Small-x Helicity Phenomenology

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Proton Spin Puzzle

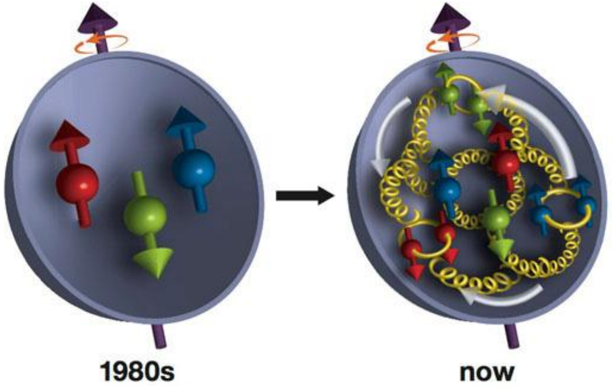
Jaffe-Manohar Spin Sum Rule:

$$\frac{1}{2} = S_q + L_q + S_g + L_g$$

$S_{q,g}$ = Helicity of quarks and gluons

$L_{q,g}$ = Orbital angular momentum

S_q ~ 30% of proton spin!

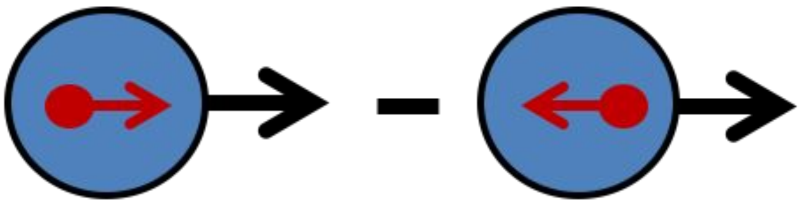


Quark Helicity Parton Distribution Functions

$$S_q(Q^2) = \frac{1}{2}\int_0^1 dx \sum_q (\Delta q(x,Q^2) + \Delta\bar{q}(x,Q^2))$$

Helicity PDFs:

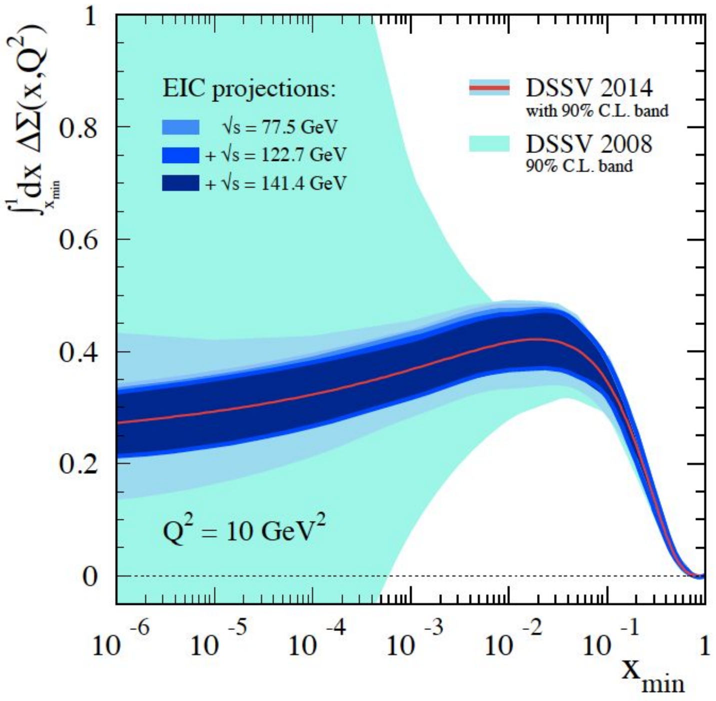
$$\Delta q =$$



- Q^2 = resolution at which we probe the proton
- Bjorken $x \sim \frac{1}{s}$. We need theory to extrapolate to x=0

Quark hPDF - DGLAP extraction

2 x (quark spin)



$$\Delta\Sigma = \sum_q (\Delta q + \Delta\bar{q})$$

- E. Aschenauer et al, **arXiv:1509.06489 [hep-ph]**, (DSSV = de Florian, Sassot, Stratmann, Vogelsang, DGLAP-based helicity PDF extraction from data)
- Large uncertainty at small-x!

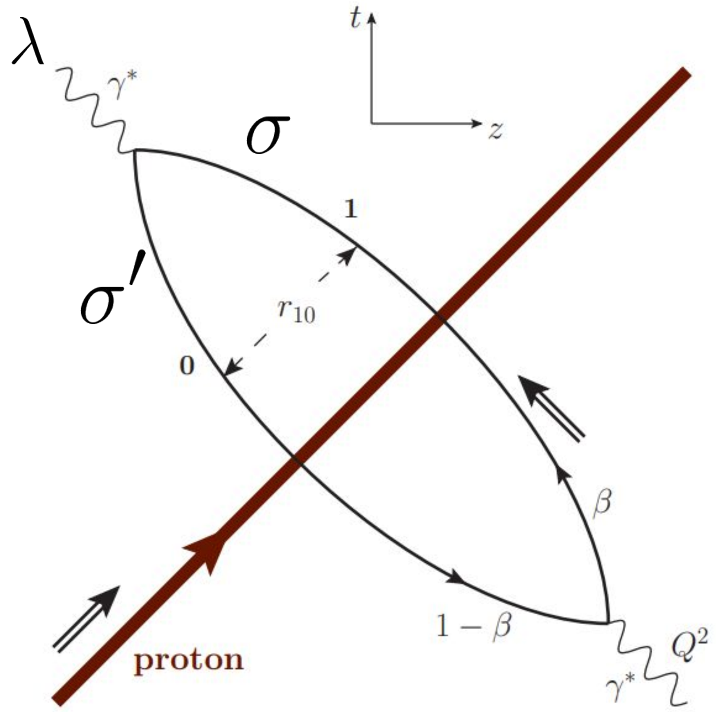
The Plan

Any complete description of quark and gluon helicity needs to

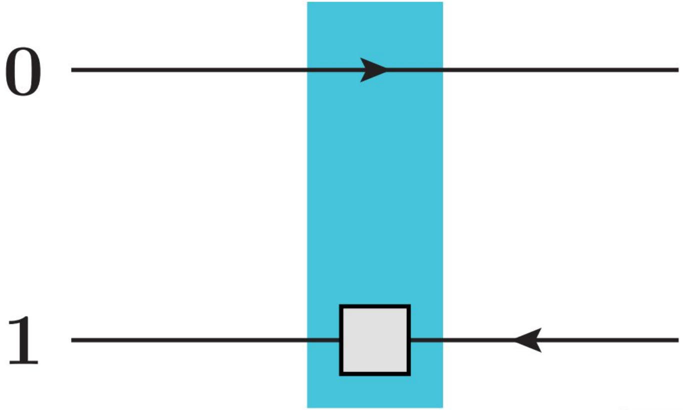
- Describe existing data $(5 \times 10^{-3} < x < 0.7)$
- Predict future, e.g EIC, data $(4 \times 10^{-3} < x < 5 \times 10^{-3})$
- Compare with said data
- Extrapolate down to $x = 0$
- While maintaining good control over theoretical uncertainty

(Polarized) DIS in the (Polarized) Dipole Picture

$$g_1 \propto |\psi|^2 \otimes (Q + 2G_2)$$



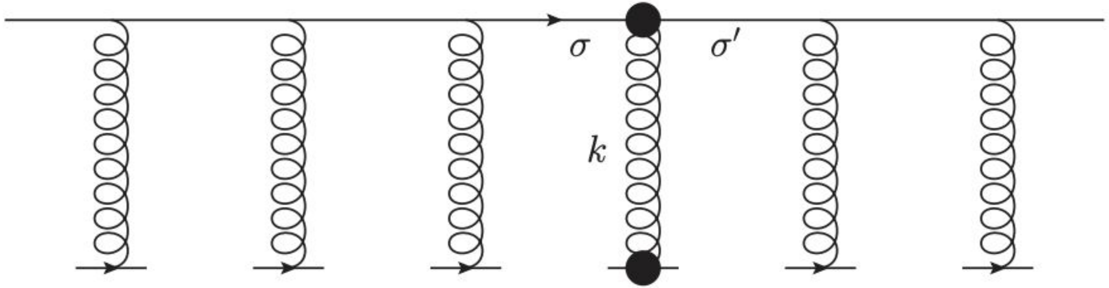
(Polarized) DIS in the (Polarized) Dipole Picture



