

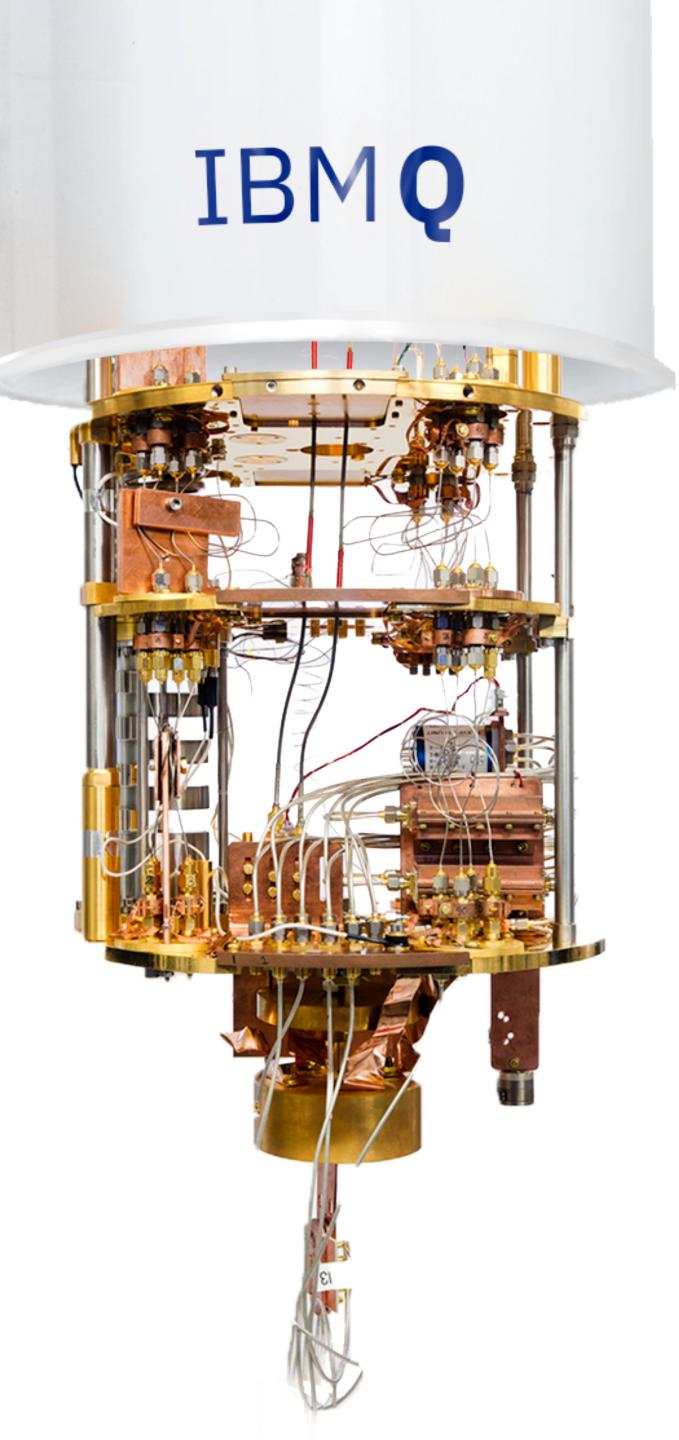
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Collider Events on a Quantum Computer

Simon Williams

Resummation, Evolution, Factorization 2022



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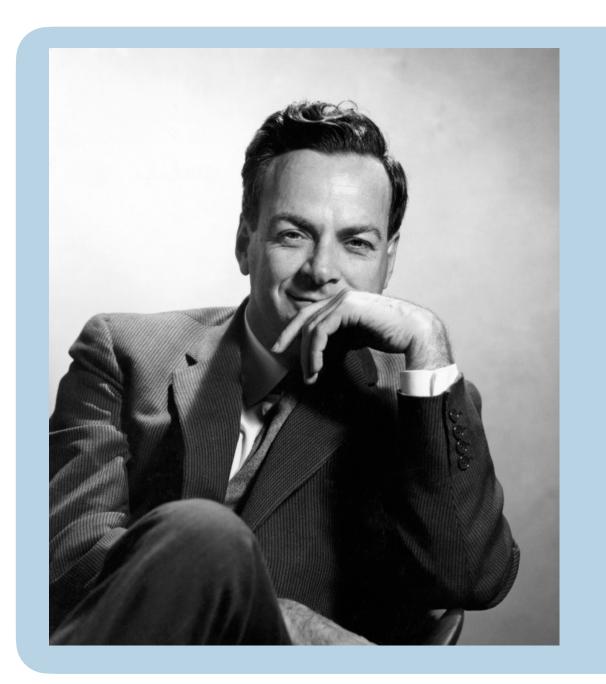
- Quantum Computing The Power of the Qubit
- Why are we interested in High Energy Physics?
- The Parton Shower
 - Discretising QCD
- Collider Events on a Quantum Computer

G. Gustafson, S. Prestel, M. Spannowsky and S. Williams, Collider Events on a Quantum Computer, <u>arXiv:2207.10694</u>





Quantum Computing - The Power of the Qubit!

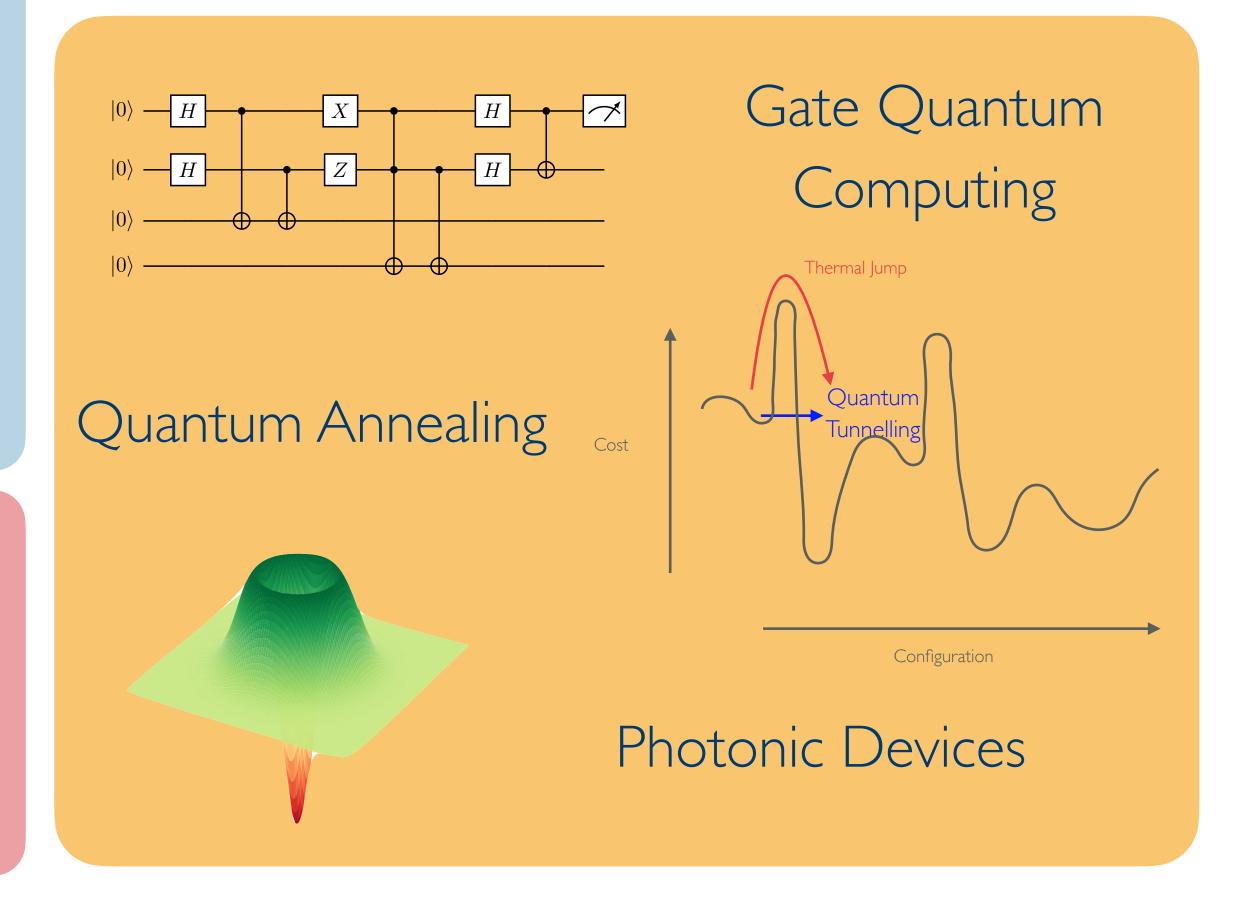


"Nature is quantum [...] so if you want to simulate it, you need a quantum computer"

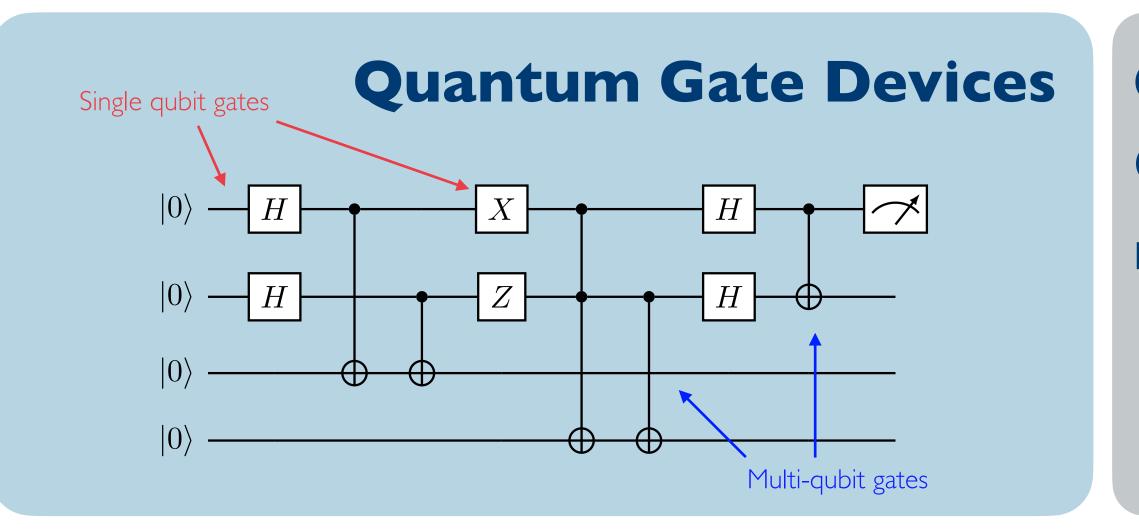
- Richard Feynman (1982)

Quantum Computing has had a lot of successes since - most recently with Shor and Deutsch winning the Breakthrough Prize and the 2022 Nobel Prize going to Quantum Information

Types of Quantum Device:



Types of Quantum Computing Devices



Qubit model:

Qubit is the quantum analogue of the classical bit, not restricted only to $|0\rangle$ and $|1\rangle$

$$|\psi\rangle = \cos\frac{\theta}{2}|0\rangle + e^{i\phi}\sin\frac{\theta}{2}|1\rangle$$

Advantages:

- Highly controllable qubits
- Universal computation

Disadvantages:

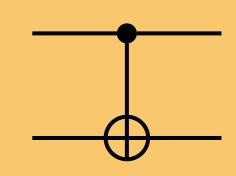
- Small number of qubits, not very fault tolerant

Single qubit gates:

$$-U_3$$

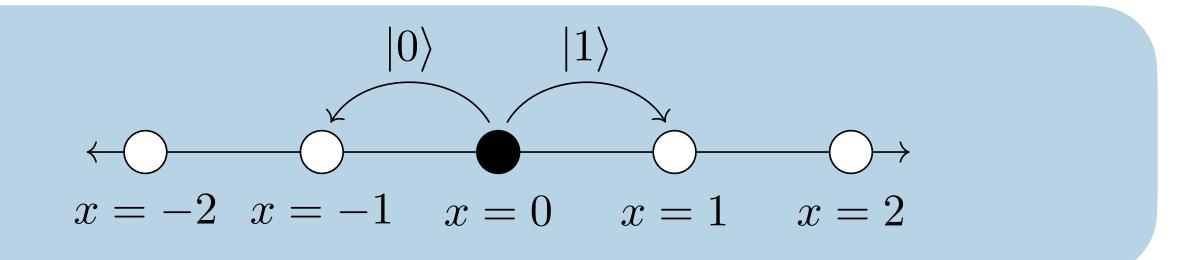
$$U_3|0\rangle \rightarrow \cos\frac{\theta}{2}|0\rangle + e^{i\phi}\sin\frac{\theta}{2}|1\rangle$$

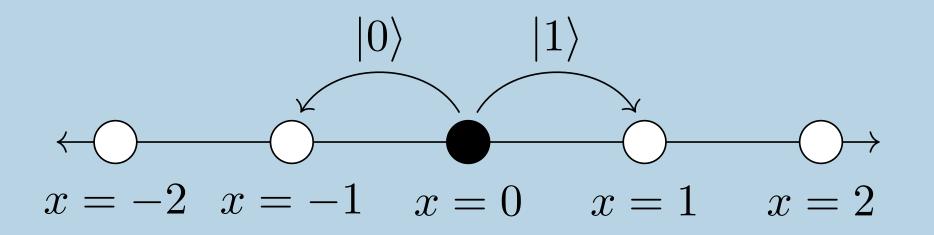
Multi-qubit gates:



$$CNOT |00\rangle \rightarrow |00\rangle, CNOT |10\rangle \rightarrow |11\rangle,$$

 $CNOT | 01 \rangle \rightarrow | 01 \rangle$, $CNOT | 11 \rangle \rightarrow | 10 \rangle$

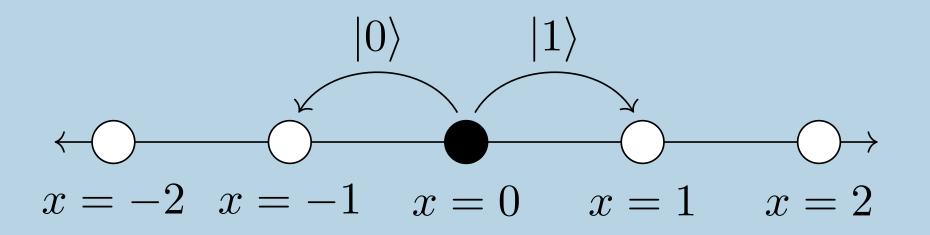




$$\mathcal{H}_{P} = \{ | i \rangle : i \in \mathbb{Z} \}$$

$$\mathcal{H}_{C} = \{ | 0 \rangle, | 1 \rangle \}$$

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$$\mathcal{H}_{P} = \{ | i \rangle : i \in \mathbb{Z} \}$$

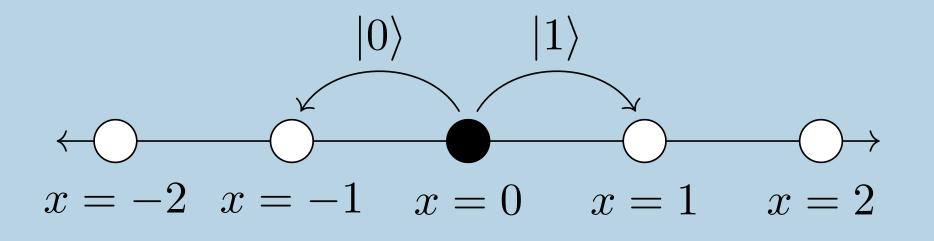
$$\mathcal{H}_{C} = \{ | 0 \rangle, | 1 \rangle \}$$

$$\mathcal{H}_{C} = \mathcal{H}_{C} \otimes \mathcal{H}_{P}$$

Unitary

Transformation:

$$U = S \cdot (C \otimes I)$$



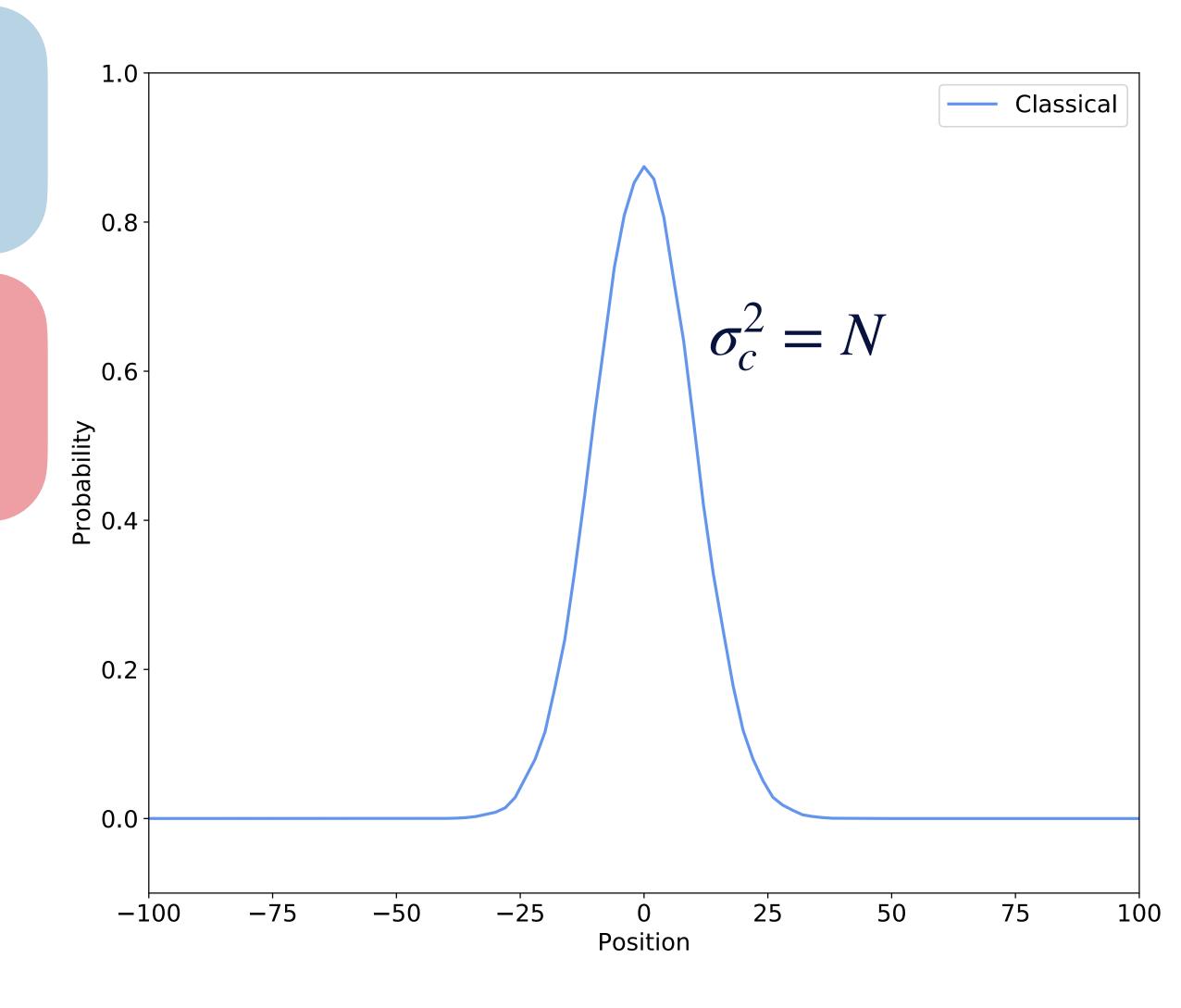
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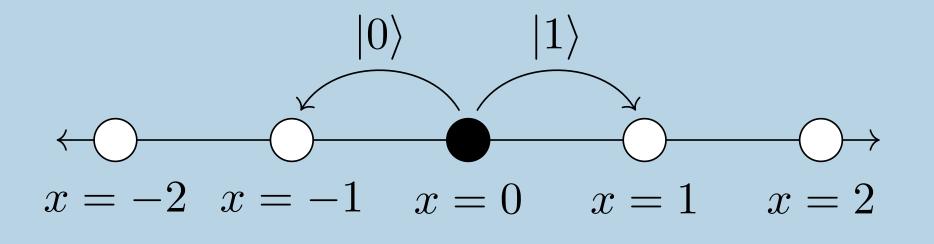
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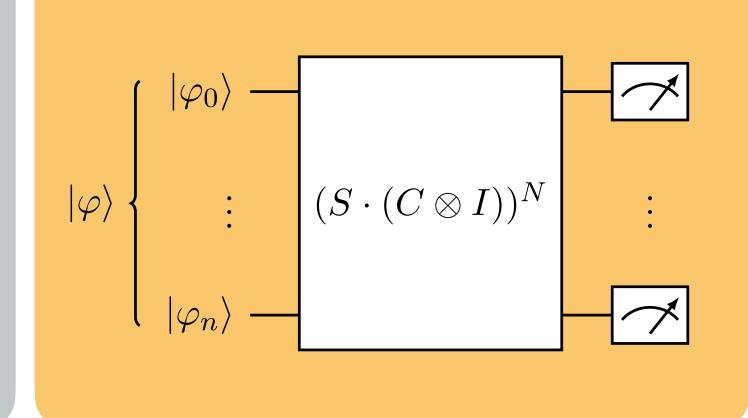
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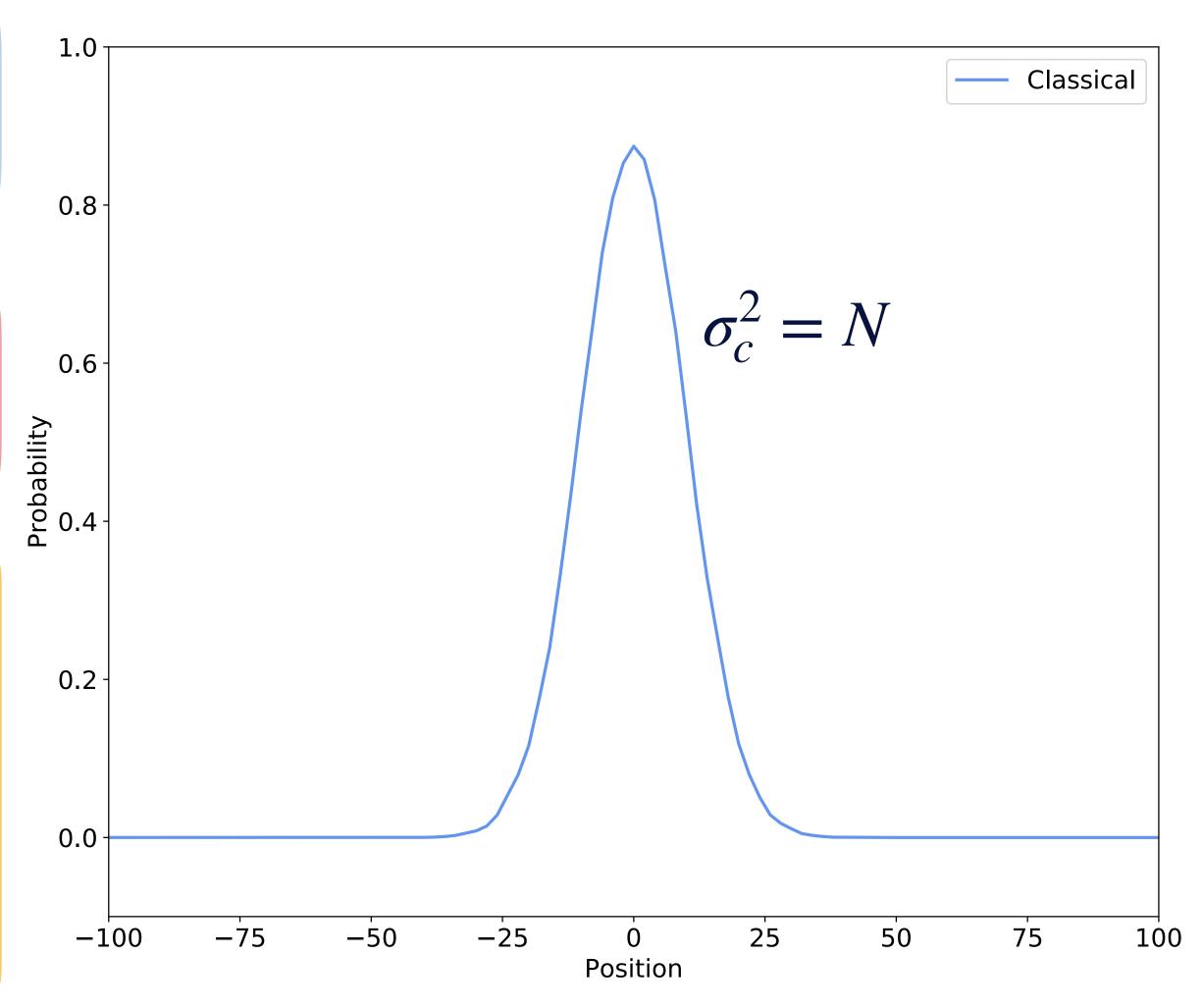
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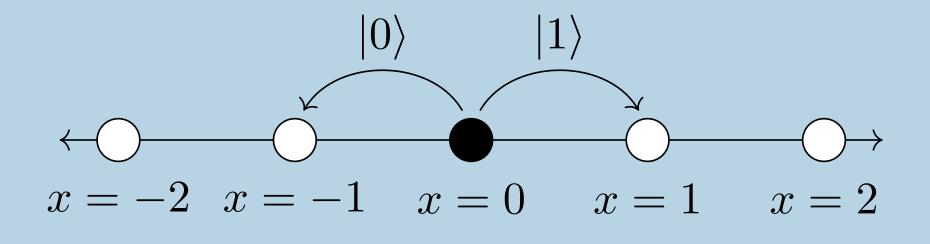
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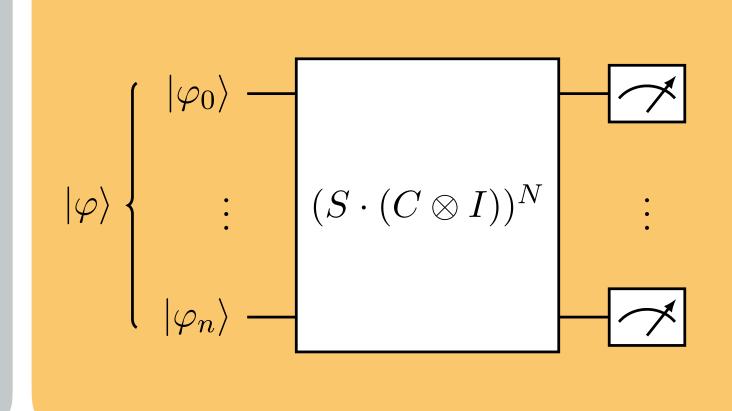
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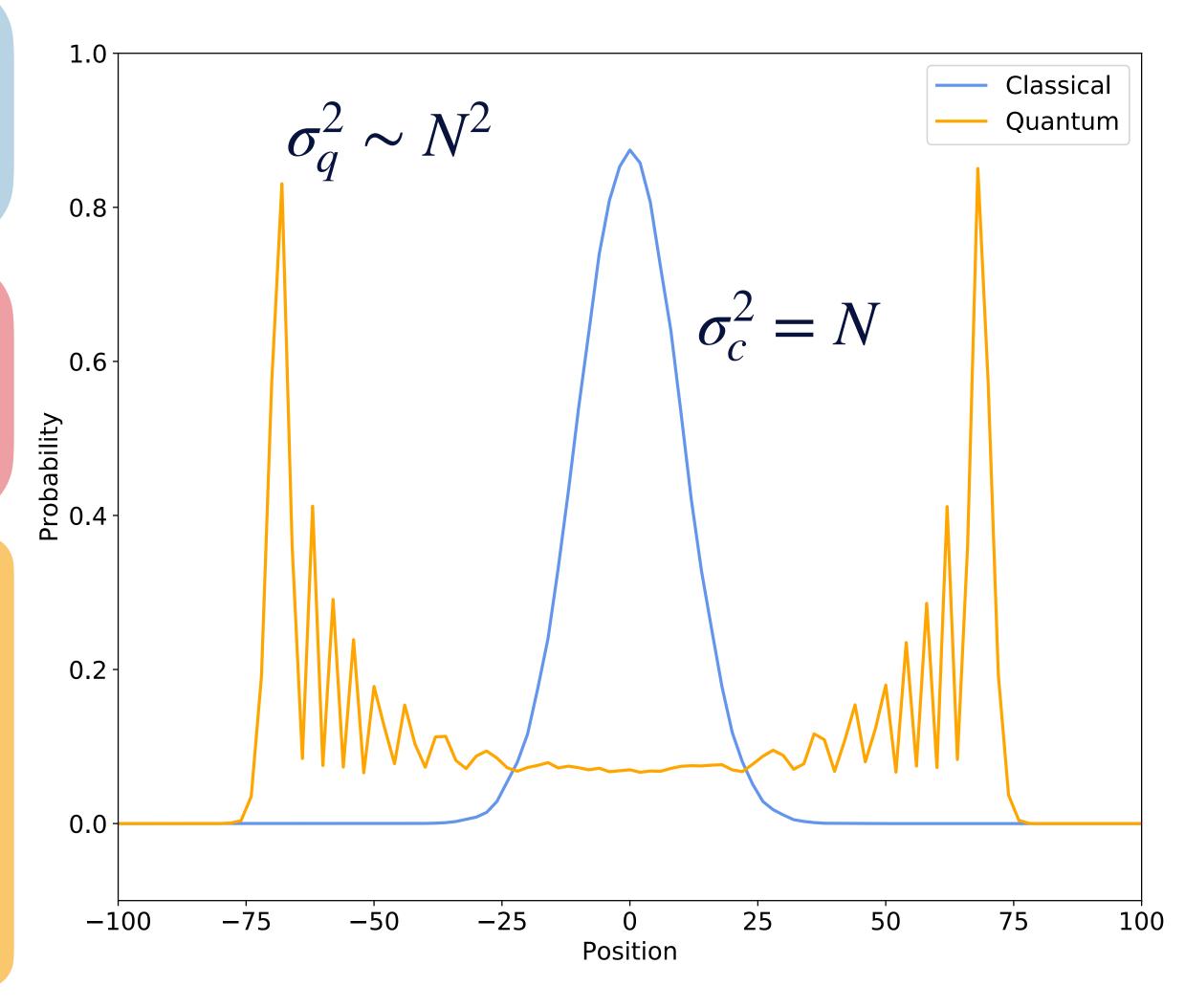
$$\mathcal{H}_{C} = \{ | 0 \rangle, | 1 \rangle \}$$

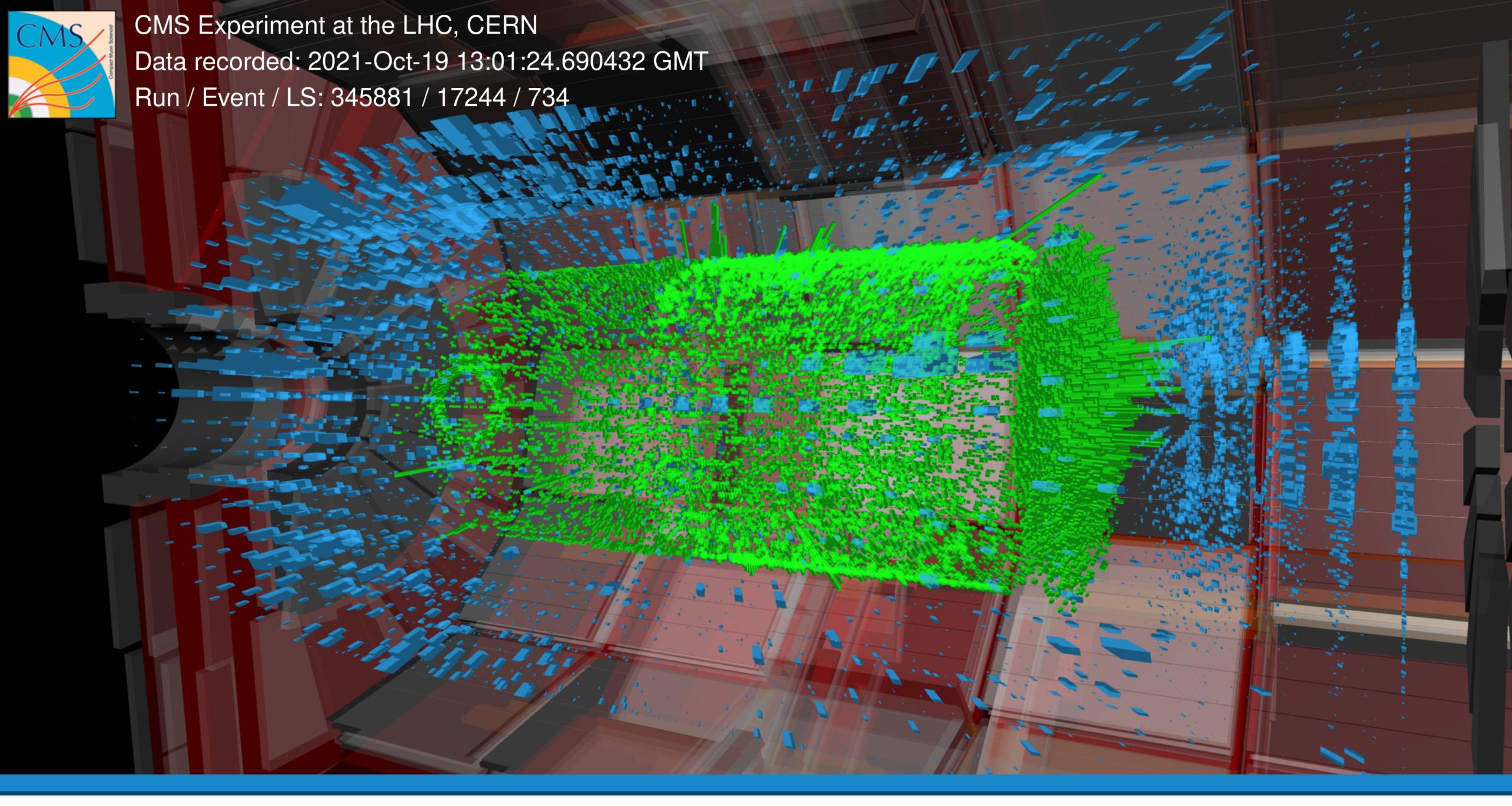
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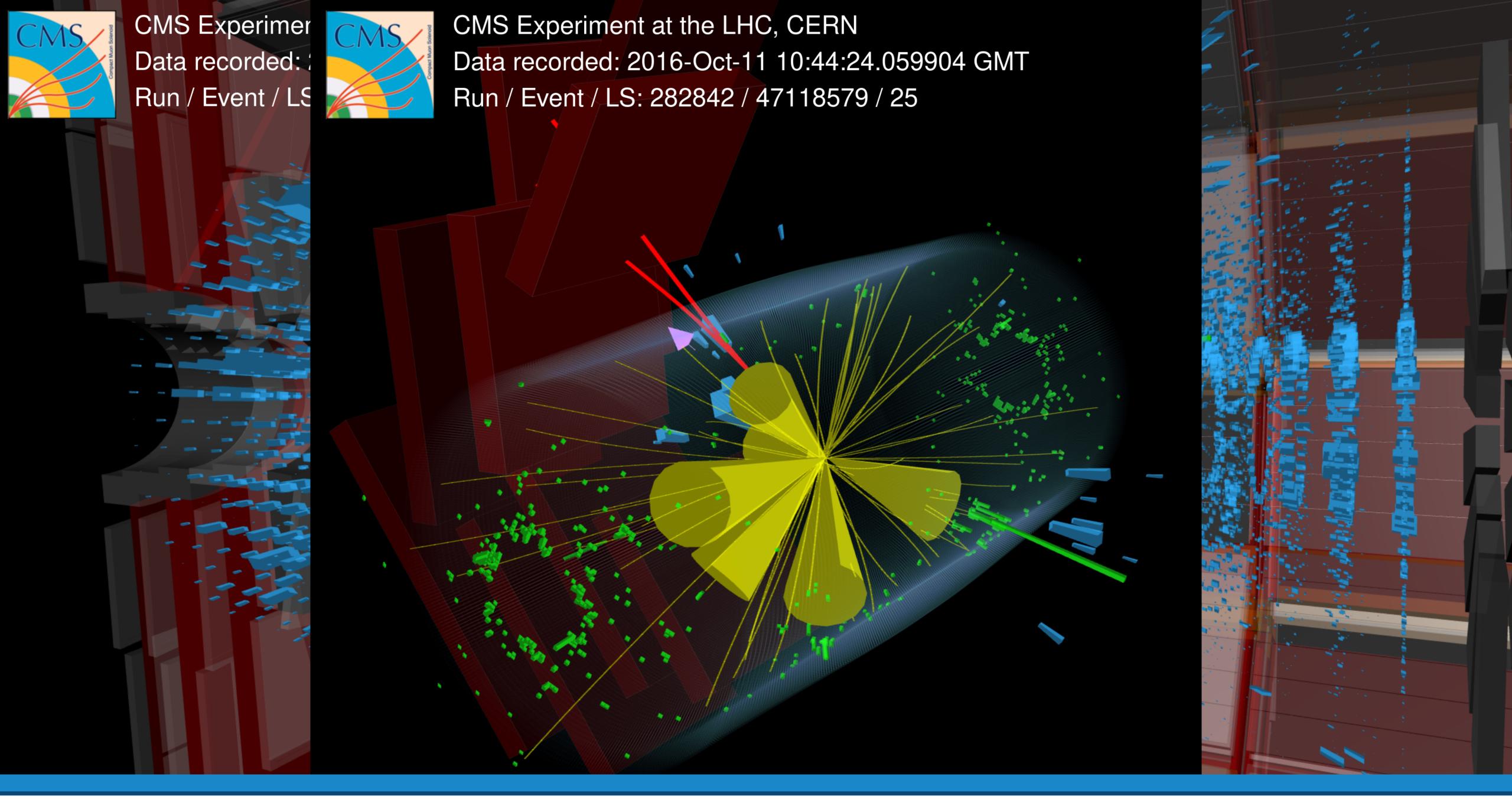
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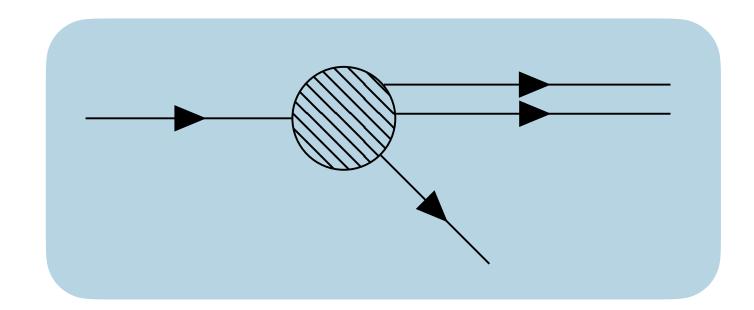


