

# Top quark contribution to two-loop helicity amplitudes for W/Z boson pair production in gluon fusion

Christian Brønnum-Hansen | 26th of April | Loops and Legs 2022 Based on work with Chen-Yu Wang







#### Outline

#### 1. Motivation

#### 2. Two-loop $gg \rightarrow WW/ZZ$ with full $m_t$ dependence

- Numerical IBP reduction
- Master integrals evaluated with the Auxiliary Mass Flow method

#### 3. Conclusion

Motivation 00000 

#### Motivation

- Higher experimental precision at the LHC requires more accurate theoretical predictions.
  - $\rightarrow$  more loops
- As  $\sqrt{s}$  increases, massive particles in the loops become more important.
  - $\rightarrow$  more masses
- Precision wish list Amoroso et al. 2020

process	known	desired
:	:	:
$pp  ightarrow H + t\overline{t}$	$NLO_{QCD} + NLO_{EW}$	NNLO <sub>QCD</sub>
:	:	:
$nn \rightarrow 1/1//$	$NNLO_{QCD} + NNLO_{EW}$	NNLO <sub>QCD</sub>
$pp \rightarrow vv$	$+ \text{ NLO}_{\text{QCD}} (gg \text{ channel})$	(gg channel with massive loops)
:	:	:
	$NNLO_{QCD} + NLO_{EW}$	
$pp  ightarrow t \overline{t}$	NLO <sub>QCD</sub> (with decays)	NNLO <sub>QCD</sub> (with decays)
	NLO <sub>EW</sub> (with decays)	Czakon, Mitov, et al. 2021
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Two-loop  $gg \rightarrow WW/ZZ$  with full  $m_t$  dependence

#### **Motivation**

• background to  $pp \rightarrow H^* \rightarrow WW/ZZ$ , interference  $\rightarrow$  Higgs width Campbell et al. 2011a; Kauer and Passarino 2012; Caola and

Melnikov 2013; Campbell et al. 2014; Azatov et al. 2015

- anomalous gauge couplings
- gg channel: loop-induced: enters σ<sub>ppVV</sub> at NNLO, enhanced by gluon flux & event selection Binoth et al. 2006; Campbell et al.
   2011b
- massless NLO contribution: Caola, Melnikov, Röntsch, et al. 2015, 2016  $\geq 50\% \text{ to } \sigma_{ggZZ}, 50\% \text{ to } \sigma_{ggWW}$   $6 - 8\% \text{ to } \sigma_{ppZZ}, 2\% \text{ to } \sigma_{ppWW}$
- 3rd generation increases massless LO  $\sigma_{ggWW}$  by 10 13%  $_{Binoth\ et\ al.\ 2006;\ Campbell\ et\ al.\ 2011a}$
- dominant contribution for high p<sub>T</sub> Campbell et al. 2011a
- Motivation ○●○○○











#### $gg \rightarrow WW/ZZ$ : Progress

- LO: one-loop Glover and Bij 1989; Kao and Dicus 1991; Duhrssen et al. 2005; Binoth et al. 2006
- NLO real: one-loop Agrawal and Shivaji 2012; Melia et al. 2012; Campanario et al. 2013
- NLO virtual: two-loop massless: gg 
  ightarrow VV Caola, Henn, et al. 2015; Manteuffel and Tancredi 2015
- NLO virtual: two-loop massive:
  - region expansions: gg 
    ightarrow ZZ Melnikov and Dowling 2015; Gröber et al. 2019; Davies, Mishima, Steinhauser, and Wellmann 2020
  - full  $m_t$  dependence: gg 
    ightarrow ZZ Agarwal et al. 2020; Brønnum-Hansen and Wang 2021b
  - full  $m_t$  dependence: gg 
    ightarrow WW Brønnum-Hansen and Wang 2021a

Two-loop  $gg \rightarrow WW/ZZ$  with full  $m_t$  dependence



#### Analytic or alternative

#### Analytic

- deeper understanding (e.g. Parke-Taylor)
- fast, precise evaluation (e.g.  $\text{Li}_n$ ,  $_pF_q$ ,  $K(\lambda)$ , MPL, ...)
- wider applications (e.g. changing parameters)

#### Alternatives

- region expansions:  $gg \rightarrow ZZ$  Melnikov and Dowling 2015; Davies, Mishima, Steinhauser, and Wellmann 2020,  $gg \rightarrow ZH$  Davies, Mishima, and Steinhauser 2020,  $gg \rightarrow HH$  Davies, Herren, et al. 2022 [Joshua Davies]
- numerical evaluation:
  - sector decomposition: gg 
    ightarrow HH Borowka, Greiner, et al. 2016, gg 
    ightarrow ZZ Agarwal et al. 2020, gg 
    ightarrow ZH Chen, Heinrich, et al. 2020
  - differential equations:  $gg \to t\bar{t}$  Chen, Czakon, et al. 2018,  $gg \to H$  Czakon and Niggetiedt 2020 [Marco Niggetiedt]

• combined:  $gg \rightarrow HH$  Davies, Heinrich, et al. 2019,  $gg \rightarrow ZH$  Chen, Davies, et al. 2022 [Stephen Jones]

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### Analytic or numeric?