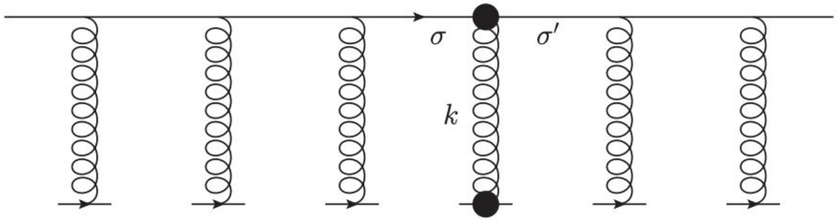
- In pDIS, the electron and proton have their helicity specified
- Cross-section now dependent on **Polarized Dipole Amplitudes**:

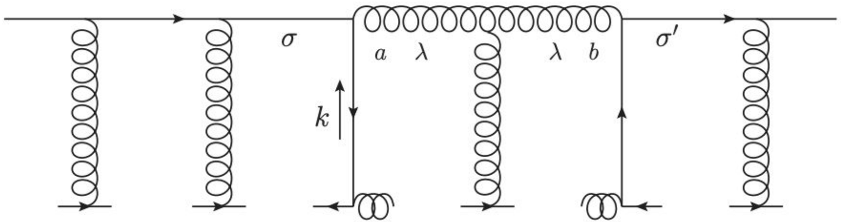
$$Q_q, G_2, \tilde{G}$$

- Quark line undergoes one extra helicity exchange, which is **sub-eikonal**



Polarized Wilson Lines





$\vec{\mu} \cdot \vec{B}$

- Chromo-magnetic field

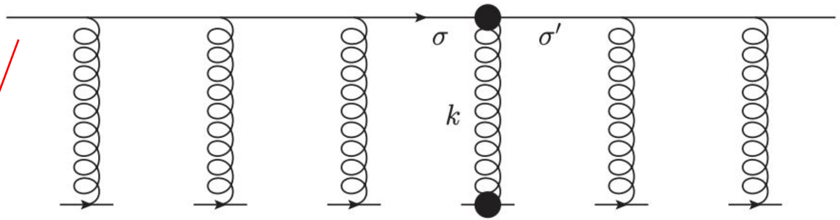
$\mu B_z \sim F_{12}$

$\bar{\psi}\gamma^+\gamma^5\psi$

- Axial Current

Polarized Dipole Amplitudes: "Q_q" , "$\tilde{G}$"

Polarized Wilson Lines



- Quark propagator

$$\int \frac{dk^+}{2\pi} e^{ikx} \frac{\not{k}}{2k^+k^- - k_\perp^2}$$

- Sub-eikonal phase expansion

$$e^{-ix^- \frac{k_\perp^2}{2k^-}} \approx 1 - ix^- \frac{k_\perp^2}{2k^-} \Rightarrow \partial_\perp^2 \rightarrow D_\perp^2 \rightarrow \text{“}G_2\text{”}$$

$$A_\perp$$

- Polarized gluon vertex

Calculating Helicity Distributions

$$\Delta q + \Delta \bar{q} = \frac{1}{N_c} \int_0^{\eta_{max}} d\eta \int_{s_{10}^{min}}^{\eta} ds_{10} \frac{1}{\alpha_s(s_{10})} \left(Q_q(s_{10}, \eta) + 2G_2(s_{10}, \eta) \right)$$

- We incorporate running coupling that runs with size of the dipole

- η ~ Longitudinal momentum
- s_{10} ~ Transverse separation of Dipole

Large Nc&Nf Helicity Evolution

In the large Nc&Nf, Nc/Nf fixed limit, the evolution equations for the polarized dipole amplitudes close:

$$Q_q(s_{10},\eta) = Q_q^{(0)}(s_{10},\eta) + \int_{s_{10}+y_0}^{\eta} d\eta' \int_{s_{10}}^{\eta'-y_0} ds_{21} \Big[Q_q(s_{21},\eta') + 2\tilde{G}(s_{21},\eta') + 2\tilde{\Gamma} s_{10}, s_{21},\eta') \\ -\bar{\Gamma}_f(s_{10},s_{21},\eta') + 2G_2(s_{21},\eta') + 2\Gamma_2(s_{10},s_{21},\eta') \Big] \\ + \frac{1}{2}\int_{y_0}^{\eta} d\eta' \int_{\max\{0,s_{10}+\eta'-\eta\}}^{\eta'-y_0} ds_{21} \Big[Q_q(s_{21},\eta') + 2G_2(s_{21},\eta') \Big]$$

+ 9 more

- 5 Polarized dipole amplitudes mix under evolution: $Q_{u,d,s}, \tilde{G}, G_2$
- With 5 auxiliary dipoles: $\Gamma_{u,d,s}, \tilde{\Gamma}, \Gamma_2$ - which impose lifetime ordering
- Small-*x* cutoff, $y_0 \propto \ln 1/x_0$

Large Nc&Nf Helicity Evolution

- 5 **Polarized dipole amplitudes** mix under evolution: $Q_{u,d,s}, \tilde{G}, G_2$
- With 5 auxiliary dipoles: $\Gamma_{u,d,s}, \tilde{\Gamma}, \Gamma_2$
- For a total of 10 equations that form a **closed system**
- Undetermined initial conditions: $Q^{(0)}_{u,d,s}, \tilde{G}^{(0)}, G^{(0)}_2$

Recap:

$$\frac{1}{2} = S_q + L_q + S_g + L_g$$

$$S_q(Q^2) = \frac{1}{2}\int_0^1 dx \sum_q (\Delta q(x,Q^2) + \Delta\bar{q}(x,Q^2))$$

$$S_g(Q^2) = \int_0^1 dx \Delta G(x,Q^2)$$

$$\Delta q + \Delta\bar{q} = \frac{1}{N_c}\int_0^{\eta_{max}} d\eta \int_{s_{10}^{min}}^{\eta} ds_{10} \frac{1}{\alpha_s(s_{10})}\left(Q_q(s_{10},\eta) + 2G_2(s_{10},\eta)\right)$$

$$\Delta G(x,Q^2) = \frac{2N_c}{\alpha_s(Q^2)} G_2\left(\sqrt{\frac{N_c}{2\pi}}\ln\frac{Q^2}{\Lambda^2}, \sqrt{\frac{N_c}{2\pi}}\ln\frac{Q^2}{x\Lambda^2}\right)$$

Large $N_c \& N_f$ Helicity Evolution

$$Q_q^{(0)}, \tilde{G}^{(0)}, G_2^{(0)}$$

Describing Observables - pDIS

What enters into observables are linear combinations of hPDFs

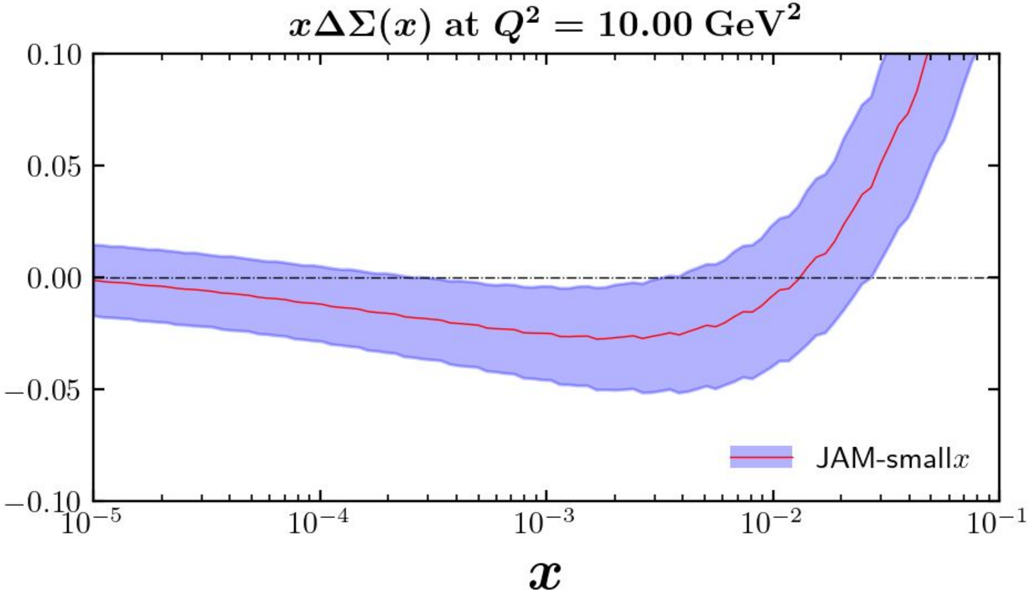
$$\Delta q^+ = \Delta q + \Delta \bar{q}$$
$$\Delta q^- = \Delta q - \Delta \bar{q}$$

