

# High quality GW relics

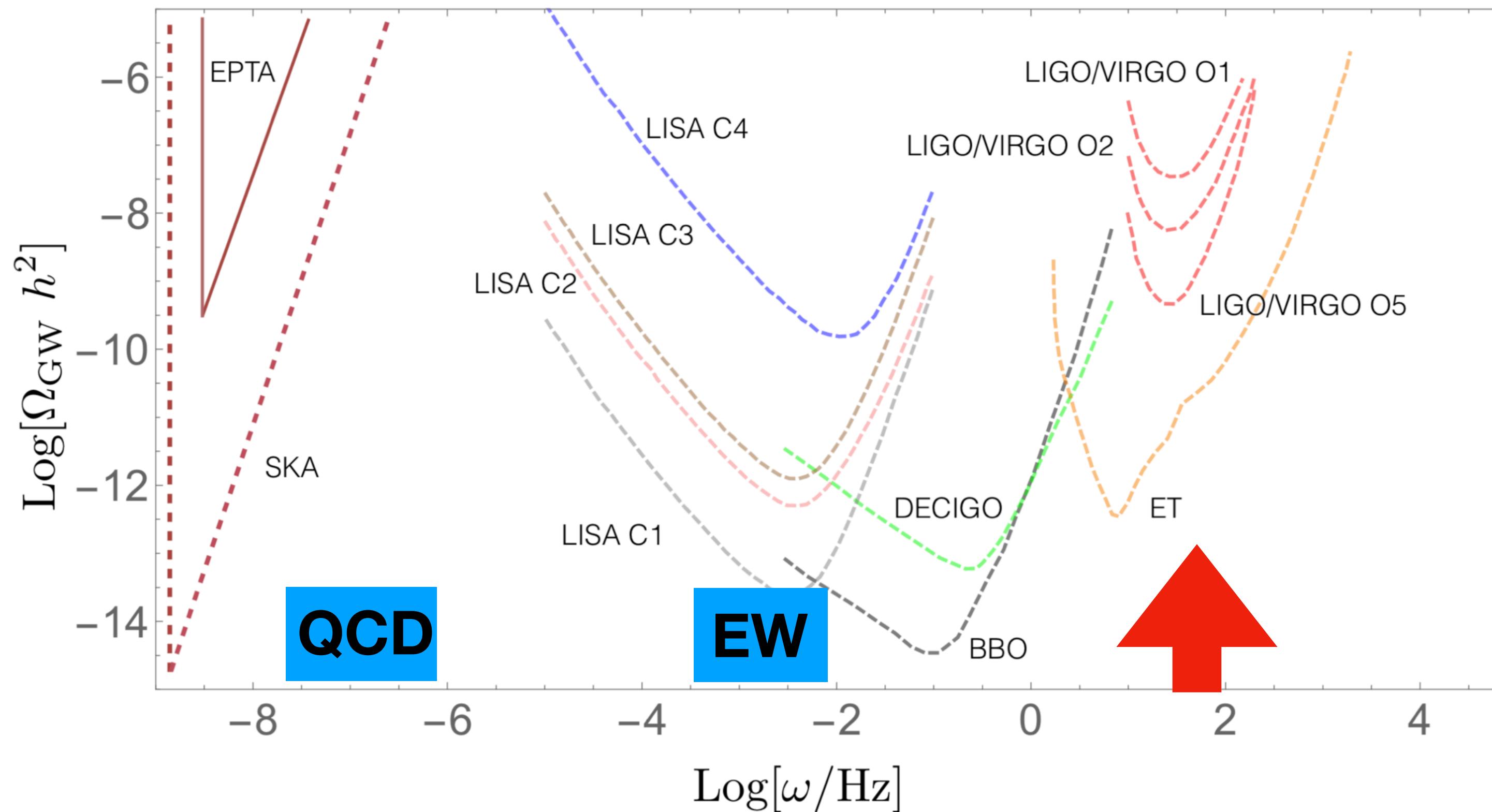
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The Wolfgang Pauli Centre  
8-12 november 2021

# GWs

- LIGO/Virgo/Kagra (and coming): new window, with impressive views!
- From HEP perspective: **relic stochastic GWB**



# GWs

- LIGO/Virgo/Kagra (and coming): new window, with impressive views!
- From HEP perspective: **relic stochastic GWB**
- E.g., LVK sensitive to cosmological phase transitions at  **$T \sim 10^8 \text{ GeV}$** 
  - > close to a *usual suspect*: **Peccei-Quinn scale,  $10^9 - 10^{11} \text{ GeV}$**

**Can we probe axion physics with LVK, already??**

# plan

- Intro

- I) GWs from a 1st order Peccei Quinn transition

- II) GWs from *heavy axions*

# plan

- Intro

- I) GWs from *Peccei Quinn bubbles*

- II) GWs from *high quality defects*

# plan

- Intro

I) GWs from *Peccei Quinn bubbles*

II) GWs from *high quality defects*

→ post-inflationary PQ – most constrained, but *still* viable  
 $10^9 \text{ GeV} < f_a < 10^{11} \text{ GeV}$

## I) GWs from *Peccei Quinn bubbles*

# Can the PQ transition be 1st order?

## Peccei-Quinn Phase Transition at LIGO

Benedict von Harling<sup>a</sup>, Alex Pomarol<sup>a,b</sup>, Oriol Pujolàs<sup>a</sup>  
and Fabrizio Rompineve<sup>c</sup>

also JHEP 04 (2020) 025  
Delle Rose, Panico, Redi, Tesi

**arXiv:1912.07587**  
**JHEP 04 (2020) 195**

# Can the PQ transition be 1st order?

Yes, depending on the type of PQ model

## Peccei-Quinn Phase Transition at LIGO

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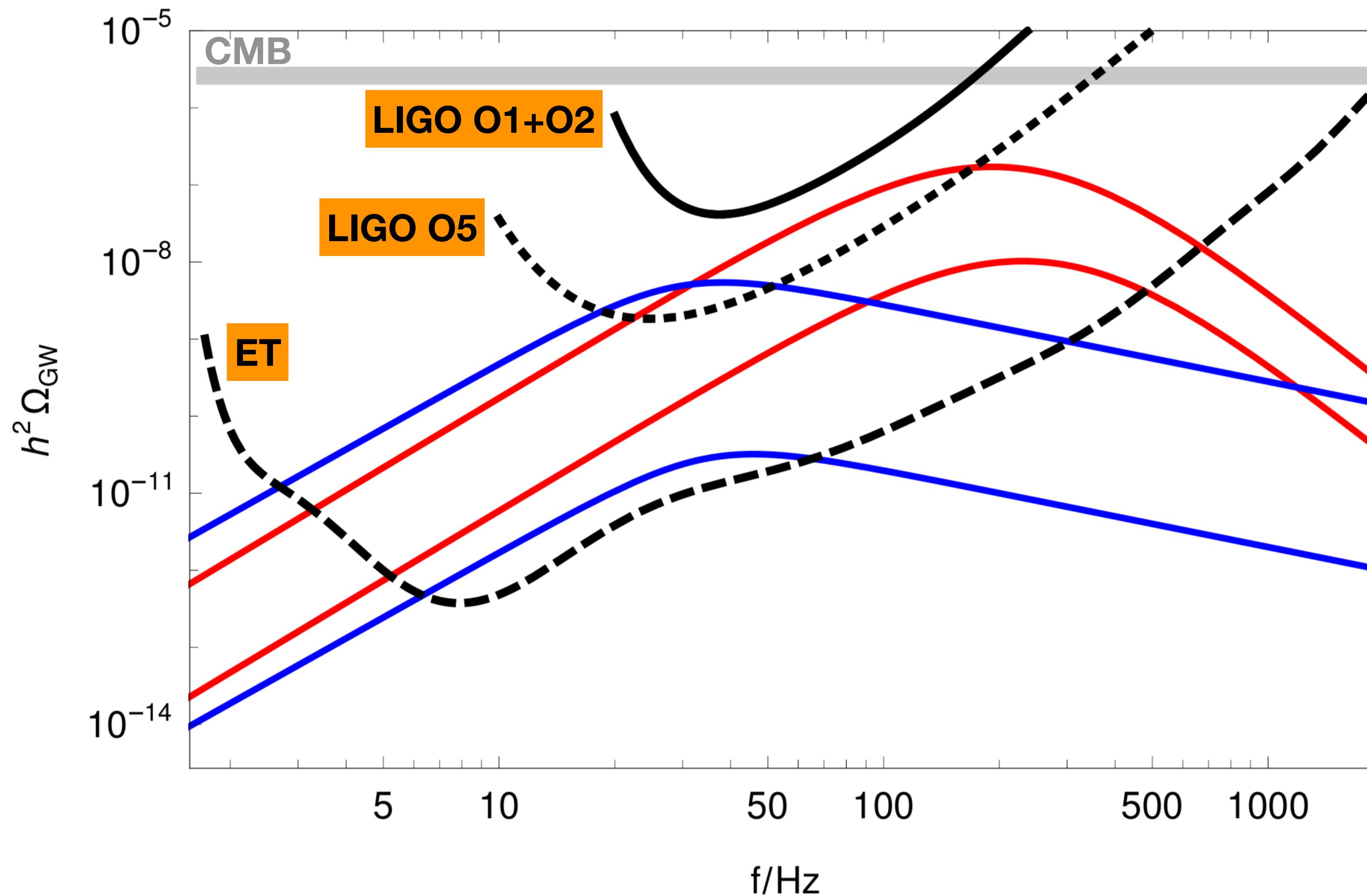
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## “ABC” of (1st order) cosmological phase transitions

$\alpha = \frac{\Delta V}{\rho_\gamma(T_n)}$	<b>latent heat</b>
$\beta = \frac{\dot{\Gamma}}{\Gamma} \Big _{T_n}$	<b>inverse duration</b>
$\Gamma = \mathcal{A} e^{-S_B}$	<b>nucleation rate</b>
$T_n$	<b>nucleation temperature</b>
$T_*$	<b>transition completed</b>

# 1) sensitivity reach to PTs

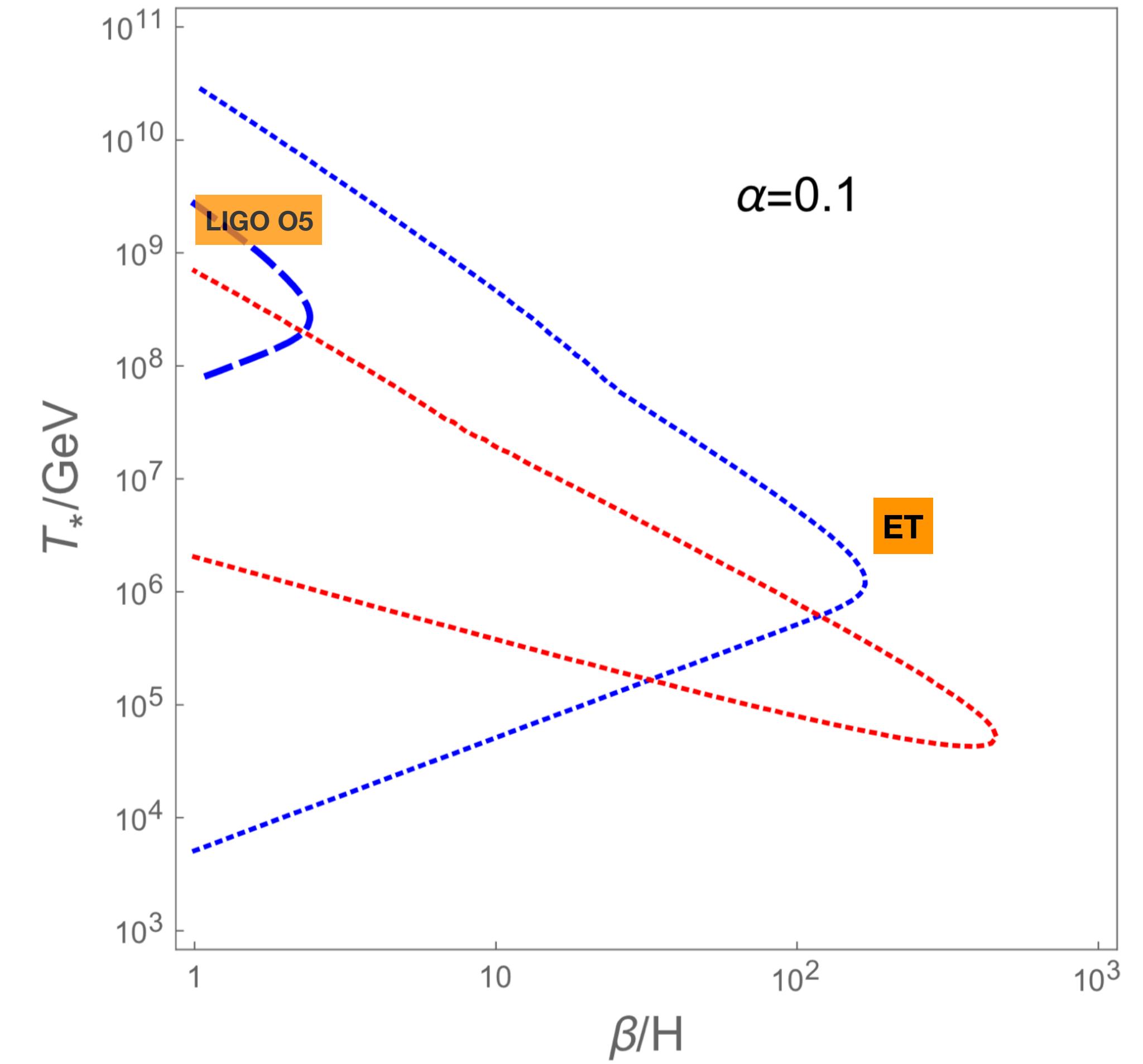
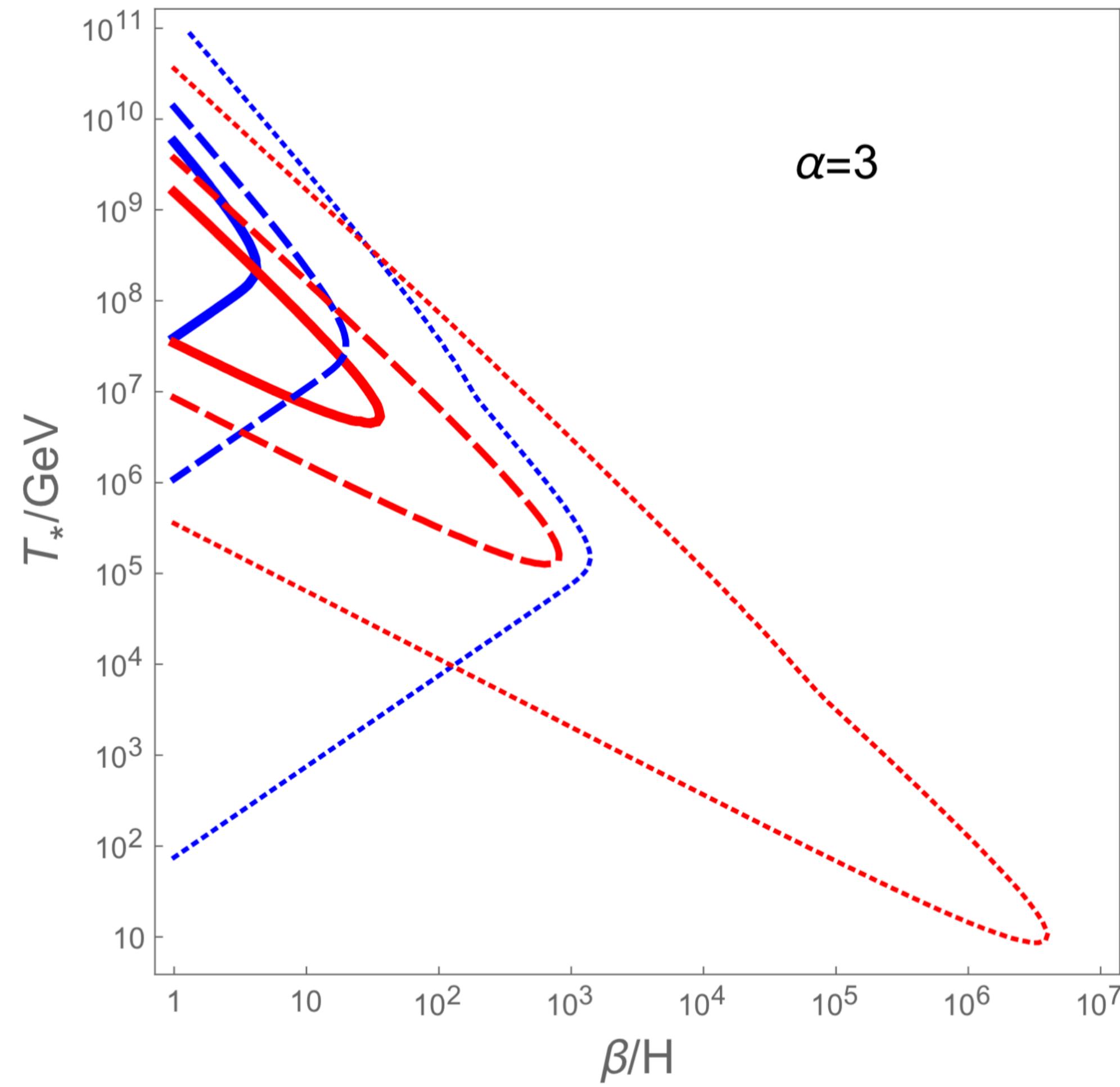


