Kinematic fitting for ParticleFlow Detectors at Future Higgs Factories

14th Annual Meeting of Helmholtz Alliance on "Physics at the Terascale"

Yasser Radkhorrami

DESY, Hamburg Universität Hamburg, Hamburg

November 24, 2021



CLUSTER OF EXCELLENCE QUANTUM UNIVERSE





Kinematic fitting for ParticleFlow Detectors at Future Higgs Factories | Yasser Radkhorrami | November 24, 2021

Page 2/15 👢





Conclusions

International Large Detector (ILD)

foreseen for the International Linear Collider (ILC)

- Polarized e^+e^- beams colliding at $\sqrt{s} = 250$ GeV (upgrades: 500GeV and 1TeV)
 - detailed knowledge of initial states
 - clean collision environment









Kinematic Fitting and Jet Error Parametrisation

Concept of ν -correction in a semi-leptonic decay

- ▶ Find heavy-quark jets: Identify b or c jet \rightarrow flavour tag
- ► Find semi-leptonic decay(s): Identify lepton in jet if present → possible using detector's high granularity
- Estimate neutrino energy from decay kinematics:
 - Assign B^0 or D^0 meson mass to mother hadron (m_X) .
 - Reconstruct flight direction of mother hadron (||,⊥) from position of primary and secondary vertex.
 - Calculate neutrino momentum: up to a sign ambiguity.

$$E_{\nu} = E_X - E_{vis} = \frac{E_{vis}E'_{vis} - \vec{p}_{vis\parallel} \cdot \vec{p}'_{vis\parallel}}{m_{vis}^2 + \vec{p}_{vis\perp}^2} m_X - E_{vis}$$

$$E'_{vis} = \frac{m_X^2 + m_{vis}^2}{2m_X} \qquad \vec{p}'_{vis\parallel} = \pm \sqrt{(\frac{m_X^2 - m_{vis}^2}{2m_X})^2 - \vec{p}_{vis\perp}^2}$$

$$(E'_{vis}, \vec{p}'_{vis}): \text{ visible 4-momentum in the rest frame of } X$$

As proof-of-principle: CHEAT from MC truth

The neutrino momentum can be determined up to a two-fold ambiguity

Can we use overall event kinematics to decide between solutions? \Rightarrow kinematic fit!





ν-correction with reconstructed information

reconstructed 4-momentum

Energy of parent hadron is distributed between:

- Charged decay products: reconstructed using tracker information:
 ⇒ rather easy to assign to 2nd vertex
 - lepton from semi-leptonic decay
 - other charged decay products
- Neutral Particles: reconstructed using calorimeter information:

 \Rightarrow difficult to assign to 2^{nd} vertex Many decays **without** extra neutral!

Neutrinos: missing energy, needs ν-correction

Charged particles profit excellent momentum resolution in ILD

Neutral particles have least contribution to the 4-momentum of parent hadron





Conclusion

ν -correction with reconstructed information

reconstructed flight direction of parent hadron

Probable scenarios:

- lepton is in 2nd vertex: direction from primary vertex to secondary vertex (~30% of SLD's)
- ▶ lepton + 3^{rd} vertex: lepton is not in 2^{nd} vertex, but a vertex exists in the jet (~40% of SLD's) \Rightarrow intersecting $track_{lep}$ and $\vec{p}_{3^{rd}vtx}$
- options for no vertex in jet (~17% of SLD's):
 - use jet axis as flight direction
 - use flight direction of leading particle in jet
 - use flight direction of lepton from semi-leptonic decay



