

Demystifying Black Holes.

(Saturons: Black-hole-like bound states in ordinary theories)

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2003.05546 [hep-th]

1907.07332

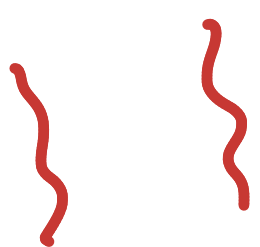
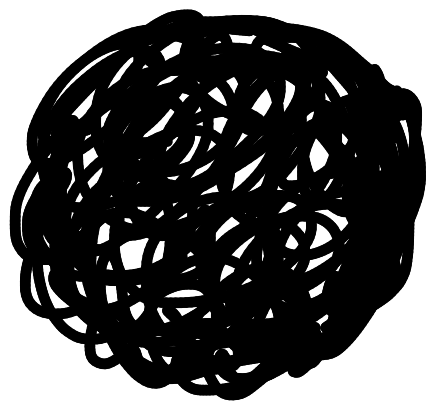
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+ with O. Kaikov, J. Valbuena '21

+ with O. Sakhelashvili

2111.03620 [hep-th]

Black holes are considered
to be mysterious



⊛ Hawking temperature

$$T = \frac{1}{R}$$

⊛ Bekenstein-Hawking entropy

$$S_{BH} = \frac{\text{Area}}{G_N}$$

⊛ Long time-scale of information
retrieval. (Page):

$$t \sim S R \sim \frac{R^3}{G_N}$$

We wish to show that all these "mysterious" properties are fully shared by objects that have maximal entropy compatible with unitarity.

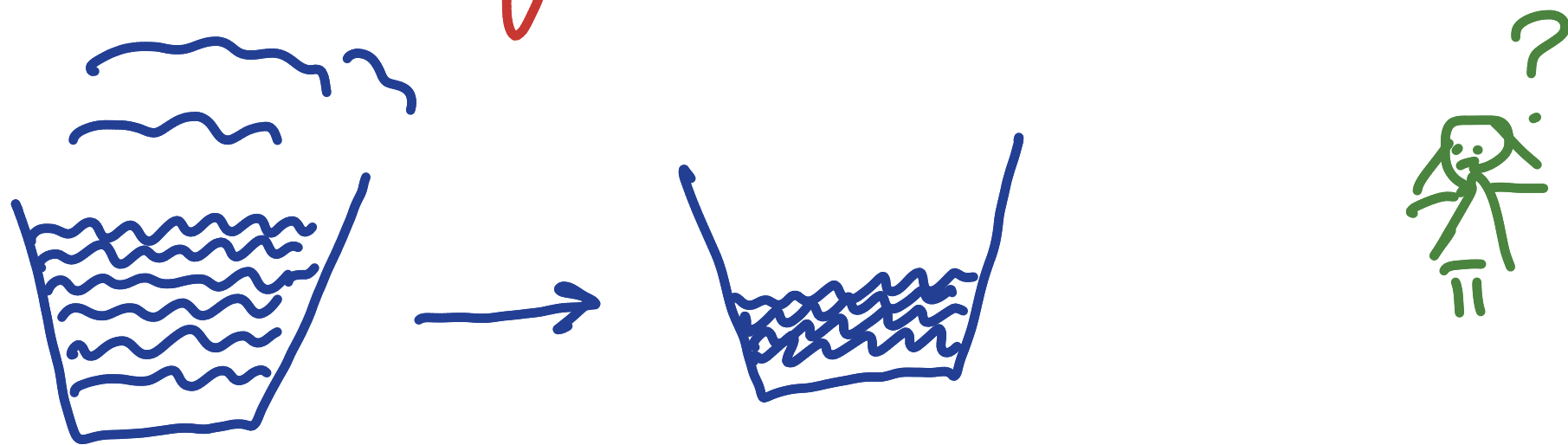
We call them "Saturns".

In particular, we shall demonstrate their existence in renormalizable calculable theories.

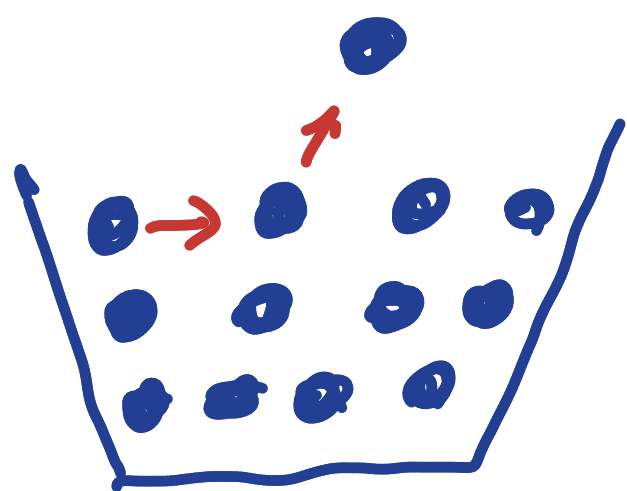
This indicates that:

Black hole = Saturn

When encountering a mysterious phenomenon of nature:



* Create a microscopic theory
(corpuscular resolution)