- **Three** relevant hPDFs in DIS: $\Delta u^+, \Delta d^+, \Delta s^+$, involving **five** amplitudes
- Data exist for **two** observables that contain these hPDFs in linearly independent combinations: g_1^p and g_1^n

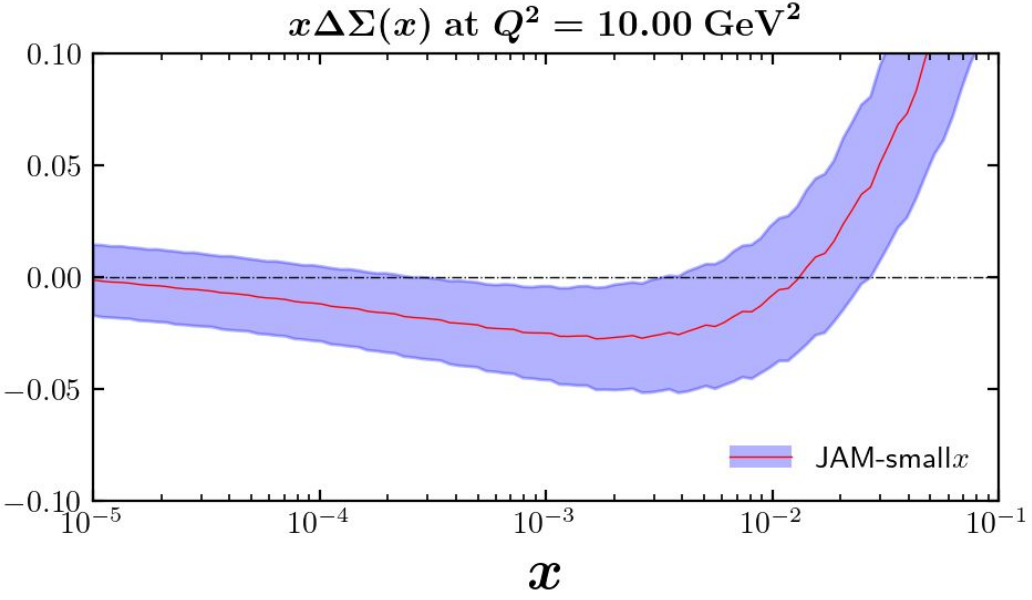
$$g_1^p(x, Q^2) = \frac{1}{2} \sum_q Z_q^2 \Delta q^+(x, Q^2)$$

- Z_q is the quark charge fraction

Contribution from Quark Spin $\Delta\Sigma = \sum_q \Delta q^+(x, Q^2)$



Contribution from Quark Spin $\Delta\Sigma = \sum_q \Delta q^+(x, Q^2)$

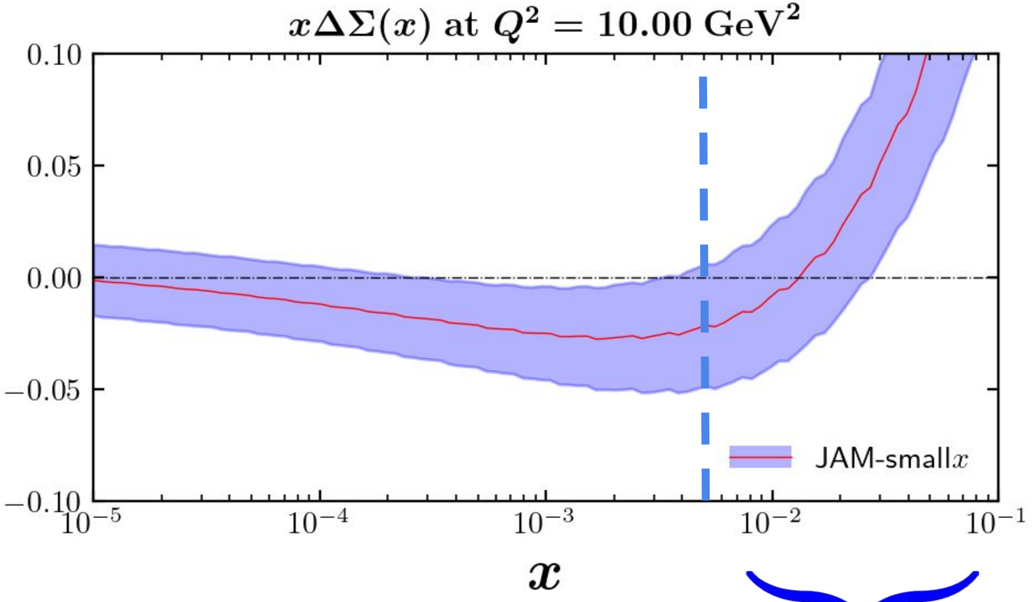


Large-x region

$$\int_{0.01}^{0.7} dx\Delta\Sigma(x) = \pm 0.36$$

Contribution from Quark Spin

$$\Delta\Sigma = \sum_q \Delta q^+(x, Q^2)$$



Large-*x* region

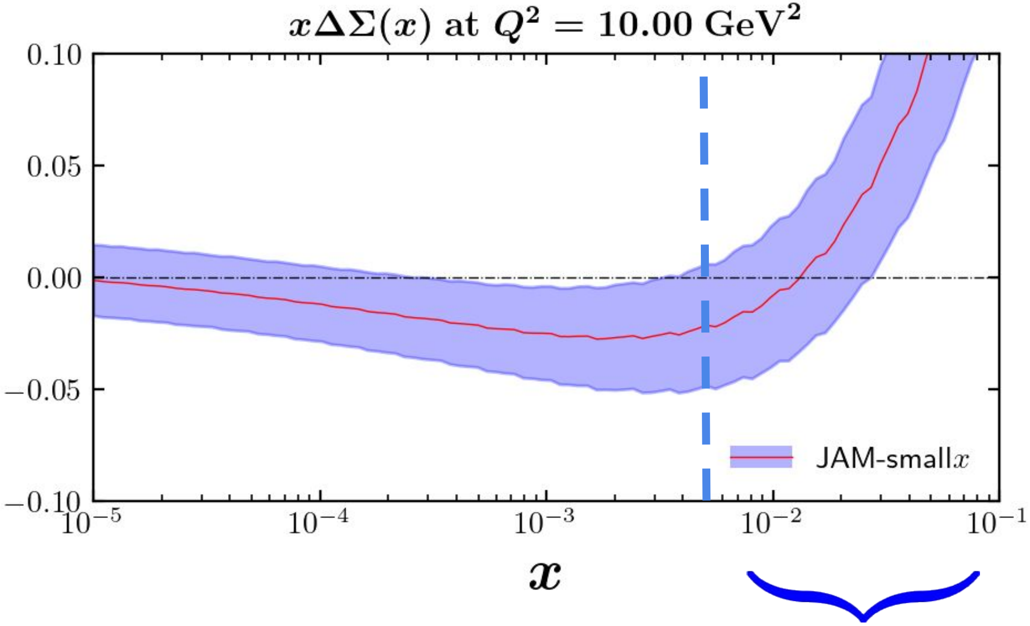
$$\int_{0.01}^{0.7} dx \Delta\Sigma(x) = \pm 0.36$$