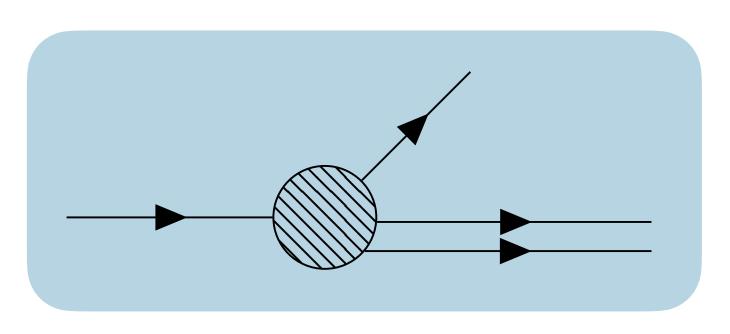






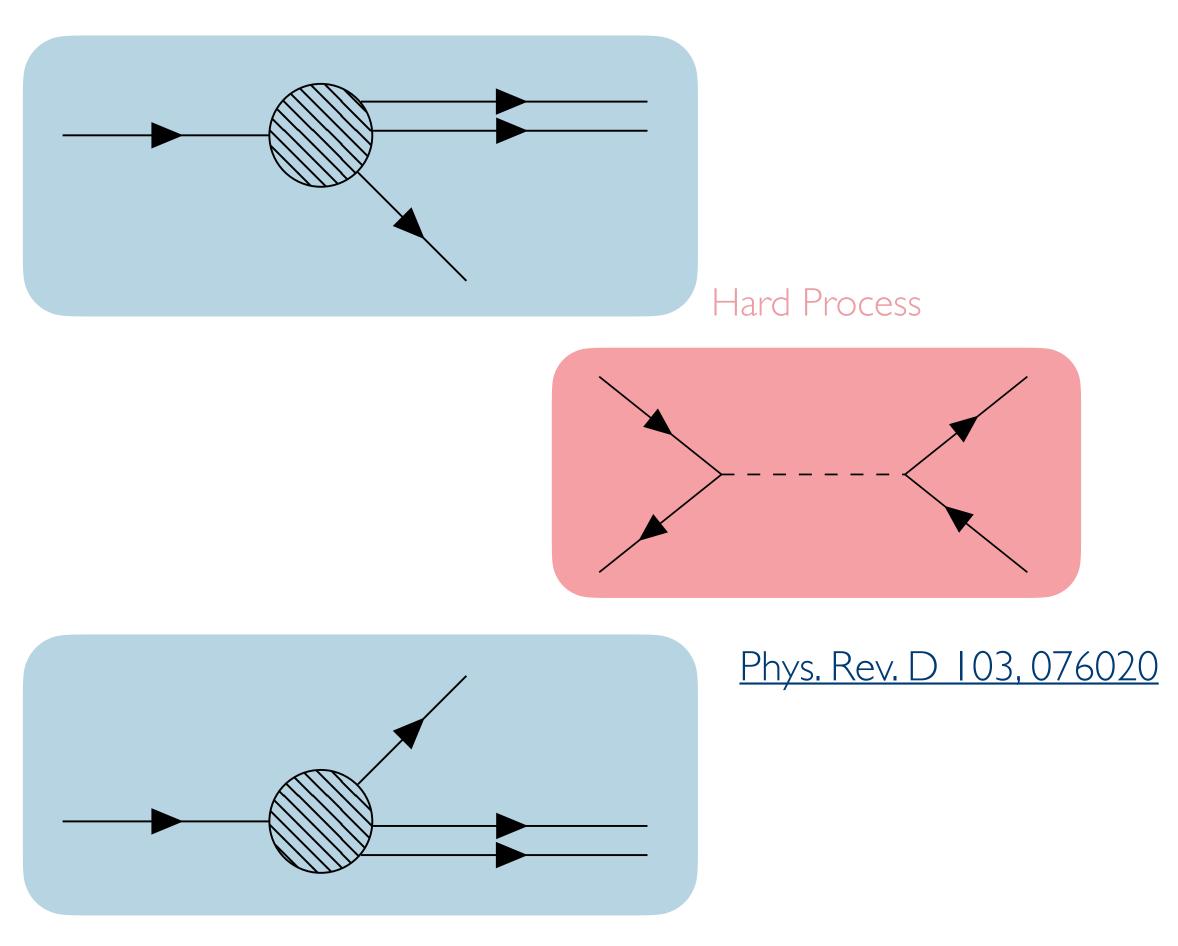
Parton Density Functions





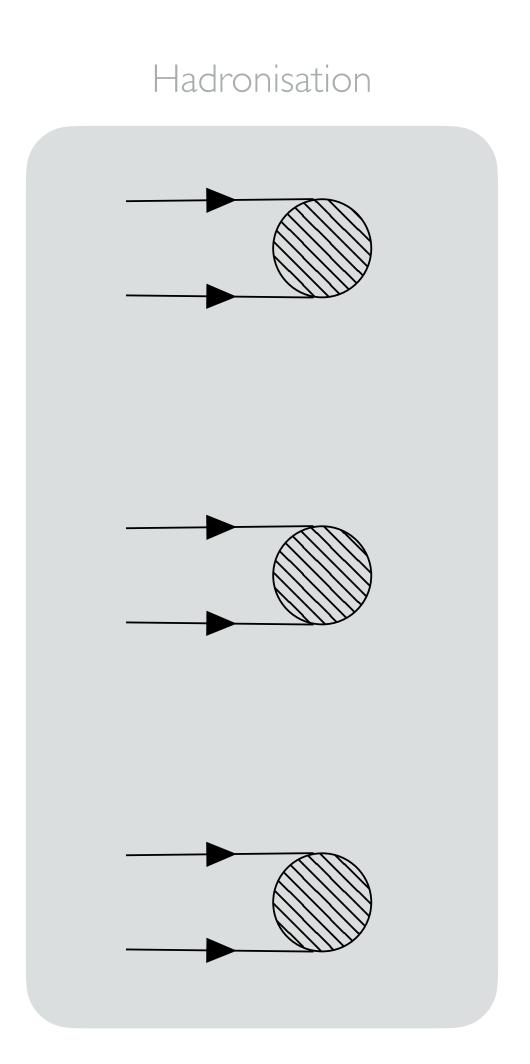
Phys. Rev. D 103, 034027

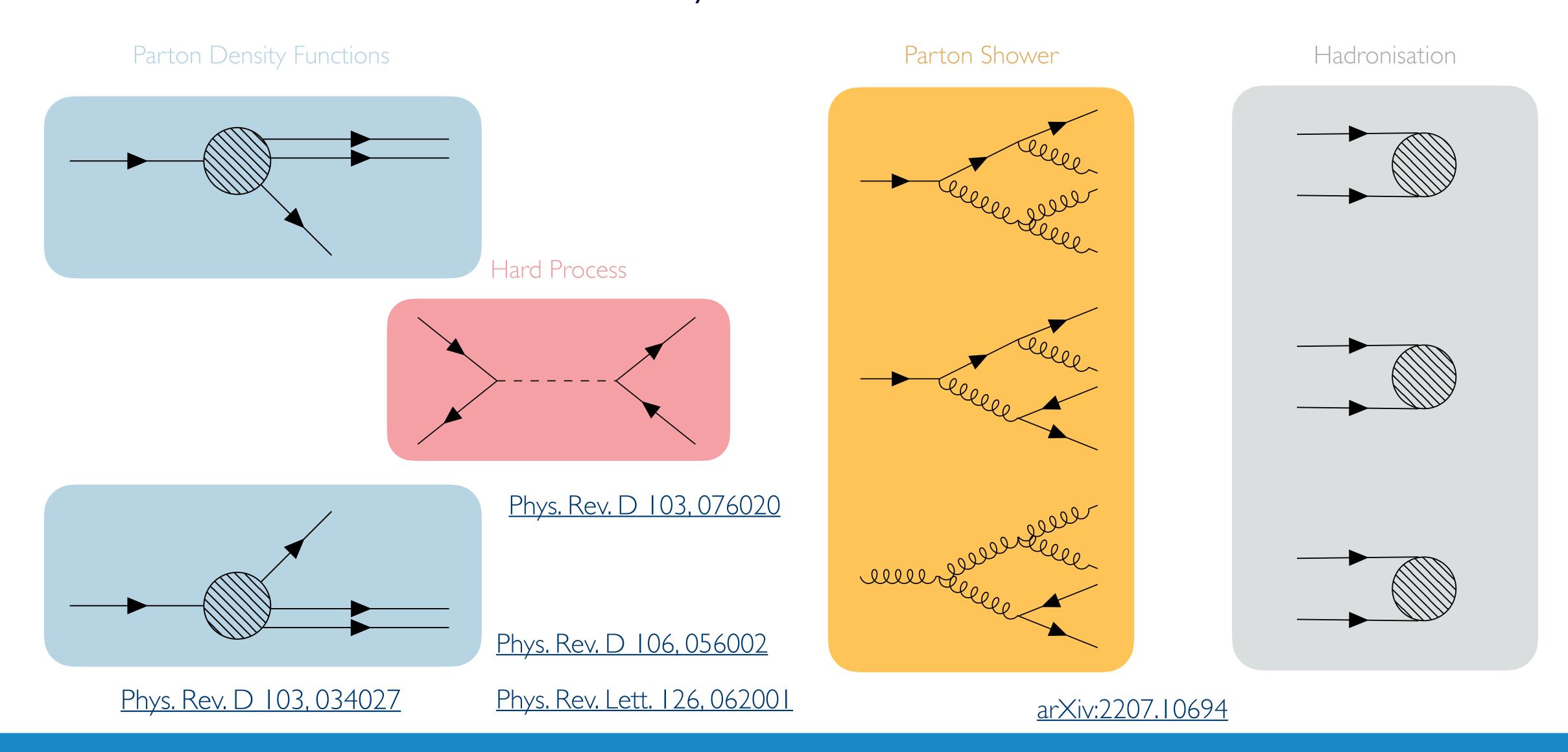
Parton Density Functions



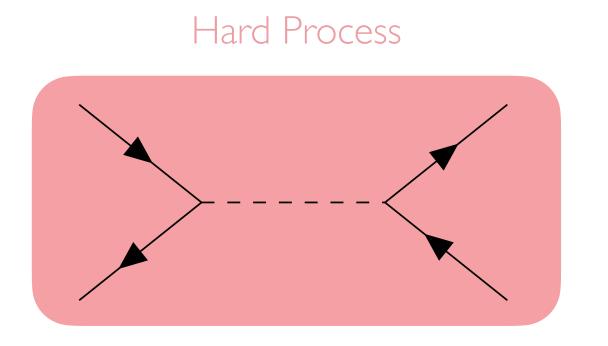
Phys. Rev. D 103, 034027

Parton Density Functions Hard Process Phys. Rev. D 103, 076020 Phys. Rev. D 103, 034027





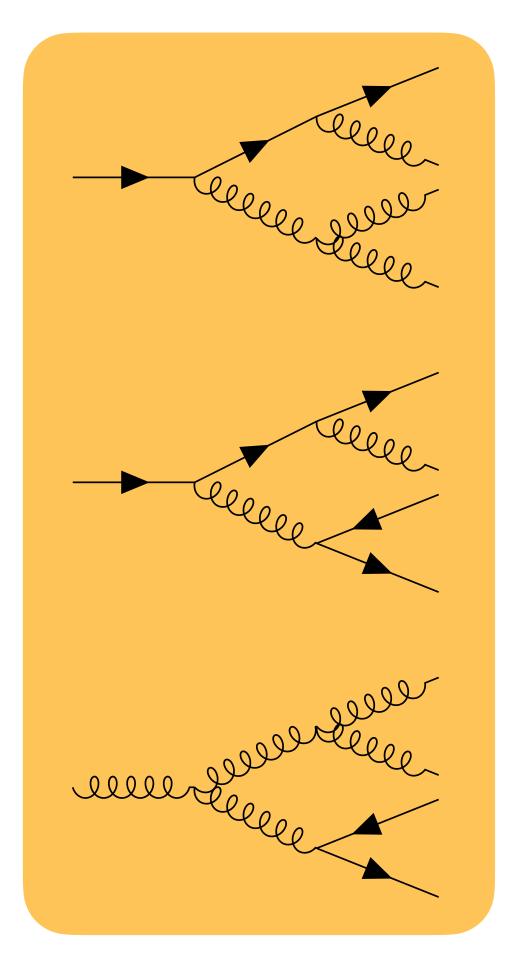
Parton Shower



Phys. Rev. D 103, 076020

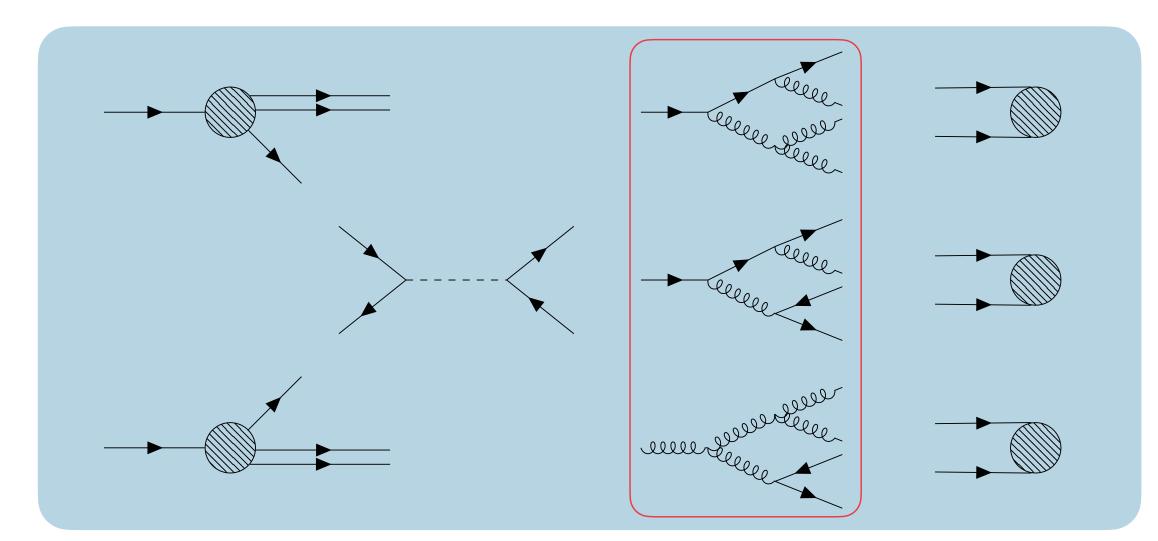
Phys. Rev. D 106, 056002

Phys. Rev. Lett. 126, 062001



arXiv:2207.10694

The Parton Shower

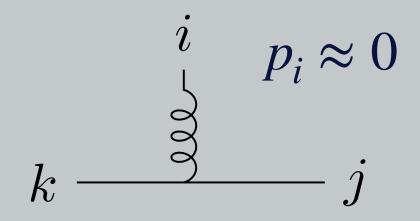


Collinear mode:

$$k \xrightarrow{\overrightarrow{P}} \underbrace{\qquad \qquad }_{j} \qquad p_{i} = zP, \quad p_{j} = (1-z)P$$

Successive decay steps factorise into independent quasi-classical steps

Soft mode:



Interference effects only allow for partial factorisation

Leading contributions to the decay rate in the collinear limit are included in the soft limit

In this limit, the decay from high energy to low energy proceeds as a colour-dipole cascade.

This interpretation allows for straightforward interference patterns and momentum conservation

The Parton Shower - The Veto Algorithm

The choice of the variables ξ and t is known as the phase space parameterisation

Non-Emission Probability

$$\Delta(t_n, t) = \exp\left(-\int_t^{t_n} dt d\xi \frac{d\phi}{2\pi} C \frac{\alpha_s}{2\pi} \frac{2s_{ik}(t, \xi)}{s_{ij}(t, \xi)s_{jk}(t, \xi)}\right)$$

$$\mathcal{F}_{n}(\Phi_{n}, t_{n}, t_{c}; O) = \Delta(t_{n}, t_{c}) O(\Phi_{n})$$

$$+ \int_{t_{c}}^{t_{n}} dt d\xi \frac{d\phi}{2\pi} C \frac{\alpha_{s}}{2\pi} \frac{2s_{ik}(t, \xi)}{s_{jk}(t, \xi)} \Delta(t_{n}, t) \mathcal{F}_{n}(\Phi_{n+1}, t, t_{c}; O)$$

Inclusive Decay Probability

$$d\mathcal{P}\left(q(p_{\mathrm{I}})\bar{q}(p_{\mathrm{K}}) \to q(p_{i})g(p_{j})\bar{q}(p_{k})\right) \simeq \frac{ds_{ij}}{s_{\mathrm{IK}}} \frac{ds_{jk}}{s_{\mathrm{IK}}} C \frac{\alpha_{s}}{2\pi} \frac{2s_{\mathrm{IK}}}{s_{ij}s_{jk}}$$

Current interpretations of the veto algorithm treat the phase space variables ξ and t as **continuous**

Collider Events on a Quantum Computer

Gösta Gustafson, a Stefan Prestel, a Michael Spannowsky, b Simon Williams c

ABSTRACT: High-quality simulated data is crucial for particle physics discoveries. Therefore, Parton shower algorithms are a major building block of the data synthesis in event generator programs. However, the core algorithms used to generate parton showers have barely changed since the 1980s. With quantum computers' rapid and continuous development, dedicated algorithms are required to exploit the potential that quantum computers provide to address problems in high-energy physics. This paper presents a novel approach to synthesising parton showers using the Discrete QCD method. The algorithm benefits from an elegant quantum walk implementation which can be embedded into the classical toolchain. We use the <code>ibm_algiers</code> device to sample parton shower configurations and generate data that we compare against measurements taken at the ALEPH, DELPHI and OPAL experiments. This is the first time a Noisy Intermediate-Scale Quantum (NISQ) device has been used to simulate realistic high-energy particle collision events.

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^bInstitute for Particle Physics Phenomenology, Department of Physics, Durham University, Durham DH1 3LE, U.K.

^cHigh Energy Physics Group, Blackett Laboratory, Imperial College, Prince Consort Road, London, SW7 2AZ, United Kingdom

I. Parameterise phase space in terms of gluon transverse momentum and rapidity:

$$k_{\perp}^2 = rac{s_{ij}s_{jk}}{s_{\mathrm{IK}}}$$
 and $y = rac{1}{2}\ln\left(rac{s_{ij}}{s_{jk}}
ight)$

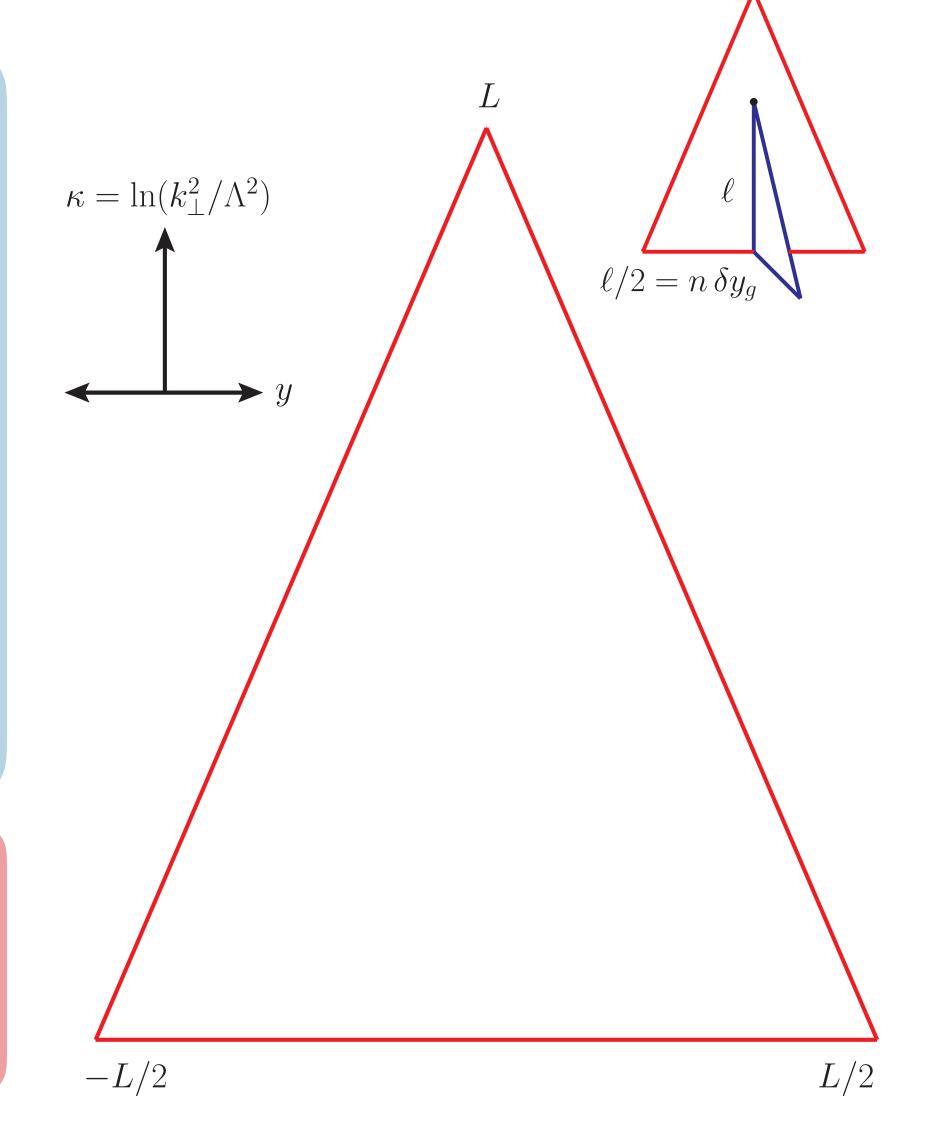
which leads to the inclusive probability:

$$d\mathcal{P}\left(q(p_{\mathrm{I}})\bar{q}(p_{\mathrm{K}})\to q(p_{i})g(p_{j})\bar{q}(p_{k})\right)\simeq =\frac{C\alpha_{s}}{\pi}d\kappa dy$$

where $\kappa = \ln\left(\frac{k_{\perp}^2}{\Lambda^2}\right)$ and Λ is an arbitrary mass scale

Due to the colour charge of emitted gluons, the rapidity span for subsequent dipole decays is increased. This is interpreted as

"folding out"



2. Neglect $g \to q \overline{q}$ splittings and examine transverse-momentum-dependent running coupling

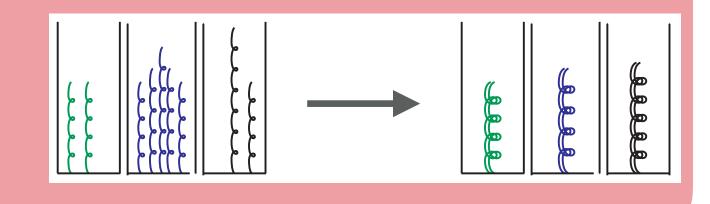
$$\alpha_s(k_{\perp}^2) = \frac{12\pi}{33 - 2n_f} \frac{1}{\ln(k_{\perp}^2/\Lambda_{\text{QCD}}^2)}$$

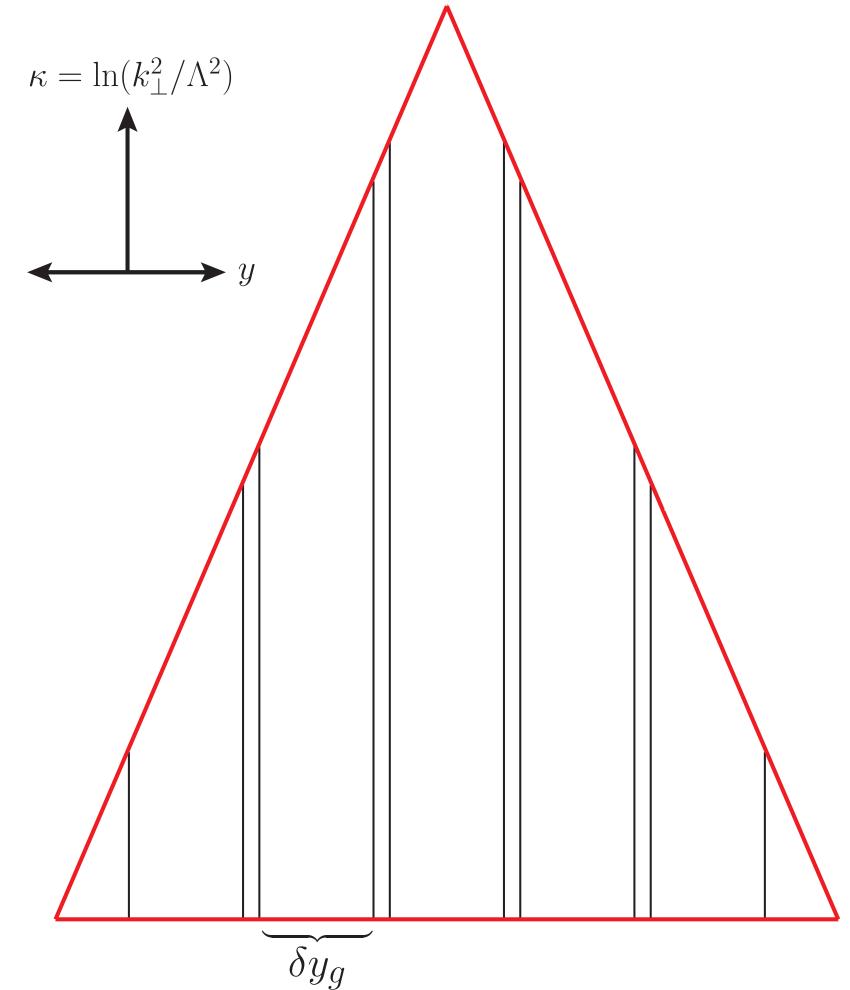
leads to the inclusive probability

$$d\mathcal{P}\left(q(p_{\mathrm{I}})\bar{q}(p_{\mathrm{K}}) \to q(p_{i})g(p_{j})\bar{q}(p_{k})\right) \simeq = \frac{d\kappa}{\kappa} \frac{dy}{\delta y_{g}} \quad \text{with} \quad \delta y_{g} = \frac{11}{6}$$

Interpreting the running coupling renormalisation group as a gainloss equation:

Gluons within δy_g act coherently as one effective gluon





2. Neglect $g \rightarrow q\overline{q}$ splittings and examine transverse-momentum-dependent running coupling

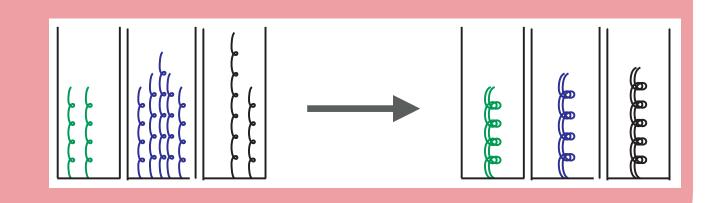
$$\alpha_s(k_\perp^2) = \frac{12\pi}{33 - 2n_f} \frac{1}{\ln(k_\perp^2/\Lambda_{\rm QCD}^2)} = \frac{\text{const.}}{\kappa}$$

