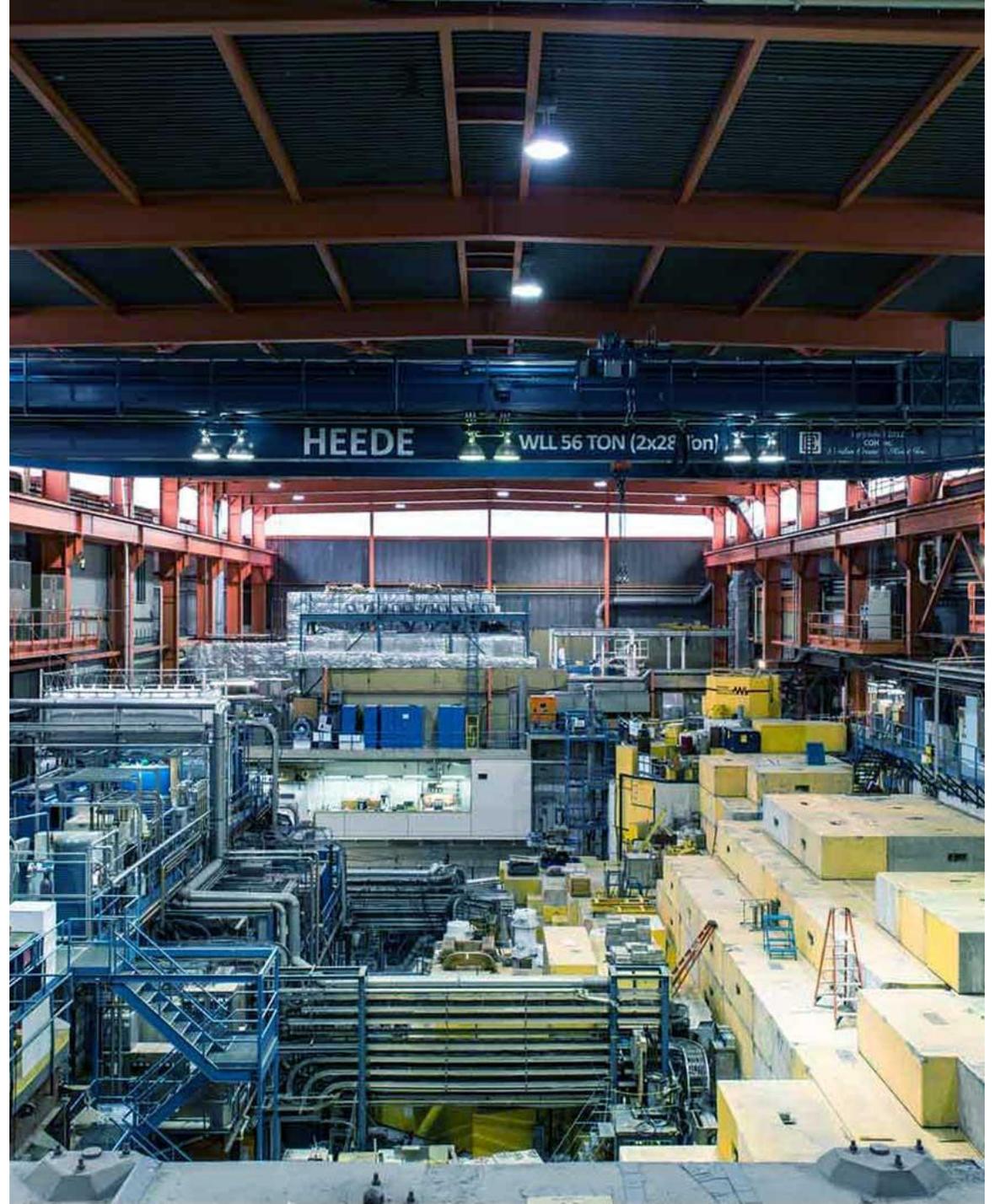


CaloDVAE : Discrete VAEs for Fast Calorimeter Simulation

A. Abhishek, E. Drechsler, W.
Fedorko, B. Stelzer

abhishek@myumanitoba.ca

2019-03-27

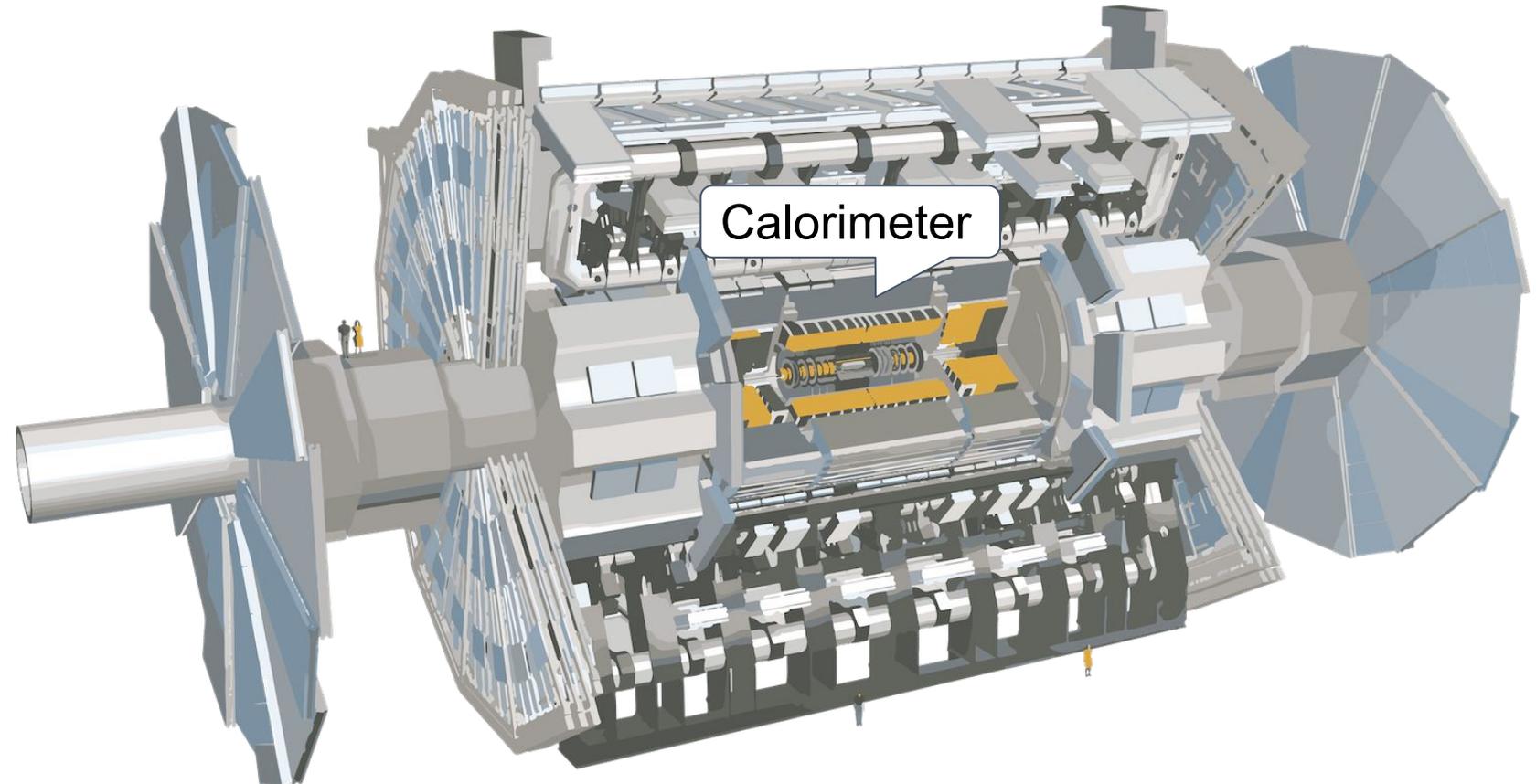


ATLAS at Large Hadron Collider

2

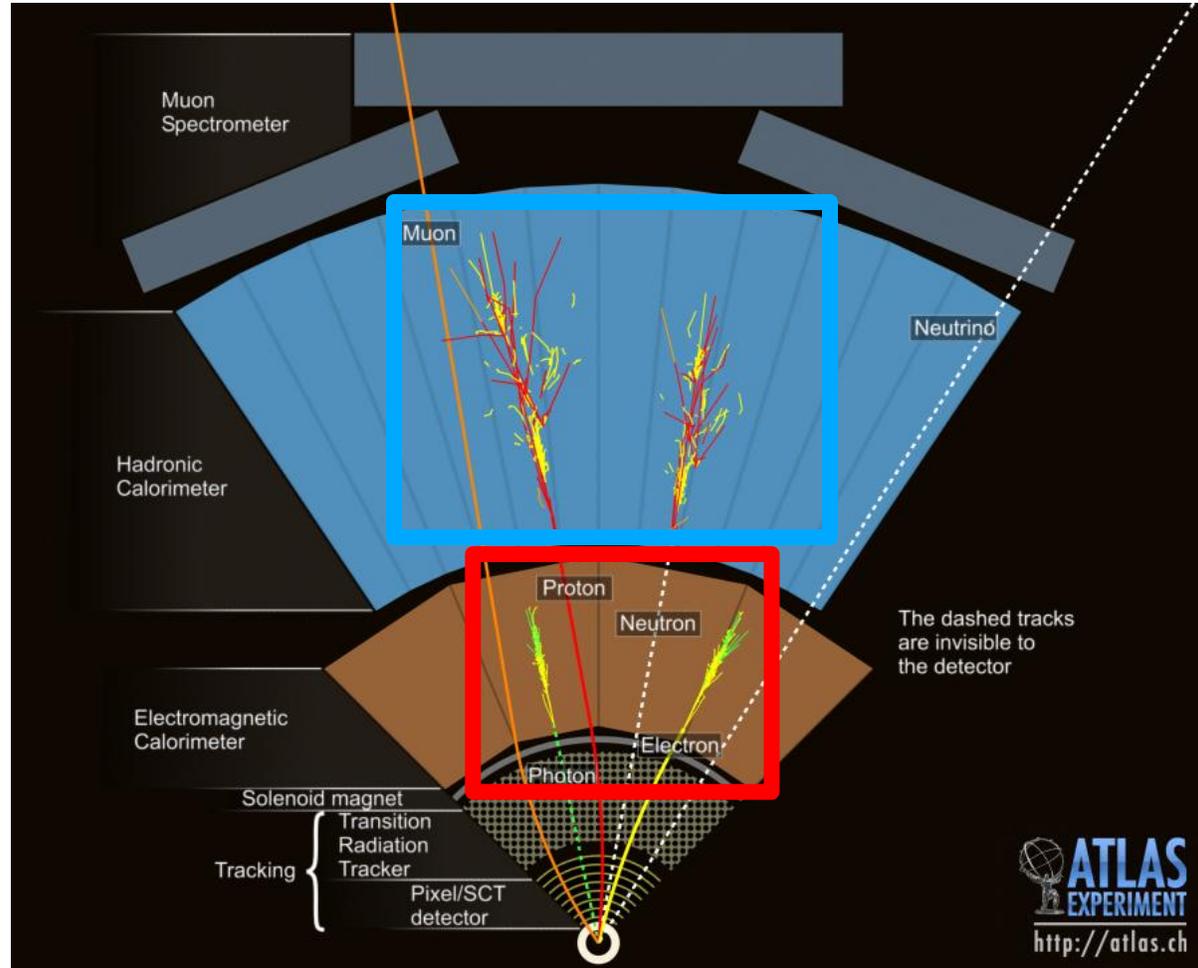
Physics goals :

1. Higgs Boson measurements
2. Physics beyond the Standard Model (BSM) e.g. Dark matter
3. Rare processes e.g. WVV production