Conclusion



Conclusion

v-correction with reconstructed information

- use reco 4-momentum mainly affected by:
 - σ_{p_T} for charged particles
 - $\sigma_{E_{CAL}}$ for neutral particles
- use reco flight direction mainly affected by:
 - ▶ l in 2^{nd} Vtx. / $l+3^{rd}$ Vtx.: vertex finding algorithm and detector IP resolution $(\sigma_{d_0} \text{ and } \sigma_{z_0})$
 - ▶ no $2^{nd}/3^{rd}$ Vtx.: hadrons decay so close to the IP, no chance to get the vertex! ⇒ use jet axis
- overall uncertainty $(\sigma_{E_{\nu}}, \sigma_{\theta_{\nu}} \text{ and } \sigma_{\phi_{\nu}})$ can be parameterized step-by-step
- still cheated: lepton ID, particle assignment (ongoing)





DESY

Kinematic fit

- > Kinematic fit: adjustment of measured quantities under certain kinematic constraints:
 - Energy and momentum conservation
 - Invariant masses of particles



Exploit well-known initial state in e^+e^- colliders

 \Rightarrow need error parametrization, in particular for jets

• Minimize χ^2 :

$$\chi^2(\boldsymbol{a},\boldsymbol{\xi},\boldsymbol{f}) = (\boldsymbol{\eta}-\boldsymbol{a})^T \boldsymbol{V}^{-1}(\boldsymbol{\eta}-\boldsymbol{a}) - 2\boldsymbol{\lambda}^T \boldsymbol{f}(\boldsymbol{a},\boldsymbol{\xi})$$

Conclusions

- η : vector of measured kinematic variables (x)
- a: vector of fitted quantities
- $\pmb{\xi}$: vector of unmeasured kinematic variables
- V: covariance matrix
- λ : Lagrange multipliers
- $f(a, oldsymbol{\xi})$: vector of constraints



Conclusion

Jet specific energy resolution

Parametrize sources of uncertainties (assumed uncorrelated) in jet energy measurements (ErrorFlow):

 $\sigma_{E_{jet}} = \sigma_{Det} \oplus \sigma_{Conf} \oplus \sigma_{\nu} \oplus \sigma_{Clus} \oplus \sigma_{Had}$

 \triangleright σ_{Det} : Detector resolution using track and cluster parameters

- ► σ_{Conf} : Particle confusion in Particle Flow Algorithm Estimated based on jet energy and neutral hadron / photon energy fractions
- > σ_{ν} : Semi-leptonic decays: error propagation from neutrino correction currently none since cheating
- \blacktriangleright σ_{Clus} : Misassignment of particles in the jet clustering, has not been included yet
- \triangleright σ_{Had} : Mismodeling of QCD effects in parton shower and hadronization, has not been included yet



DESY-THESIS-2017-045

Conclusions

Uncertaities in jet-level: Energy

Propagation of errors from PFOs to jets:

- Transform the covariance matrix of each PFO (*E*,*x*,*y*,*z* for clusters, track parameters for charged) to (*E*,*p_x*,*p_y*,*p_z*)
- Add up covariance matrices of all PFOs
- Add confusion term for jet energy
 - calculate using jet energy composition
- Transform to (E,θ,ϕ,m)





Confusion term improves the estimate of the jet energy uncertainty, but not quite enough \Rightarrow need adjustment \Rightarrow use scaling factor 1.2 in Kinematic fit



Introduction

Application of kinematic fit to $e^+e^- \rightarrow ZH \rightarrow \mu \bar{\mu} b \bar{b}$ events

Parameters of jets and leptons are variated within their uncertainties to satisfy 5 constraints: Conservation of momentum (hard constraints):

- ► p_x : e^+e^- crossing angle: 14 mrad $\Sigma p_x = \sqrt{s} \times \sin 0.007 \approx 1.75$ GeV
- $\triangleright p_y: \ \Sigma p_y = 0$

$$\blacktriangleright p_z: \ \Sigma p_z = 0$$

Conservation of total energy (hard constraint):

►
$$E_{lab} = 2\sqrt{(\frac{\sqrt{s}}{2})^2 + (\Sigma p_x)^2}$$

Constrain di-muon mass to agree with m_Z within its natural width (soft constraint):

$$\blacktriangleright$$
 $m_Z=91.2~{
m GeV}$, $\sigma_{m_Z}=rac{2.5}{2}$



Conclusion





DESY.