Existing small-*x* data

Contribution from Quark Spin

$$\Delta\Sigma = \sum_q \Delta q^+(x, Q^2)$$

Saturation region



Existing small-*x* data

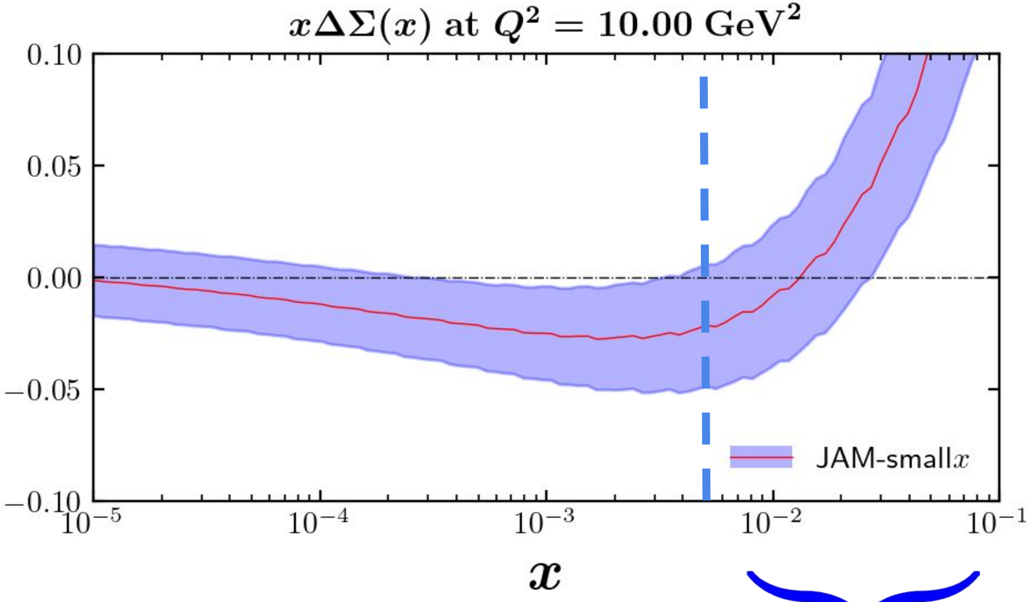
Large-*x* region

$$\int_{0.01}^{0.7} dx \Delta\Sigma(x) = \pm 0.36$$

Contribution from Quark Spin

$$\Delta\Sigma = \sum_q \Delta q^+(x, Q^2)$$

Saturation region



Existing small-x data

Large-x region

$$\int_{0.01}^{0.7} dx \Delta\Sigma(x) = \pm 0.36$$

Compare with:

$$\int_{10^{-5}}^{10^{-3}} dx \Delta\Sigma(x) = -0.1 \pm 0.1$$

Describing Observables - pSIDIS

- 2 observables are not enough to describe 3 hPDFs.
- Expand our horizons to Semi-Inclusive DIS - all hPDFs are relevant here, both singlet, Δq^+ and non-singlet, Δq^-
- **Non-singlet distributions obey their own small-*x* evolution that has been solved**

$$\Delta q^- = \frac{N_c}{2\pi^3} \int d\eta \int ds_{10} Q_q^{NS}(s_{10}, \eta)$$

- Q_q^{NS} is the non-singlet Polarized Dipole Amplitude - obeys its own evolution equation
- pSIDIS grants us access to the semi-inclusive, spin dependent structure functions g_1^h

g_1^h Structure Functions

$$g_1^h(x,z,Q^2) = \frac{1}{2}\sum_q Z_q^2 \Delta q(x,z,Q^2) D_q^h(z,Q^2)$$

- D_q^h are fragmentation functions - giving the probability quark *q* fragments into hadron *h*
- z Is the fraction of the virtual photons momentum carried by the hadron
- The flavour hPDF is obtained via $\Delta q = \frac{1}{2}(\Delta q^+ + \Delta q^-)$
- In pSIDIS, we are able to scatter on 2 targets (proton, neutron), tag 2 outgoing hadrons (pion, kaon) that each have 2 charges - 2x2x2=8 new observables

Constraining the rest of the Polarized Dipole Amplitudes

$$g_1^{p,n} \sim Q_u, Q_d, Q_s, G_2$$

- 2 observables, 4 polarized dipole amplitudes. Under constrained system

$$g_1^h \sim Q_q, G_2, Q_q^{NS}$$

- 8 new observables, 3 new polarized dipole amplitudes. Exactly constrained - but $\tilde{G}$ does not enter directly into observables

$$pp \to jets \sim G_2, \tilde{G}$$

- Particle production might provide final constraints

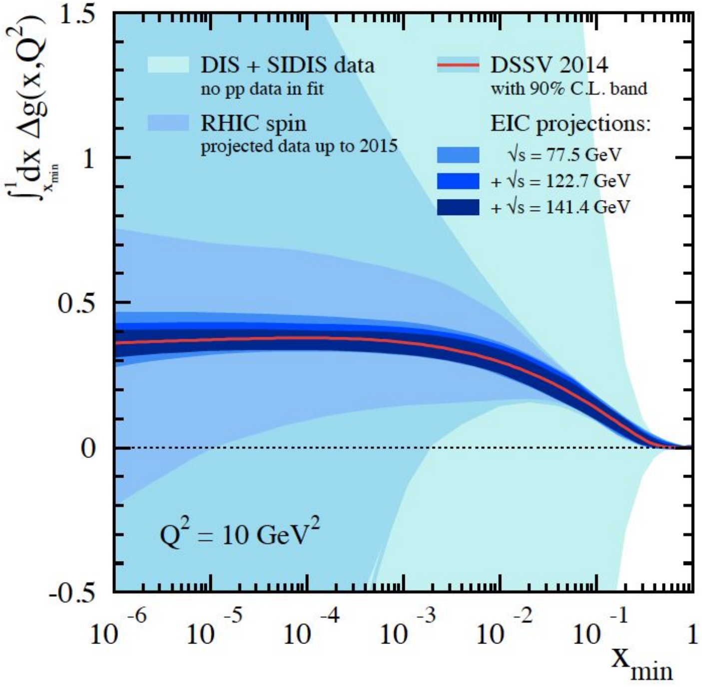
Conclusions

- In order to resolve the spin puzzle, the small-*x* behaviour of the hPDFs need to be understood
- This is accomplished using small-*x* evolution
- Along with fitting to data
- Potentially a significant amount if spin is hiding in the small-*x* region
- More work needs to be done to constrain small-*x* behavior of the various polarized dipoles - especially G_2 and $\tilde{G}$
- Could be constrained by studying particle production in *pp* collisions

Backup Slides

Gluon Helicity Parton Distributions Function

gluon
spin



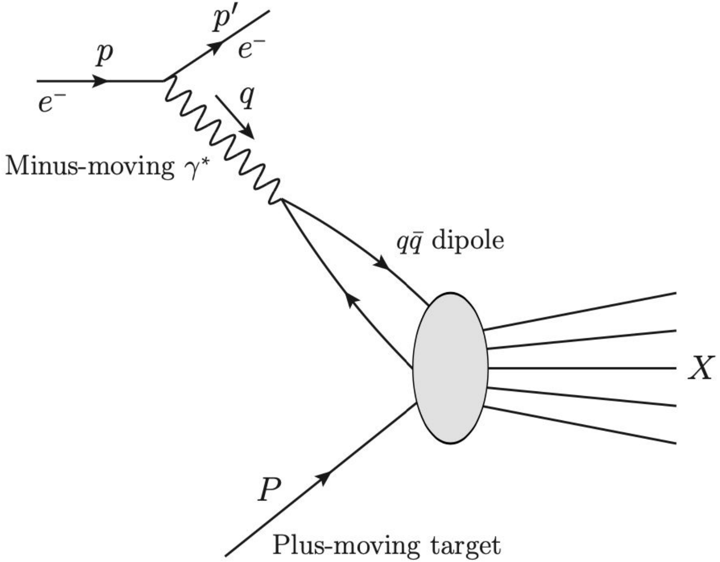
$$S_g(Q^2) = \int_0^1 dx \Delta G(x, Q^2)$$

ΔG = Gluon Helicity PDF

- Uncertainty consistently blows up when extrapolating beyond data

Deep-Inelastic Scattering (DIS)