$F_a = 10^8 \text{ GeV}$  with  $\alpha \approx 3.5$

$F_a = 10^9 \text{ GeV}$  with  $\alpha \sim 10^6$

# 1) sensitivity reach to PTs

already  
ruling-out  
some  
parameter  
space!  
(solid lines)



## 2) Axion models with strong PT



- **KSVZ** (Kim, Shifman, Vainshtein, Zakharov)  
=> *extra heavy quarks with PQ-charge*

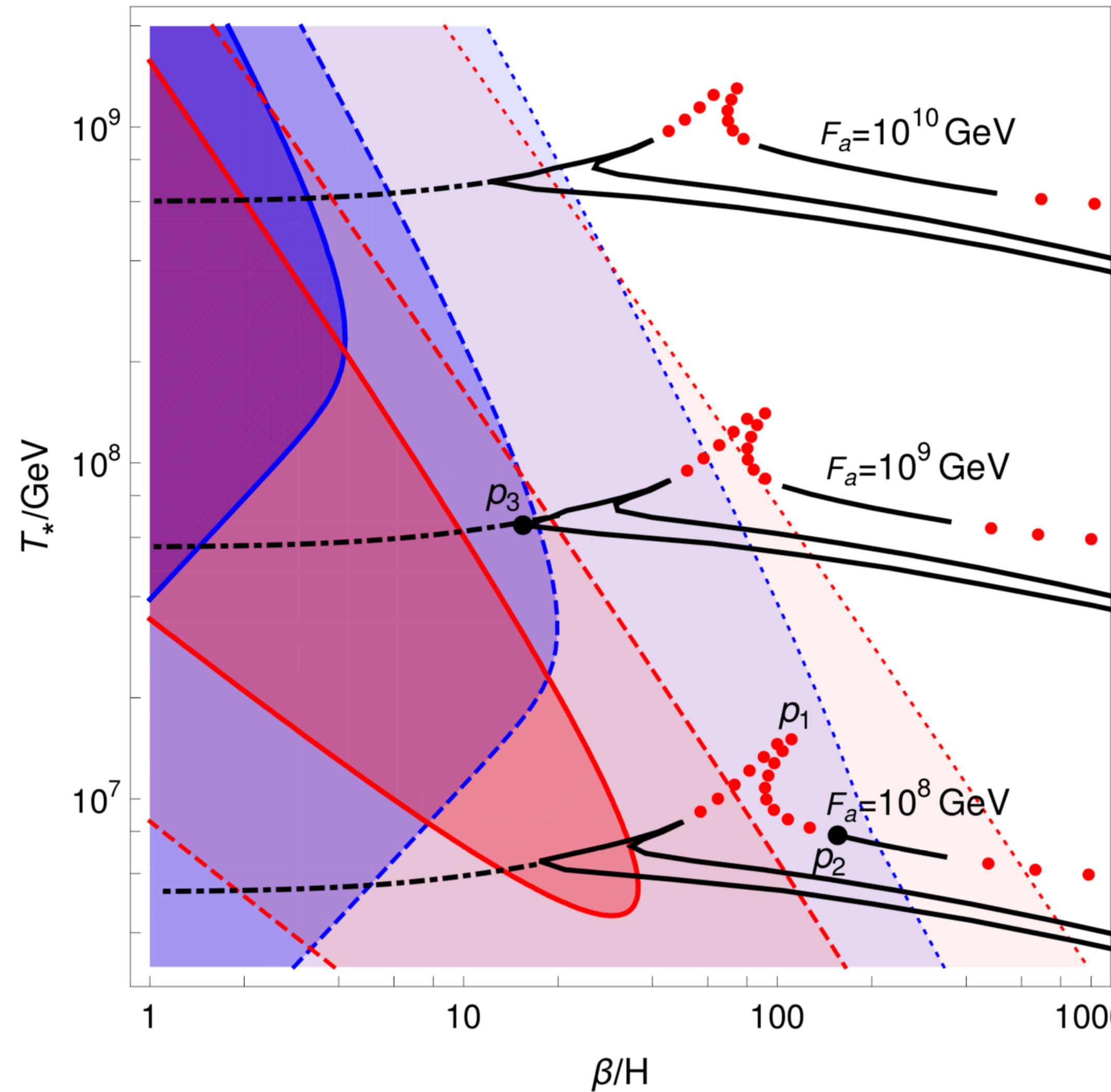


- **DFSZ** (Dine, Fischler, Srednicki, Zhitnitsky)  
=> *extra Higgs doublet with PQ-charge*

## 2) Axion models with strong PT

**DFSZ can do it !**

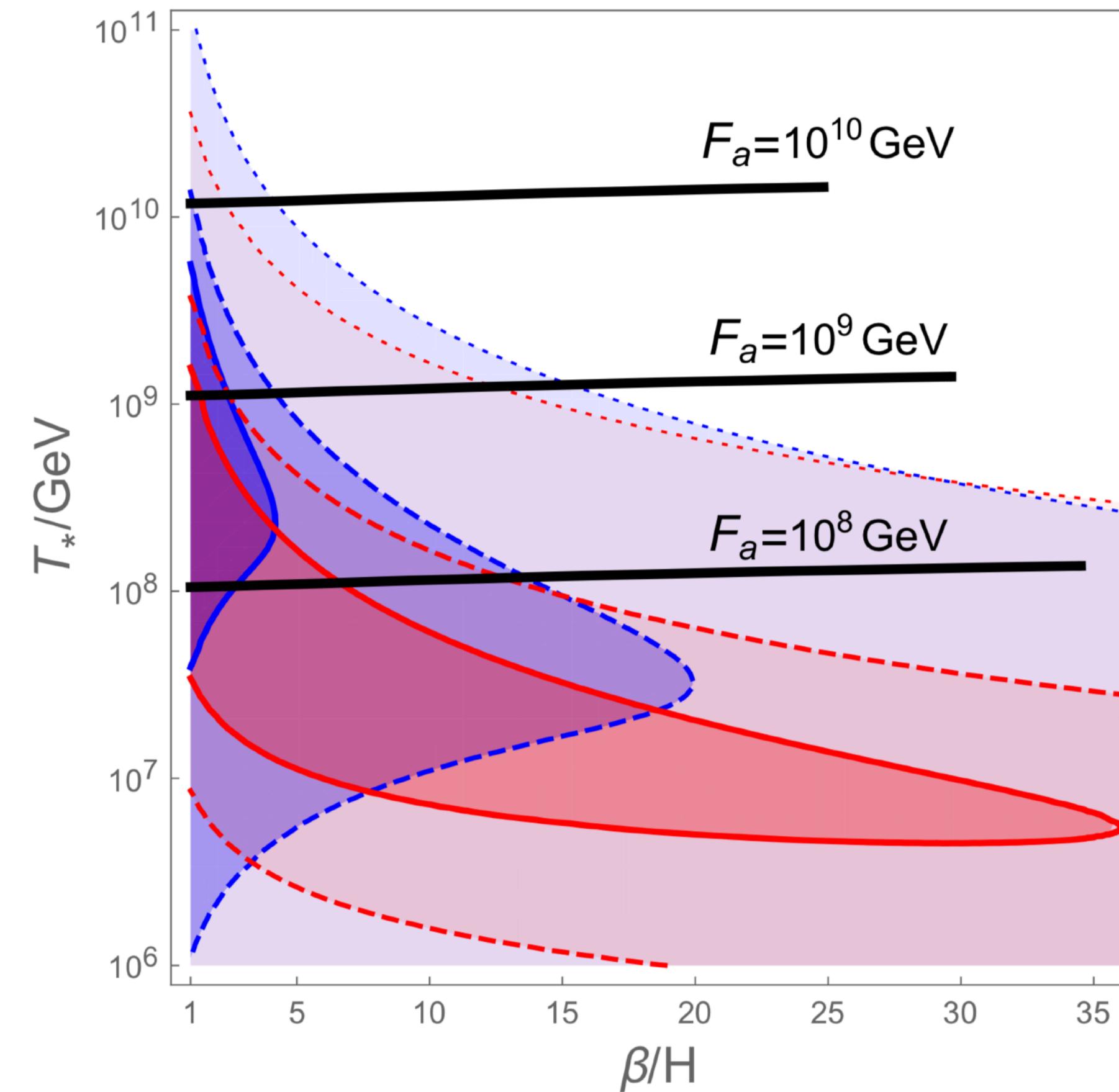
$T_*$  is  $\sim 1$  order of magnitude below  $F_a$



## 2) Axion models with strong PT

**KSVZ/DFSZ can be embedded in larger new physics sectors**

- SUSY version of KSVZ also works
- Strong dynamics as well



<b>Phenomenological model (bubble collisions)</b>					
		$\Omega_{\text{coll}}^{95\%} \text{ (25 Hz)}$			
$\beta/H_{\text{pt}} \setminus T_{\text{pt}}$		$10^7 \text{ GeV}$	$10^8 \text{ GeV}$	$10^9 \text{ GeV}$	$10^{10} \text{ GeV}$
1		$1.0 \times 10^{-8}$	$8.4 \times 10^{-9}$	$5.0 \times 10^{-9}$	—
10		$4.0 \times 10^{-9}$	$6.3 \times 10^{-9}$	—	—

TABLE III: The 95% CL upper limits on  $\Omega_{\text{coll}}^{95\%}(25 \text{ Hz})$  for fixed values of  $\beta/H_{\text{pt}}$  and  $T_{\text{pt}}$ , and  $v_w = \kappa_\phi = 1$ . The dashed lines denote no sensitivity for exclusion.

=> some models are already ruled out

# Conclusions

- LIGO/Virgo are sensitive to cosmological phase transitions in the early universe at around  $10^8 \text{ GeV}$  – HEP!
- Detectable signals arise if sufficient *supercooling*,  $\alpha \gtrsim 1$
- Prospects for the ET are very good
- Very relevant for axion physics: PQ phase transition  $T \sim 10^8 - 10^{12} \text{ GeV}$
- Detectable PQ models: DFSZ, strong dynamics, SUSY-KSVZ...
- Complementary to other gravitational axion tests from superradiance

## **II) GWs from *high quality defects***

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High Quality QCD Axion at Gravitational Wave Observatories

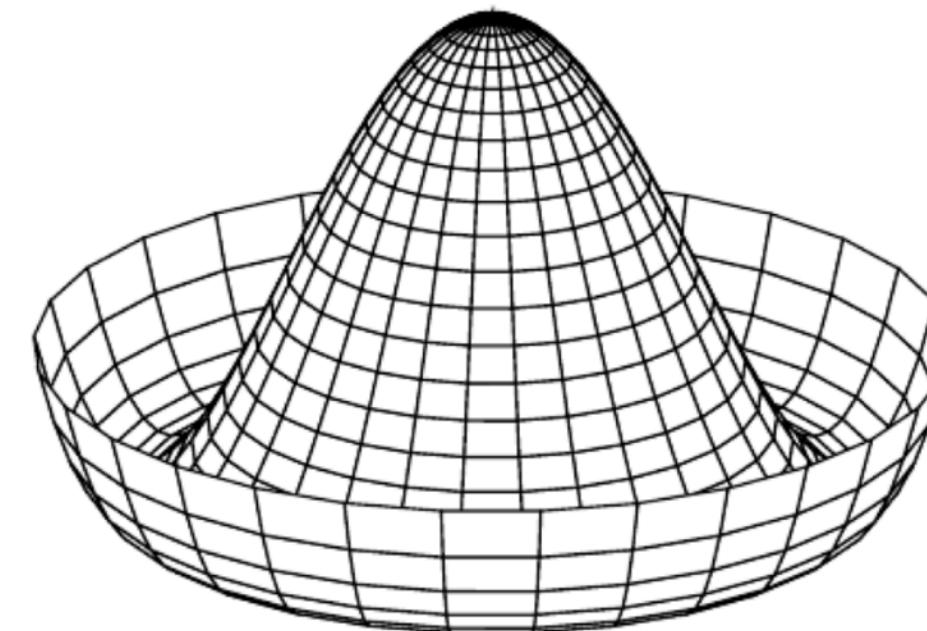
Ricardo Z. Ferreira,<sup>1,\*</sup> Alessio Notari,<sup>2,†</sup> Oriol Pujolàs,<sup>1,‡</sup> and Fabrizio Rompineve<sup>3,§</sup>

arXiv:2107.07542

# Axionic defects

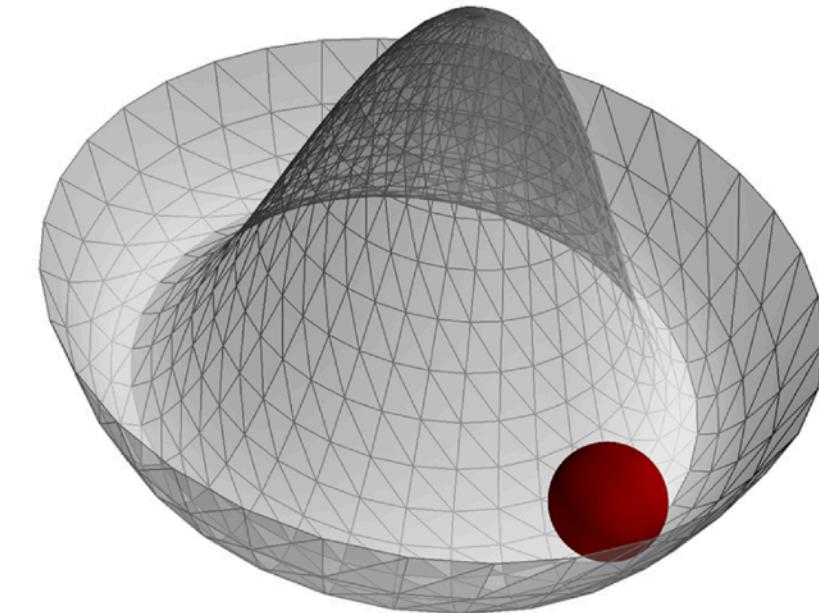
Peccei-Quinn mech: introduce a  $U(1)_{PQ}$  symmetry  
with:

i) SB and



spont.  $\cancel{U(1)_{PQ}}$  @  $f_a$

ii)  $U(1)_{PQ} \text{SU}(3)_c \text{SU}(3)_c$  anomaly



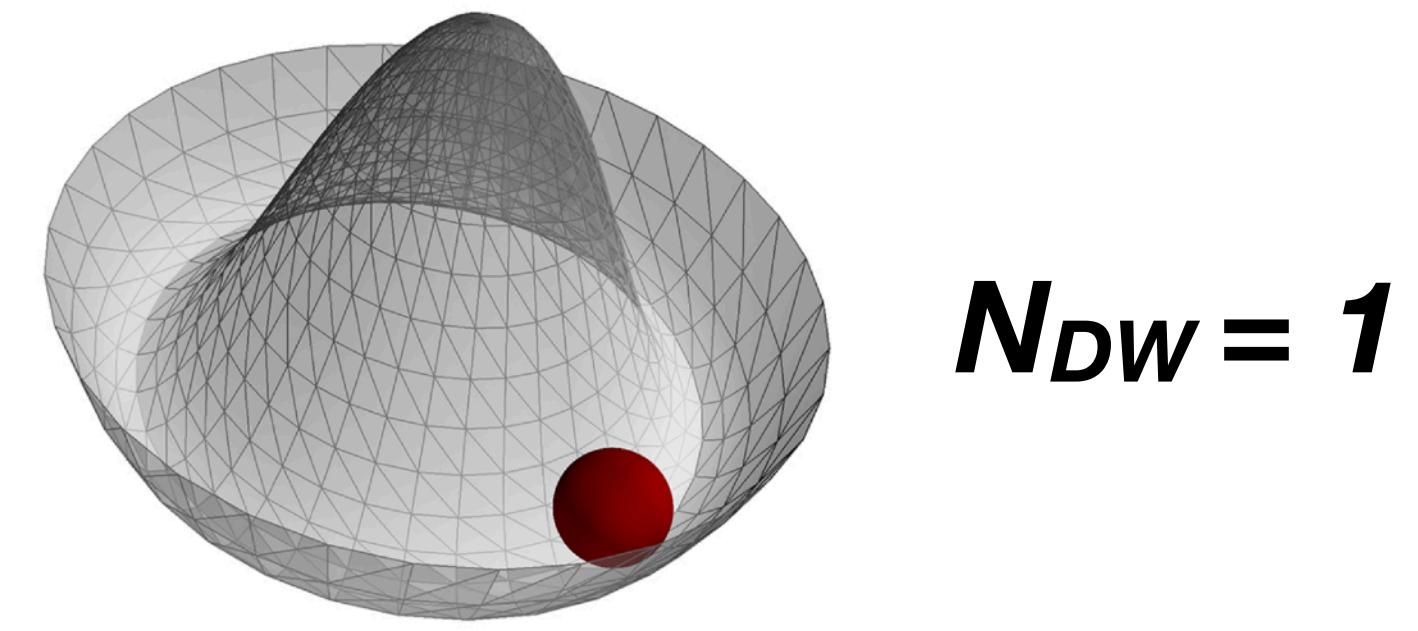
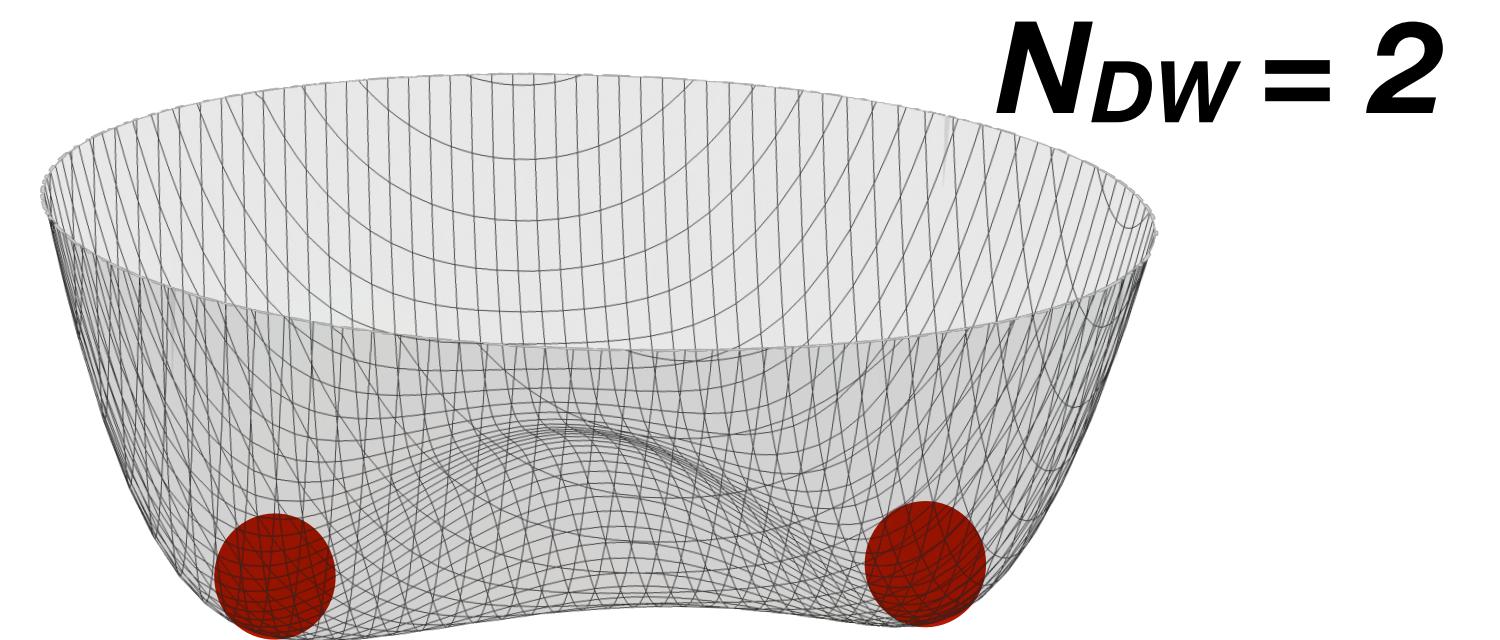
expl.  $\cancel{U(1)_{PQ}}$  @  $\Lambda_{QCD}$

# Axionic defects

→  $N_{DW}$  "Domain wall number"

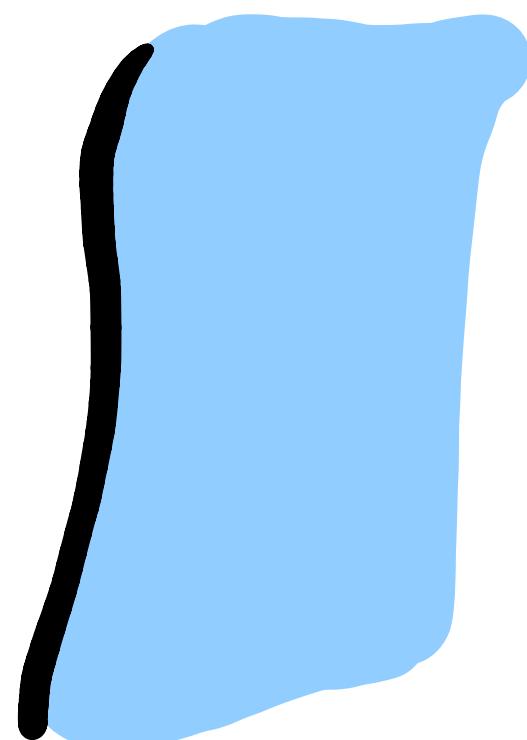
$$U(1)_{PQ} \rightarrow Z_{N_{DW}}$$

$$V_{QCD}(a) = \Lambda_{QCD}^4 \left( 1 - \cos \left( N_{DW} \frac{a}{f} \right) \right)$$

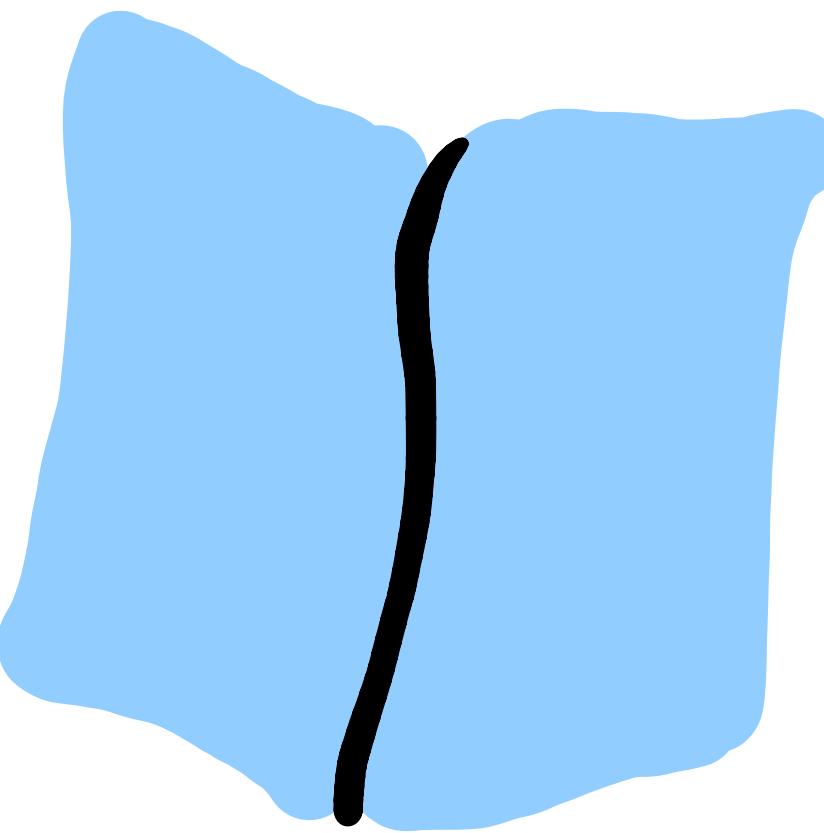


# Axionic defects

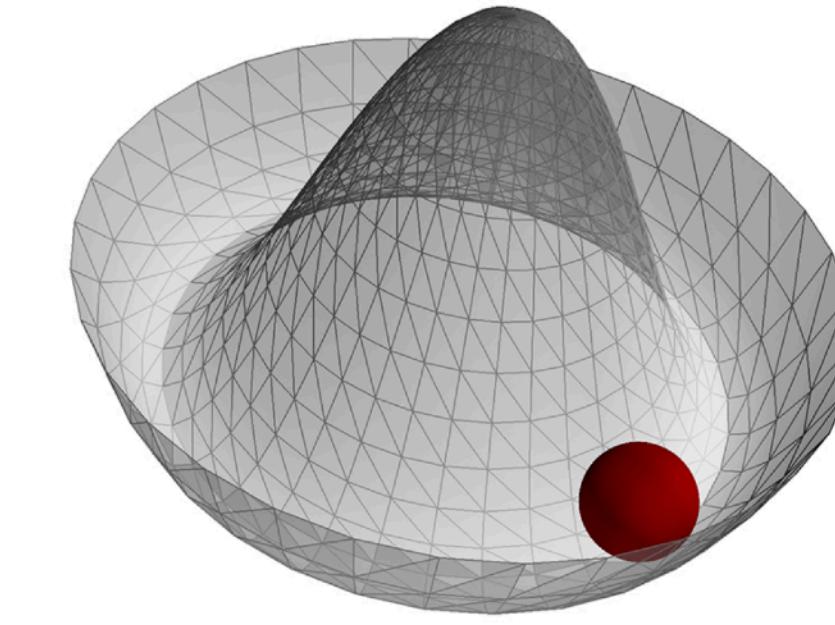
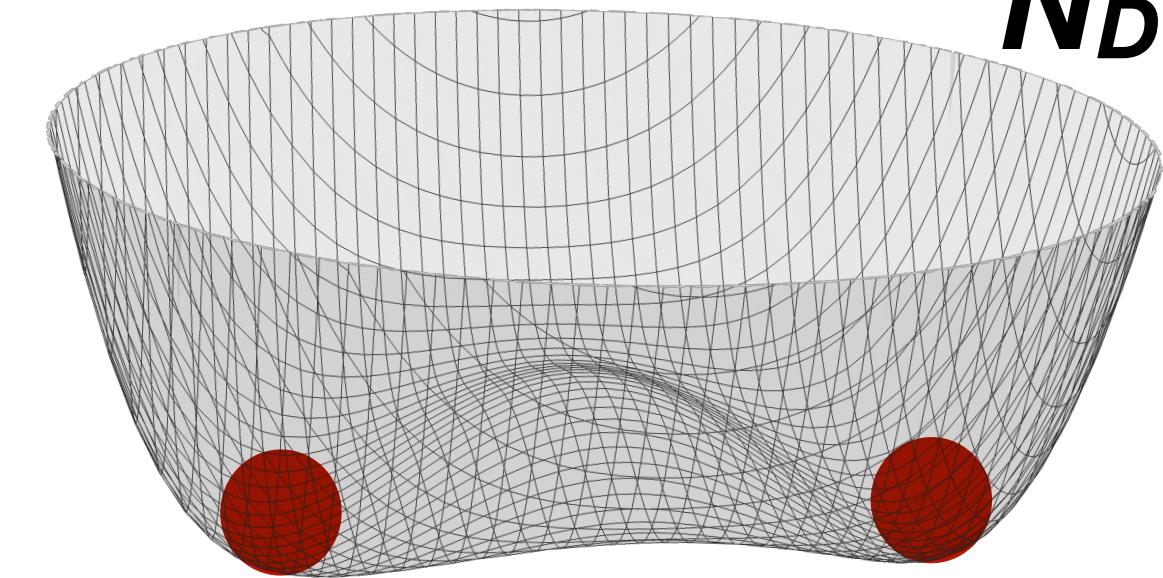
$N_{DW} = 1$



$N_{DW} = 2$



$N_{DW} = 2$



$N_{DW} = 1$

# Axion 'quality problem'

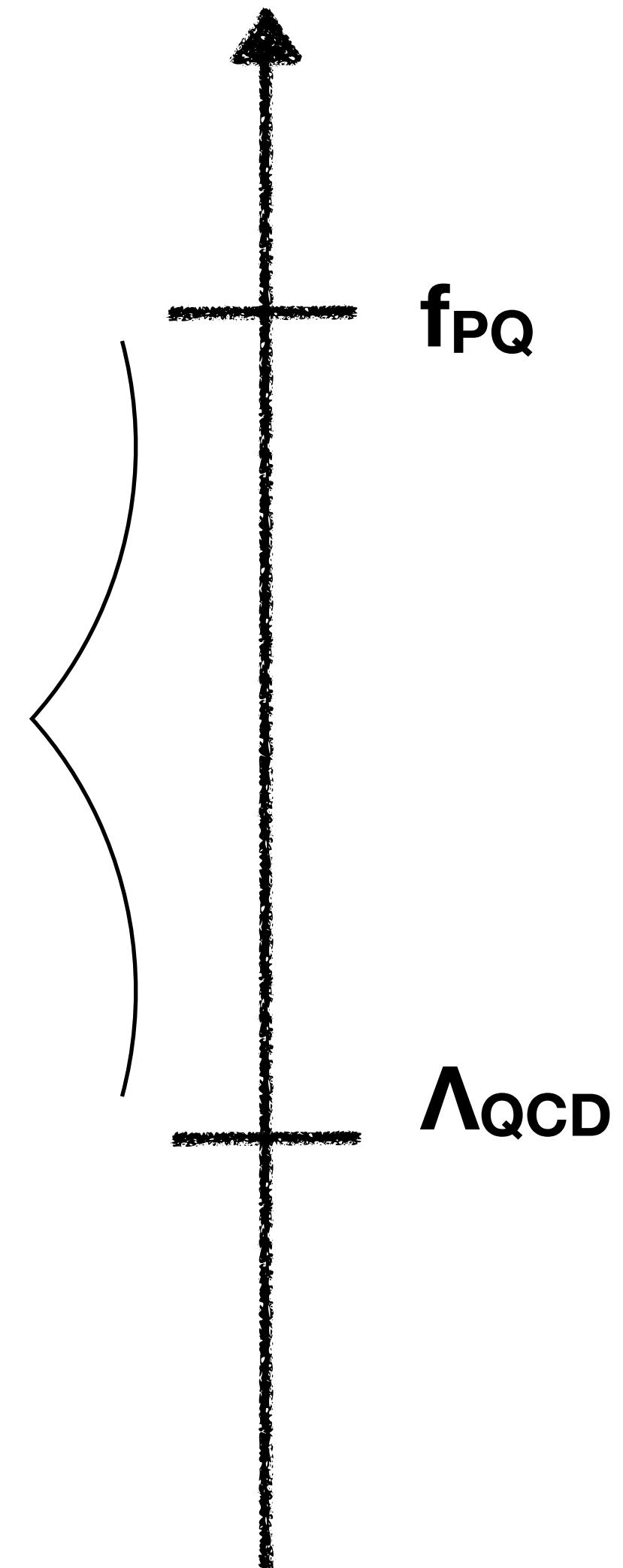
# Axion 'quality problem'

Implicitly, only QCD contributes to  $V(a)$

Solution can be **spoiled** by *any* other "misaligned" sector

$$V(a) = \Lambda_{QCD}^4 \left(1 - \cos(a)\right) + \Lambda_{mis}^4 \left(1 - \cos(a + \delta)\right)$$

strong CP solved =>  $\Delta\theta = \sin\delta \frac{\Lambda_{mis}^4}{\Lambda_{QCD}^4} \lesssim 10^{-10}$