Conclusion

Kinematic fit performance in $e^+e^- \rightarrow ZH \rightarrow \mu\bar{\mu}b\bar{b}$ at $\sqrt{s} = 250$ GeV

without semi-leptonic decays

pull distribution



fit probability



Improved kinematic fit performance with full CovMat of jets + scaled jet energy uncertainty



Conclusions

Kinematic fit performance in $e^+e^- \rightarrow ZH \rightarrow \mu\bar{\mu}b\bar{b}$ at $\sqrt{s} = 250$ GeV (cntd.)

without semi-leptonic decays



Improved kinematic fit performance with full CovMat of jets + scaled jet angular uncertainties



Conclusion

Z/Higgs mass reconstruction in presence of SLDs

- ν -correction alone recovers Z/H mass:
 - resolve sign ambiguity by kinematic fit
 - pre-fit (E_{ν}, \vec{p}_{ν}) is used
- striking improvement from new jet error parametrisation fed to kinematic fit even without v-correction
- significant further improvement by combined kinematic fit and *v*-correction (fully cheated), especially for the Higgs peak
- less powerful ν -correction \Rightarrow performance is expected between Green and Red







Conclusions

- \blacktriangleright Higgs mass reconstruction essential eg in ZZH vs ZHH separation (Higgs self-coupling measurement)
- Heavy flavour jets are essential for Higgs physics
- Correction of semi-leptonic decays of heavy flavour jets is important for Higgs mass reconstruction
 - Neutrino momentum can be reconstructed up to a sign ambiguity
 - Ambiguity can be resolved by kinematic fit
 - Next: remove the partial cheating from the neutrino correction
- Kinematic fit exploits well-known initial state in e⁺e⁻ colliders and requires excellent understanding of jet measurement
- ▶ ILD as a Particle Flow detector provides full detail for estimating jet measurement uncertainties
- \blacktriangleright huge improvement on $Z/H \rightarrow b\bar{b}$ di-jet mass reconstruction has been achieved







BACKUP







DESY

International Large Detector (ILD)

Momentum Resolution



arXiv:2003.01116

▶ Jet Energy Resolution $(E_{PFO} + E_{\nu}^{MC})$





DESY

International Large Detector (ILD)





 \blacktriangleright Impact Parameter Resolution, z_0

arXiv:2003.01116





Mis-reconstruction of bb invariant mass due to missing neutrino energy from semi-leptonic decays

Backup

Can the missing momentum be retrieved from event and decay kinematics?





correcting neutrino energy

4-vector based approach

 \blacktriangleright (*E*, \vec{p})-based approach



$$\begin{split} \vec{p}_{\nu,\perp} &= -\vec{p}_{vis,\perp} \\ \vec{p}_{\nu}, \parallel &= \frac{1}{2D} (-A \pm \sqrt{A^2 - BD}) \hat{n} \\ A &= p_{vis,\parallel} (2p_{vis,\perp}^2 + m_{vis}^2 - m_X^2) \\ B &= 4p_{vis,\perp}^2 E_{vis}^2 - (2p_{vis',\perp}^2 + m_{vis}^2 - m_X^2)^2 \\ D &= E_{vis}^2 - p_{vis,\parallel}^2 \end{split}$$

$$\hat{n} = \frac{p_{v\vec{is},\parallel}}{|\vec{p_{v\vec{is},\parallel}}|}$$

The neutrino momentum can be determined up to a two-fold ambiguity

Kinematic fitting for ParticleFlow Detectors at Future Higgs Factories | Yasser Radkhorrami | November 24, 2021

 closure test: apply correction with fully cheated information and compare with true neutrino energy





Correcting neutrino energy

Rapidity based approach

Rapidity under Lorents-transformations ~ velocity under Galileo-transformations: $\omega = \omega_X + \omega'$; $\omega = \frac{1}{2} ln \frac{E + p'_{\parallel}}{E - p'_{\parallel}}$

 ω : rapdity in lab frame , ω' : rapdity in rest frame of X , ω_X : rapdity of X in lab frame





DESY.