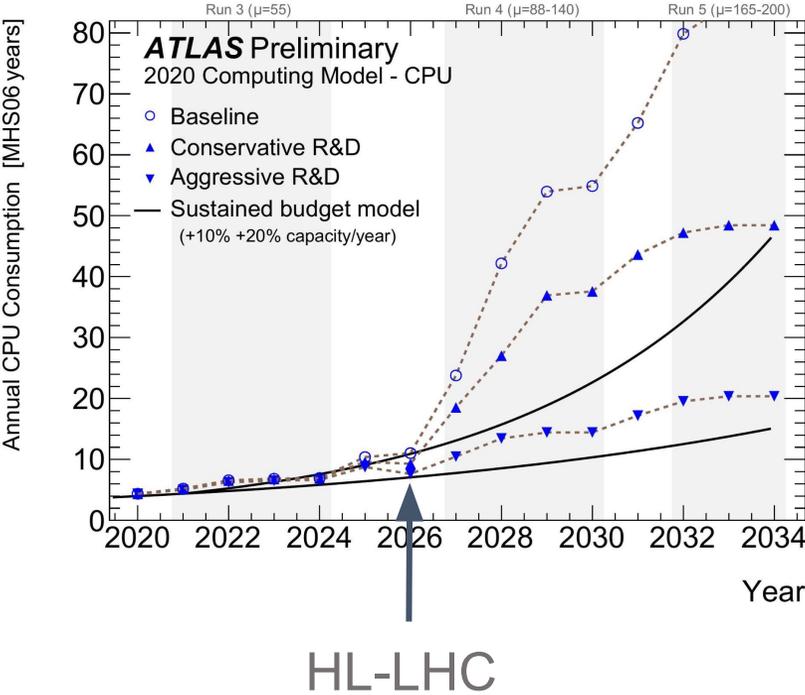
Calorimeters

Measure energy of the particles through **Electromagnetic** or **Hadronic** interactions

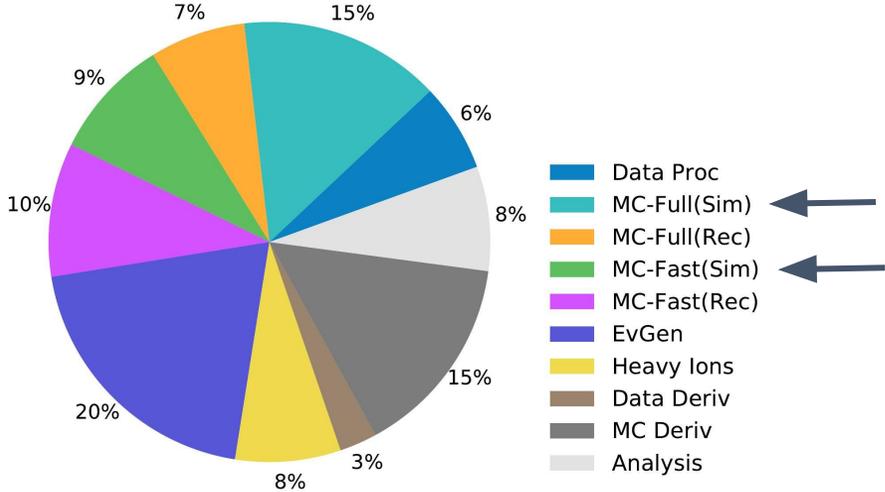


Motivation

Current techniques for Calorimeter shower simulation are computationally expensive

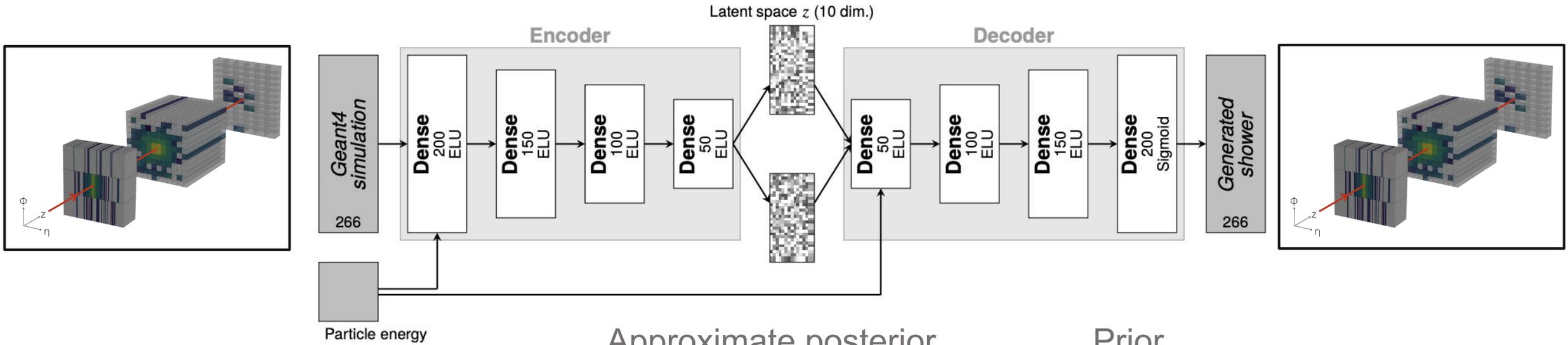


ATLAS Preliminary
2020 Computing Model -CPU: 2030: Baseline



[1] <https://twiki.cern.ch/twiki/bin/view/AtlasPublic/ComputingandSoftwarePublicResults>

Previous Work : DGMs for Fast Calorimeter Simulation



Approximate posterior

$$z_i \sim \mathcal{N}(\mu_i(\mathbf{x}, e), \sigma_i^2(\mathbf{x}, e))$$

Prior

$$\mathbf{z}_i \sim \mathcal{N}(0, 1)$$

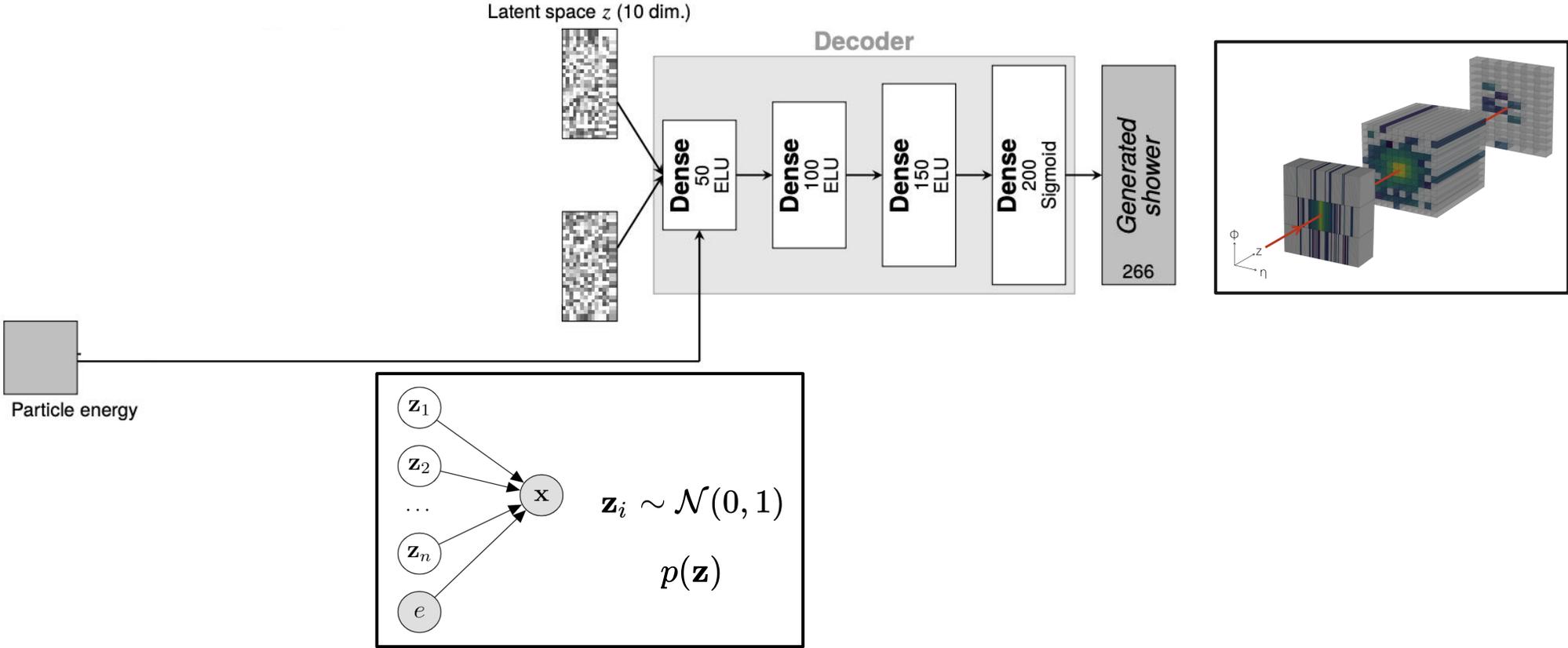
$$\text{Loss}(\mathbf{x}, \mathbf{x}'), \text{ e.g. } \frac{1}{n} \sum_{i=1}^n (\mathbf{x}_i - \mathbf{x}'_i)^2 \quad + \quad D_{KL}[q_\phi(\mathbf{z}|\mathbf{x}, e)||p(\mathbf{z})]$$

Autoencoding loss

KL divergence loss

Upper bound to the Negative log-likelihood of the data under the model distribution

Previous Work : DGMs for Fast Calorimeter Simulation

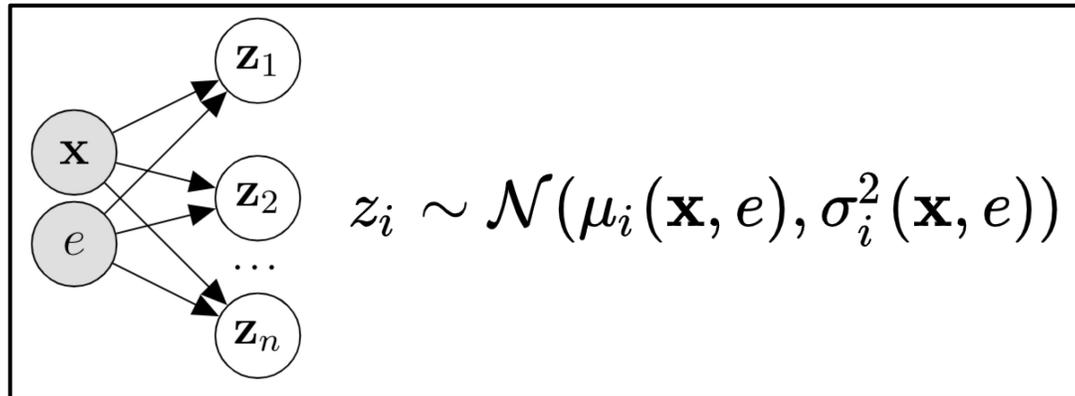


[1] Paganini, de Oliveira and Nachman (2018), arXiv:1712.10321 [2] <https://atlas.web.cern.ch/Atlas/GROUPS/PHYSICS/PUBNOTES/ATL-SOFT-PUB-2018-001>.

