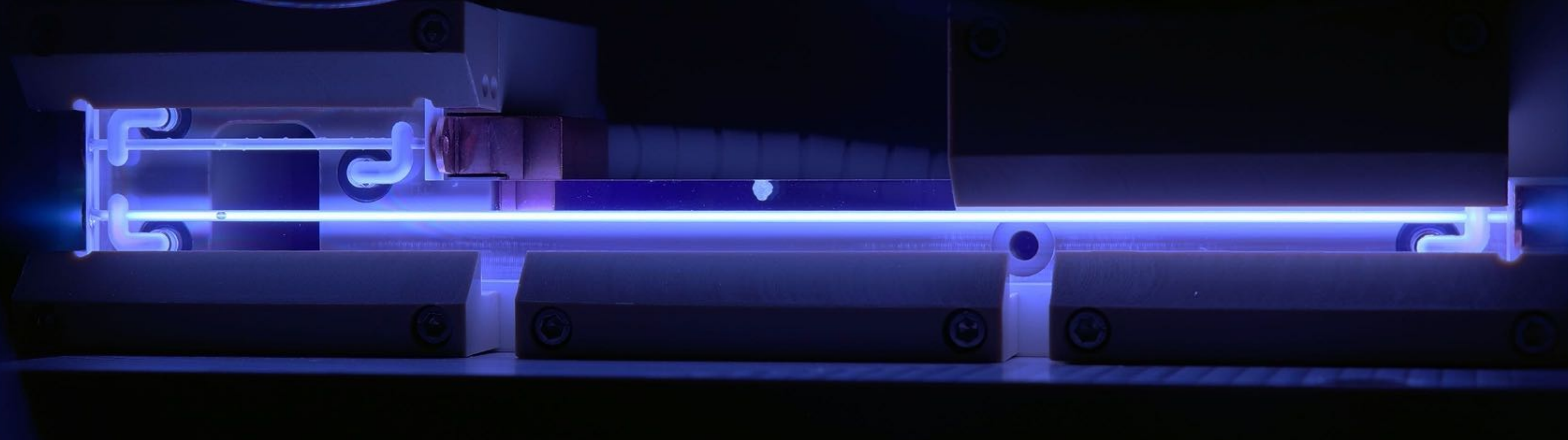


Plasma Wakefield Acceleration

An introduction to laser- and beam-driven concepts



Jens Osterhoff

Head of Plasma Accelerator R&D

DESY. Accelerator Division

DESY Summer School

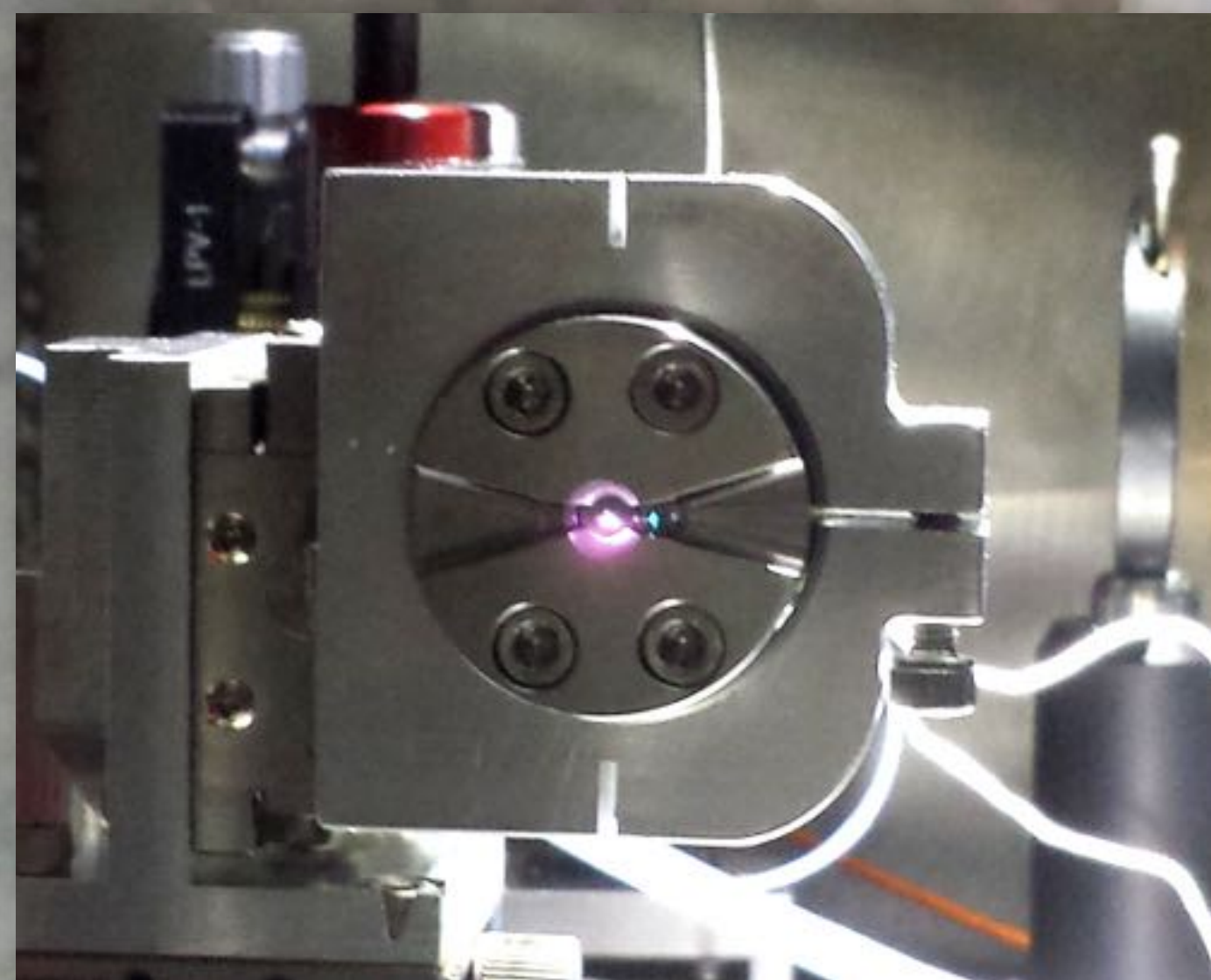
August 19th, 2021

HELMHOLTZ RESEARCH FOR
GRAND CHALLENGES

MIT | **ARD**

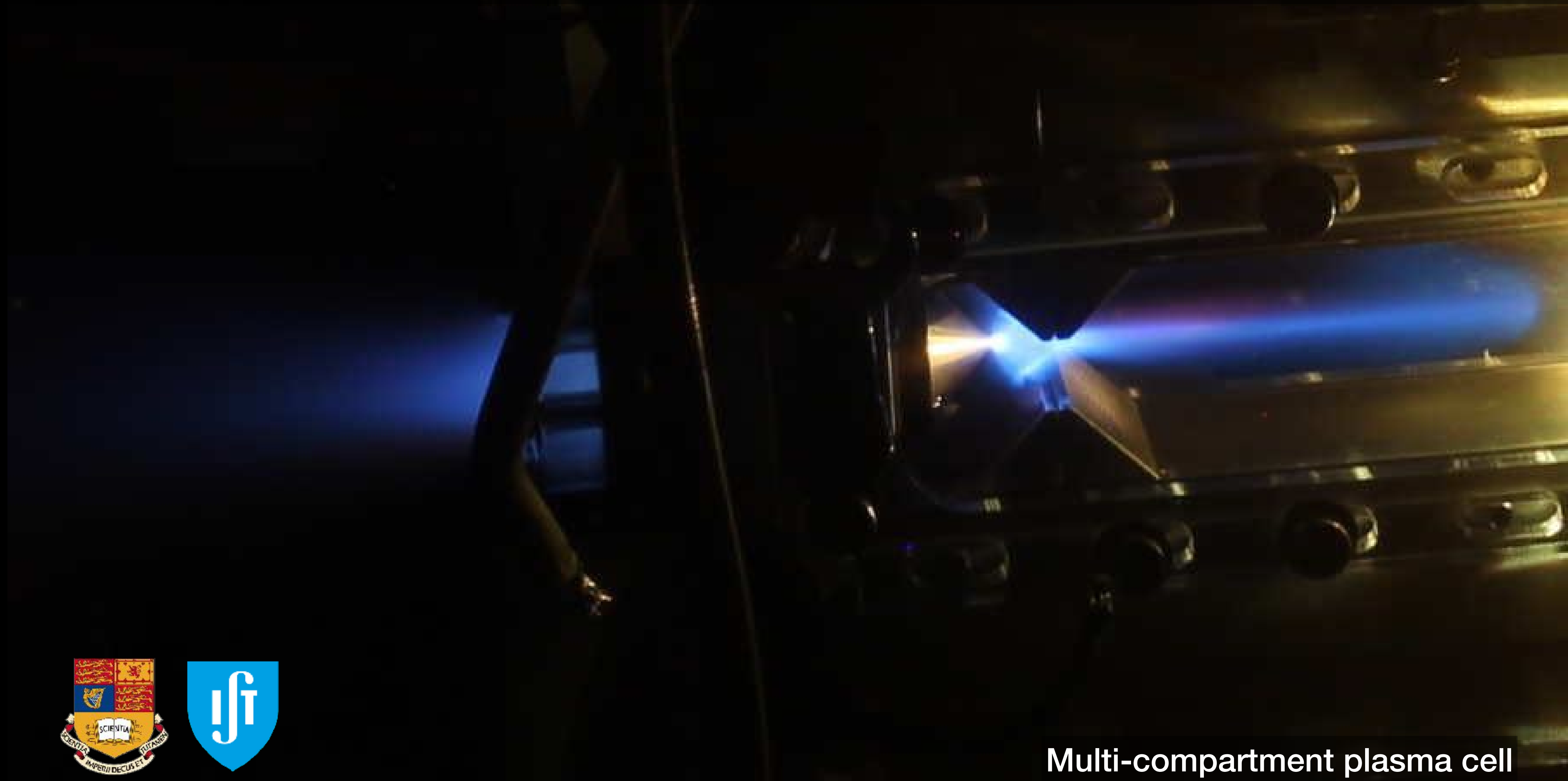


Plasmas for accelerators come in many flavors...



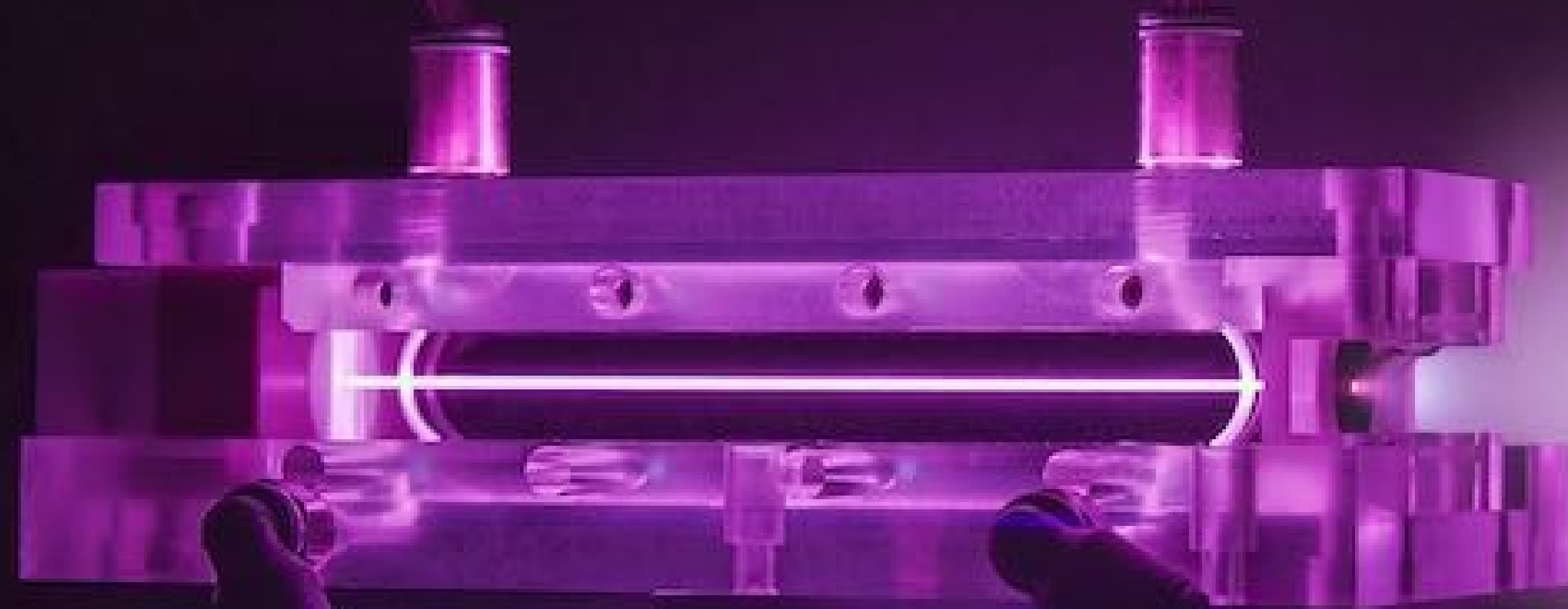
Gas jet

Plasmas for accelerators come in many flavors...



Multi-compartment plasma cell

Plasmas for accelerators come in many flavors...



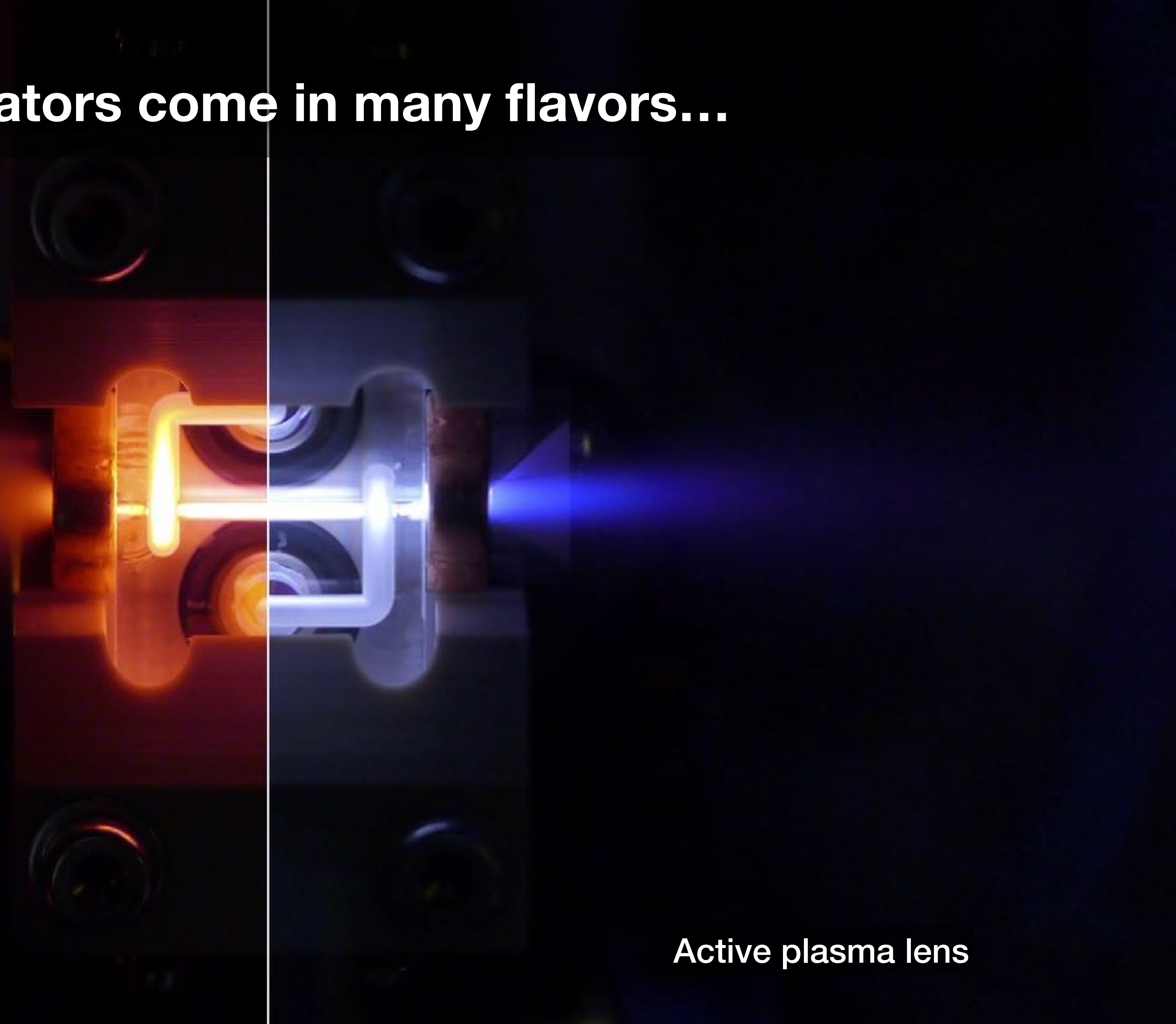
Capillary discharge waveguide

Plasmas for accelerators come in many flavors...



Alkali vapor oven

Plasmas for accelerators come in many flavors...



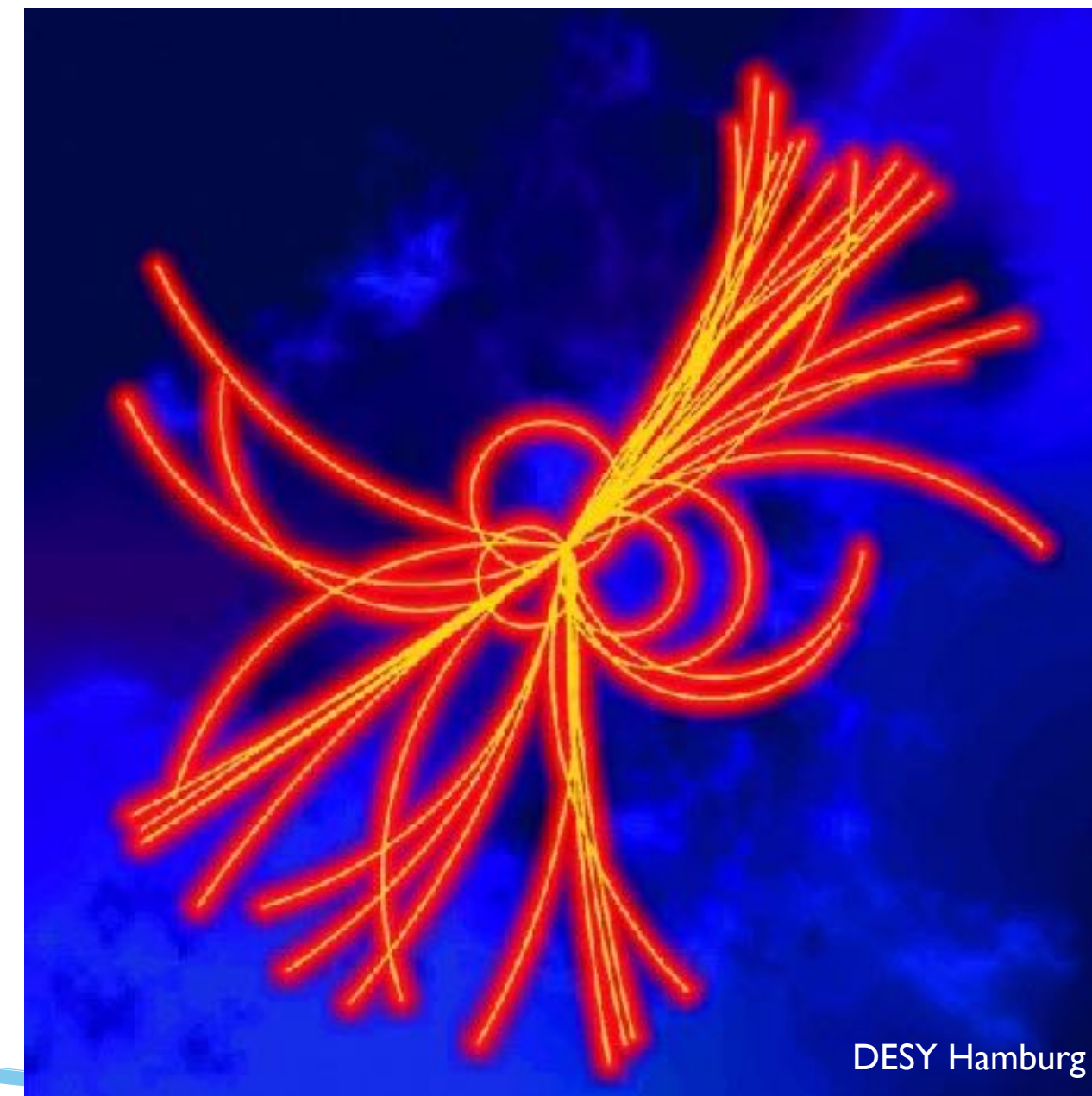
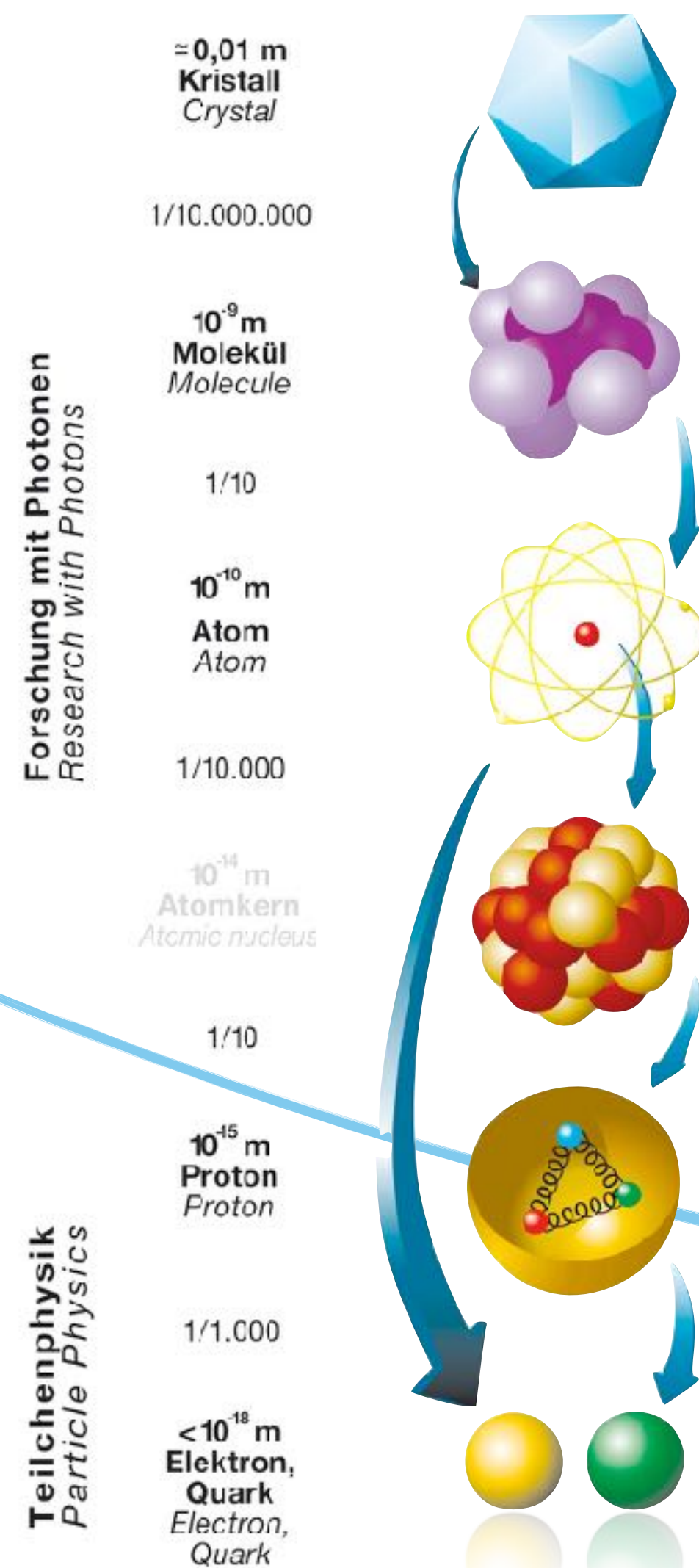
Active plasma lens

Outline

- > Introduction to laser-driven plasma wakefield accelerators: why do we care?
- > Properties of plasma wakefields
- > Example of current research: FLASHForward at DESY

Accelerators are the heart of high-energy photon sources & colliders

Cutting-edge, high-end slow-motion-cameras and microscopes to study the structure of matter



DESY Hamburg

Simulation of the decay of a Higgs Boson (LHC, CERN)

Particle colliders

investigation of the fundamental forces and constituents of matter



Accelerators are the heart of high-energy photon sources & colliders

Cutting-edge, high-end slow-motion-cameras and microscopes to study the structure of matter