#### Analytic

- deeper structure (e.g. Parke-Taylor)
- fast, precise evaluation (e.g.  $Li_n$ ,  ${}_{p}F_{q}$ ,  $K(\lambda)$ , MPL, ...)
- wider applications (e.g. changing parameters)

#### Alternatives

fast numerical method that produces **arbitrary precision** result good **interpolation** algorithm (or fixed grid [Jérémie Quarroz])  $\Rightarrow$  fast, precise evaluation in the parameter space

- Numeric IBP reduction: set masses and kinematic variables to rational numbers
- Numeric DE evaluation: simple boundary condition + arbitrary precision Lee et al. 2018; Liu, Ma, and Wang 2018; Abreu et al.

2020; Hidding 2020; Moriello 2020

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### $gg \rightarrow VV$ : Diagrams



Two-loop  $gg \rightarrow WW/ZZ$  with full  $m_t$  dependence

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### $gg \rightarrow VV$ : Diagrams



Two-loop  $gg \rightarrow WW/ZZ$  with full  $m_t$  dependence

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### $gg \rightarrow VV$ : Projection





- $\gamma_5$  scheme:
  - *WW* and  $\triangle^2$  diagrams: Larin scheme Larin and Vermaseren 1991; Moch et al. 2015.
  - two-loop ZZ: naive scheme, thanks to Furry's theorem.
- Projection onto 36 tensor structures Binoth et al. 2006



Two-loop $gg \rightarrow WW/ZZ$	with full n	nt dependence
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### $gg \rightarrow VV$ : IBP reduction

Find symmetry relations with REDUZE 2 Manteuffel and Studerus 2012.





Reduction performed at each phase space point with KIRA 2.2: Klappert et al. 2020; Lange et al. 2021:

- parametric only in d
  - $m_t = 173 \text{ GeV}, \ m_W = 80 \text{ GeV}, \ m_Z = 91 \text{ GeV}$
  - rational values for s, t

avoid non-factorisable denominators Smirnov and Smirnov 2020; Usovitsch 2020

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### $gg \rightarrow VV$ : IBP reduction



Iower the memory/storage consumption significantly

straightforward parallelisation

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## Based on the auxiliary mass flow method Liu, Ma, and Wang 2018; Liu, Ma, Tao, et al. 2020; Liu and Ma 2022

 $I \propto \lim_{\eta \to 0^+} \int \prod_{i=1}^2 \mathrm{d}^d k_i \prod_{a=1}^9 \frac{1}{[q_a^2 - (m_a^2 - i\eta)]^{\nu_a}}$ 

• Add an imaginary part to the top quark mass

$$m_t^2 \rightarrow m_t^2 - i\eta$$

Solve differential equations w.r.t the mass

$$\partial_x I = MI, \quad x \propto -i\eta$$

with boundary condition at  $x \to -i\infty$ . Physical mass at  $x \to 0$ .

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$$gg \rightarrow VV$$
: Master integral evaluation



Motivation





### $gg \rightarrow VV$ : Master integral evaluation: boundary



Boundary example: double box with one internal massive line

• Collect and reduce boundary integrals for all masters

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### $gg \rightarrow VV$ : Master integral evaluation: boundary



#### • All boundary integrals known analytically

't Hooft and Veltman 1979; Chetyrkin et al. 1980; Scharf and Tausk 1994; Gehrmann and Remiddi 2000; Gehrmann, Huber, et al. 2005

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#### $gg \rightarrow VV$ : Master integral evaluation



• Expand *I* around **boundary** in variable  $y = x^{-1} = 0$ :

$$\boldsymbol{I} = \sum_{j}^{M} \epsilon^{j} \sum_{k}^{N} \sum_{l} \boldsymbol{c}_{jkl} y^{k} \ln^{l} y + \dots$$

• Evaluate and expand around regular points:

$$\boldsymbol{I} = \sum_{j}^{M} \epsilon^{j} \sum_{k=0}^{N} \boldsymbol{c}_{jk} x^{\prime k} + \dots$$

- Evaluate at the physical point.  $x = 0 \leftarrow$  regular point
- Path is fixed by singularities and desired precision.