VAEs to Discrete VAEs

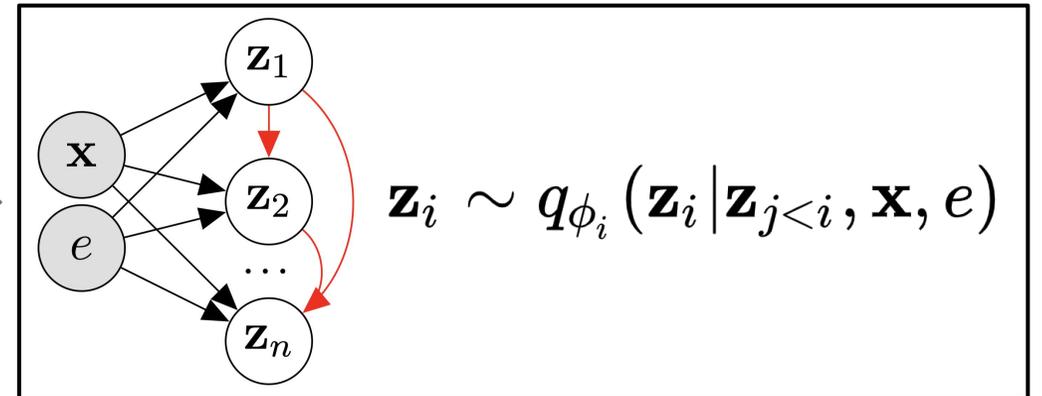
Approximate posterior distribution using during autoencoding $q_\phi(\mathbf{z}|\mathbf{x}, e)$

Continuous, Factorial Gaussian



Independence assumption

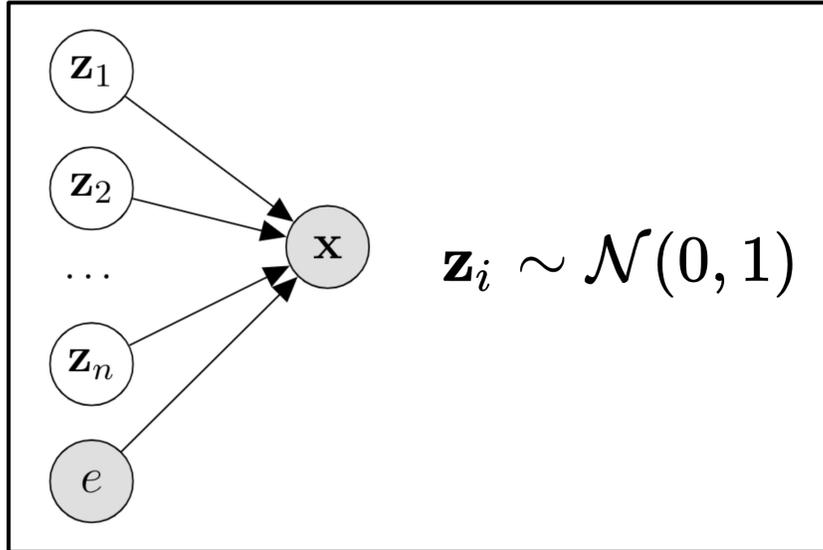
Discrete, Hierarchical Bernoulli



No independence assumption

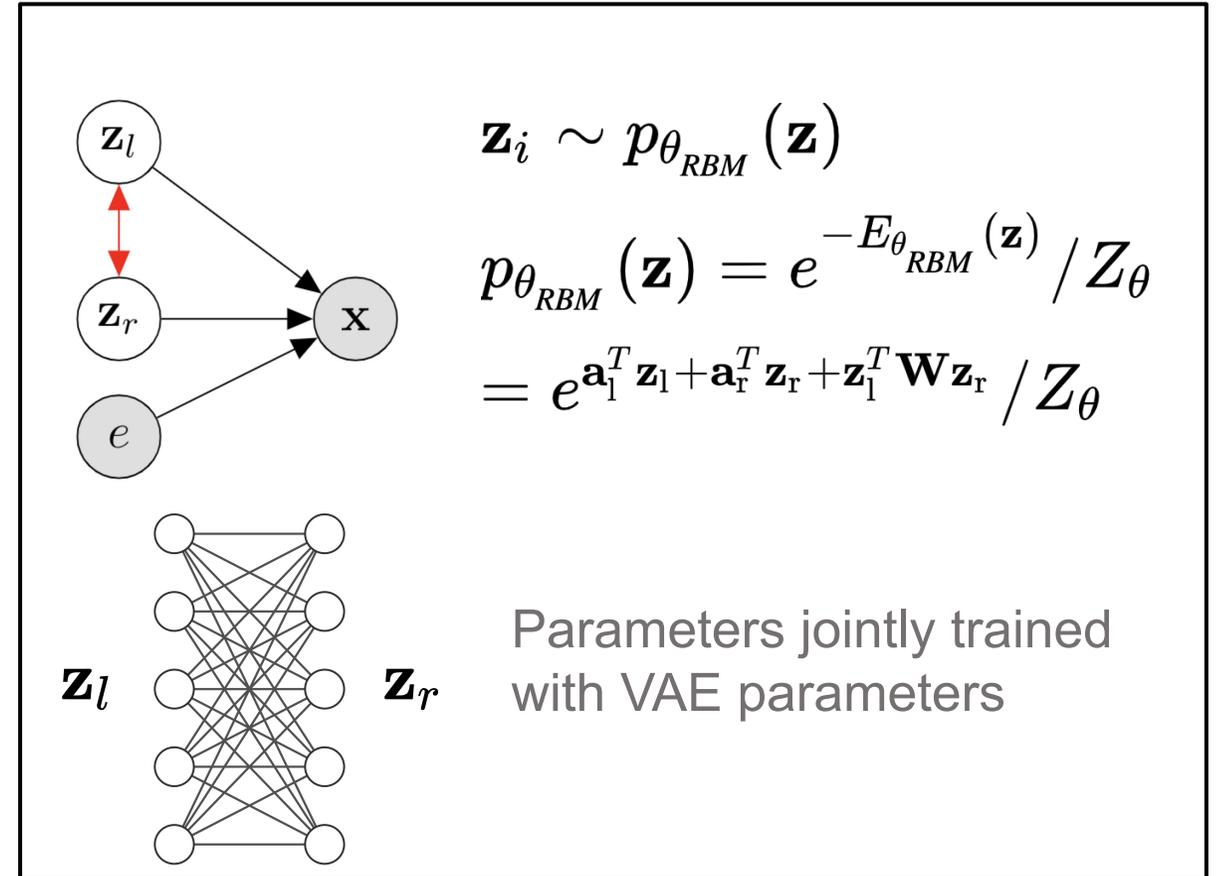
VAEs to Discrete VAEs

Factorial Normal Gaussian prior



- Computationally efficient sampling
- Independence assumption

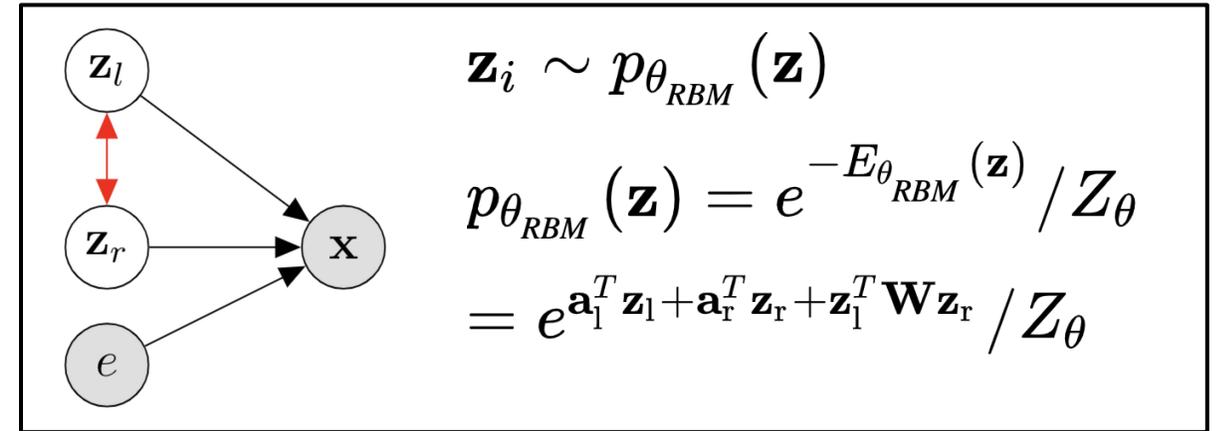
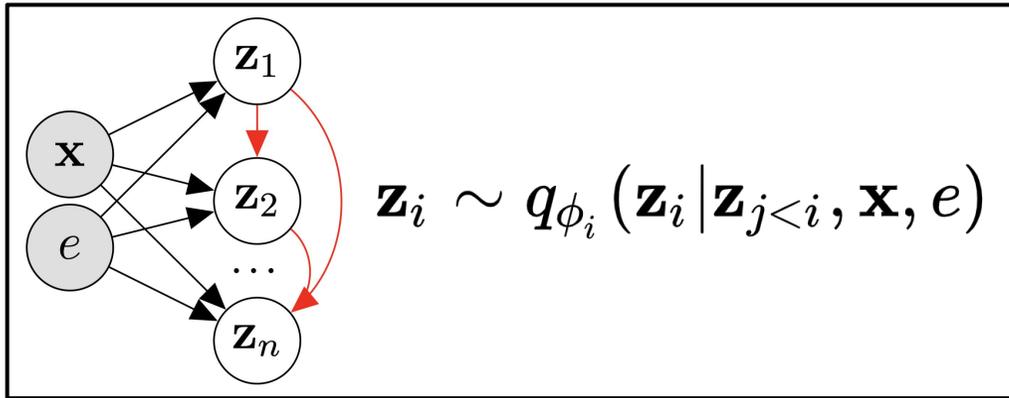
Trainable Restricted Boltzmann Machine prior



- Computationally expensive block Gibbs sampling

Why Discrete VAEs ?

1. More expressive latent space \longrightarrow Better performing generative model



2. Prospect of Quantum VAEs

- Replace the classical Restricted Boltzmann Machine (RBM) prior with a Quantum Boltzmann Machine (QBM) prior
- Use Quantum Annealing to generate latent variable samples instead of Block Gibbs sampling

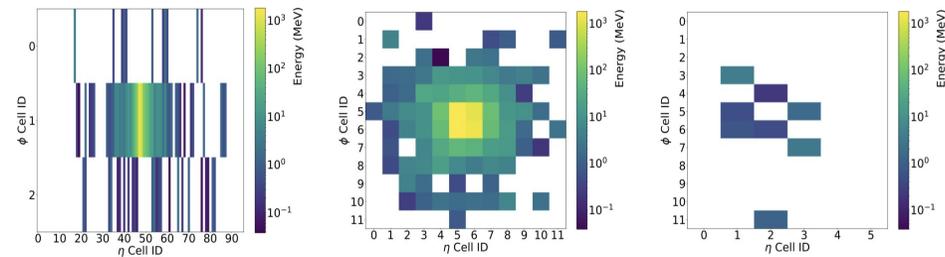
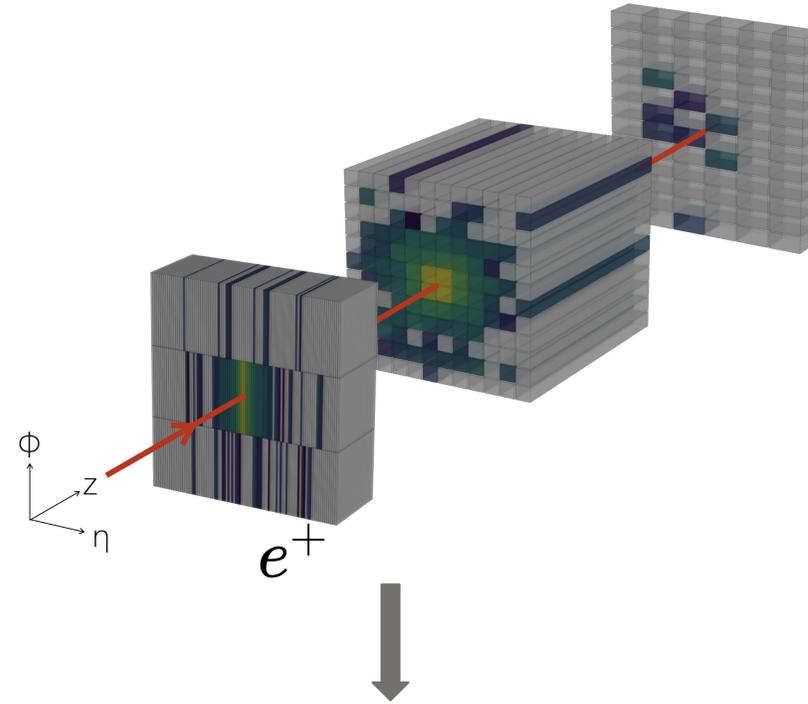
EM Calorimeter Shower Dataset

