w basis as $\{f_i^{(w)}\}$.

Start from w = 0: integrals are rational numbers $\iff \{f_i^{(0)}\} \coloneqq \{1\}.$

- Consider subspace $\mathbf{G}^{(w)} \coloneqq \bigcup_{\tau,\sigma} \vec{g}_{\tau,\sigma}^{(w)}$ in vector space spanned by all iints in $\mathbf{G}^{(w)}$
- Add all products of lower-weight functions $\{f_i^{(w' < w)}\}$ to $\mathbf{G}^{(w)},$

$$\mathbf{F}^{(w)} \coloneqq \mathbf{G}^{(w)} \bigcup_{\vec{w}, \vec{i}} f_{\vec{i}}^{(\vec{w})}, \qquad f_{\vec{i}}^{(\vec{w})} = f_{i_1}^{(w_1)} \cdots f_{i_2}^{(w_n)}, \qquad \sum_k w_k = w$$

Use shuffle algebra of iints to linearize identities in $\mathbf{F}^{(w)}$.

Use linear algebra to choose a basis in F^(w), preferring products of lower-weight functions. This basis is functions {f_i^(w)}.

Explicit real-analytic \log, Li_2, Cl_2 functions, e.g.

$$g_{1,4}^{(1)} \equiv [W_4] = \log(s_{45})$$

$$g_{2,10}^{(1)} = [W_{31}] + \frac{i\pi}{2} + \frac{\log(3)}{2} = \log(\sqrt{\Delta})$$

$$g_{1,3}^{(2)} \equiv -[W_2, W_2] + [W_2, W_{14}] + [W_4, W_2] - [W_4, W_{14}] + ([W_2] - [W_4])\log(2) + \frac{\pi^2}{12}$$

$$= -\operatorname{Li}_2\left(\frac{s_{45}}{s_{23}}\right) - \log\left(-\frac{s_{45}}{s_{23}}\right)\log\left(1 - \frac{s_{45}}{s_{23}}\right)$$

Positivity properties of the alphabet are important!

Weight 3 functions are one-fold integrals of weight 2 functions by definition

$$f_i^{(3)}(X) = \sum_{j,k} c_{i,j,k} \int_0^1 \mathrm{d}\log W_j(t) \, h_k^{(2)}(t) + \tau_i^{(3)}$$

Integrands analytic on the integration domain \implies integration well-defined.

- Efficient numerical integration possible
- Some care to avoid numerical cancellations if $d \log(W_j)$ can vanish along the path

Weight 4: one-fold integral representation

Change order of integration [Caron-Huot, Henn '14]

$$I_{\gamma}[\omega_{1}, \dots, \omega_{n}] = \int_{0}^{1} (\gamma^{\star} \circ \omega_{n})(t) \int_{0}^{t} (\gamma^{\star} \circ \omega_{n-1})(u) I_{\gamma(u)}[\omega_{1}, \dots, \omega_{n-2}] = \int_{0}^{1} (\gamma^{\star} \circ \omega_{n-1})(u) \left(\int_{u}^{1} (\gamma^{\star} \circ \omega_{n})(t) \right) I_{\gamma(u)}[\omega_{1}, \dots, \omega_{n-2}] \qquad 0 \qquad 1 \qquad u$$

For logarithmic forms the last integration is trivial

$$\int_{u}^{1} (\gamma^{\star} \circ \omega_n)(t) = \int_{u}^{1} \mathrm{d}\log\left(W_n(t)\right) = \log(W_n(1)) - \log(W_n(u))$$

Integrands are analytic, except possibly at loci of vanishing spurious letters, where integrable logarithmic divergences can occur.

- Linear letters: only at X₀
- Quadratic: anywhere along the path \implies efficient numerical integration over multiple line segments