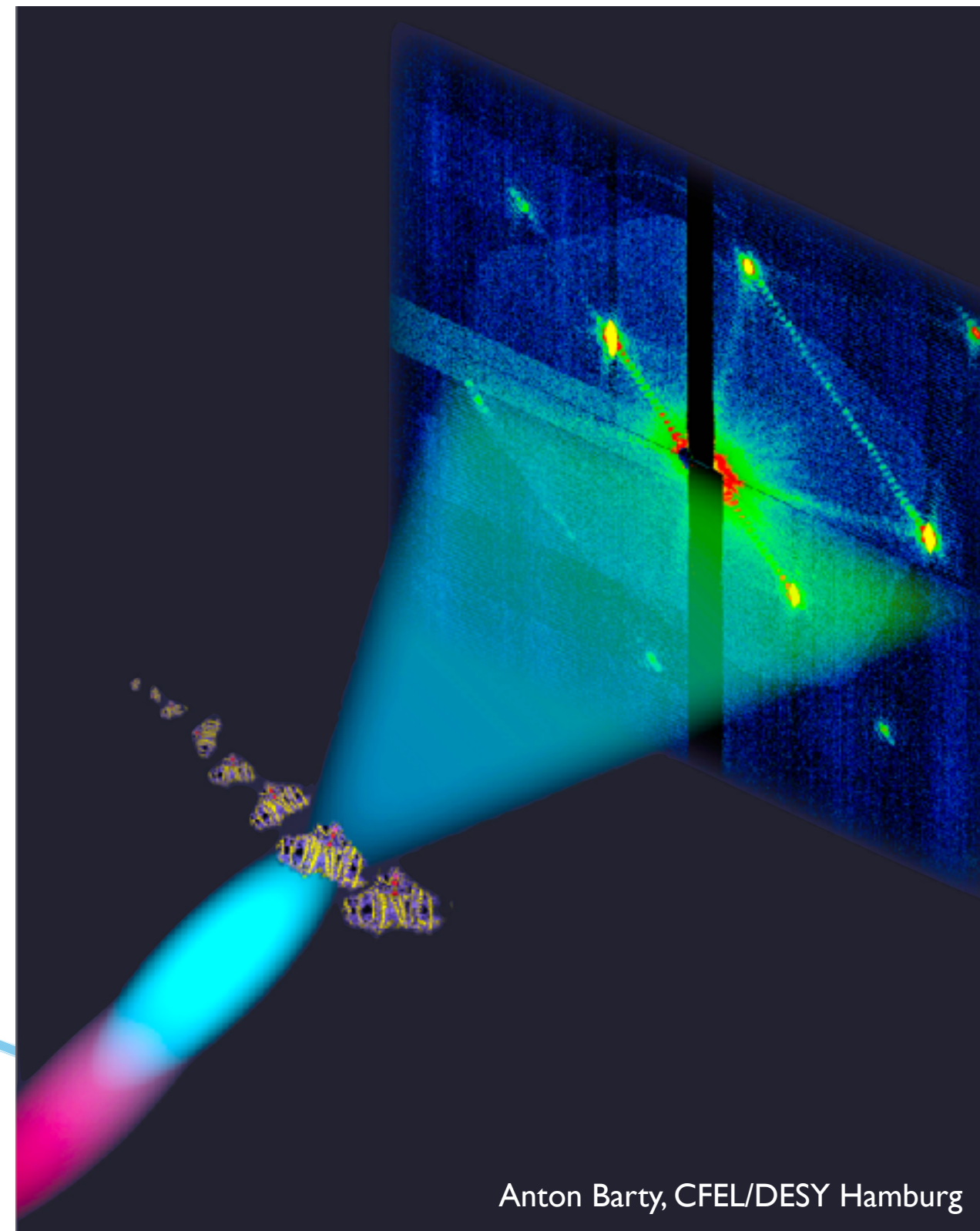
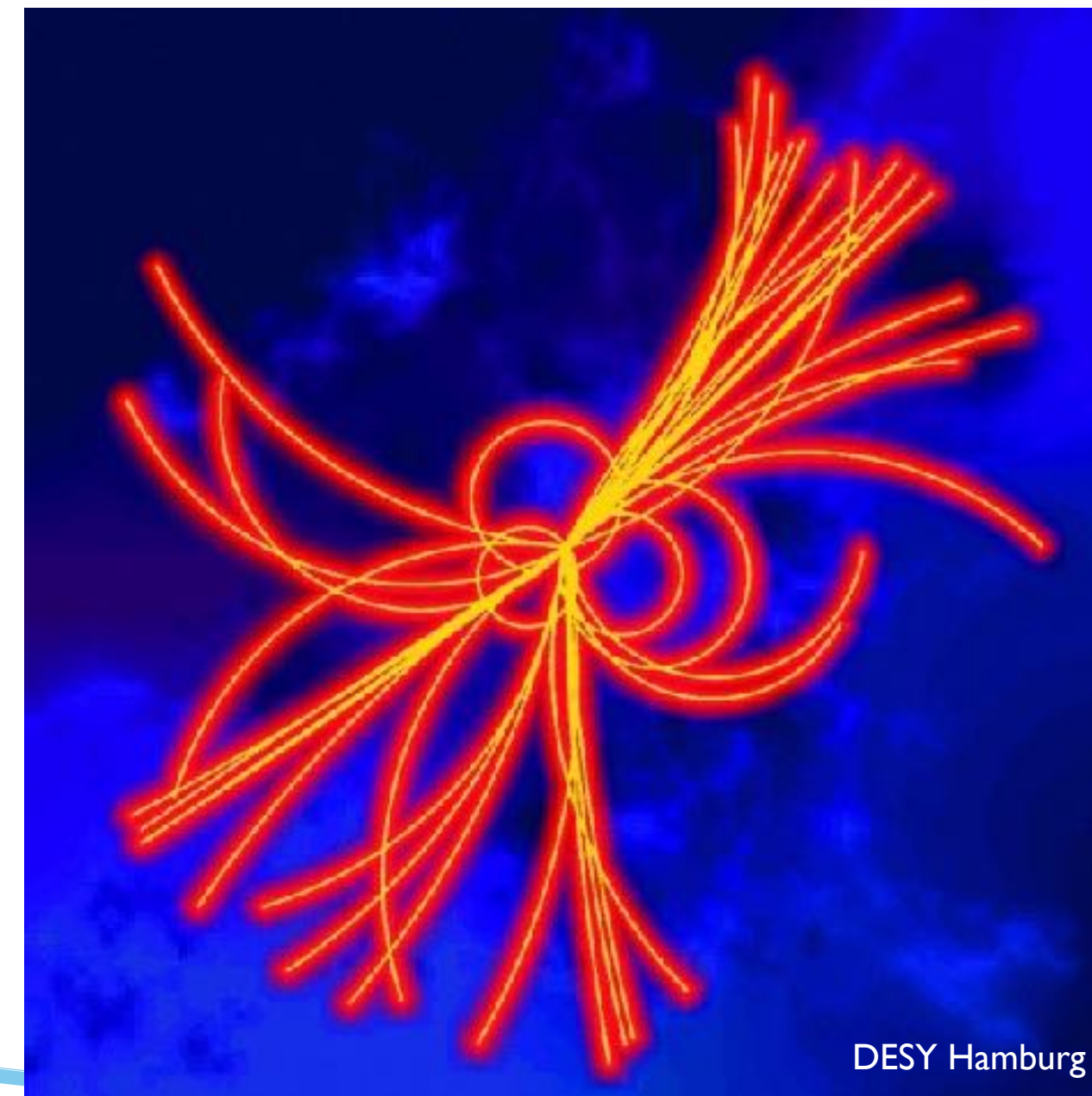


Illustration of an FEL-pulse diffracting off a protein (XFEL, DESY)

Synchrotron photon sources, e.g. Free-Electron Lasers (FELs)
investigation of processes on atomic and molecular scales

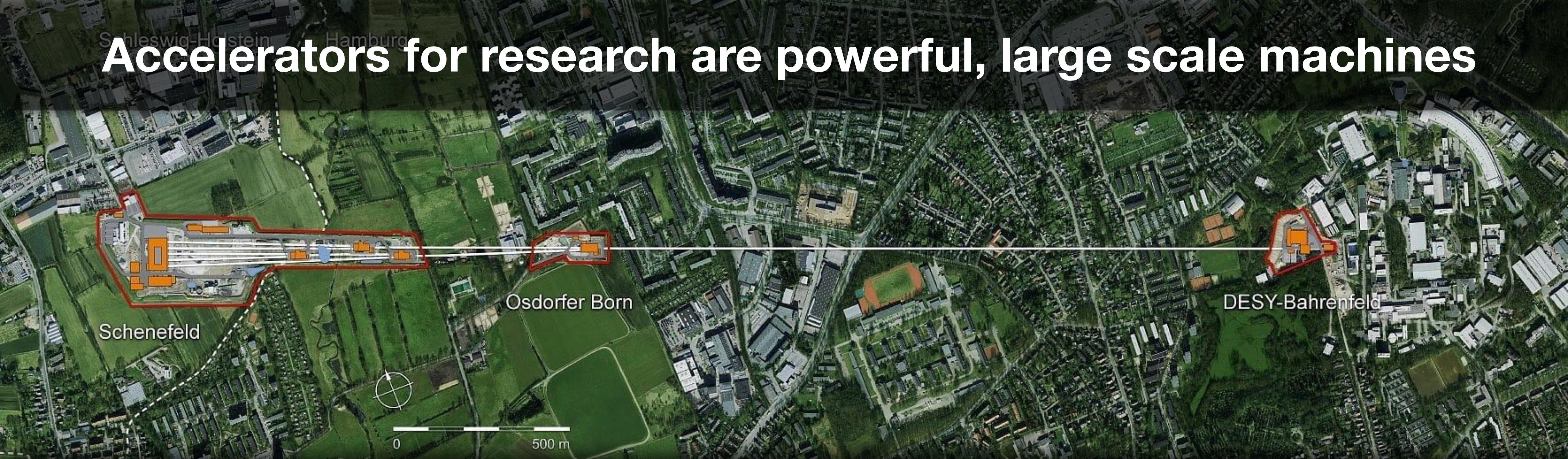


Simulation of the decay of a Higgs Boson (LHC, CERN)

Particle colliders
investigation of the fundamental forces and constituents of matter



Accelerators for research are powerful, large scale machines



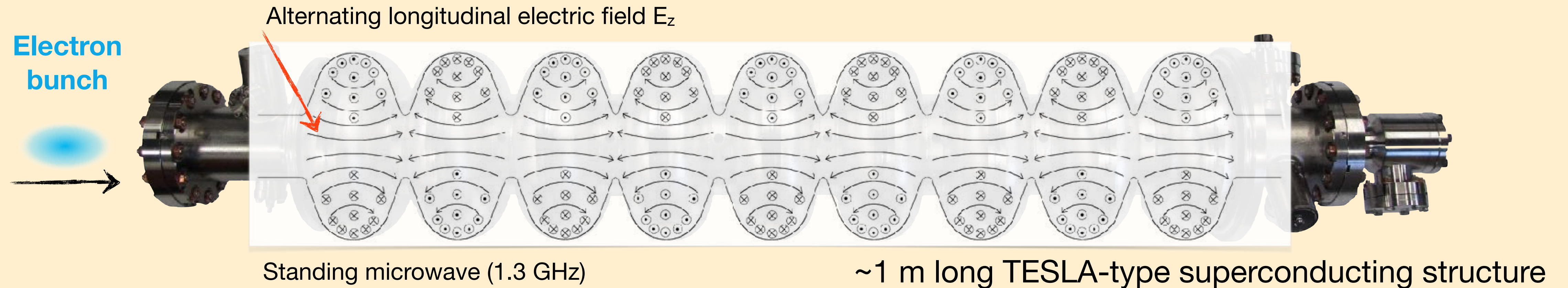
European XFEL 17 GeV electron accelerator



What defines the scale length of the accelerator?

Limits of conventional technology

Working principle of an RF-cavity



The goal: electrons with well defined energy gain

$$\text{Kinetic energy gain } \Delta W_{kin} = eE_z d$$

Accelerating field strength limited to
~50 MV/m by electrical breakdown

Energy increase can only be achieved
by longer acceleration distances!

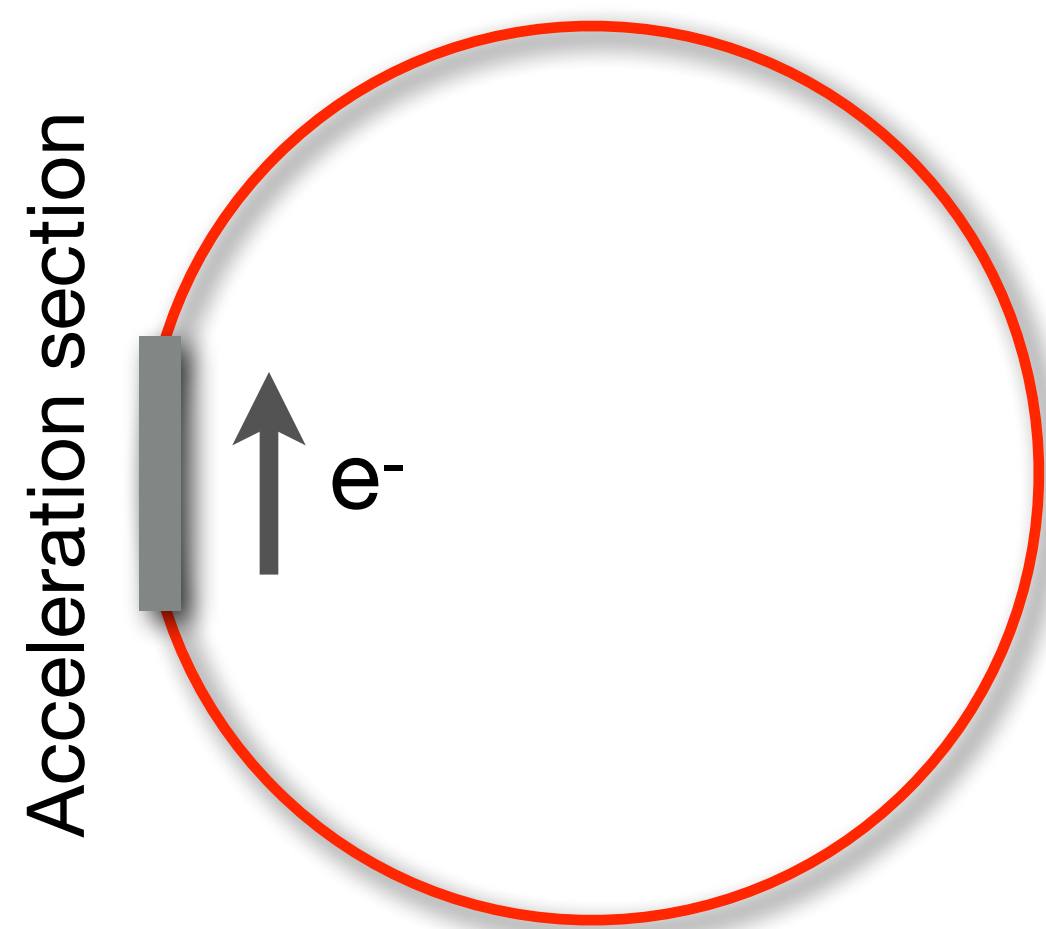
What defines the scale length of the accelerator?

Limits of conventional technology

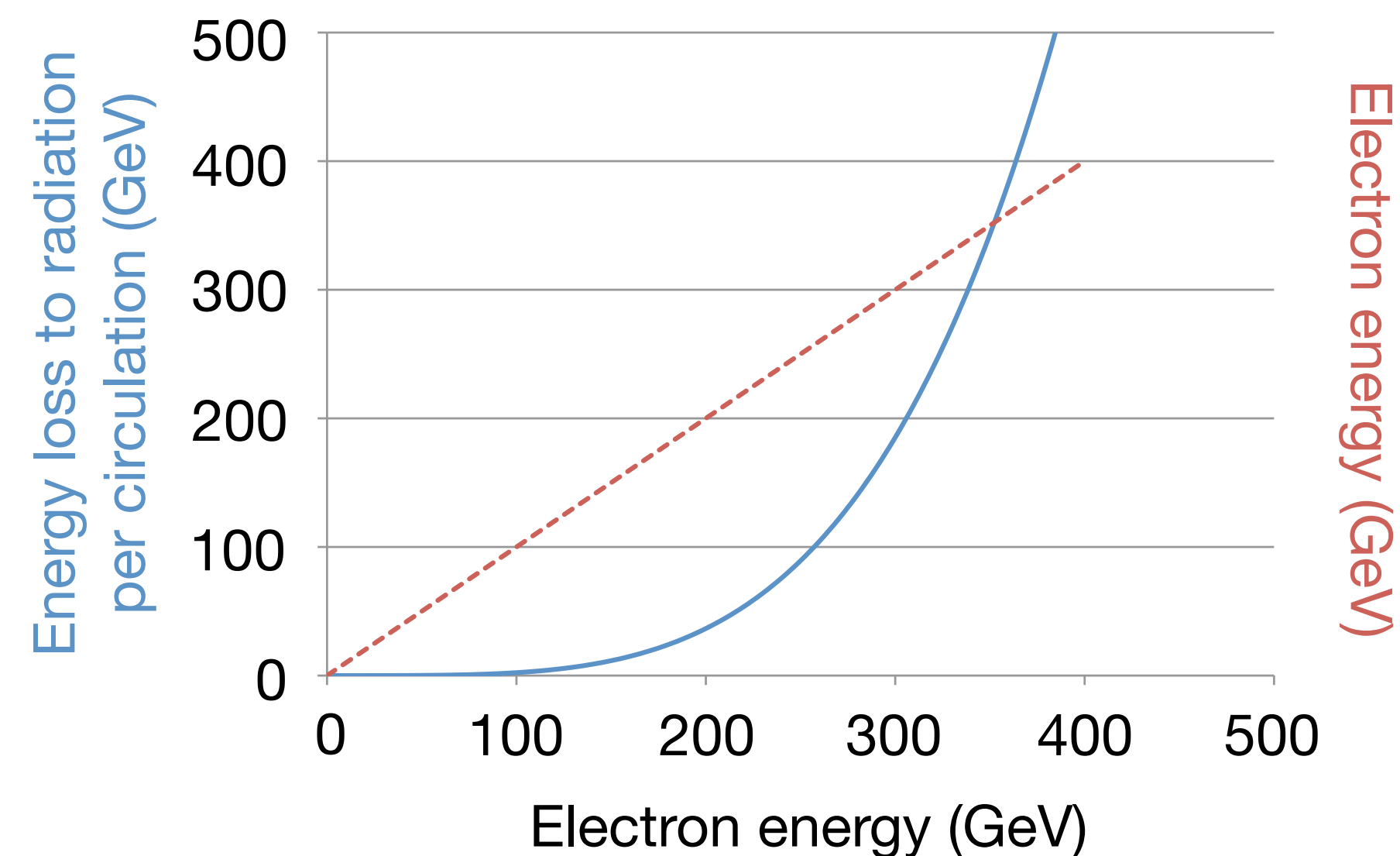
Ring accelerators for electrons?

Advantage: the same (short) acceleration section may be used multiple times

Disadvantage: the energy loss by synchrotron radiation limits the maximum energy
(and achievable beam quality: insufficient for X-ray FELs...)

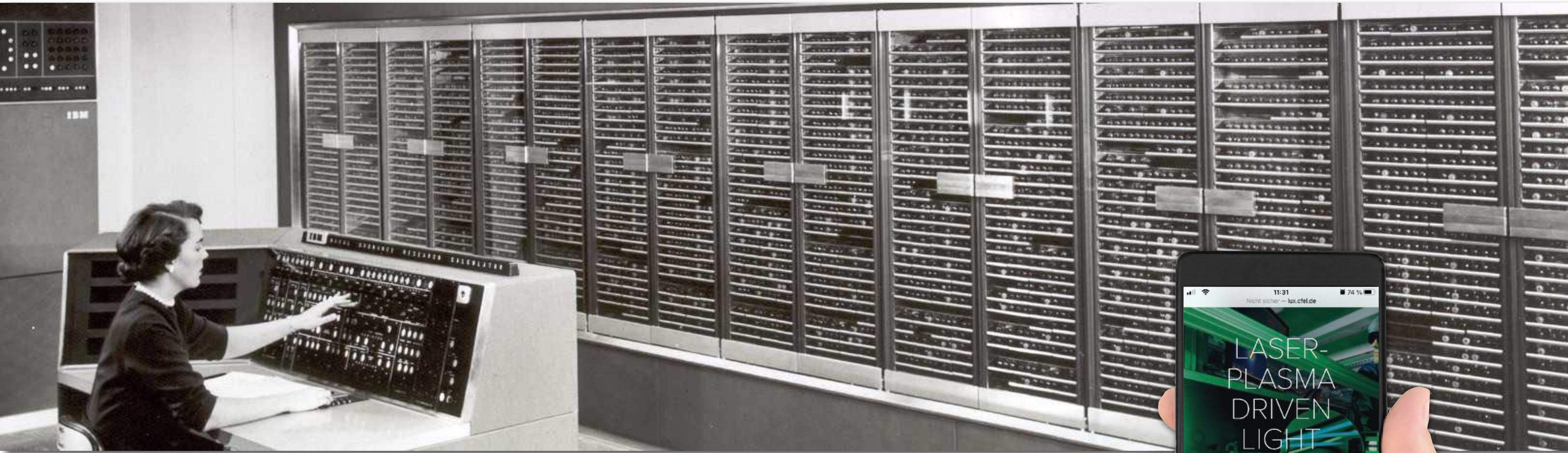


Example: utilize the LEP/LHC ring (27 km circumference) for electrons



$$\Delta W_{syn} \propto \frac{1}{r^2} \frac{W_{kin}^4}{(m_e c^2)^4}$$

Computers shrunk and became super powerful in the past 60 years



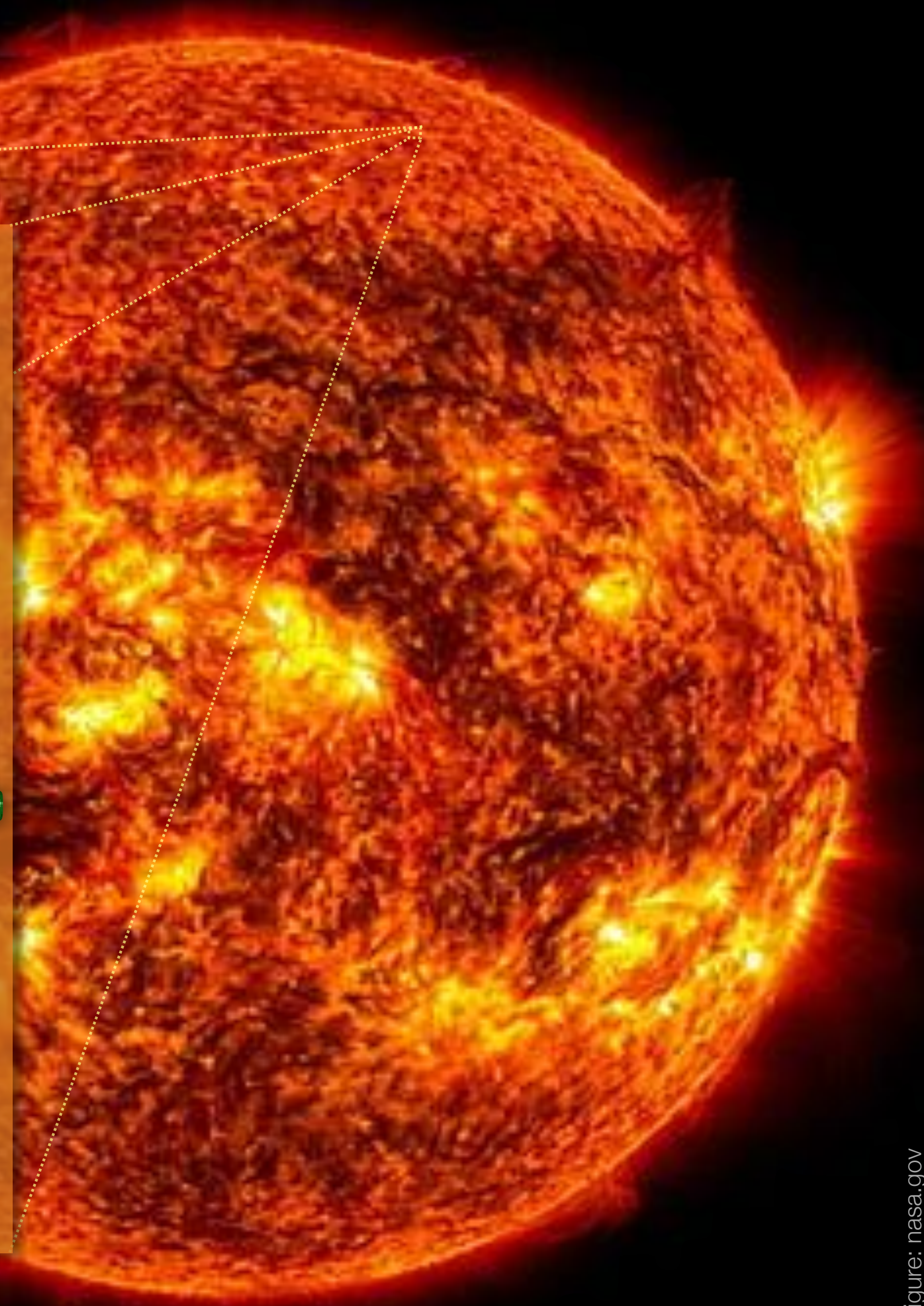
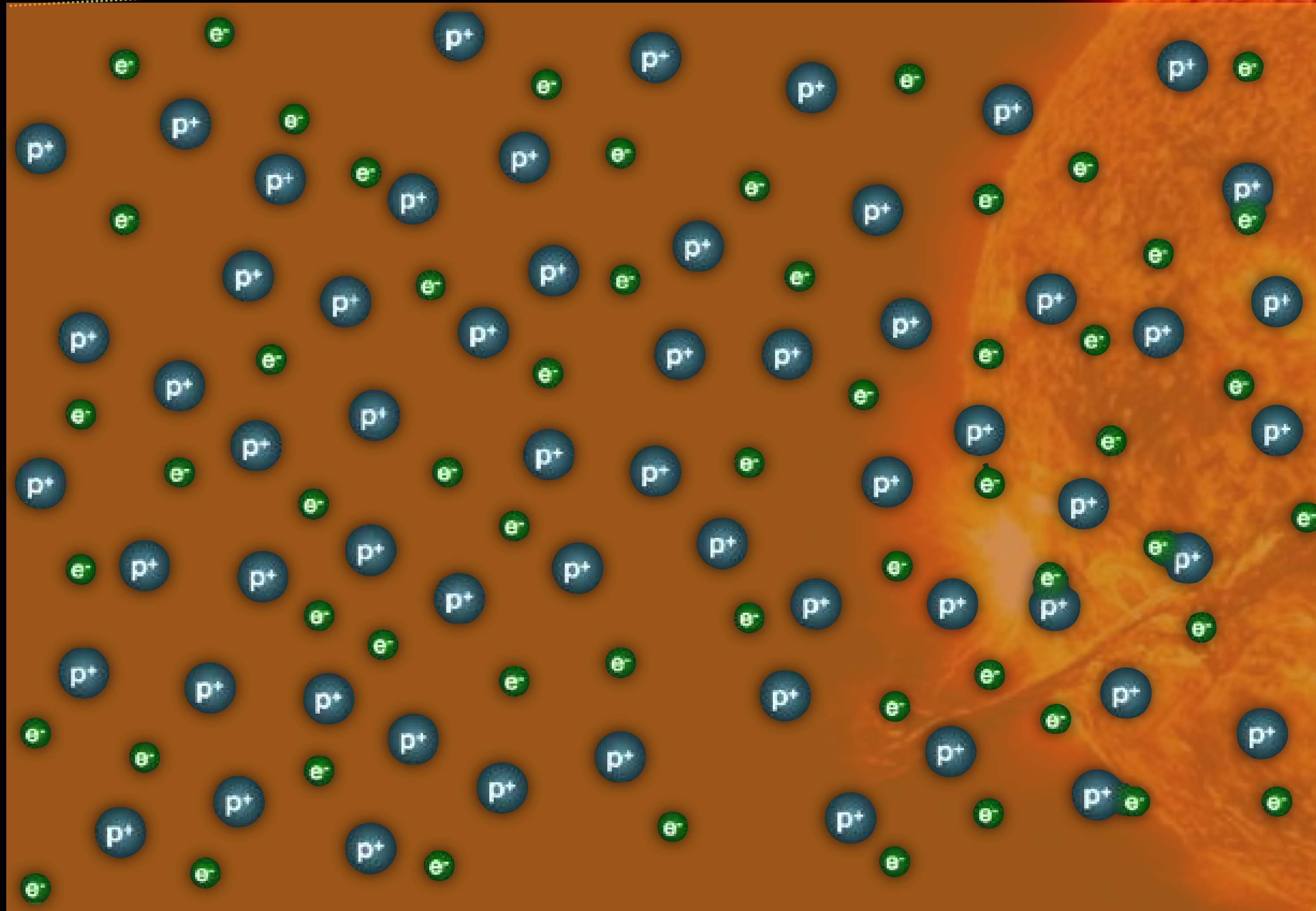
IBM NORC supercomputer (1954)

Can the same be done with particle accelerators?