e^+, γ, π^+

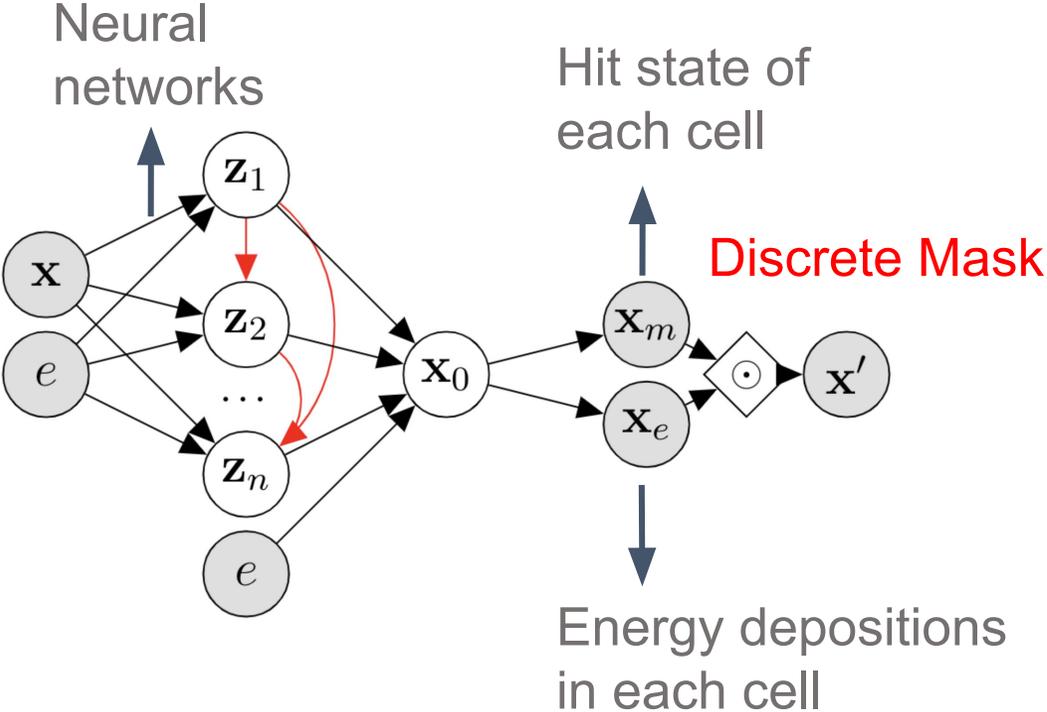
$1 < E < 100 \text{ GeV}$

100,000 events per particle type

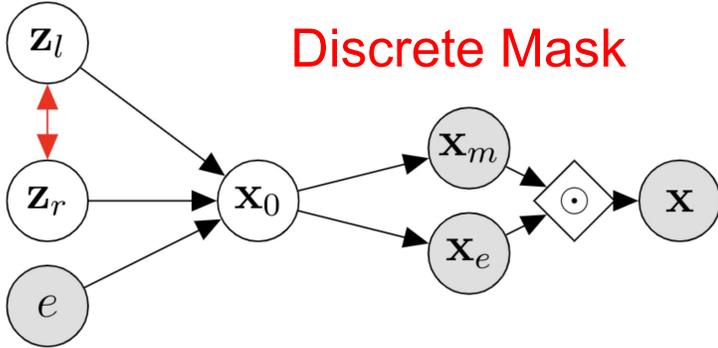
Inspired by ATLAS LAr EM Barrel
Calorimeter Geometry



DVAEs for Calorimeter Simulation



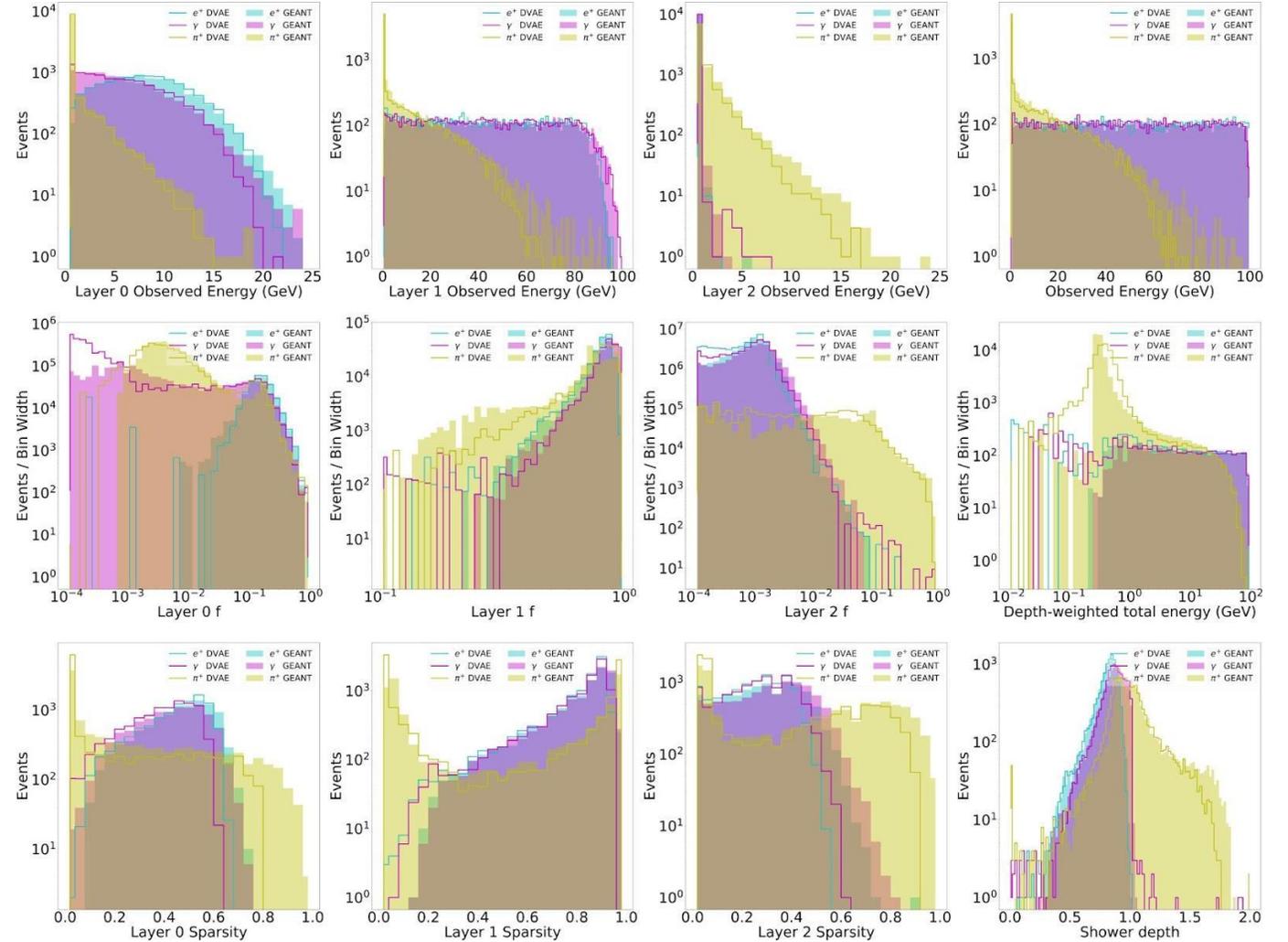
Autoencoding model



Generative model

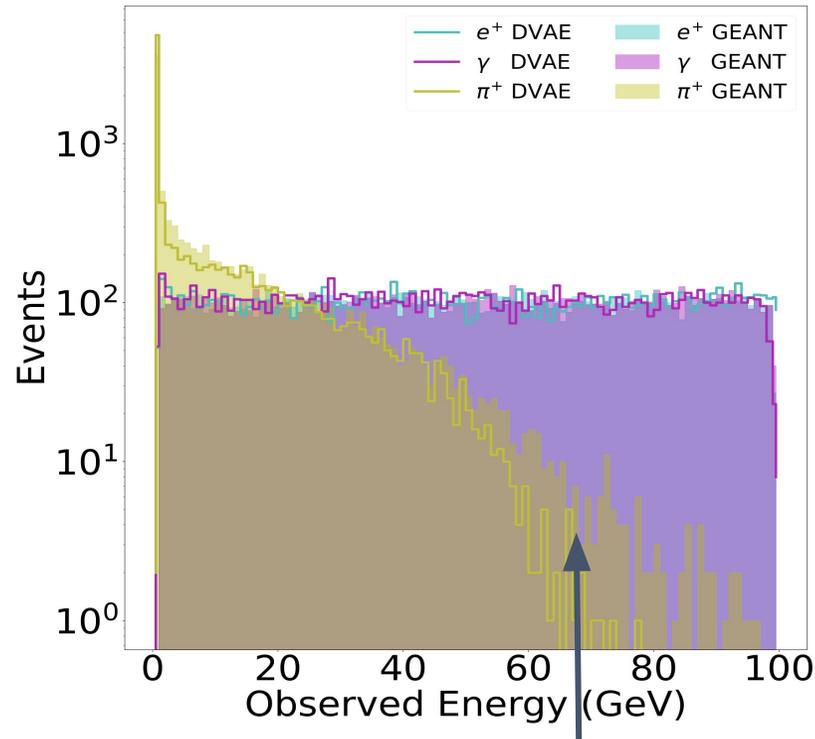
Results : Shower shape variables

Shower shape variable	Notes
$E_i = \sum_{\text{pixels}} \mathcal{I}_i$	Energy deposited in the i^{th} layer of calorimeter
$E_{\text{tot}} = \sum_{i=0}^2 E_i$	Total energy deposited in the electromagnetic calorimeter
$f_i = E_i/E_{\text{tot}}$	Fraction of measured energy deposited in the i^{th} layer of calorimeter
Depth-weighted total energy, $l_d = \sum_{i=0}^2 i \cdot E_i$	The sum of the energy per layer, weighted by layer number
Shower Depth, $s_d = l_d/E_{\text{tot}}$	The energy-weighted depth in units of layer number