 $m_t^2 
ightarrow m_t^2(1+x)$ 

Conclusion

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### $gg \rightarrow VV$ : Master integral evaluation





• Steps for a double-box integral with five massive propagators

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lackstyle Numerical evaluation of integrals point by point,  $\sim 1$  CPU hour per phase space point

$$\begin{tabular}{|c|c|c|c|c|} $gg \rightarrow WW & gg \rightarrow ZZ \\ \hline $accurate digits $ $\sim 15 digits $ $\sim 20 digits $ \end{tabular} \end{tabular}$$

- Expected relative error is (<sup>△</sup>/<sub>R</sub>)<sup>N</sup> ⇒ arbitrary precision attainable for step size △ respective to convergence radius *R* and expansion order *N*
- Self-consistency check with DEs in s and t
- Cross-checks with pySecDec Borowka, Heinrich, Jahn, Jones, Kerner, Schlenk, and Zirke 2018; Borowka, Heinrich, Jahn, Jones, Kerner, and Schlenk 2019

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#### $gg \rightarrow VV$ : Helicity amplitudes





• Phase space parametrised in terms of the relative velocity  $\beta$  and scattering angle  $\theta$ 

$$s=rac{4m_W^2}{1-eta^2}, \qquad t=m_W^2-rac{s}{2}\left(1-eta\cos heta
ight).$$

Massive boson polarisation vectors written in terms of decay currents

$$\epsilon^{*\mu}_{3,L} = \langle 5|\gamma^{\mu}|6], \qquad \epsilon^{*\mu}_{4,L} = \langle 7|\gamma^{\mu}|8]$$

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gg 
ightarrow WW

•  $\sqrt{s} \approx 367 \text{ GeV}$ 

•  $\theta \approx 36.9^{\circ}$ 

•  $\mu = m_W$ 

CA		$\epsilon^{-2}$	$\epsilon^{-1}$	
	$A^{(2)}/A^{(1)}$	$1.0000000023 - 3.0 \cdot 10^{-11}i$	-4.94945452453 + 3.89380807761 <i>i</i>	
LLLL	IR pole	1.0000000000	-4.94945452593 + 3.89380807754i	
	$A^{(2)}/A^{(1)}$	$0.99999999815 - 1.6 \cdot 10^{-9}i$	-4.29712348534 + 9.47440879823 <i>i</i>	
LNLL	IR pole	1.0000000000	-4.29712347965+9.47440880940 i	

 $gg \rightarrow ZZ$ 

•  $\sqrt{s} \approx 210 \text{ GeV}$ 

 $\bullet \ \theta \approx 114^\circ$ 

•  $\mu = m_Z$ 

C <sub>A</sub>		$\epsilon^{-2}$	$\epsilon^{-1}$	
	$A^{(2)}/A^{(1)}$	$1.000000000008 - 7.6 \cdot 10^{-13}i$	0.8304916142577 + 3.229874368770	
	IR pole	1.000000000000	0.8304916142539 + 3.229874368771 i	
	$A^{(2)}/A^{(1)}$	$1.000000000009 - 1.4 \cdot 10^{-12}i$	0.2359507533 <mark>599</mark> + 2.885154863850 <i>i</i>	
LKLL	IR pole	1.000000000000	0.2359507533772 + 2.885154863852 <i>i</i>	

$$\mathbf{I}^{(1)}(\epsilon,\mu)A^{(1)}(\epsilon,\mu) = \frac{e^{\epsilon\gamma_E}}{\Gamma(1-\epsilon)} \left(\frac{C_A}{\epsilon^2} + \frac{\frac{11}{6}C_A}{\epsilon}\right) \left(\frac{\mu^2 e^{i\pi}}{s}\right)^{\epsilon} A^{(1)}(\epsilon,\mu)$$

gg 
ightarrow WW

•  $\sqrt{s} \approx 367 \text{ GeV}$ 

•  $\sqrt{s} \approx 210 \text{ GeV}$ 

 $\theta \approx 114^{\circ}$ 

 $\mu = m_7$ 

•  $\theta \approx 36.9^{\circ}$ 

 $\mu = m_W$ 

CA		$\epsilon^{-2}$	$\epsilon^{-1}$	
	$A^{(2)}/A^{(1)}$	$1.0000000023 - 3.0 \cdot 10^{-11}i$	-4.94945452453 + 3.89380807761 <i>i</i>	
LLLL	IR pole	1.0000000000	-4.94945452593 + 3.89380807754 <i>i</i>	
	$A^{(2)}/A^{(1)}$	$0.99999999815 - 1.6 \cdot 10^{-9}i$	-4.29712348534 + 9.47440879823 <i>i</i>	
LNLL	IR pole	1.0000000000	-4.29712347965 + 9.47440880940 <i>i</i>	

 $gg \rightarrow ZZ$ 





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### Conclusion

- We calculated two-loop helicity amplitudes for  $gg \rightarrow WW/ZZ$  with full  $m_t$  dependence.
- Numeric IBP reduction is efficient in practice for complicated multi-scale processes.
- Auxiliary mass flow method provides an efficient and precise way to evaluate multi-loop integrals.
- Future work: Impact on the cross section as well as interference with Higgs.