# Axion 'quality problem'

**From the 'PQ-EFT' perspective, misaligned terms**

$$c_N \frac{\Phi^{4+N}}{M^N}$$

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$$c_N \frac{\Phi^{4+N}}{M^N} \rightarrow \Lambda_{mis} = c_N^{1/4} \left( \frac{f}{M} \right)^{N/4} \lesssim 10^{-2.5} \Lambda_{QCD}$$

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Either first nonzero  $c_N$  with large N

Or ~~PQ~~ is nonperturbative ( $c_N$ 's exponentially suppressed)

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Either first nonzero  $c_N$  with large N

Or ~~PQ~~ is nonperturbative (c<sub>N</sub>'s exponentially suppressed)

In any case, robustness against this ~~PQ~~ is desirable →

push  $\Lambda_{QCD}$  up  
Heavy axion

# Heavy axions

**SM extended with (or embedded into) new sector with  $\Lambda \gg \Lambda_{QCD}$**

$$V(a) = [\Lambda_{QCD}^4 + \Lambda^4] (1 - \cos(a))$$

- ***Mirror SM with softly broken  $Z_2$***
- ***Small instantons***
- ***Color - Dark Color unification***

# Heavy axions

**In practice:**

$$V(a) = \Lambda^4 \left( 1 - \cos(N_{DW} a) \right) + \Lambda_{mis}^4 \left( 1 - \cos(N'_{DW} a + \delta) \right)$$

# Heavy axions

**In practice:**

$$V(a) = \Lambda^4 \left( 1 - \cos(N_{DW} a) \right) + \Lambda_{mis}^4 \left( 1 - \cos(N'_{DW} a + \delta) \right)$$

$$f_a m_a = \Lambda^2 \gg \Lambda_{QCD}^2$$

$$m_a \simeq 10^8 \text{ GeV} \left( \frac{10^{12} \text{ GeV}}{f} \right) \left( \frac{\Lambda_H}{10^{10} \text{ GeV}} \right)^2$$

# Heavy axions

In practice:

$$V(a) = \Lambda^4 \left( 1 - \cos(N_{DW} a) \right) + \Lambda_{mis}^4 \left( 1 - \cos(N'_{DW} a + \delta) \right)$$

$$\Lambda_{mis} \lesssim 10^{-2.5} \Lambda$$

$$N_{DW}, \quad N'_{DW} \quad \text{co-prime}$$

# Heavy axions

**Heavy => decays quickly**

**Axions  $\neq$  DM**

**Heavy and weakly coupled => difficult to probe directly!**

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**presence of the  
misaligned  
sector =>  $\theta \neq 0$**

**improving precision in neutron EDM  $d_n$**

**nEDM@SNS (1908.09937)**

**->**

**$\Delta\theta \gtrsim 10^{-12}$**

**pEDM (2007.10332)**

**->**

**$\Delta\theta \gtrsim 10^{-13}$**

# Heavy axions

**Heavy => decays quickly**

**Axions  $\neq$  DM**

**Heavy and weakly coupled => difficult to probe directly!**

—> cosmology!

**Axionic defects (and  $N_{DW}$ ) play a crucial role**

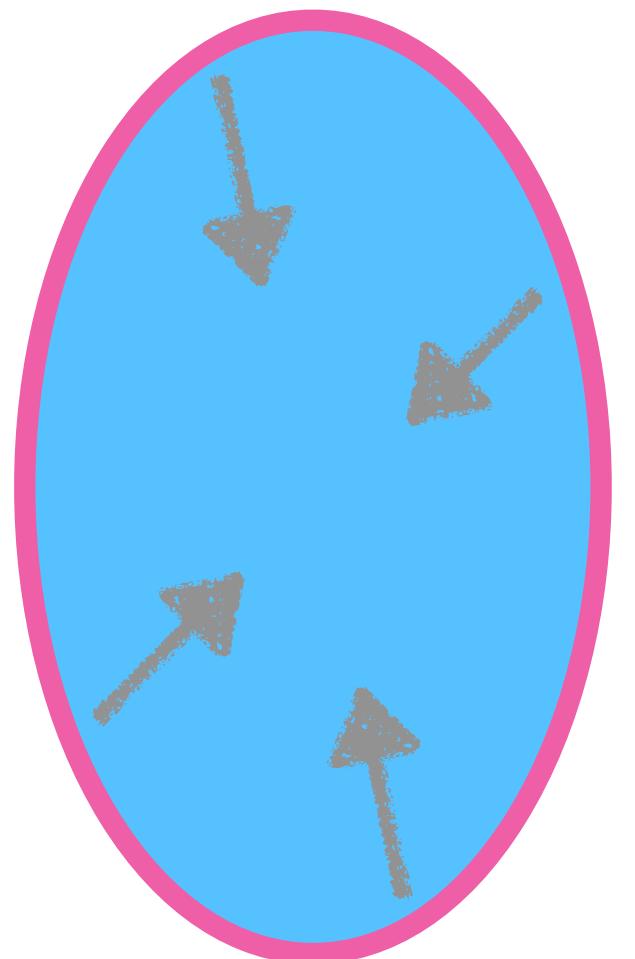
# High quality axionic defects

$$\sigma_{DW} \sim m_a f_a^2 \quad \Rightarrow \text{ DWs are also } \textit{heavy}$$

DW motions induce ***strong*** GWs

$$\rho_{GW} \sim \frac{\sigma_{DW}^2}{M_P^2}$$

$$N_{DW} = 1$$

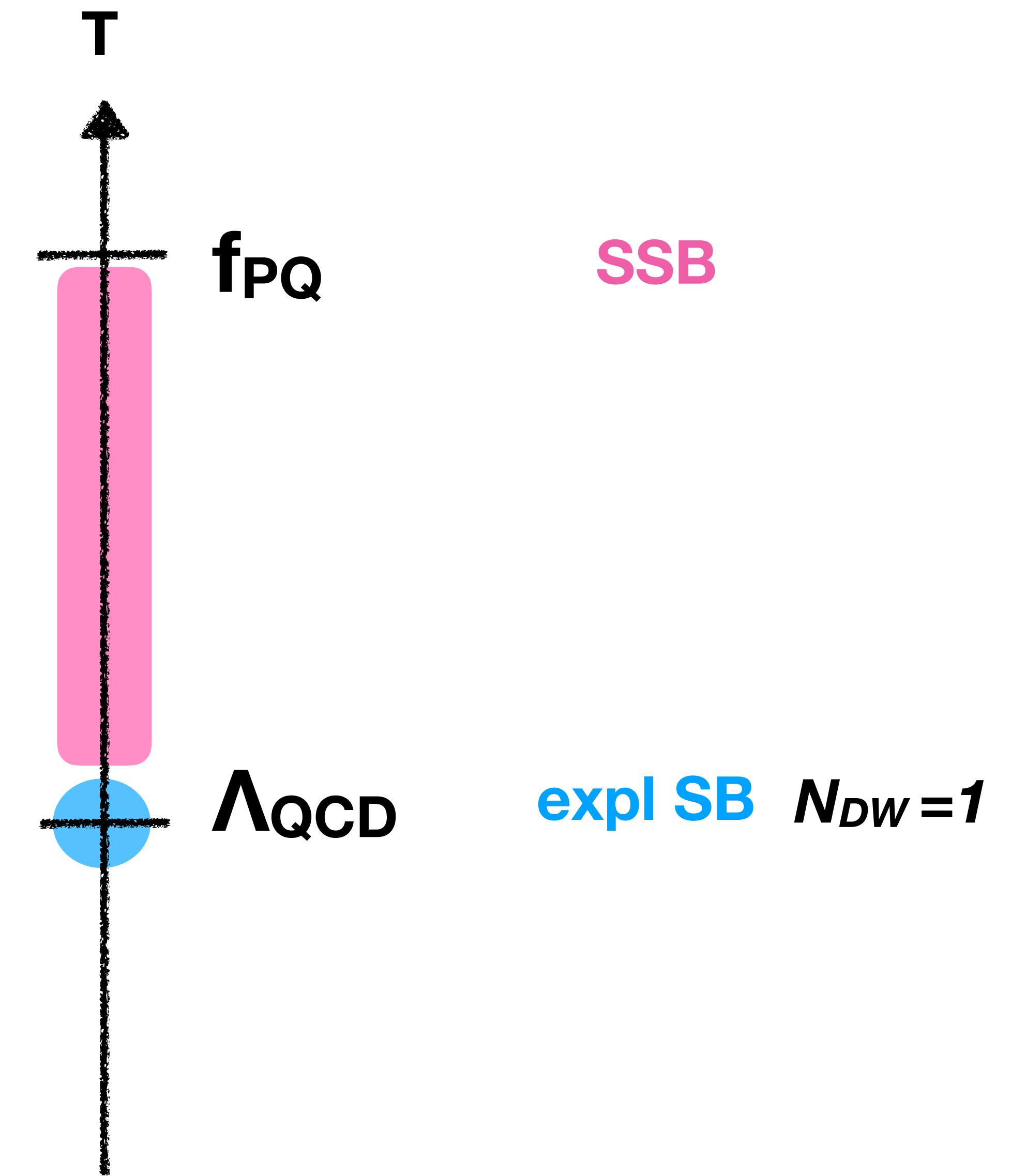


1 DW per string

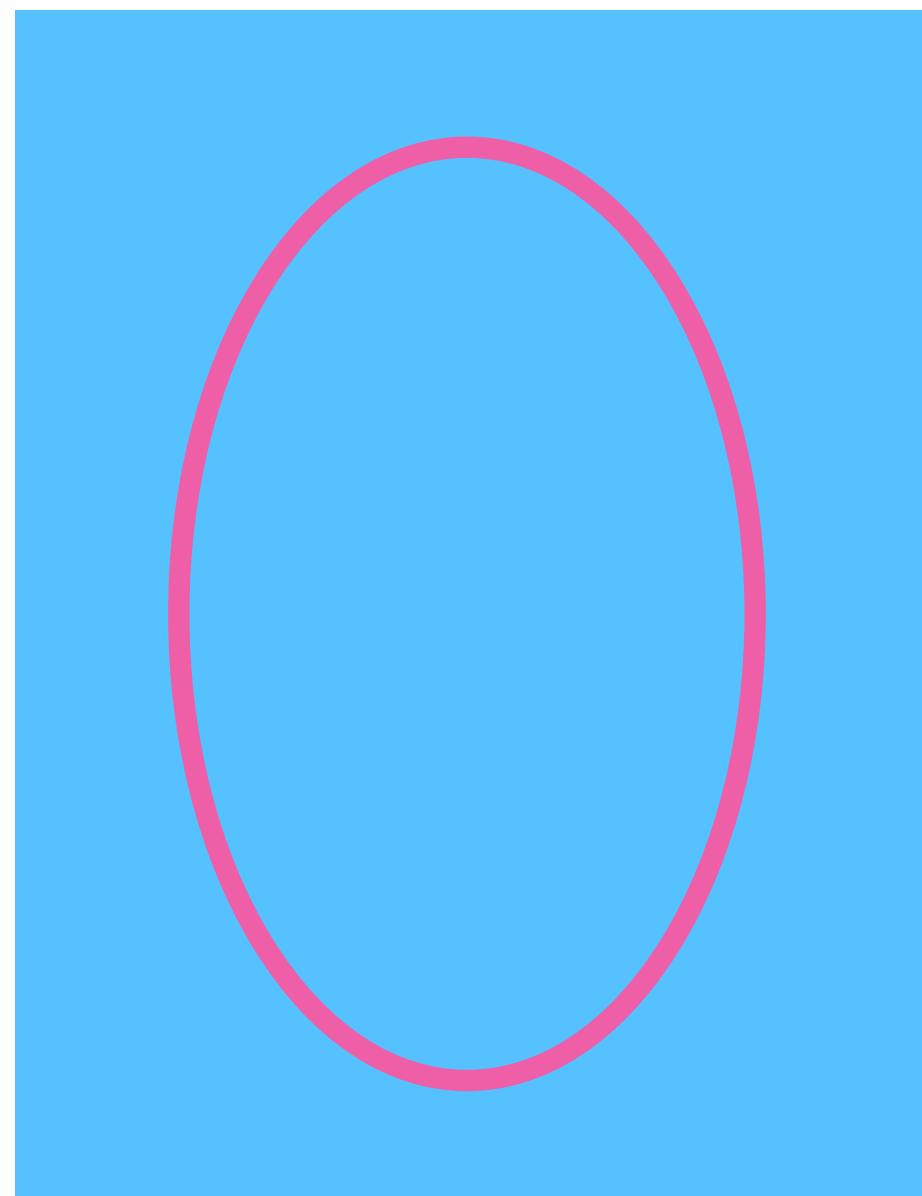
DWs make the whole network disappear quickly