Energy spectra

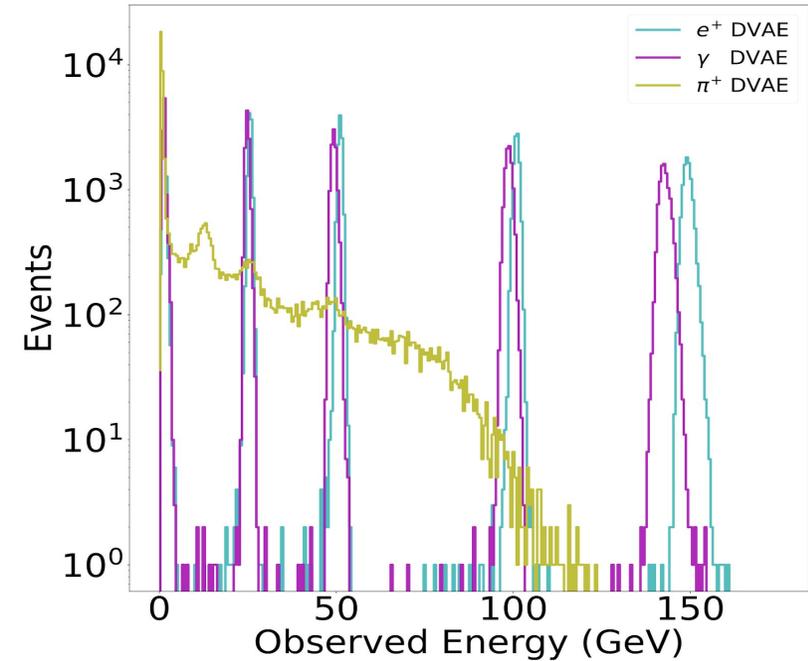
$1 < E < 100 \text{ GeV}$



Energy spectra of high energy charged pions underestimated

Energy conditioning

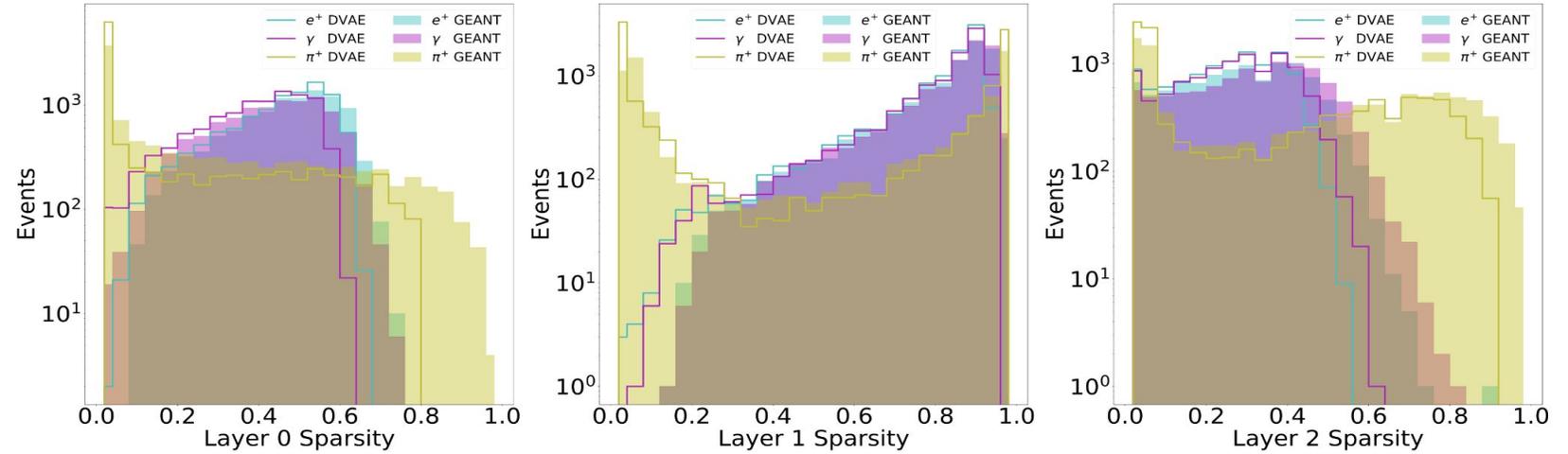
$E \in [1, 25, 50, 100, 150] \text{ GeV}$



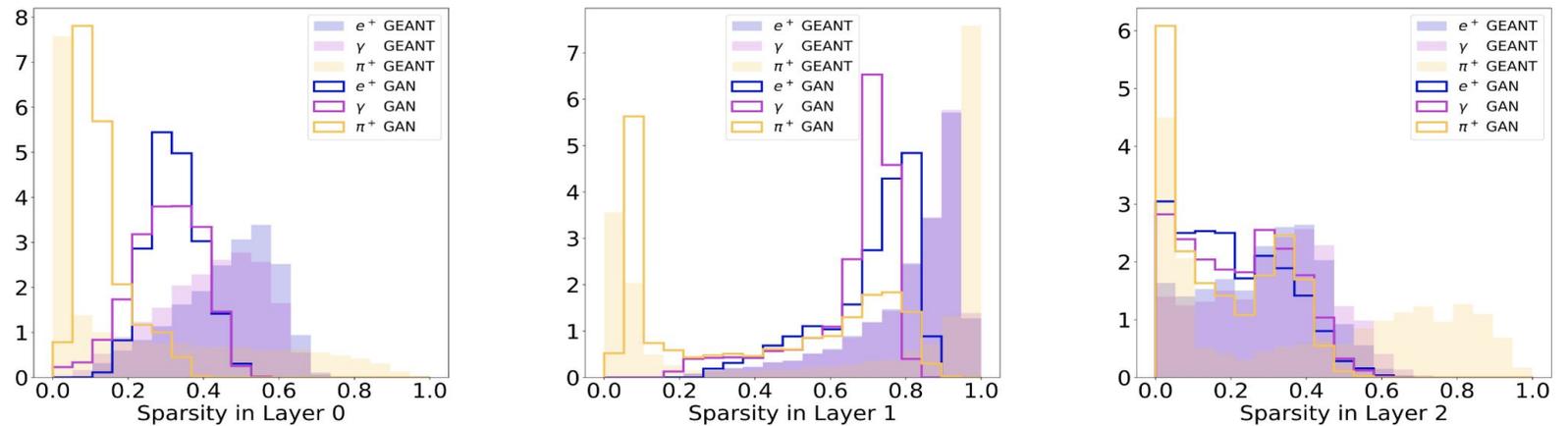
- Compelling peak at 150 GeV for positrons and photons
- Uncontained charged pions

Sparsity (Fraction of cells hit) per layer

CaloDVAE

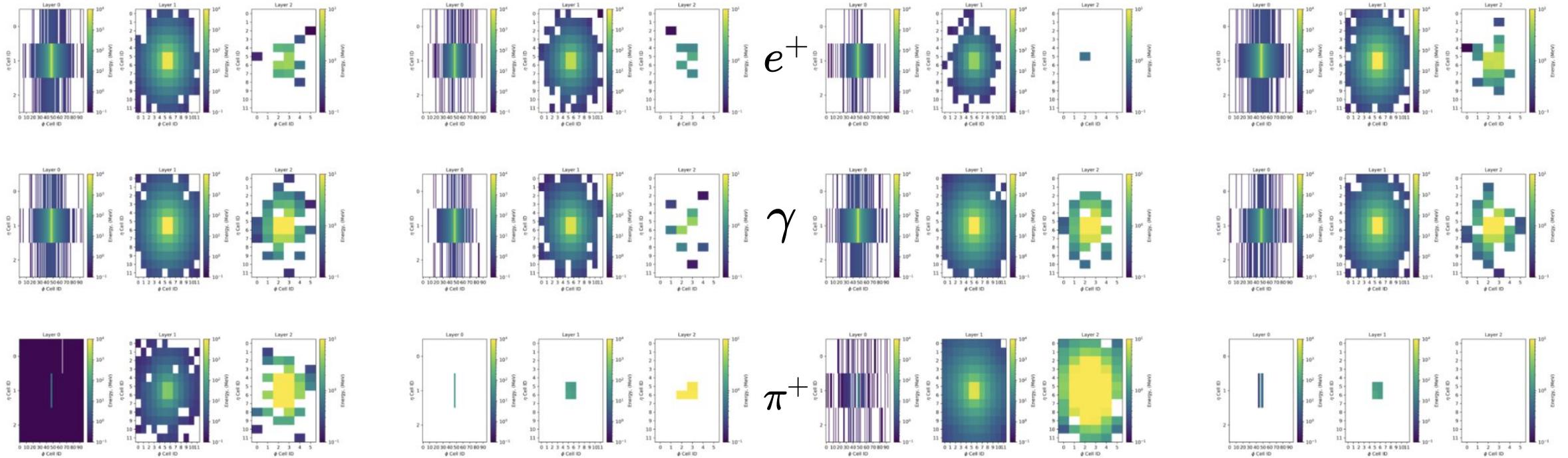


CaloGAN



- Previously observed to be challenging to reproduce
- Bimodal distribution for charged pions matches well

Shower Images



Generated samples recover :

- Wide variety of patterns of activated and non-activated cells
- **Centrality** and **lateral width** of the clusters
- Longitudinal behaviour of the shower - Most of the energy deposited in the middle layer

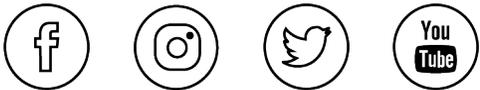
Summary and Future Outlook

- **DVAEs** show promising results for Calorimeter shower simulation
 - Able to match distributions of “shower shape” variables
- Latent generative process modelled by a **Restricted Boltzmann Machine (RBM)** is still computationally expensive
- Prospect of using **QVAEs** with latent generative process modelled by a **Quantum Boltzmann Machine (QBM)** allow us to use **quantum annealers** as sampling devices for the latent variables

Thank you
Merci

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Backup

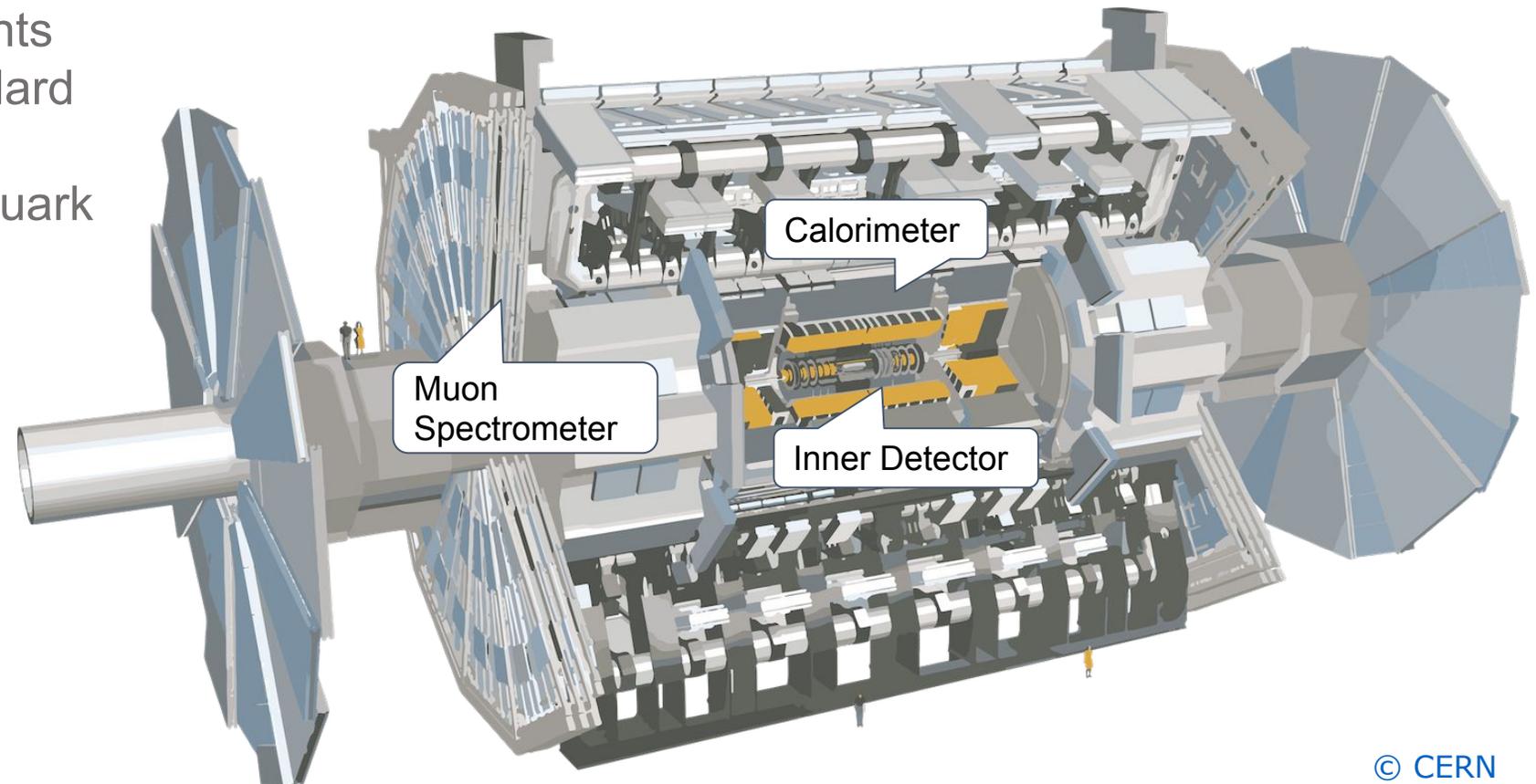
ATLAS at Large Hadron Collider

19

General-purpose detector

Scientific goals :

