

Precise predictions for the trilinear Higgs self-coupling in the Standard Model and beyond

predicting κ_λ in *any* model at the one-loop order

Henning Bahl, Johannes Braathen, **Martin Gabelmann**, Georg Weiglein
ECFA, 06.10.2022



The SM scalar sector

> V_{SM} fixed at tree-level by $m_h \approx 125 \text{ GeV}$ and $v = (\sqrt{2}G_F)^{-1/2} \approx 246 \text{ GeV}$:

$$\begin{aligned} V_{\text{SM}}(h) &= \frac{m_h^2}{2} h^2 + 3 \frac{m_h^2}{v} h^3 + 3 \frac{m_h^2}{v^2} h^4 \\ &= \frac{m_h^2}{2} h^2 + \lambda_{hhh} h^3 + \lambda_{hhhh} h^4 \end{aligned}$$

> However: λ_{hhh} (and λ_{hhhh}) experimentally unknown.

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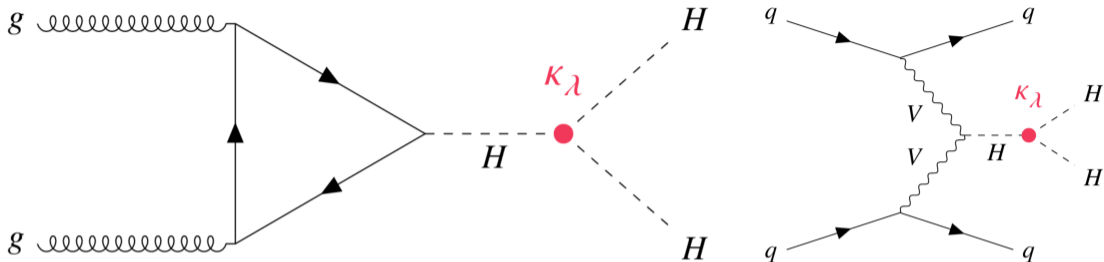
- > However: λ_{hhh} (and λ_{hhhh}) experimentally unknown.
> BSM case: deformation of the scalar potential possible!

$$V_{\text{BSM}}(h, \dots) = \frac{m_h^2}{2} h^2 + 3 \frac{m_h^2}{v} \kappa_\lambda^{\text{BSM}} h^3 + 3 \frac{m_h^2}{v^2} \kappa_{2\lambda} h^4 + \dots$$

$\kappa_\lambda^{\text{BSM}}$: describes deviation from SM: $\kappa_\lambda^{\text{BSM}} = \frac{\lambda_{hhh}^{\text{BSM}}}{\lambda_{hhh}^{\text{SM}}}$ in model "BSM".

- > $|\kappa_\lambda| \gg 1$ possible even though other (known) couplings behave very SM-like
> Example: "BSM"=THDM (talk by Johannes Braathen)
→ **higher-order corrections important**

Higgs pair production (in the alignment limit)



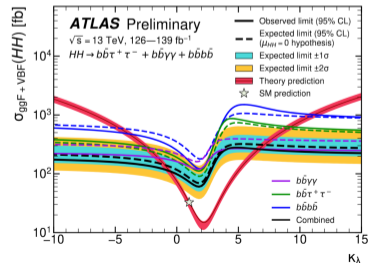
$pp \rightarrow hh: -0.4 < \kappa_\lambda < 6.3$ [ATLAS-CONF-2022-050]

When to apply the κ_λ -constraint to BSM models?

> only κ_λ is *significantly* modified by BSM physics

> all other couplings SM-like

→ in many models given by experimental constraints



Higher-order corrections to $\lambda_{hhh}^{\text{BSM}}$

A few studies of the trilinear Higgs self-couplings (i.e. κ_λ) already exists:

- > SM [Kanemura et al. '04]
- > SM + singlet [Kanemura et al. '16]
- > THDM [Kanemura et al. '04][Basler et al. '17][Braathen et al. '19]
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of which many can have sizeable deviations from $\kappa_\lambda = 1$ and hence also from $\sigma_{hh}^{\text{SM}}(\kappa_\lambda = 1)$ [Abouabid et al. '21]

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However, many more models to explore:

- > THDM-variants
- > singlet-variants
- > extended gauge sectors
- > extended fermion sectors
- > non-minimal SUSY, e.g. seesaw extensions, dirac gauginos, SplitSUSY, ...
- > (+ combinations)
- > ...