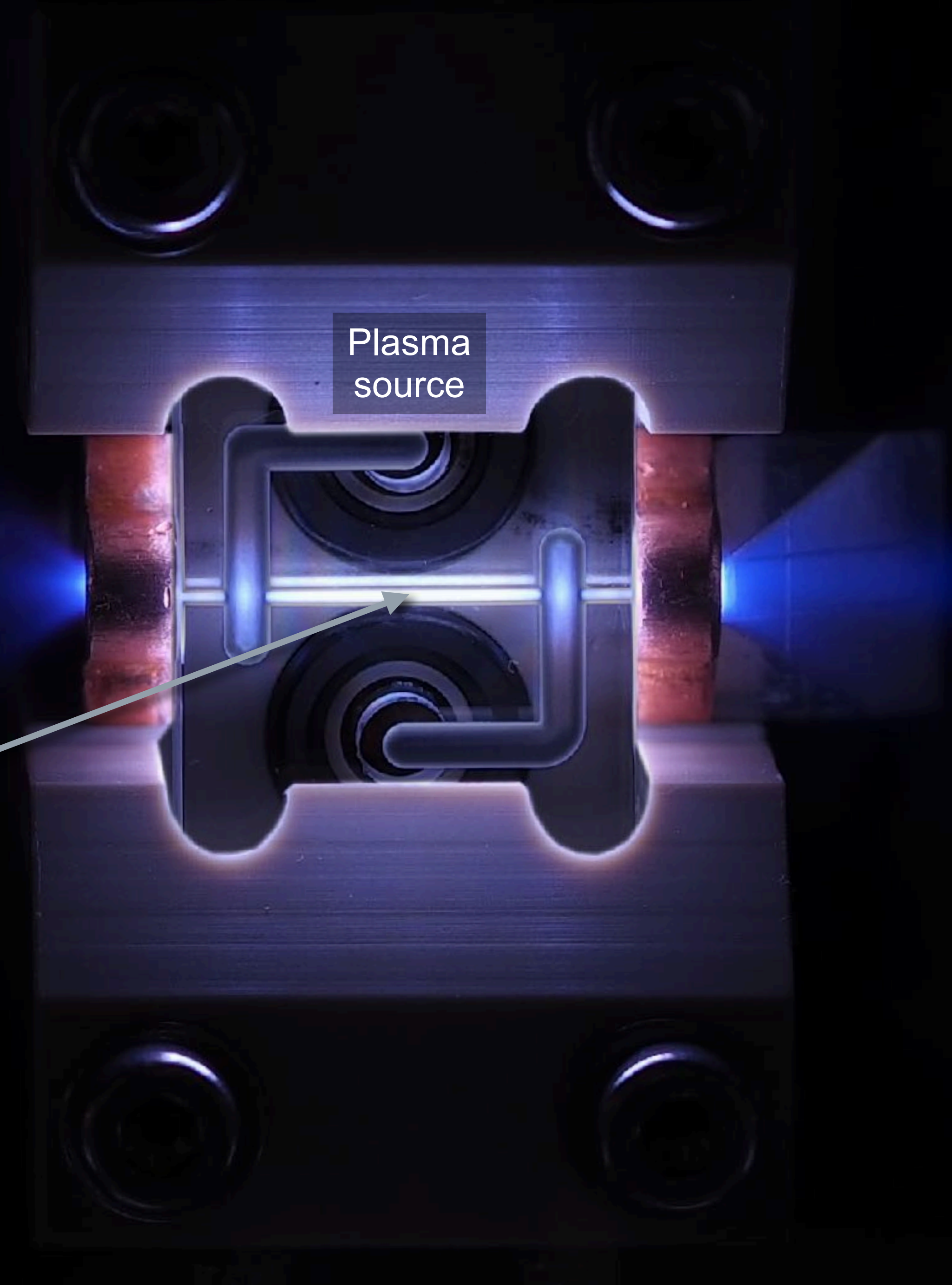
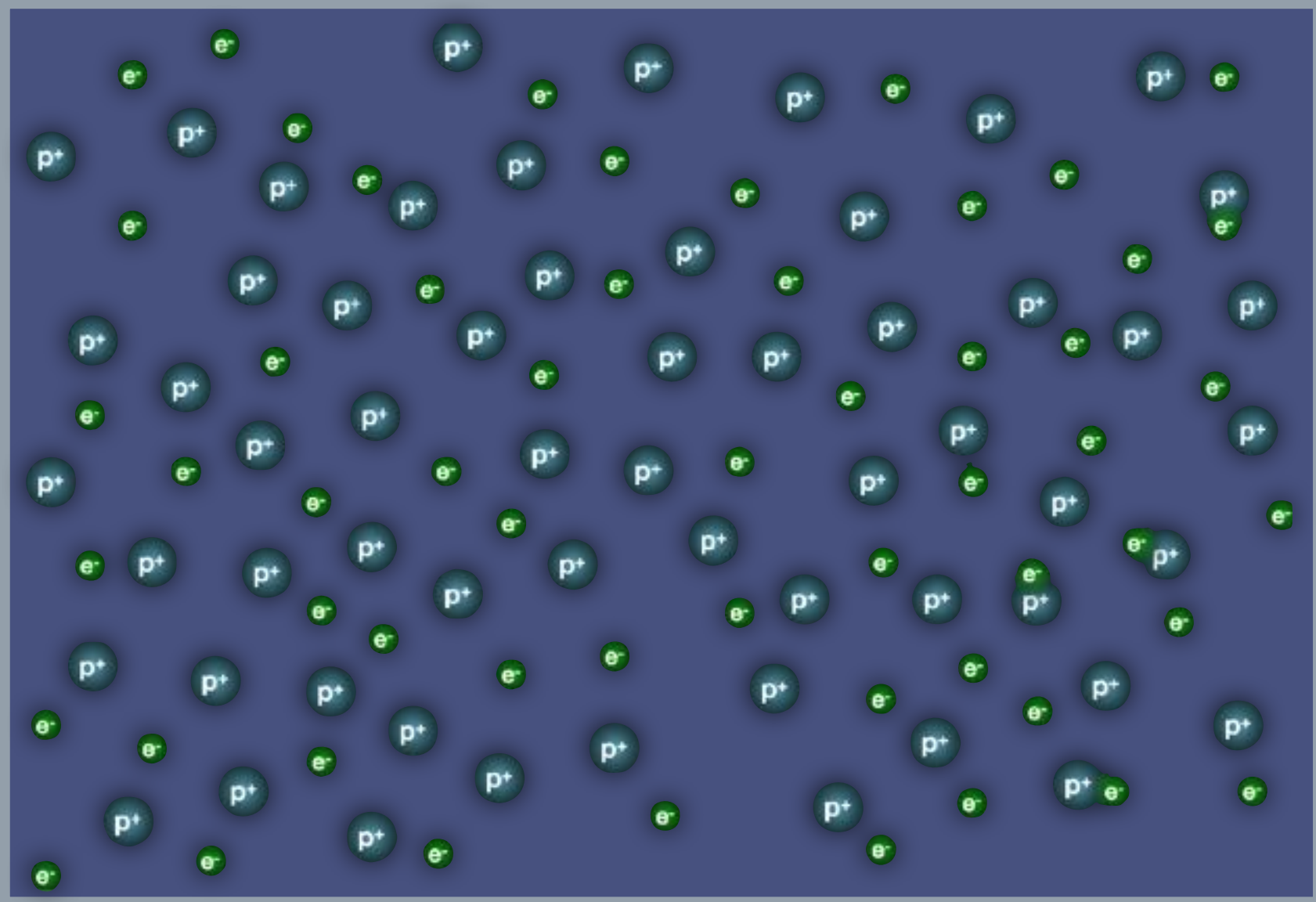
What new applications will this enable?

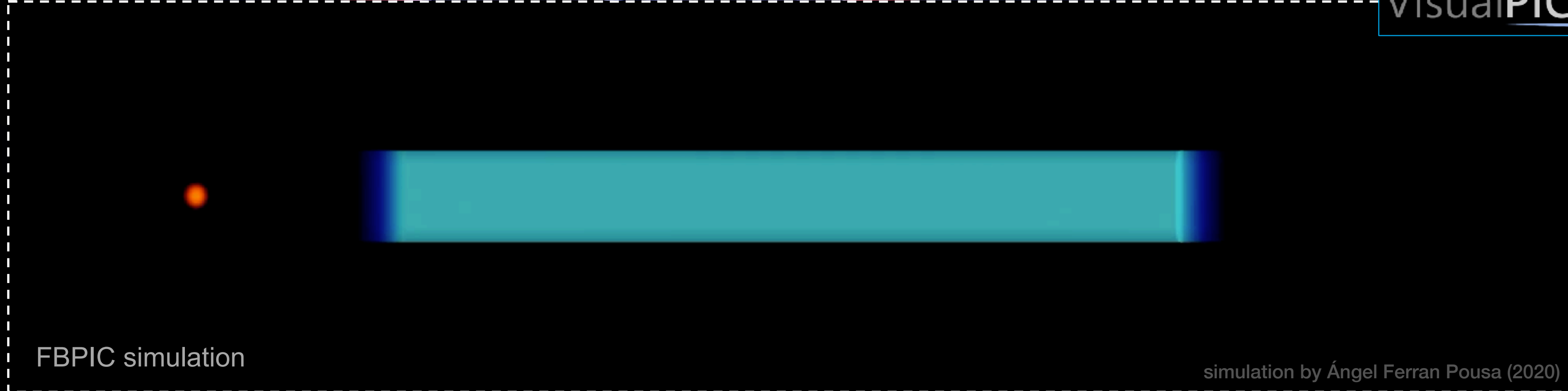
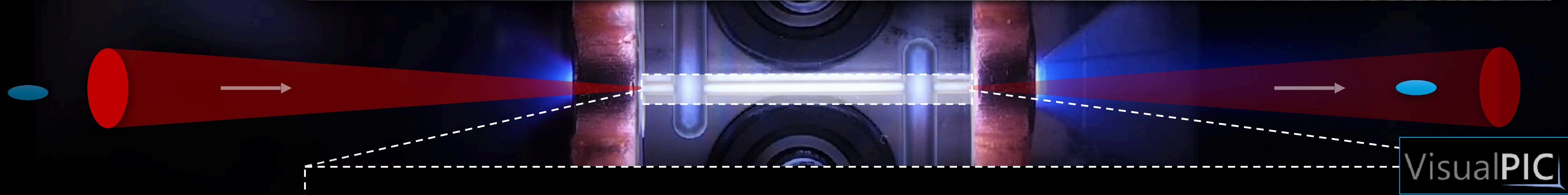
Plasma is everywhere

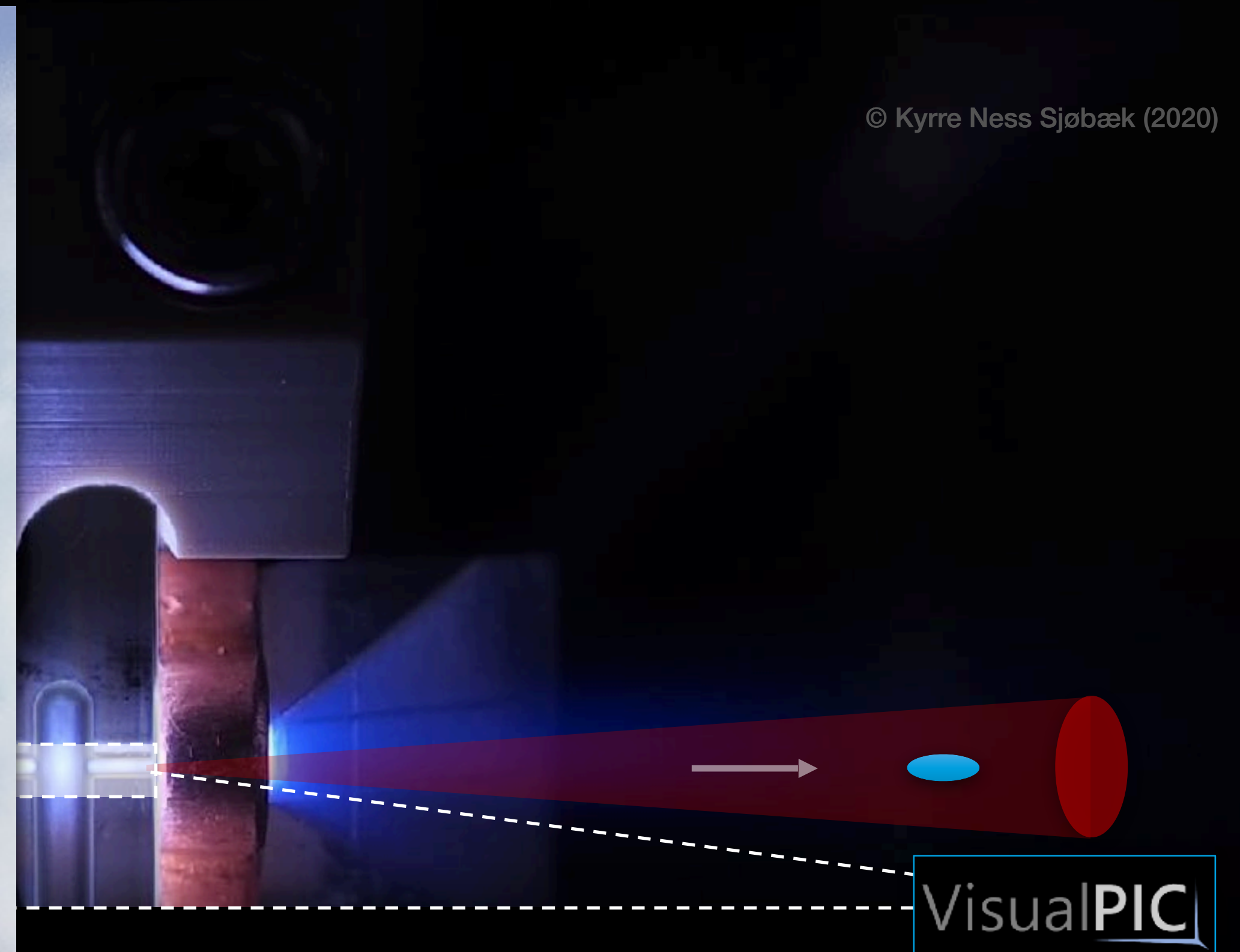
A soup of ionized (broken-down) matter



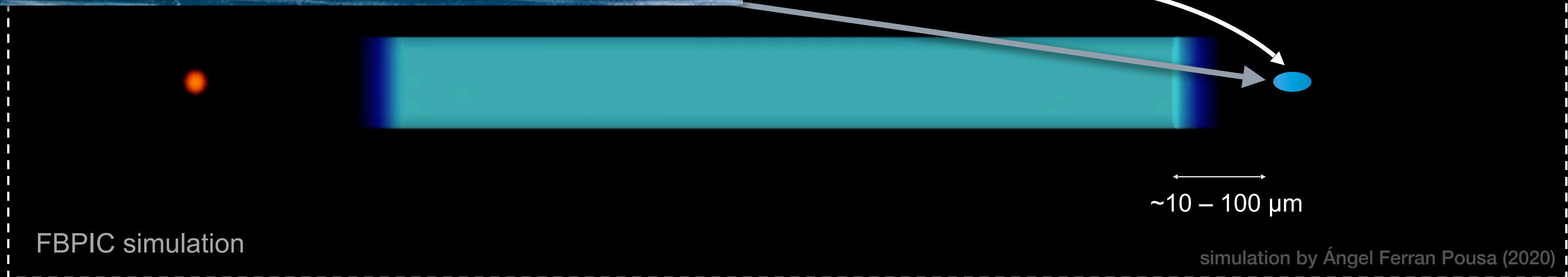
Plasma
source

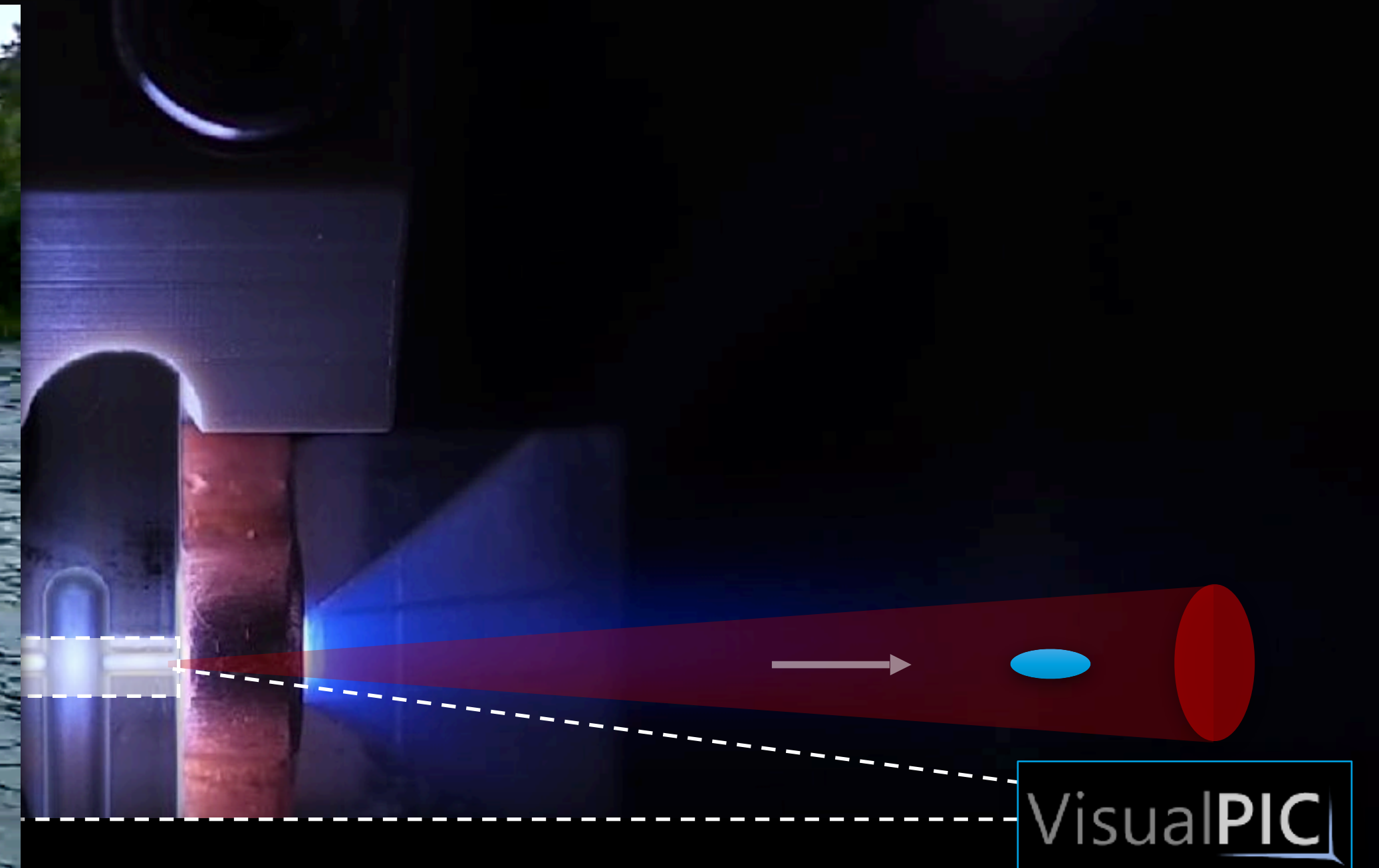




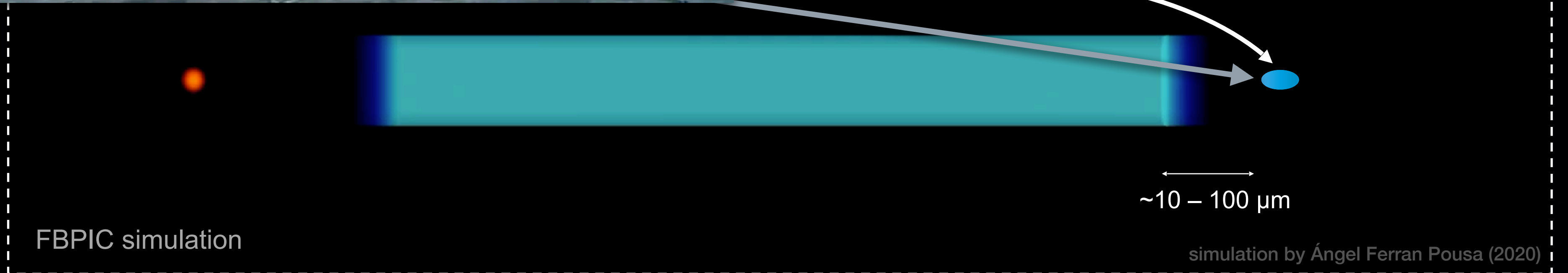


Electron beam can be **externally injected** or formed from trapped plasma electrons (**internal injection**)