1. Higgs Boson Measurements
2. Physics beyond the Standard Model (BSM)
3. Rare processes e.g. top quark production

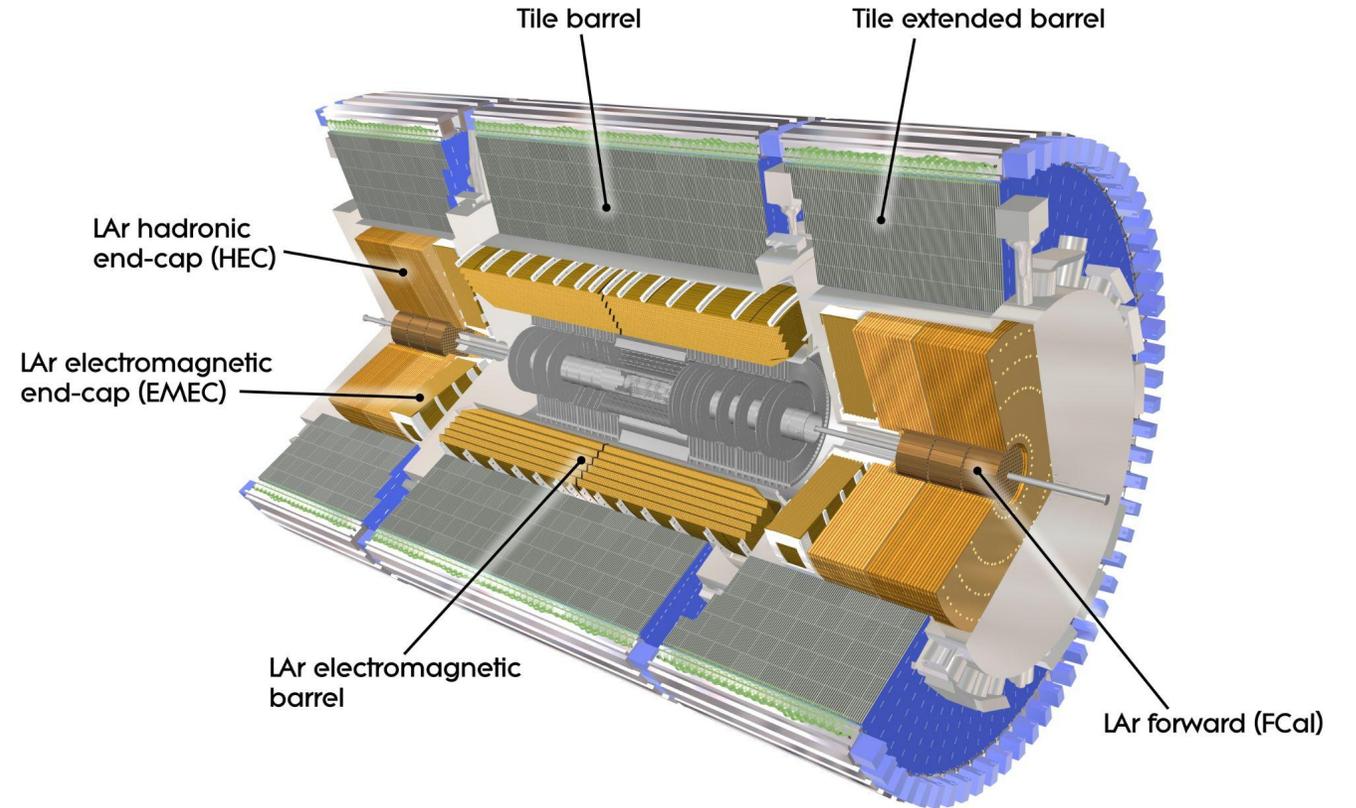


Calorimeters

Measure the energy of the particles through :

1. Electromagnetic showers
2. Hadronic showers

Cascading process



Discrete VAEs (in practice)

$$\frac{\partial}{\partial \phi} \mathbb{E}_{q(z|x, \phi)} [\log p(x|z, \theta)] \approx \frac{1}{N} \sum_{\rho \sim \mathcal{D}} \frac{\partial}{\partial \phi} \log p(x|f(x, \rho, \phi), \theta)$$

$$\mathbf{z}_i \sim \mathcal{N}(m(x, \phi), \sqrt{v(x, \phi)})$$

$$z = f(x, \phi, \rho) = m(x, \phi) + \sqrt{v(x, \phi)} \cdot \rho$$
$$\rho \sim \mathcal{N}(0, 1)$$

$$\mathbf{z}_i$$

$$p(z = 1) = \bar{q}$$

$$\bar{q} = \sigma(l)$$

$$z = f(x, \phi, \rho) = \mathcal{H}(\rho - (1 - \bar{q}(x, \phi)))$$
$$= \mathcal{H}(l(x, \phi) + \sigma^{-1}(\rho)), \rho \sim \mathcal{U}(0, 1)$$

Non-differentiable $\downarrow \frac{d\mathcal{H}(x)}{dx} = \delta(x)$

Gumbel-Softmax Trick $z = f(x, \phi, \rho) = \sigma\left(\frac{l(x, \phi) + \sigma^{-1}(\rho)}{\tau}\right)$

Towards QVAE

Classical RBM

$$p_{\boldsymbol{\theta}}(\mathbf{z}) \equiv e^{-E_{\boldsymbol{\theta}}(\mathbf{z})}/Z_{\boldsymbol{\theta}}, \quad Z_{\boldsymbol{\theta}} \equiv \sum_{\mathbf{z}} e^{-E_{\boldsymbol{\theta}}(\mathbf{z})},$$

$$E_{\boldsymbol{\theta}}(\mathbf{z}) = \sum_l z_l h_l + \sum_{l < m} W_{lm} z_l z_m, \quad \mathbf{h}, \mathbf{W} \in \{\boldsymbol{\theta}\}$$

Markov Chain Monte Carlo sampling

$$b_l = \beta_{eff}^* h_l, \quad W_{lm} = \beta_{eff}^* J_{lm}, \quad \Gamma_l = \beta_{eff}^* \Gamma^*,$$

$$\beta_{eff} \equiv B(s^*)/\beta_{phys}, \quad \Gamma^* \equiv A(s^*)/B(s^*). \quad (17)$$

Quantum BM

$$p_{\boldsymbol{\theta}}(\mathbf{z}) \equiv \text{Tr}[\Lambda_{\mathbf{z}} e^{-\mathcal{H}_{\boldsymbol{\theta}}}] / Z_{\boldsymbol{\theta}}, \quad Z_{\boldsymbol{\theta}} \equiv \text{Tr}[e^{-\mathcal{H}_{\boldsymbol{\theta}}}],$$

$$\mathcal{H}_{\boldsymbol{\theta}} = \sum_l \sigma_l^x \Gamma_l + \sum_l \sigma_l^z h_l +$$

$$+ \sum_{l < m} W_{lm} \sigma_l^z \sigma_m^z, \quad \boldsymbol{\Gamma}, \mathbf{h}, \mathbf{W} \in \{\boldsymbol{\theta}\},$$

Quantum Annealing

$$\mathcal{H}(s) = A(s) \sum_l \sigma_l^x + B(s) \left[\sum_l \sigma_l^z h_l + \sum_{l < m} J_{lm} \sigma_l^z \sigma_m^z \right]$$

DVAE to QVAE Challenges : Effective temperature estimation

Effective temperature β_{eff}^*

$$\begin{aligned} b_l &= \beta_{eff}^* h_l, & W_{lm} &= \beta_{eff}^* J_{lm}, & \Gamma_l &= \beta_{eff}^* \Gamma^*, \\ \beta_{eff} &\equiv B(s^*)/\beta_{phys}, & \Gamma^* &\equiv A(s^*)/B(s^*). \end{aligned} \quad (17)$$

Quantum annealer operates in a thermal environment

β_{eff}^* is not fixed

Need to perform real-time estimation

Quantum BM

$$\begin{aligned} p_{\theta}(\mathbf{z}) &\equiv \text{Tr}[\Lambda_{\mathbf{z}} e^{-\mathcal{H}_{\theta}}] / Z_{\theta}, & Z_{\theta} &\equiv \text{Tr}[e^{-\mathcal{H}_{\theta}}], \\ \mathcal{H}_{\theta} &= \sum_l \sigma_l^x \Gamma_l + \sum_l \sigma_l^z h_l + \\ &+ \sum_{l < m} W_{lm} \sigma_l^z \sigma_m^z, & \Gamma, \mathbf{h}, \mathbf{W} &\in \{\theta\}, \end{aligned}$$

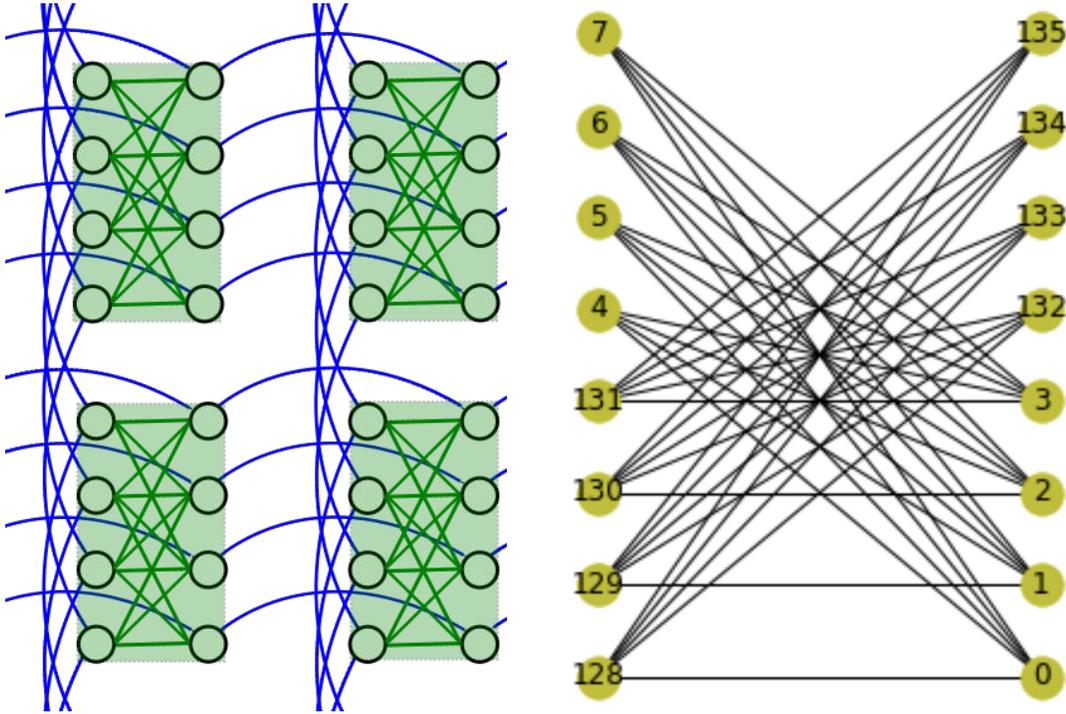


Quantum Annealing

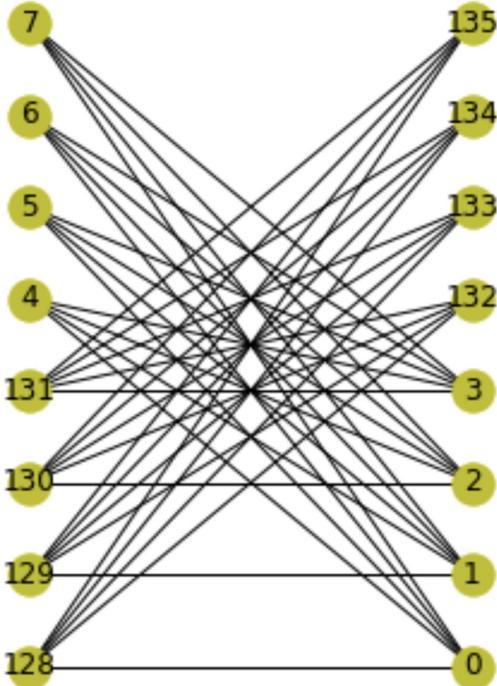
$$\mathcal{H}(s) = A(s) \sum_l \sigma_l^x + B(s) \left[\sum_l \sigma_l^z h_l + \sum_{l < m} J_{lm} \sigma_l^z \sigma_m^z \right]$$