Probing the proton at small *x*



- Electron of momentum *p* scatters off proton of momentum *P*
- Transverse size given by virtuality of photon:

$$\frac{1}{x_\perp^2} \propto Q^2 = -q^2$$

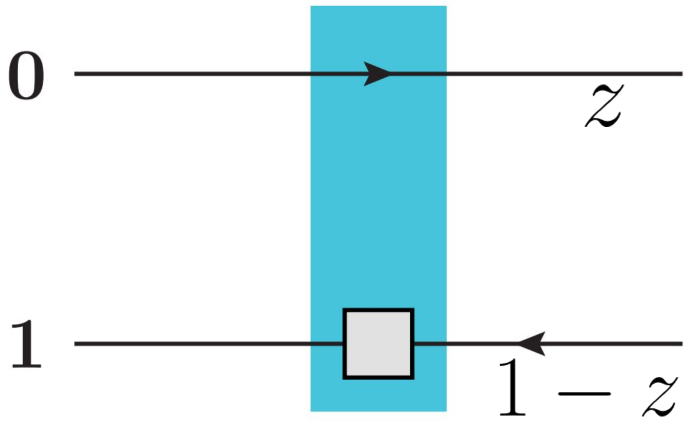
- Bjorken-*x*: $x = \frac{Q^2}{2P \cdot q} \approx \frac{Q^2}{s}$

Calculating Helicity Distributions

$$\Delta G(x, Q^2) = \frac{2N_c}{\alpha_s(Q^2)} G_2\left(\sqrt{\frac{N_c}{2\pi}} \ln \frac{Q^2}{\Lambda^2}, \sqrt{\frac{N_c}{2\pi}} \ln \frac{Q^2}{x\Lambda^2}\right)$$

- Jaffe-Manohar Gluon Helicity Distribution
- Λ^2 Infrared cutoff

Polarized Dipole Amplitude - Degrees of Freedom



$$Q_q(s_{10}, \eta)$$

Polarized Dipole Amplitudes are functions of

- Transverse separation:

$$x_{10}^2 = (\underline{x}_1 - \underline{x}_0)^2$$

- Momentum Fraction times center of mass energy: zs
- Rescaled variables:

$$\eta = \sqrt{\frac{N_c}{2\pi}} \ln \frac{zs}{\Lambda^2} \qquad s_{10} = \sqrt{\frac{N_c}{2\pi}} \ln \frac{1}{x_{10}^2 \Lambda^2}$$

Helicity Evolution

Using Light-Cone Operator Treatment, we need to resum all gluon exchanges that exchange helicity information

$\Rightarrow$ Resumming all terms containing:

$$\alpha_s \int_x^1 \frac{dz}{z} \int_{1/s}^{1/Q^2} \frac{d^2 x_{21}}{x_{21}^2}$$

Longitudinal part. Present in un-polarized evolution

Transverse part. UV exactly cancelled in un-polarized evolution

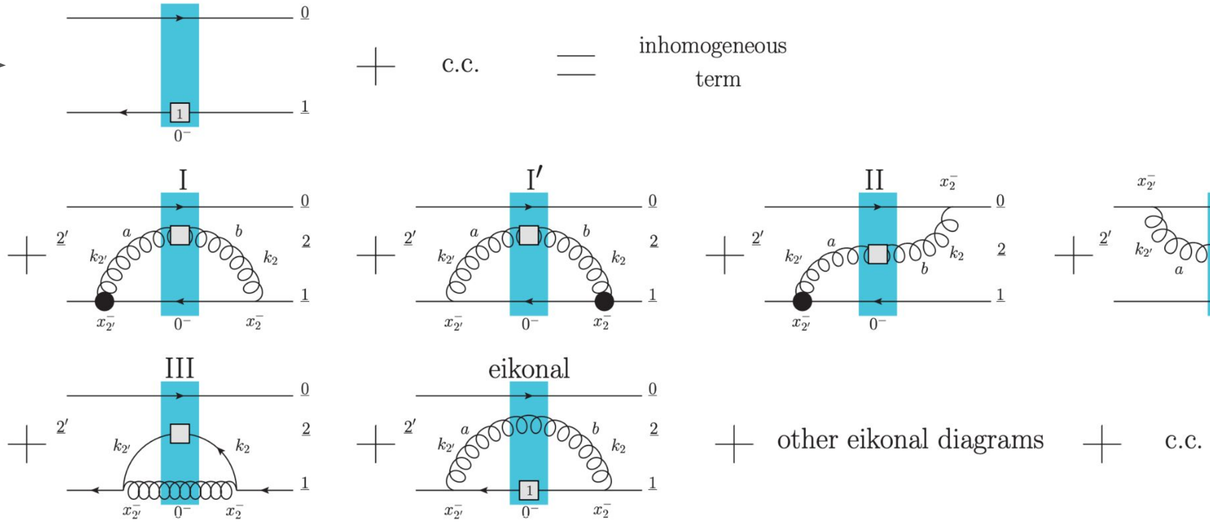
Resum double log (DLA) terms:

$$\alpha_s \ln^2(1/x)$$

Helicity Evolution

- Relate Polarized Dipole Amplitude to themselves at higher energies by resumming emission diagrams - resumming Double Log (DLA) contributions: $\alpha_s \ln^2(1/x)$

$Q_q(x_{10}^2, zs) \longrightarrow$



Sub-eikonal Expansion

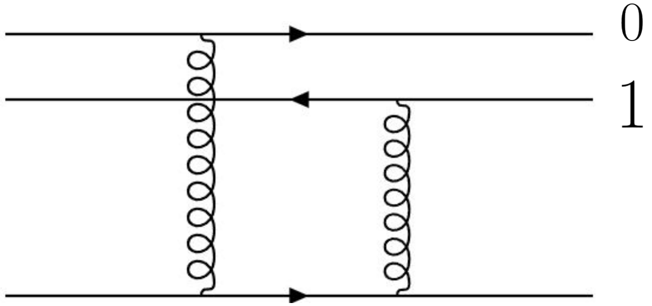
- Expansion in energy or in *x*

Eikonal	Sub-Eikonal	Sub-Sub-Eikonal
$1/x,$	$x^0,$	x^1
F_1, F_2	$g_1^{p,n}, \Delta q, \Delta \bar{q}$	Transversity

- No eikonal terms contain any helicity information - Wilson lines are helicity independent
- Must calculate sub-eikonal terms to access helicity

Inhomogeneous term

The inhomogeneous term is given by a Born-inspired ansatz:



$$\propto \int_0^s \frac{dk_\perp^2}{k_\perp^2}(1 - e^{-\underline{k}\cdot\underline{x}_{10}}) = \pi \ln(s x_{10}^2)$$

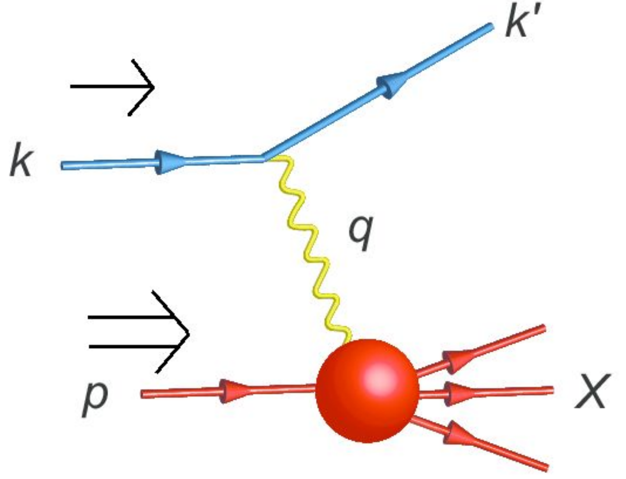
$$\propto \eta - s_{10}$$

$$\Gamma_q^{(0)} = Q_q^{(0)} = a\eta + b s_{10} + c$$

- Same form of the other Dipole Amplitudes
- Parameters a,b,c need to be extracted from data

Observables - Double Spin Asymmetries in DIS

$$A_{||} = \frac{\sigma^{\uparrow\Downarrow} - \sigma^{\uparrow\Uparrow}}{\sigma^{\uparrow\Downarrow} + \sigma^{\uparrow\Uparrow}} \propto A_1 \propto g_1^{p,n}$$

