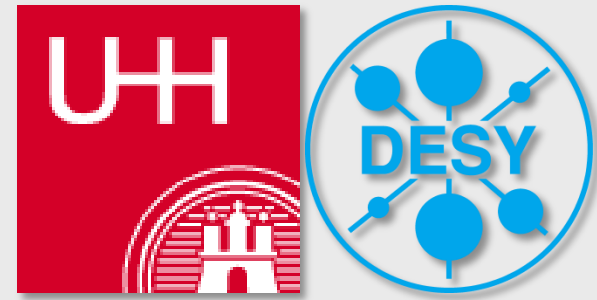


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for the Gfitter group*



Advanced Statistics WS 2010
Göttingen
20th October 2010



Gfitter – Global Electroweak Fits in the Standard Model and Beyond

main paper **Eur. Phys. J. C 60, 543 (2009)**

<http://cern.ch/Gfitter>

*) M. Baak (CERN), M. G. (Univ. Hamburg, DESY), J. Haller (Univ. Göttingen), A. Höcker (CERN),
D. Ludwig (Univ. Hamburg, DESY), K. Mönig (DESY), M. Schott (CERN) J. Stelzer (DESY/Michigan)

A Generic Fitter Project for HEP Model Testing

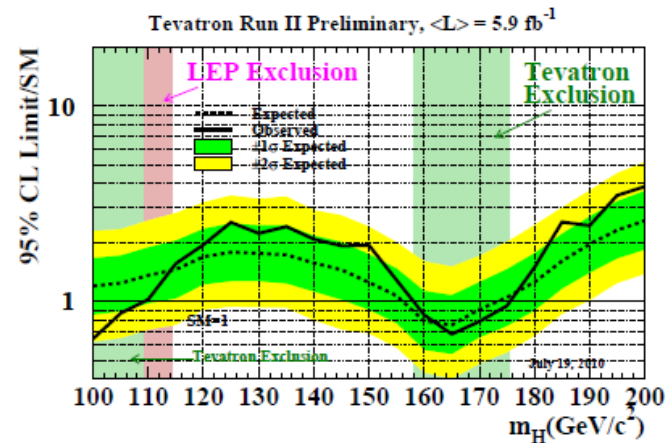
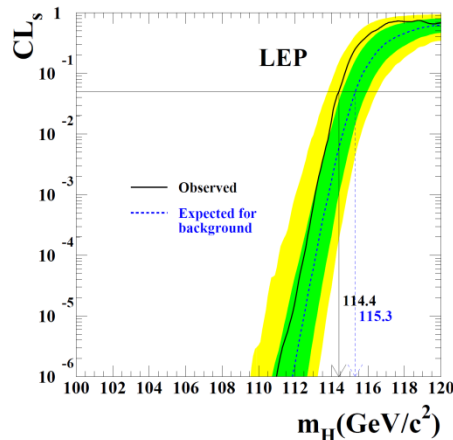
- SUSY / BSM Fit Working Group of the Helmholtz Alliance
- Goal: provide a state-of-the-art model testing tool for the LHC/ILC era
- modular, object-oriented C++, relying on ROOT, XML, python
- core package for data handling, fitting and statistics tools
 - various fitting tools: Minuit, Genetic Algorithm and Simulated Annealing
 - full statistics analysis: parameter scans, p -values, MC analyses, goodness-of-fit tests, ...
 - coherent treatment of statistical, systematic errors, and correlations
 - theo. uncertainties included in χ^2 with flat likelihood in allowed ranges
- physics plug-in packages
 - Library for the Standard Model fit to the electroweak precision data → this talk
 - Library for SM extensions via the oblique parameters → this talk
 - Library for the super-symmetric extension of the SM
 - Library for the 2HDM extension of the SM



A Gfitter Package for the Global Electroweak Fit

- state-of-the art calculations (OMS scheme); in particular:
 - M_W and $\sin^2\theta_{\text{eff}}^f$: full two-loop + leading beyond-two-loop correction
[M. Awramik et al., Phys. Rev D69, 053006 (2004)][M. Awramik et al., JHEP 11, 048 (2006), (M. Awramik et al., Nucl.Phys.B813:174-187 (2009)]
 - **radiator functions**: N³LO of the massless QCD Adler function
[P.A. Baikov et al., Phys. Rev. Lett. 101 (2008) 012022]
- theoretical uncertainties: M_W ($\delta M_W=4-6$ GeV), $\sin^2\theta_{\text{eff}}^l$ ($\delta\sin^2\theta_{\text{eff}}^l=4.7 \cdot 10^{-5}$), truncation of higher QCD orders
 - included in χ^2 with flat likelihood \rightarrow vary within uncertainties without contribution to χ^2

- usage of latest experimental results:
 - **Z-pole observables:** LEP/SLD results [ADLO+SLD, Phys. Rept. 427, 257 (2006)]
 - **M_W and Γ_W :** LEP/Tevatron $M_W = 80.399 \pm 0.023$ GeV [ADLO, hep-ex/0612034] [CDF, Phys. Lett. 100, 071801 (2008)] [CDF&D0, Phys. Rev. D 70, 092008 (2004)] [CDF&D0, arXiv:0908.1374v1]
 - **m_{top} :** $m_{\text{top}} = 173.3 \pm 1.1$ GeV [D0&CDF, arXiv:1007.3178]
 - **$\Delta\alpha_{\text{had}}^{(5)}(M_Z^2)$:** including α_s dependency [Davier, Hoecker, Malaescu, Zhang]
 - **\bar{m}_c, \bar{m}_b :** world averages [PDG, J. Phys. G33,1 (2006)]
- floating fit parameters: $M_Z, M_H, m_t, \Delta\alpha_{\text{had}}^{(5)}(M_Z^2), \alpha_s(M_Z^2), \bar{m}_c, \bar{m}_b$
- fits are performed in two versions:
 - all data except results from direct Higgs searches
 - all data including results from direct Higgs searches at LEP [ADLO: Phys. Lett. B565, 61 (2003)] and **Tevatron** [CDF+D0: arXiv: 1007.4587]



$$\alpha(s) = \frac{\alpha(0)}{1 - \Delta\alpha(s)}, \quad \Delta\alpha(s) = \Delta\alpha_{\text{lep}}(s) + \Delta\alpha_{\text{had}}^{(5)}(s) + \Delta\alpha_{\text{top}}(s)$$

- leptonic (top) contribution to running of α precisely known (small)
- new value for hadronic contribution [Davier, Hoecker, Malaescu, Zhang]
 - several improvements: new $\pi\pi$ cross-section data from KLOE, all available multi-hadron data from BABAR, reevaluation of the continuum contributions from perturbative QCD at four loops, ...

$$\Delta\alpha_{\text{had}}^{(5)}(M_Z^2) = (274.2 \pm 1.0) \cdot 10^{-4}$$

- error includes uncertainty of α_s ($0.37 \cdot 10^{-04}$), which is a free fit parameter and has therefore no uncertainty in a certain fit step
- ⇒ subtract α_s uncertainty from total error
- variation of α_s needs to be included in the central value (Gfitter rescaling mechanism)

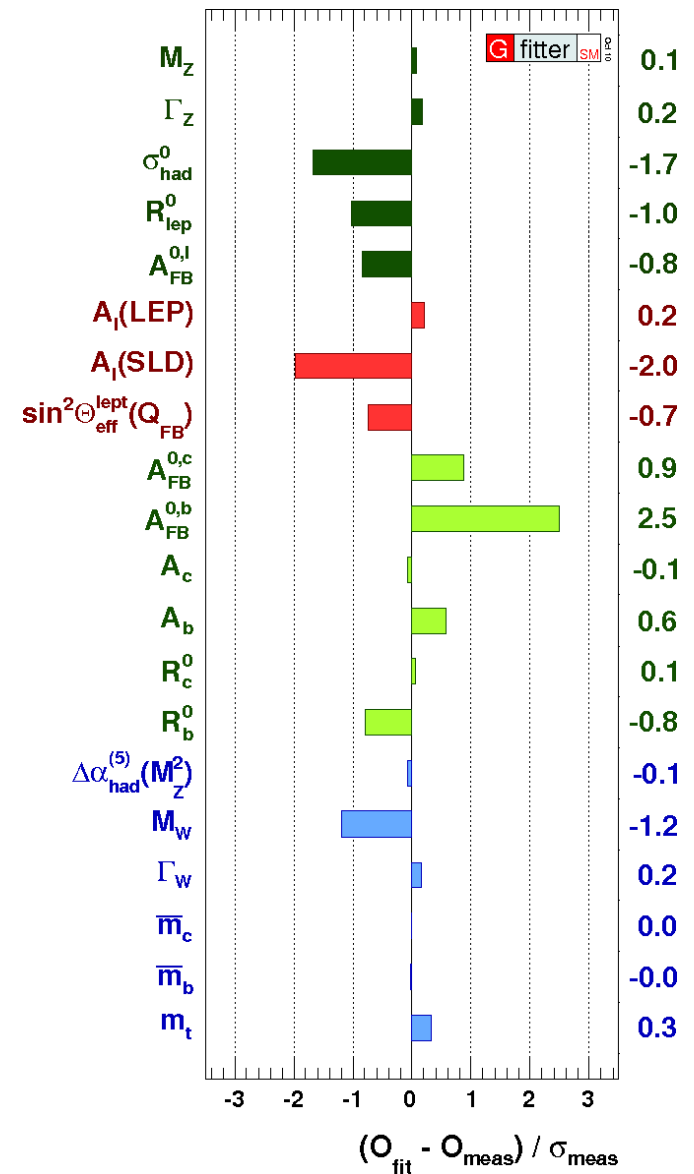
$$\Delta\alpha_{\text{had}}^{(5)}(M_Z^2) = \left(274.2 \pm 0.97 \pm 0.37 \cdot \left(\frac{\alpha_s(M_Z^2) - 0.1193}{0.0028} \right) \right) \cdot 10^{-4}$$