Strings

Strings & DWs



$$N_{DW} > 1$$

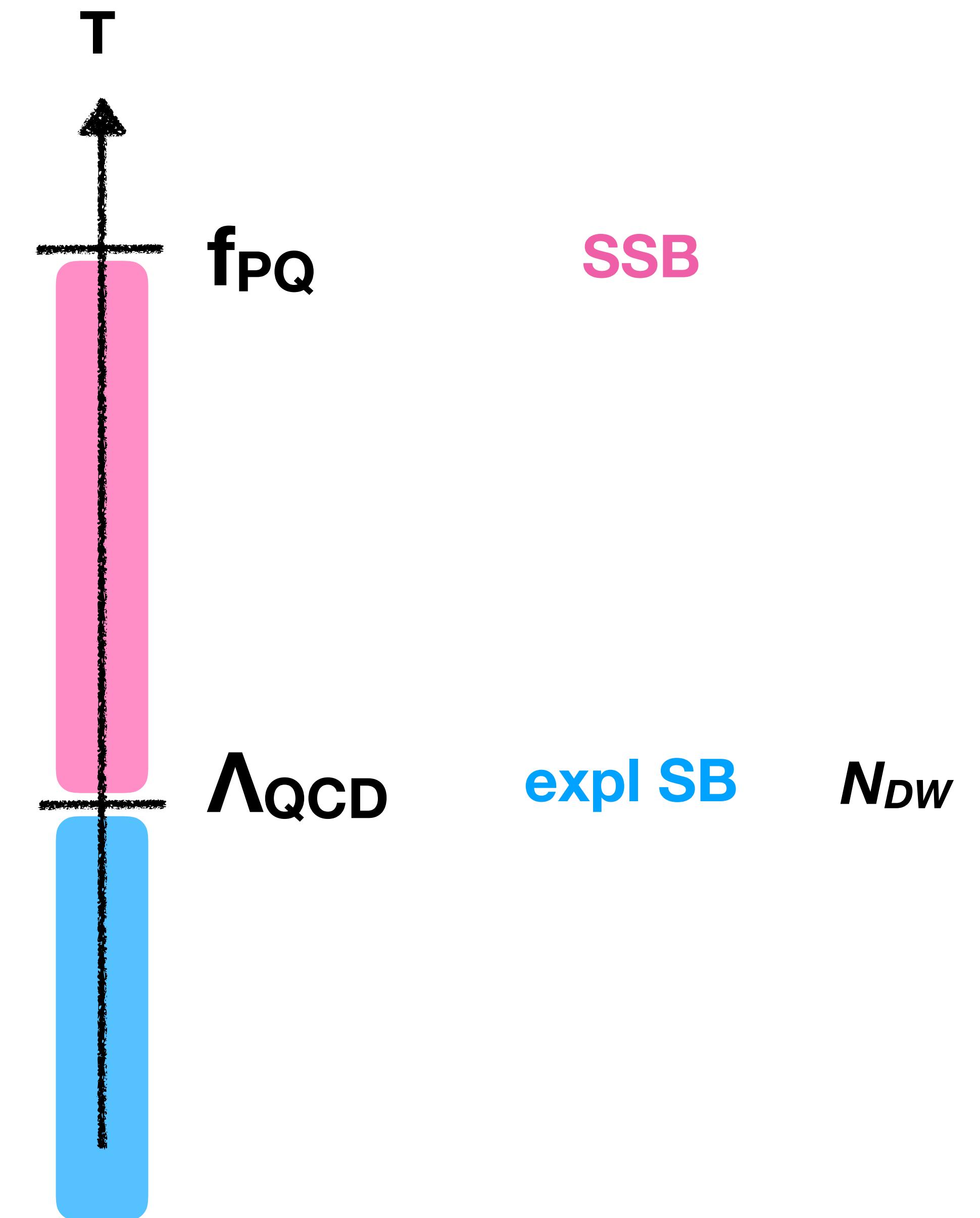


**Network stable**

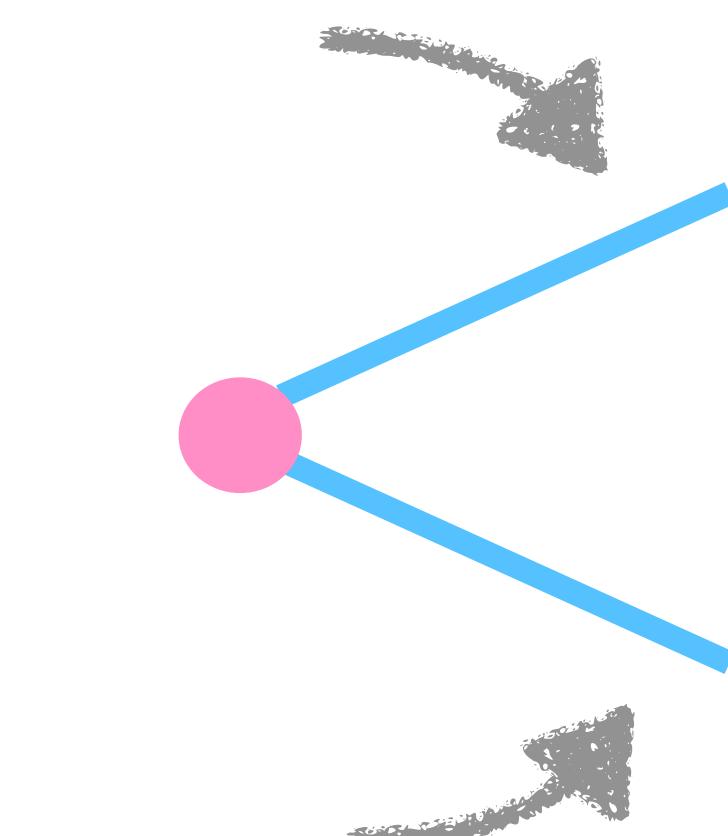
**DW problem!**

**Strings**

**Strings & DWs**



$N_{DW} > 1$



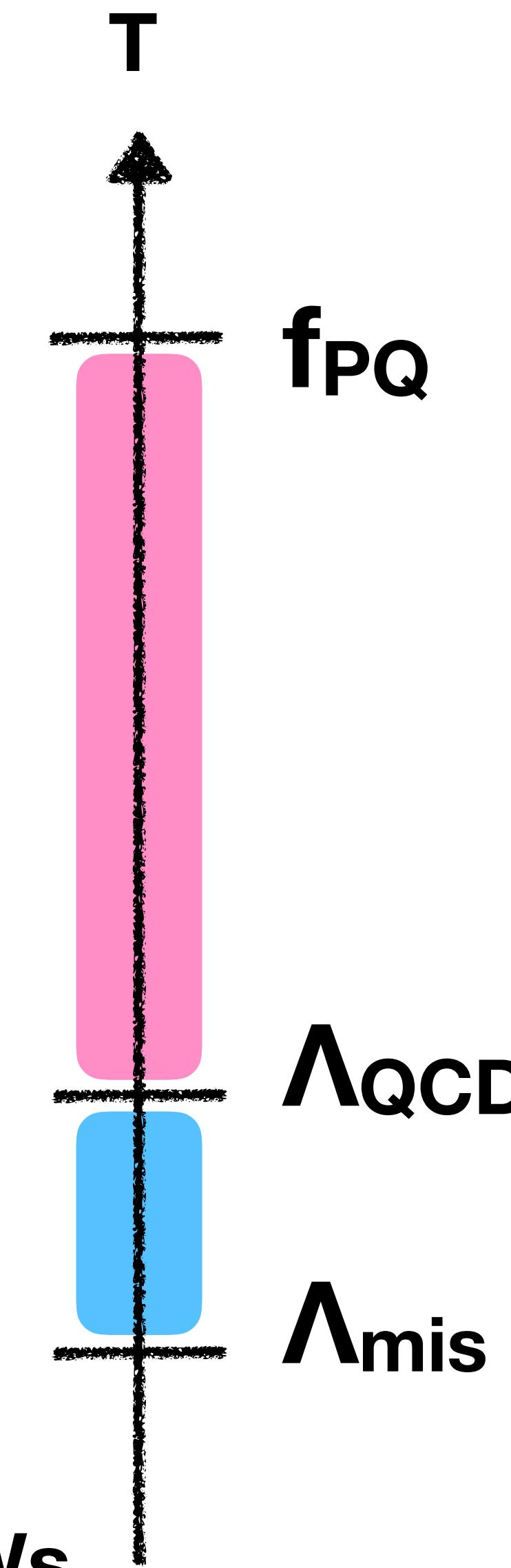
Delayed  
network  
annihilation

long DW  
epoch

Strings & DWs

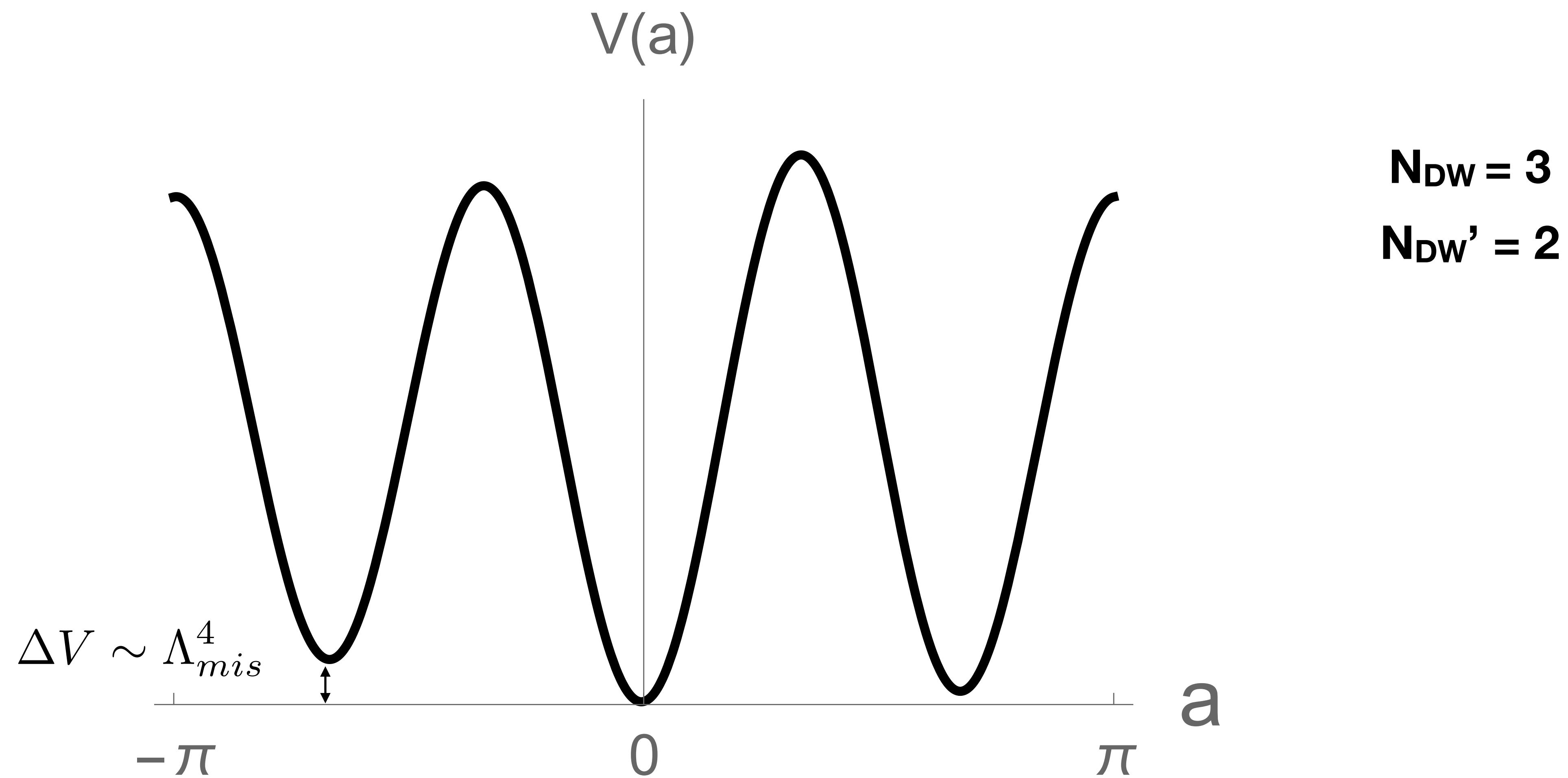
Strings & DWs &  
pressure difference on the DWs