"N-Portrait" of Black hole
G.D. Gomez '11

* Try to capture universal properties



Water

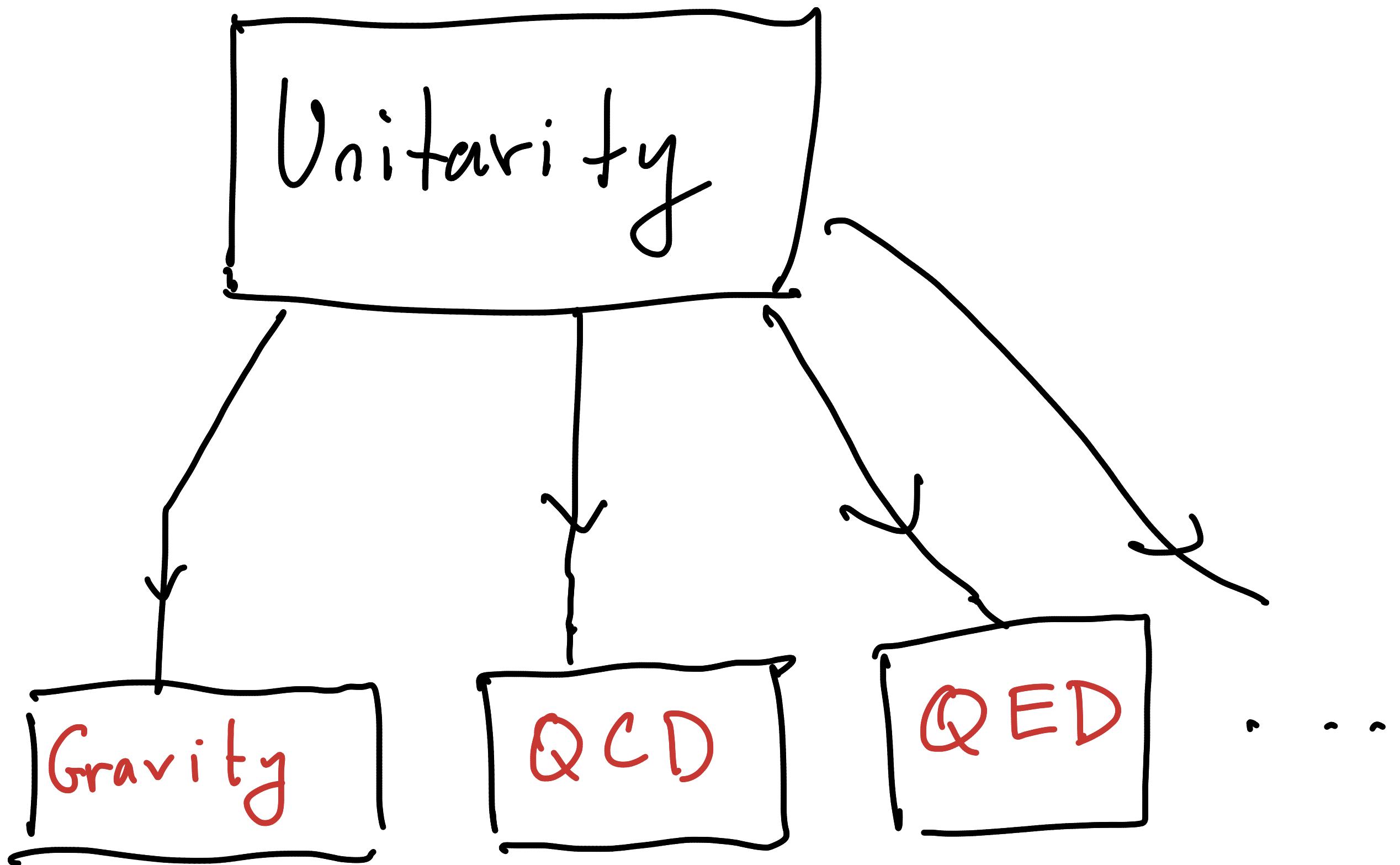


Mercury

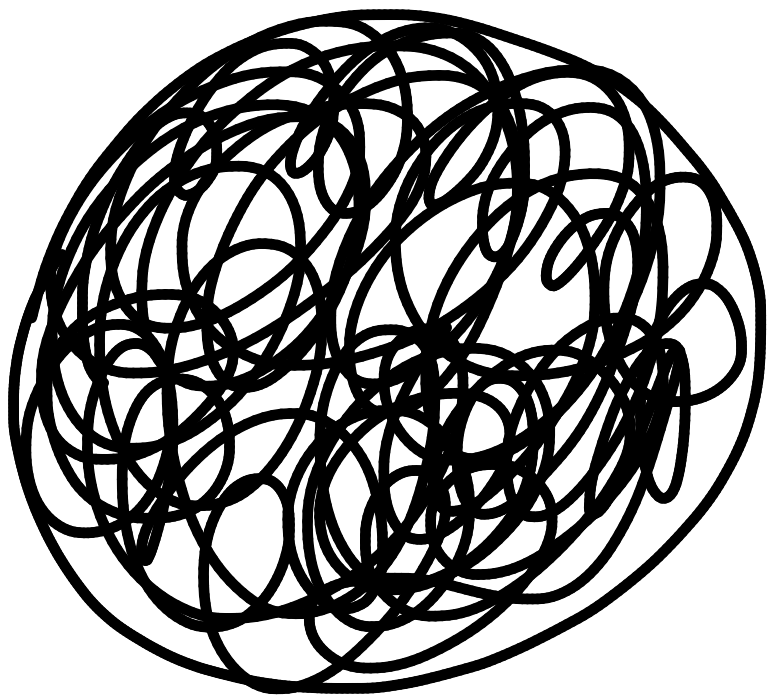


Ethanol

Most prominent constraint:



Unitarity bound on
entropy:



$k \sim R \rightarrow 1$

For any
self-sustained
object of size R

$$S \leq \frac{1}{\alpha}$$

$\alpha \equiv$ running quantum coupling
evaluated at scale $\frac{1}{R}$

Goldstone form of entropy bound

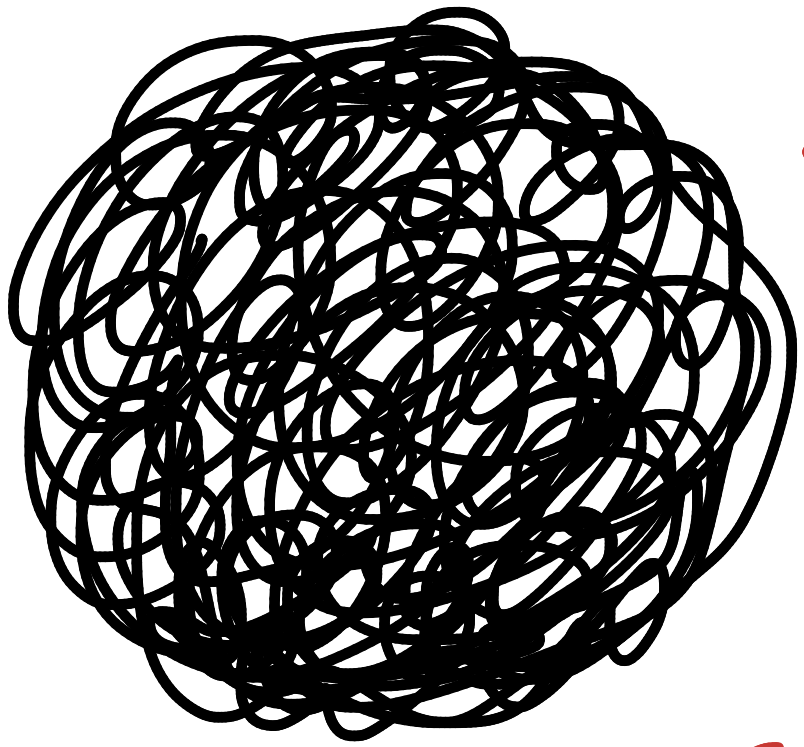
$$S \leq \frac{\text{Area}}{G_{\text{Gold}}}$$

$$\text{Area} \equiv R^{d-2}$$

$G_{\text{Gold}} \equiv$ Coupling of Goldstone

$$G_{\text{Gold}} \equiv \frac{1}{f^2}$$

$f \equiv$ Goldstone decay constant



← Any self-sustained
object breaks
Poincare symmetry
spontaneously.

G_{Gold} is unambiguously defined

For a boundstate of N
quanta of wavelength $\sim R$

$$G_{\text{Gold}} \equiv \frac{1}{f^2} \equiv \frac{1}{N} \cdot R^{d-2}$$

Note: States of high entropy contain other Goldstone bosons (see later), but Poincare Goldstone is universal.

Objects of maximal entropy:

$$S_{\text{MAX}} = \frac{1}{\alpha} = \frac{\text{Area}}{G_{\text{Gold}}}$$

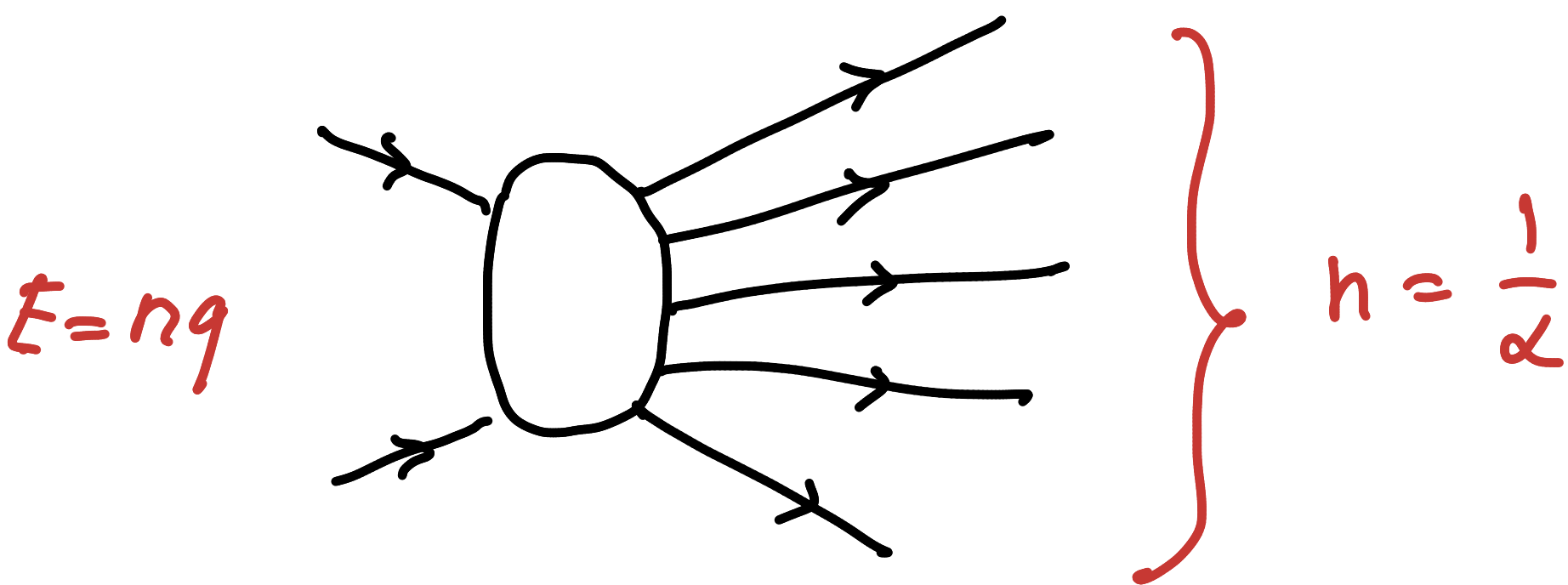
are called "Saturons".