naïve p-value

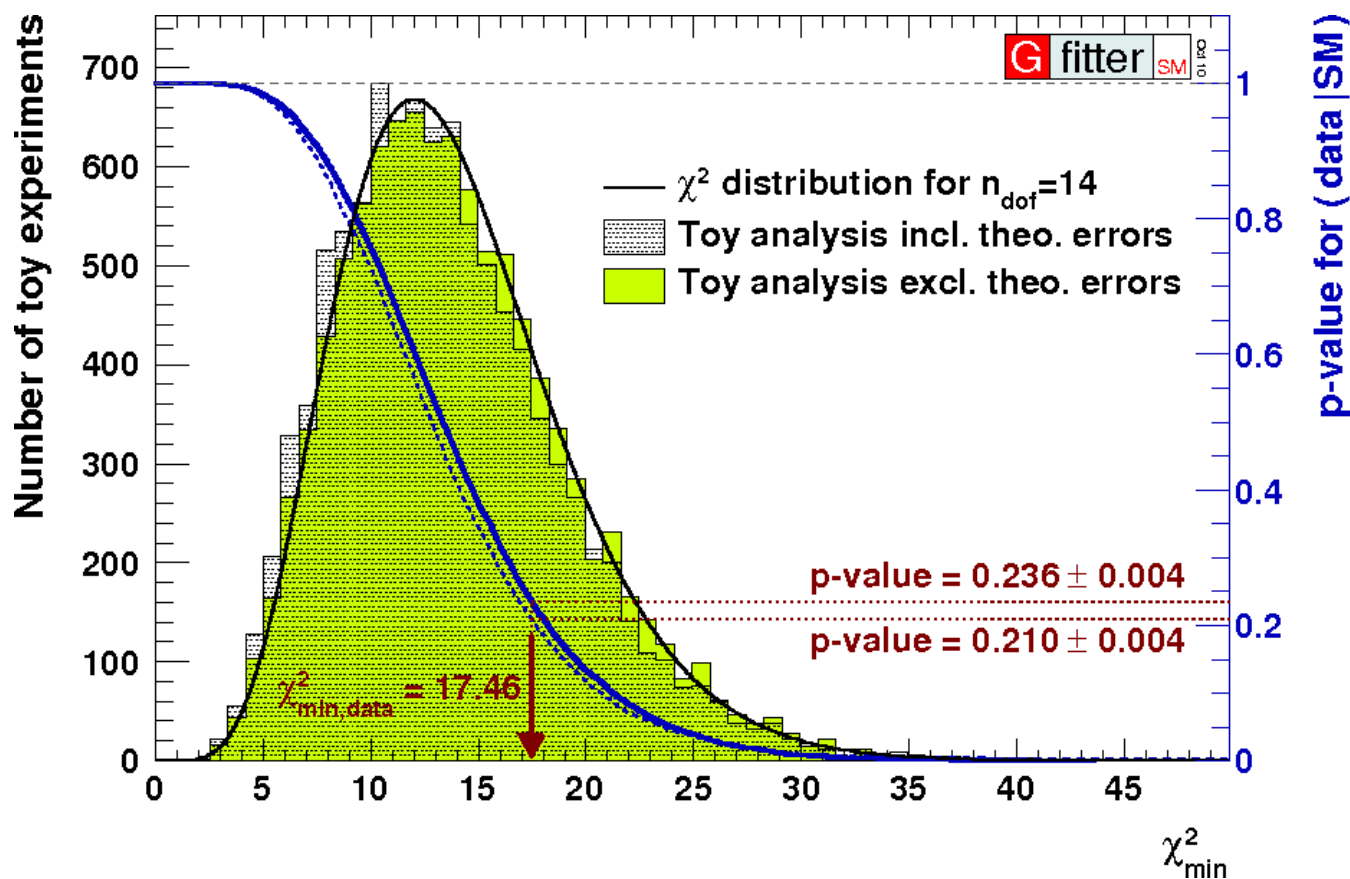
- w/o direct Higgs searches:
 $\chi^2_{\min} = 16.6 \rightarrow \text{Prob}(\chi^2_{\min}, 13) = 0.22$
- with direct Higgs searches:
 $\chi^2_{\min} = 17.5 \rightarrow \text{Prob}(\chi^2_{\min}, 14) = 0.23$

pull-values for the fit with Higgs searches (right figure \rightarrow)

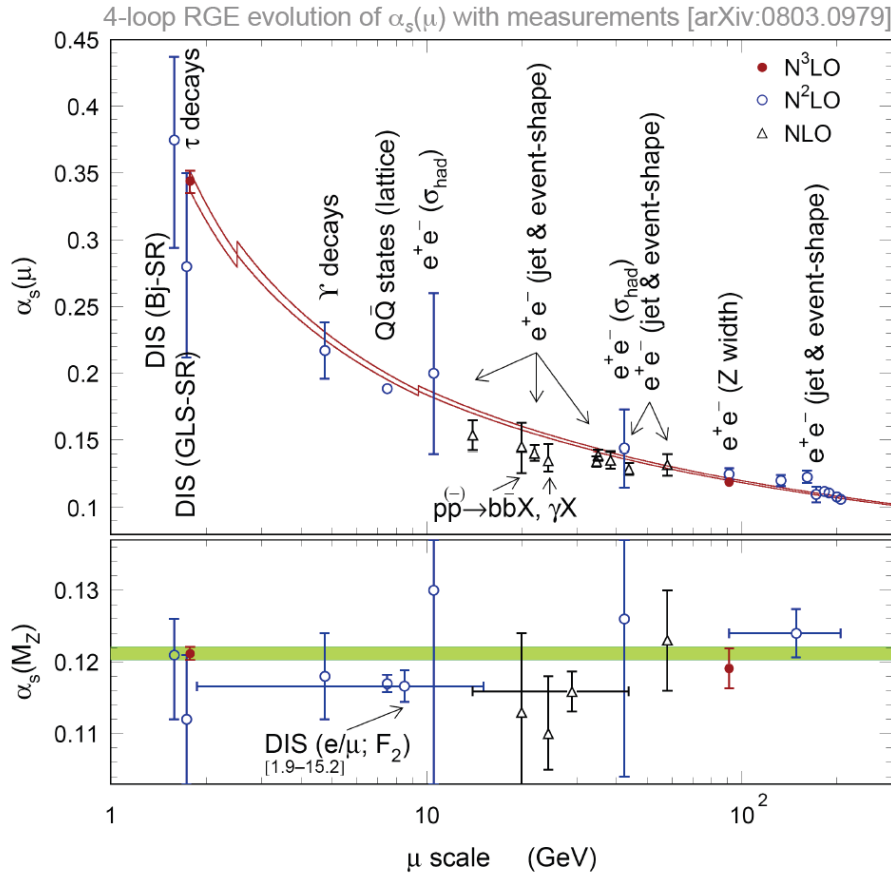
- FB asymmetry of bottom quarks
 \rightarrow largest contribution to χ^2
 - no value exceeds 3σ
 - small contributions from M_Z , $\Delta\alpha_{\text{had}}^{(5)}$, m_c , and m_b indicate that their input accuracies exceed fit requirements
- \Rightarrow no significant requirement for new physics



- p-value: probability for wrongly rejecting the SM
- p-value: probability for getting a $\chi^2_{\min, \text{toy}}$ larger than the $\chi^2_{\min, \text{data}}$ from data
- p-value for fit with Higgs searches $0.24 \pm 0.03 - 0.02_{\text{theo}}$



Determination of Strong Coupling



- R_1 observable most sensitive to α_s
- N^3LO (massless Adler function)
determination of α_s from complete fit:

$$\alpha_s(M_Z) = 0.1193 \pm 0.0028$$

$$\pm 0.0001$$

- first error experimental
- second error theoretical
[incl. variation of renorm. scale from $M_Z/2$ to $2M_Z$ and massless terms of order/beyond $a_s^5(M_Z)$ and massive terms of order/beyond $a_s^4(M_Z)$]
- excellent agreement with N^3LO result from hadronic τ decays [M. Davier et al., arXiv:0803.0979]

$$\alpha_s(M_Z) = 0.1212 \pm 0.0005_{\text{exp}}$$

$$\pm 0.0008_{\text{theo}}$$

$$\pm 0.0005_{\text{evol}}$$

w/o direct Higgs searches:

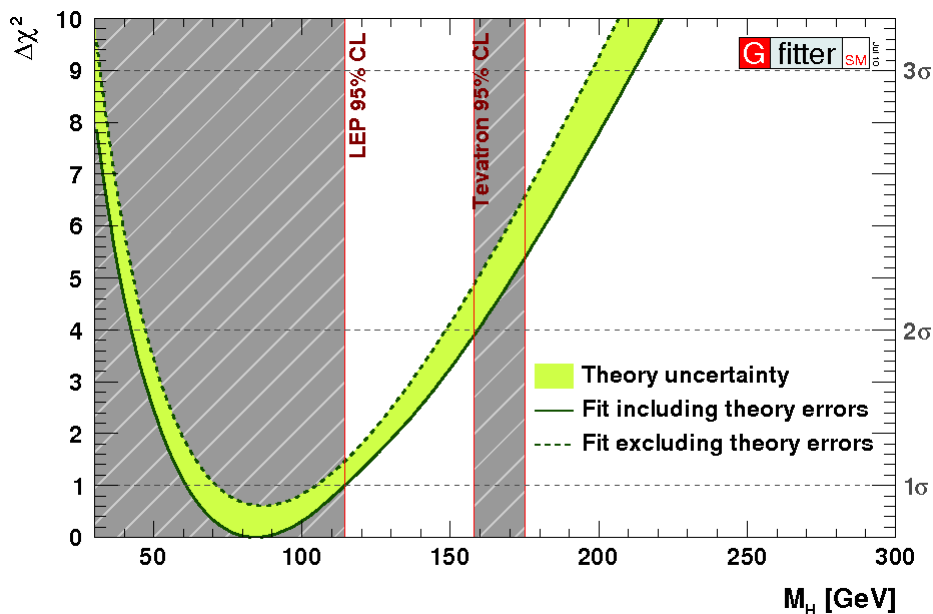
- value at minimum $\pm 1\sigma$:

$$M_H = 83^{+30}_{-23} \text{ GeV}$$

- 2σ interval: [42, 159] GeV

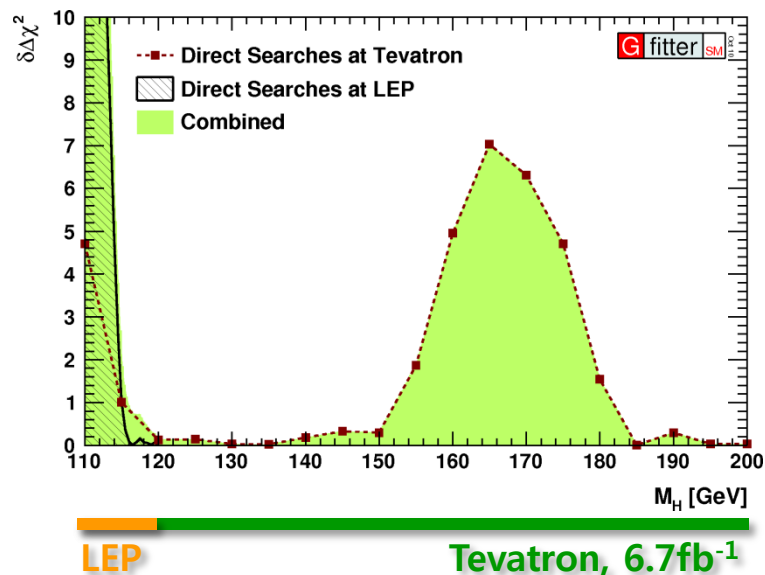
$$\Delta\alpha_{had}^{(5)}(M_Z^2) = (276.8 \pm 2.2) \cdot 10^{-4}$$