Strings

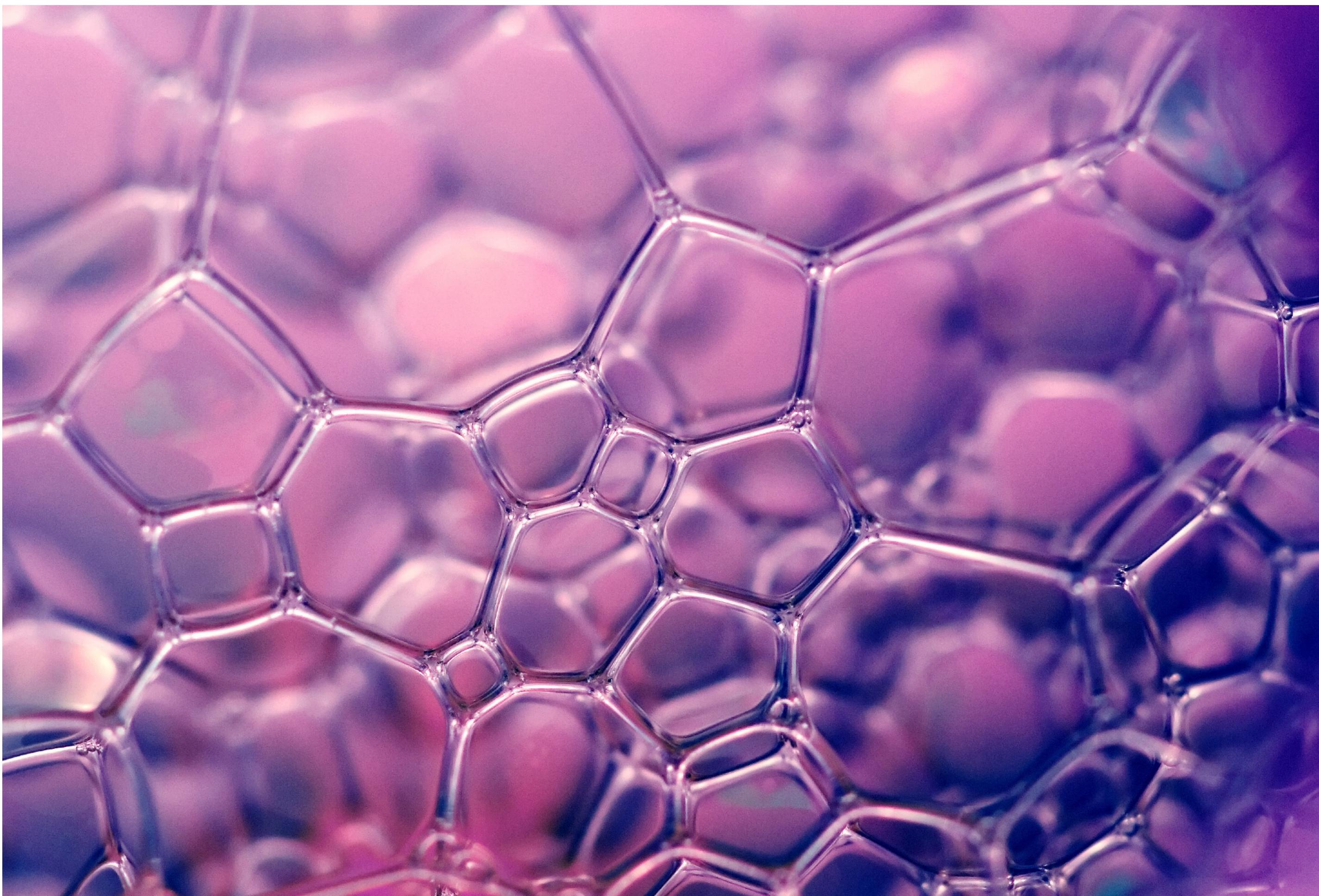


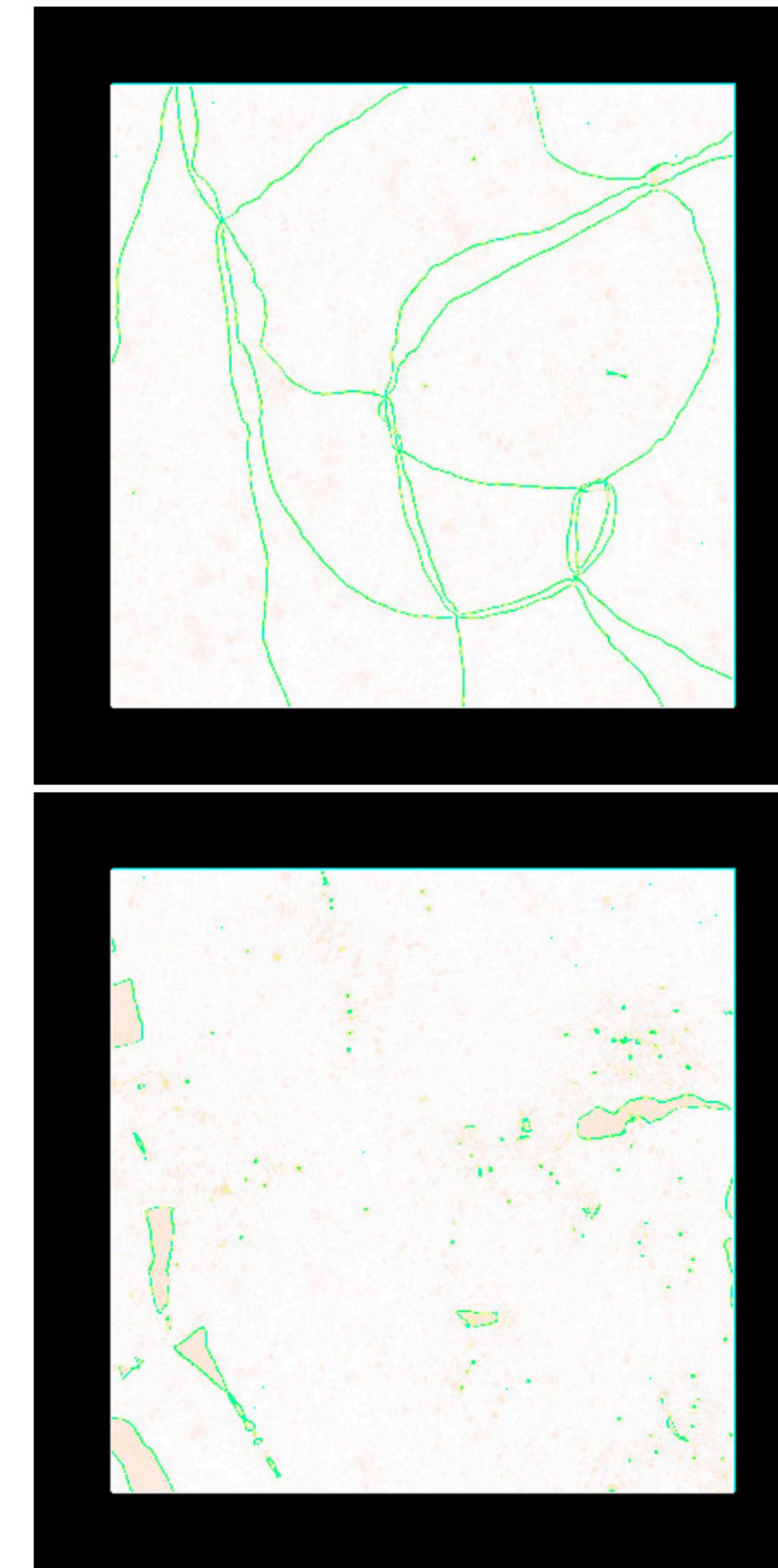
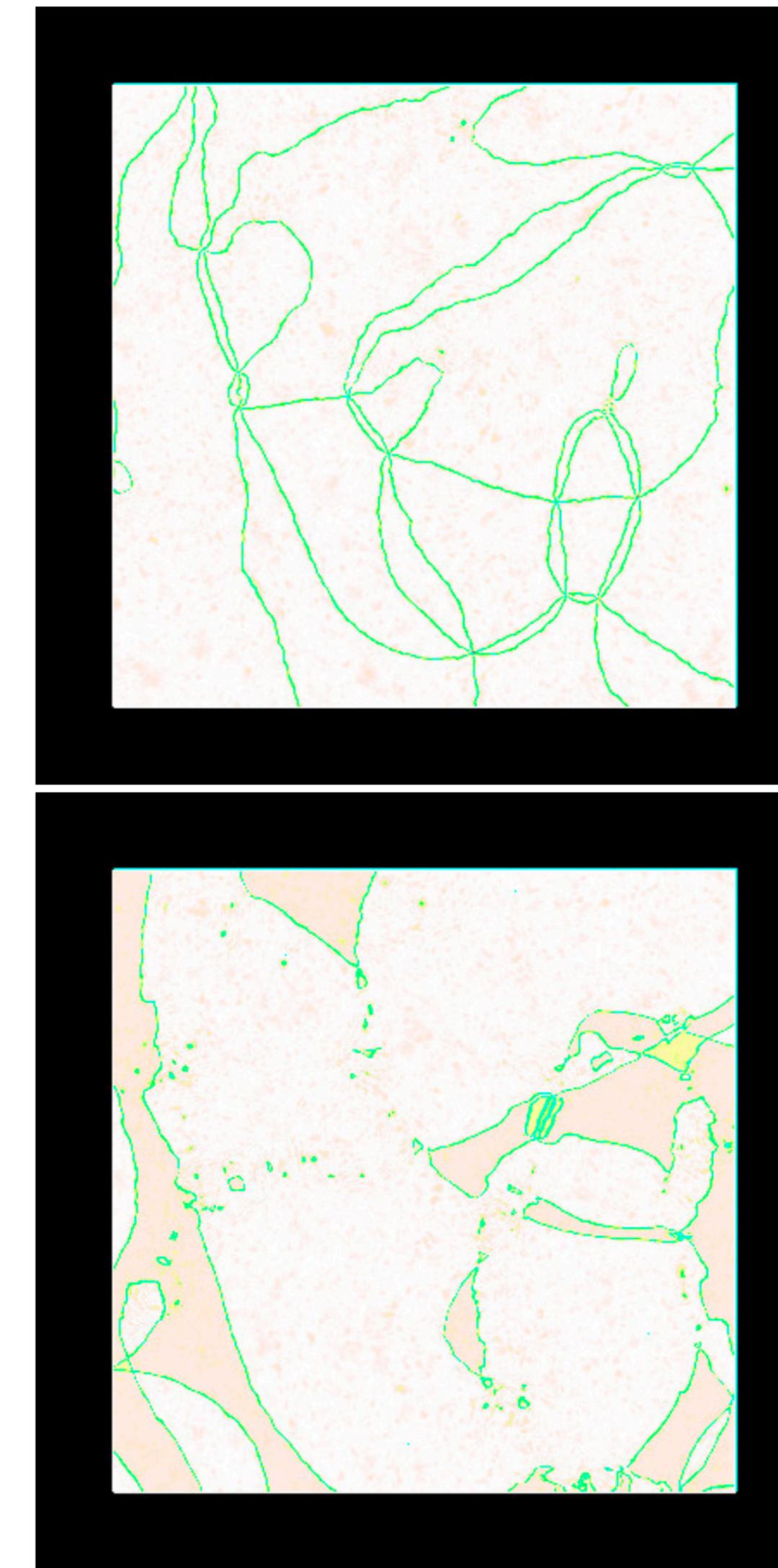
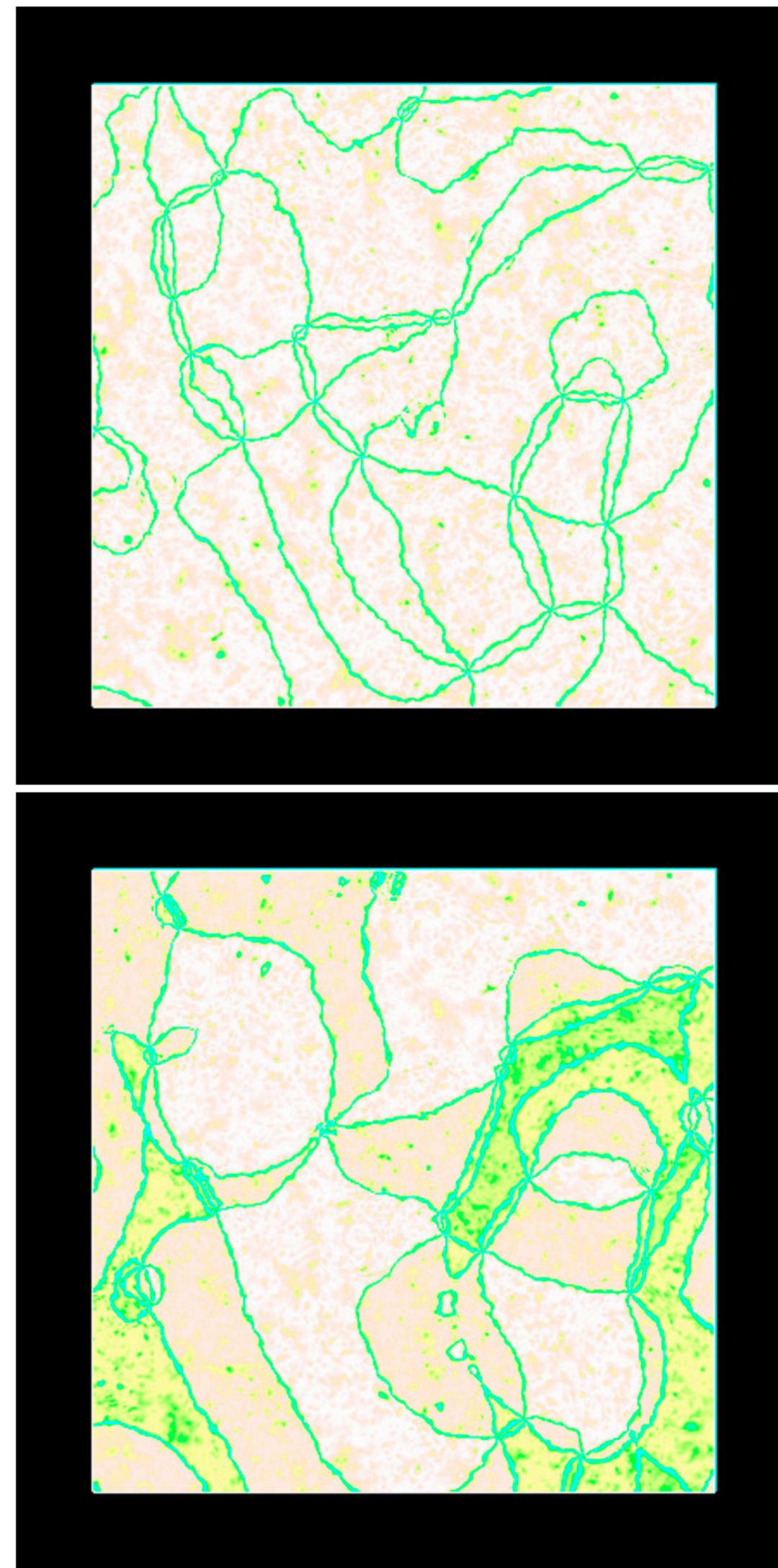
expl SB       $N_{DW}$   
expl SB       $N_{DW}'$

$$V(a) = \Lambda^4 \left(1 - \cos(N_{DW} a)\right) + \Lambda_{mis}^4 \left(1 - \cos(N'_{DW} a + \delta)\right)$$



# Wall-string network dynamics





$(N_{DW} > 1)$

**without  $\Lambda_{mis}$   
network persists**

**with  $\Lambda_{mis}$   
network  
annihilates**

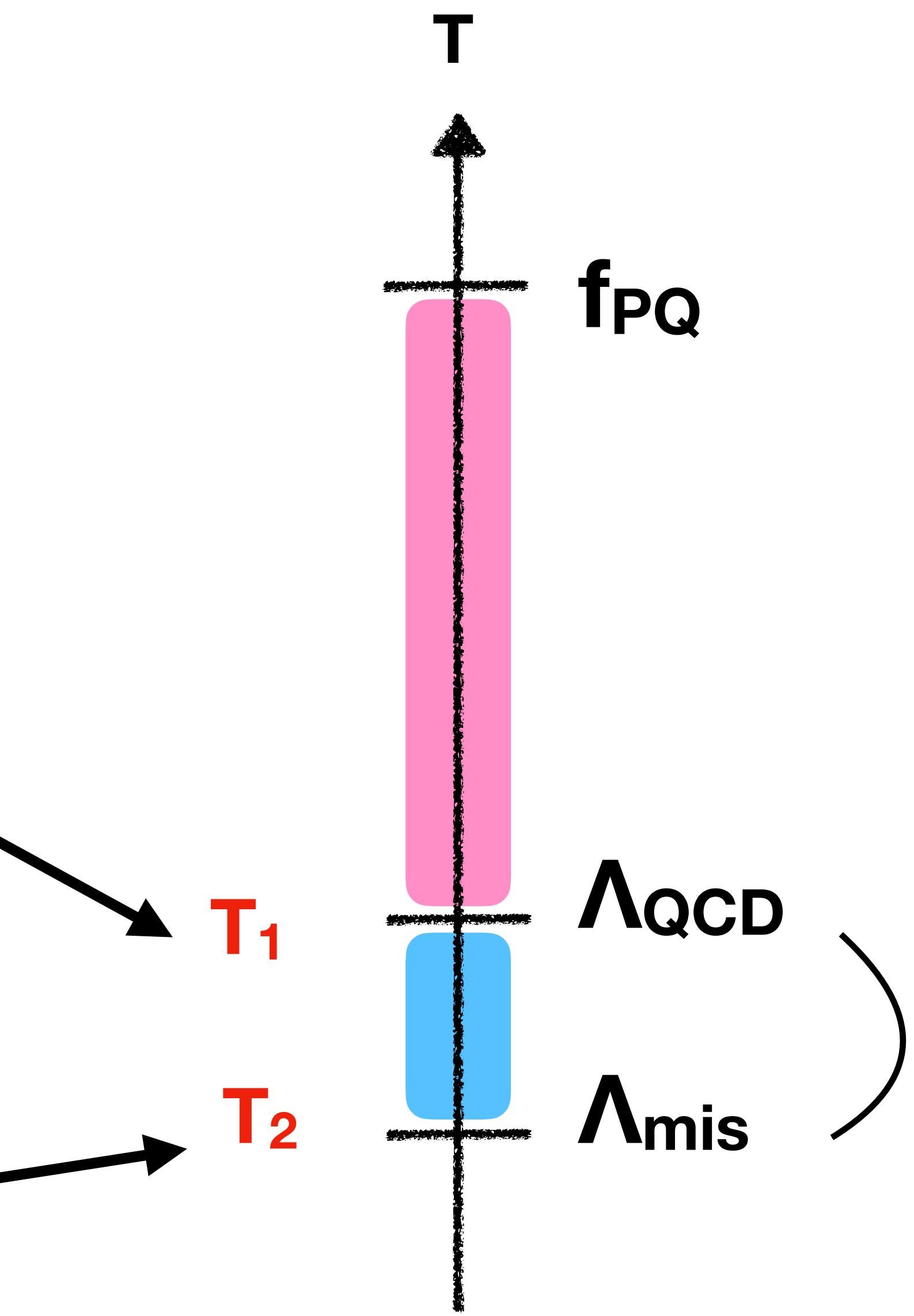
from [Kawasaki et al. '14]