Saturation of these entropy bounds is in one-to-one correspondence with saturation of unitarity by

$2 \rightarrow n$ scattering amplitudes

for $n = \frac{1}{\alpha}$

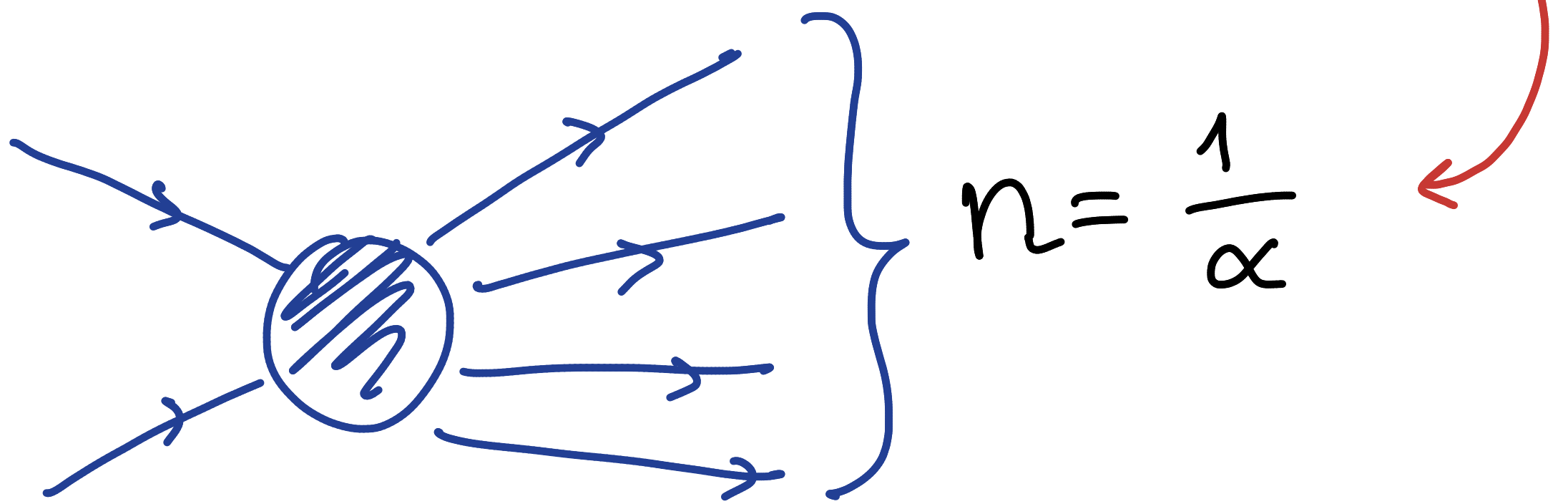
The point of optimal truncation.



momentum-transfer $q = \frac{1}{R}$

Cross section of $2 \rightarrow n$ scattering:

Optimal truncation



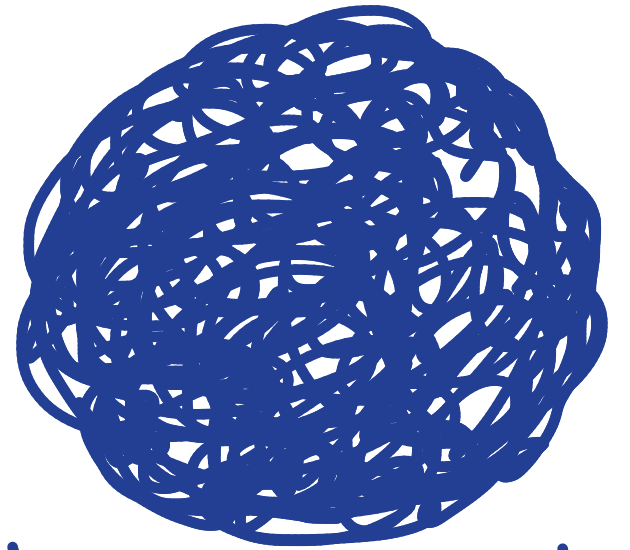
$$\sigma_{2 \rightarrow n} = e^{-\frac{1}{\alpha} + S}$$

Unitarity bound on entropy:

$$S \leq S_{\text{MAX}} = \frac{1}{\alpha}$$

Universal bound on time-scale
of start of information retrieval:

$$t_{\min} = \frac{R}{\alpha}$$



$| \leftarrow R \rightarrow |$

Due to saturation of
entropy bound can be written
as

$$t_{\min} = \rho R = R^3 f^2$$

$$t_{\min} \sim \text{Volume} \cdot f^2$$

All saturons have properties very similar to black holes:

- ① * Area-law entropy;
- ① * Thermal evaporation with $T = \frac{1}{R}$;
- ① * Information horizon;
- ① * Time of information retrieval
 $t_{\text{min}} \sim \frac{R}{\alpha} \sim \frac{1}{\alpha} R$
- ① * Saturate scattering amplitudes.

Black-hole-like bound states in $d=4$

A model

G.D., '20; G.D., Kaikov
& Valbuena '21

$SU(N)$ with adjoint Φ_β^α + fundamental ξ_α
 $\alpha, \beta = 1 \dots N$

Lagrangian (most general
renormalizable)

$$\mathcal{L} = \frac{1}{2} \text{Tr} \partial_\mu \Phi \partial^\mu \Phi + \frac{1}{2} \partial_\mu \xi^\dagger \partial^\mu \xi - V(\Phi, \xi)$$

$$V = \frac{\alpha}{2} \text{Tr} \left[f \Phi - \Phi^2 + \frac{\hat{1}}{N} \text{Tr} \Phi^2 \right]^2 +$$
$$+ \alpha_\xi \xi^\dagger \Phi^2 \xi + \tilde{\alpha}_\xi (\xi^\dagger \xi)^2 \text{Tr} \Phi^2 +$$
$$+ \bar{\alpha}_\xi (\xi^\dagger \xi)^2$$