[Hagiwara et al., Phys. Lett. B649, 173 (2007)]



with direct Higgs searches:

- direct Higgs searches from LEP and Tevatron
- resulting contribution added to χ^2 during fit



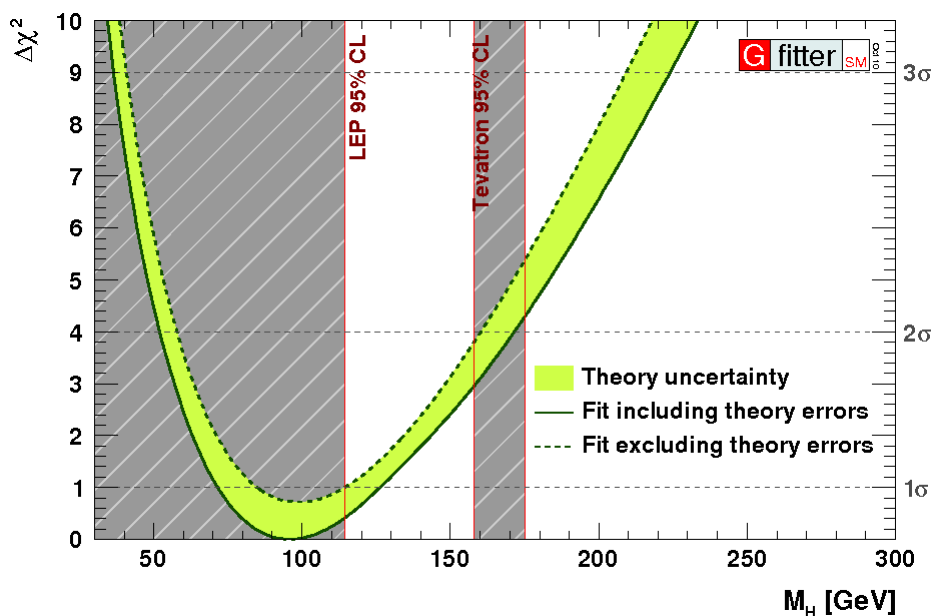
w/o direct Higgs searches:

- value at minimum $\pm 1\sigma$:

$$M_H = 96^{+31}_{-24} \text{ GeV}$$

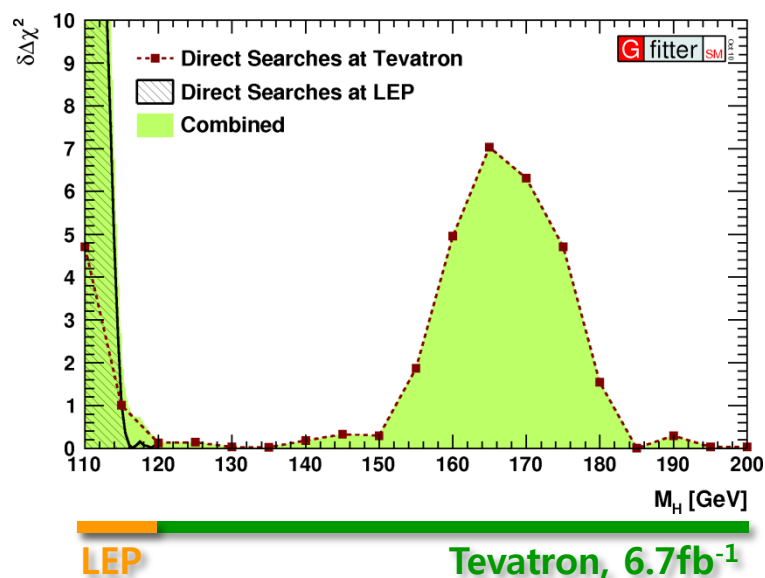
- 2σ interval: [52, 172] GeV

$$\Delta\alpha_{had}^{(5)}(M_Z^2) = (274.9 \pm 1.0) \cdot 10^{-4}$$



with direct Higgs searches:

- direct Higgs searches from LEP and Tevatron
- resulting contribution added to χ^2 during fit



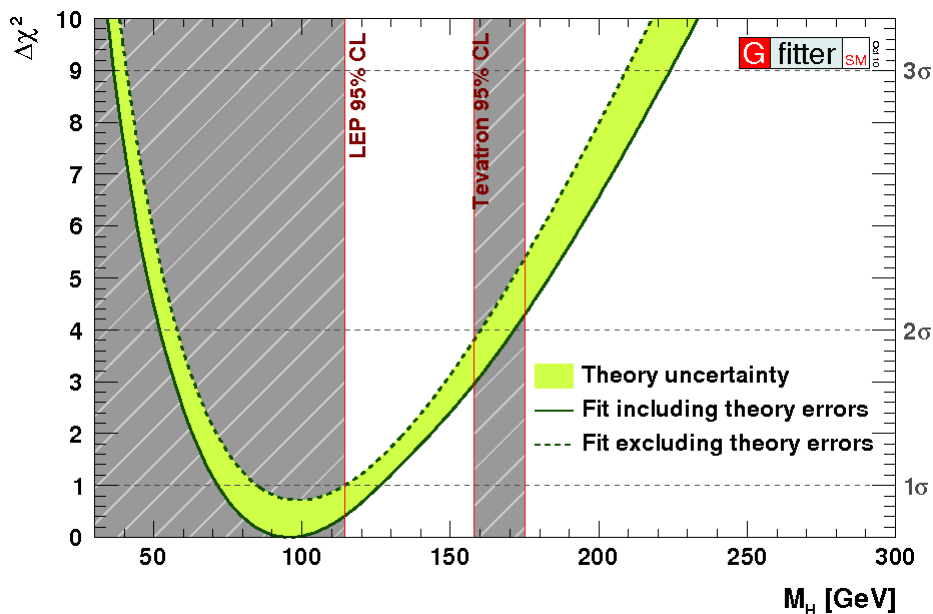
w/o direct Higgs searches:

- value at minimum $\pm 1\sigma$:

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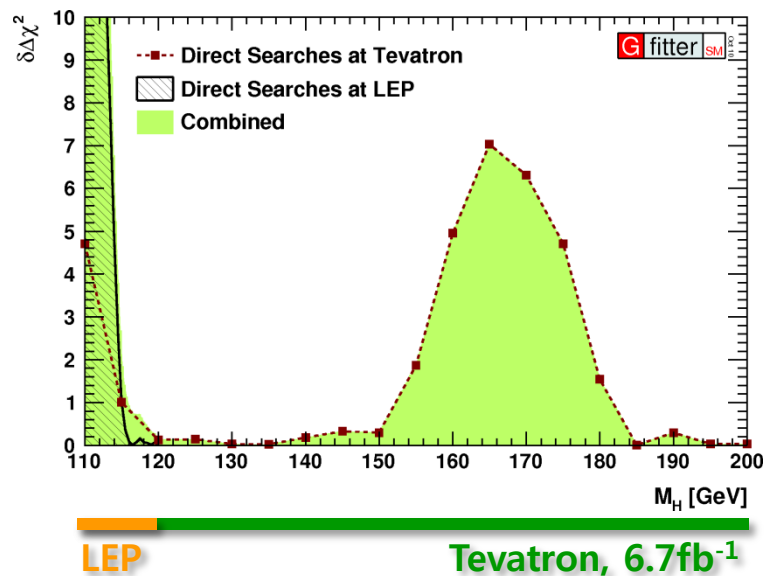
- 2σ interval: [52, 172] GeV

green error band: theory uncertainties directly included in χ^2 ("flat likelihood")



with direct Higgs searches:

- direct Higgs searches from LEP and Tevatron
- resulting contribution added to χ^2 during fit



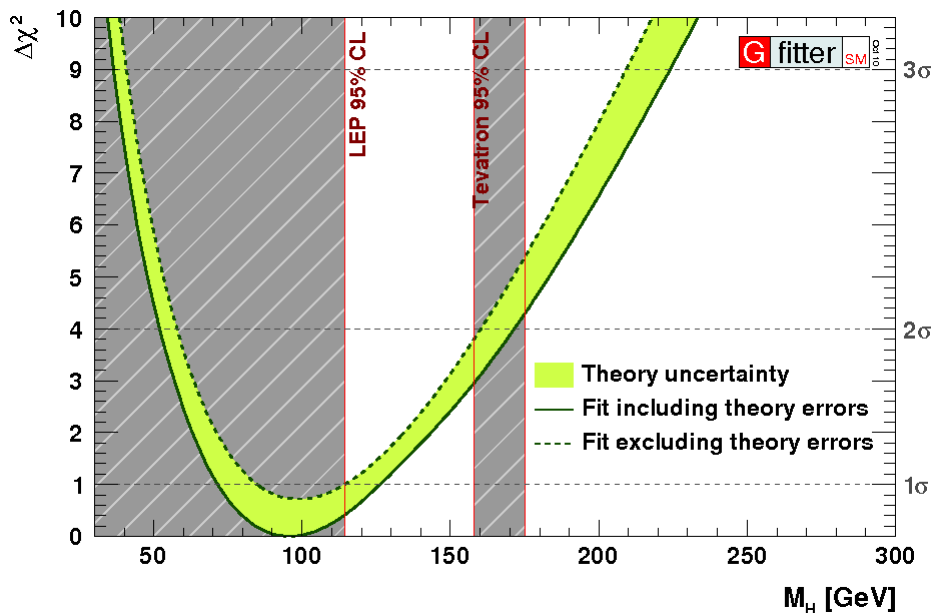
w/o direct Higgs searches:

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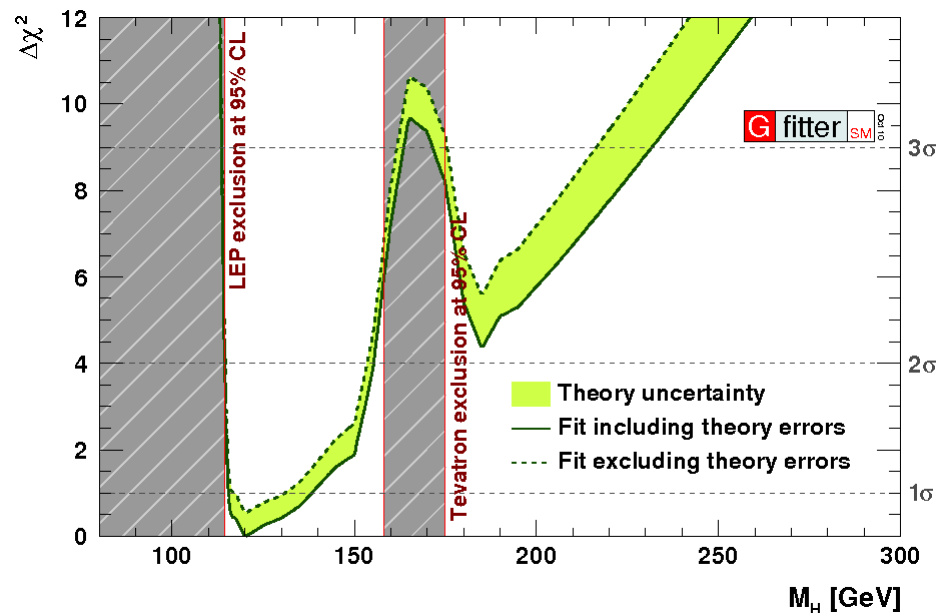


with direct Higgs searches:

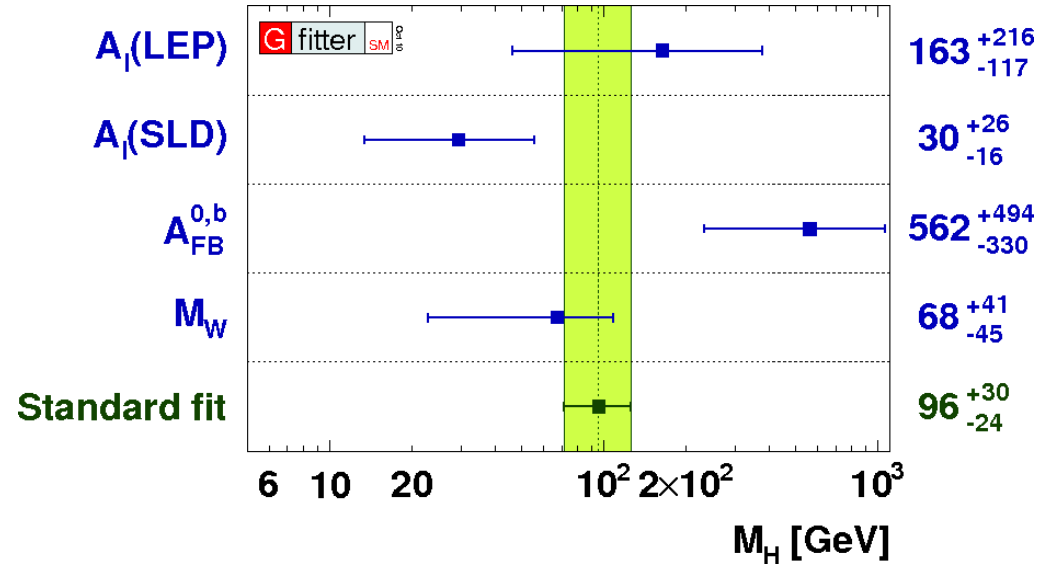
- value at minimum $\pm 1\sigma$:

$$M_H = 120.2^{+18.1}_{-4.7} \text{ GeV}$$

- 2σ interval: [114, 155] GeV



- Higgs mass constraints from most sensitive observables
 - known tension between M_W , $A_l(\text{SLD})$, and $A_{\text{FB}}^{0,b}$
 - including measurements of floating fit parameters



- compatibility of these measurements:
 - MC toy analysis (“look-elsewhere-effect”)
 - compare the χ^2_{min} of the full fit with χ^2_{min} of a fit without the least compatible measurement (here $A_{\text{FB}}^{0,b}$) $\rightarrow \Delta\chi^2_{\text{min}} = 7.9$
 - Generate toy sample around fitted values and repeat procedure by calculating the $\Delta\chi^2_{\text{min}} \rightarrow \Delta\chi^2_{\text{min}}^{\text{toy-distribution}}$
 - 1.6% (2.4σ) of toys show a result worse than the $\Delta\chi^2_{\text{min}}$ of the data

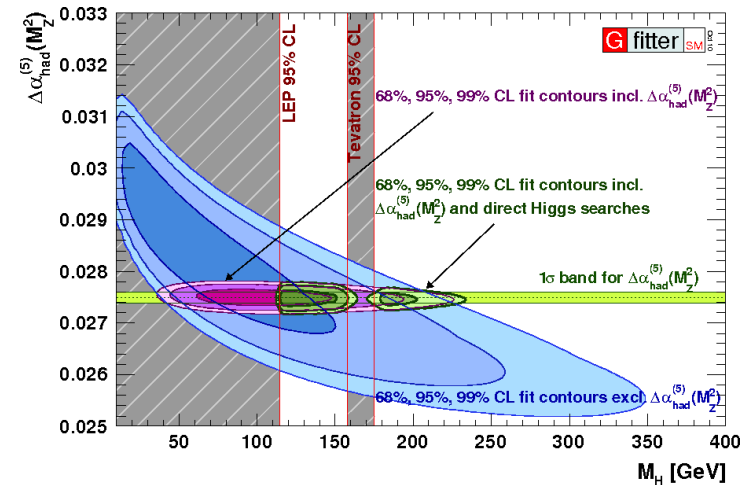
Fit Correlation between Fit Parameters

Parameter	$\ln M_H$	$\Delta\alpha_{\text{had}}^{(5)}(M_Z^2)$	M_Z	$\alpha_s(M_Z^2)$	m_t	\bar{m}_c	\bar{m}_b
$\ln M_H$	1	-0.17	0.13	0.03	0.32	-0.00	-0.01
$\Delta\alpha_{\text{had}}^{(5)}(M_Z^2)$		1	-0.01	0.35	0.01	0.00	0.02
M_Z			1	-0.01	-0.01	-0.00	-0.00
$\alpha_s(M_Z^2)$				1	0.03	0.01	0.05
m_t					1	0.00	-0.00
\bar{m}_c						1	0.00

due to heavier Higgs

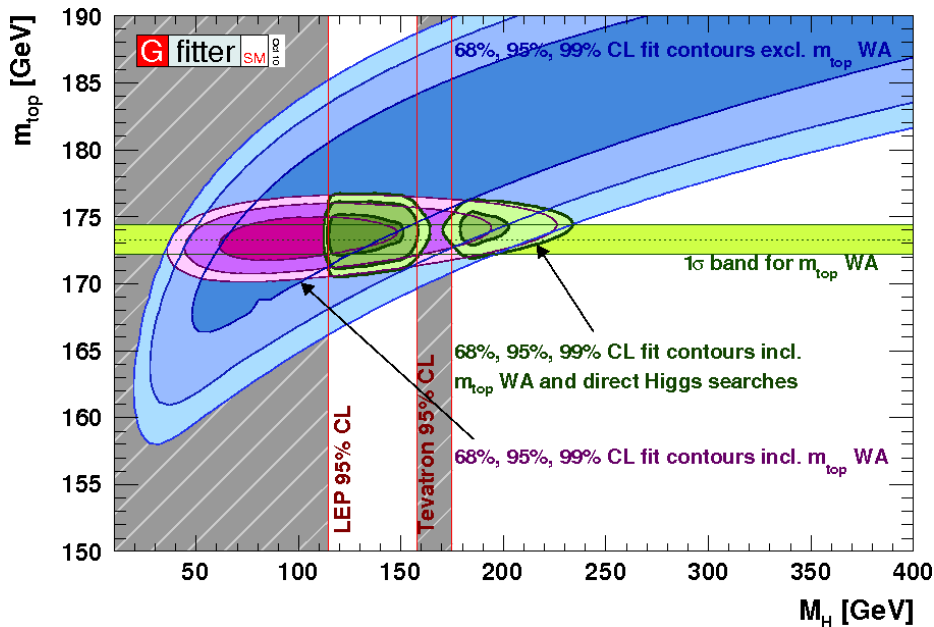
for comparison matrix determined by using old value for hadronic contribution to $\alpha_{\text{QED}} \rightarrow$ lighter Higgs

Parameter	$\ln M_H$	$\Delta\alpha_{\text{had}}^{(5)}(M_Z^2)$	M_Z	$\alpha_s(M_Z^2)$	m_t	\bar{m}_c	\bar{m}_b
$\ln M_H$	1	-0.395	0.113	0.041	0.309	-0.001	-0.006
$\Delta\alpha_{\text{had}}^{(5)}(M_Z^2)$		1	-0.006	0.101	-0.007	0.001	0.003
M_Z			1	-0.019	-0.015	-0.000	0.000
$\alpha_s(M_Z^2)$				1	0.021	0.011	0.043
m_t					1	0.000	-0.003
\bar{m}_c						1	0.000