$$\Delta V \approx 10^{-3} V_{QCD}$$

**DWs form when**

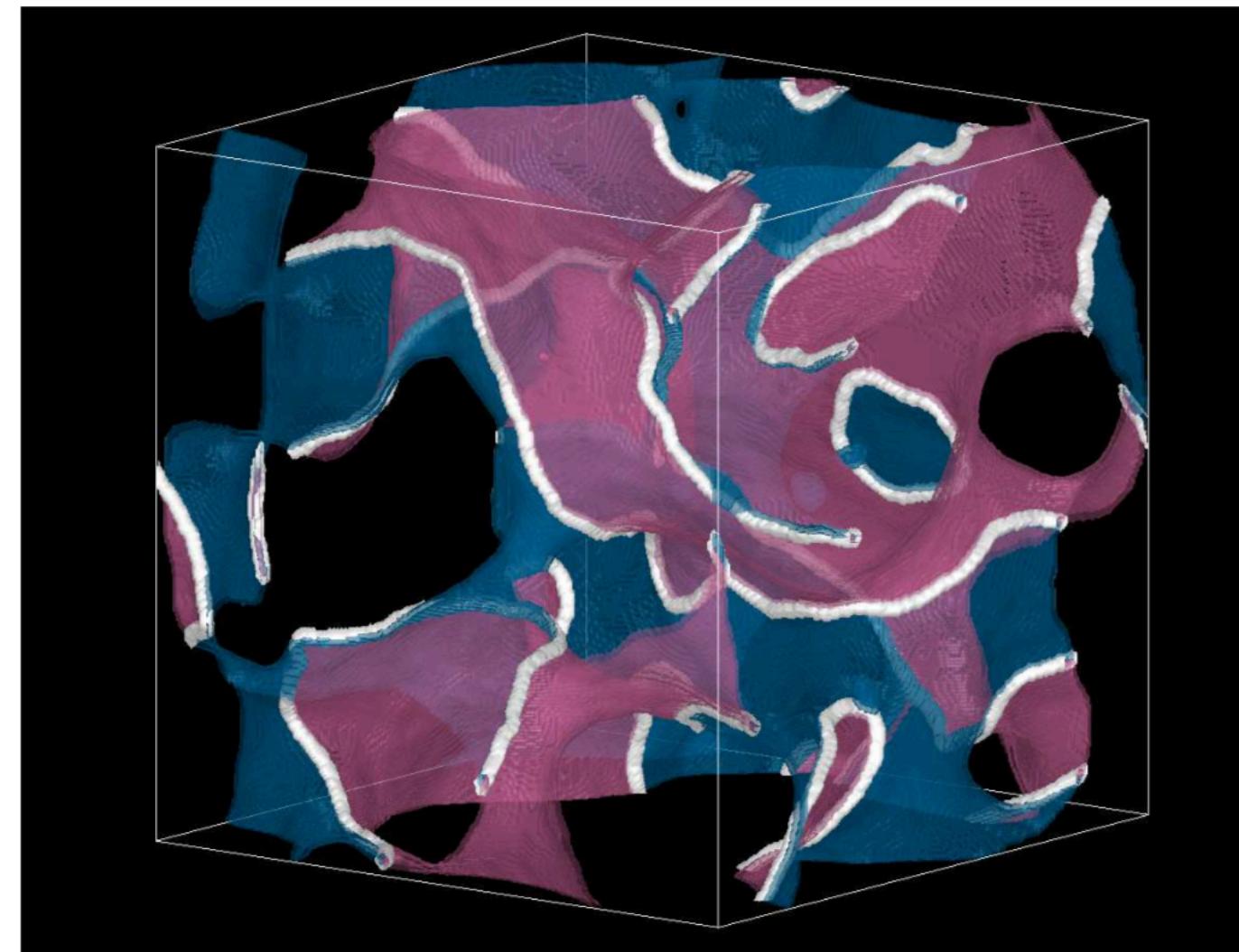
$$3H(T_1) = m_a(T_1)$$

**DW+string network ends when**

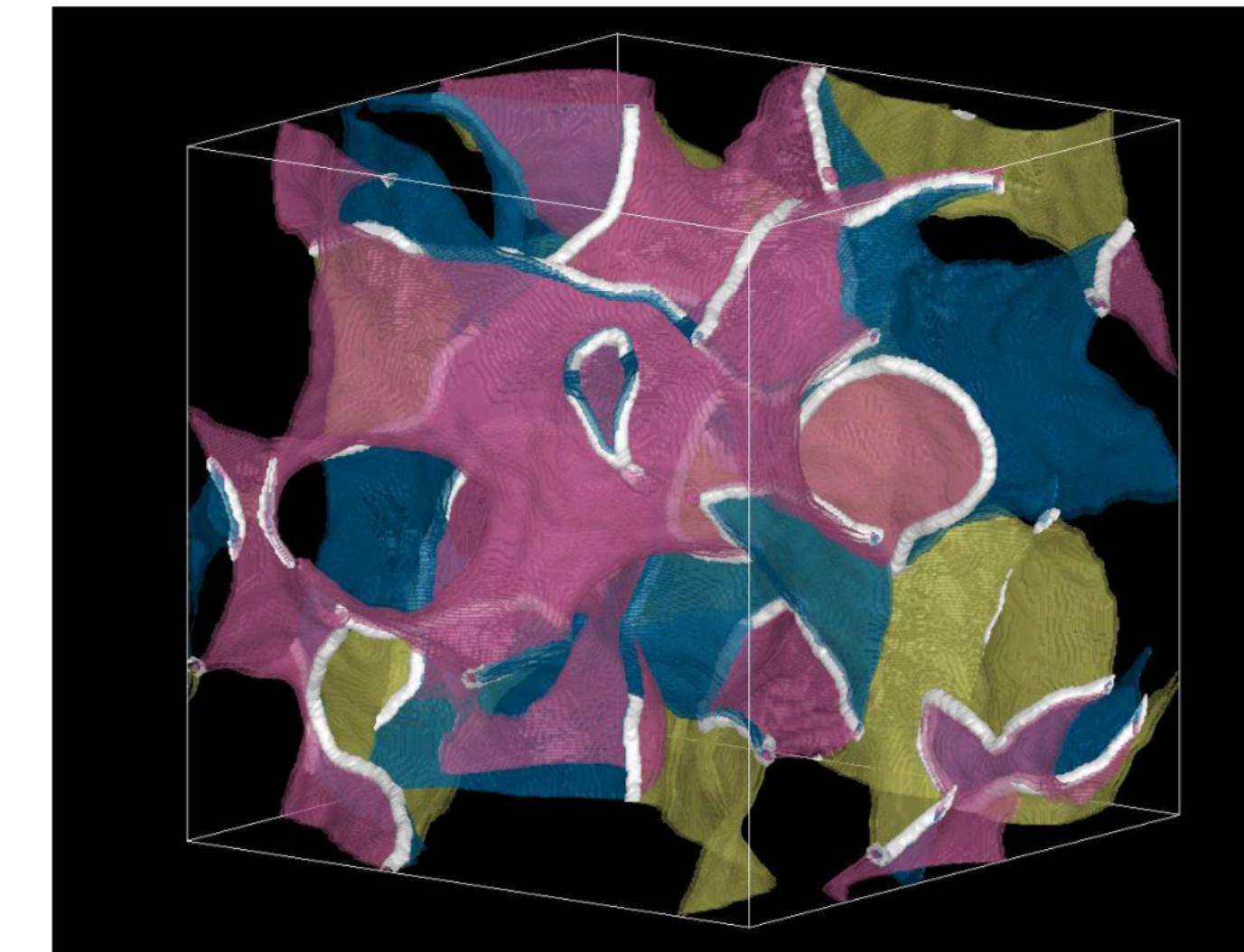
$$a_{DW} = \frac{\Delta V}{\sigma} \sim H(T_2)$$


**DWs quickly  
reach *scaling***

$$\rho_{DW} \propto \sigma_{DW} H$$



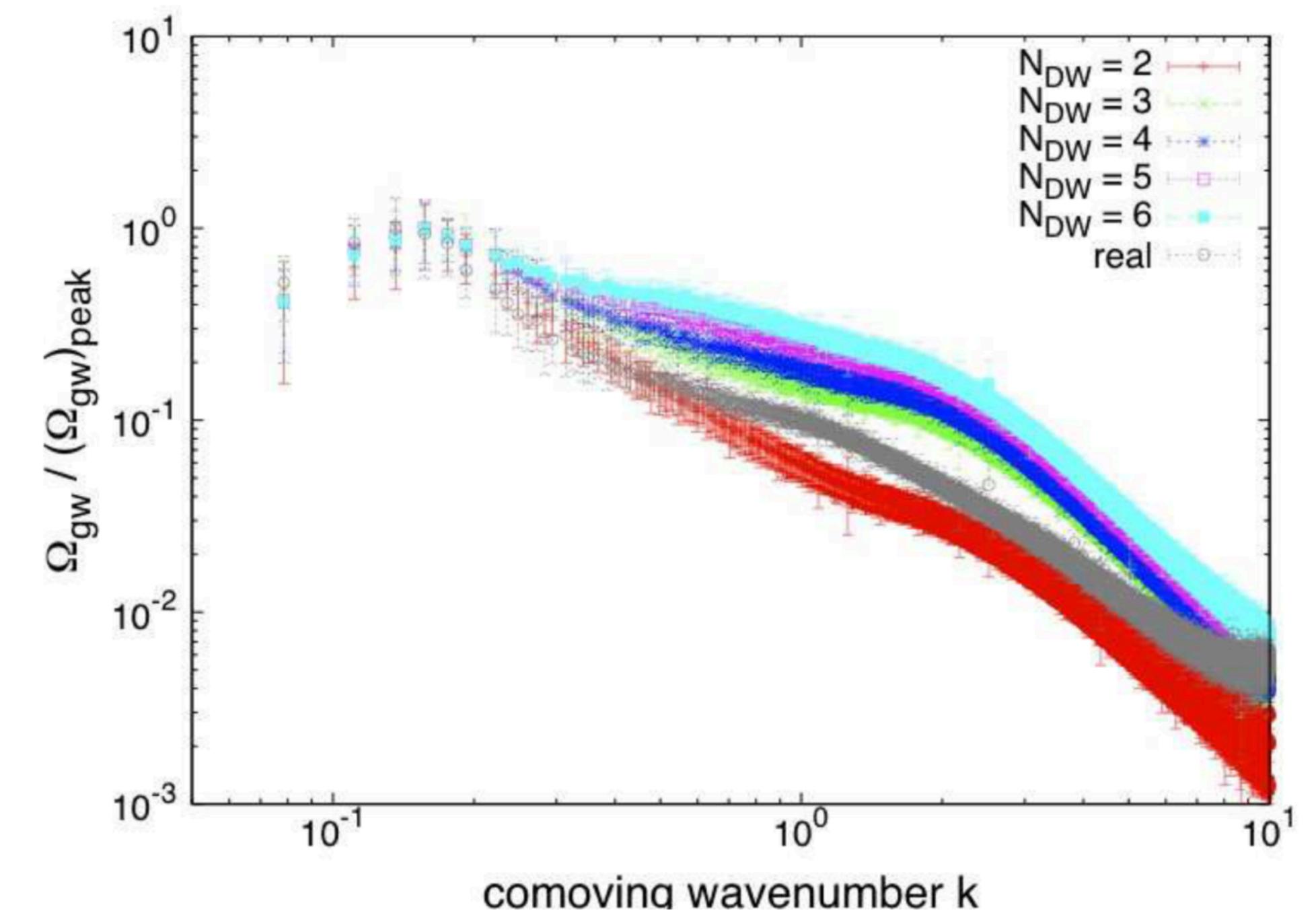
(a)  $N_{\text{DW}} = 2$



(b)  $N_{\text{DW}} = 3$

**GW spectrum  $\sim$   
broken power law  
 $w^3$  ,  $w^{-1}$**

**Hiramatsu et al arXiv:  
1207.3166**



**Peak in the GW spectrum from network annihilation corresponds to temperature**

$$T_2 \simeq \frac{10^7 \text{ GeV}}{\sqrt{\kappa_H}} \sqrt{\frac{10^{12} \text{ GeV}}{f}} \left( \frac{\Lambda_H}{10^{10} \text{ GeV}} \right) \left( \frac{r}{0.005} \right)^2$$

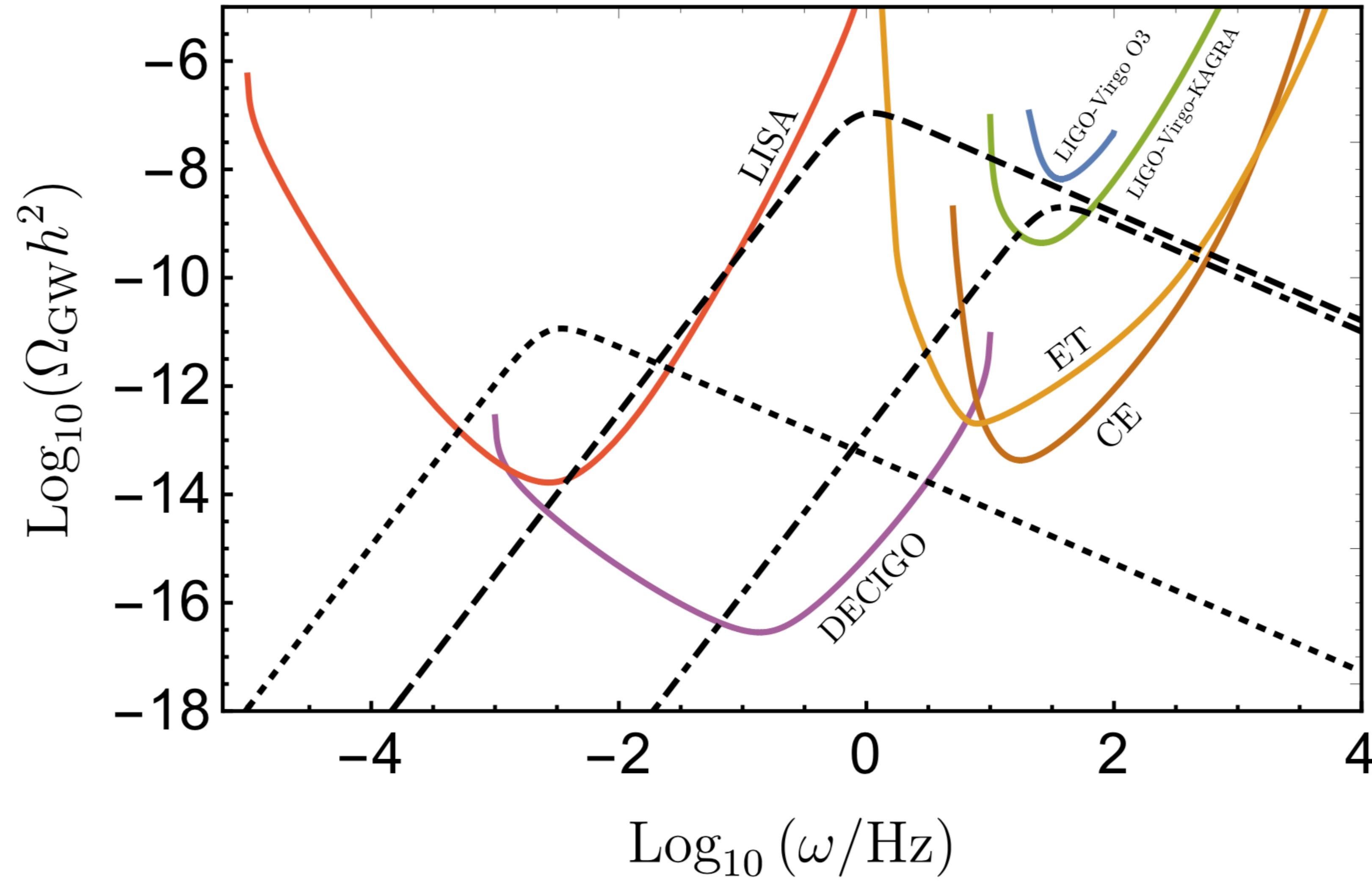
$$r = \Lambda_{mis}/\Lambda_H$$

$$\Lambda^4 \equiv \kappa_H^2 \Lambda_H^4$$

**Signal:**

$$\omega_{\text{peak}} \simeq \frac{5 \text{ Hz}}{\sqrt{\kappa_H}} \left( \frac{r}{0.005} \right)^2 \left( \frac{\Lambda_H}{10^{10} \text{ GeV}} \right) \sqrt{\frac{10^{11} \text{ GeV}}{f}}.$$

$$\Omega_{GW} \sim 10^{-7} \left( \frac{\rho_{DW}}{\rho_{rad}} \right)_{T=T_2}^2$$