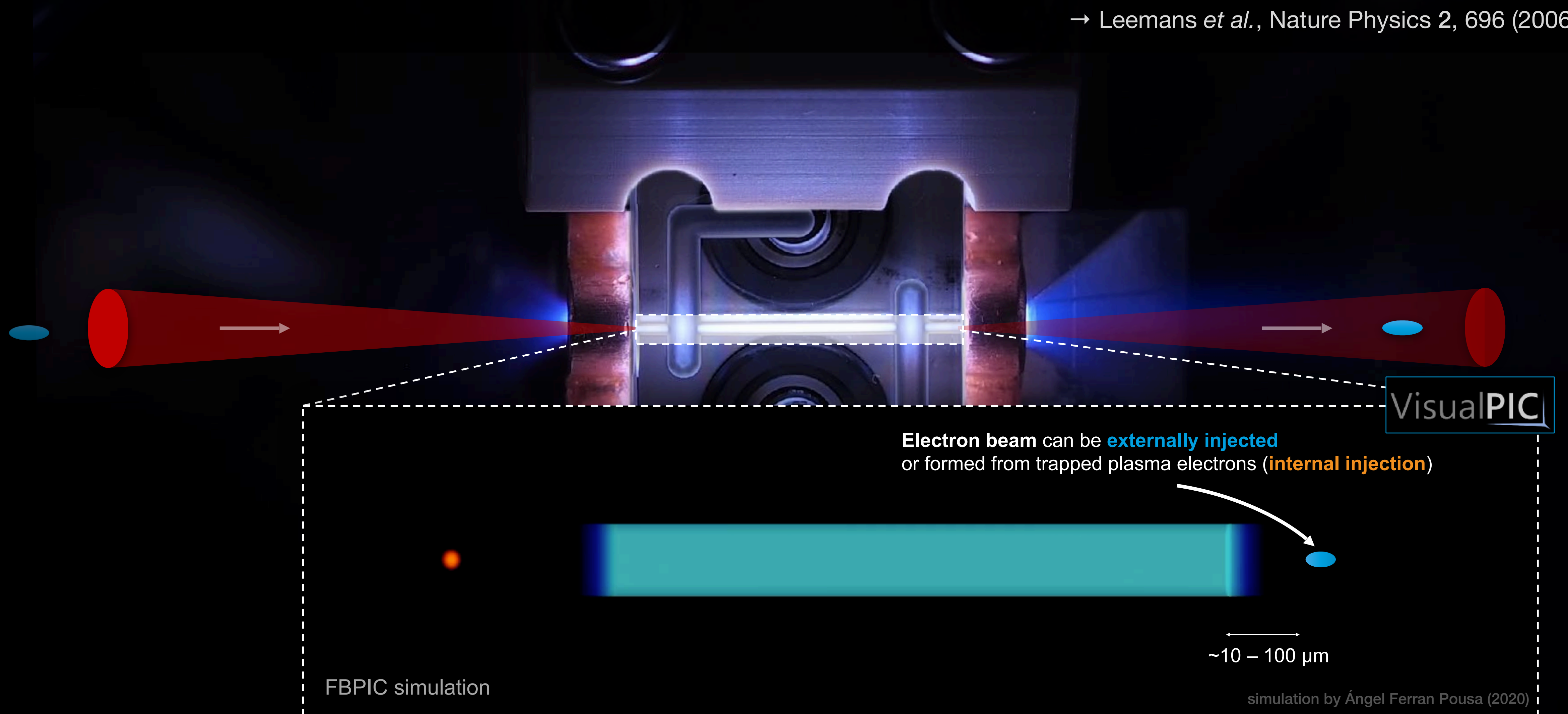


Electron beam can be **externally injected** or formed from trapped plasma electrons (**internal injection**)



Plasma accelerators are a centimeter-scale source of GeV beams

→ Leemans *et al.*, Nature Physics 2, 696 (2006)



Plasma wakefields can sustain accelerating fields of up to $\sim 1\text{-}100 \text{ GV/m}$
with focusing gradients above $\sim 1 \text{ MT/m}$

x1000 more than
RF technology

Simple fluid model for plasma-wave excitation

1D in space / 3D in momentum

Transverse electron momentum
(from equation of motion)

$$\gamma\beta_y = a_0$$

Longitudinal electron momentum
(from equation of motion)

$$\frac{d}{dt}(\gamma\beta_x) = c \left(\frac{\partial\phi_0}{\partial x} - \frac{1}{2\gamma} \frac{\partial a_0^2}{\partial x} \right)$$

Continuity equation

$$\frac{\partial n_e}{\partial t} + c \frac{\partial}{\partial x} (n_e \beta_x) = 0$$

Electro-magnetic wave equation

$$\frac{\partial^2 a_0}{\partial t^2} - c^2 \frac{\partial^2 a_0}{\partial y^2} = -\omega_p^2 \frac{n_0 a_0}{\gamma}$$

Poisson's equation

$$\frac{\partial^2 \phi_0}{\partial x^2} = \frac{\omega_p^2}{c^2} (n_0 - 1)$$

→ Transformation into a co-moving frame with
 $\tau = t$ and $\xi = x - v_g t$ and *quasi-static approximation*



Resulting differential equation for scalar potential

$$\begin{aligned} \frac{\partial^2 \phi_0}{\partial \xi^2} &= \frac{\omega_p^2}{c^2} (n_0 - 1) \\ &= \frac{\omega_p^2}{c^2} \gamma_g^2 \left[\frac{\beta_g (1 + \phi_0)^2}{\sqrt{(1 + \phi_0)^2 - \frac{1+a_0^2}{\gamma_g^2}}} - 1 \right] \end{aligned}$$

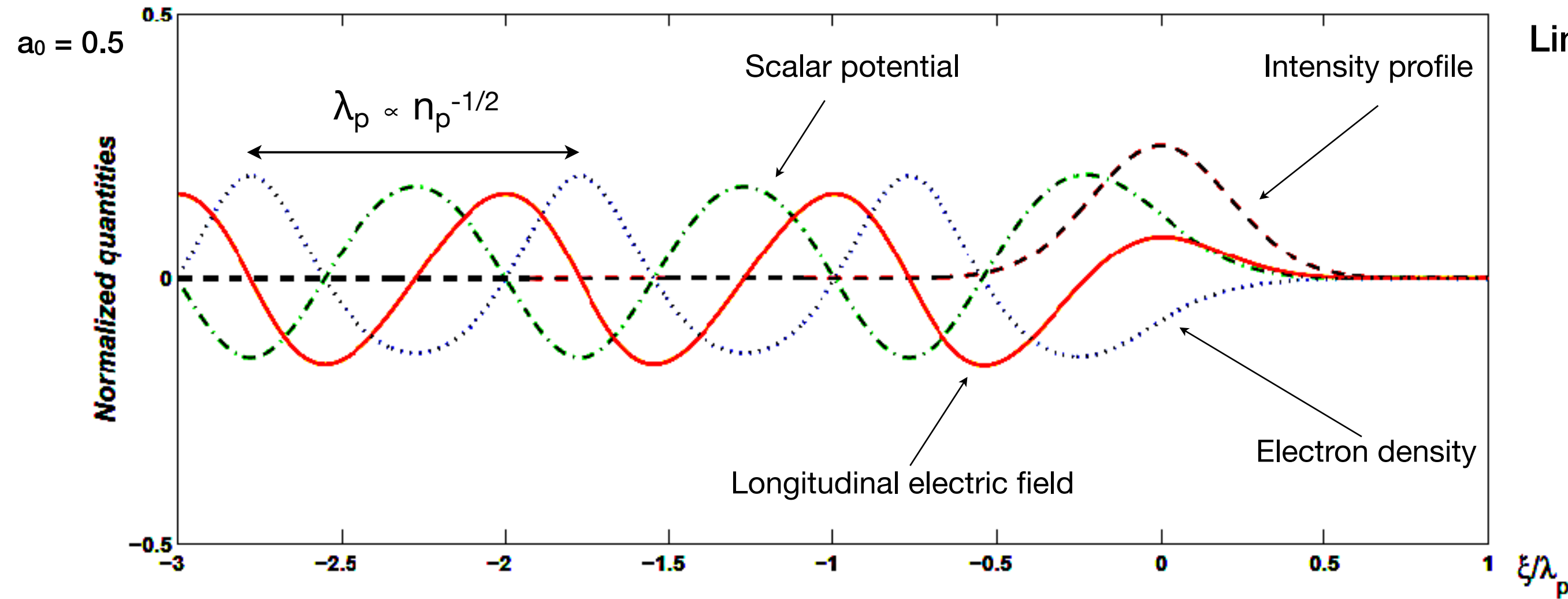


with $\beta_g \rightarrow 1$

$$\frac{\partial^2 \phi_0}{\partial \xi^2} = \frac{\omega_p^2}{2c^2} \left[\frac{1 + a_0^2}{(1 + \phi_0)^2} - 1 \right]$$

Solution for the scalar wake potential for a Gaussian laser pulse

1D in space / 3D in momentum

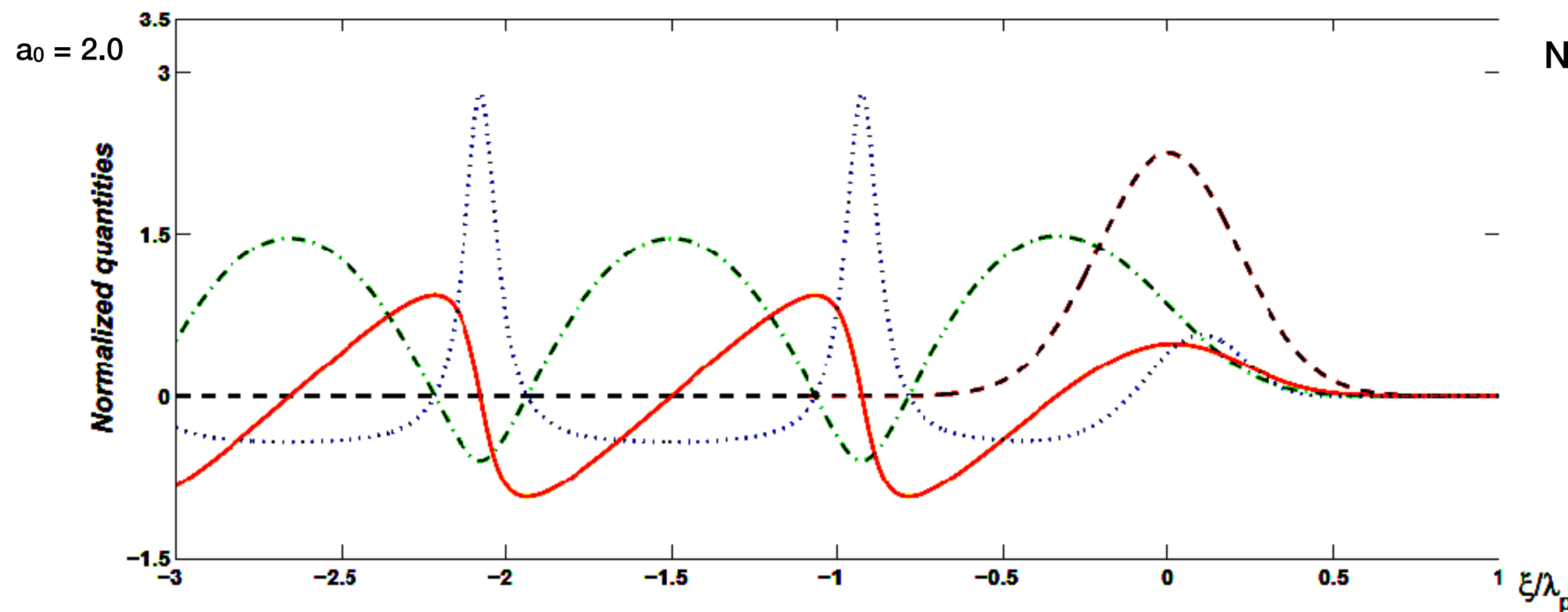


Linear regime

$$\frac{\partial^2 \phi_0}{\partial \xi^2} = \frac{\omega_p^2}{2c^2} \left[\frac{1 + a_0^2}{(1 + \phi_0)^2} - 1 \right]$$

Temporal laser-pulse shape in intensity

$$a^2(\xi) = a_0^2 \exp \left[- \left(\frac{\xi - \xi_0}{L} \right)^2 4 \log(2) \right]$$



Nonlinear regime

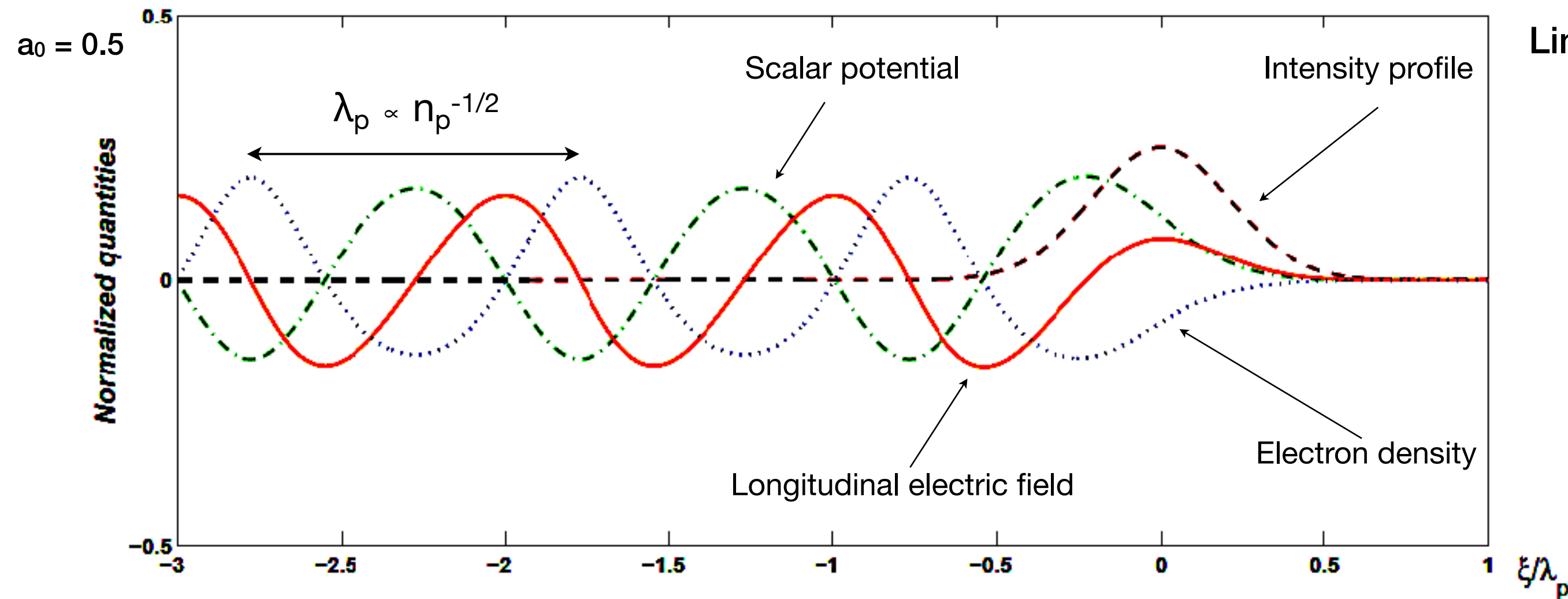
Accelerating field strength

$$E \sim \left(\frac{mc\omega_p}{e} \right) \frac{a^2}{(1 + a^2/2)} \approx (96 \text{ V/m}) \sqrt{n_0 [\text{cm}^{-3}]} \frac{a^2}{(1 + a^2/2)}$$

e.g. $E \approx 100 \text{ GV/m}$ (for $n \approx 10^{18} \text{ cm}^{-3}$, $a \approx 1$)

Solution for the scalar wake potential for a Gaussian laser pulse

1D in space / 3D in momentum

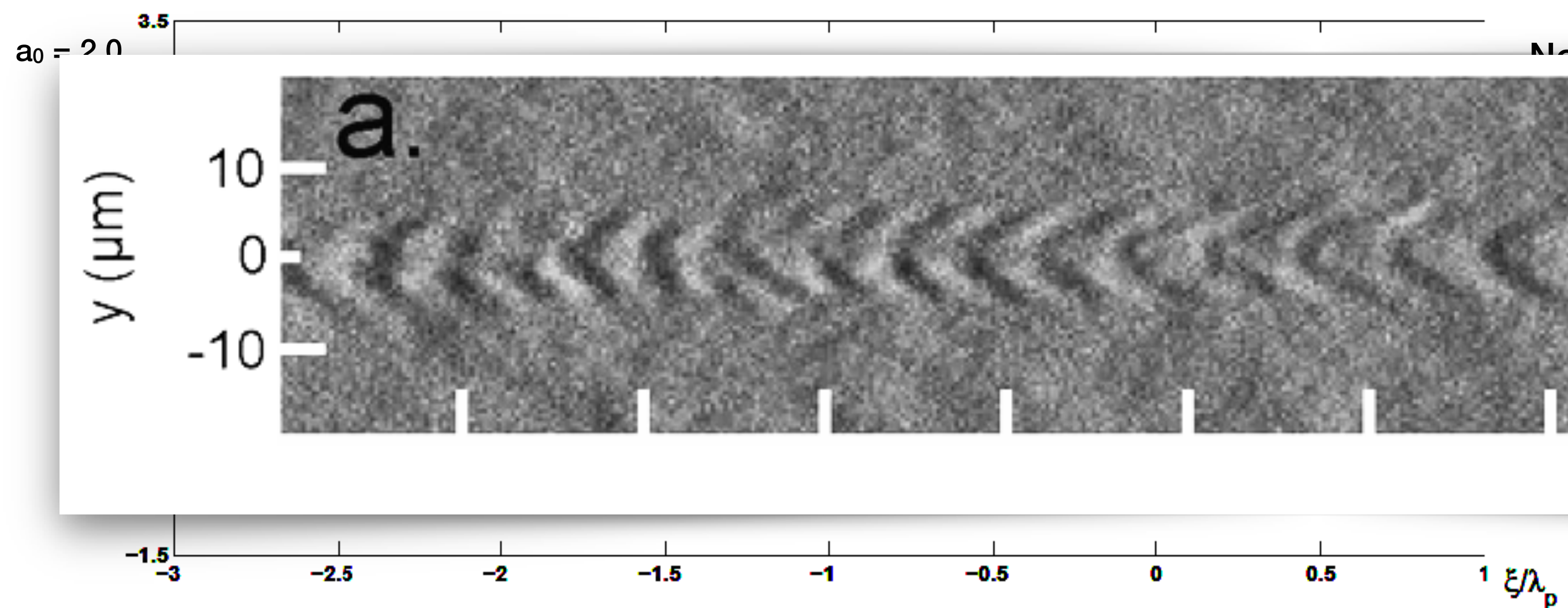


Linear regime

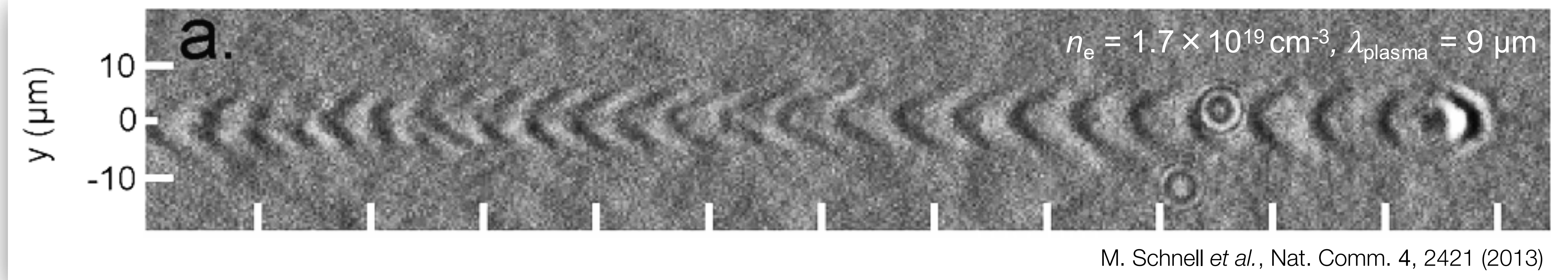
$$\frac{\partial^2 \phi_0}{\partial \xi^2} = \frac{\omega_p^2}{2c^2} \left[\frac{1 + a_0^2}{(1 + \phi_0)^2} - 1 \right]$$

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Nonlinear regime



M. Schnell *et al.*, Nat. Comm. 4, 2421 (2013)

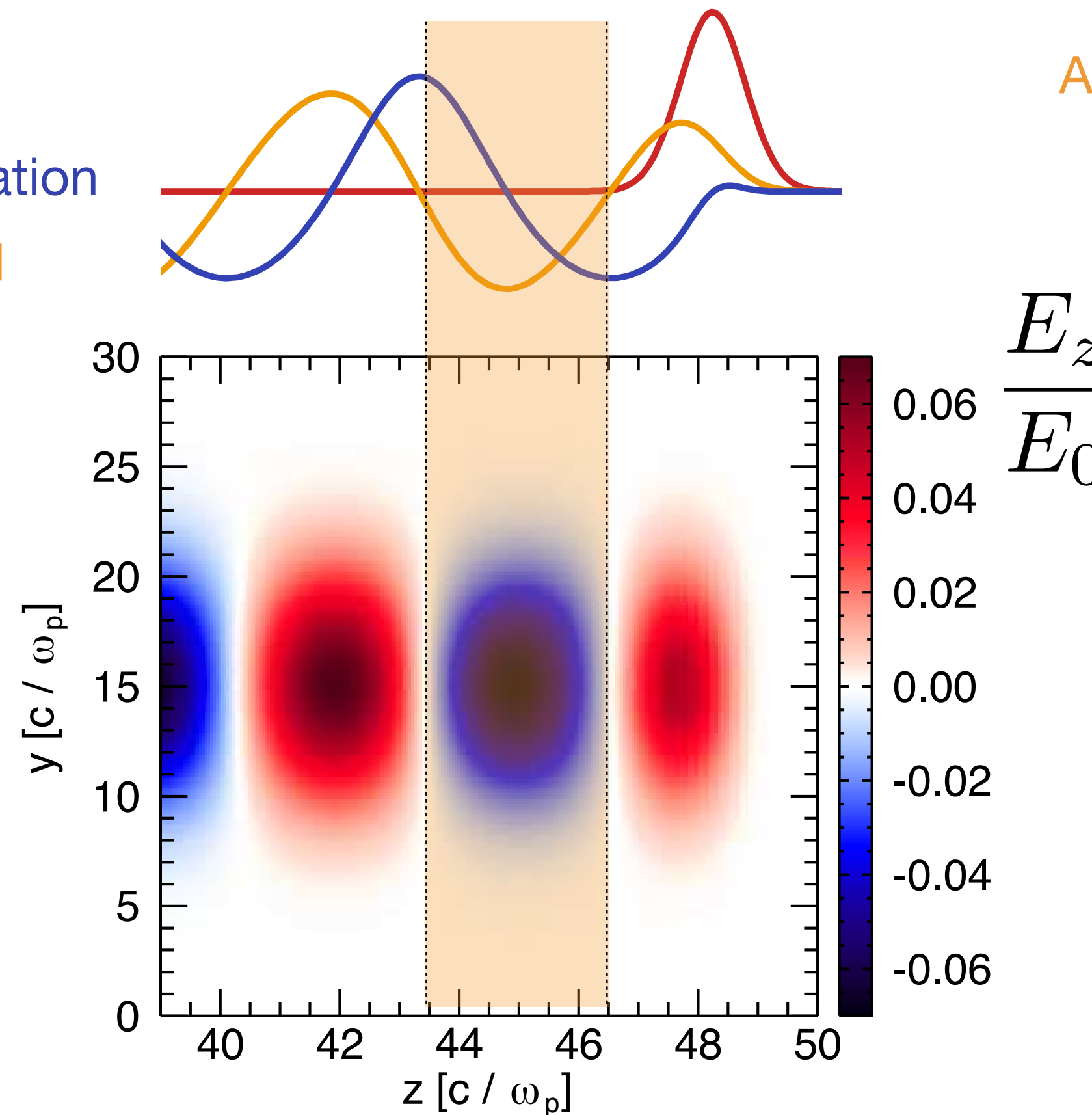
Wakefield properties in transverse dimensions