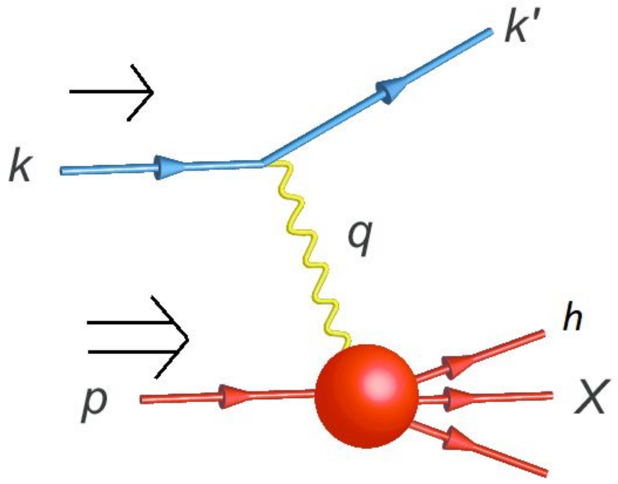


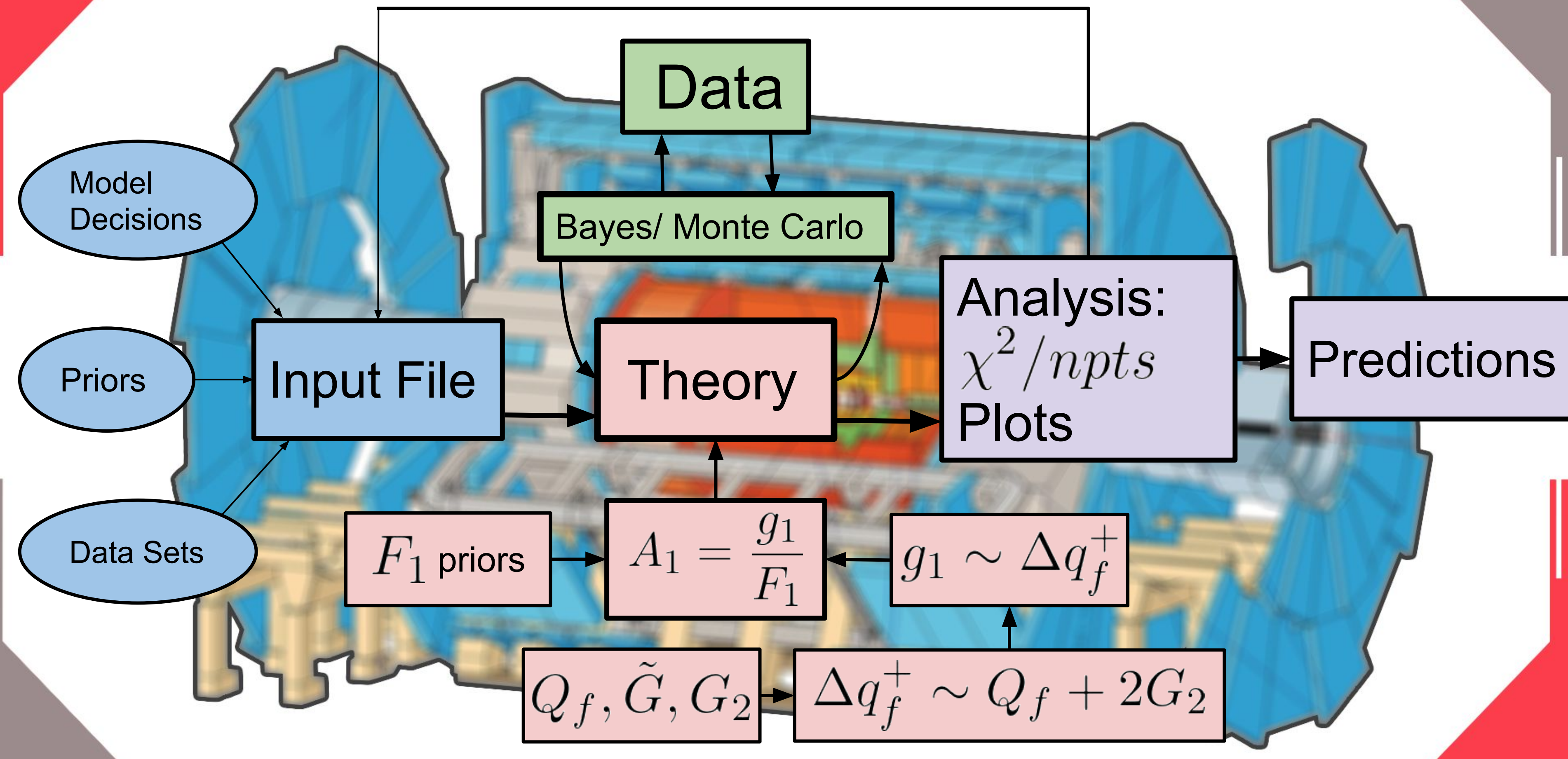
- $\uparrow (\downarrow)$ is positive (negative) helicity electron
- $\Uparrow (\Downarrow)$ is positive (negative) helicity proton
- A_1 is virtual photoproduction asymmetry

Observables - Double Spin Asymmetries in SIDIS

$$A_{||}(z) = \frac{\sigma^{\uparrow\Downarrow} - \sigma^{\uparrow\Uparrow}}{\sigma^{\uparrow\Downarrow} + \sigma^{\uparrow\Uparrow}} \propto g_1^h(z)$$

- *h* is the tagged hadron
- *z* is the momentum fraction of the virtual photon carried by the tagged hadron

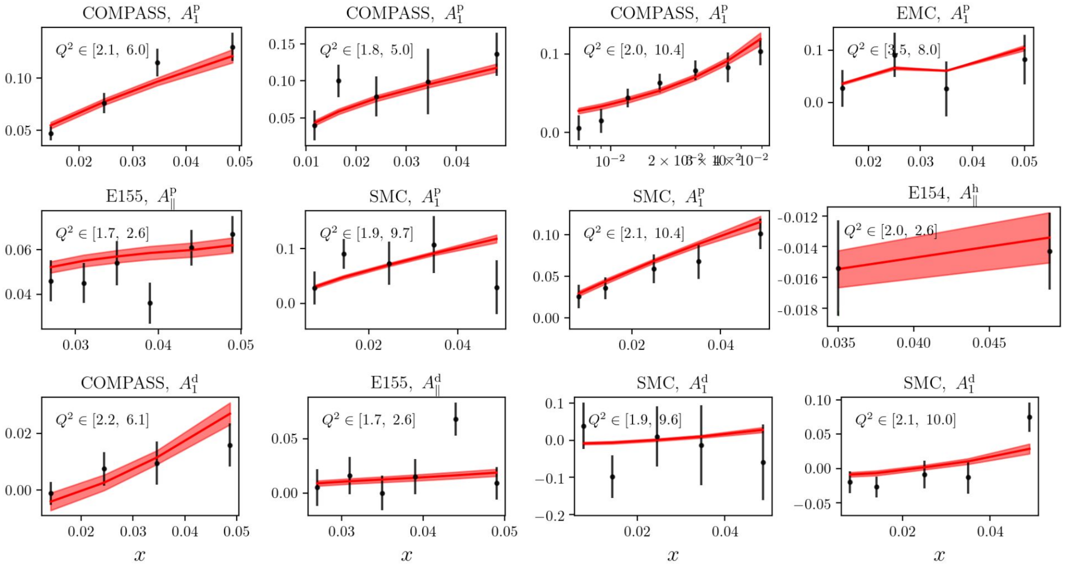


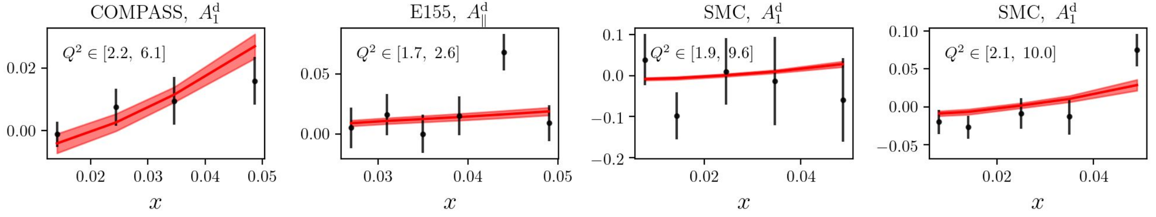

Data
Bayes/ Monte Carlo
Model Decisions
Priors
Data Sets
Input File
Theory
Analysis: $\chi^2/npts$ Plots
Predictions
F_1 priors
$A_1 = \frac{g_1}{F_1}$
$g_1 \sim \Delta q_f^+$
$Q_f, \tilde{G}, G_2$
$\Delta q_f^+ \sim Q_f + 2G_2$