Many degenerate vacua with spontaneously broken symmetry

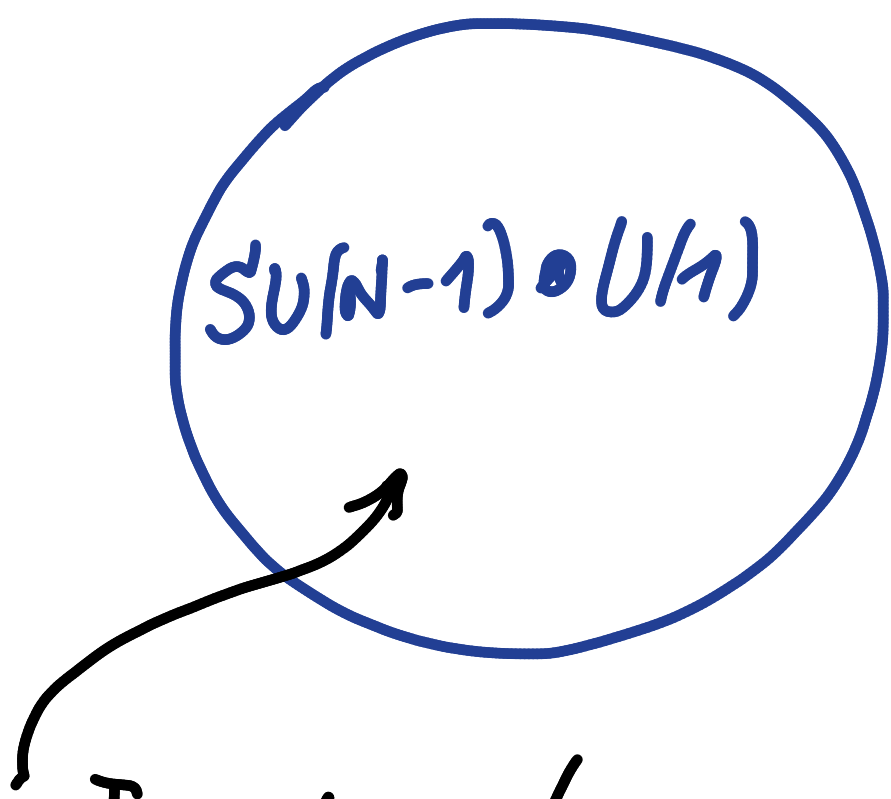
$$SU(N) \rightarrow SU(N-k) \otimes SU(k) \otimes U(1)$$

$$0 < k < N$$

We choose $SU(N)$ invariant vacuum as asymptotic S -matrix vacuum.

In this vacuum, there exist an infinite tower of bound states:

Bubbles of $SU(N-1) \otimes U(1)$ -vacuum:



$SU(N)$



$$\Phi = \text{diag}(N-1, -1, -1, \dots, -1)$$

$$S \sim (N-1) \cdot U(N)$$

In outer vacuum there is a mass gap

$$m^2 = \alpha f^2$$

$N_G \sim N$ Goldstone species Θ^a , $a=1,2,\dots, 2N$

$f \equiv$ Goldstone decay constant

$$G_{\text{odd}} = \frac{1}{f^2} \leftarrow \text{Goldstone coupling}$$

There exists an infinite tower of classically - stable bubbles.

They represent bound states of Goldstones with total occupation number n .

The spectrum of bound-states

$$E_n = n \sqrt{\frac{m}{R}}$$

One can think of them as sort of non-topological solitons, or Q-balls.

However, they are not ordinary solitons, due to degeneracy:

Level E_n has microstate degeneracy:

$$N_{st} = \left(1 + \frac{2N}{n}\right)^n \left(1 + \frac{n}{2N}\right)^{2N}$$

Correspondingly, they carry
a microstate entropy

$$S = \ln(\Omega_{st}) = 2N \ln \left(1 + \frac{2N}{n}\right)^{\frac{n}{2N}} \left(1 + \frac{n}{2N}\right)$$

Now, theory saturates unitarity for

$$N\alpha \sim 1$$

Thus, entropy is saturating unitarity
bound $S \sim \frac{1}{\alpha}$ for

$$n \sim N \sim \frac{1}{\alpha} \leftarrow \text{saturation point}$$

Thus, entropy of saturation is

$$S \sim \frac{1}{\alpha} \sim N \sim n$$

Notice, Goldstone decay constant
both for Poincare and for $SU(N)$
is

$$f^2 = \frac{N}{R^2}$$

Goldstone coupling:

$$G_{\text{Gold}} = \frac{1}{f^2} = \frac{R^2}{N}$$

$$\alpha_{\text{Gold}} = \frac{1}{N}$$

Parameters of the saturated bound state:

① * Mass $M = \frac{N}{R} = \frac{l}{\alpha R} = \frac{S}{R}$

① * Entropy

$$S = \frac{1}{\alpha} = \frac{\text{Area}}{G_{\text{Gold}}} = N$$

① * Size

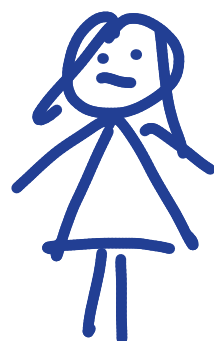
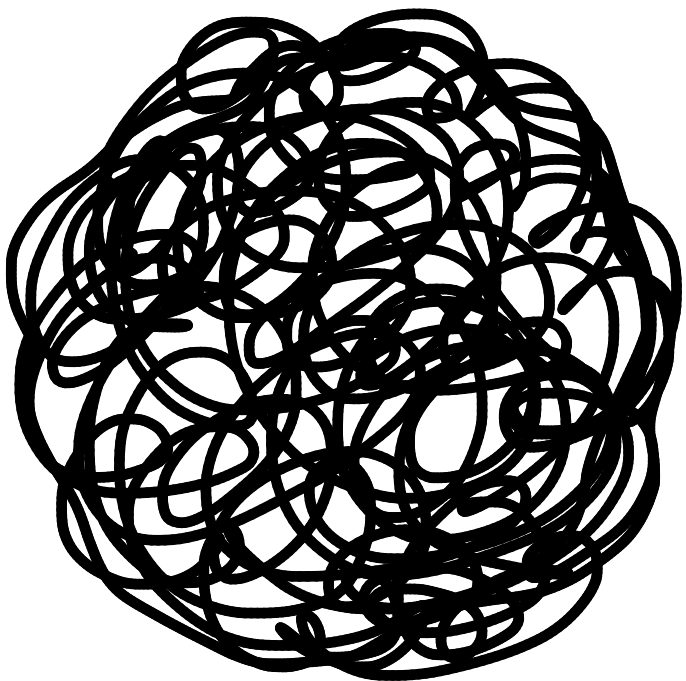
$$R = m^{-1}$$

In quantum theory bound state decays.

The decay rate of a saturated bubble:

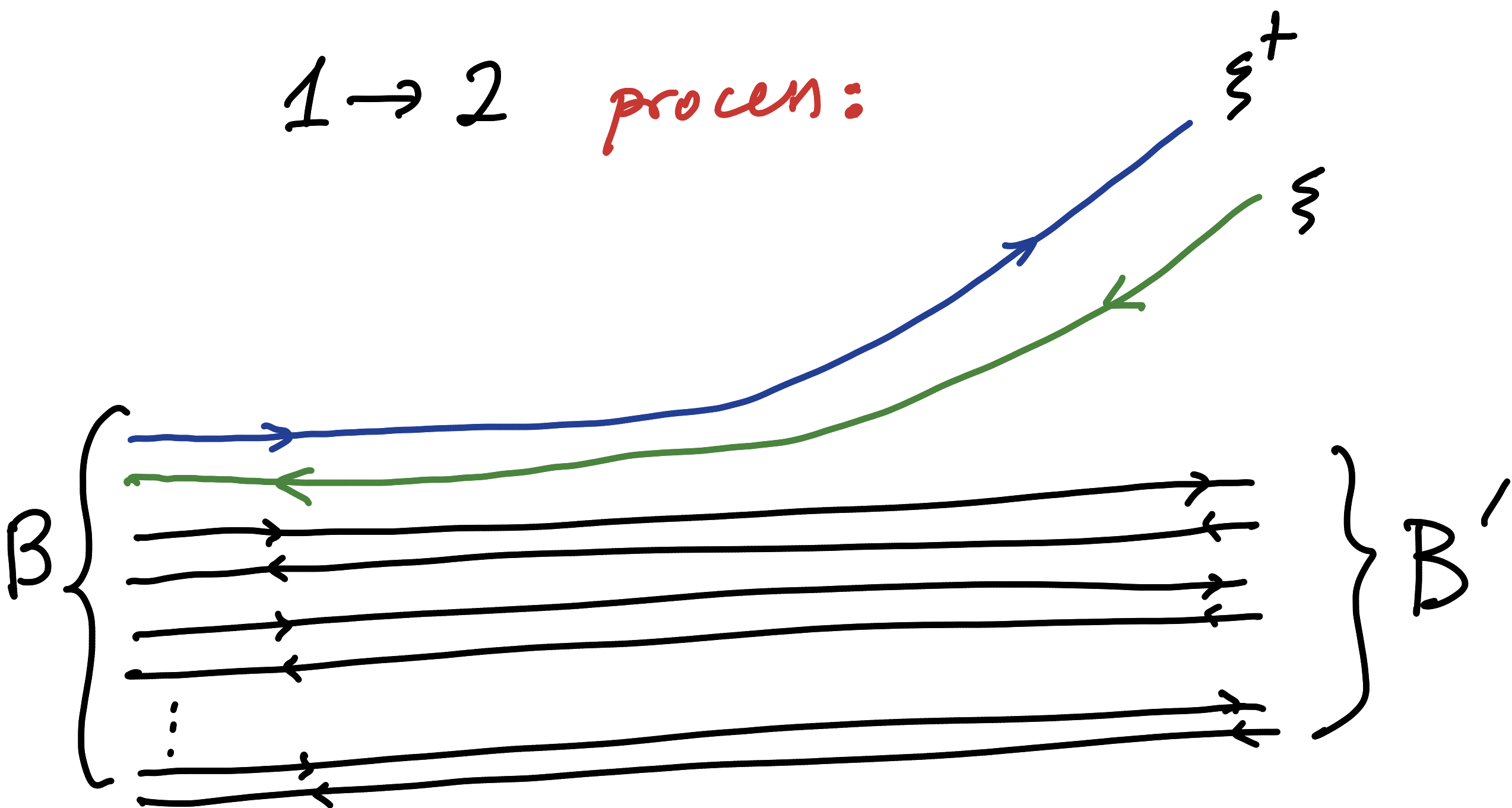
$$\Gamma \sim \frac{1}{R} \leftarrow \text{Hawking rate with temperature}$$