3D in space / 3D in momentum

Laser envelope

Electron density modification

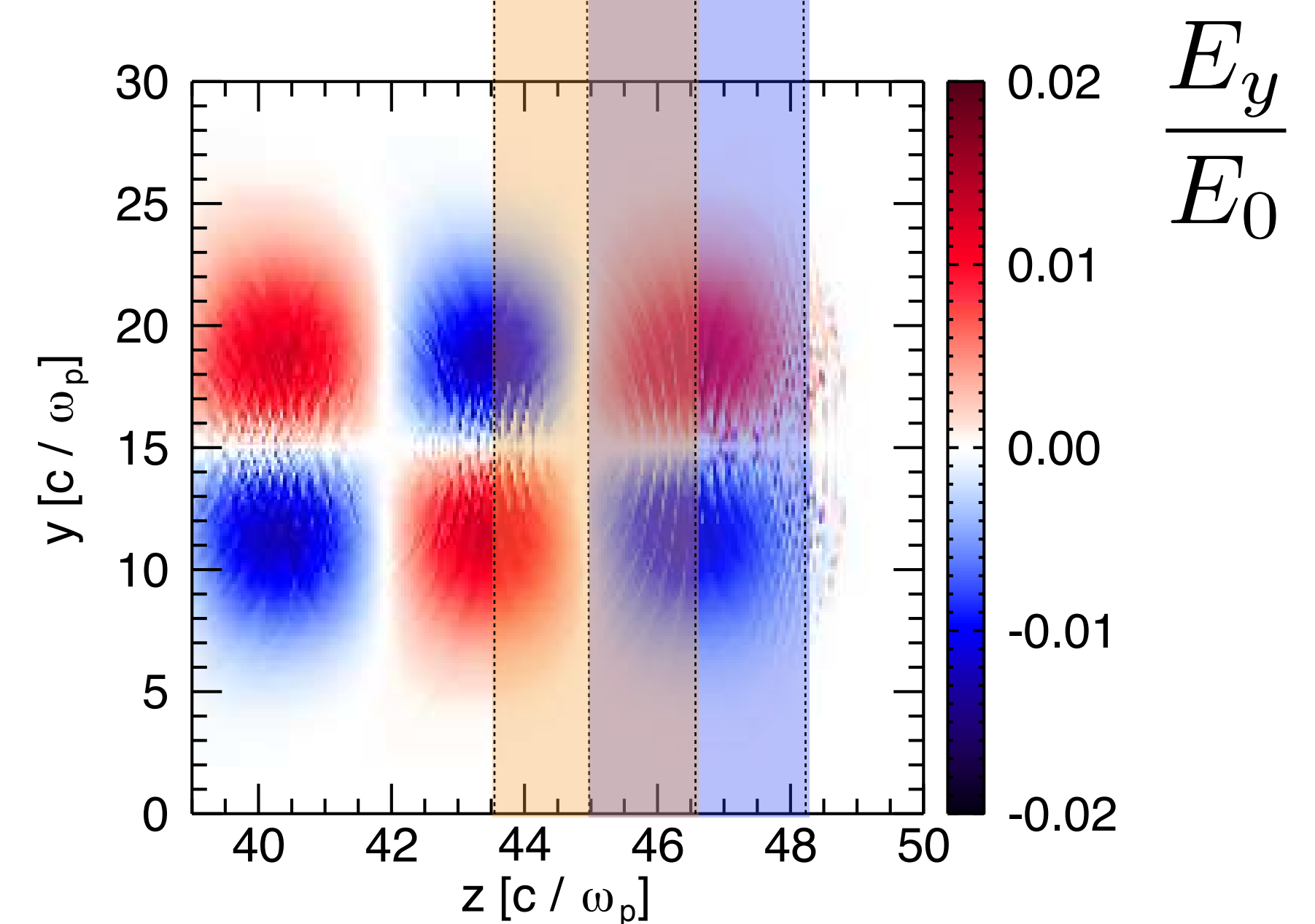
Longitudinal electric field



Longitudinal fields
of a quasi-linear plasma wave

Accelerating phase

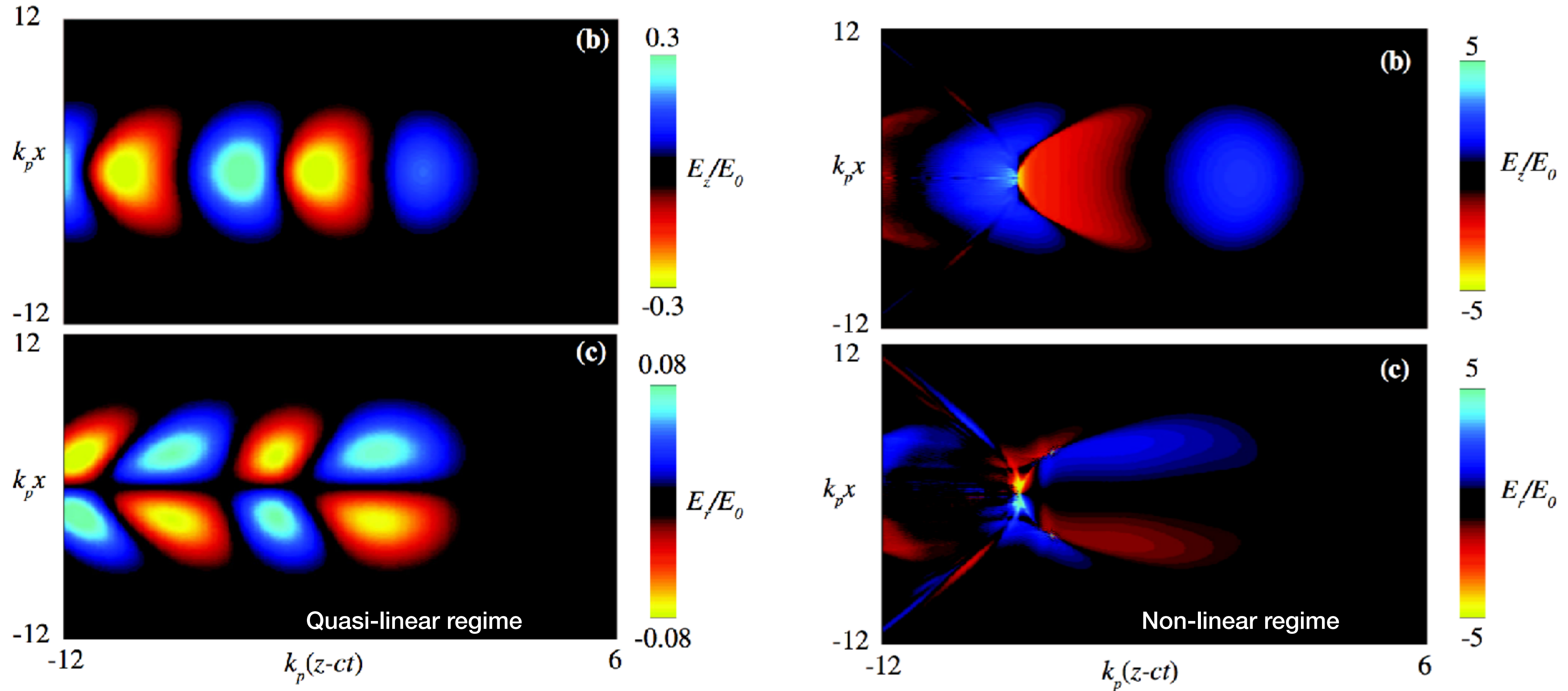
Focusing phase



Transverse fields
of a quasi-linear plasma wave

Wakefield properties in transverse dimensions

3D in space / 3D in momentum



from C.B.Schroeder *et al.*, PRSTAB 13, 101301 (2010)

The LWFA process can be complex

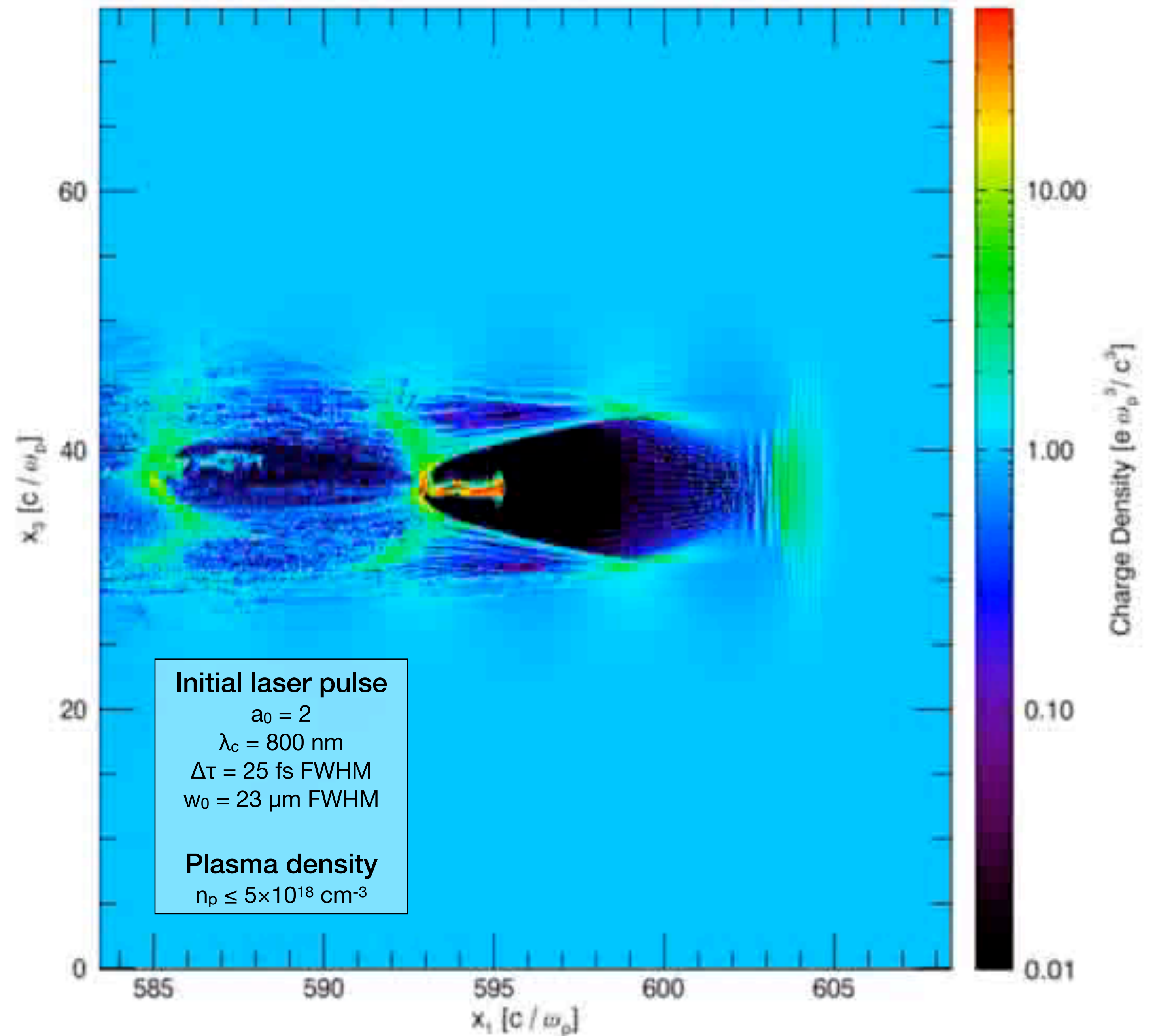
- laser self-focussing
- laser self-compression
- **wave breaking**
- beam hosing
- beam loading
- ...



3D particle-in-cell (PIC) simulation



Time = 583.44 [1 / ω_p]



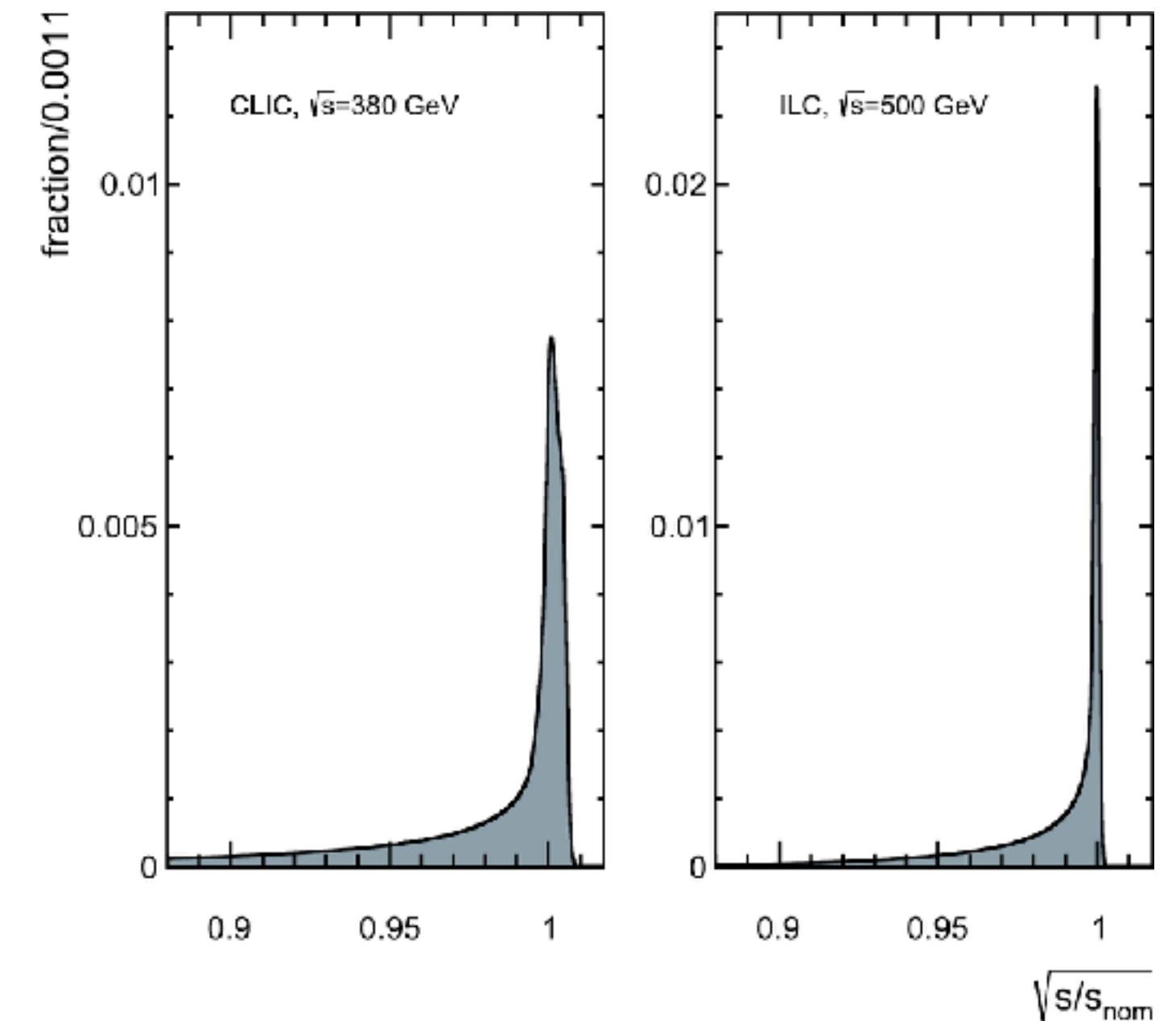
Our customers: high-energy physics and photon science

- > High-energy physics and photon science demand high(est) energy at low cost.
 - > *Solution:* Plasma accelerators — significantly higher acceleration gradients.
- > Simultaneously, particle colliders have strict demands for luminosity: (FELs have similar demands for brightness)

$$\mathcal{L} = \frac{H_D}{8\pi m_e c^2} \frac{P_{\text{wall}}}{\sqrt{\beta_x \beta_y}} \frac{\eta N}{\sqrt{\epsilon_{nx} \epsilon_{ny}}}$$

High repetition rate (points to P_{wall})
 High energy efficiency (points to η)
 Low energy spread (luminosity spectrum, final focusing) (points to $\beta_x \beta_y$)
 Low emittance (points to $\epsilon_{nx} \epsilon_{ny}$)

- > Energy efficiency motivates use of beam-driven plasma acceleration.



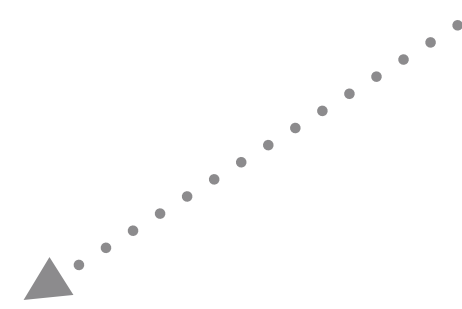
Luminosity distribution across collision energies.
 Source: M. Boronat *et al.*, Phys. Lett. B 804, 135353 (2020).

$$\eta = \eta_{\text{wall} \rightarrow \text{DB}} \times \eta_{\text{DB} \rightarrow \text{WB}}$$

↑
 Beam-drivers are orders of magnitude more efficient than laser-drivers (for now)

Primary goal of FLASHFORWARD▶▶

Develop a self-consistent plasma-accelerator stage
with high-efficiency, high-quality, and high-average-power



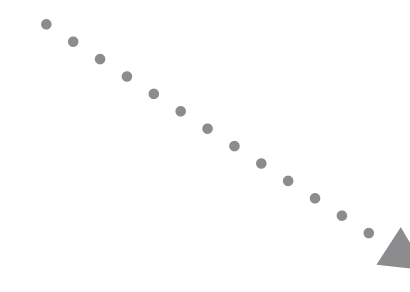
High efficiency

Transfer efficiency
Driver depletion



High beam quality

Energy-spread preservation
Emittance preservation

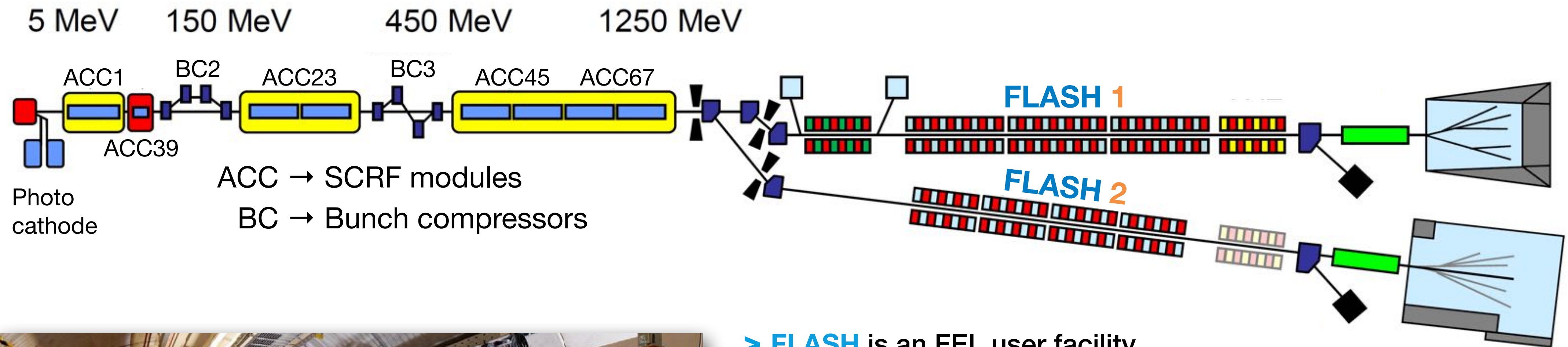


High average power

High repetition rate

FLASHFORWARD ►► utilizes FLASH superconducting accelerator

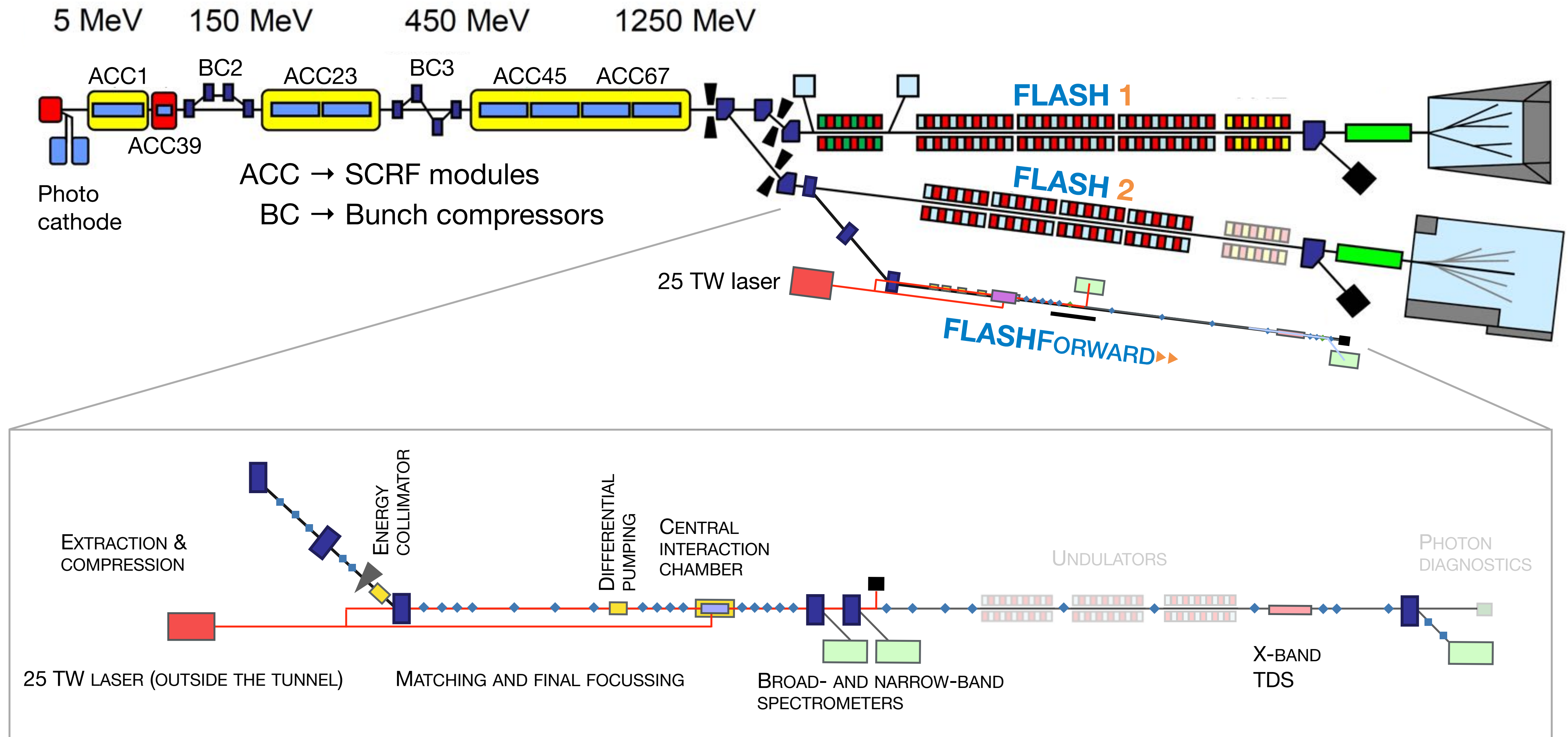
Plasma accelerator tightly integrated into facility and benefits from Free-Electron Laser beam quality



- > **FLASH** is an FEL user facility
 - 10% of beam time dedicated to generic accelerator research
- > Superconducting accelerator based on ILC/XFEL technology
 - ≈ 1.25 GeV energy with \sim nC charge at few 100 fs bunch duration
 - ~ 2 μ m trans. norm. emittance
 - ~ 10 kW average beam power, MHz repetition rate in 10 Hz bursts
 - exquisite stability by advanced feedback/feedforward systems
- > Unique opportunities for plasma accelerator science

FLASHFORWARD ►► utilizes FLASH superconducting accelerator

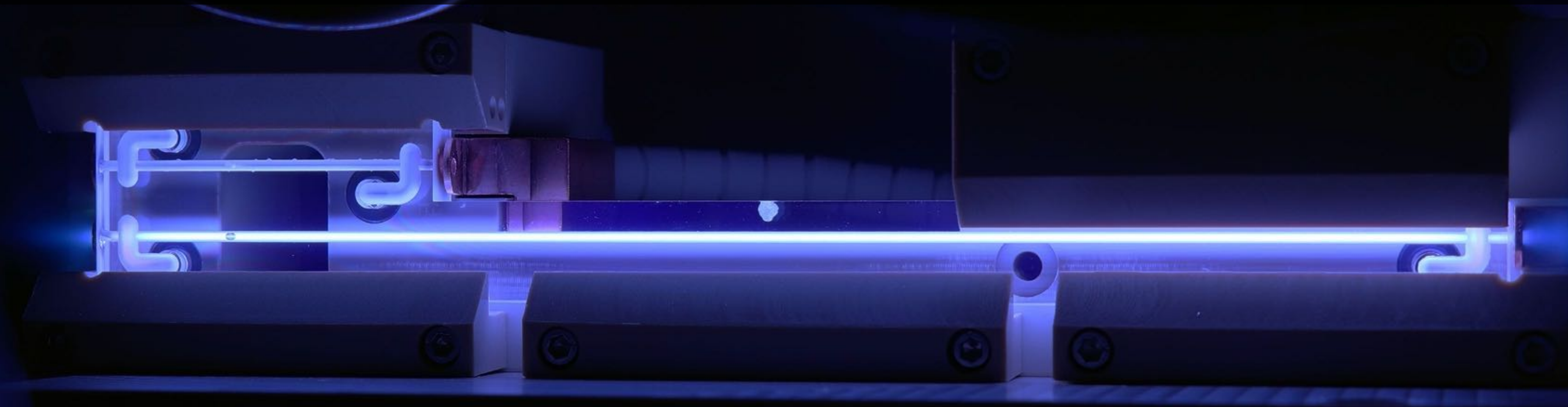
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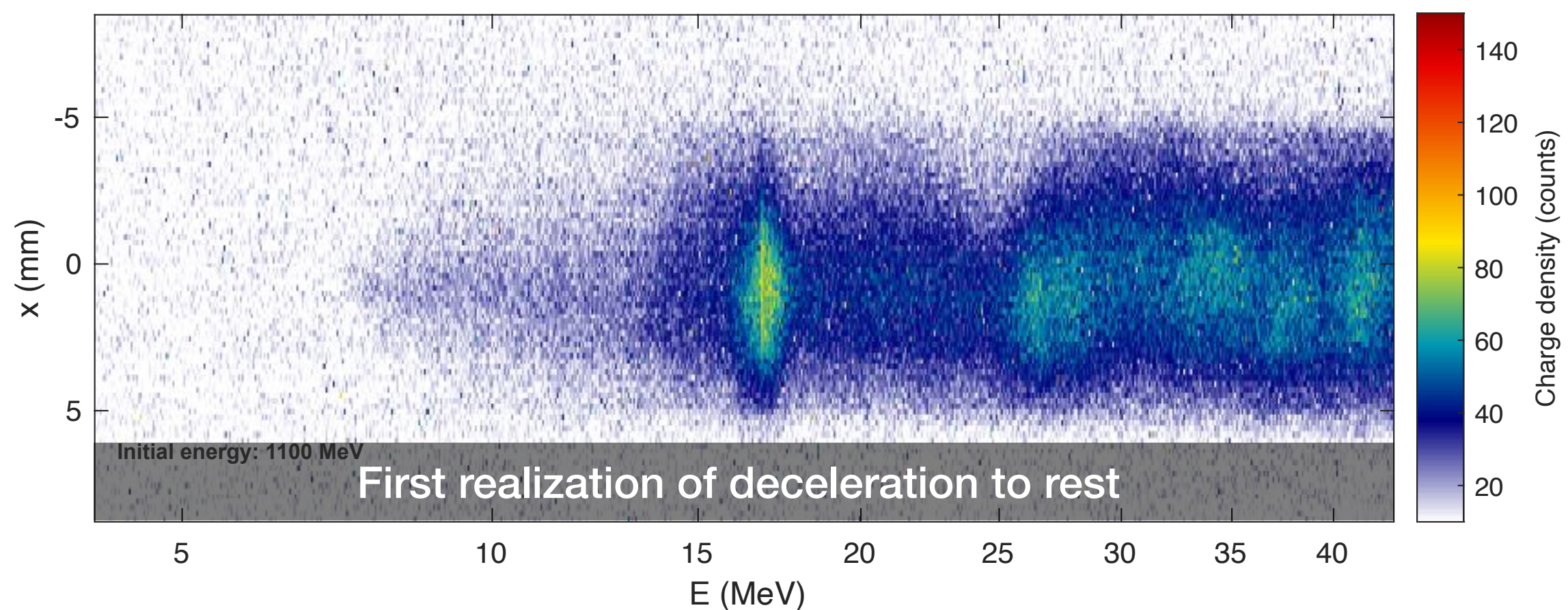
R. D'Arcy *et al.*, *Phil. Trans. R. Soc. A* **377**, 20180392 (2019)