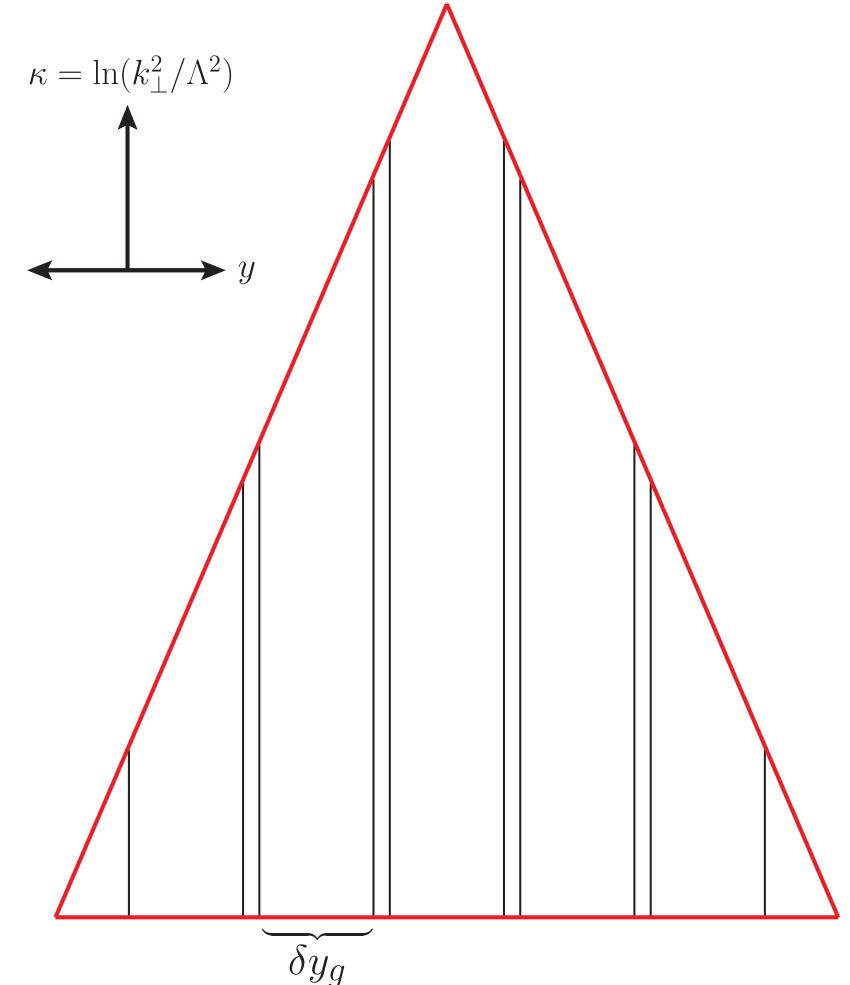
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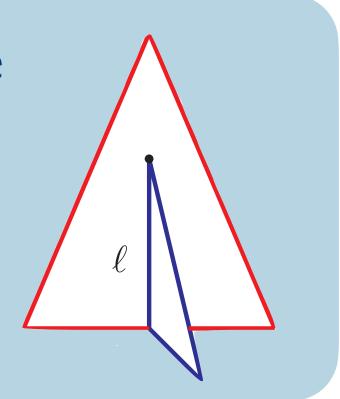
Interpreting the running coupling renormalisation group as a gainloss equation:

Gluons within δy_g act coherently as one effective gluon





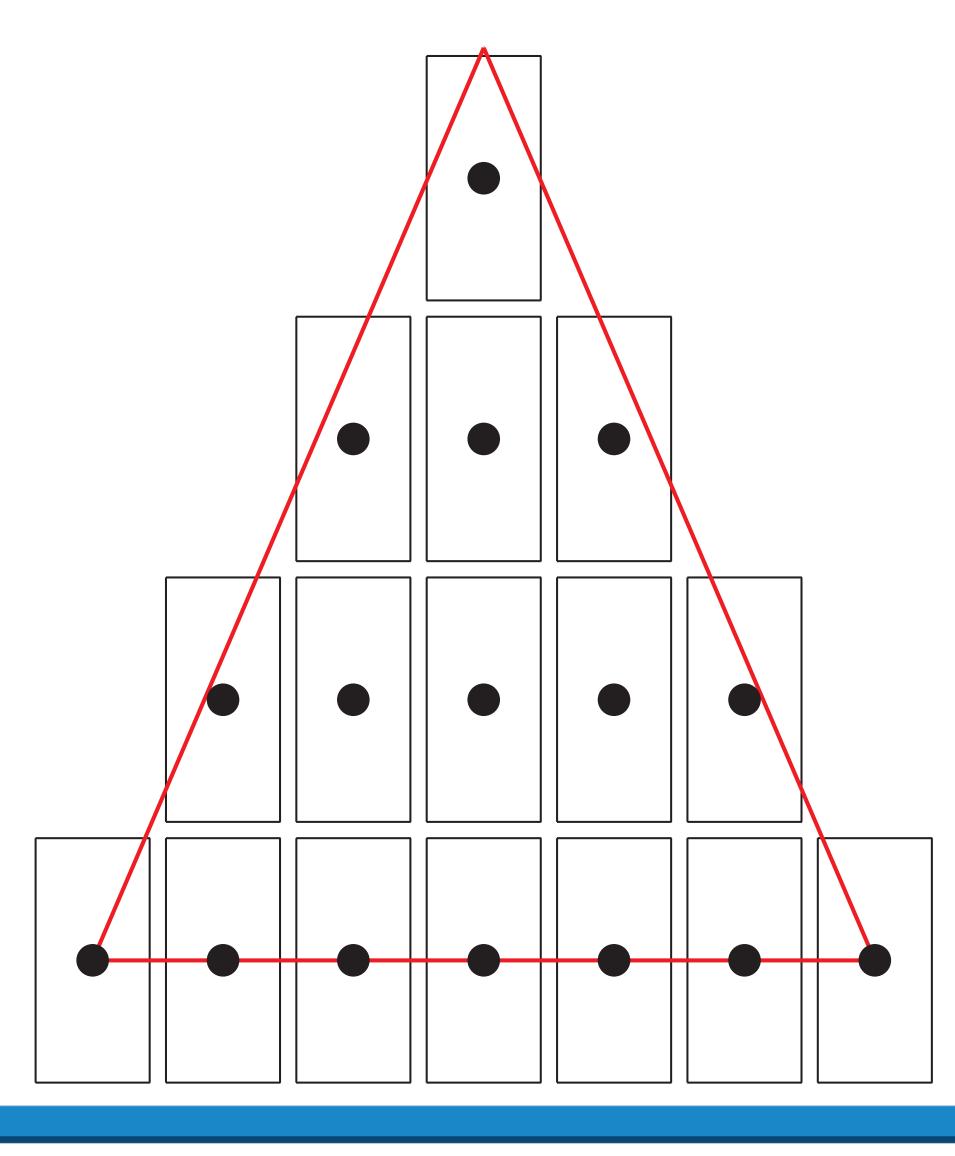
Folding out extends the baseline of the triangle to positive y by $\frac{l}{2}$, where l is the height at which to emit effective gluons



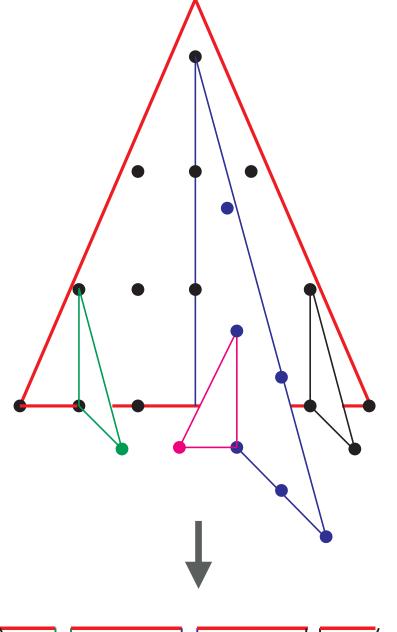
A consequence of folding is that the κ axis is quantised into multiples of $2\delta y_g$

Each rapidity slice can be treated independently of any other slice. The exclusive rate probability takes the simple form:

$$\frac{d\kappa}{\kappa} \exp\left(-\int_{\kappa}^{\kappa_{max}} \frac{d\bar{\kappa}}{\bar{\kappa}}\right) = \frac{d\kappa}{\kappa_{max}}$$

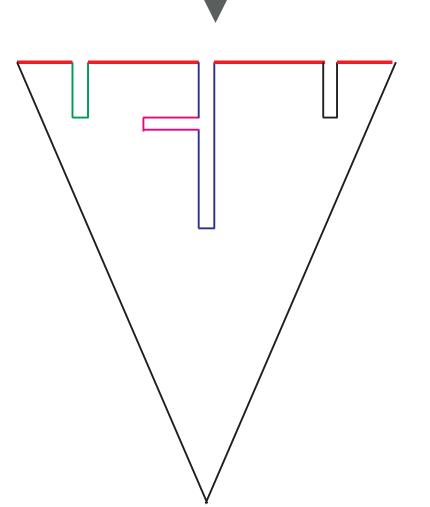


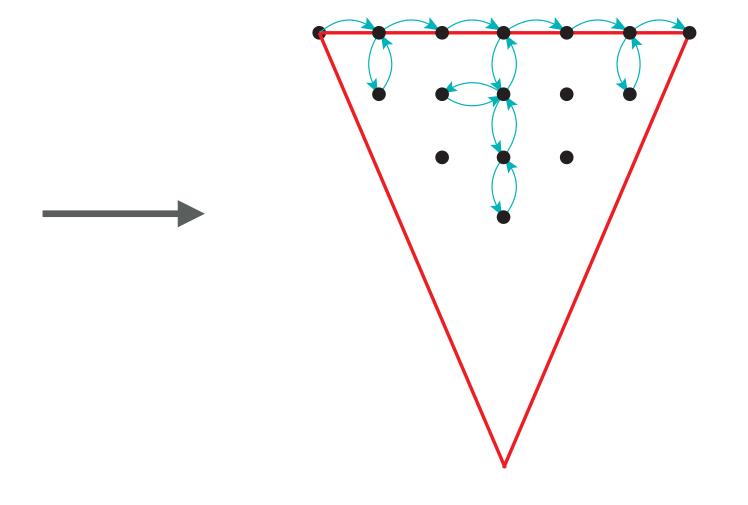
Discrete QCD as a Quantum Walk



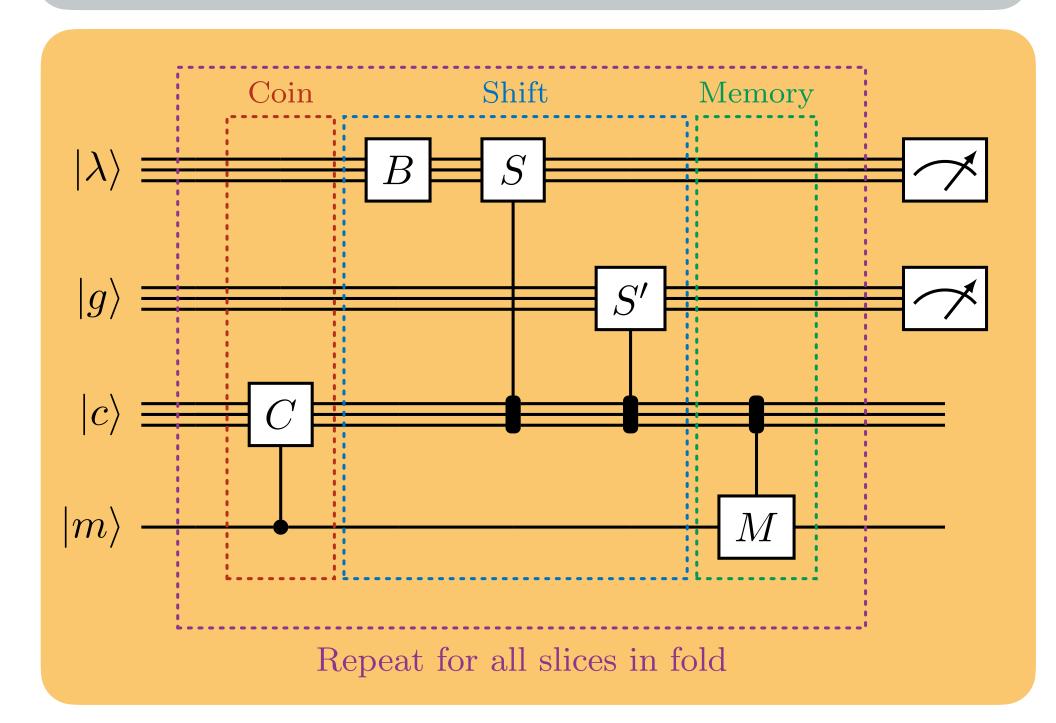
The **baseline** of the grove structure contains all kinematics information

For LEP data there are **24 unique grove** structures for $\Lambda_{\text{OCD}} \in [0.1,1]$ GeV

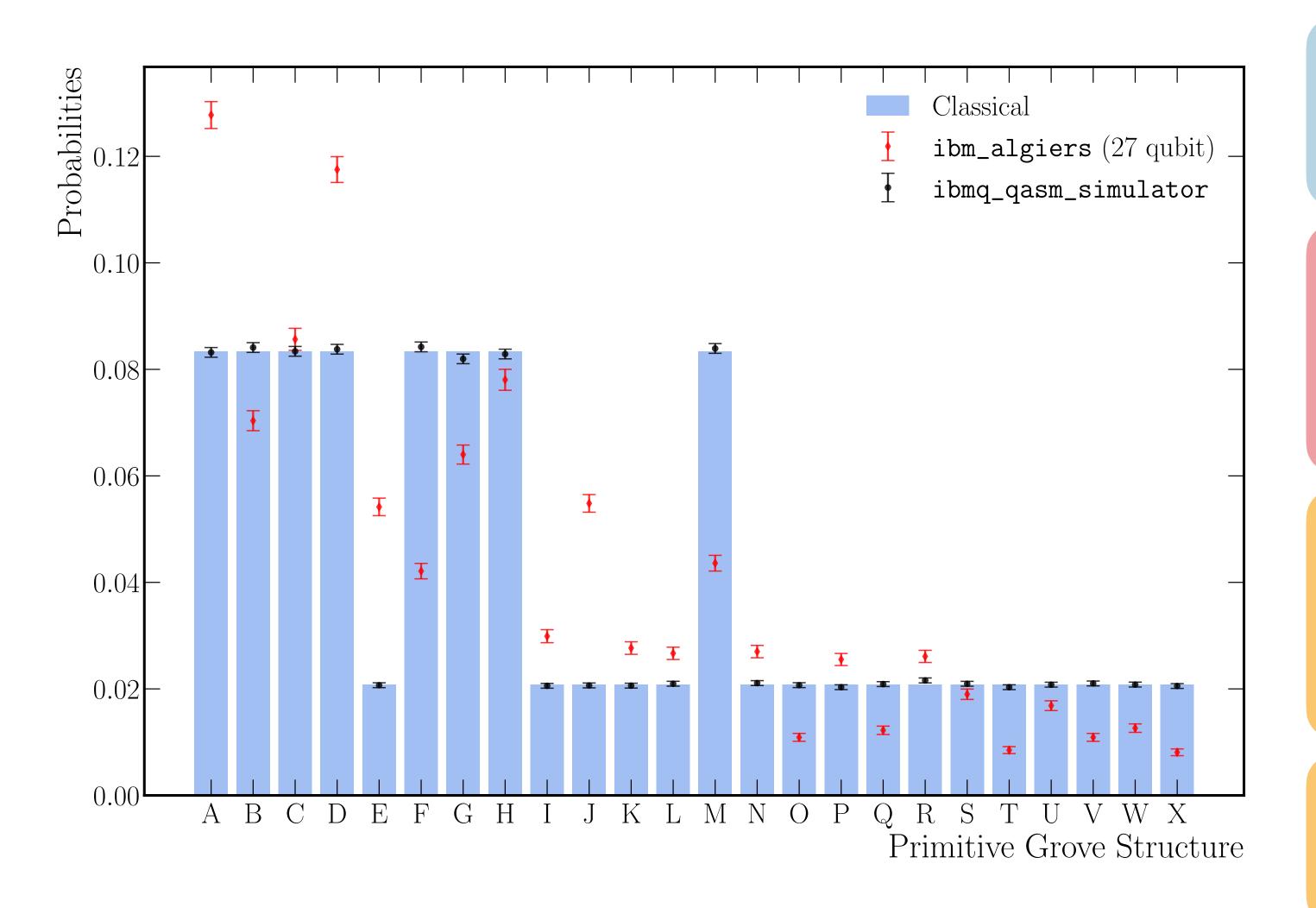




The Discrete-QCD dipole cascade can therefore be implemented as a simple Quantum Walk



Discrete QCD as a Quantum Walk - Raw Grove Simulation



The algorithm has been run on the IBM Falcon 5.1 Ir chip

The figure shows the uncorrected performance of the **ibm_algiers** device compared to a simulator

The 24 grove structures are generated for a $E_{CM}=91.2$ GeV, corresponding to typical collisions at LEP.

Main source of error from CNOT errors from large amount of SWAPs

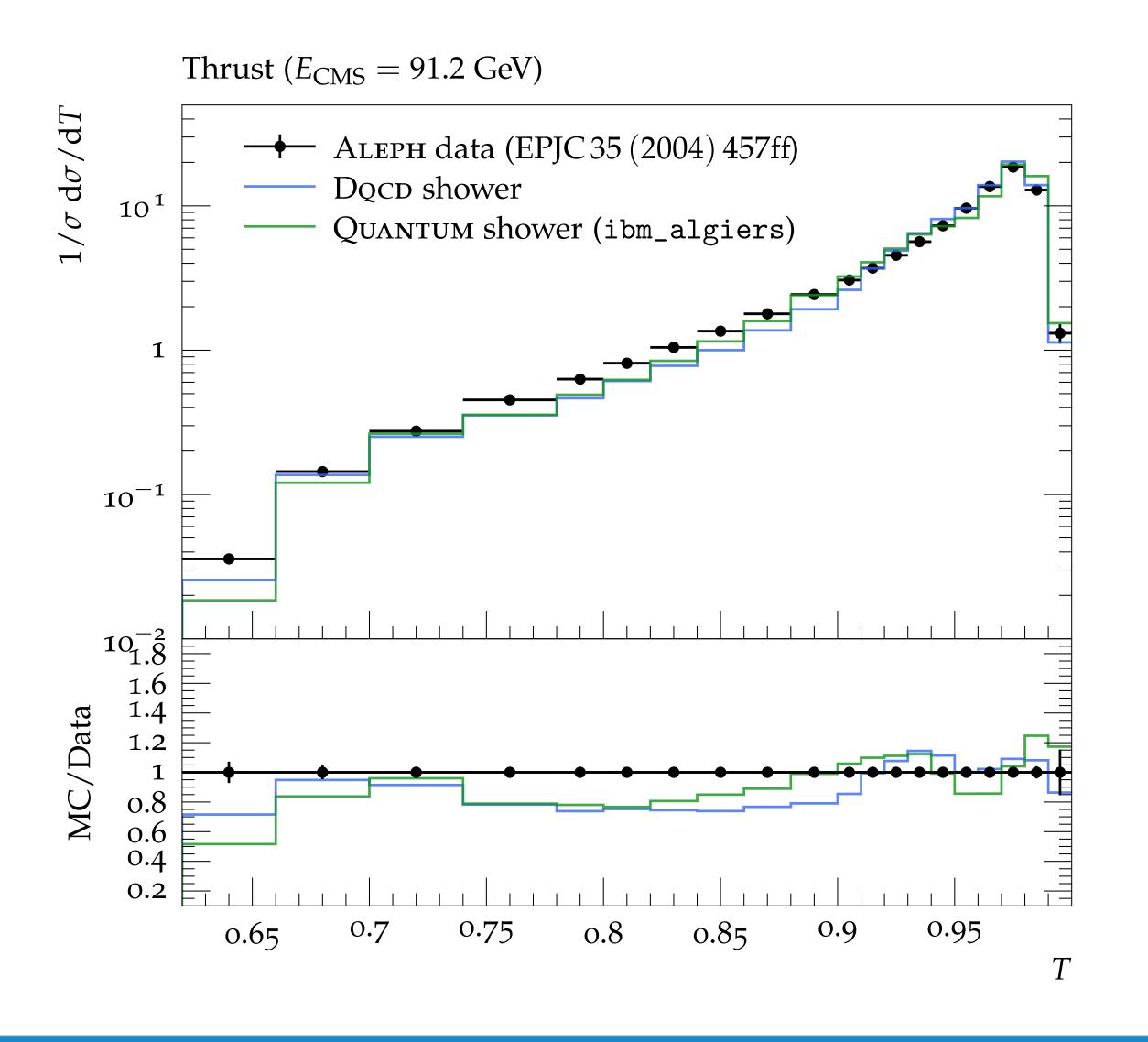
Generating Scattering Events from Groves