$\nu\text{-}correction$ with reconstructed information

reconstructed flight direction of parent hadron





Backup

$\nu\text{-}correction$ with reconstructed information

- use reco 4p mainly affected by:
 - σ_{p_T} for charged particles
 - $\sigma_{E_{CAL}}$ for neutral particles
- use reco flight direction mainly affected by vertex finding algorithm and detector IP resolution (σ_{d_0} and σ_{z_0})
- overall uncertainty $(\sigma_{E_{\nu}}, \sigma_{\theta_{\nu}} \text{ and } \sigma_{\phi_{\nu}})$ can be parameterized step-by-step
- still cheated: lepton ID, particle assignment (ongoing)





ErrorFlow: Jet Error Parametrisation from Particle Flow Objects (PFO) Energy

Error estimation in PFO level:

- Photons: energy error is perfectly modeled. (sigma ~ 1)
- Charged PFOs: uncertainties propagated from track fit covariance matrix
 - uncertainties 30% too small
 - possible future improvement from track refitting with specific mass hypothesis after particle ID
- Neutral Hadrons: energy and energy error are significantly overestimated. work on improvement in progress.









ErrorFlow: Jet Error Parametrisation from Particle Flow Objects (PFO) Angles

The angular uncertainties obtained directly from track parameters / cluster position errors



 \Rightarrow Scale σ_{θ} and σ_{ϕ} by factor ~ 1.3 (for photons) and ~ 1.8 (for neutral hadrons)





DESY

Uncertaities in jet-level: θ & ϕ



Jet angular uncertainties need scaling factor ${\sim}1.6$





Event selection

Select $e^+e^- \rightarrow ZH \rightarrow \mu \bar{\mu} b \bar{b}$ events at $\sqrt{s} = 250$ GeV with (exactly) 2-leptons + 2-jets final state:

- IsolatedLeptonTagging Training for the IDR 500 GeV samples is used,
 - 1. Lepton ID: μ^{\pm} Deposited energy in subdetectors
 - 2. Vertex: primary or secondary Significance of impact parameters (d_0, z_0)
 - 3. Isolated: not belong to jets

FastJetProcessor

- Exclusive k_t (Durham) algorithm (no overlay)
- Find smallest of (d_{ij}, d_{iB}) $d_{ij} = 2 \min(E_i^2, E_j^2)(1 - \cos \theta_{ij})$ i,j: particles, B: Beam
- $d_{ij} < d_{iB}$: combine i&j as pseudojet(p): $p_i + p_j$
- $d_{iB} < d_{ij}$: remove particle *i* from list
- Repeat iteration until d_{ij} or $d_{iB} > d_{cut}$ (threshold)



IsolatedLeptonTagging has not been trained for new software at 250 GeV yet!



Page 12/27



event selection

separate Higgs decay modes: $H \to b \bar{b}$, cheat from MCTruth



 $\frac{2}{3}$ of $b\bar{b}$ jets contain at-least one semi-leptonic decay \Rightarrow Frequent $H \rightarrow b\bar{b}$ needs neutrino correction. Kinematic fitting for ParticleFlow Detectors at Future Higgs Factories | Yasser Badkhorrami | November 24, 2021 | Page 13/22



DESY

Neutral PFO identification by Pandora



Majority of identified photons are true photons.

No explicit decision for mass of identified neutral hadrons due to their multiplicity.