DVAE to QVAE Challenges : Qubit Connectivity

Chimera connectivity (D'Wave 2000Q)

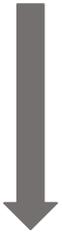


Maximum qubit connectivity : 6



Quantum BM

$$\begin{aligned}
 p_{\theta}(\mathbf{z}) &\equiv \text{Tr}[\Lambda_{\mathbf{z}} e^{-\mathcal{H}\theta}] / Z_{\theta}, \quad Z_{\theta} \equiv \text{Tr}[e^{-\mathcal{H}\theta}], \\
 \mathcal{H}_{\theta} &= \sum_l \sigma_l^x \Gamma_l + \sum_l \sigma_l^z h_l + \\
 &\quad + \sum_{l < m} W_{lm} \sigma_l^z \sigma_m^z, \quad \Gamma, \mathbf{h}, \mathbf{W} \in \{\theta\},
 \end{aligned}$$



Quantum Annealing

$$\mathcal{H}(s) = A(s) \sum_l \sigma_l^x + B(s) \left[\sum_l \sigma_l^z h_l + \sum_{l < m} J_{lm} \sigma_l^z \sigma_m^z \right]$$

DVAE to QVAE Challenges : QPU Resources

	Positron e^+	Photon γ	Pion π^+
Model Type	Model II	Model IV	Model IV
Learning Rate	10^{-4}	0.5×10^{-4}	10^{-4}
Epochs	100	100	100
Batch Size	100	100	100
Gibbs Steps	50	60	50
Latent smoothing temperature, τ_z	1/5	1/7	1/5
Output mask smoothing temperature, τ_{x_m}	1/5	1/5	1/9

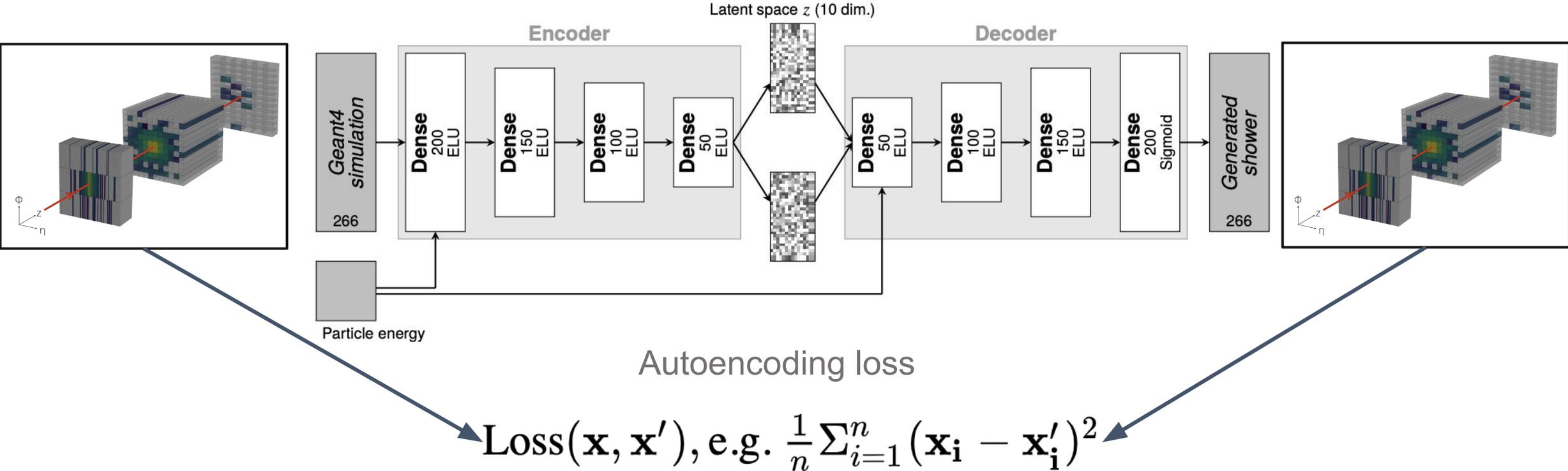
Parameter	Value
Model	Advantage
Graph Size	P16
Qubits	5436
Couplers	37440
Qubit Temperature (mK)	15.8 ± 0.5
M_{AFM}^1 (pH)	1.951
Average Single Qubit Thermal Width (Ising units)	0.196
FM Problem Freezeout (scaled time)	0.066
Single Qubit Freezeout (scaled time)	0.609
Annealing Time Range (μs)	1.0 to 2000.0
Readout Time Range ² (μs)	18.0 to 131.0
Programming Time ³ (μs)	~ 25100
Readout Error Rate ⁴	≤ 0.001

NAME (CHIP ID)	DESCRIPTION	DEFAULT ANNEALING TIME (μs)	Target - Samples - 20
Advantage_system11	Advantage system	DEFAULT PROGRAMMING THERMALIZATION (μs)	1000
QUBITS	SUPPORTED PROBLEM TYPES	DEFAULT READOUT THERMALIZATION (μs)	0
5760	ising, qubo	MAX ANNEAL SCHEDULE POINTS	12
TOPOLOGY	TAGS	MAX H GAIN SCHEDULE POINTS	20
[16] pegasus		NUMBER OF READS RANGE	1 to 10000
VFYC		PER QUBIT COUPLING RANGE	-18 to 15
false		PROBLEM RUN DURATION RANGE (μs)	0 to 1000000
ANNEAL OFFSET STEP		PROGRAMMING THERMALIZATION RANGE (μs)	0 to 10000
-0.00017565852000668507		READOUT THERMALIZATION RANGE (μs)	0 to 10000
ANNEALING TIME RANGE (μs)		H RANGE	-2 to 2
1 to 2000		H GAIN SCHEDULE RANGE	-4 to 4
ANNEAL OFFSET STEP PHIO		J RANGE	-1 to 1
0.00001486239425109832		EXTENDED J RANGE	-2 to 1

Per batch : 25.1 ms (Programming time) + [0.02 (Default annealing time) + 0.131 (Readout time)]*100 (Batchsize) ms = 40.2 ms (no delay)
 100000 events / particle type = 1000 batches

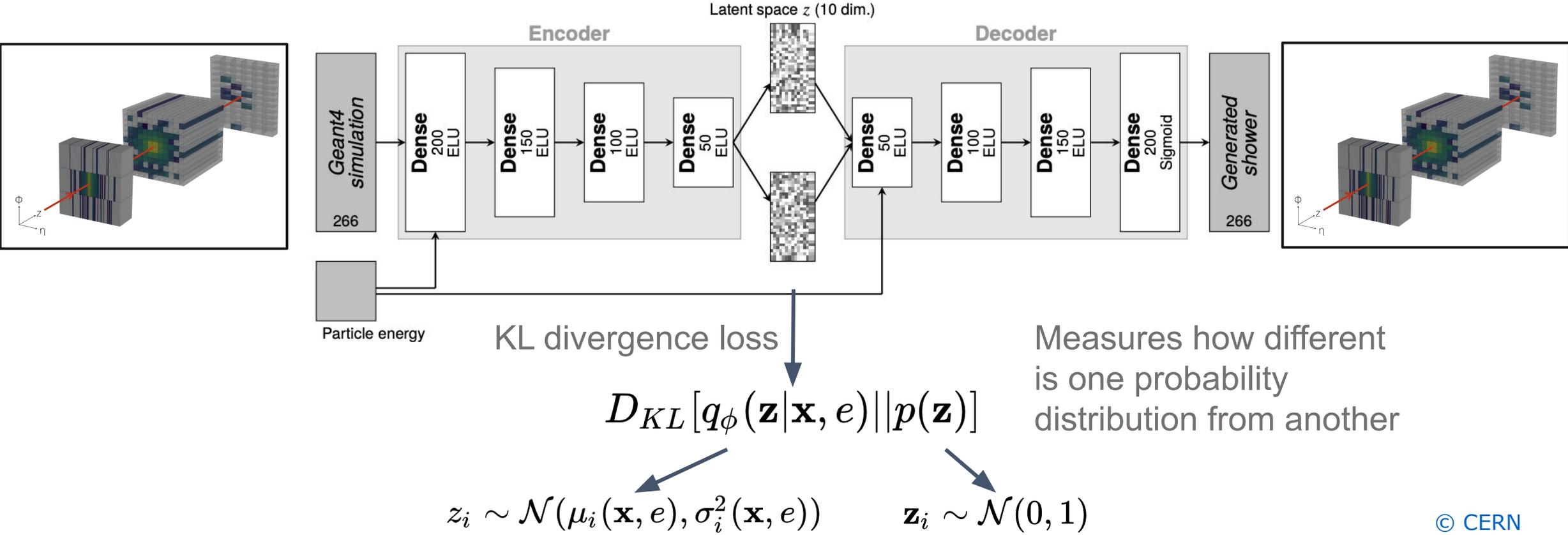
Per epoch : 40.2 s (no delay)
 Per particle type : 4020 s (no delay)

Previous Work : DGMs for Fast Calorimeter Simulation



[1] Paganini, de Oliveira and Nachman (2018), arXiv:1712.10321 [2] <https://atlas.web.cern.ch/Atlas/GROUPS/PHYSICS/PUBNOTES/ATL-SOFT-PUB-2018-001>.