$$\Lambda_{mis} = 10^7 \text{ GeV}$$

$$f = 10^{10} \text{ GeV}$$

$$\theta = 8 \cdot 10^{-13}$$

$$10^{10} \text{ GeV}$$

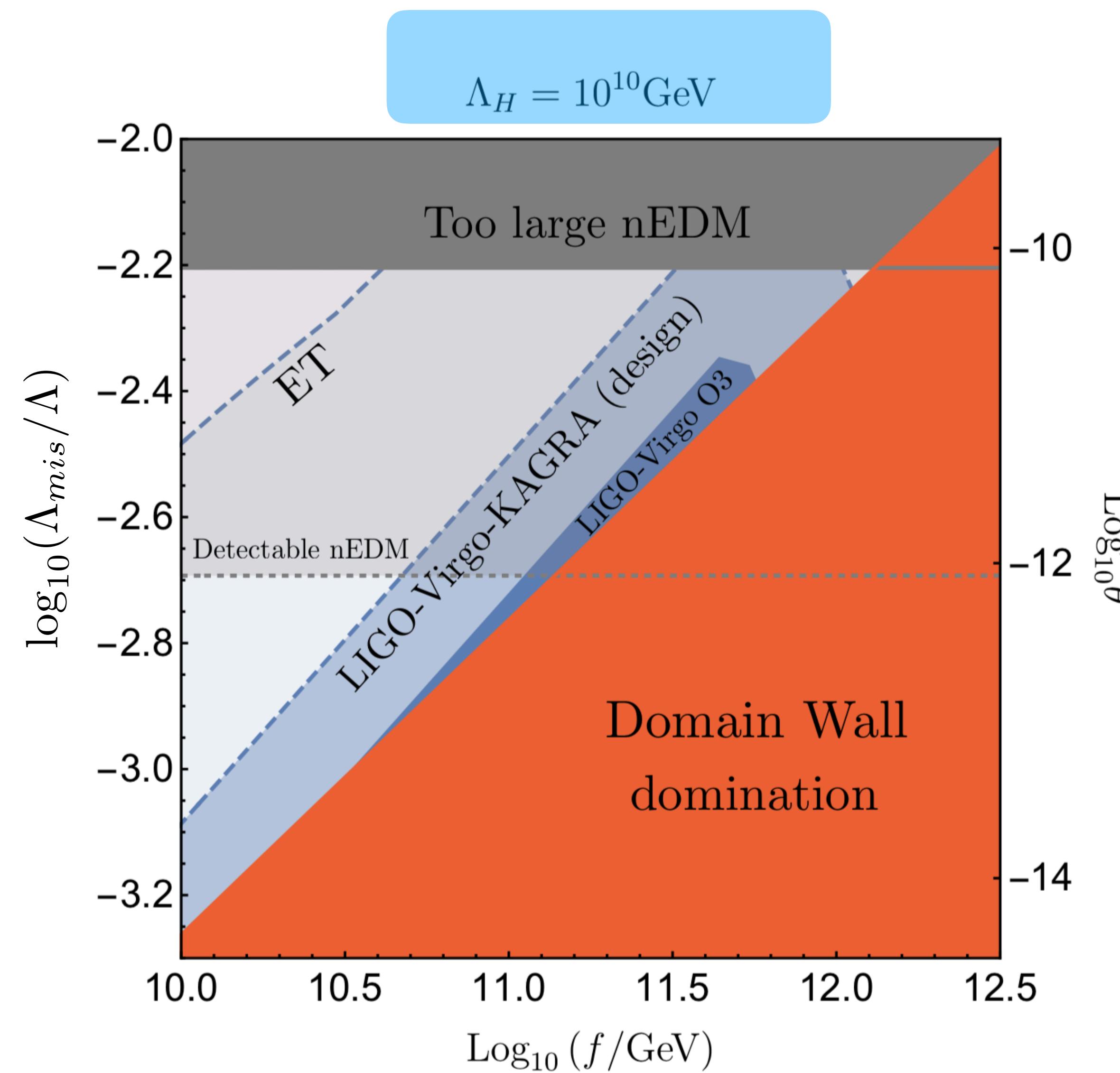
$$10^{11} \text{ GeV}$$

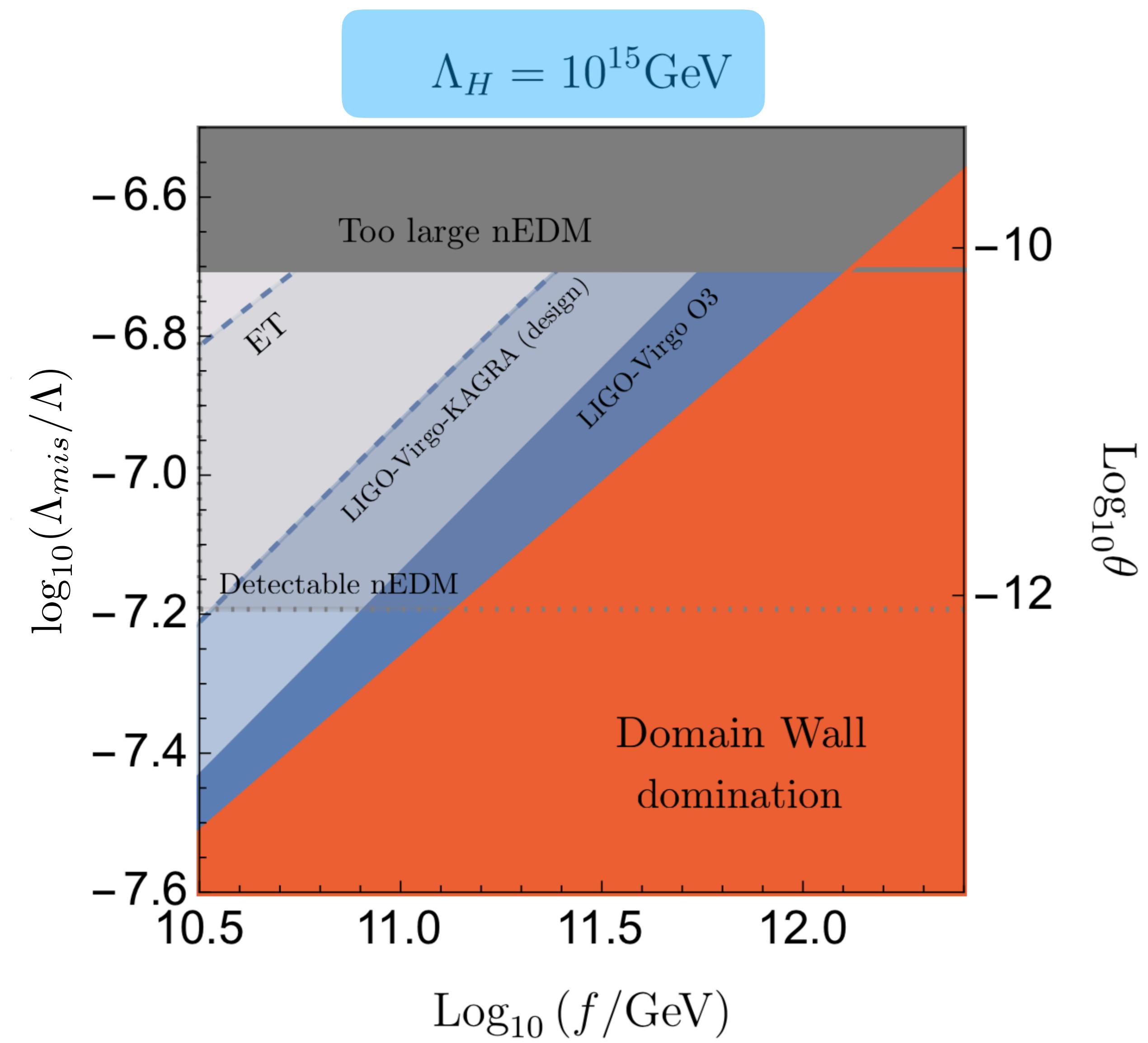
$$8 \cdot 10^{-13}$$

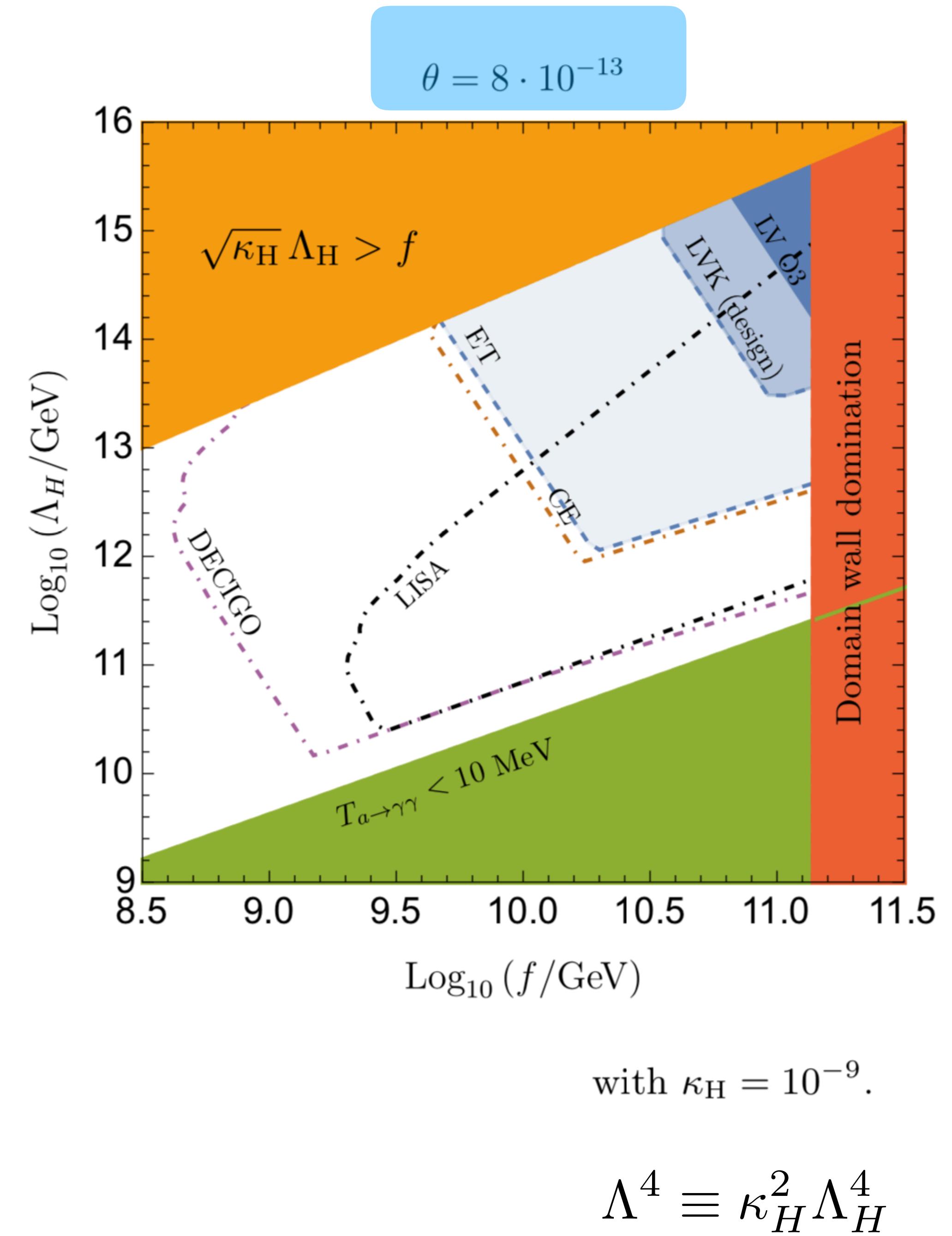
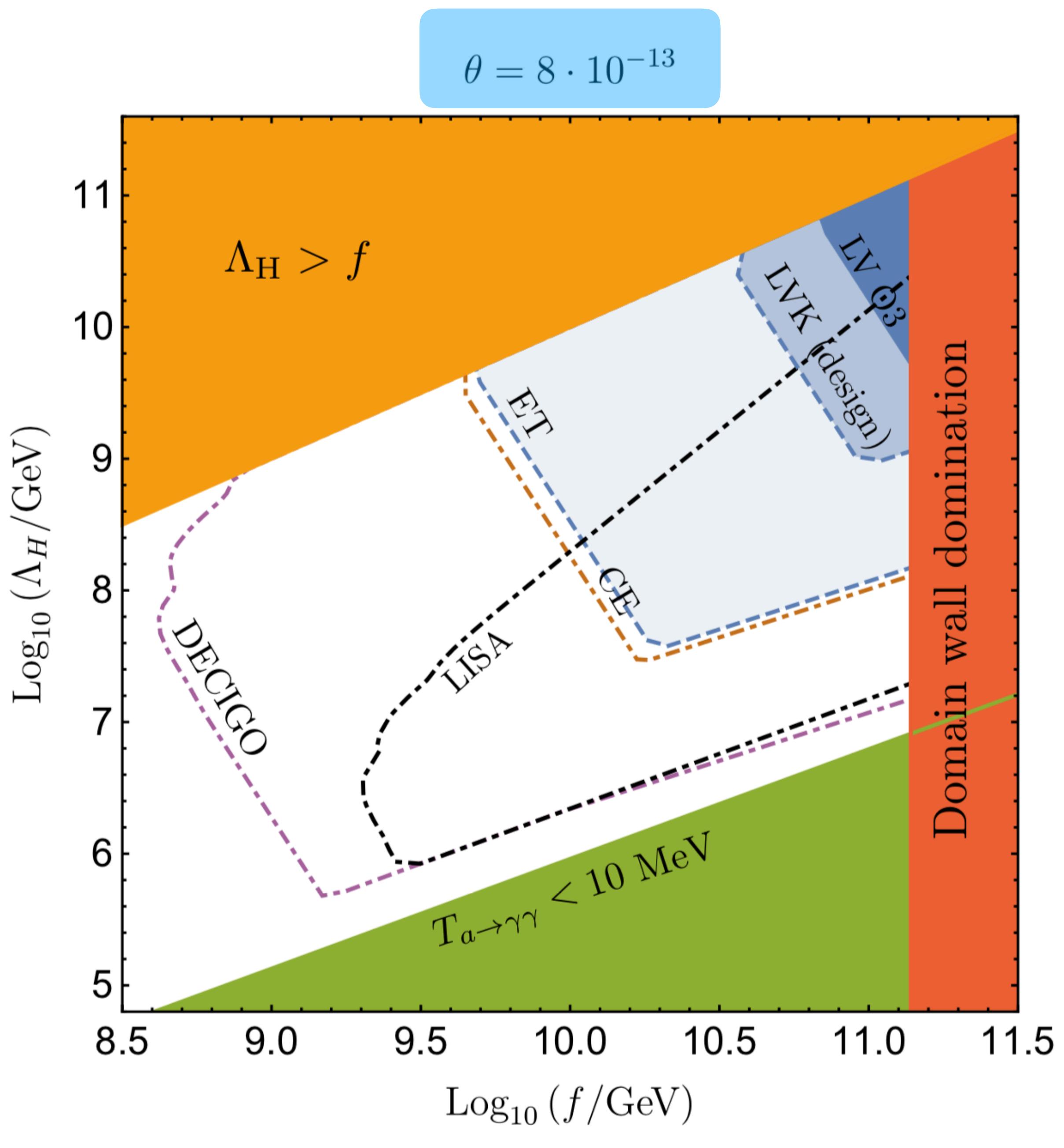
$$10^{11} \text{ GeV}$$

$$1.6 \cdot 10^{11} \text{ GeV}$$

$$1.5 \cdot 10^{-11}$$







Broken power law model			
	$f_* = 1 \text{ Hz}$	$f_* = 25 \text{ Hz}$	$f_* = 200 \text{ Hz}$
$n_2 = -1$	$3.3 \times 10^{-7}$	$3.5 \times 10^{-8}$	$2.8 \times 10^{-7}$
$n_2 = -2$	$8.2 \times 10^{-6}$	$6.0 \times 10^{-8}$	$3.7 \times 10^{-7}$
$n_2 = -4$	$5.2 \times 10^{-5}$	$1.8 \times 10^{-7}$	$3.7 \times 10^{-7}$

TABLE II: Upper limits for the energy density amplitude,  $\Omega_*^{95\%}$ , in the broken power law model for fixed values of the peak frequency,  $f_*$ , and negative power law index,  $n_2$ .

$$\Omega_{\text{bppl}}(f) = \Omega_* \left( \frac{f}{f_*} \right)^{n_1} \left[ 1 + \left( \frac{f}{f_*} \right)^\Delta \right]^{(n_2 - n_1)/\Delta}$$

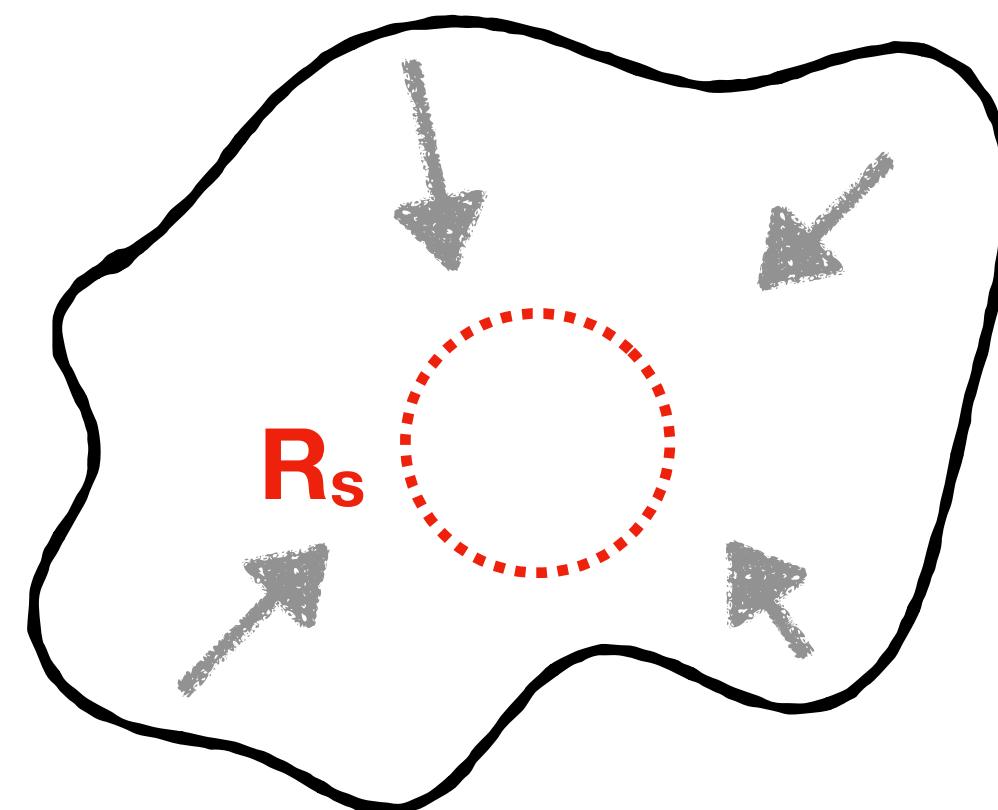
=> some models are already ruled out...

$$n_1 = 3 \quad n_2 = -1$$

# Primordial Black Holes

The same network collapse can lead to PBHs

Ferrer Masso Panico OP Rompineve  
Phys.Rev.Lett. 122 (2019) 10, 101301  
1807.01707



Collapse condition

$$R_S = \frac{\sigma}{H^2 M_P^2} \gtrsim \frac{1}{H}$$

Tension & pressure aid collapse

coincides with DW domination

near DW domination region, PBHs are expected

# Conclusions

**Heavy axion models naturally contain long DW epoch — very different from bubbles**

**DW network collapse generates strong GW signal (and PBHs)**

**Combined signature both in GWs and in nEDM**

**LIGO/Virgo/KAGRA are already testing interesting axion models!**

**LVK = extremely high-energy machines**

**Better simulations of bubble collisions and string/wall networks are needed to chart the excluded parameter space**

# **Thank you**

σ

## **and congratulations, Valery!**