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Thank you for your attention!

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Figure 10: The ratio of NLO to LO (upper) and the percentage of the NLO cross section from the gg initial state (lower) for  $WW \rightarrow e^+\mu^-\nu_e \overline{\nu}_\mu$  production, as a function of the jet-veto  $p_T^{veto}$ . Results are shown using the basic cuts of Eq. (7.6) (upper, blue curves) and the Higgs search cuts of Eq. (7.7) (lower, red curves). The NLO to LO ratio and gluon percentage with no veto applied Backup econcurves.

 $\overline{q}_1 q_2 W$  vertex introduces axial current

$$i\overline{q}_1\gamma_\mu \frac{1-\gamma_5}{2}q_2W^\mu.$$
 (1)

 $\gamma_5$  in *d*-dimensions through Larin-scheme Larin and Vermaseren 1991; Moch et al. 2015

$$\gamma_{\mu}\gamma_{5} = -\frac{1}{3!}\varepsilon_{\mu\nu\rho\sigma}\gamma^{\nu}\gamma^{\rho}\gamma^{\sigma}$$

$$\varepsilon^{\mu\nu\rho\sigma}\varepsilon_{\alpha\beta\kappa\lambda} = \begin{vmatrix} \delta^{\mu}_{\alpha} & \delta^{\mu}_{\beta} & \delta^{\mu}_{\kappa} & \delta^{\mu}_{\lambda} \\ \delta^{\nu}_{\alpha} & \delta^{\nu}_{\beta} & \delta^{\nu}_{\kappa} & \delta^{\nu}_{\lambda} \\ \delta^{\rho}_{\alpha} & \delta^{\rho}_{\beta} & \delta^{\rho}_{\kappa} & \delta^{\rho}_{\lambda} \\ \delta^{\sigma}_{\alpha} & \delta^{\sigma}_{\beta} & \delta^{\sigma}_{\kappa} & \delta^{\sigma}_{\lambda} \end{vmatrix}$$

$$(3)$$

finite renormalisation to restore axial Ward identities Larin 1993

$$J_{\mu}^{A} = Z_{5} J_{\mu, b}^{A} = \left[ 1 - \frac{\alpha_{s}}{2\pi} 2C_{F} + \mathcal{O}(\alpha_{s}^{2}) \right] J_{\mu, b}^{A}$$
(4)

Backup 00000 projection onto 38 tensor structures  $T_I$  and  $S_I$  Binoth et al. 2006; Manteuffel and Tancredi 2015

$$A(\{p_i\}, \{\epsilon_j\}, m_t) = \sum_{l=1}^{20} A_l(s, t, m_W, m_t) T_l(\{p_i\}, \{\epsilon_j\}) + \sum_{l=21}^{38} A_l(s, t, m_W, m_t) S_l(\{p_i\}, \{\epsilon_j\}).$$

define projectors to obtain form factors  $A_I$ 

e

$$P_{J} = \sum_{l=1}^{20} a_{Jl} T_{l}^{*}(\{p_{i}\}, \{\epsilon_{j}\}) + \sum_{l=21}^{38} b_{Jl} S_{l}^{*}(\{p_{i}\}, \{\epsilon_{j}\})$$

$$A_{I} = \sum_{\text{pol}} P_{I} A(\{p_{i}\}, \{\epsilon_{j}\}, m_{t})$$
(6)
$$\text{.g.} \sum_{\text{pol}} \epsilon_{1}^{\mu} \epsilon_{1}^{\nu*} = -\eta^{\mu\nu} + \frac{p_{1}^{\mu} p_{2}^{\nu} + p_{1}^{\nu} p_{2}^{\mu}}{p_{1} \cdot p_{2}}$$

Backup 00●00 (5)

expected relative error:  $\left(\frac{\Delta}{R}\right)^N$ 

- *N*, order of expansion
- $\Delta/R$  size of step  $\Delta$  in terms of radius of convergence R

table of correct digits compared to analytic result:

$\Delta/R, \# - N$	4	8	16	32
$1, \ 1$	0	1	1	3
$\frac{1}{2}, 3$	2	3	5	11
$\frac{1}{4}, 9$	3	6	10	20
$\frac{1}{8}$ , 23	4	8	17	31
$\frac{1}{16}, 58$	5	11	20	41
$\frac{1}{32}$ , 139	7	13	25	49

# number of intermediate steps required

Backup 000●0



### $gg \rightarrow ZZ$ : comparison to series expansion results



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