χ^2 and Data Cuts

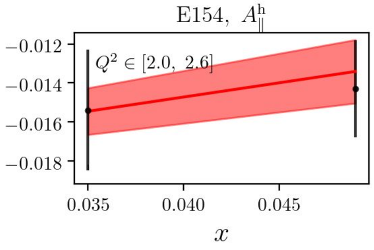
- First simultaneous fit of small-*x* theory to polarized DIS & SIDIS data
- Cut of 0.005< x < 0.1
- Cut of $1.0 GeV^2 < Q^2 < 10.4 GeV^2$
- Cut of 0.2 < z <1.0
- Describing 234 data points
- With a $\chi^2/npts$ = 1.01

Global fit of DIS - Data vs Theory

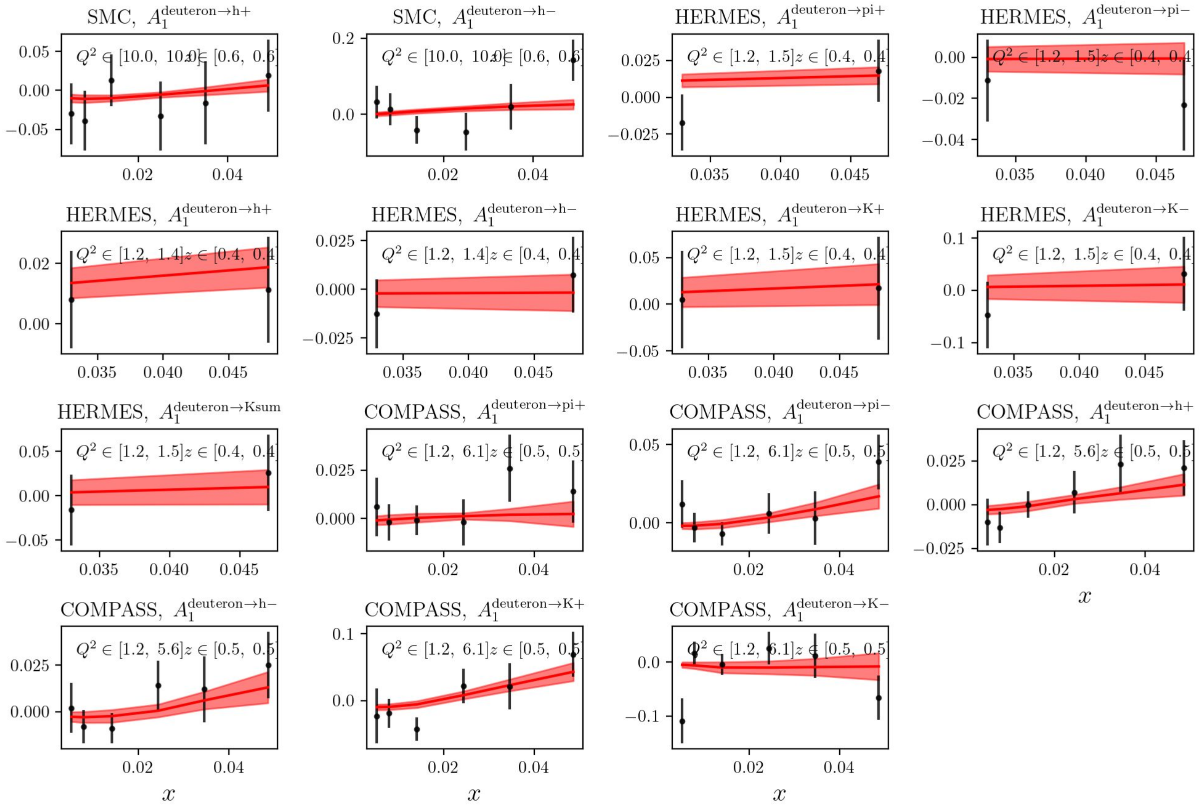




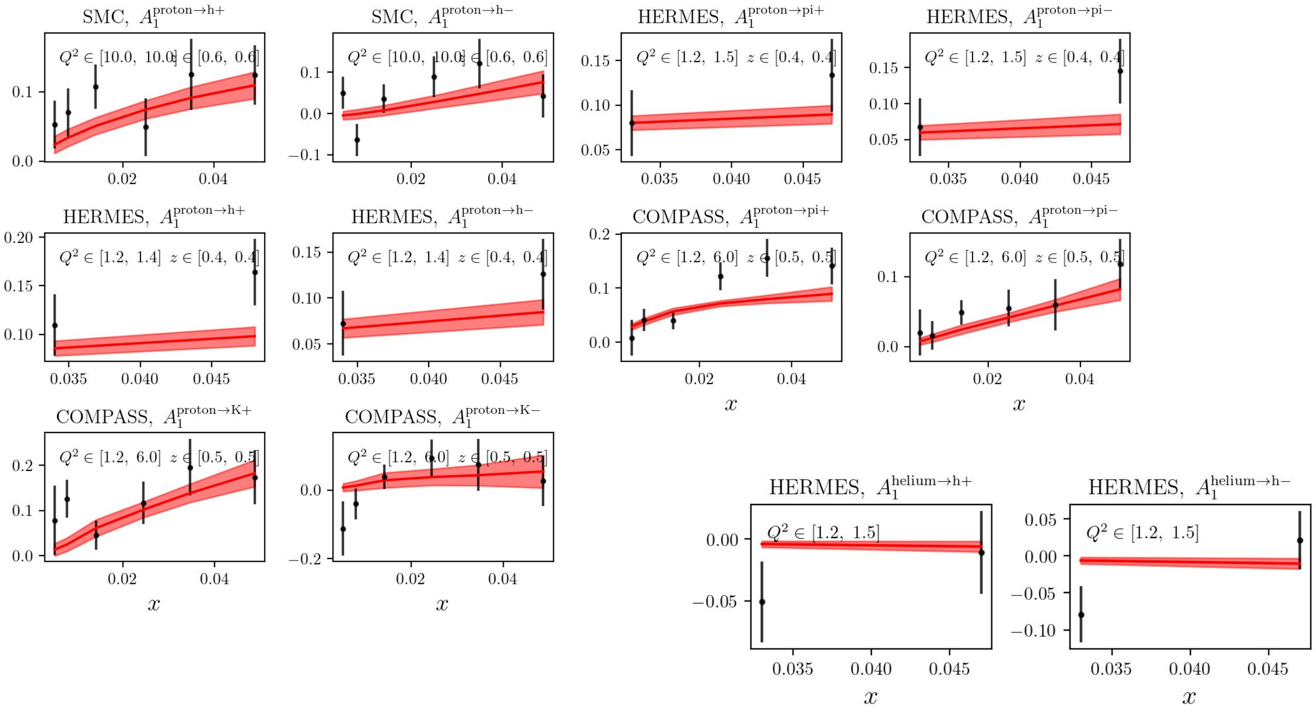
- Red curves - our theory
- Black dots - data
  - COMPASS
  - EMC
  - SMC
  - SLAC
  - HERMES
- Preliminary results