1.1 GeV energy gain and loss achieved in a 195 mm plasma module

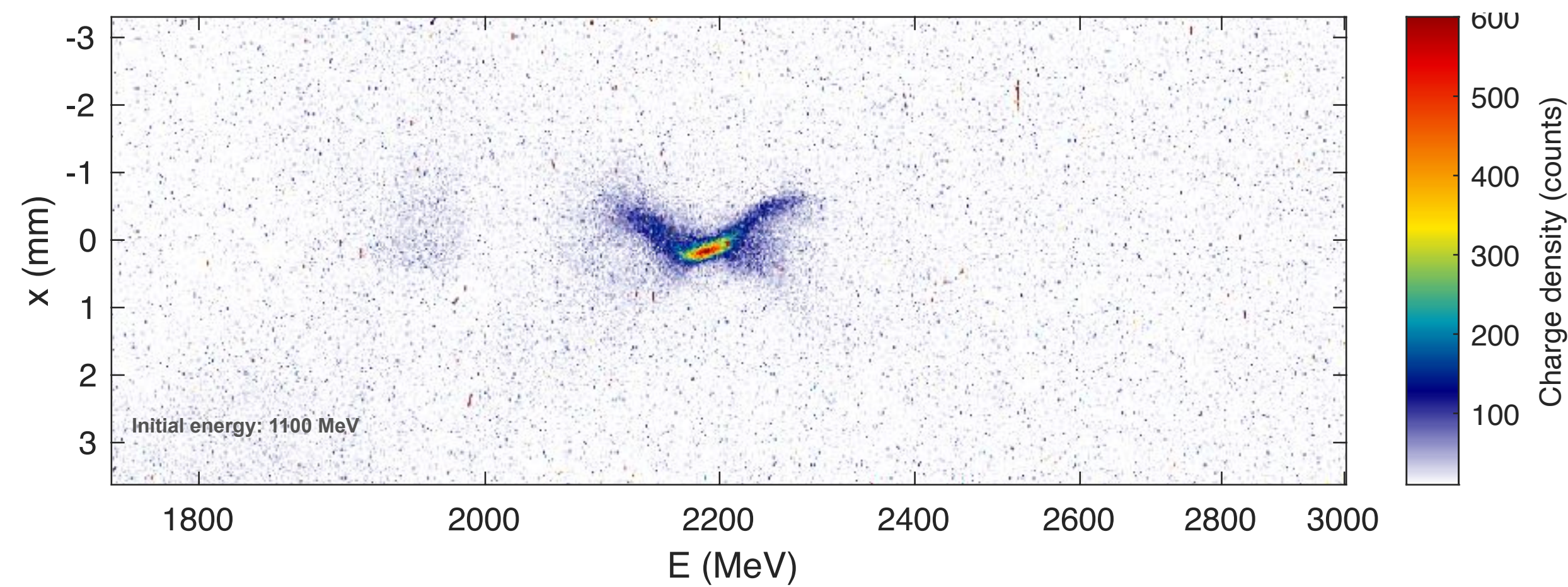
Plasma accelerator essentials — demonstrating 6 GV/m field strength



Energy extraction → plasma beam dump (+ efficiency)



Energy doubling to 2.2 GeV → plasma booster



Optimal beam loading enables uniform and efficient acceleration

- > *Problem 1:* Compared to RF cavities ($Q \sim 10^4\text{--}10^{10}$), the electric fields in a plasma decay very rapidly ($Q \sim 1\text{--}10$).
- > The energy needs to be extracted very rapidly
—ideally within the first oscillation.

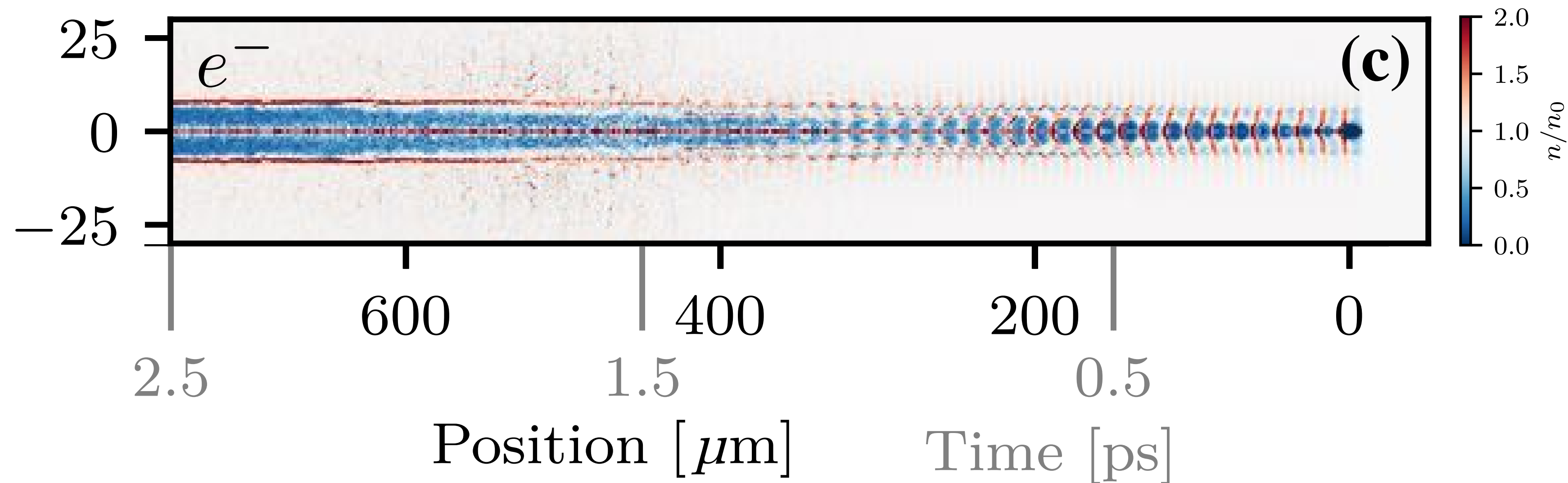


Image source: M. F. Gilljohann *et al.*, Phys. Rev. X **9**, 011046 (2019)

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 - > The energy needs to be extracted very rapidly —ideally within the first oscillation.
 - > *Solution:* Beam loading
The trailing-bunch wakefield “destructively interferes” with the driver wakefield—extracting energy.
- > *Problem 2:* to extract a large fraction of the energy, the beam will cover a large range of phases (~ 90 degrees or more).
 - > Large energy spread is induced.
 - > *Not (easily) possible:*
 Dechirping

R. D'Arcy *et al.*,
 PRL **122**, 034801 (2019)

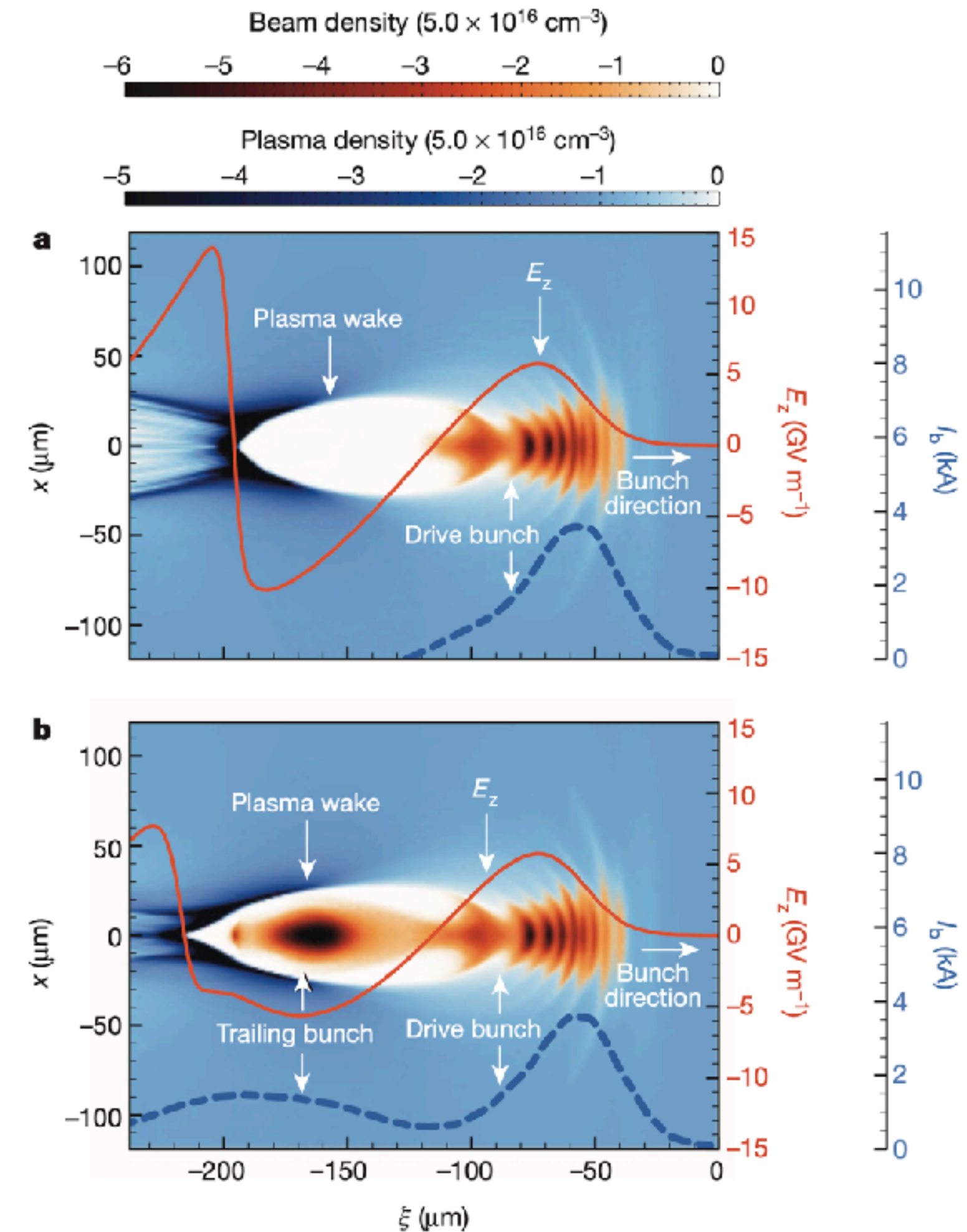
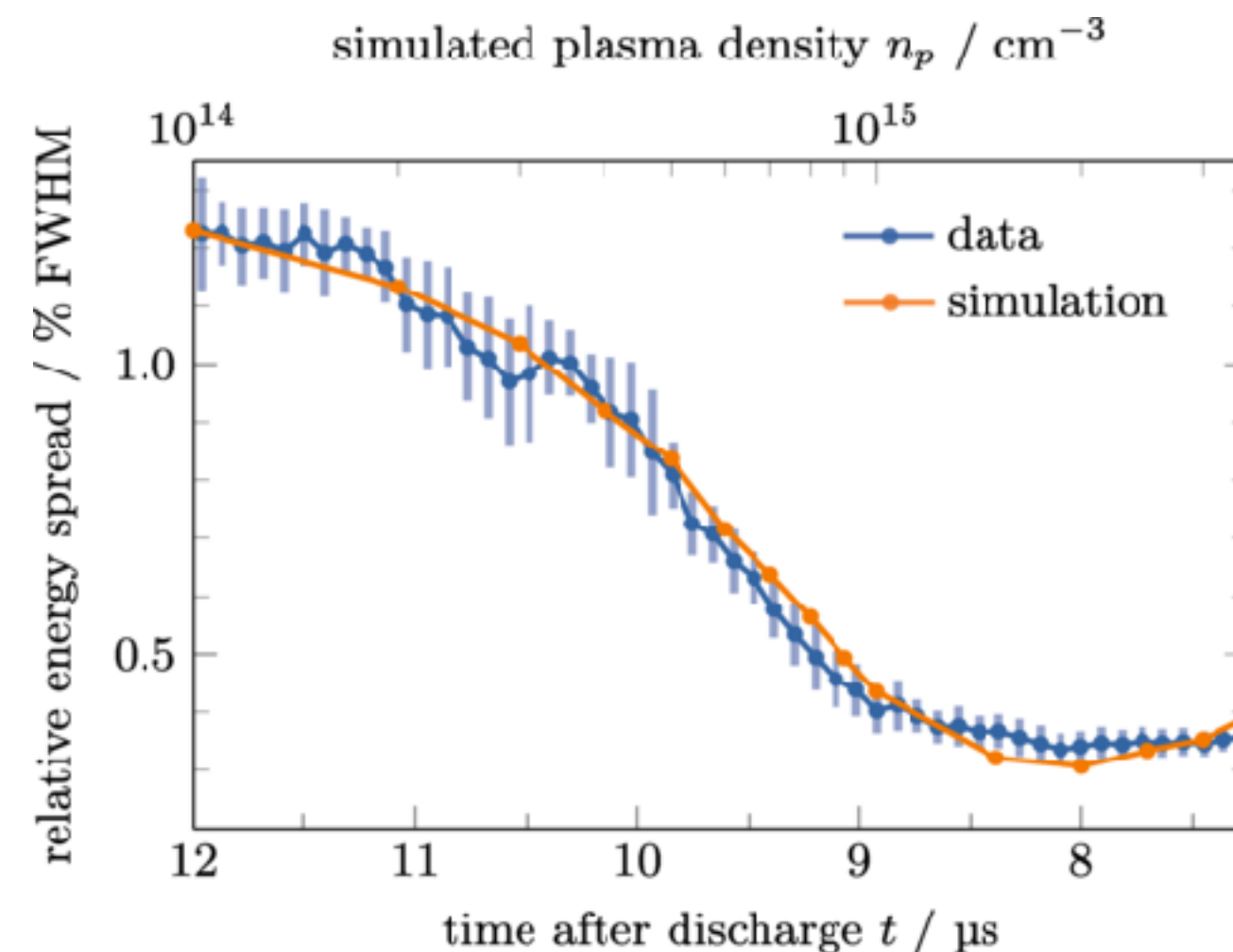


Image credit: M. Litos *et al.*, Nature **515**, 92 (2014)

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 - > Large energy spread is induced.
 - > *Solution:* Optimal beam loading
The current profile of the trailing bunch is *precisely tailored* to exactly flatten the wakefield.
- > This requires extremely precise control of the current profile.
 - > **FLASHForward provides the tools to do that.**

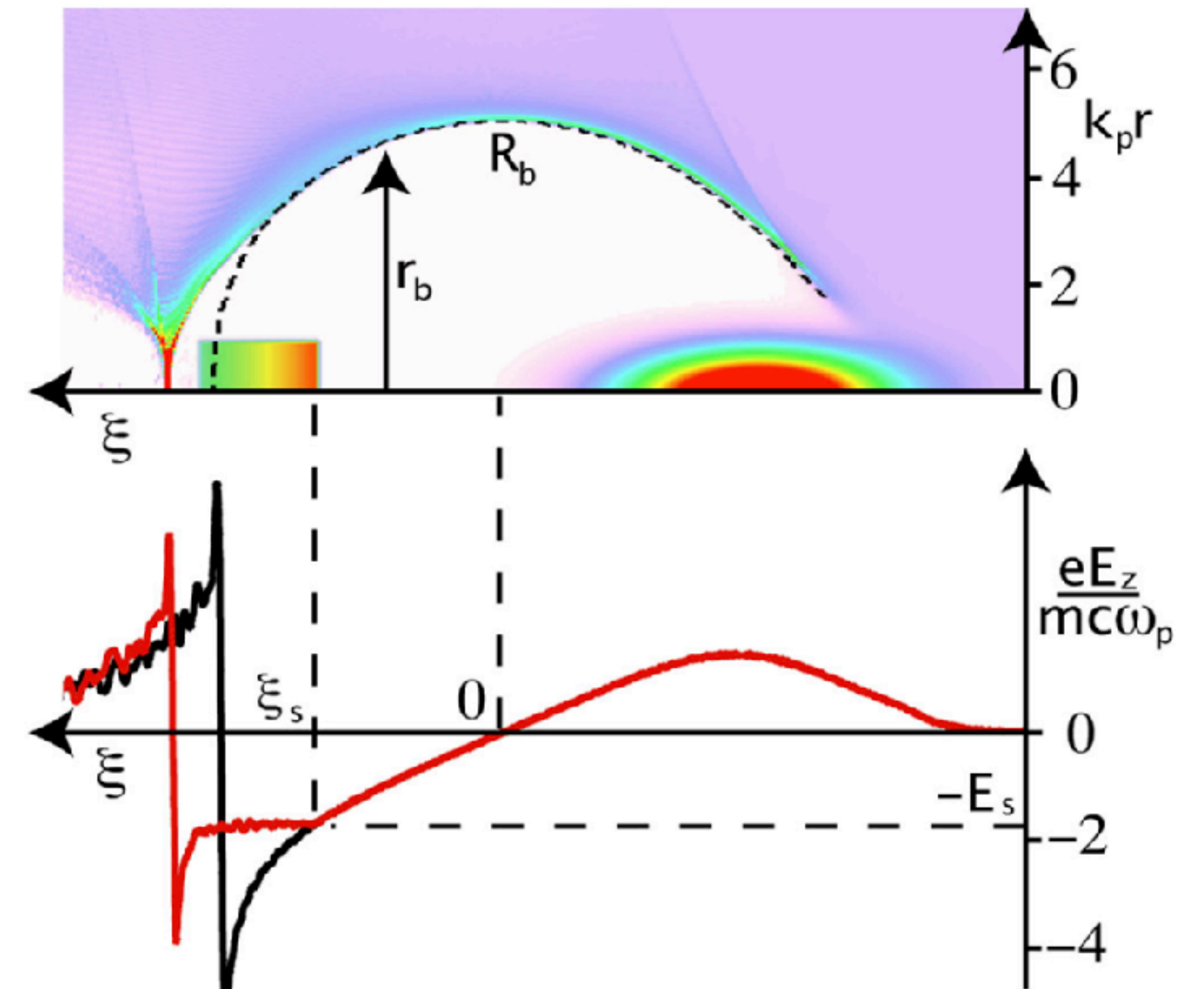


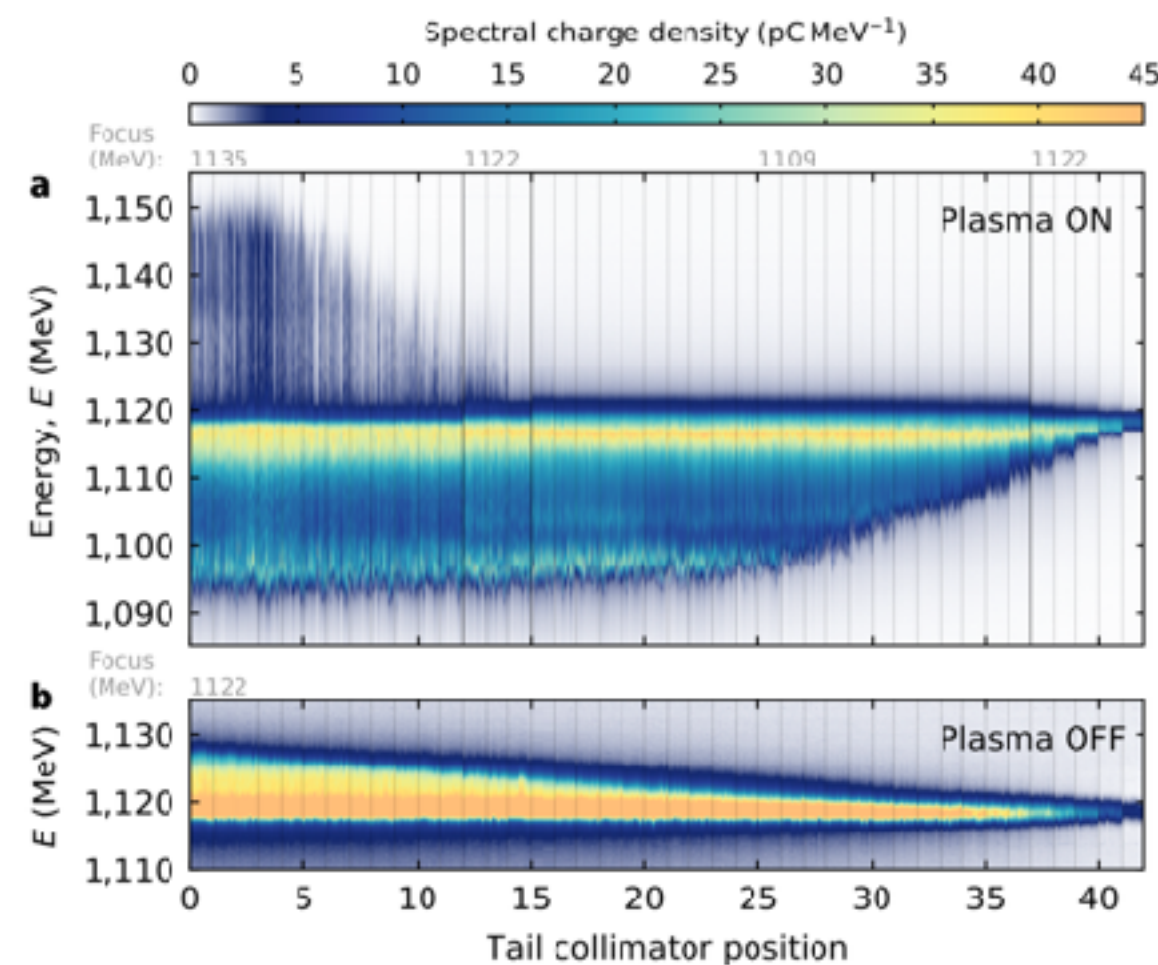
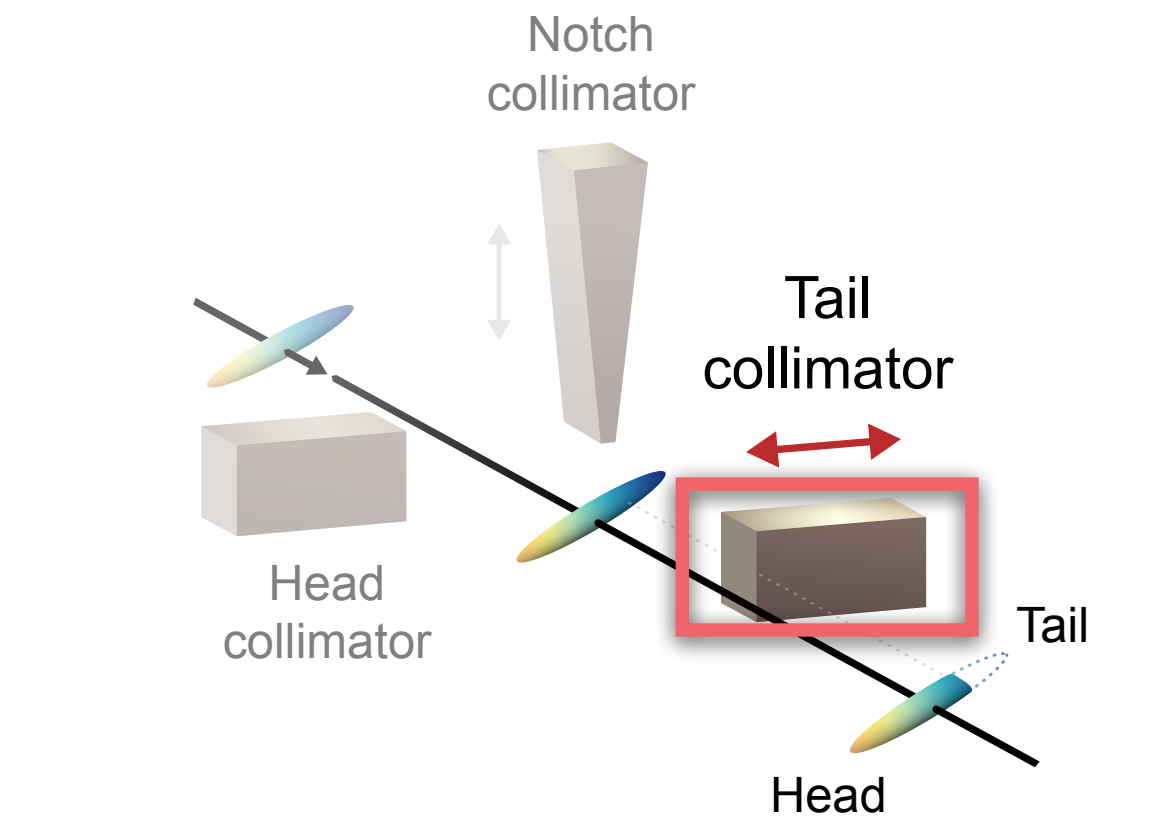
Image credit: M. Tzoufras *et al.*, Phys. Rev. Lett. **101**, 145002 (2008)

