Isolated photon-hadron production in high energy *pp* and *pA* collisions at RHIC and LHC REF workshop 2022

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Based on Phys. Rev. D 105, 114052 (2022)

Introduction •000 Parton evolution at high energy

Motivation: The photon as a tool in pp and pA collisions

Results

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- $p + A \rightarrow \gamma + h$ as a probe of cold nuclear matter effects
- Complements hh production (see todays talk by F. Salazar)
- γh vs. hh as a probe:
 - 1 Better theoretical control
 - **2** Downside: smaller cross sections by α_e vs. α_s
- Isolated photons exclusion of fragmentation photons via isolation cone around the photon, $R = \sqrt{\Delta\phi^2 + \Delta\eta^2}$

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Introduction ○●○○ $pA
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Parton evolution at high energy

Gluon saturation



- At high energy, the parton density becomes large
- Gluon emission and recombination processes balance out
- Emergent saturation scale:

$$Q_s^2(x) \sim A^{1/3}/x^{0.3}$$

 See also talk by J.Jalilian-Marian on Tuesday

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Parton evolution at high energy

Phase space diagram of QCD



 Dependence of saturation on rapidity, transversal impulse and A of the target

Venugopalan, J.Phys.G 35 (2008) 104003.

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- - High energy eikonal scattering of the nucleus on partons
 - Representation of gluon shockwave through effective vertex
 - Gluon distributions are correlators of Wilson lines:

$$\begin{split} \tilde{U}(x_{\perp}) &= \mathcal{P} \exp\left[ig \int_{x^{+}} A_{a}^{-}(x)t_{a}\right] \\ \tilde{\mathcal{N}}_{\mathcal{A},Y_{\mathcal{A}}}(k_{\perp}) &= \frac{1}{N_{c}} \int_{\mathbf{y}_{\perp}} e^{ik_{\perp}y_{\perp}} \langle \tilde{U}(\mathbf{y}_{\perp})\tilde{U}^{\dagger}(0) \rangle_{x_{\mathcal{A}}} \end{split}$$

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Refinements on Compton scattering

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 cross section

$$\begin{aligned} \frac{d\sigma_{CGC}^{pA\to\gamma h}}{d^{2}\mathbf{k}_{\gamma\perp}d\eta_{\gamma}d^{2}\mathbf{P}_{h\perp}d\eta_{h}} &= (\pi R_{A}^{2})\sum_{q}\int_{0}^{1}\frac{dz_{h}}{z_{h}^{2}}D_{q}(z_{h},\mu^{2})\times \\ &\times \frac{e_{q}^{2}N_{c}}{8\pi^{4}}x_{p}f_{q}(x_{p},\mu^{2})k_{\perp}^{2}\tilde{\mathcal{N}}_{A,Y_{A}}(\mathbf{k}_{\perp})\hat{\sigma} \\ \hat{\sigma} &= \frac{\alpha_{e}}{2N_{c}}\frac{P_{q\gamma}}{q\cdot k_{\gamma}}\frac{z^{2}}{\mathbf{k}_{\gamma\perp}^{2}}, \quad P_{q\gamma} = \frac{1+(1-z)^{2}}{z}, \quad z = \frac{k_{\gamma}^{+}}{q^{+}+k_{\gamma}^{+}}. \end{aligned}$$

• Imbalance momentum: $\mathbf{k}_{\perp} = \mathbf{k}_{\gamma \perp} + \mathbf{q}_{\perp}$, $q = P_h/z_h \implies \mathbf{k}_{\perp}^2 \sim Q_s^2 \sim A^{1/3}$

(Gelis, Jalilian-Marian, Phys.Rev.D 66 014021 (2002))

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Target distributions

 MV^{γ} model - evolution of the following initial condition via the rcBK equation:

$$ilde{\mathcal{N}}_{Y_0}(\mathbf{x}_{\perp}) = \exp\left\{-rac{(x^2Q_{s0}^2)^{\gamma}}{4}\ln\left(rac{1}{x_{\perp}\Lambda_{IR}}+e
ight)
ight\}$$

- Ñ_{Y_A}(k_⊥) from AAMQS (Albacete, Armesto, Milhano, Quiroga-Arias, Salgado, EPJC 71, 1705 (2011))
- D_q(z_h, μ²) from DSS (de Florian, Sassot, Stratmann, PRD 75, 114010 (2007))
- x_pf_q(x_p, μ²) from CTEQ6M (Pumplin, Stump, Huston, Lai, Nadolsky, Tung, JHEP 07, 012 (2002))

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 $pA o \gamma h^{\pm}$ cross section $\circ \circ \circ \circ \circ$

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Refinements on Compton scattering