Pandora treatment with Neutral Hadrons

What Pandora does:

- Cluster energy is assigned to PFO(massless) energy $E_{PFO} = |\vec{p}_{PFO}| = E_{cluster}$
- Neutral Hadrons are identified as neutron
- neutron mass is set for PFO \Rightarrow incosistent 4-momentum!
- ► CovMat of Neutral PFO is calculated (using inconsistent 4-momentum): CovMat $(\vec{p}, E) = J^T$ CovMat $(\vec{x}_{clu}, E_{clu}) J$

$$I = \begin{pmatrix} \frac{\partial p_x}{\partial x_c} & \frac{\partial p_y}{\partial x_c} & \frac{\partial p_z}{\partial x_c} & \frac{\partial E}{\partial x_c} \\ \frac{\partial p_x}{\partial y_c} & \frac{\partial p_y}{\partial y_c} & \frac{\partial p_z}{\partial y_c} & \frac{\partial E}{\partial y_c} \\ \frac{\partial p_x}{\partial z_c} & \frac{\partial p_y}{\partial z_c} & \frac{\partial p_z}{\partial z_c} & \frac{\partial E}{\partial z_c} \\ \frac{\partial p_x}{\partial E_c} & \frac{\partial p_y}{\partial E_c} & \frac{\partial p_z}{\partial E_c} & \frac{\partial E}{\partial E_c} \end{pmatrix}$$

 $CovMat(\vec{p}, E)$ of Neutral PFOs depend on the mass assumption.

Suggestion: Take consistent 4-momentum of massive neutral hadrons for CovMat calculations.

۰.





CovMat of Neutral PFOs

- Current CovMat calculation (MarlinReco/Analysis/AddClusterProperties) $E_{PFO} = |\vec{p}_{PFO}| = E_{clu} , p_x = E_{clu} \frac{x}{r} , p_y = E_{clu} \frac{y}{r} , p_z = E_{clu} \frac{z}{r}$
- Alternative CovMat calculation (taking consistent 4-momentum of neutral hadrons)

$$E_{PFO} = \sqrt{|\vec{p}_{PFO}|^2 + m_{PFO}^2} = \sqrt{E_{clu}^2 + m_{r}^2}$$

$$J = \begin{pmatrix} E_{clu} \frac{r^2 - x^2}{r^3} & -E_{clu} \frac{xy}{r^3} & -E_{clu} \frac{xz}{r^3} & 0\\ -E_{clu} \frac{xy}{r^3} & E_{clu} \frac{r^2 - y^2}{r^3} & -E_{clu} \frac{yz}{r^3} & 0\\ -E_{clu} \frac{xz}{r^3} & -E_{clu} \frac{yz}{r^3} & E_{clu} \frac{r^2 - z^2}{r^3} & 0\\ \frac{x}{r} & \frac{y}{r} & \frac{z}{r} & 1 \end{pmatrix} \rightarrow J = \begin{pmatrix} E_{clu} \frac{xy}{r^3} & -E_{clu} \frac{xy}{r^3} & -E_{clu} \frac{xz}{r^3} & 0\\ -E_{clu} \frac{xy}{r^3} & E_{clu} \frac{r^2 - z^2}{r^3} & -E_{clu} \frac{yz}{r^3} & 0\\ -E_{clu} \frac{xz}{r^3} & -E_{clu} \frac{yz}{r^3} & E_{clu} \frac{r^2 - z^2}{r^3} & 0\\ \frac{E_{clu} \frac{xz}{r}}{r} & \frac{E_{clu}}{r} \frac{x}{r} & \frac{z}{r} & 1 \end{pmatrix}$$

using error propagation, PFO angular uncertainties are calculated directly from cluster position error: $\sigma_{\theta}^{2} = \left(\frac{\partial\theta}{\partial x}\right)^{2}\sigma_{x}^{2} + \left(\frac{\partial\theta}{\partial y}\right)^{2}\sigma_{y}^{2} + \left(\frac{\partial\theta}{\partial z}\right)^{2}\sigma_{z}^{2} + \frac{\partial\theta}{\partial x}\frac{\partial\theta}{\partial y}\sigma_{xy} + \frac{\partial\theta}{\partial x}\frac{\partial\theta}{\partial z}\sigma_{xz} + \frac{\partial\theta}{\partial y}\frac{\partial\theta}{\partial z}\sigma_{yz}$ $\sigma_{\phi}^{2} = \left(\frac{\partial\phi}{\partial x}\right)^{2}\sigma_{x}^{2} + \left(\frac{\partial\phi}{\partial y}\right)^{2}\sigma_{y}^{2} + \frac{\partial\phi}{\partial x}\frac{\partial\phi}{\partial y}\sigma_{xy}$

MUST: angular and energy uncertainties remain unchanged!