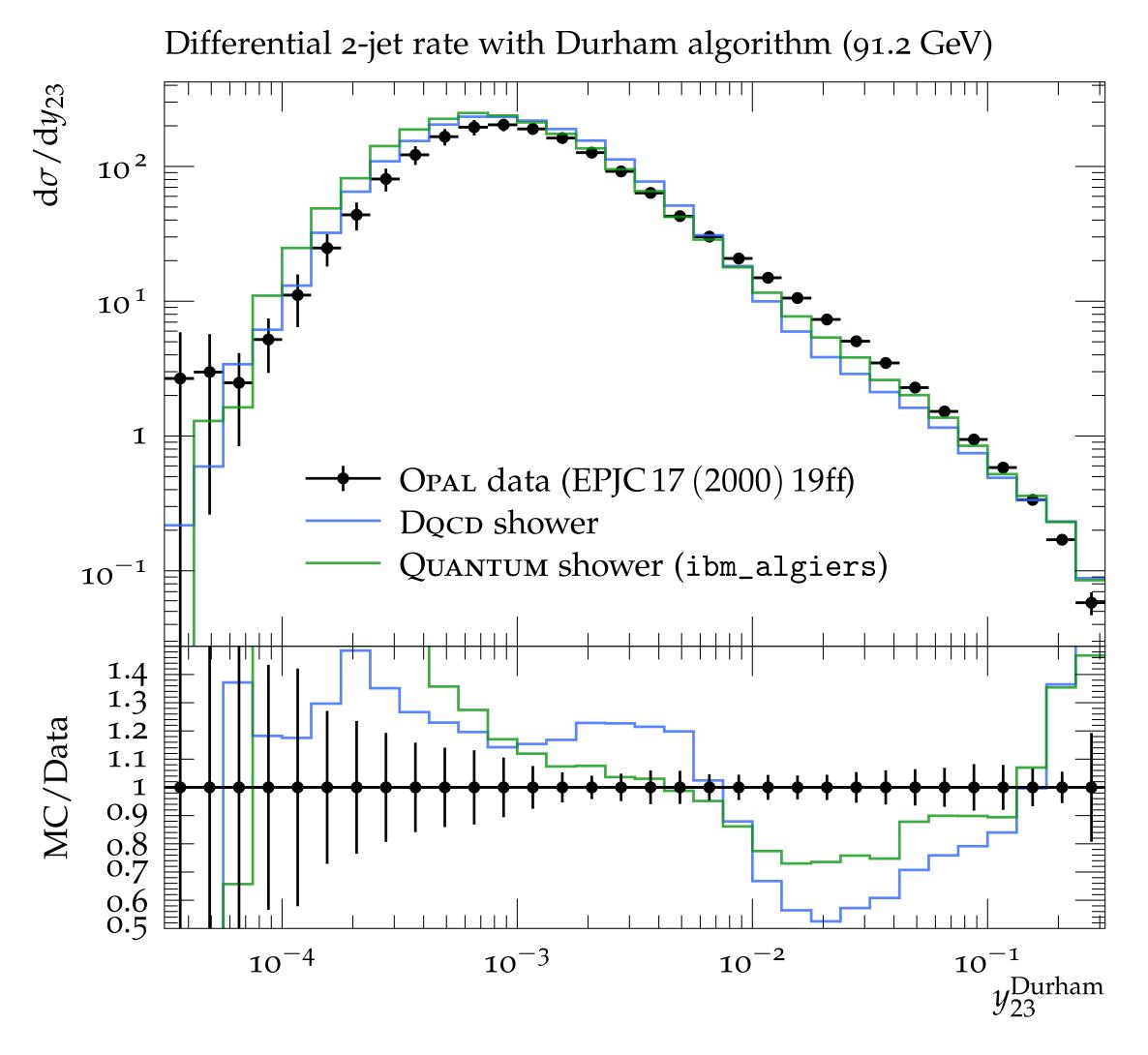
Once the grove structure has been selected, event data can be synthesised in the following steps using the baseline:

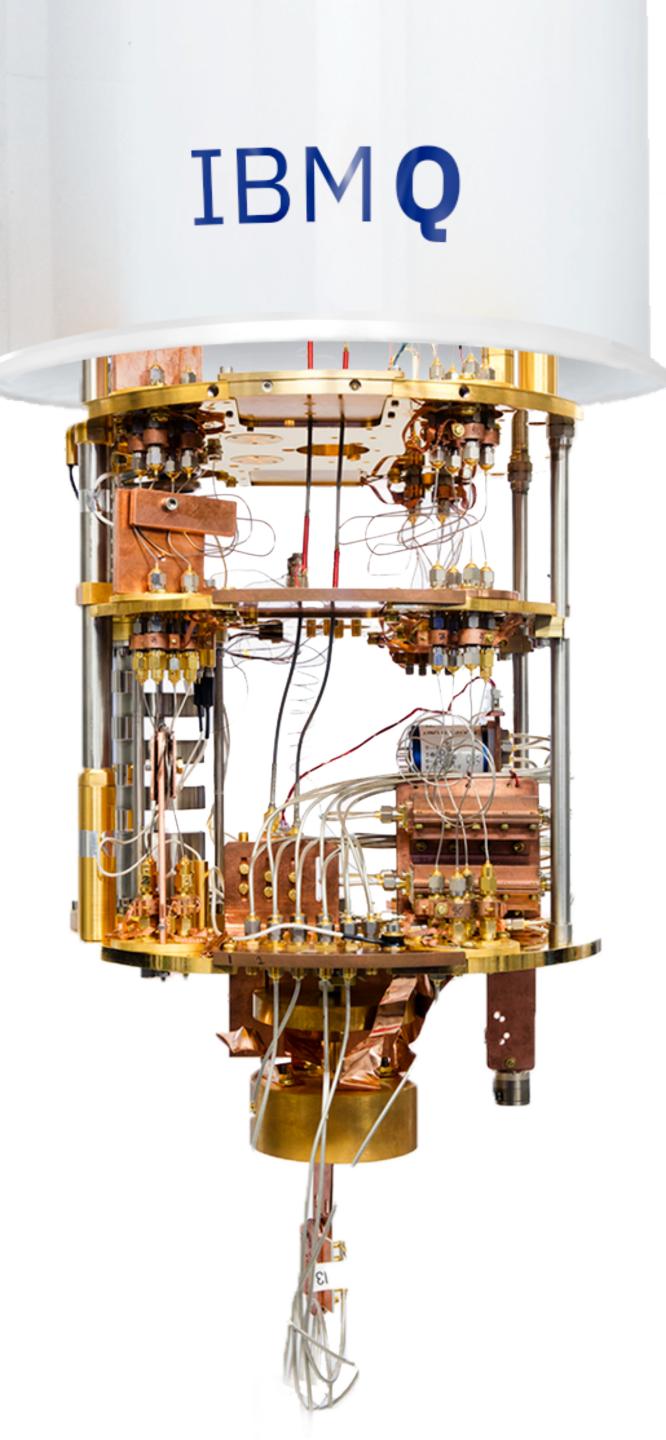
- 1. Create the highest κ effective gluons first (i.e. go from top to bottom in phase space)
- 2. For each effective gluon j that has been emitted from a dipole IK, read off the values s_{ij} , s_{jk} and s_{IK} from the grove
- 3. Generate a uniformly distributed azimuthal decay angle ϕ , and then employ momentum mapping (here we have used Phys. Rev. D 85, 014013 (2012), 1108.6172) to produce post-branching momenta

This has been done using the ibm_cloud 27 qubit device ibm_algiers, with 20,000 shots on the device. A comparison with a like-for-like classical parton shower algorithm has been made.

Collider Events on a Quantum Computer







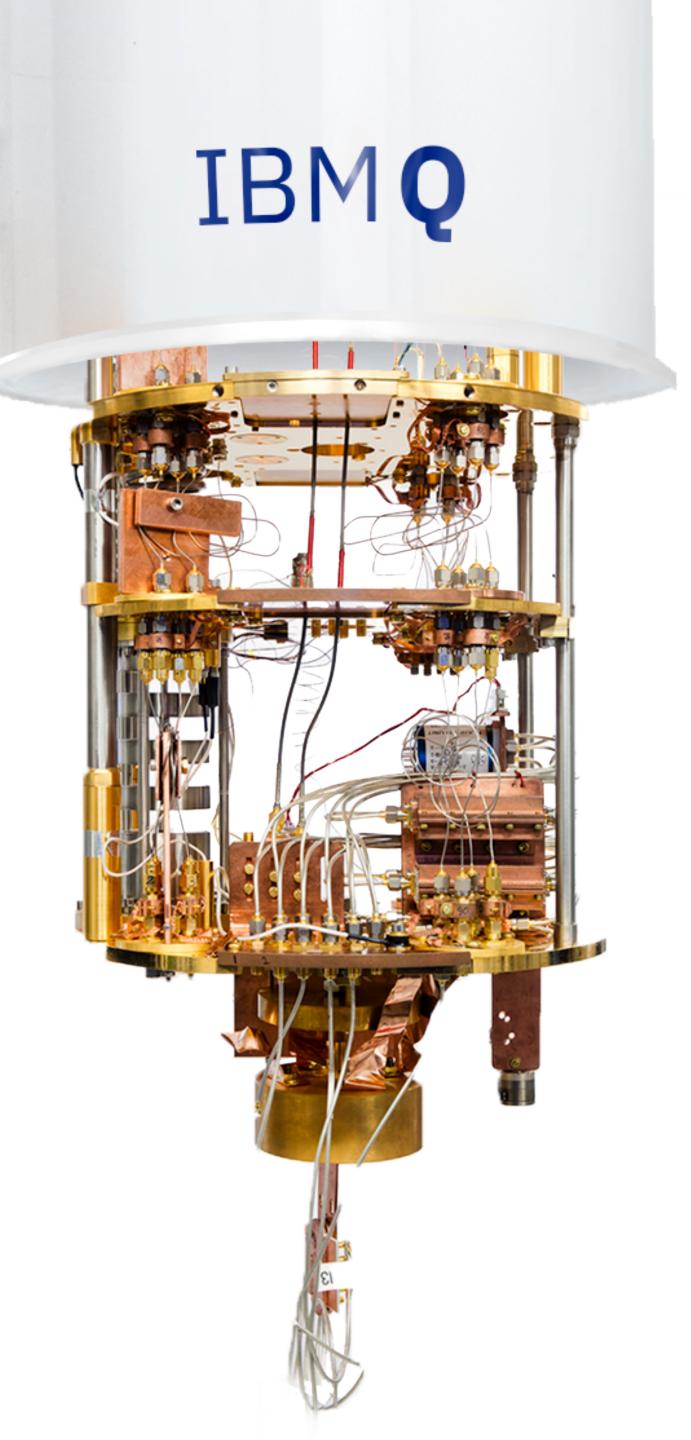
Summary

High Energy Physics is on the edge of a computational frontier, the High Luminosity Large Hadron Collider and FCC will provide unprecedented amounts of data

Quantum Computing offers an impressive and powerful tool to combat computational bottlenecks, both for theoretical and experimental purposes

The first realistic simulation of a high energy collision has been presented using a compact quantum walk implementation, allowing for the algorithm to be run on a NISQ device

Future Work: A dedicated research effort is required to fully evaluate the potential of quantum computing applications in HEP



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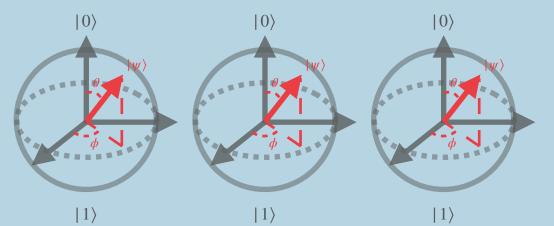
Backup Slides

Simon Williams

Resummation, Evolution, Factorization 2022

The Future of Quantum Computing

More qubits?



A lot of emphasis on more qubits, but without fault tolerance, large qubit devices become impractical

Be better architects?

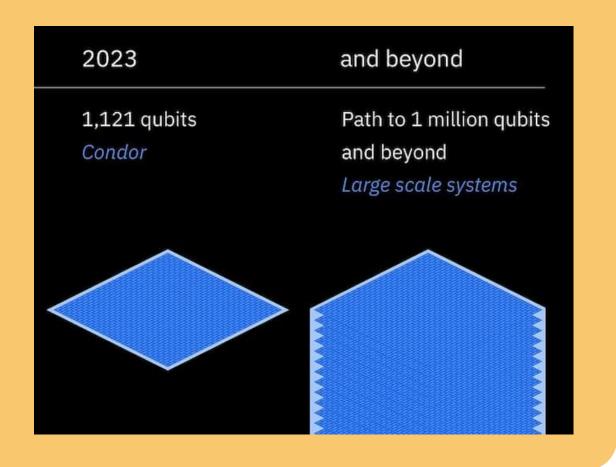
Realistic algorithms are already being created for NISQ devices. Efficient architectures allow for **practical algorithms** on NISQ devices.

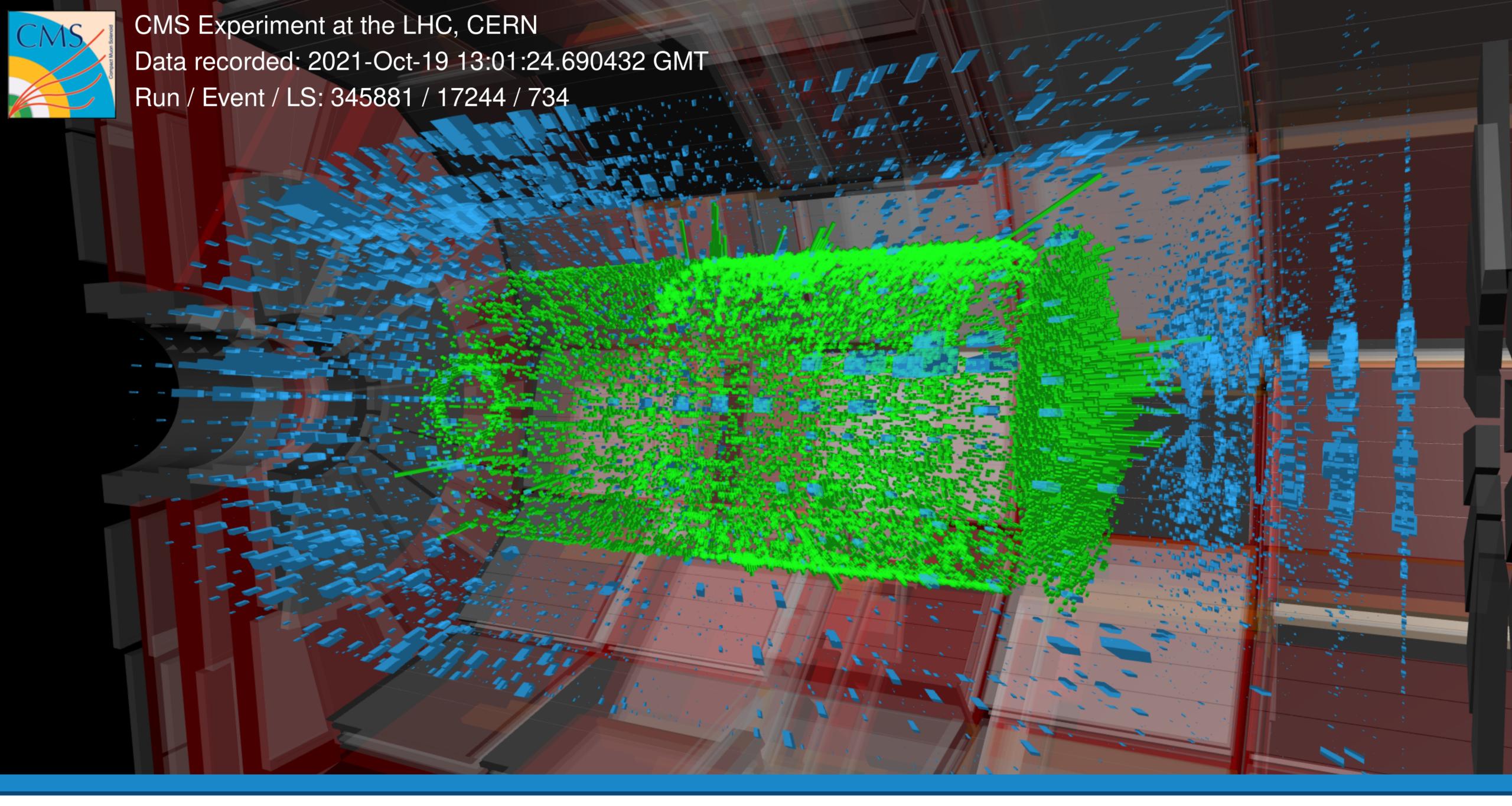
Better technology?

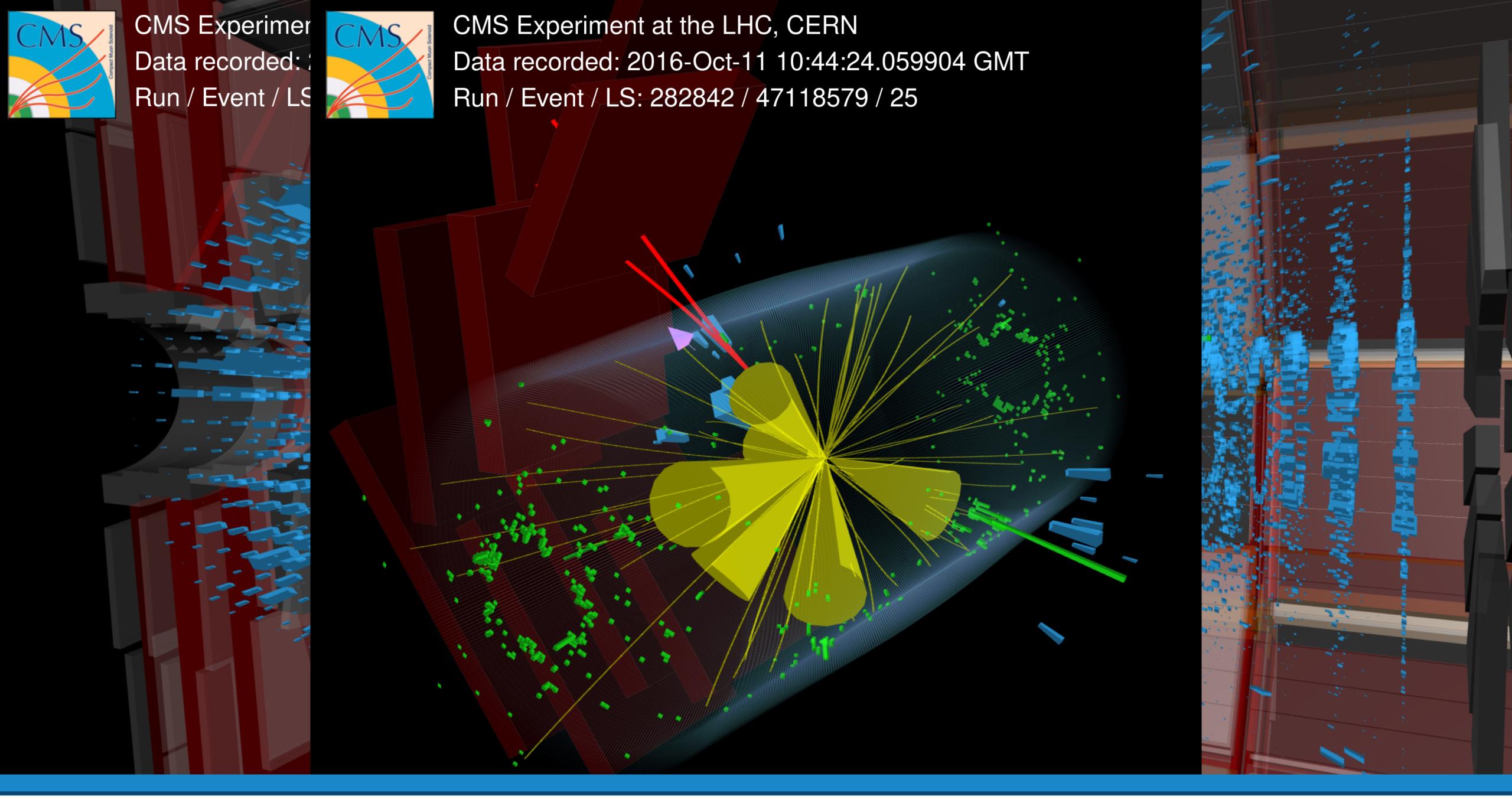
New technology could be the answer - will new qubit hardwares be more **fault tolerant?**