Fit Correlation between Fit Parameters

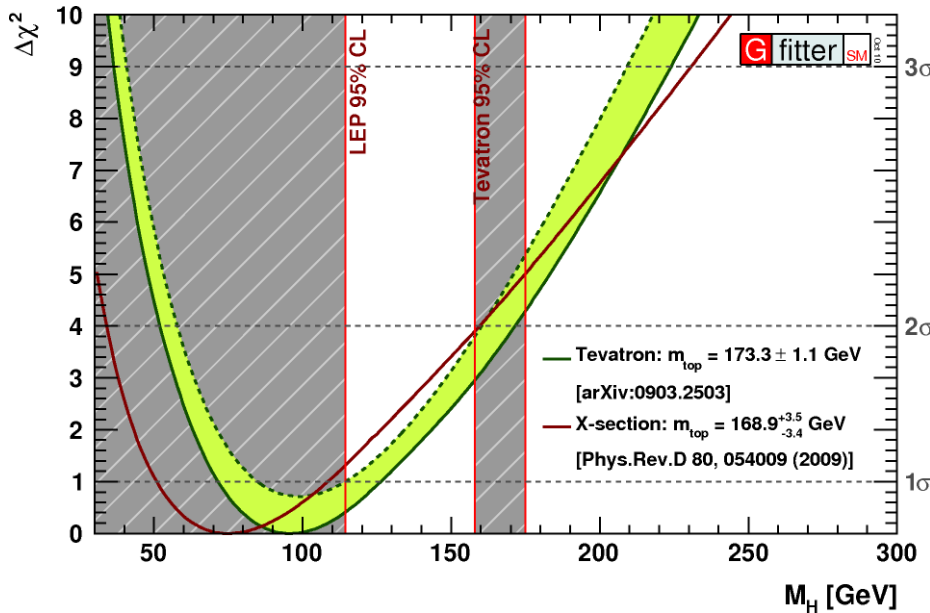
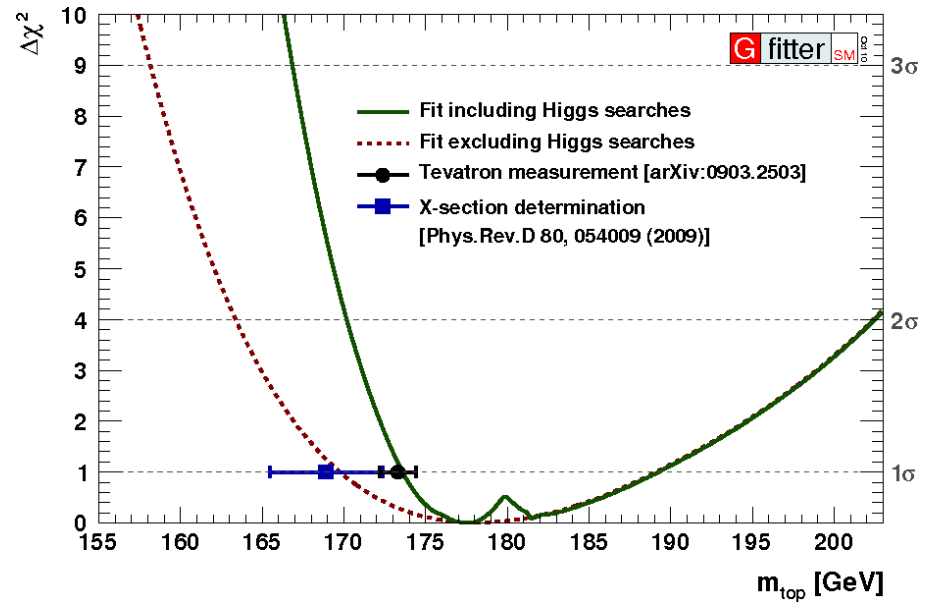
Parameter	$\ln M_H$	$\Delta\alpha_{\text{had}}^{(5)}(M_Z^2)$	M_Z	$\alpha_s(M_Z^2)$	m_t	\overline{m}_c	\overline{m}_b
$\ln M_H$	1	-0.17	0.13	0.03	0.32	-0.00	-0.01
$\Delta\alpha_{\text{had}}^{(5)}(M_Z^2)$		1	-0.01	0.35	0.01	0.00	0.02
M_Z			1	-0.01	-0.01	-0.00	-0.00
$\alpha_s(M_Z^2)$				1	0.03	0.01	0.05
m_t					1	0.00	-0.00
\overline{m}_c						1	0.00



⇒ top mass crucial input parameter for electroweak fit

Top Mass Determination

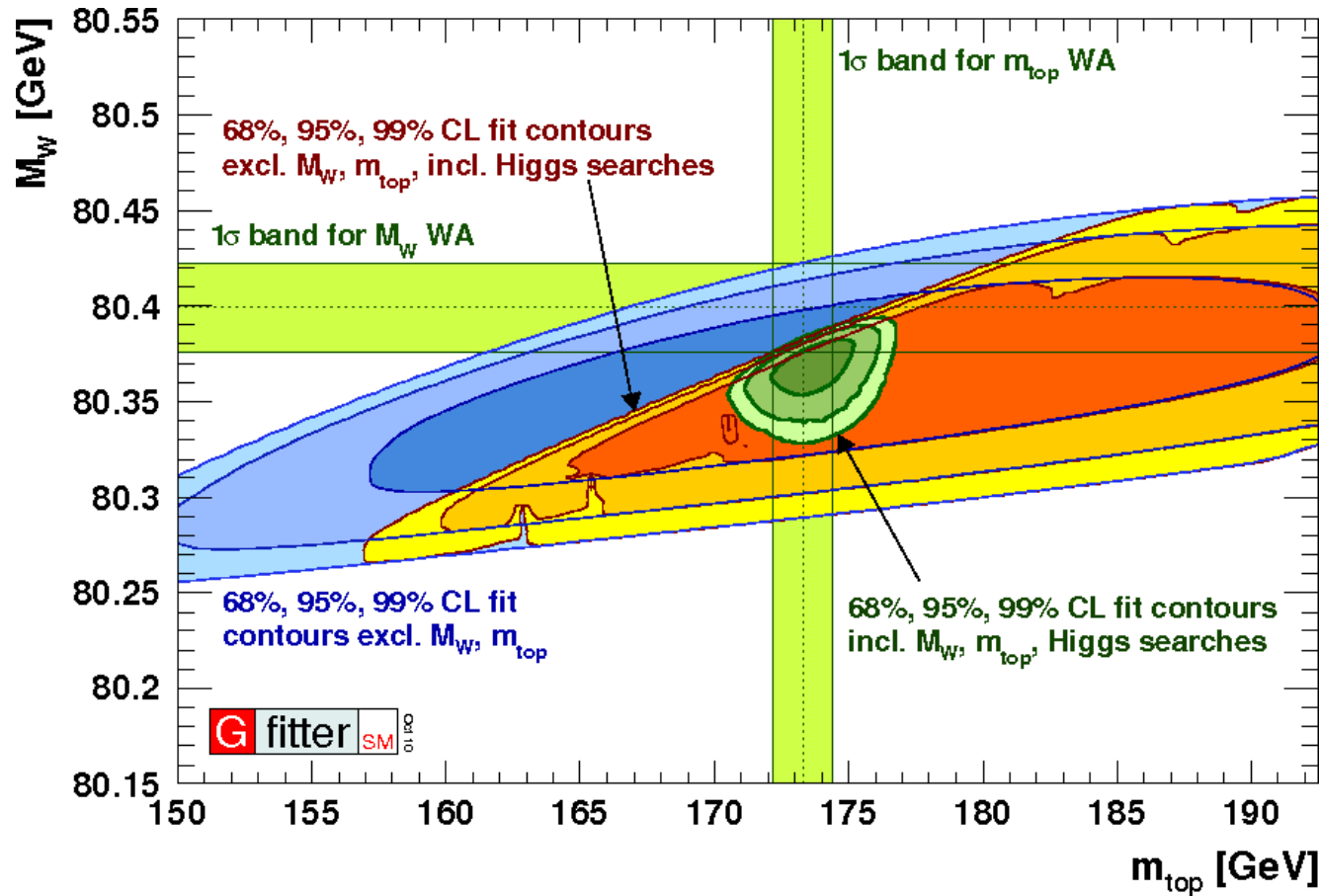
- theoretical predictions use top pole mass
- unclear definition of top mass at Tevatron: "MC" or pole mass?
[Hoang & Steward., Nucl.Phys.Proc.Suppl.185:220-226,2008]
⇒ additional uncertainty?
- alternative: extract top mass from total top pair cross-section
[Langenfeld, Moch, Uwer, Phys.Rev.D80:054009,2009]



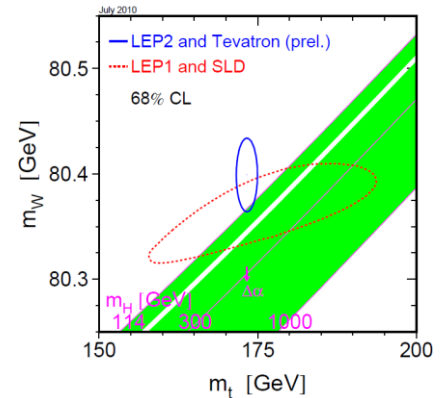
- direct: $m_{top} = 173.3 \pm 1.1$ GeV
- X-section: $m_{top} = 168.9^{+3.5}_{-3.4}$ GeV

SM Fit:

- w Higgs searches:
 $m_{top} = 177.4^{+11.8}_{-3.5}$ GeV
- w/o Higgs searches:
 $m_{top} = 178.2^{+10.9}_{-8.8}$ GeV



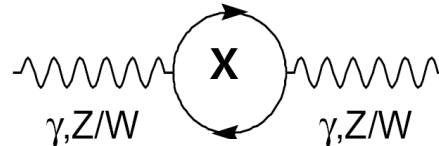
for comparison:



- indirect fit results agree with experimental values
- results from Higgs searches significantly reduce the allowed parameter space
- illustrative probe of SM (if M_H is measured at LHC and/or ILC)



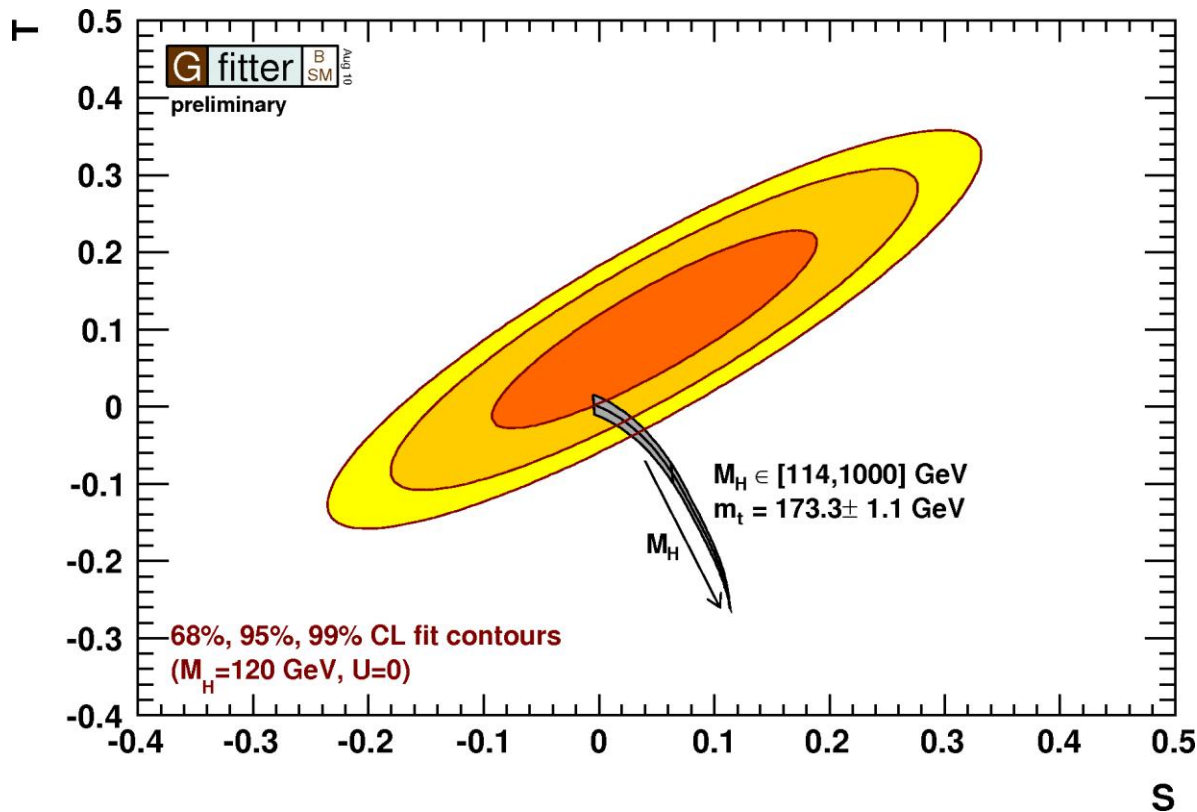
A Gfitter Package for SM Extensions



- oblique electroweak corrections to SM observables (physics beyond SM appear only through vacuum polarizations)
- STU parameters [Peskin and Takeuchi, Phys. Rev. D46, 1 (1991)]
 - $O_{\text{measurement}} = O_{\text{SM}}(M_{H,\text{ref}}, m_{t,\text{ref}}) + c_S S + c_T T + c_U U$
 - $S=T=U=0$ if data are equal to SM_{ref} prediction
 - S : new physics contribution to neutral current processes
 - $(S+U)$: new physics contribution to charged current processes
 - U only sensitive to W mass and width
 - usually very small in new physics models (often: $U=0$)
 - T : difference between neutral and charged current processes (sensitive to isospin violation)
 - also implemented corrections to Zbb couplings
[Burgess et al., Phys. Lett. B326, 276 (1994)] [Burgess et al., Phys. Rev. D49, 6115 (1994)]