CovMat of Jets

- ▶ AddClusterProperties/FourMomentumCovMat: $CovMat(cluster/track) \rightarrow CovMat(\vec{p}, E)$
 - Current CovMat calculation (inconsistent 4-momentum of neutral hadrons):

 $E_{PFO}=|\vec{p}_{PFO}|=E_{clu}$, $p_x=E_{clu}\frac{x}{r}$, $p_y=E_{clu}\frac{y}{r}$, $p_z=E_{clu}\frac{z}{r}$, $m_{PFO}=m_n$

Alternative CovMat calculation (taking consistent 4-momentum of neutral hadrons)

$$E_{PFO} = \sqrt{|\vec{p}_{PFO}|^2 + m_{PFO}^2} = \sqrt{E_{clu}^2 + m_n^2} \frac{J_{(wrong)} \rightarrow J_{(right)}}{J_{(right)}}$$

$$\begin{pmatrix} E_{clu} \frac{r^2 - x^2}{r^3} & -E_{clu} \frac{xy}{r^3} & -E_{clu} \frac{xz}{r^3} & 0 \\ -E_{clu} \frac{xy}{r^3} & E_{clu} \frac{r^2 - y^2}{r^3} & -E_{clu} \frac{yz}{r^3} & 0 \\ -E_{clu} \frac{xz}{r^3} & -E_{clu} \frac{yz}{r^3} & E_{clu} \frac{r^2 - z^2}{r^3} & 0 \\ -E_{clu} \frac{xz}{r^3} & -E_{clu} \frac{yz}{r^3} & E_{clu} \frac{r^2 - z^2}{r^3} & 0 \\ -E_{clu} \frac{xz}{r^3} & -E_{clu} \frac{yz}{r^3} & E_{clu} \frac{r^2 - z^2}{r^3} & 0 \\ -E_{clu} \frac{xz}{r^3} & -E_{clu} \frac{yz}{r^3} & E_{clu} \frac{r^2 - z^2}{r^3} & 0 \\ -E_{clu} \frac{xz}{r^3} & -E_{clu} \frac{yz}{r^3} & E_{clu} \frac{r^2 - z^2}{r^3} & 0 \\ -E_{clu} \frac{xz}{r^3} & -E_{clu} \frac{yz}{r^3} & E_{clu} \frac{r^2 - z^2}{r^3} & 0 \\ -E_{clu} \frac{xz}{r^3} & -E_{clu} \frac{yz}{r^3} & E_{clu} \frac{r^2 - z^2}{r^3} & 0 \\ -E_{clu} \frac{xz}{r^3} & -E_{clu} \frac{yz}{r^3} & E_{clu} \frac{r^2 - z^2}{r^3} & 0 \\ -E_{clu} \frac{xz}{r^3} & -E_{clu} \frac{yz}{r^3} & E_{clu} \frac{r^2 - z^2}{r^3} & 0 \\ -E_{clu} \frac{xz}{r^3} & -E_{clu} \frac{yz}{r^3} & E_{clu} \frac{r^2 - z^2}{r^3} & 0 \\ -E_{clu} \frac{xz}{r^3} & -E_{clu} \frac{yz}{r^3} & E_{clu} \frac{r^2 - z^2}{r^3} & 0 \\ -E_{clu} \frac{xz}{r^3} & -E_{clu} \frac{yz}{r^3} & E_{clu} \frac{r^2 - z^2}{r^3} & 0 \\ -E_{clu} \frac{xz}{r^3} & -E_{clu} \frac{yz}{r^3} & E_{clu} \frac{r^2 - z^2}{r^3} & 0 \\ -E_{clu} \frac{xz}{r^3} & -E_{clu} \frac{yz}{r^3} & E_{clu} \frac{r^2 - z^2}{r^3} & 0 \\ -E_{clu} \frac{xz}{r^3} & -E_{clu} \frac{yz}{r^3} & E_{clu} \frac{r^2 - z^2}{r^3} & 0 \\ -E_{clu} \frac{xz}{r^3} & -E_{clu} \frac{yz}{r^3} & E_{clu} \frac{r^2 - z^2}{r^3} & 0 \\ -E_{clu} \frac{xz}{r^3} & -E_{clu} \frac{yz}{r^3} & E_{clu} \frac{z}{r^3} & 0 \\ -E_{clu} \frac{xz}{r^3} & -E_{clu} \frac{yz}{r^3} & E_{clu} \frac{z}{r^3} & 0 \\ -E_{clu} \frac{xz}{r^3} & -E_{clu} \frac{yz}{r^3} & E_{clu} \frac{z}{r^3} & 0 \\ -E_{clu} \frac{xz}{r^3} & -E_{clu} \frac{yz}{r^3} & E_{clu} \frac{z}{r^3} & 0 \\ -E_{clu} \frac{x}{r^3} & -E_{clu} \frac{yz}{r^3} & E_{clu} \frac{z}{r^3} & 0 \\ -E_{clu} \frac{x}{r^3} & -E_{clu} \frac{yz}{r^3} & 0 \\ -E_{c$$