IBM Roadmap

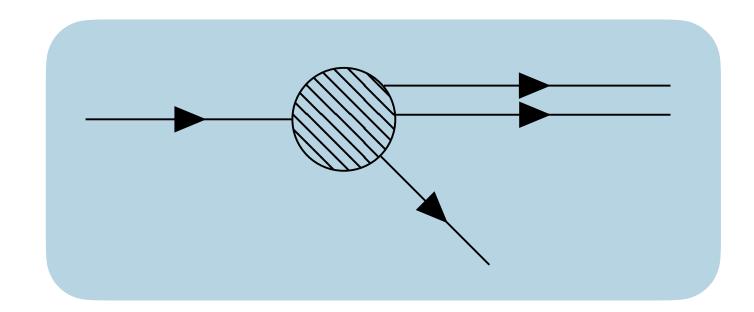
On track to deliver 1000 qubits by 2023

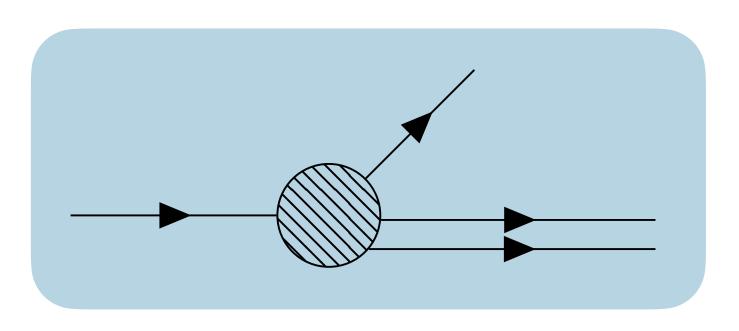






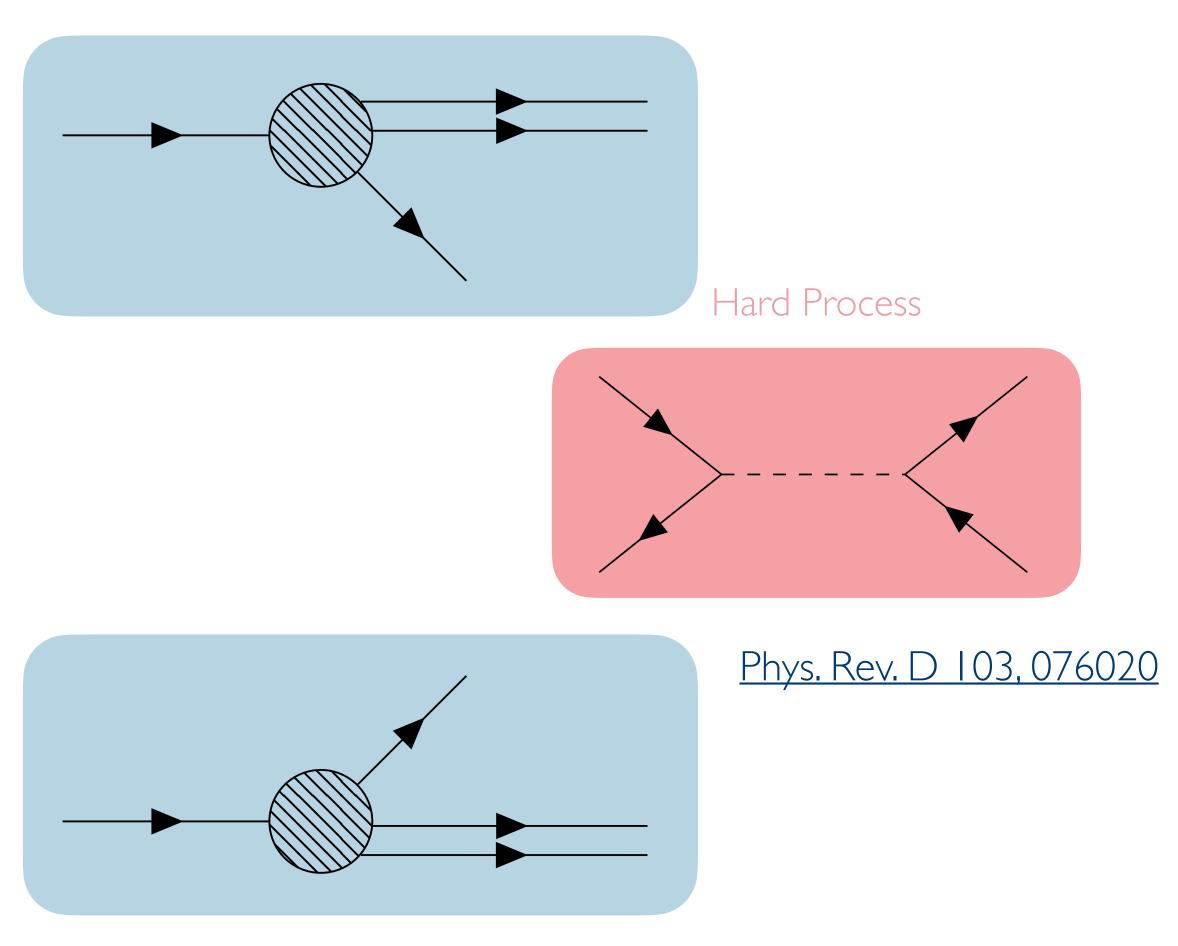
Parton Density Functions





Phys. Rev. D 103, 034027

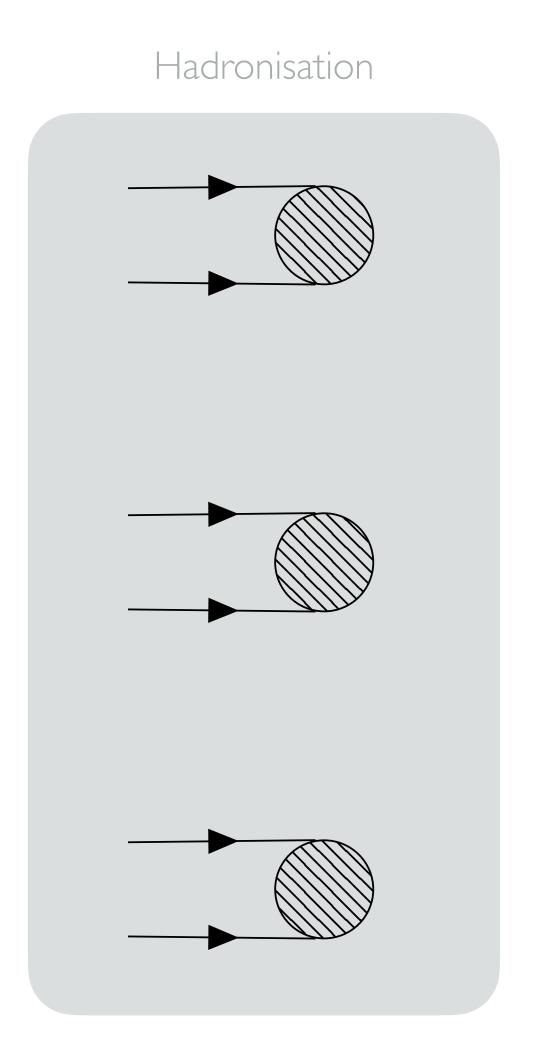
Parton Density Functions

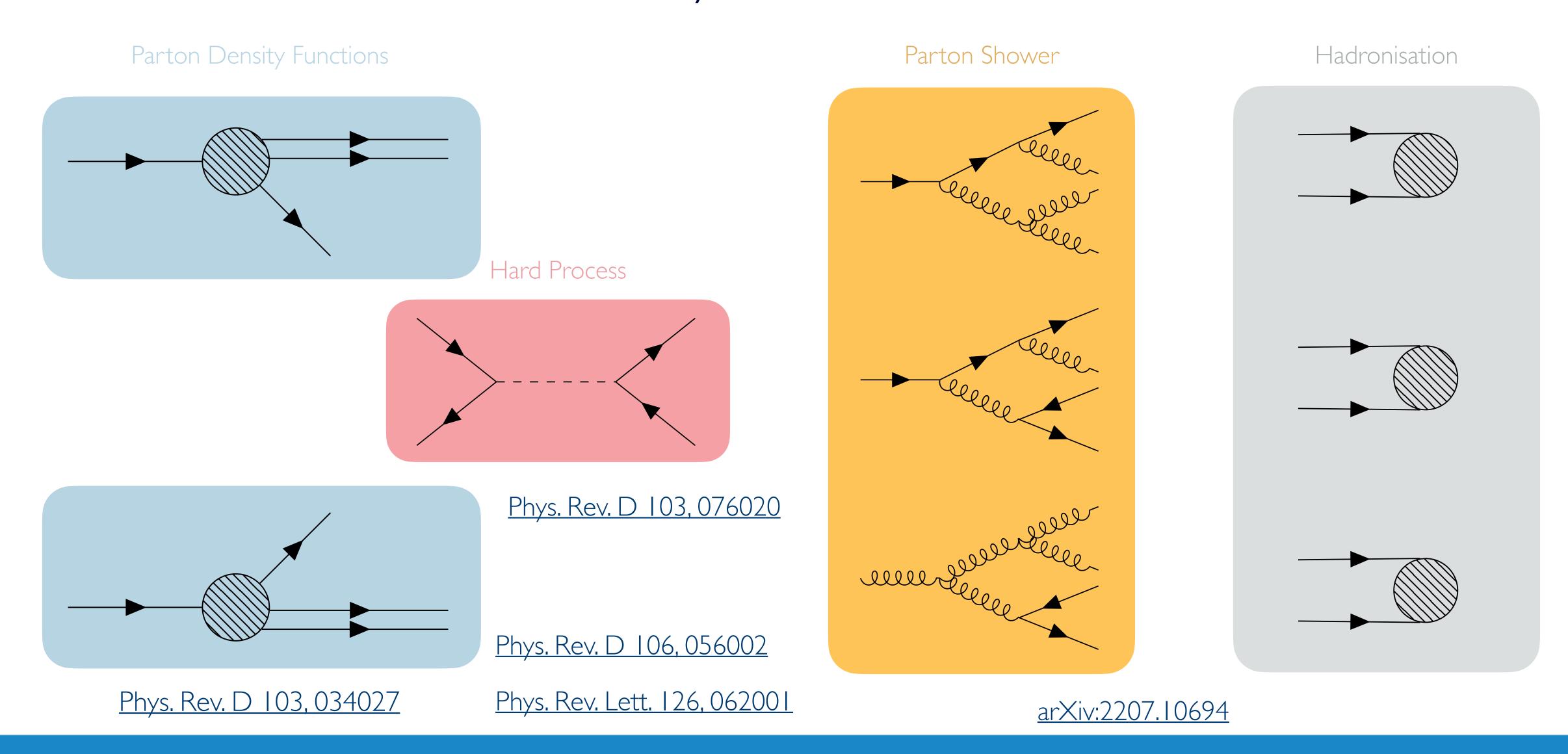


Phys. Rev. D 103, 034027

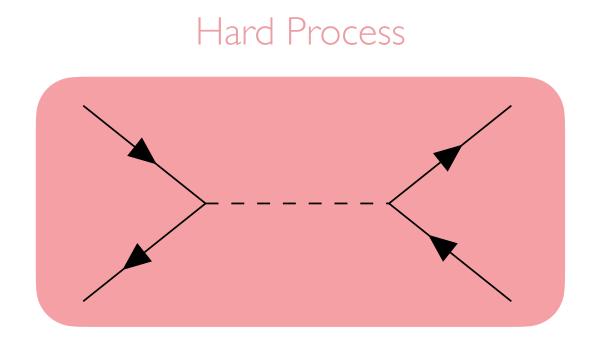
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Parton Density Functions Hard Process Phys. Rev. D 103, 076020 Phys. Rev. D 103, 034027





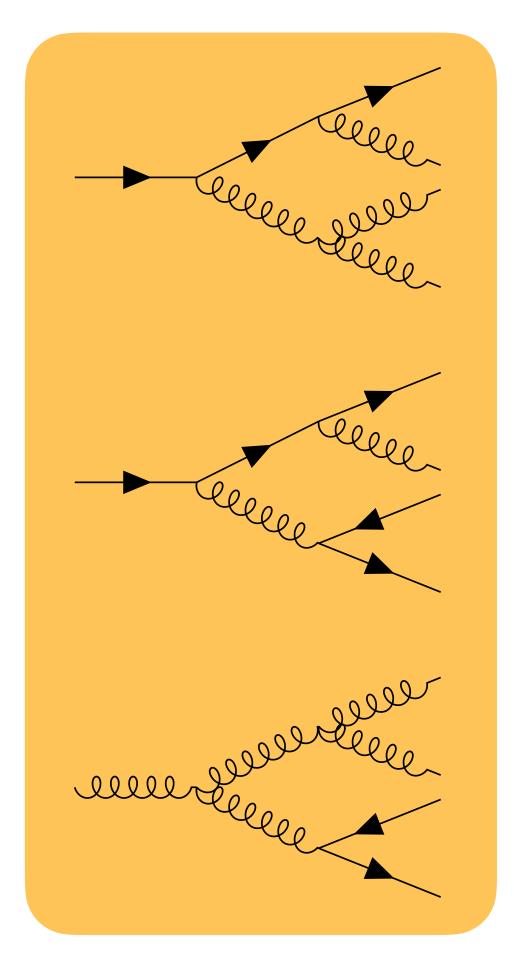
Parton Shower



Phys. Rev. D 103, 076020

Phys. Rev. D 106, 056002

Phys. Rev. Lett. 126, 062001



arXiv:2207.10694

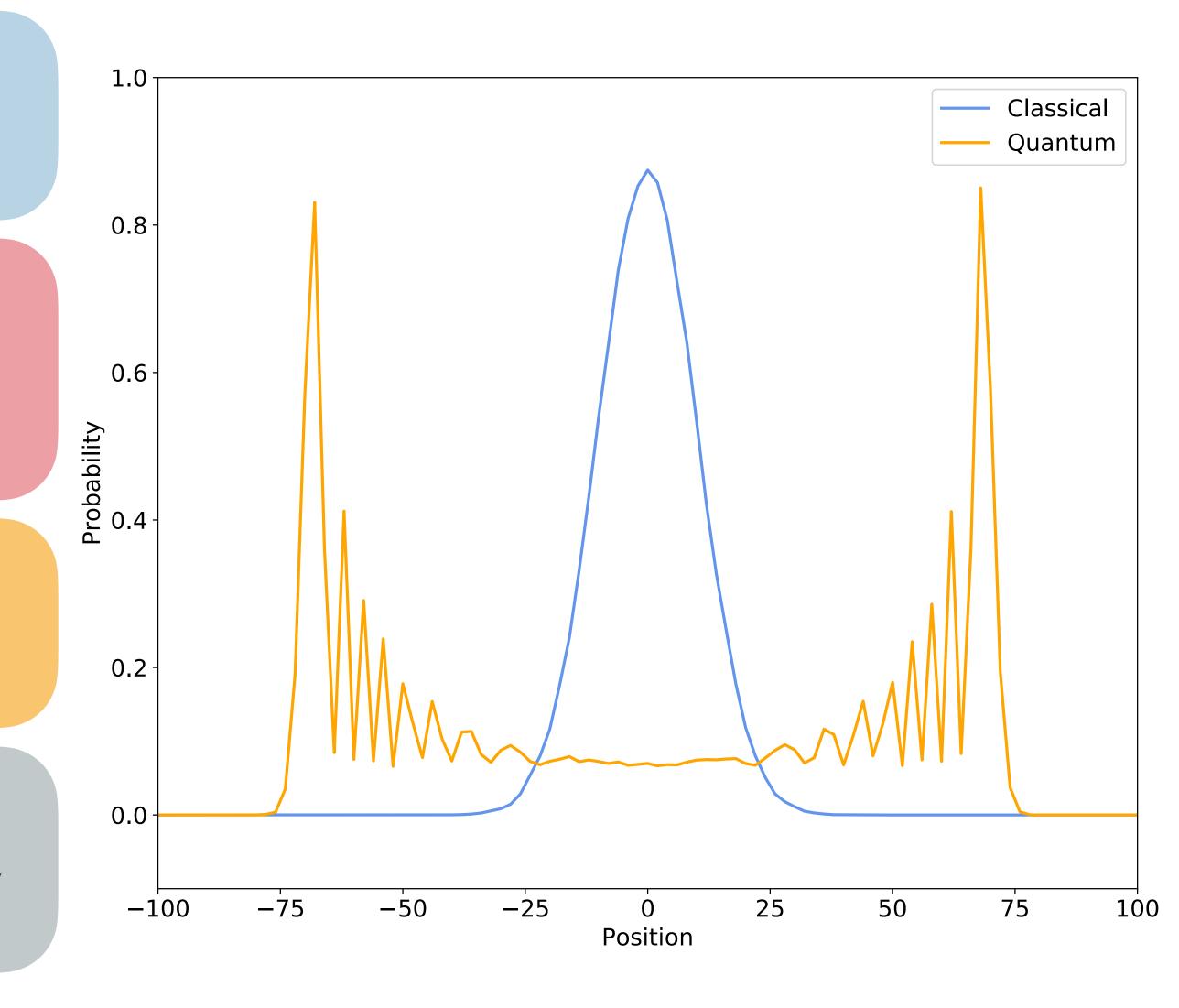
Speed up via Quantum Walks

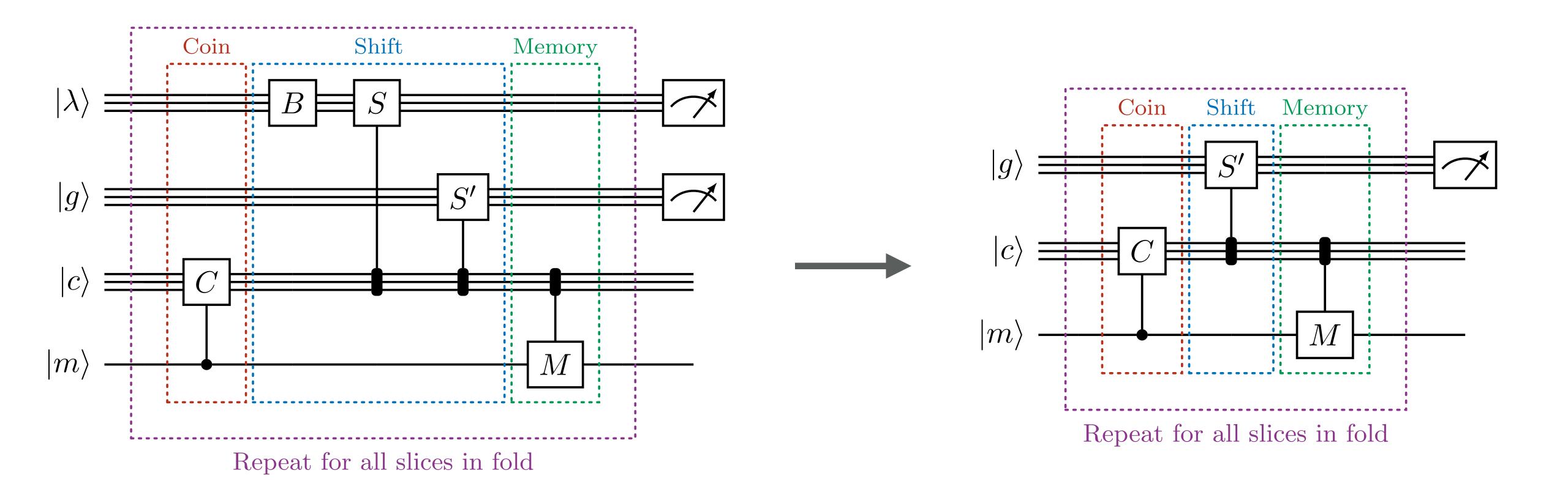
Quantum Walks have long be conjectured to achieved at least quadratic speed up

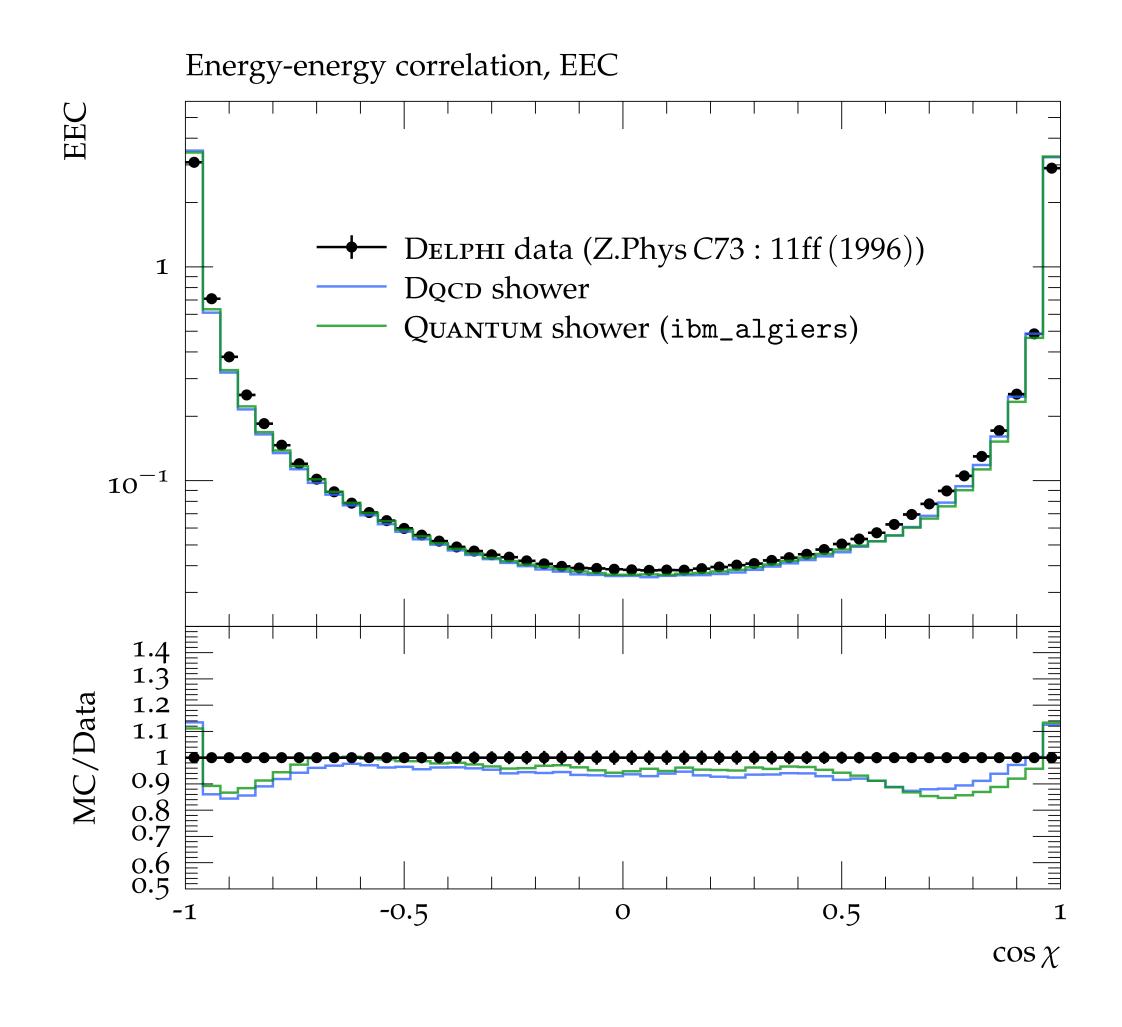
Szegedy Quantum Walks have been proven to achieve quadratic speed up for Markov Chain Monte Carlo