$$T = \frac{1}{R} !$$



Emission of massive ξ -quanta

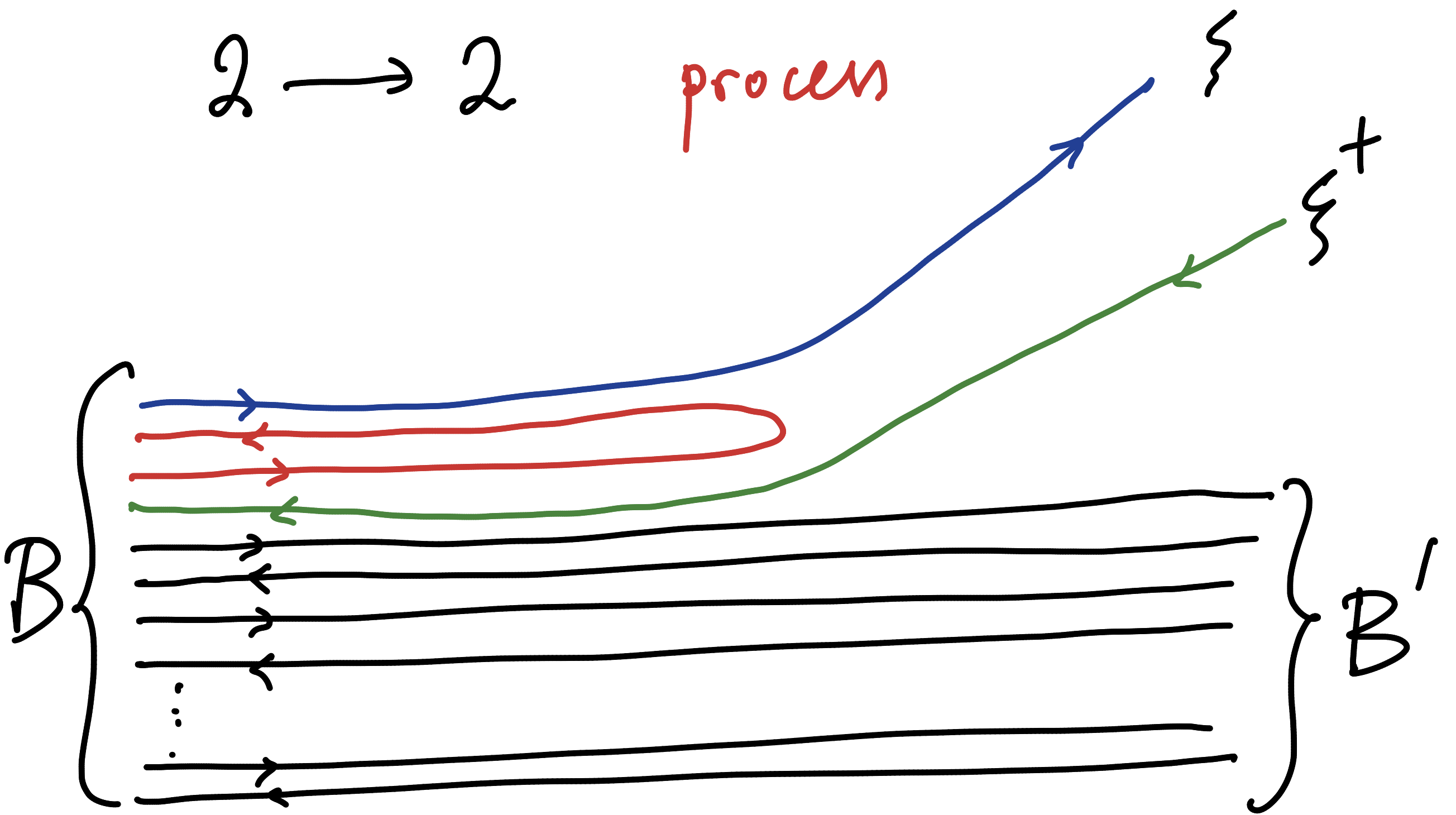
$1 \rightarrow 2$ process:



$$\Gamma \sim \frac{1}{R} \underbrace{\alpha \cdot N}_{\substack{\text{Recall,} \\ \alpha N = 1}} \sim \frac{1}{R}$$

Recall, $\alpha N = 1$

2 \rightarrow 2 process



$$\Gamma \sim \frac{1}{R} \alpha^2 N^2 \sim \frac{1}{R}$$

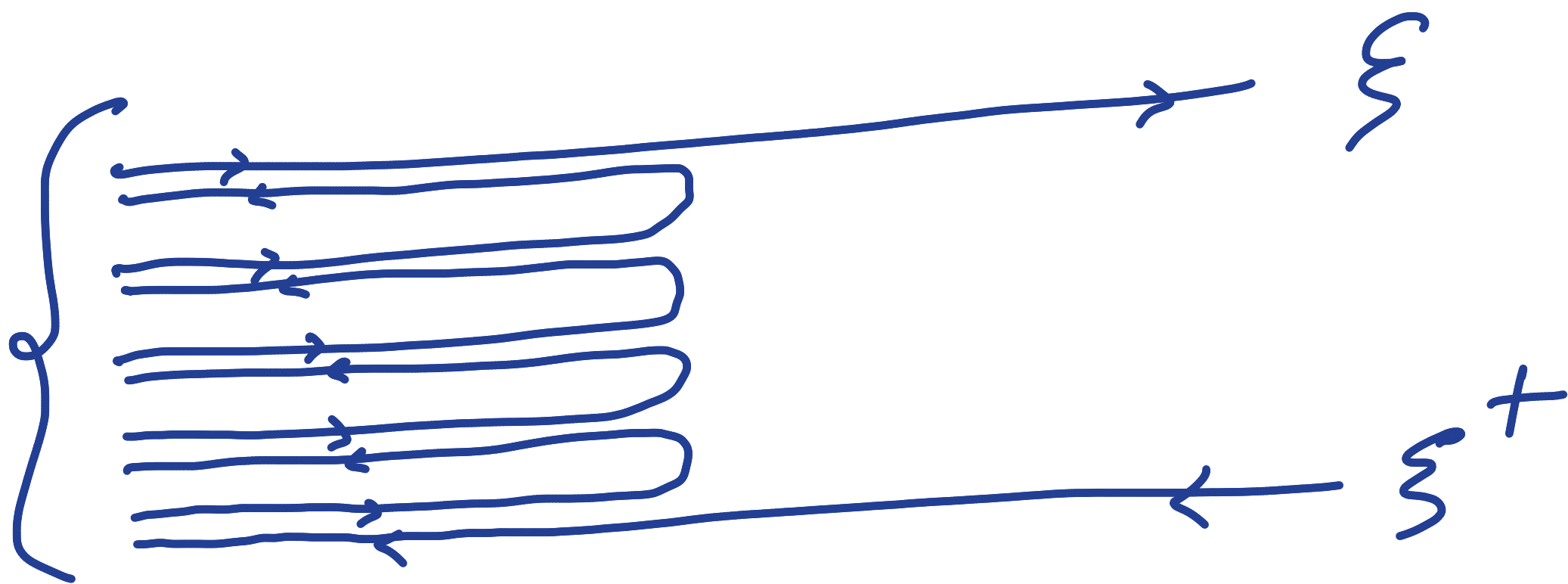
$\alpha N = 1$ \leftarrow saturation condition.