ErrorFlow:

$$\mathsf{CovMat}(\vec{p}_{jet}, E_{jet}) = \sum_{PFO} \mathsf{CovMat}(\vec{p}, E) \quad : \quad \sigma_{E_{jet}}^2 = \sigma_{conf}^2 + \sum_{PFO} \sigma_{E_{PFO}}^2$$

MarlinKinfitProcessors:

$$\mathsf{CovMat}(ec{p}_{jet}, E_{jet}) o (\sigma_{ heta_{jet}}$$
 , $\sigma_{\phi_{jet}}$, $\sigma_{E_{jet}}$





ErrorFlow: Jet Error Parametrisation from Particle Flow Objects (PFO) Angles

The angular uncertainties obtained directly from track parameters / cluster position errors



 \Rightarrow Scale σ_{θ} and σ_{ϕ} by factor ~ 1.3 (for photons) and ~ 1.8 (for neutral hadrons)





DESY

Uncertaities in jet-level: θ & ϕ



Jet angular uncertainties need scaling factor ${\sim}1.6$





Neutrino correction hypothesis

- Assign semi-leptonic decays to jets
- Add neutrino momentum to 4-momentum of assigned jet:

Test three hypothesis for neutrino energy in each semi-leptonic decay: E_{ν}^{+} , E_{ν}^{-} , 0 3^{nSLD} combination of E_{ν} 's for adding to jet 4-momentum: Number of semileptonic decays in a jet: nSLD = nSLDB + nSLDC Example:

If an event contains two jets: jet-1 contains 2 semi-leptonic decays and jet-2 contains 1 semi-leptonic decay, $27(=3^2 \times 3^1)$ combinations of E_{ν} 's are available for neutrino correction in the event:

jet-1:

DESY

comb.	1	2	3	4	5	6	7	8	9
$\vec{p}_{\nu,1}$	-	+	0	-	+	0	-	+	0
$ec{p}_{ u,2}$	-	-	-	+	+	+	0	0	0

jet-2:

comb.	1	2	3
$\vec{p}_{ u,3}$	-	+	0

 $\vec{p}_{\nu,1} + \vec{p}_{\nu,2}$ is added to 4-momentum of jet-1 and $\vec{p}_{\nu,3}$ is added to 4-momentum of jet-2. $\vec{p}_{\nu,1} + \vec{p}_{\nu,2} + \vec{p}_{\nu,3} = 0$ allows fitter to neglect neutrino correction Combination with highest fit probability is chosen as best neutrino correction.