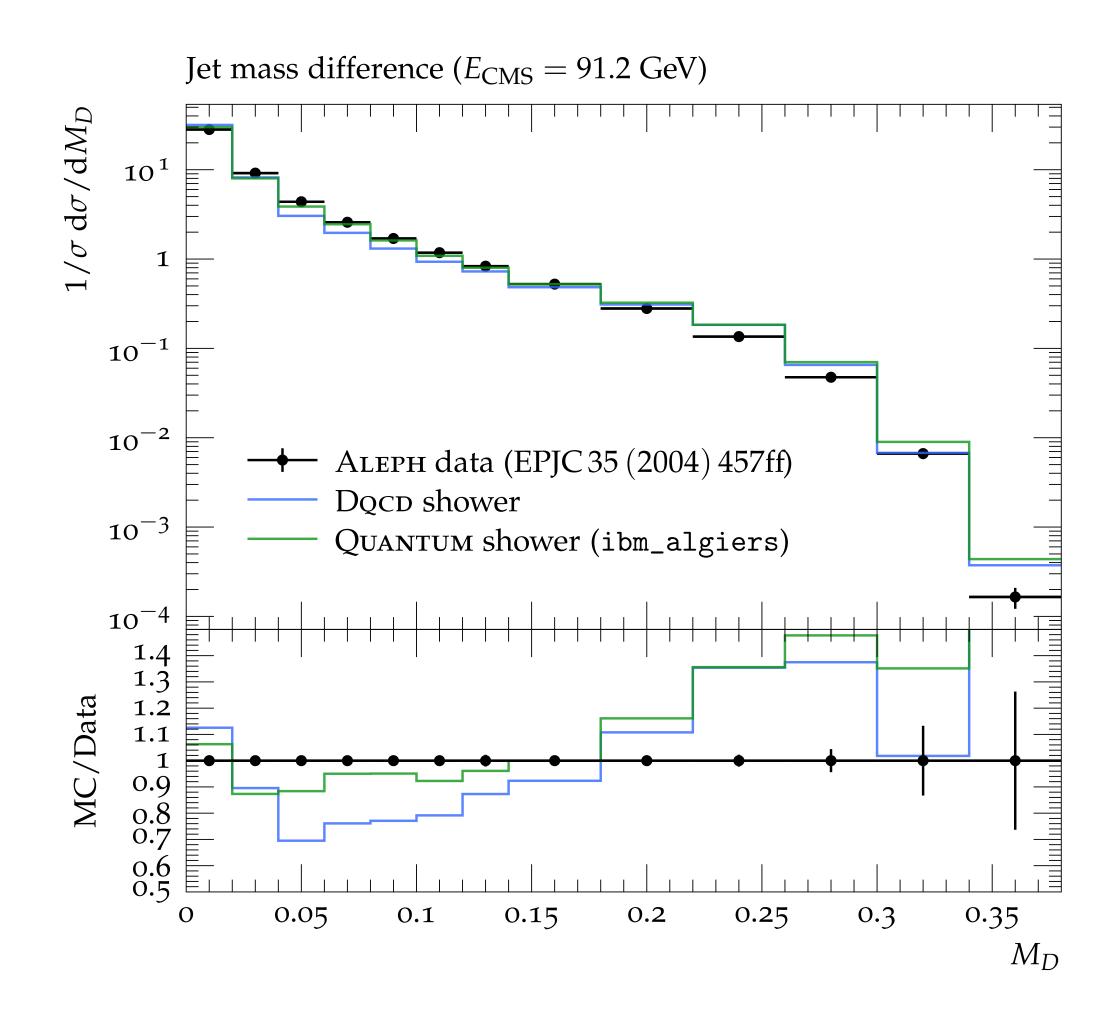
This has been proven under the condition that the MCMC algorithm is reversible and ergodic

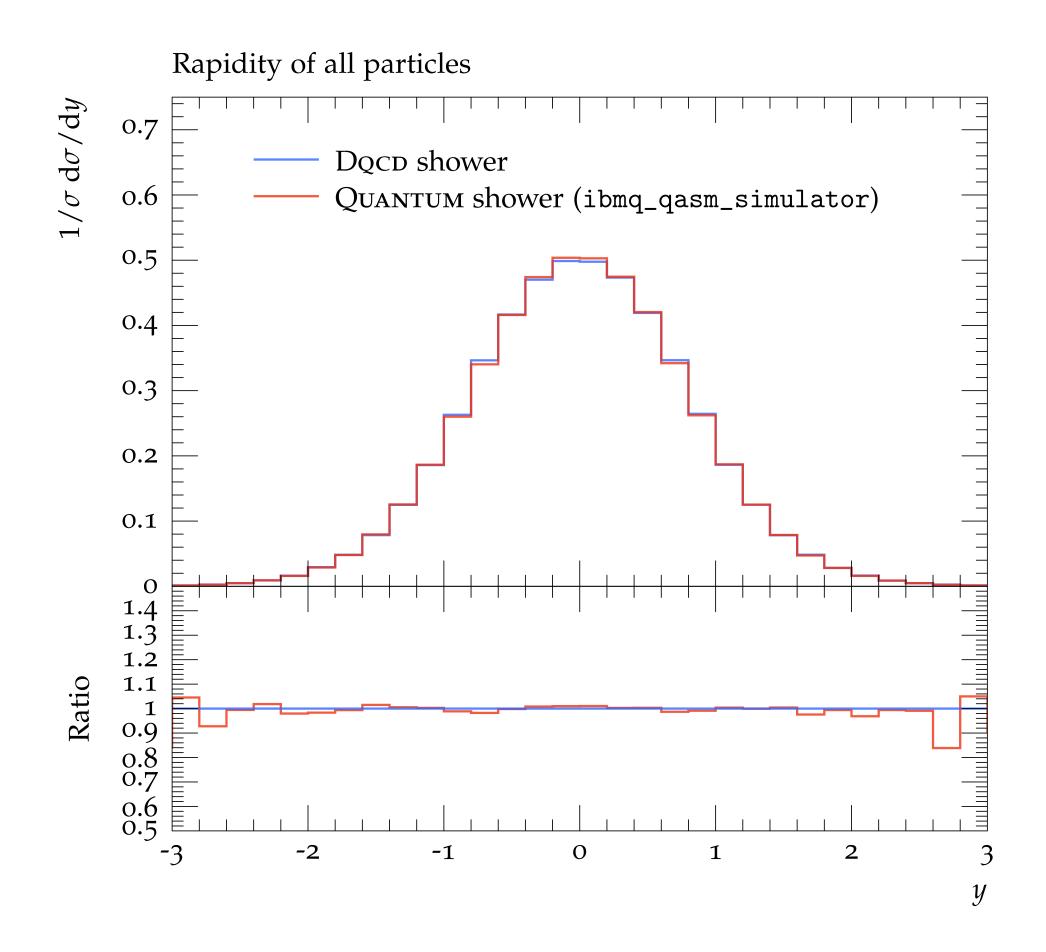
Work is ongoing to prove this is true for all QWs, but latest upper limits are on par with classical RW