Notice, emission of more energetic quanta is "Boltzmann" suppressed:

$$E_{\xi+\xi^+} \gg \frac{1}{R}$$

rate is exponentially suppressed

$$\Gamma \sim e^{-\frac{E_{\xi+\xi^+}}{T}} \quad T = \frac{1}{R}$$



Thus, a satoron evaporates
as a black hole with
thermal-like rate



In reality the state is
pure: Information is carried

by $\frac{1}{N} \sim \frac{1}{S}$ corrections.

Quantum information is stored in flavor content of the bound state.

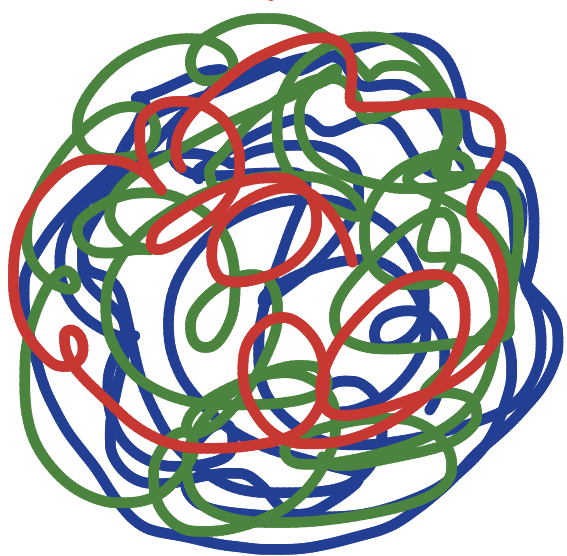
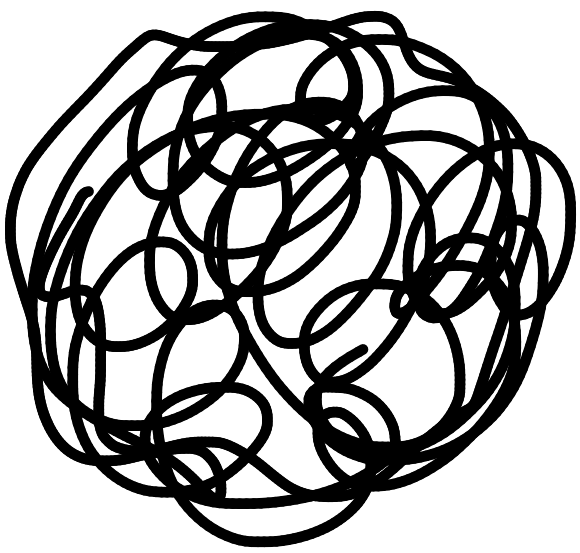
In the semi-classical limit: $N \rightarrow \infty$, $\alpha \rightarrow 0$,

$$(N\alpha) = 1 \quad R = \text{finite}.$$

It is unreadable: Satoron has information horizon.

For $N = \text{finite}$, the information is readable after some time.

Time scale of the start of
information retrieval:



$$t_{\min} \sim \frac{R}{\alpha} \sim \frac{R}{\rho R} \sim \frac{R^3}{G_{\text{Gold}}}$$

identical to
Page's time for
a black hole

Black hole / sataron
correspondence:

$$G_N \longrightarrow G_{\text{Gold}}$$

All characteristics are
identical:

$$\begin{array}{ccc} \alpha_{\text{gr}} & \longleftrightarrow & \alpha \\ S_{\text{BH}} & \longleftrightarrow & S \\ T_{\text{H}} & \longleftrightarrow & T \\ t_{\text{Page}} & \longleftrightarrow & t_{\text{min}} \end{array}$$

Saturons in $d=2$

G.A. Sakhelashvili '21

Gross-Neveu Model

$$\mathcal{L} = i\bar{\psi}\not{\partial}\psi + \frac{\alpha}{2}(\bar{\psi}\psi)^2$$

$2N$ Majorana fermions with
flavor group $SO(2N)$

The model is asymptotically
free and the spectrum of
bound states is known.

The spectrum of bound states
(exact in $N \rightarrow \infty$):

Boundstate of n fermions
($n < N$)

$$M_h = m_f \frac{2N}{\pi} \sin\left(\frac{n\pi}{2N}\right)$$

is anti-symmetric tensor of $SO(2N)$

the degeneracy:

$$h_{st} = \frac{2N!}{n! (2N-n)!}$$

The boundstate of maximal
degeneracy $n = N - 1$:

$$M_n = \frac{2N}{\sqrt{2}} m_f \quad \text{has size } R = m_f^{-1}$$

Entropy

$$S = \ln(n_{st}) = 2N \ln(2)$$

Saturates unitarity bound
for

$$N \sim \frac{1}{\alpha}$$

The cross section of

vacuum $\rightarrow 2N$

process scales as

$$\sigma = e^{-\frac{2}{\alpha}} + S$$

Can be interpreted as the process of sataron pair creation

Very similar to saturation of $2 \rightarrow N$ graviton process by black holes

G.D., Gomez, Isermann, Lüst, Stieberger '14
Addazi, Bianchi, Veneziano '16

At the saturation point the entropy of a bound state exhibits complete correspondence with Bekenstein-Hawking entropy of a black hole

$$S = \frac{1}{\alpha} = \frac{\text{Area}}{G_{\text{Gold}}} = \frac{1}{G_{\text{Gold}}}$$