Backup

Simple neutrino correction for Higgs mass reconstruction

Neutrino energy correction:

Estimating neutrino energy as a fraction of corresponding lepton energy:

$$E_{jet}^{corr} = E_{jet} + E_{\nu} = E_{jet} + (\frac{1}{x} - 1)E_{lep}$$

Uncertainty on jet energy parametrised as:

$$\sigma_{E_{jet}}^{corr} = \frac{100\%}{\sqrt{E_{jet}}} \oplus \sigma_{\nu}$$
$$\sigma_{\nu}^2 = \left(\frac{\sigma_{\langle x \rangle}}{\langle x \rangle^2}\right)^2 E_{lep}^2 + \left(\frac{1}{x} - 1\right) \Delta E_{lep}^2$$

Fixed uncertainties on angles:

$$\Delta \theta_{jet} = \Delta \phi_{jet} = 100 \, \mathrm{mrad}$$



Blue: before neutrino energy correction Orange: After neutrino energy correction

Simple correction to jet energy improves jet energy pull distribution as a measure of fit performance. Kinematic fitting for ParticleFlow Detectors at Future Higgs Factories | Yasser Radkhorrami | November 24, 2021 | Page 21/27



DESY

Simple neutrino correction for Higgs mass reconstruction



 \blacktriangleright Bias and assymetry in m_H is removed by correcting jet energy and adding ISR



Backup

Error flow and application in kinematic fit

Jet specific energy resolution for $e^+e^-\to ZH\to q\bar{q}b\bar{b}$ process at $\sqrt{s}=350~{\rm GeV}$

- Full 4×4 CovMatrix on 4-momentum of jets $\sigma(\vec{p}, E)$:
 - \blacktriangleright $\sigma_{Det}:$ computed using subdetector momentum/energy resolution
 - σ_{Conf}: computed using jet energy and particle content (charged, neutral and photon)
 - $\bullet \sigma_{\nu} = 0.73.E_l$
 - \triangleright σ_{Had} , σ_{Clus} are not accounted for error flow procedure yet.

Fixed (and wide) angular resolution: $\sigma_{\theta} = \sigma_{\phi} = 100$ mrad Kinematic fit: vary jet quantities (E, θ, ϕ) within uncertainties $(\sigma_E, \sigma_{\theta}, \sigma_{\phi})$ Improved fit probability by applying Error Flow on jet energy



DESY-THESIS-2017-045

 \Rightarrow Further improvements by error parametrization and handling sl-decays





fit constraints

momentum conservation: p_x

 ISR is initialized to satisfy momentum conservation on x direction

by error flow on jet energy



angue resolution for individual jets: improved constraint on momentum conservation Kinematic fitting for ParticleFlow Detectors at Future Higgs Factories | Yasser Radkhorrami | November 24, 2021 |





fit constraints

momentum conservation: p_y

 ISR is initialized to satisfy momentum conservation on z direction

by error flow on jet energy



angue resolution for individual jets: improved constraint on momentum conservation Kinematic fitting for ParticleFlow Detectors at Future Higgs Factories | Yasser Radkhorrami | November 24, 2021 |

Page 25/27



DESY

Fit constraints

Momentum conservation: p_z

Adding 4-momentum of neutrino improves jet fit object initialization

DBD 350/500 GeV samples



Proproperties of the proprovide constraint on momentum

Kinematic fitting for ParticleFlow Detectors at Future Higgs Factories | Yasser Radkhorrami | November 24, 2021

MC-2020 250 GeV prod. samples





fit constraints

energy conservation: E

Neutrino correction (best pre-fit \vec{p}_{ν} for succesful fits) improves start values \Rightarrow better fit object initialization

► DBD 350/500 GeV samples



By neutrino correction, initial value of constraint function closer to target \Rightarrow fit should work better! Kinematic fitting for ParticleFlow Detectors at Future Higgs Factories | Yasser Radkhorrami | November 24, 2021 | Page 27/27