Fit to Oblique Parameters

- derived from fit to electroweak observables (see global SM fit)
 - SM_{ref} : $M_H=120$ GeV, $m_t=173.3$ GeV
- comparison with SM prediction of ST parameters



Fit Results:

$$S = 0.02 \pm 0.11$$

$$T = 0.05 \pm 0.12$$

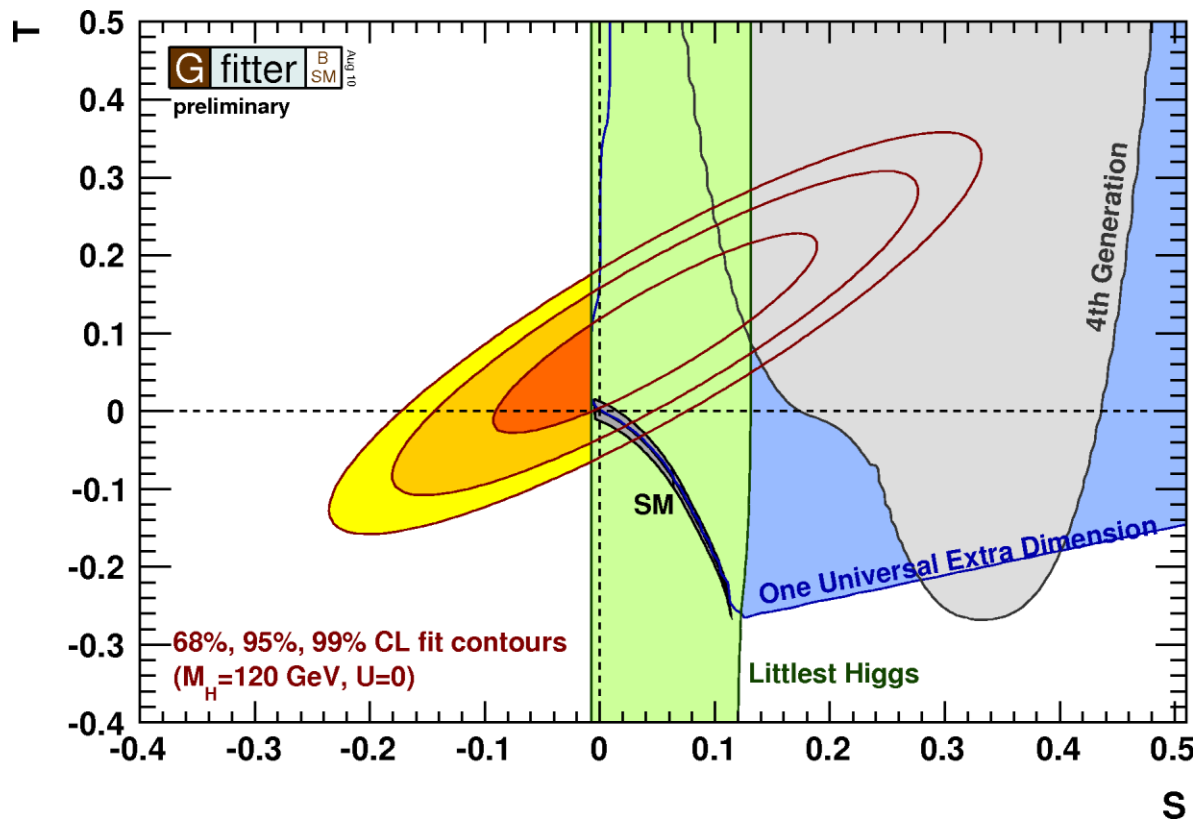
$$U = 0.07 \pm 0.12$$

Correlation:

	S	T	U
S	1	0.879	-0.469
T		1	-0.716
U			1

Fit to Oblique Parameters

- derived from fit to electroweak observables (see global SM fit)
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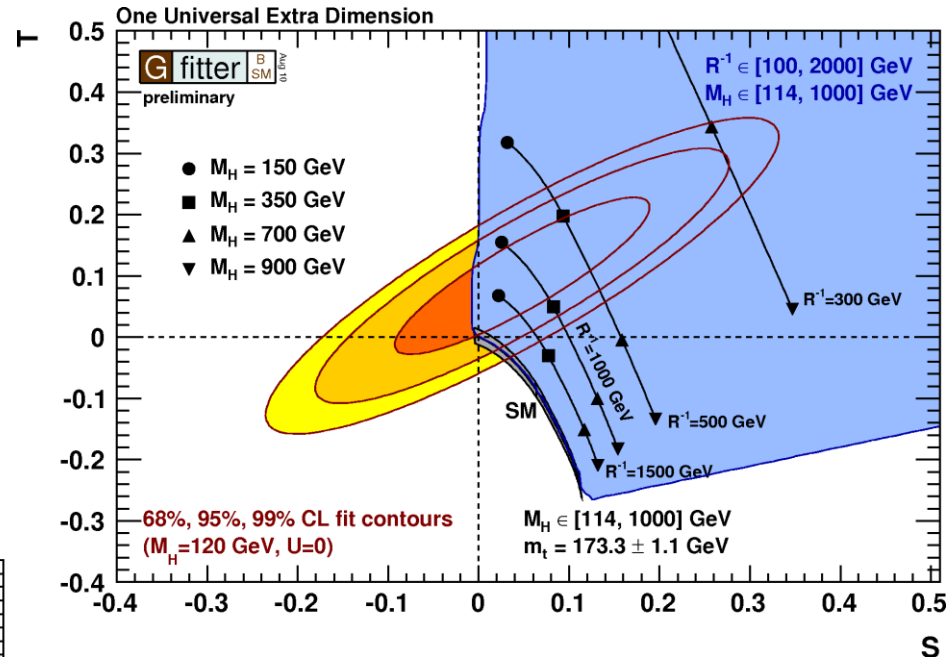
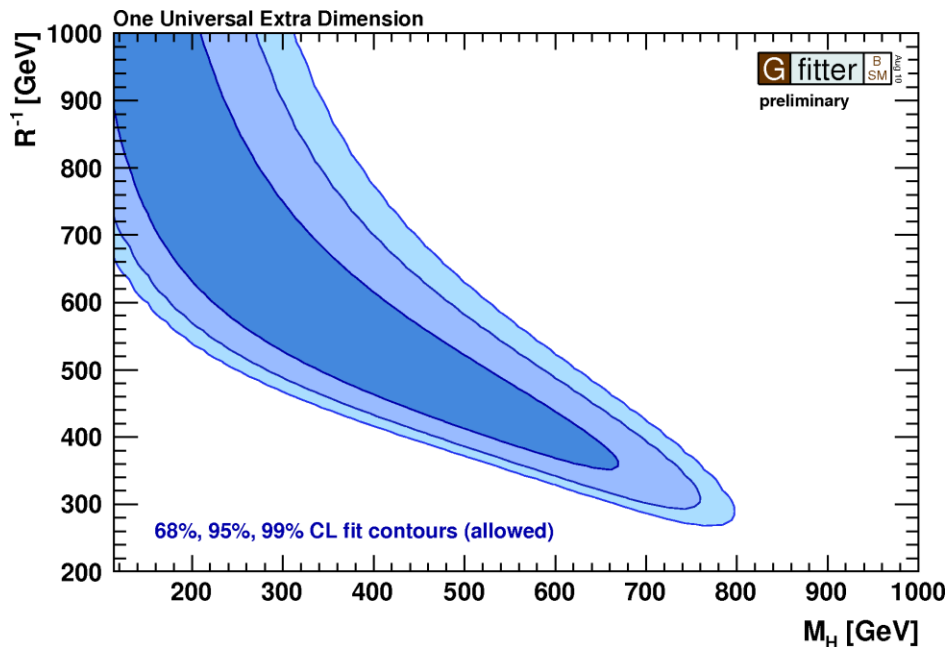
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One Universal Extra Dimension