High-resolution plasma wakefield sampling demonstrated

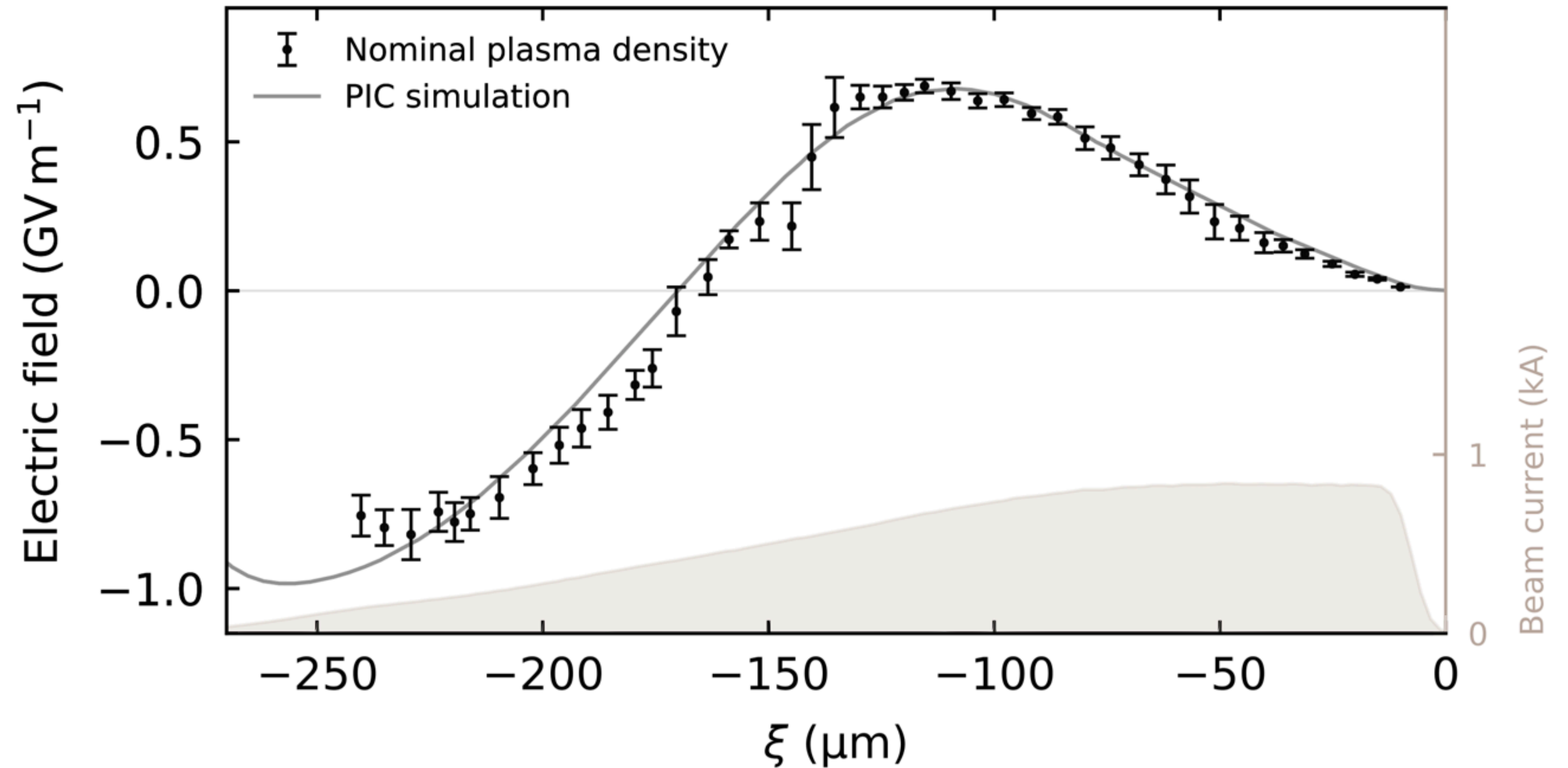
Opens a pathway to targeted and precise field manipulation

➤ Beam itself acts as a probe

→ measures in-situ (under actual operation conditions) the effective field acting on beam with μm / fs resolution



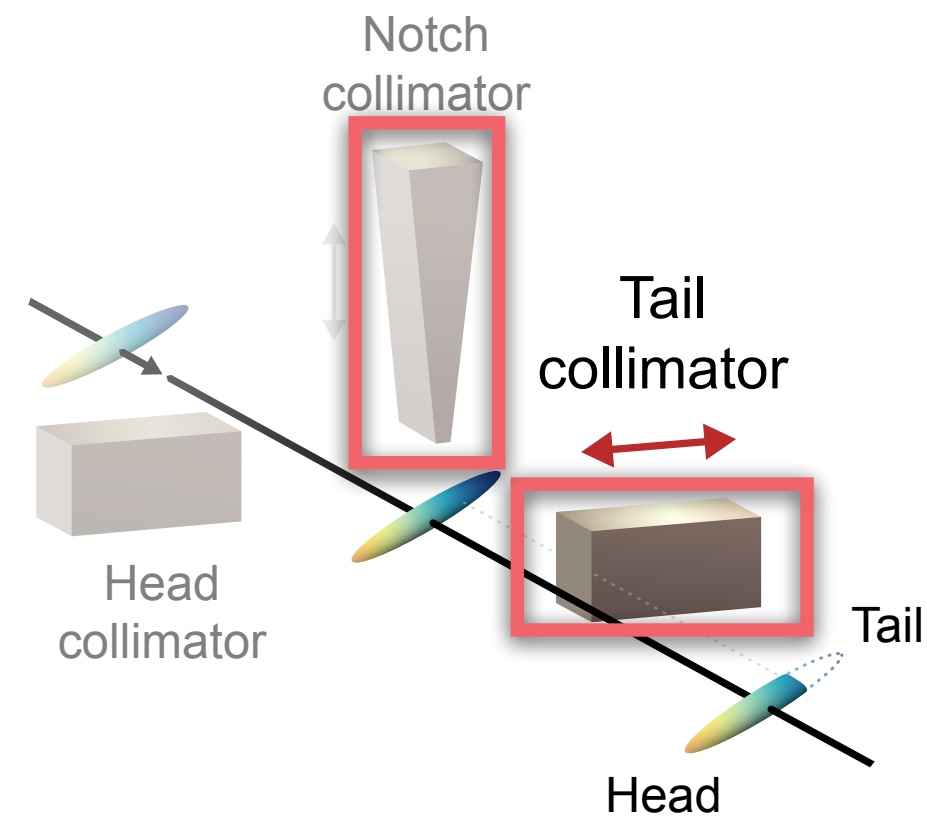
S. Schröder *et al.*, Nat. Commun. **11**, 5984 (2020)



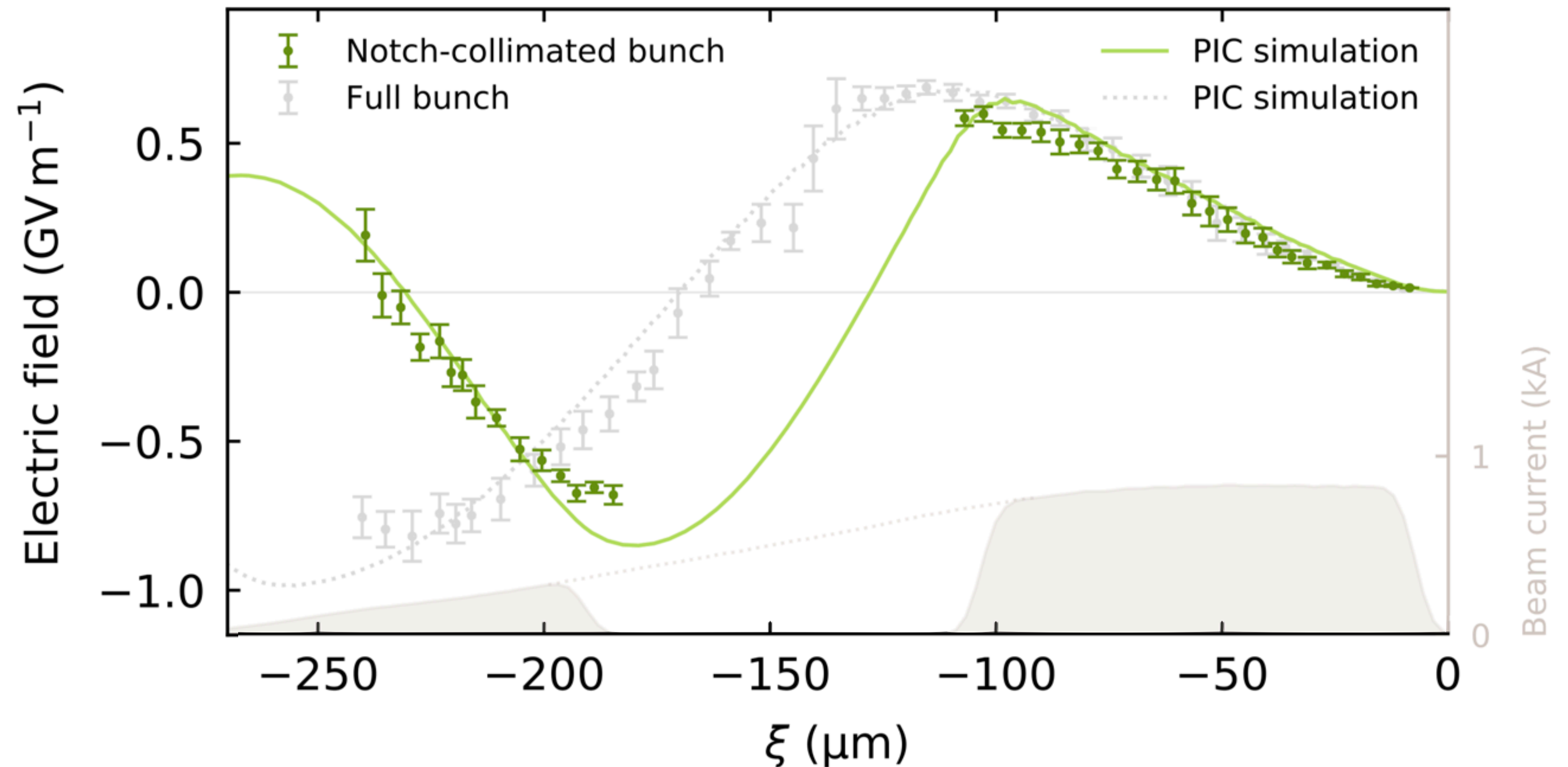
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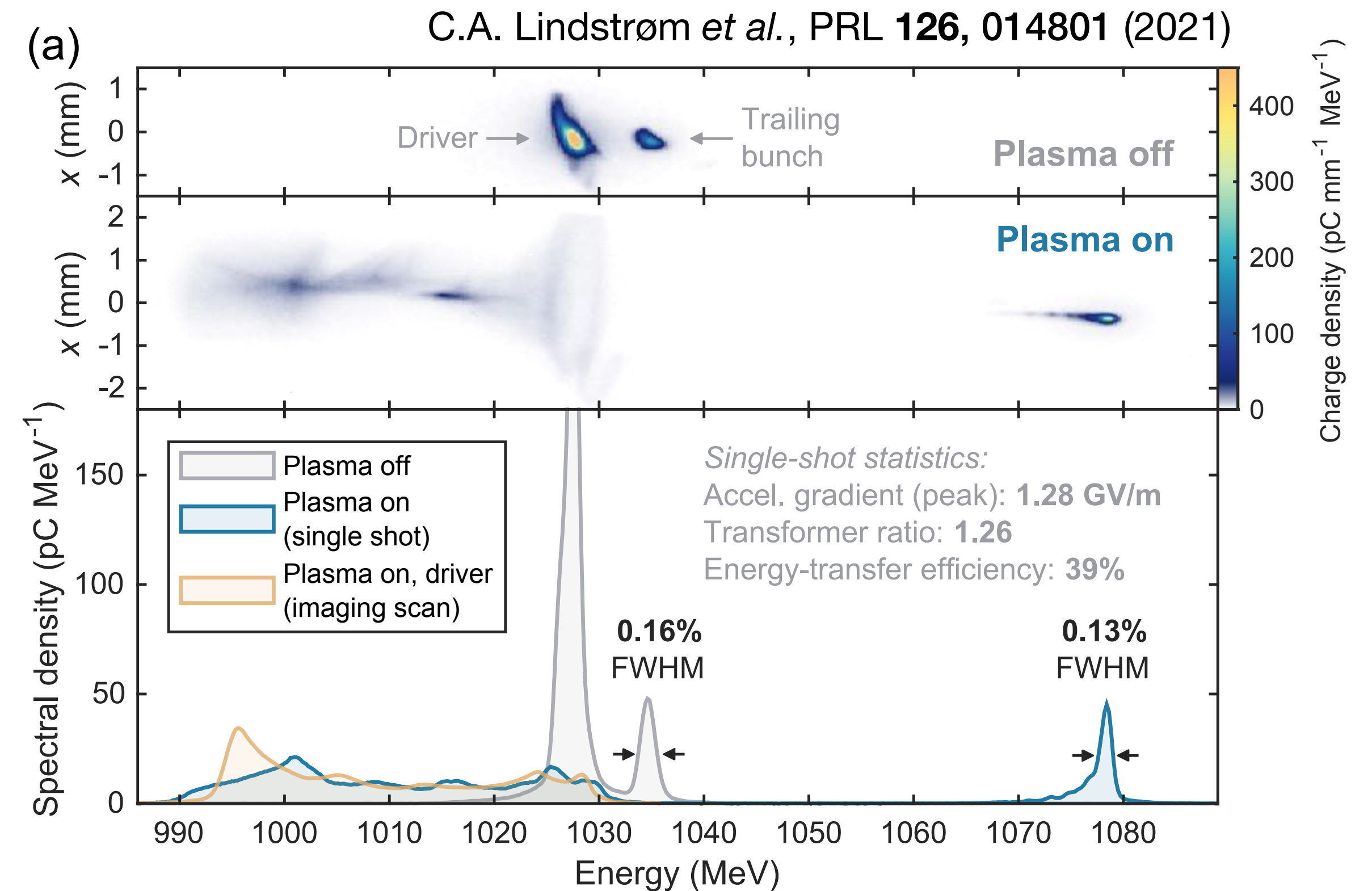
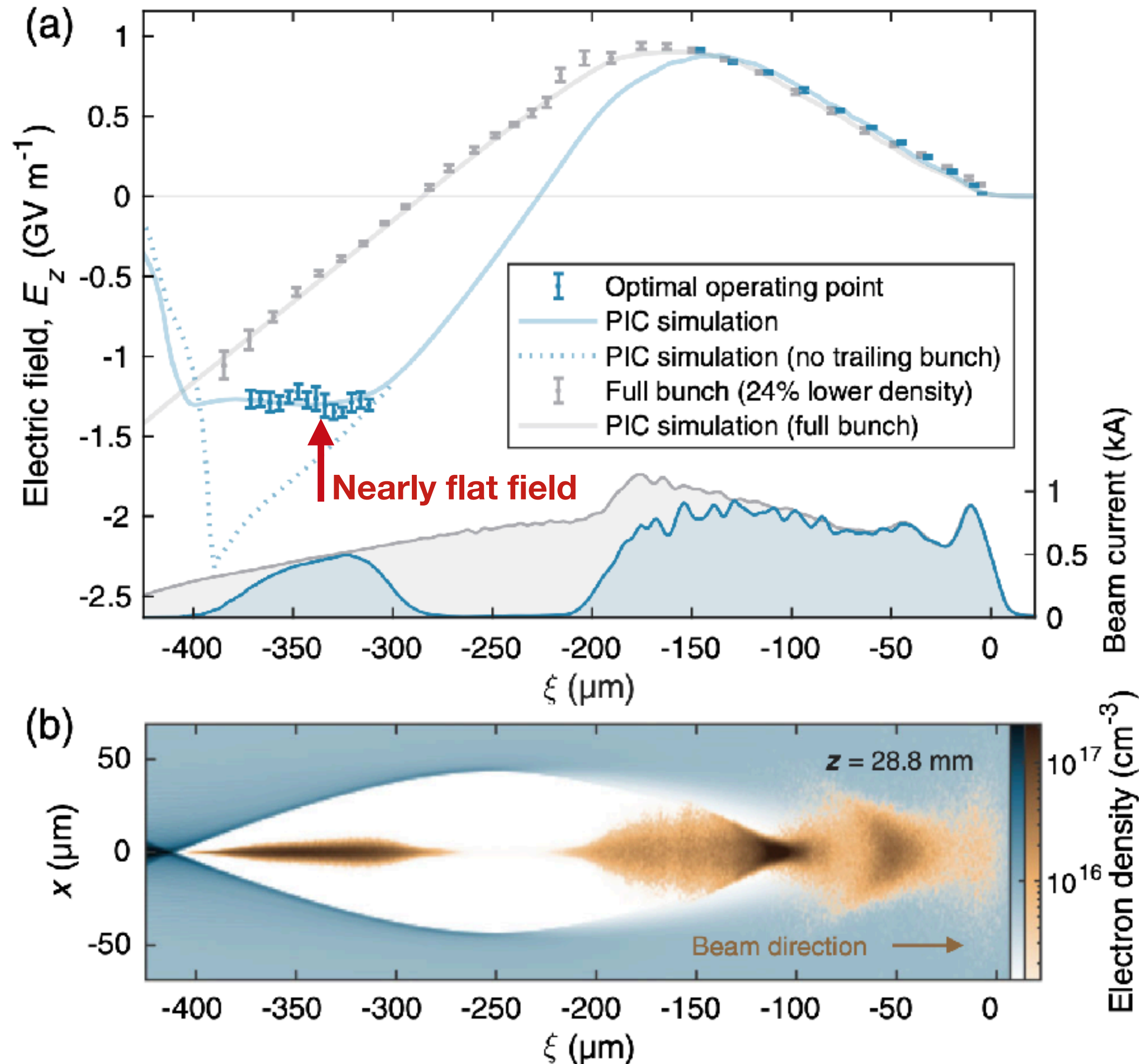


S. Schröder *et al.*, Nat. Commun. **11**, 5984 (2020)



Loading the wakefield and beam shaping flattens the gradient

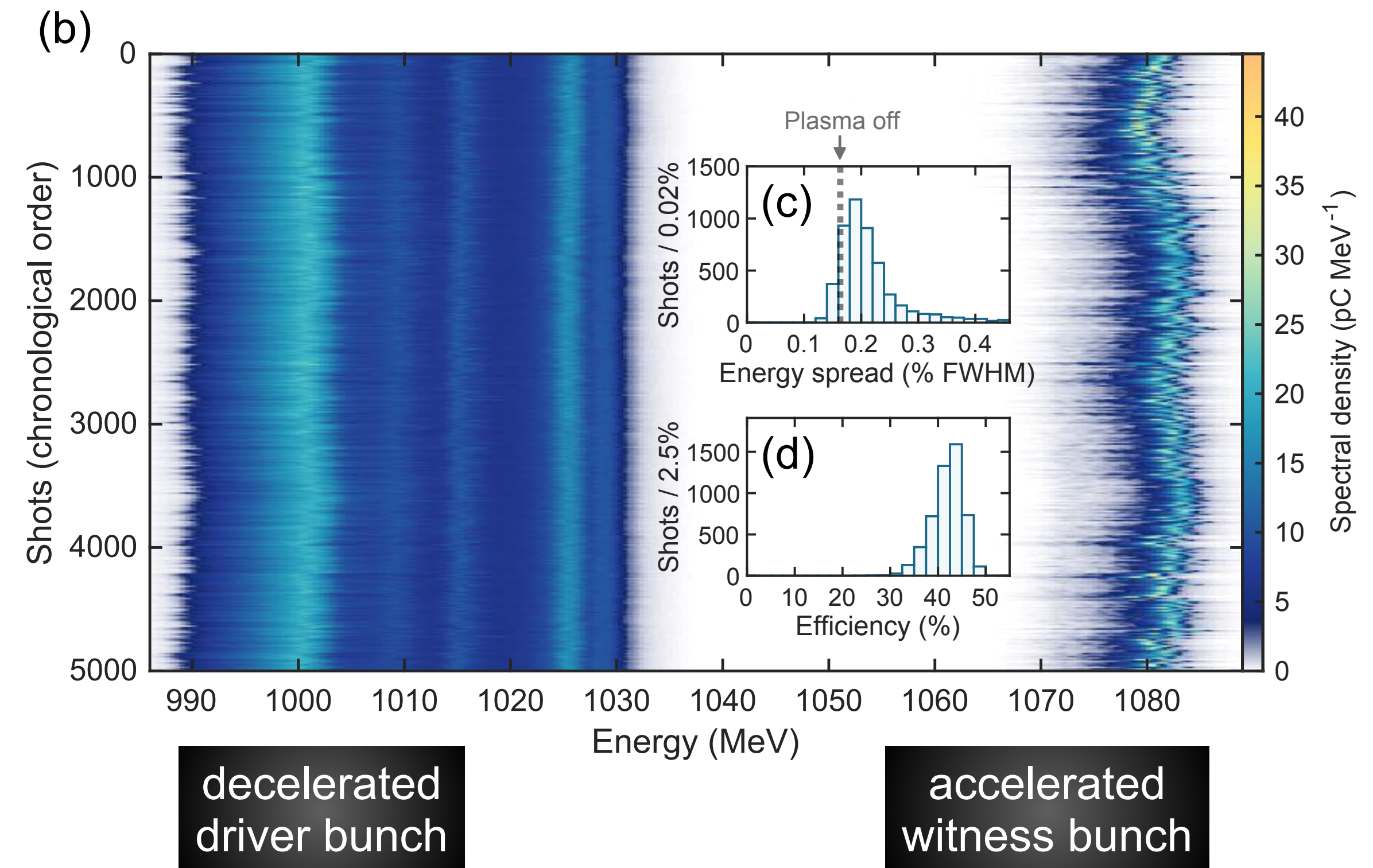
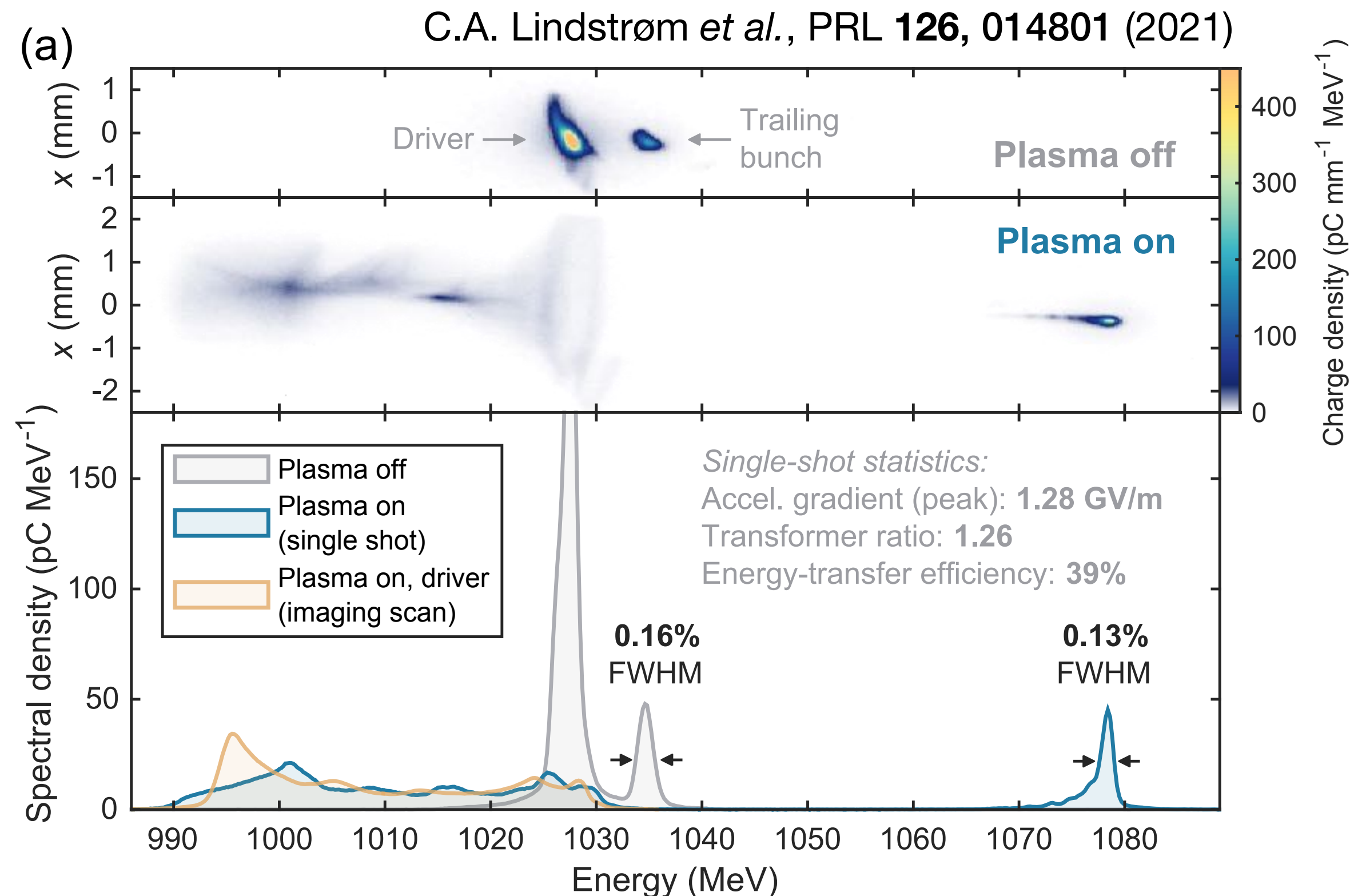
Direct visualization of electric-field control by wakefield sampling