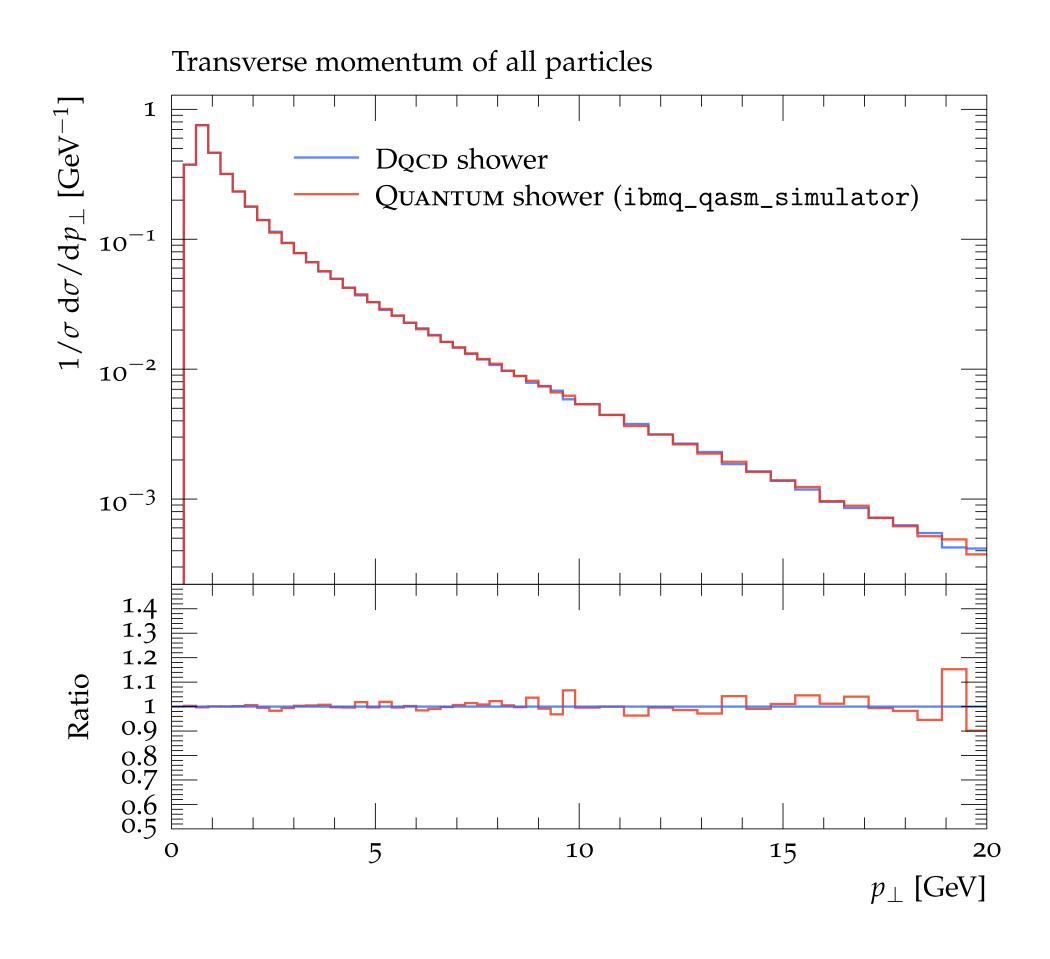


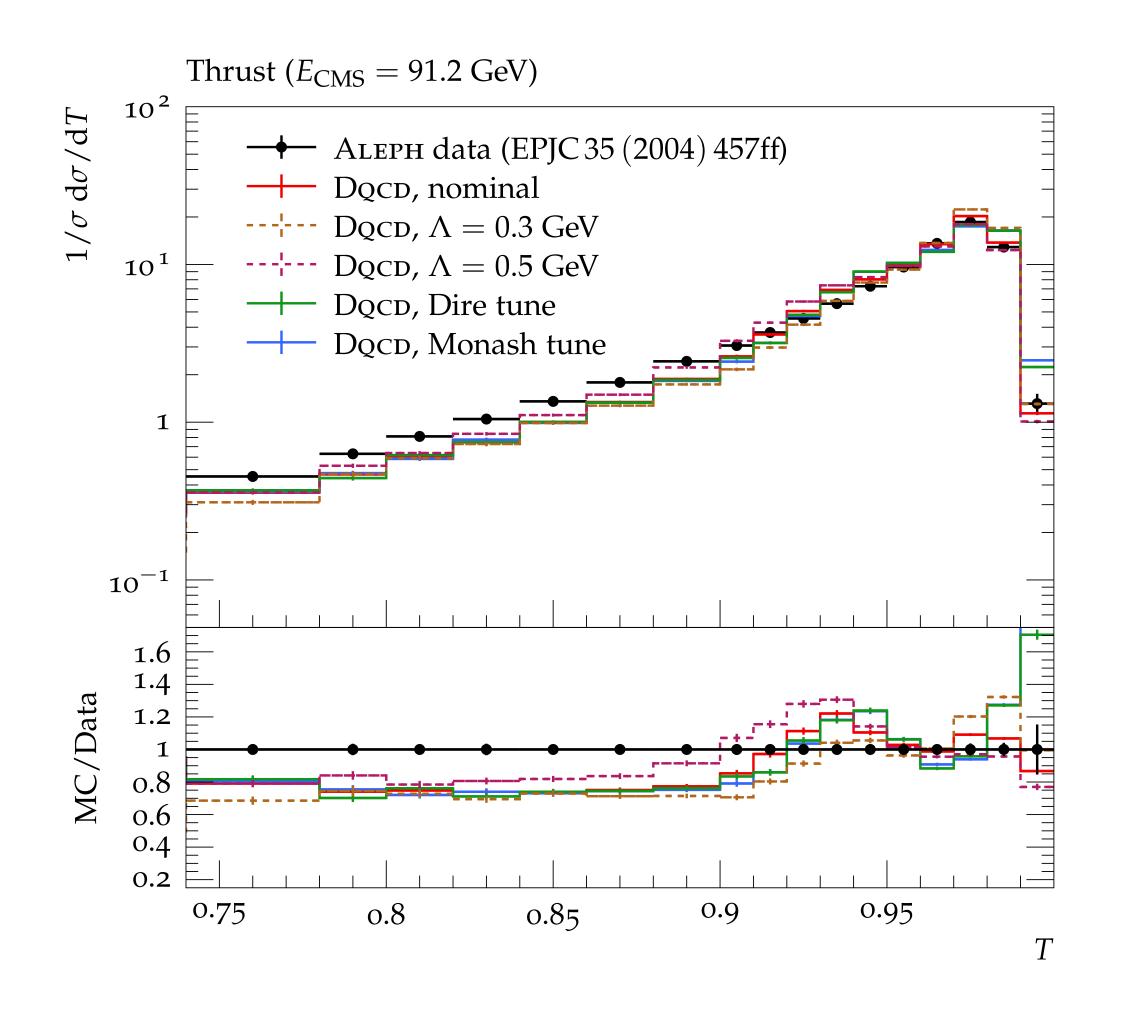


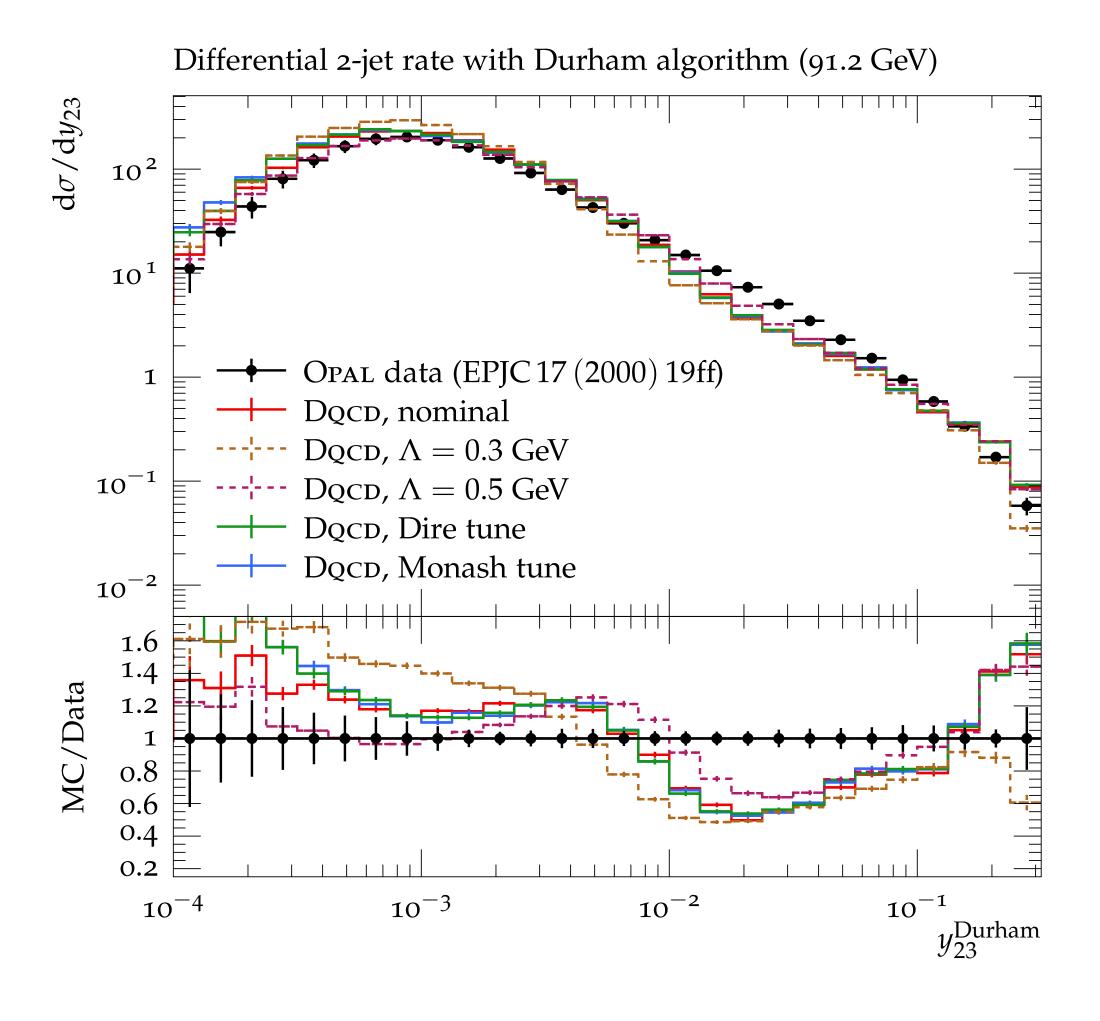


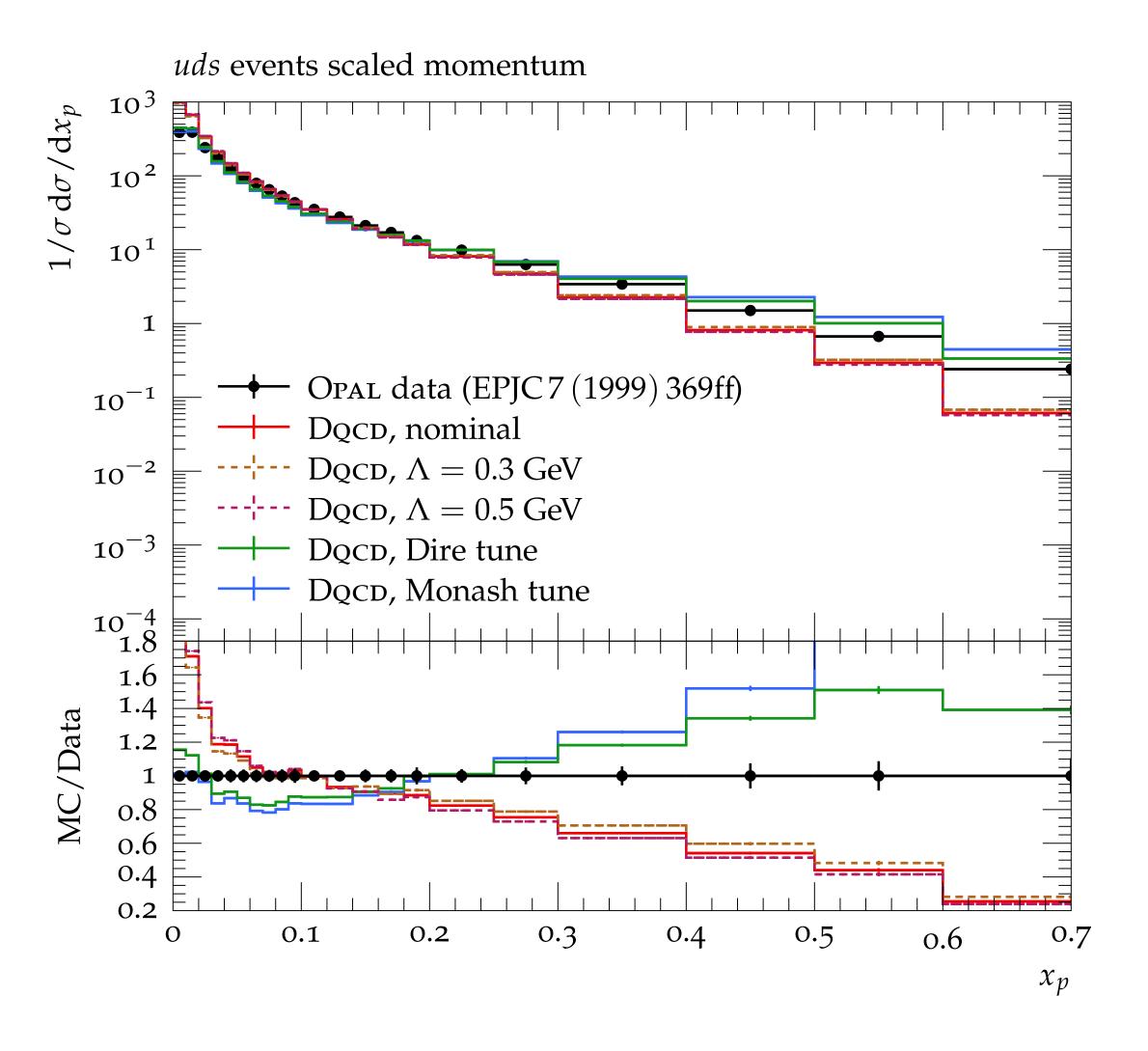


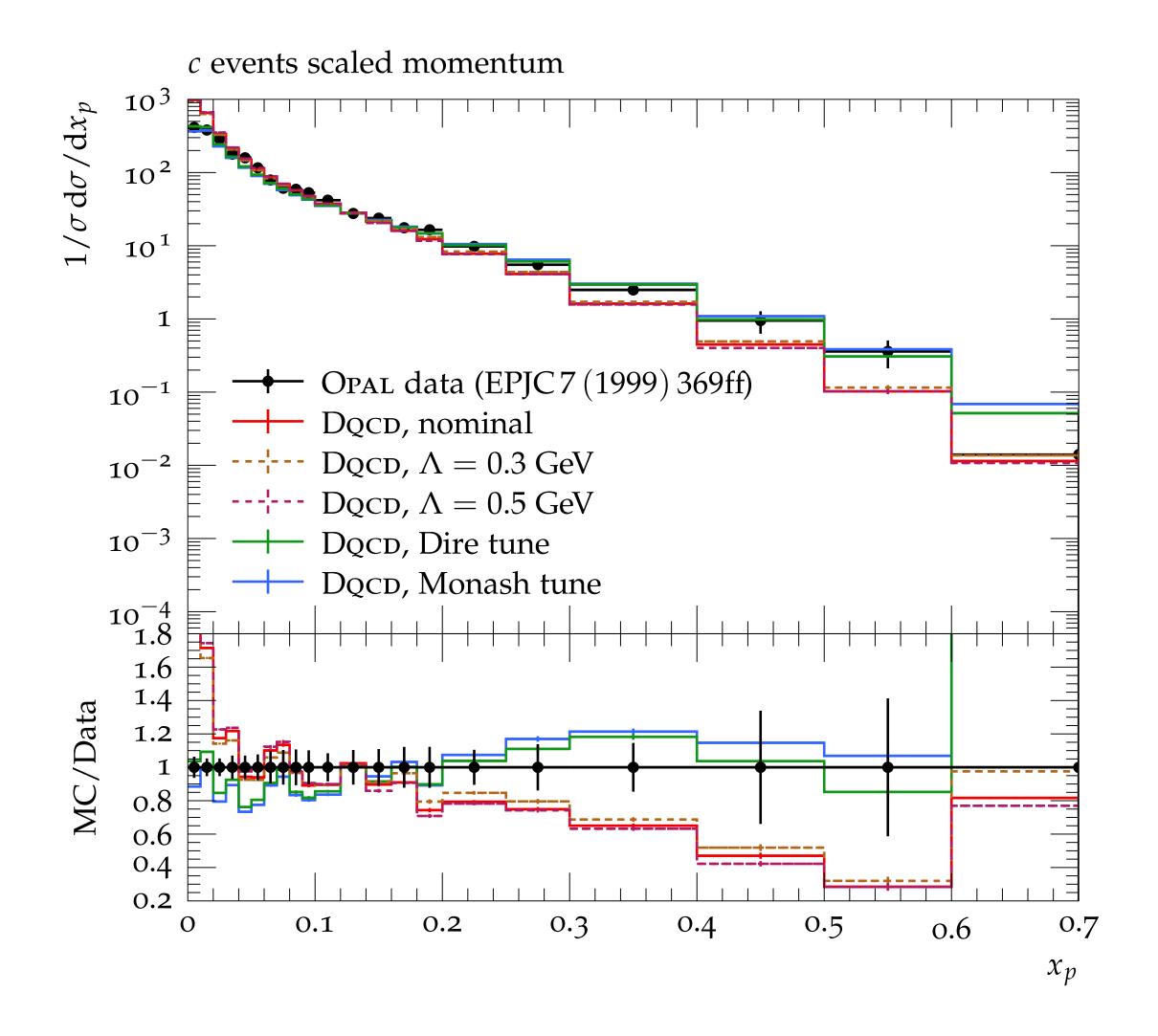


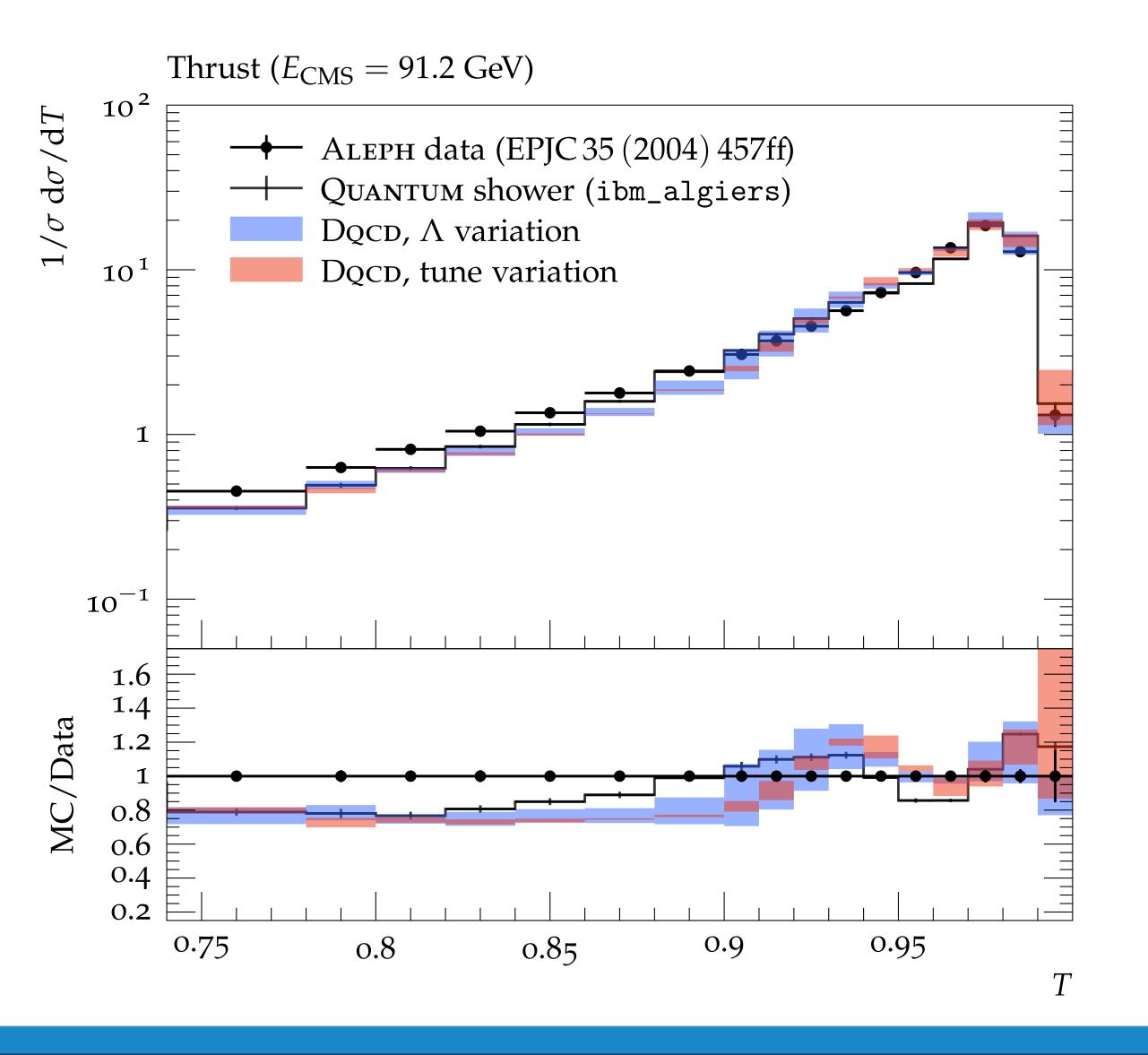


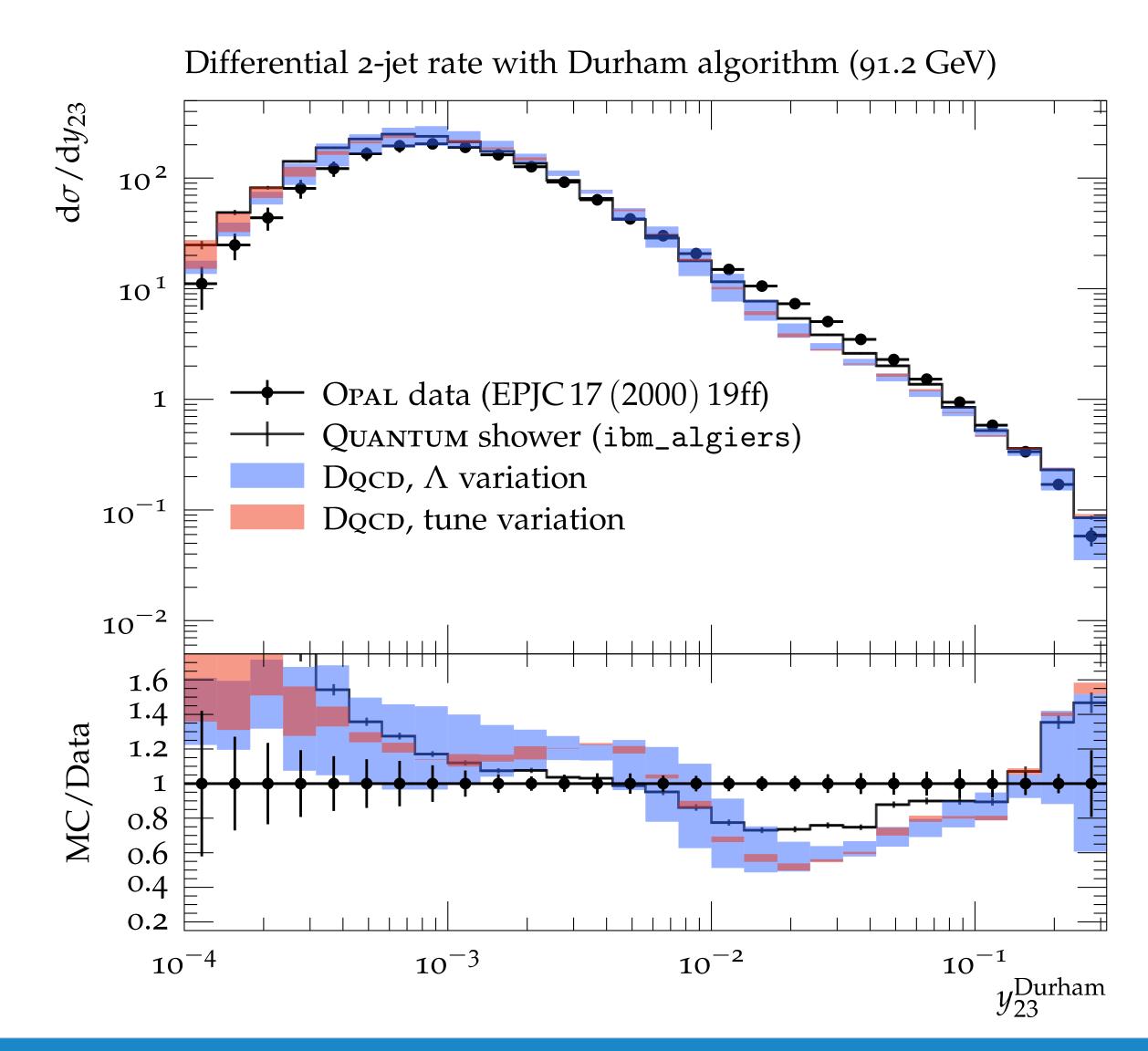


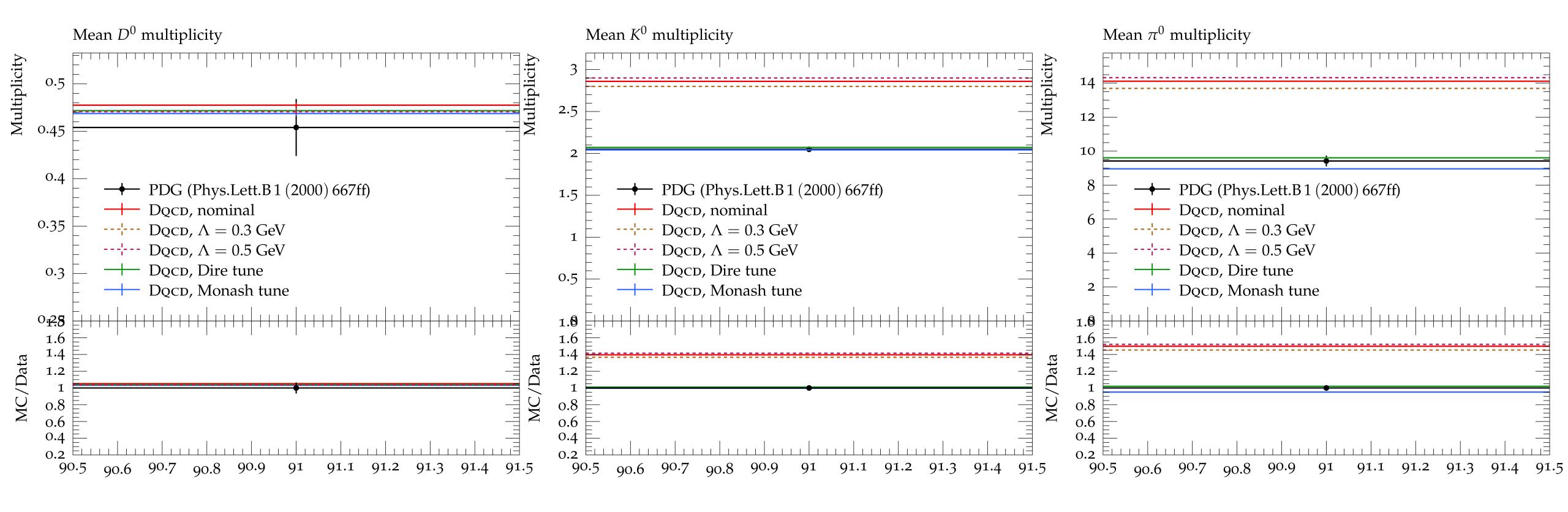






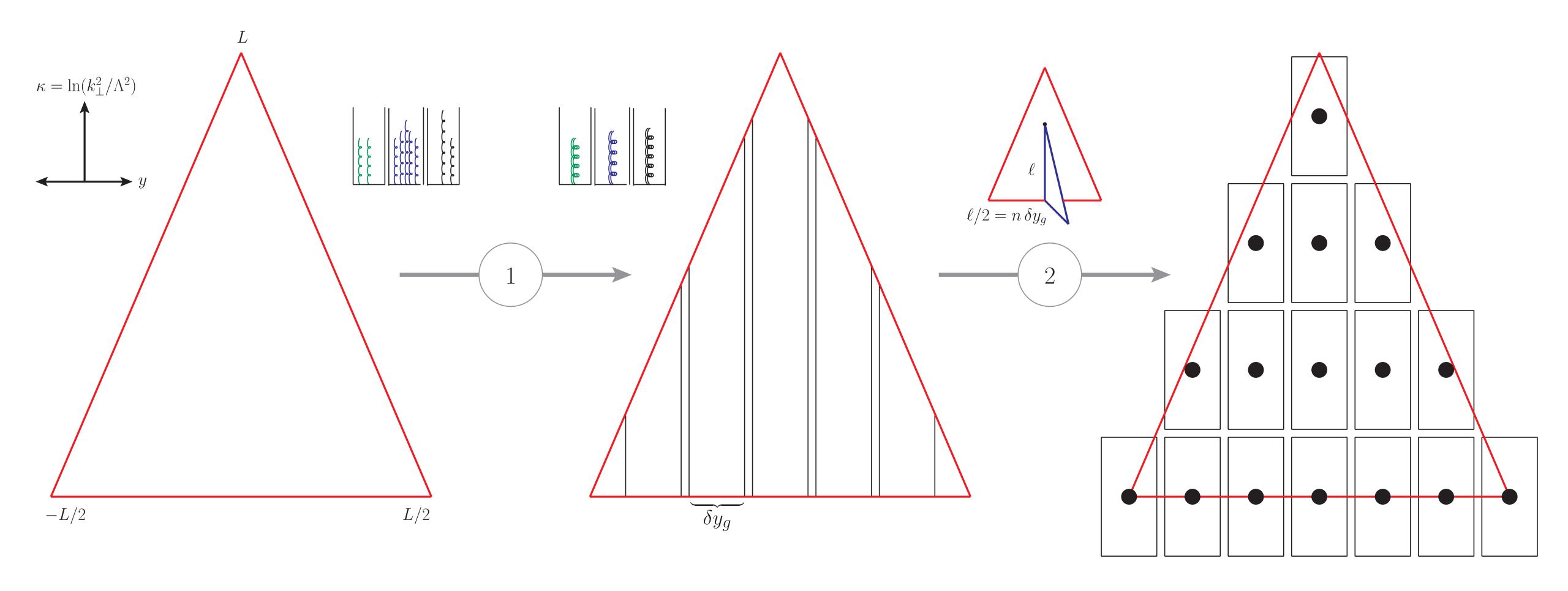


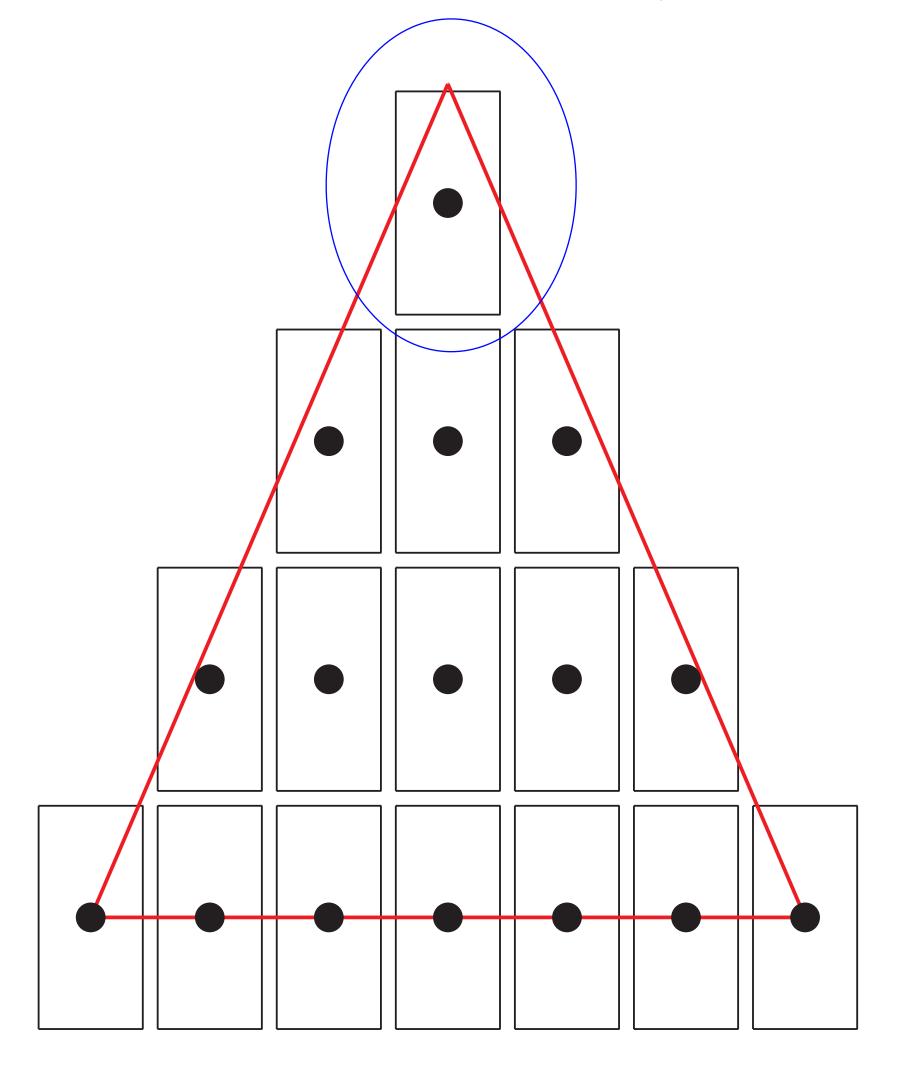


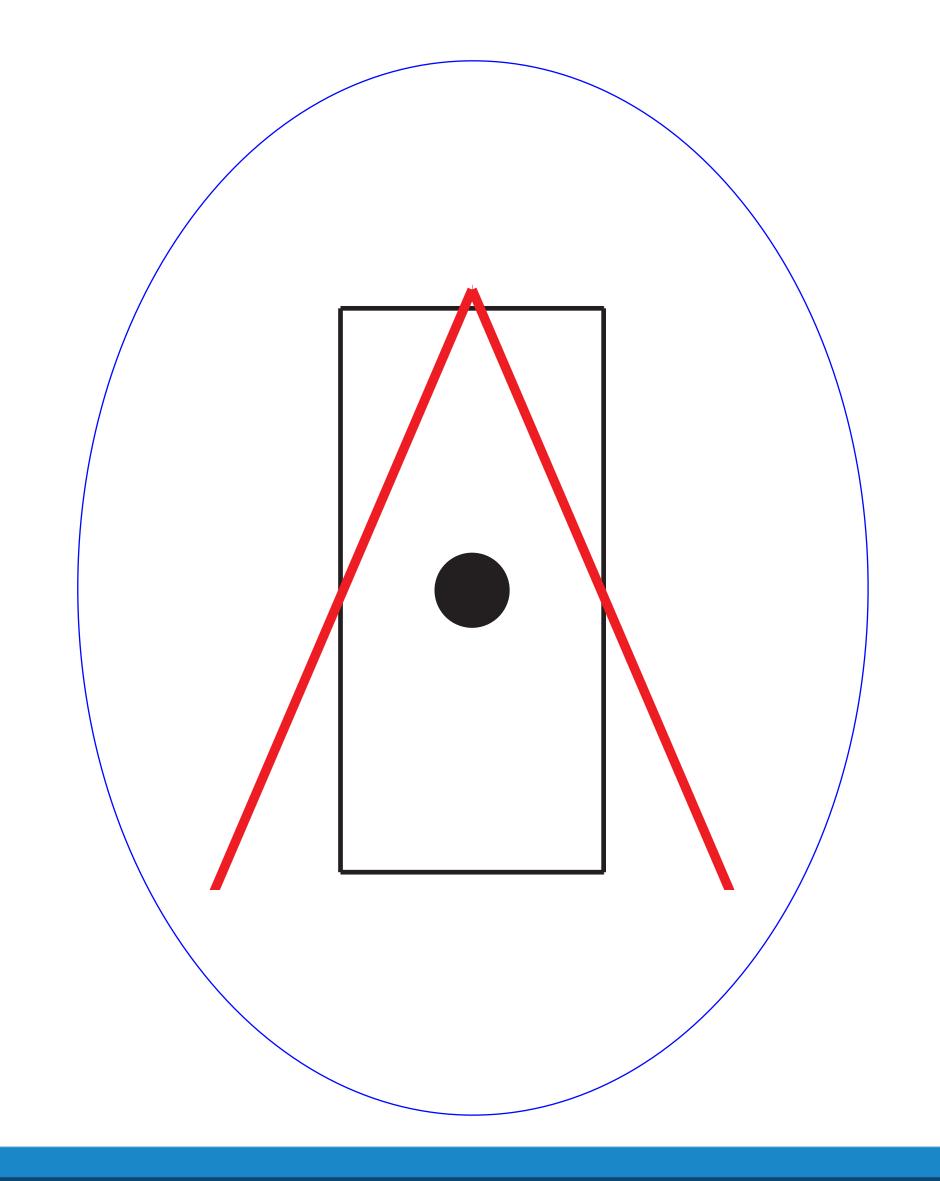


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Looking to the Future of Quantum Computers