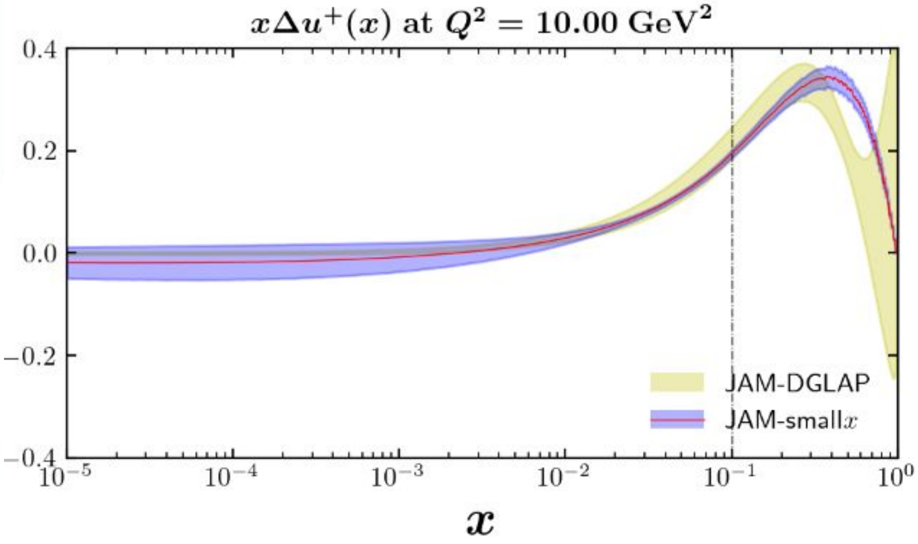
Global fit of SIDIS - Data vs Theory

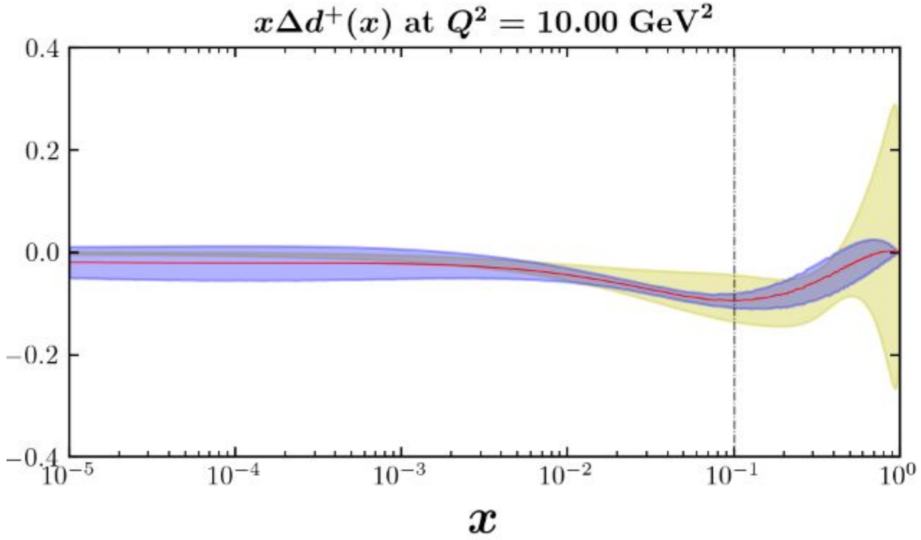


Fitting SIDIS - Data vs Theory



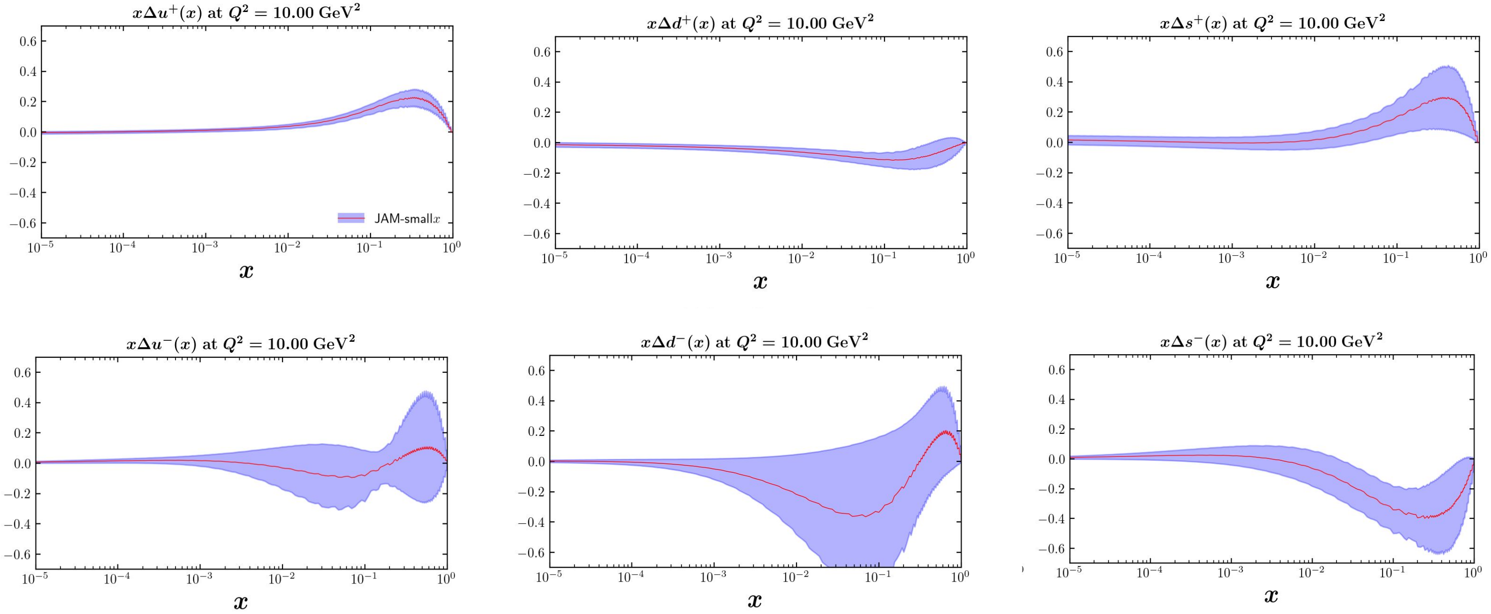
hPDFs - (Preliminary)





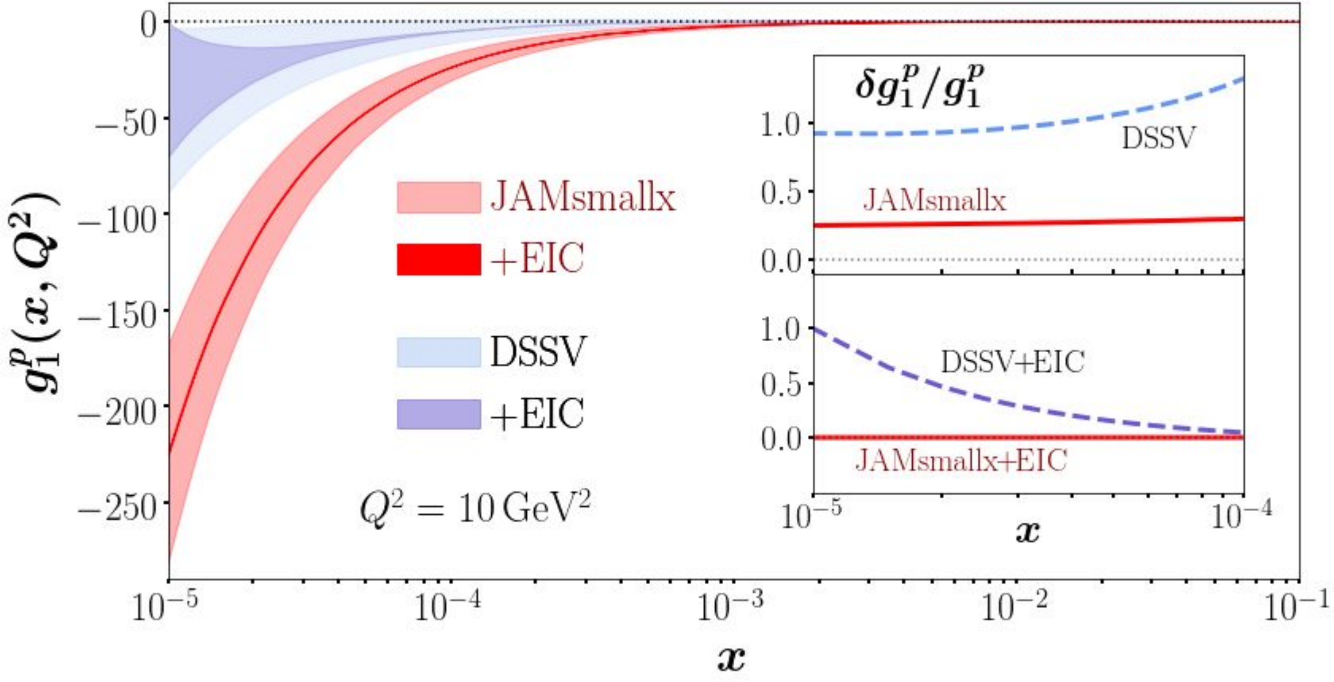
- DIS only: Strange distribution set to zero

hPDFs - Preliminary



- Old version of evolution

(Preliminary) Extraction of g_1^p



- DSSV uses DGLAP - rational function extrapolation of *x*
- We use small-*x* helicity evolution to predict the *x* behaviour
- Leads to control over uncertainty