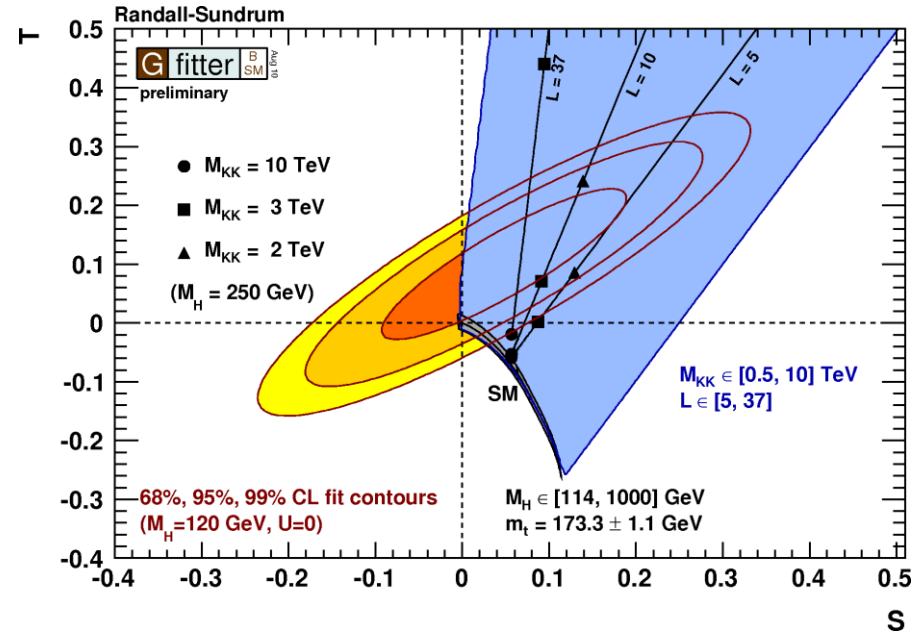
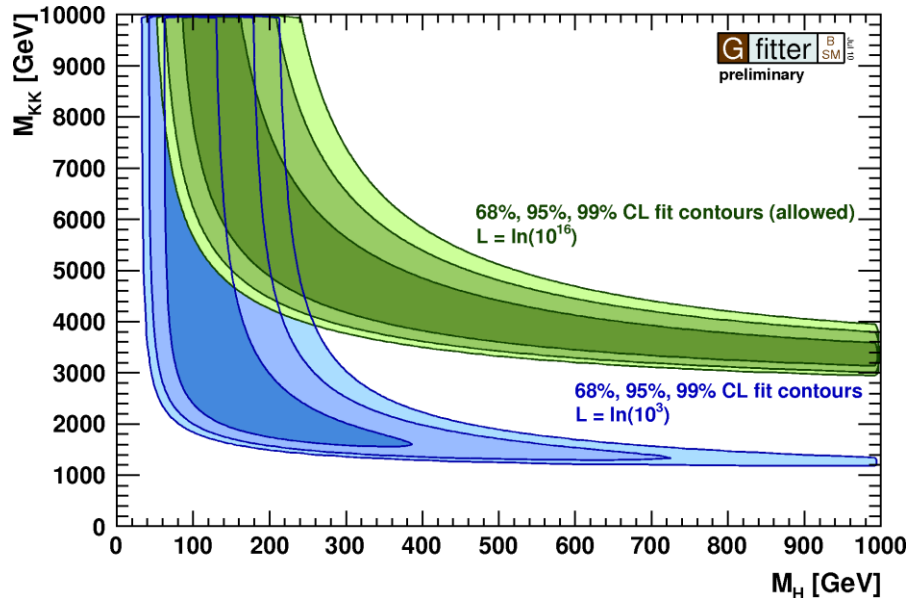
- all SM particles propagate in extra Dimension
- conservation of Kaluza-Klein (KK) parity → similar phenomenology as SUSY
- lightest KK state stable → Dark Matter candidate



- parameters of UED model
 - R^{-1} compactification scale (size of extra dimension) $m_{KK} \cong n/R$
 - oblique parameters depend on M_H
- oblique parameters replaced by corrections from UED model
 [Gogoladze et al., Phys.Rev. D 74, 093012 (2006)]
 [Appelquist et al., Phys.Rev. D67 (2003) 055002]

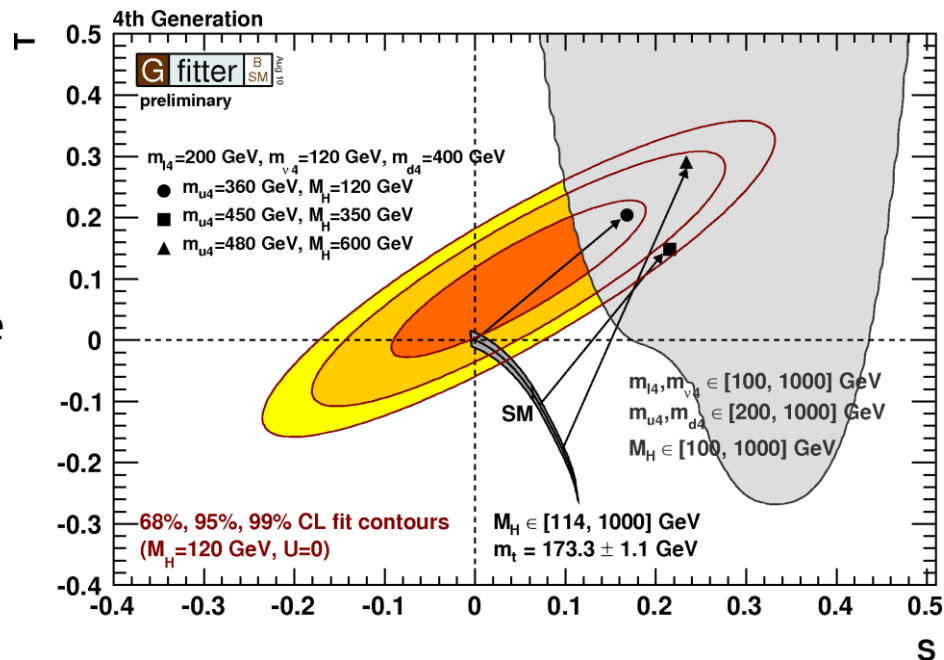
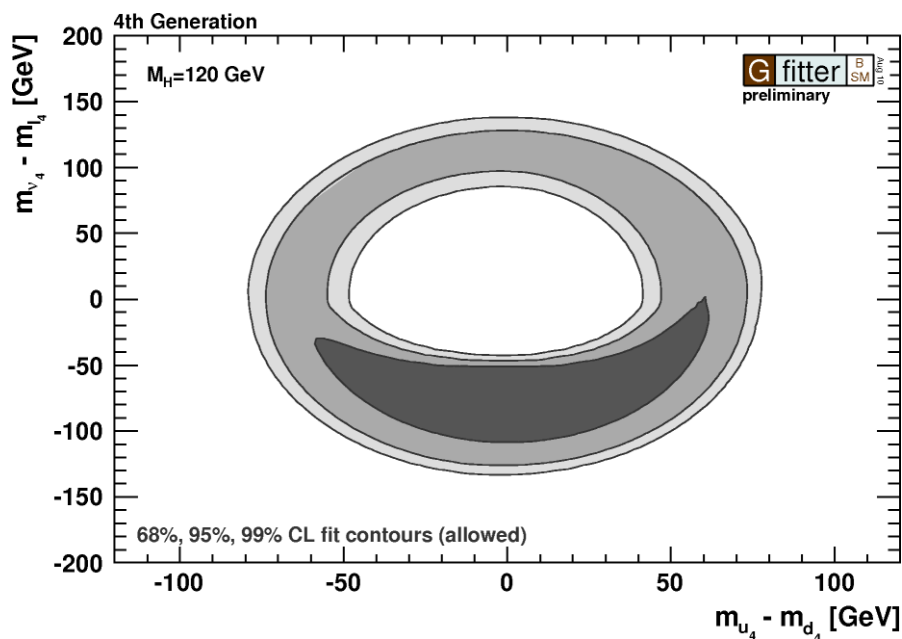
Warped Extra Dimensions (Randall-Sundrum)

- introducing one extra dimension (ED) for solving the hierarchy problem
- RS model characterized by one warped ED confined by two three-branes
- one brane contains SM particles
- extension: SM particles allowed to propagate in bulk region
- observation: heavy Kaluza-Klein modes



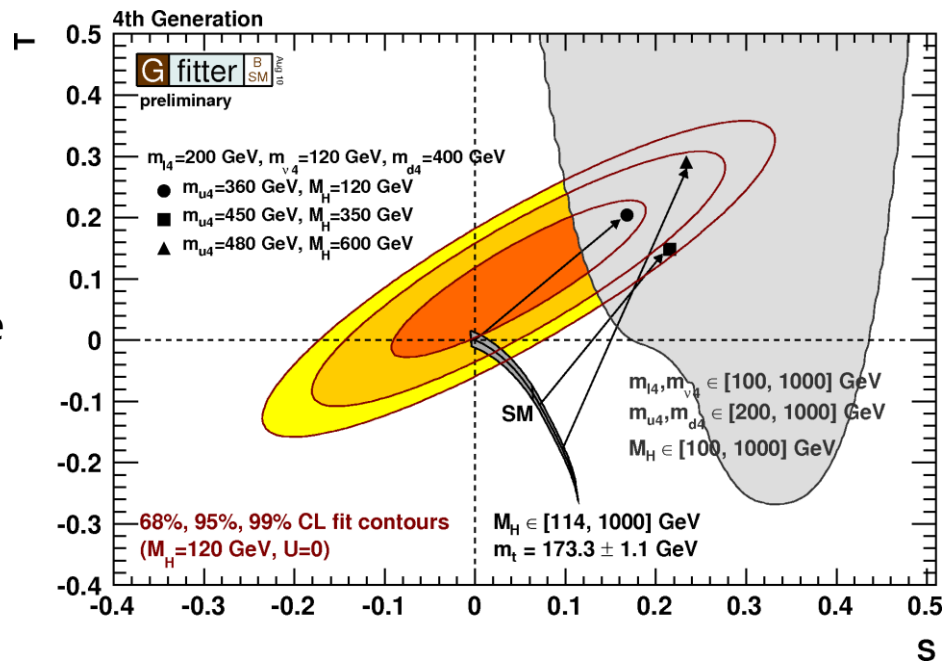
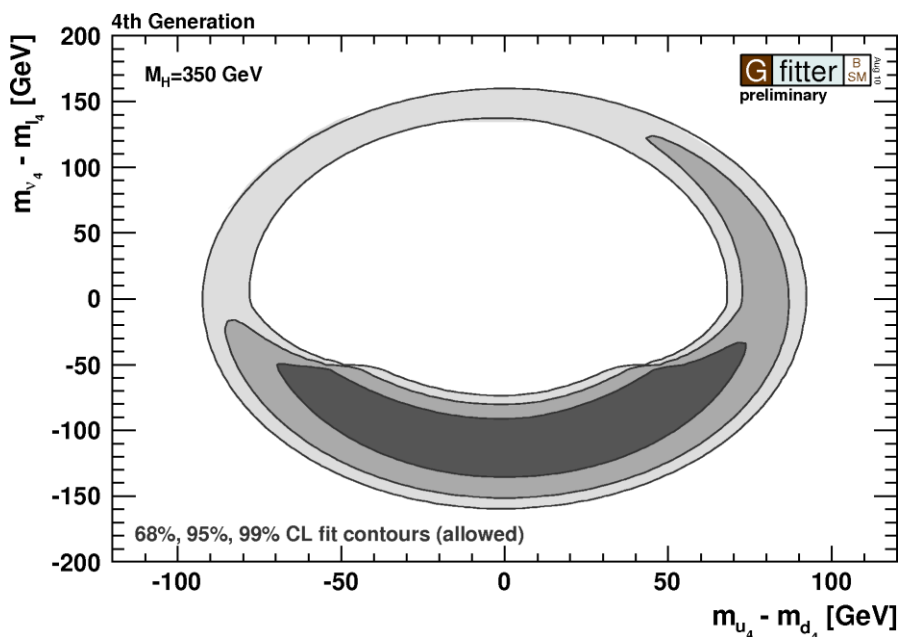
- parameters of RS model
 - M_{KK} : KK state
 - L : inverse warp factor, function of compactification radius
- oblique parameters replaced by corrections from RS model
[S. Casagrande et al., JHEP10(2008)094]