The quantity
"Area" = R^{d-2} is well
defined for $d=2$

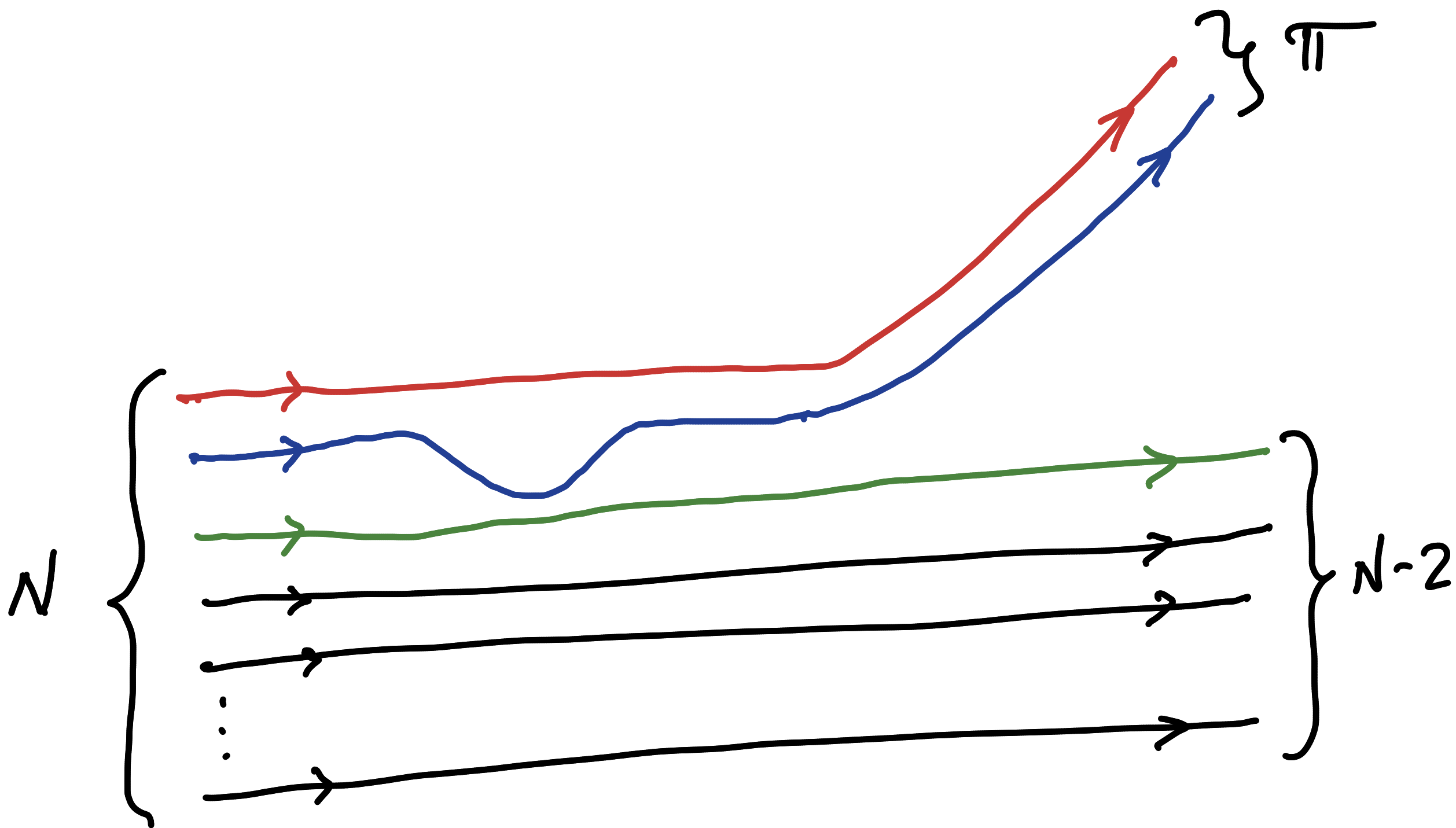
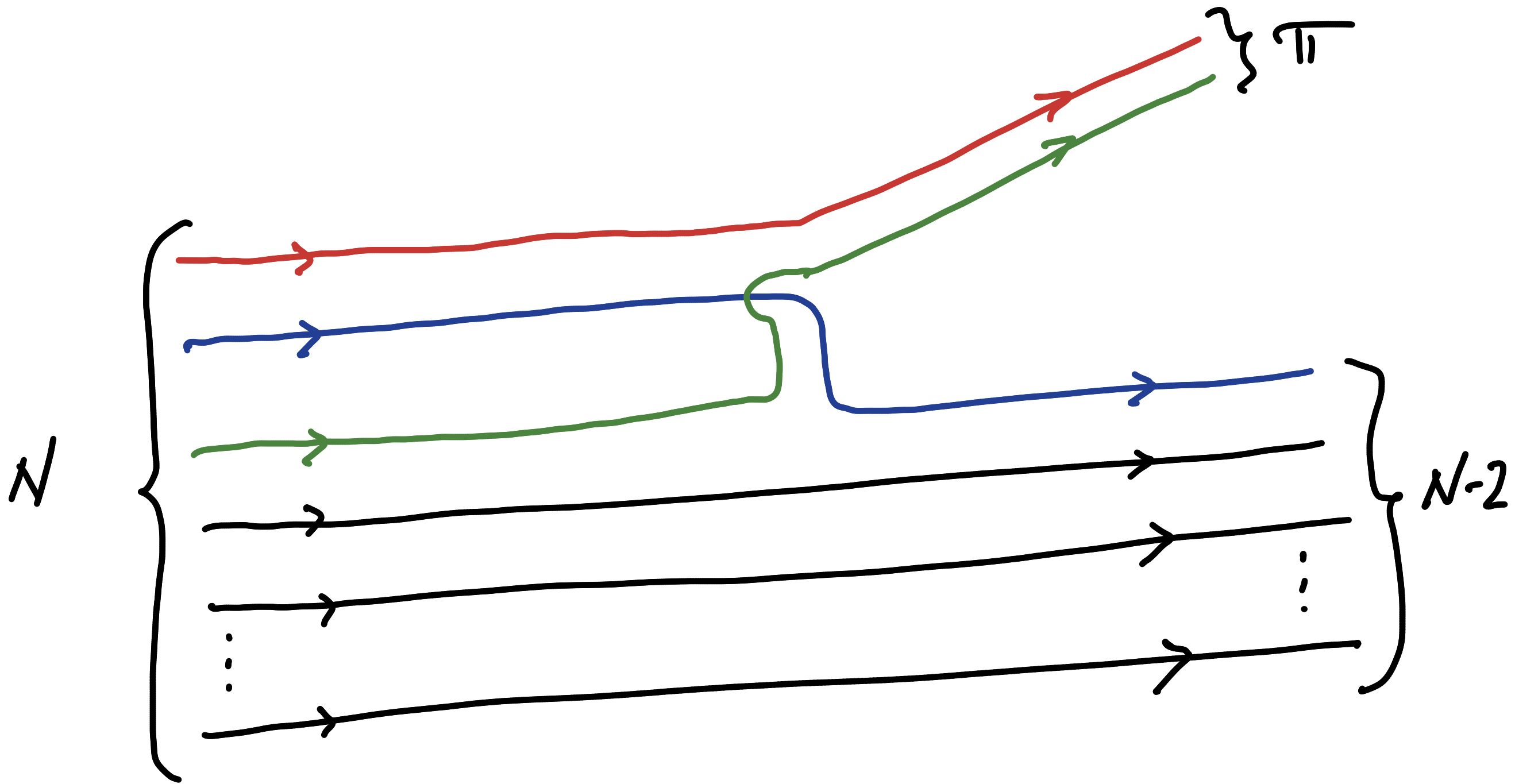
Hawking evaporation of
GN satoron.

Coupling to pion:

$$\mathcal{L}_\pi = \partial_\mu \pi_{ij} \partial^\mu \pi_{ij} + \frac{M_f}{f_\pi} \pi_{ij} \bar{\Psi}_i \gamma_5 \Psi_j$$

The satoron emits pion
and evaporates with the
rate.