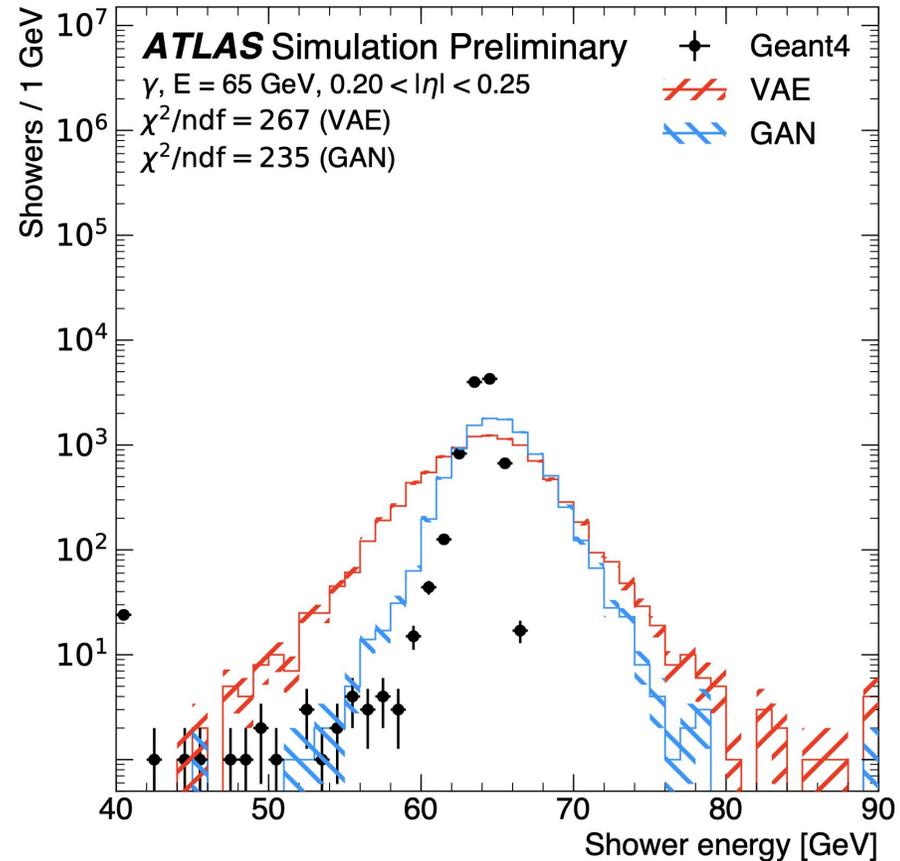
Previous Work : DGMs for Fast Calorimeter Simulation



[1] Paganini, de Oliveira and Nachman (2018), arXiv:1712.10321 [2] <https://atlas.web.cern.ch/Atlas/GROUPS/PHYSICS/PUBNOTES/ATL-SOFT-PUB-2018-001>.

Previous Work : DGMs for Fast Calorimeter Simulation

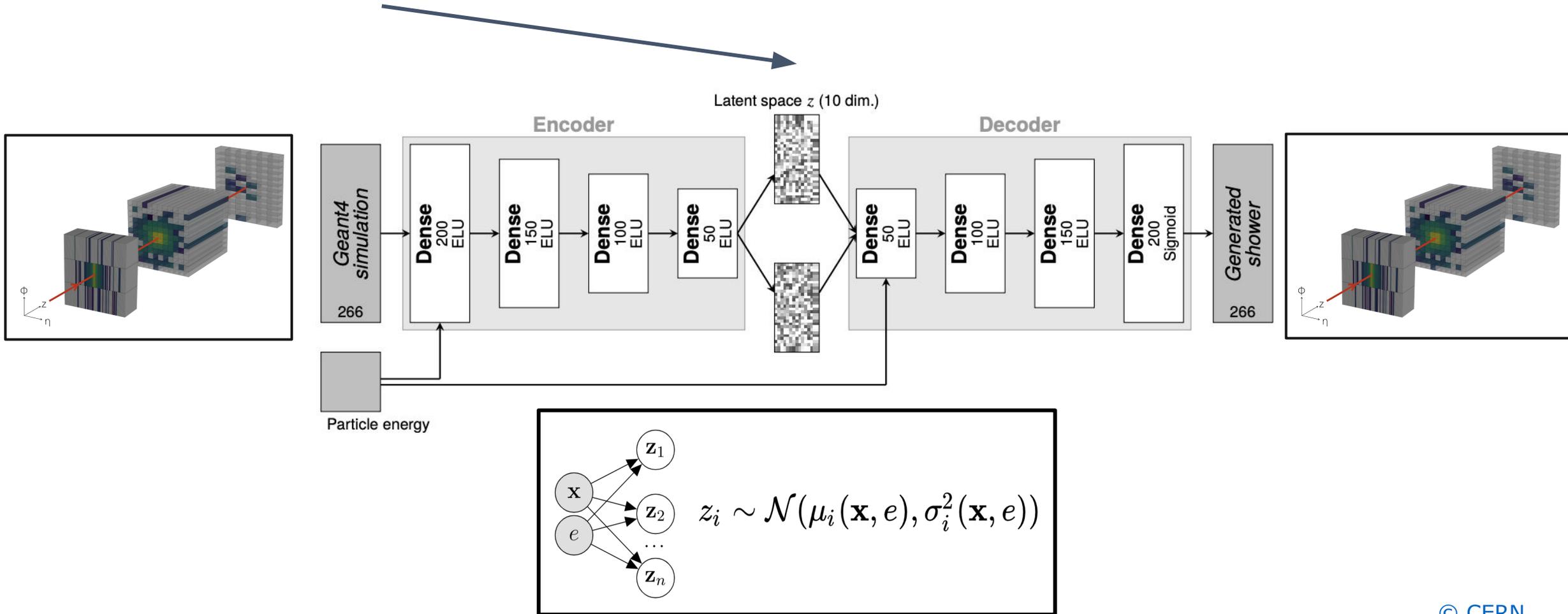
ATLAS-VAE/GAN [1]



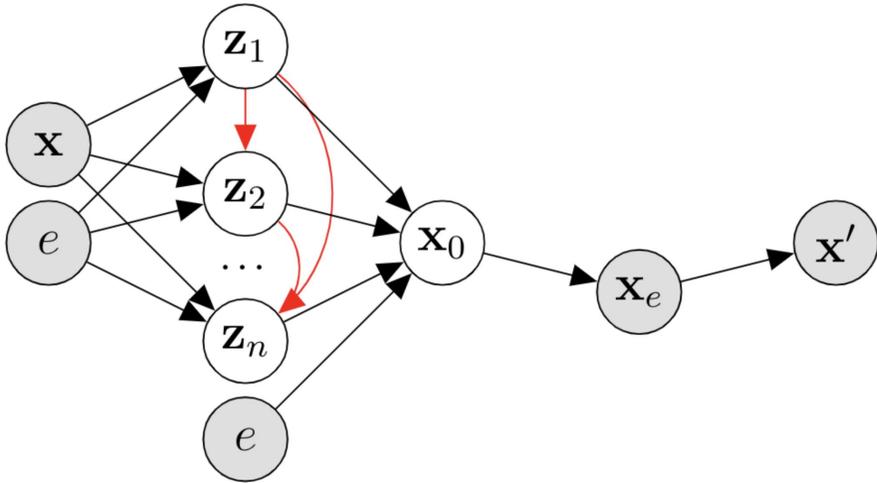
- Shape of the energy spectra of samples produced by DGMs matches
- No explicit energy conservation - DGMs produce showers with higher energy than incident particle energy

Previous Work : DGMs for Fast Calorimeter Simulation

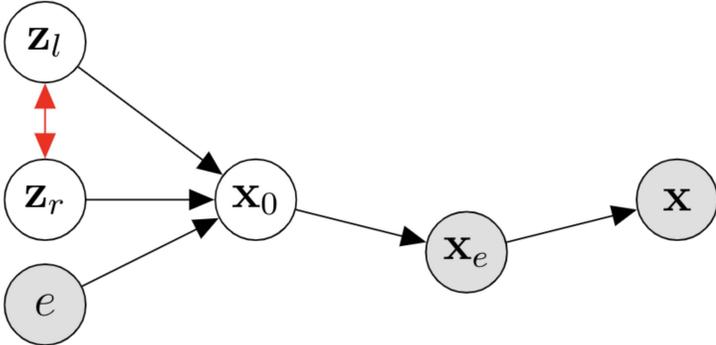
CaloGAN [1], ATLAS-VAE/GAN [2]



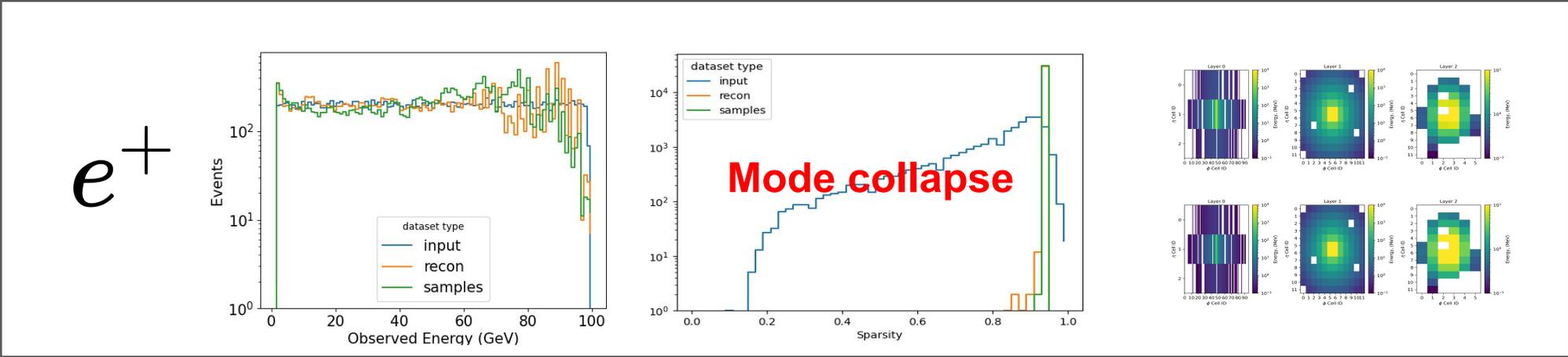
DVAEs for Calorimeter Simulation : Preliminary



Autoencoding model



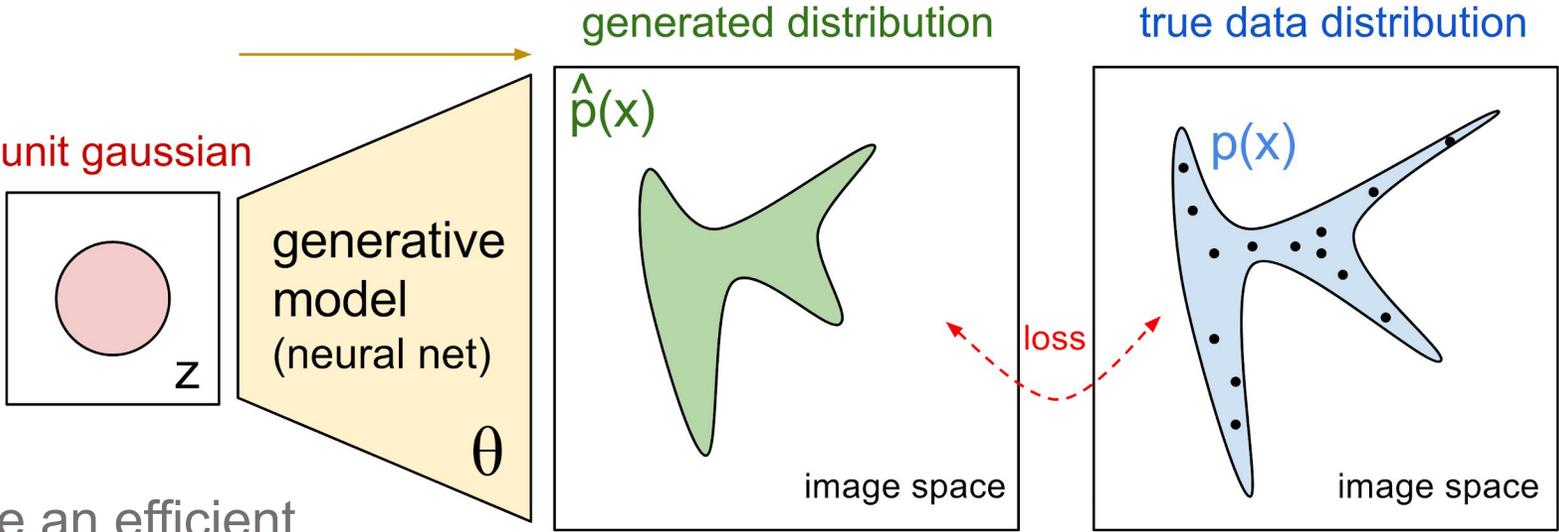
Generative model



Generative Models

1. Approximate true data distribution in a high-dimensional space

Maximize the log-likelihood of the data under the model distribution



2. Provide an efficient sampling method to sample from the model distribution

[1] <https://openai.com/blog/generative-models/>