- Accelerating gradient of 1.3 GV/m
- Energy gain 45 MeV (over 3.5 cm distance) of 100 pC witness, with energy spread of 1.4 MeV FWHM and **no charge loss**
- **Few-percent-level wakefield flattening demonstrated**

High-quality, efficient acceleration for sustainable applications

Beam-loading facilitates 42% energy-transfer efficiency, 0.2% energy spread with full charge coupling

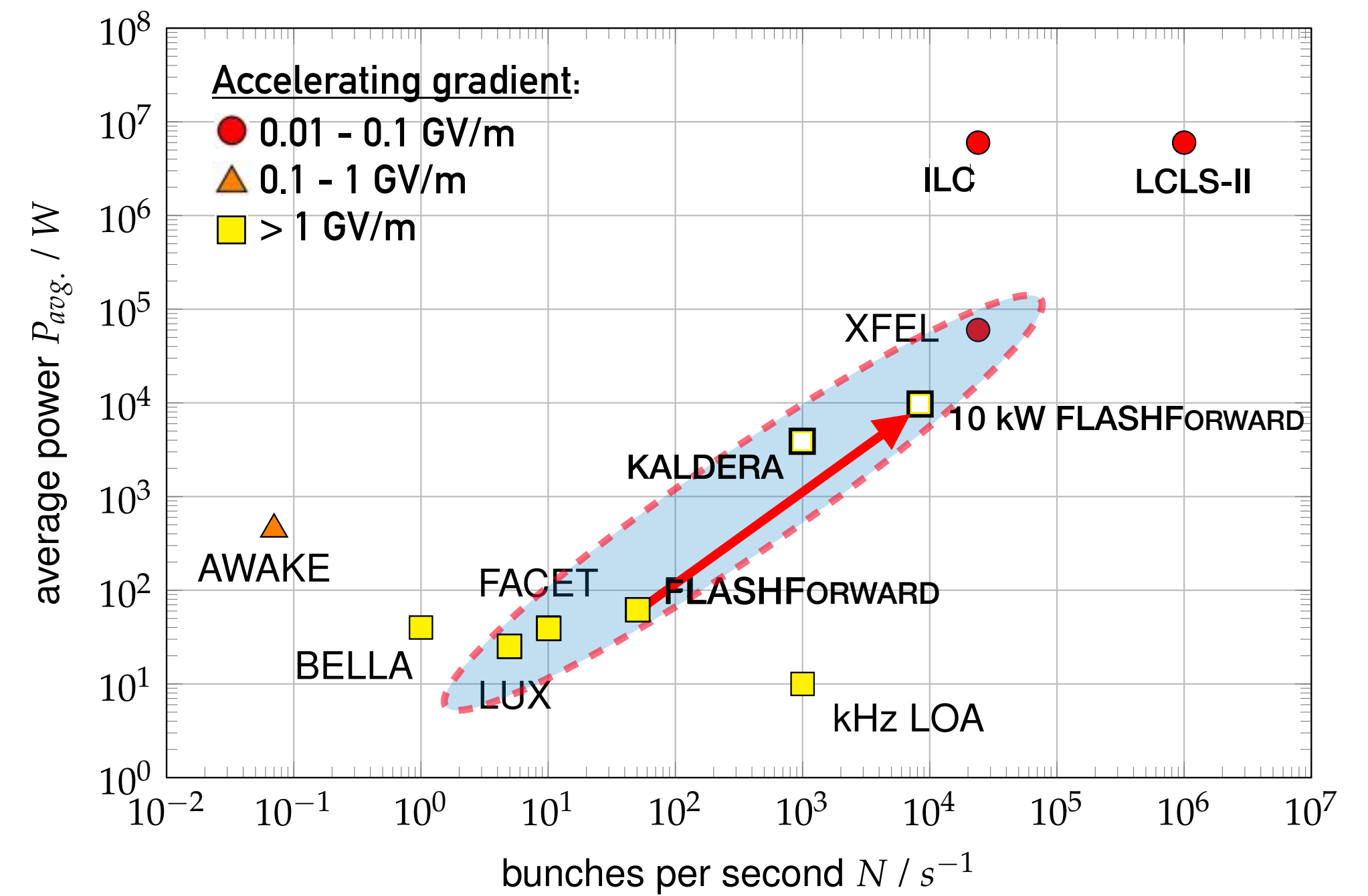
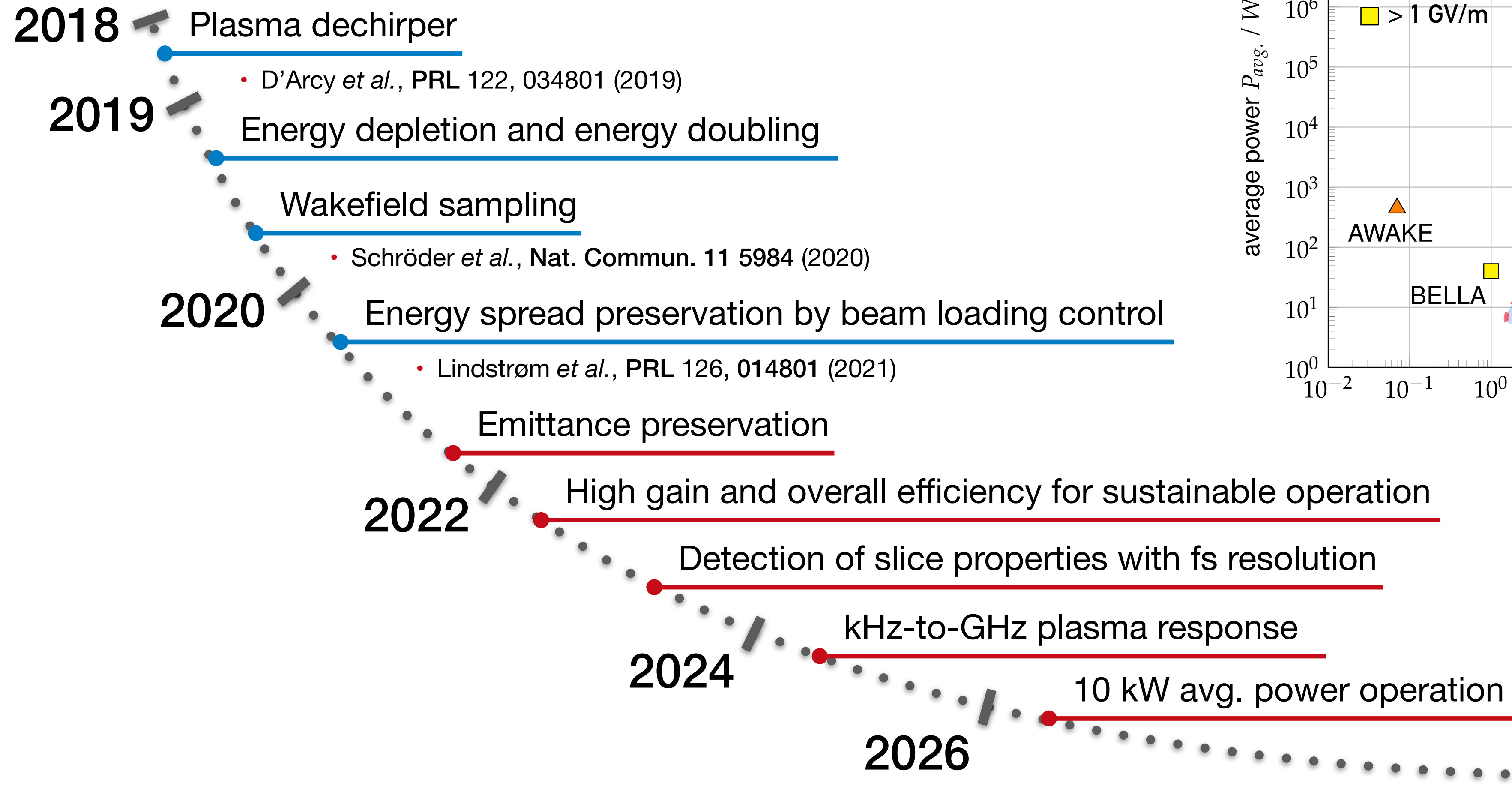


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- **0.2% energy spread (input 0.16%)**
(improvement by factor 10 over state-of-the-art)
- **(42±4)% energy transfer efficiency**
(improvement by factor 3 over state-of-the-art)

FLASHFORWARD ▶▶ roadmap aims at 10 kW with high beam quality

Plan covers major plasma accelerator challenges

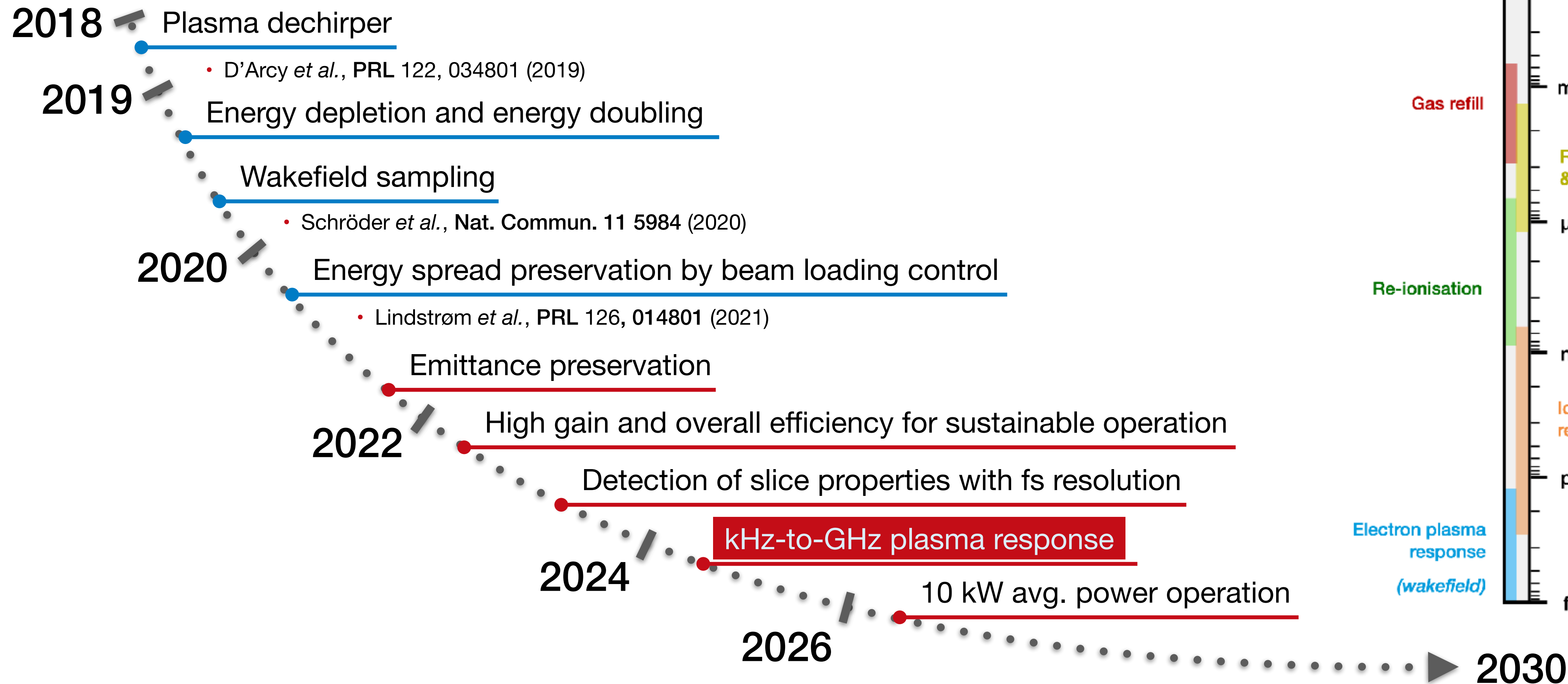


10 kW stage with 50% efficiency & beam quality conservation

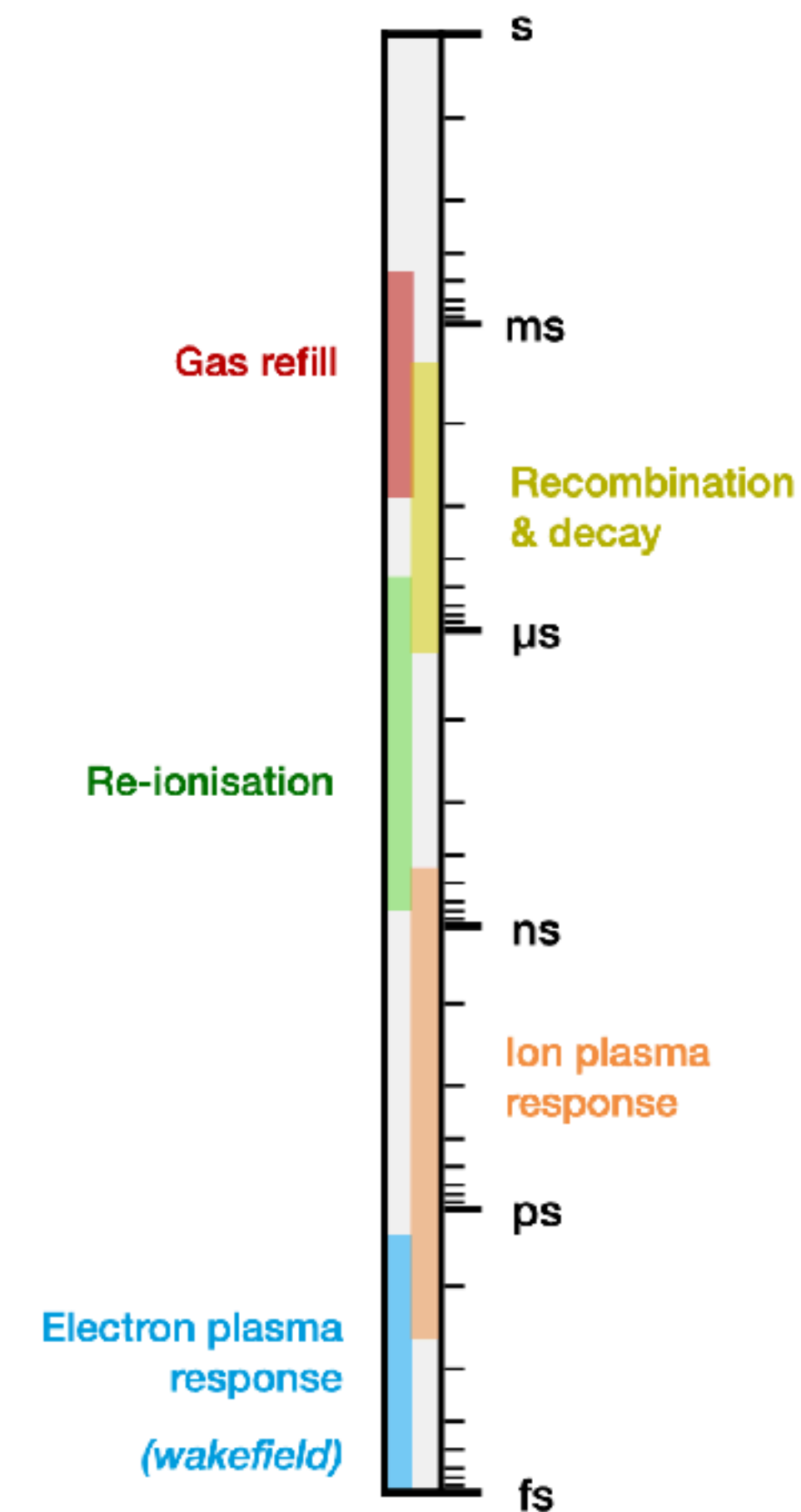
→ **FLASH:** increase FEL energies, access oxygen K-edge at 2.33 nm wavelength

FLASHFORWARD▶▶ roadmap aims at 10 kW with high beam quality

Plan covers major plasma accelerator challenges



Plasma (wakefield) timeline

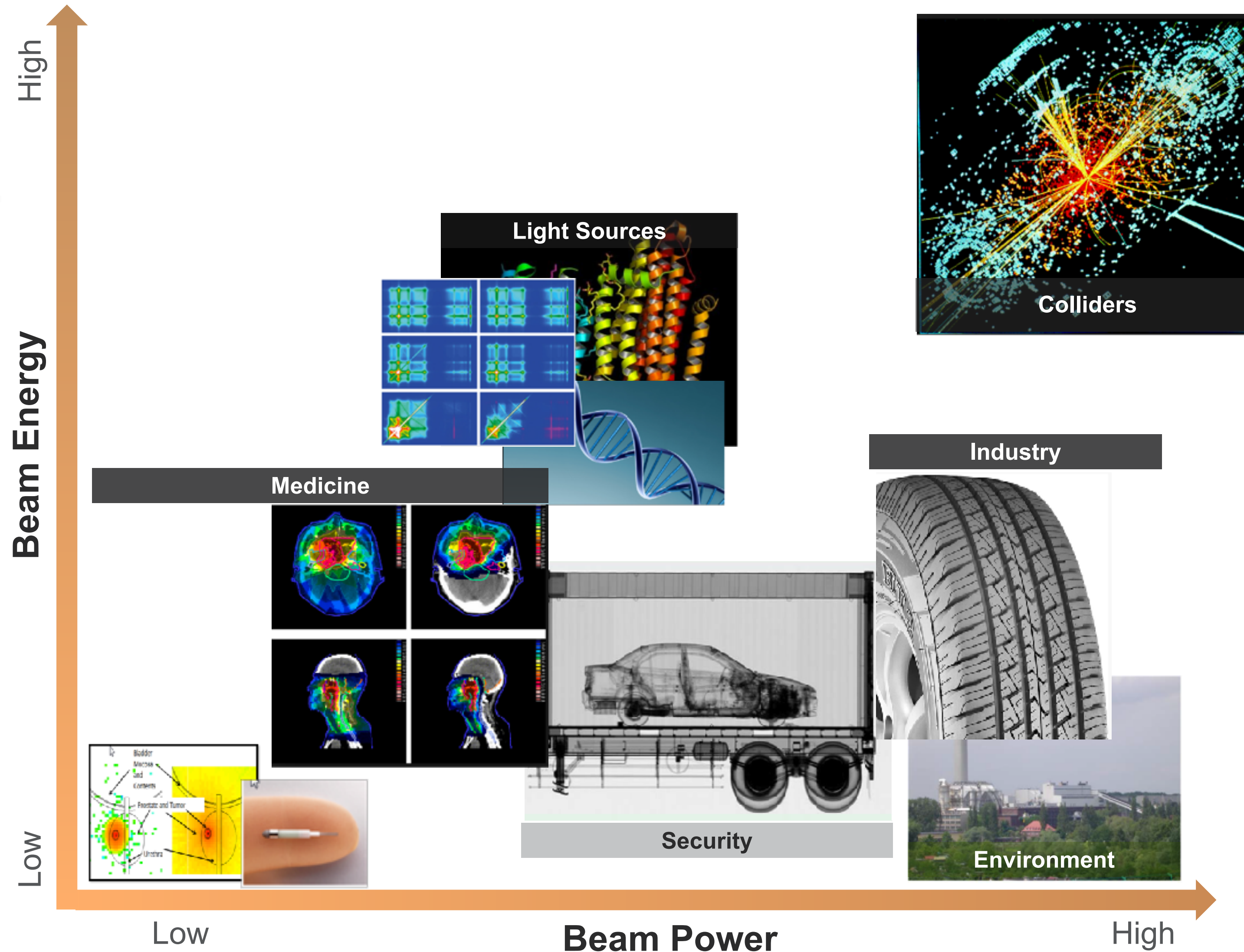


High-power plasma accelerators unlock new areas of application

Miniaturize current concepts and change the paradigm: bring the machine to the problem

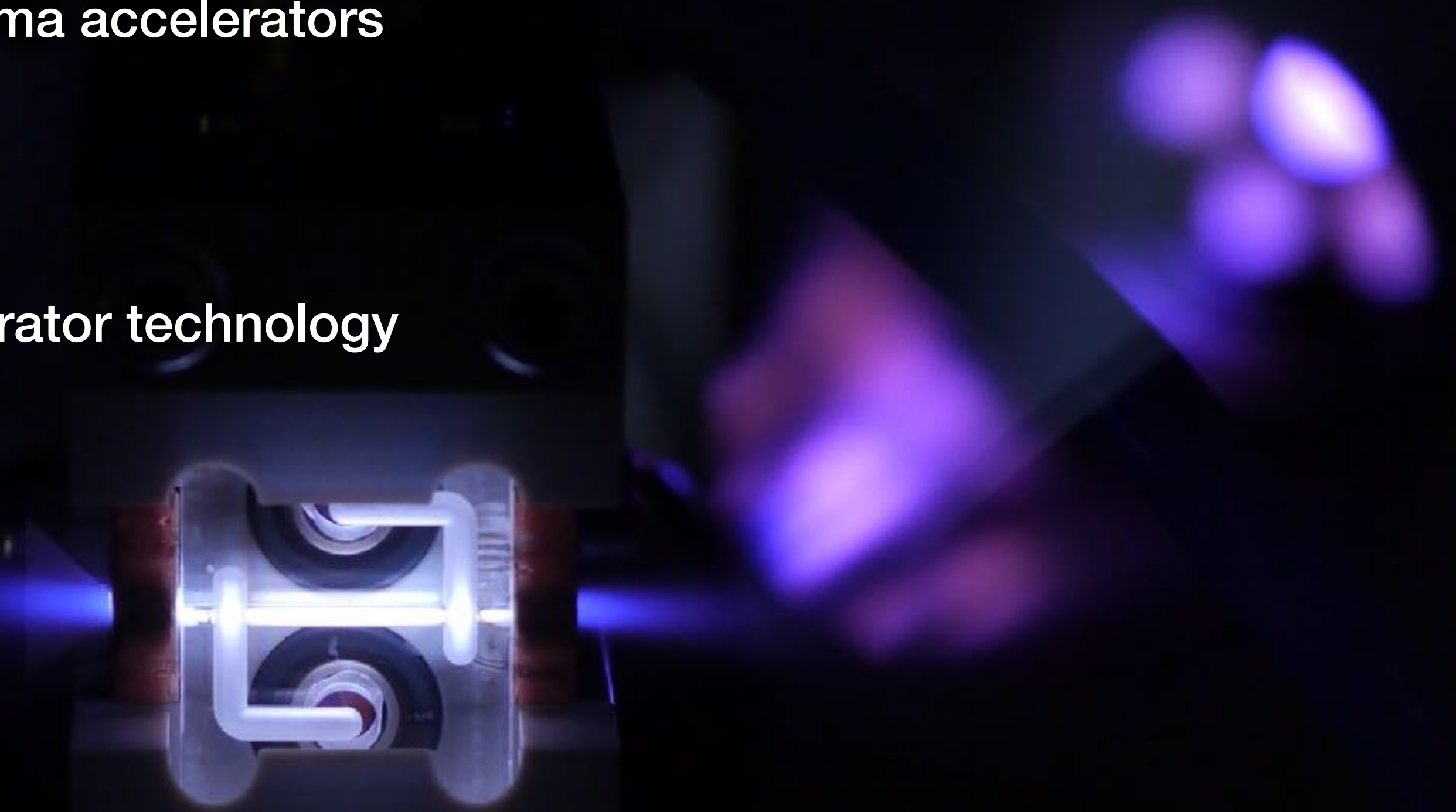


Cole et al, PNAS 115 (2018)



Summary

- Accelerators are at the heart of photon science and particle physics experiments, but are large installations
- Plasma wakefield technology offers a promising path to compact accelerators with > 1 GV/m fields
- Two alternative driver technologies: laser- and beam-excited plasma wakes
- **Common goal:**
 - plasma accelerator research → usable plasma accelerators
- **Hope:** miniaturization of accelerators leads to
 - significant cost reduction
 - widespread proliferation of compact accelerator technology
 - beams with new and extreme properties
- Plasmas may have a revolutionary influence on accelerator applications and society



Thank you for your attention!