Sudakov soft gluon resummation

We want to account for soft gluon radiations in our CGC description of γh production - Sudakov resummation (Collins, Soper, Sterman, NPB 250 199-224 (1985)) (Mueller, Xiao, Yuan, PRL 110 082301 (2013)) (Stasto, Wei, Xiao, Yuan, PLB 784 301-306 (2018))



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Implementation of the Sudakov resummation

■ Effectively we are replacing the CGC, FF and PDF distributions with a b_{\perp} integral $(k_{\perp} \text{ convolution})$: $k_{\perp}^{2}\tilde{\mathcal{N}}_{A,Y_{A}}(k_{\perp})D_{q}(z_{h},\mu^{2})f_{q}(x_{p},\mu^{2})$ $\rightarrow \int_{\mathbf{b}_{\perp}} e^{ik_{\perp}\cdot b_{\perp}}\partial_{b\perp}^{2}\tilde{\mathcal{N}}_{A,Y_{A}}(b_{\perp})D_{q}(z_{h},\mu_{b}^{2})f_{q}(x_{p},\mu_{b}^{2})e^{-S_{Sud}(b_{\perp},Q)}$

• Sudakov factor (for
$$qg \rightarrow q\gamma$$
):
 $S_{\text{Sud}}(b_{\perp}, Q) = \int_{\mu_b^2}^{Q^2} \frac{d\bar{\mu}^2}{\bar{\mu}^2} \left[A \log \left(\frac{Q^2}{\bar{\mu}^2} \right) + B \right] + S_{non-pert}(b_{\perp}, Q),$
 $A = \frac{\alpha_s(\bar{\mu}^2)}{\pi} \left(C_F + C_A/2 \right), \quad B = -\frac{3\alpha_s(\bar{\mu}^2)}{2\pi} C_F$

• $\mu_b > 2e^-\gamma_E/b_{\text{max}}$, $S_{non-pert}(\mathbf{b}_{\perp}, Q)$ prescribed by (Sun, Isaacson, Yuan, Yuan, IJMPA 33 no. 11, 1841006 (2018))

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Results

Self-normalized angular correlations

Comparison of CGC vs CGC+Sudakov angular correlations



Generic CGC prediction: double peak structure at $\Delta \phi \sim \pi$

 Adding Sudakov effects seems to broaden the distribution and destroy that structure

(Benić, Garcia-Montero, AP, Phys. Rev. D 105, 114052 (2022))

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Self-normalized angular correlations



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Self-normalized angular correlations



PHENIX $pp \rightarrow \gamma h^{\pm}$, $\sqrt{s} = 510 \text{ GeV}$

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Self-normalized angular correlations

Predictions of nuclear effects at PHENIX

- pp vs pA calculation for lowest (5-7 GeV) and highest (12-15 GeV) $k_{\gamma\perp}$ bins
- Modest nuclear effect broadening of angular distribution
- Self normalized distribution good for comparison with experimental data, but part of the physical information is lost



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 $pA
ightarrow \gamma h^{\pm}$ cross section

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Self-normalized angular correlations

ALICE pp and pA angular correlations



Barely visible nuclear effect for this kinematics - we need lower $k_{\gamma\perp}$ resolution!

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Self-normalized angular correlations

Predictions of nuclear effects at ALICE

- pp vs pA calculation for two more favorable k_{γ⊥} bins (5-7 GeV and 9-12 GeV) as well as the existing one (12-40 GeV)
- Again, moderate nuclear effect visible in $\Delta \phi$ distribution broadening



Proxy for intrinsic k_{\perp}

 $\rho A \rightarrow \gamma h^{\pm}$ cross section

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Out-of-plane momentum distributions: PHENIX

$$p_{out} = P_{h\perp} \sin(\Delta \phi), \ \ x_E = -\frac{P_{h\perp}}{k_{\gamma\perp}} \cos(\Delta \phi)$$

Close to $\Delta \phi \sim \pi$ we have $p_{out} \sim z_h k_\perp$, and $x_E \sim z_h$



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Proxy for intrinsic k_{\perp}

*p*_{out} distributions: Gaussian widths

• We extract the widths of the previous curves by fitting to a Gaussian in the range $p_{out} < 1.1 \pm 0.2$ GeV

Results

Best description with CGC+Sudakov



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 p_{out} distributions: pp vs. pA predictions

 Difference between pA and pp Gaussian widths squared more pronounced at large x_E



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Summary and future outlooks

- Benchmark results for \(\gamma h\) production in the CGC+Sudakov framework
- Good description of RHIC and LHC data
- \blacksquare Predicted nuclear effects are $\sim 10\%$ within experimental resolution?
- Further (and ongoing) inquiries:
 - 1 Study of inclusive Drell-Yan production
 - **2** Testing of systematic errors
 - 3 Comparison with future data (e.g. LHCb)

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