$$\Gamma \sim \frac{1}{R} \alpha^2 N^2 \sim \frac{1}{R}$$



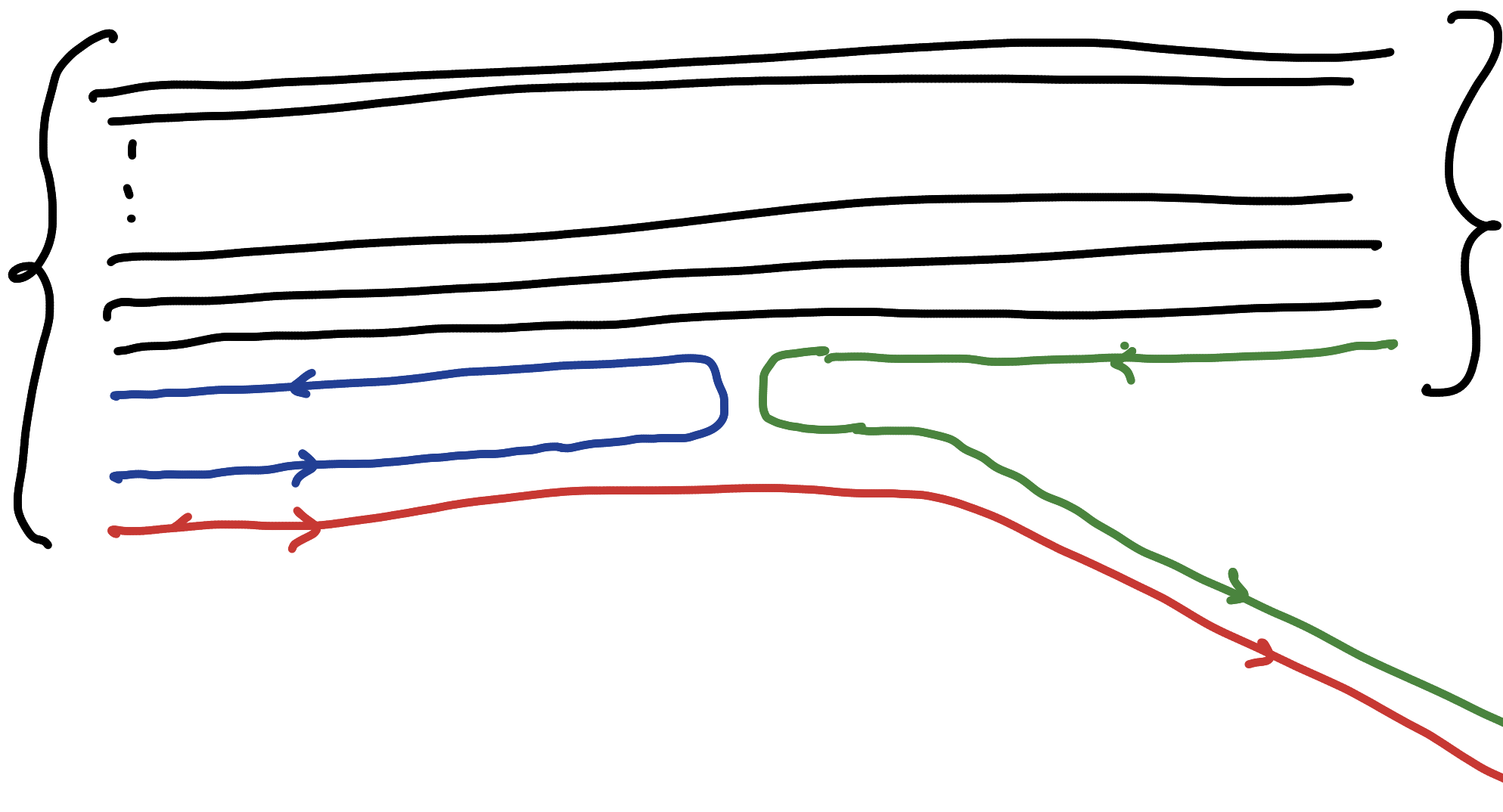
rate $\Gamma \sim \frac{1}{R} \propto \frac{1}{f_{\parallel}^2} N^3 \sim \frac{1}{R}$

We are dealing with thermal evaporation rate of temperature

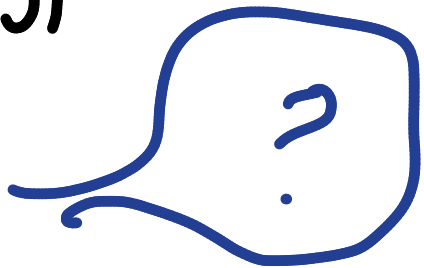
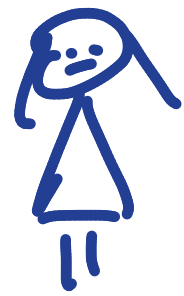
$$T = \frac{1}{R}$$

The information is hidden in $\frac{1}{N} \sim \frac{1}{S}$ corrections.

It is carried by flavor quantum numbers of pions.



π



Each pion carries a tiny fraction of Saturn's information.

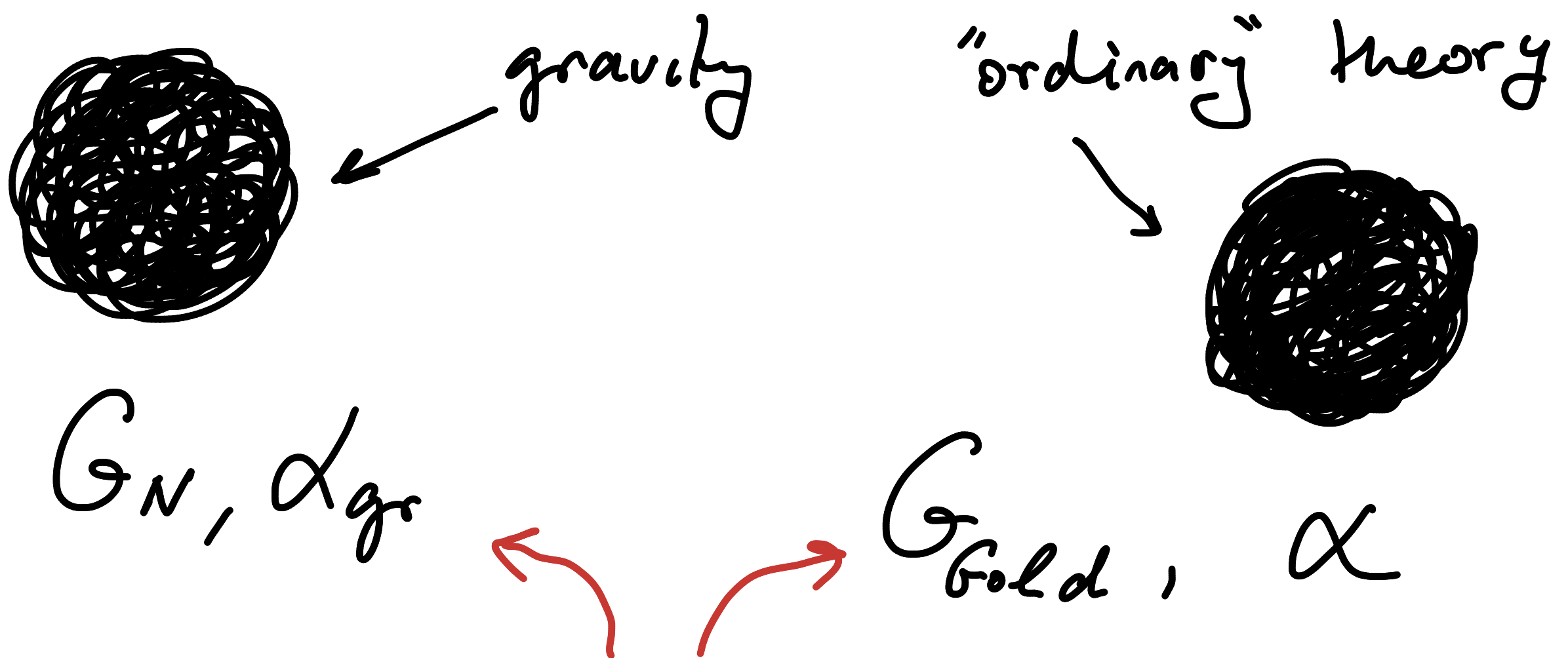
Alice has to collect $\sim N$ pions and decode their flavor.

This is only a start of information retrieval:

$$t_{\min} \sim N \cdot R \sim S'R \sim \frac{R}{\alpha} \sim \frac{R^{d-1}}{G_{\text{Gold}}}$$

Black hole/saturn correspondence
is independent of dimensionality
and other details of the
theory.

It is correspondance between
the states in different theories



Trans-theoretic parameters

We have seen that mysterious black hole properties are not rooted in gravity.

Rather, they represent generic features of states that saturate unitarity bound on entropy, saturons.

There exist many examples even in ordinary renormalizable theories, which can be reliably studied at weak coupling.

From there, we can predict new features for black holes.

Observing saturons in calculable theories, we see what are the wrong assumptions made in ordinary (semi-classical) treatment of black holes:

⊛ Evaporation is never thermal for finite S . $1/S$ -corrections break thermality.

⊛ Black hole evaporation is not self-similar.



is broken by "memory burden" effect.