- motivation for fourth generation:
 - predicted by some GUT theories
 - can play an important role in electroweak symmetry breaking
- number of light neutrino ($m_\nu < M_Z/2$) due to measurement of Z width



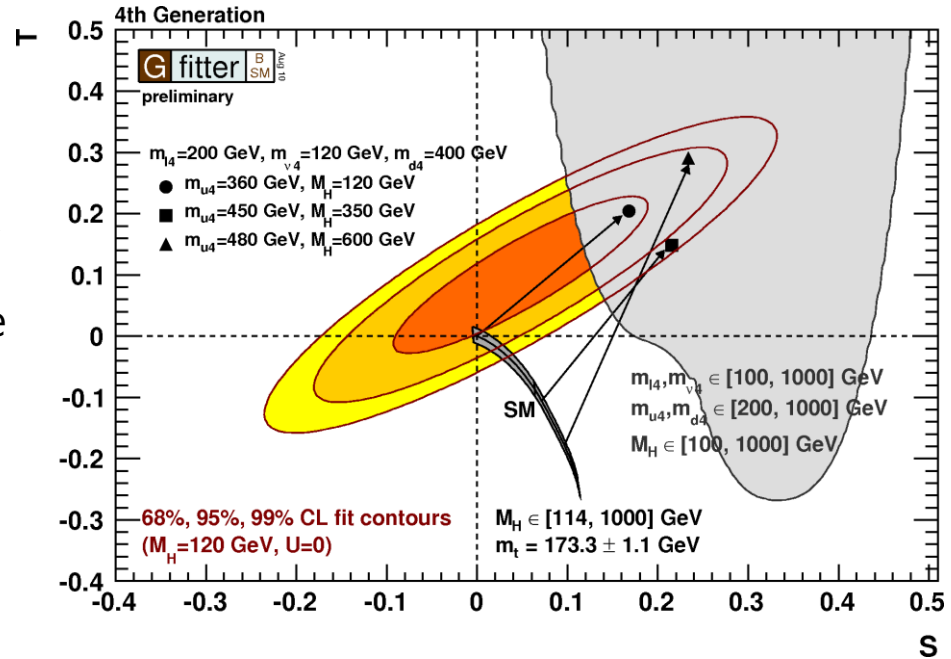
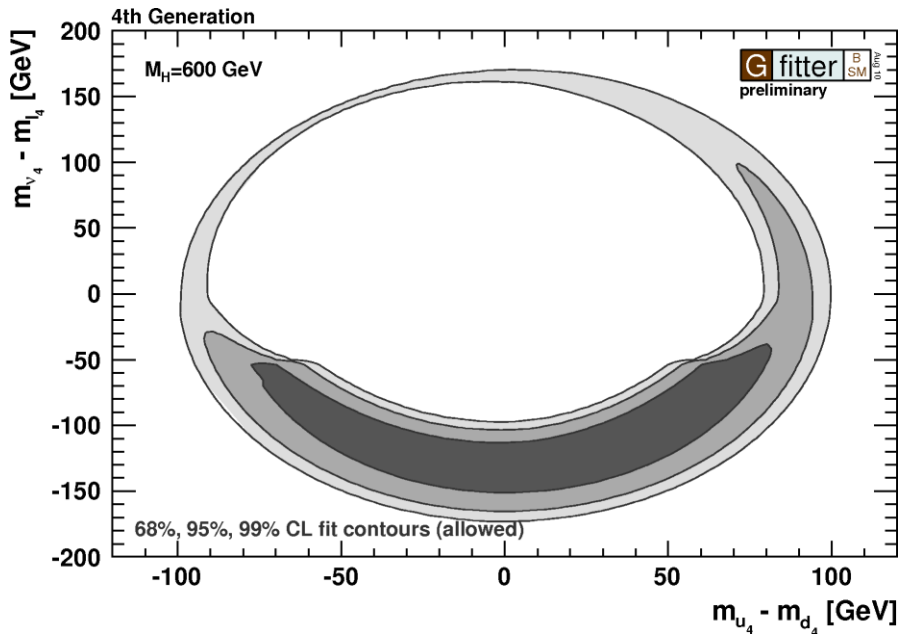
- mixing between first 3 generations and new generation neglected
- oblique parameters mostly sensitive to mass difference of new generation
 [Phys.Rev.D64, 053004 (2001)]

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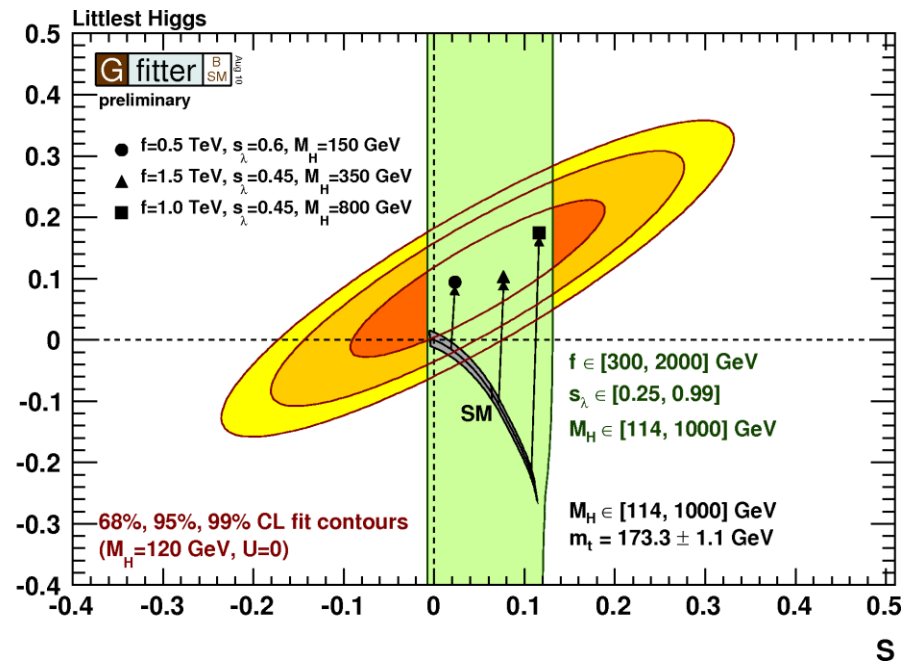
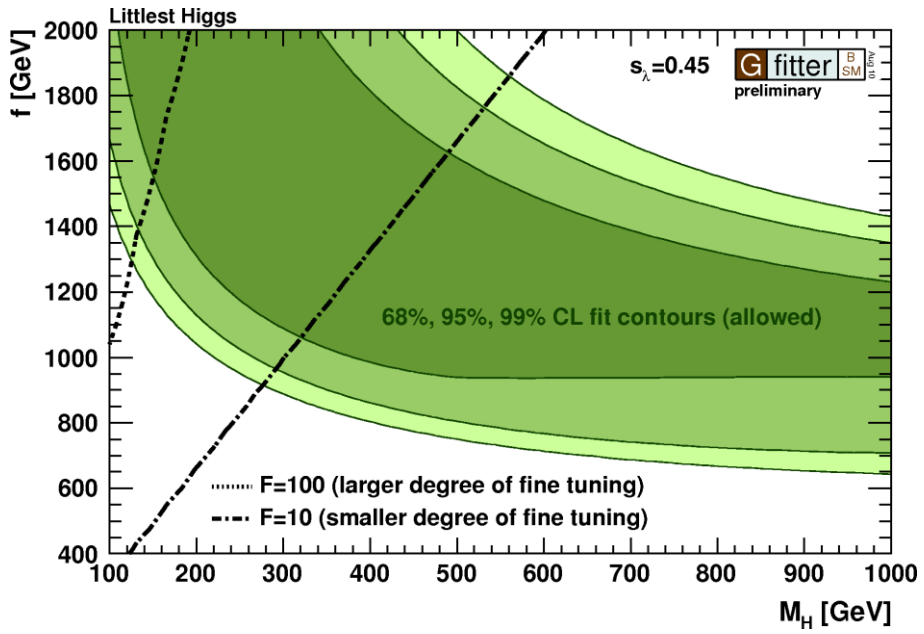
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Littlest Higgs with T-Parity

- Higgs pseudo-Nambu-Goldstone boson
- new fermions and new gauge bosons
 - two new top states (T-odd m_{T^-} and T-even m_{T^+})
 - LH solves hierarchy problem (new particles cancel SM loops)
- T-parity
 - provide dark matter candidate
 - forbids tree-level contribution from heavy gauge bosons to SM observables



- parameters of LH model
 - f symmetry breaking scale (scale of new particles)
 - $s_\lambda \cong m_{T^-} / m_{T^+}$ ratio of masses in top sector
 - order one-coefficient δ_c (exact value depends on detail of UV physics)
 - treated as theory uncertainty in fit (Rfit)
 - $\delta_c = -5 \dots 5$
- oblique parameters replaced by corrections from LH model
 [Hubisz et al., JHEP 0601:135 (2006)]

Gfitter Package

- flexible, generic C++ program including the statistical framework
- not shown SUSY, 2HDM
- results on <http://cern.ch/Gfitter>

Global SM Fit

- using state-of-the art predictions for the electroweak observables
- Toy Analysis of p-value: $p = 0.23 \pm 0.01 - 0.02$
- small Higgs masses are preferred from SM Fit
- N³LO determination of $\alpha_s(M_Z) = 0.1193 \pm 0.0028 \pm 0.0001$

Oblique Corrections

- SM extension allow heavy Higgs Boson
- more models implemented (2HDM, TechniColor, Inert Doublet Model, ADD, ...)

Interpretation of Direct Higgs Searches

- direct Higgs searches from LEP and Tevatron
 - using one-sided CL_{s+b}
 - sensitive to too few Higgs-like events
 - we are interested in any kind of deviation from “s+b” hypothesis
 - also too many Higgs-like events
 - transform one-sided CL_{s+b} into 2-sided $CL_{s+b}^{2\text{-sided}}$
 - compute contribution to χ^2 assuming symmetric PDF: $\delta\chi^2 = \text{Erf}^{-1}(1 - CL_{s+b}^{2\text{-sided}})$
- alternative: use of test statistics $-2\ln Q$
 - similar behavior, but deeper minimum \Rightarrow slightly stronger constraint