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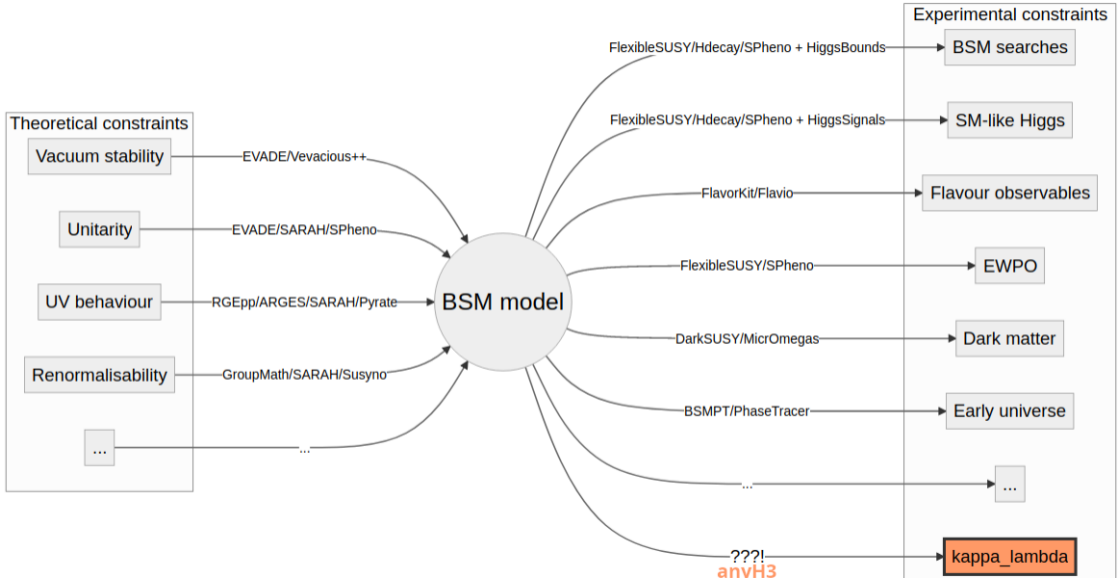
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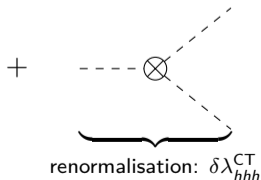
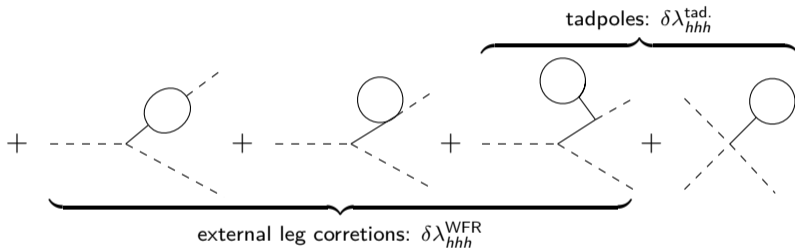
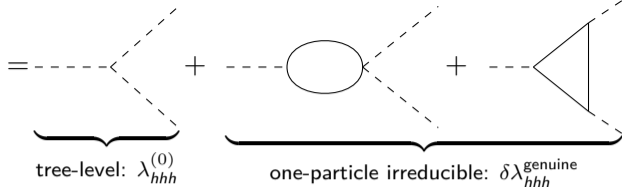
→ **framework to calculate $\lambda_{hhh}^{\text{BSM}}$ for large class of "BSM's"**



Ingredients

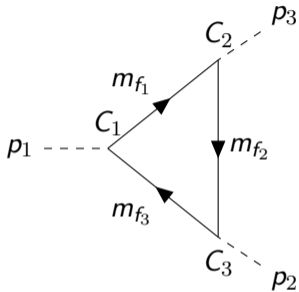
Ingredients

$(\lambda_{hhh}^{\text{BSM}})^{\text{one-loop}}$



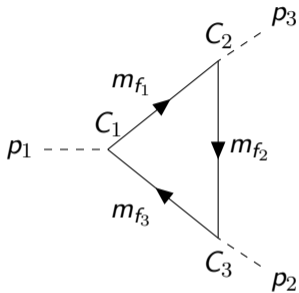
Example ingredient: generic fermion triangle

Idea: compute *generic* diagrams i.e. assume most generic



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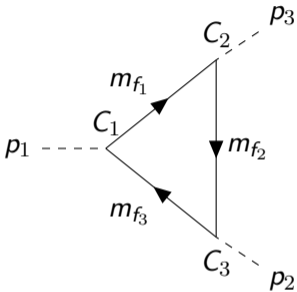
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- > couplings $C_i = P_L C_i^L + P_R C_i^R$, $P_{R/L} = \frac{1 \pm \gamma_5}{2}$
- > as well as loop-masses m_{f_i} and
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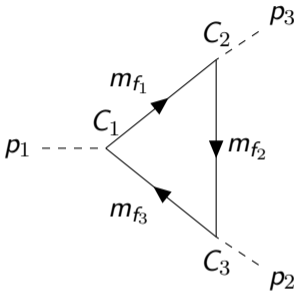


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$$\begin{aligned}
 &= 2\mathbf{B0}(p_3^2, m_2^2, m_3^2)(C_1^L(C_2^L C_3^R m_{f_1} + C_2^R C_3^R m_{f_2} + C_2^R C_3^L m_{f_3}) + C_1^R(C_2^R C_3^L m_{f_1} + \\
 &C_2^L C_3^L m_{f_2} + C_2^L C_3^R m_{f_3})) + m_{f_1} \mathbf{C0}(p_2^2, p_3^2, p_1^2, m_1^2, m_3^2, m_2^2)((C_1^L C_2^L C_3^R + \\
 &C_1^R C_2^R C_3^L)(p_1^2 + p_2^2 - p_3^2) + 2(C_1^L C_2^L C_3^L + C_1^R C_2^R C_3^R)m_{f_2} m_{f_3} + \\
 &2m_{f_1}(C_1^L(C_2^L C_3^R m_{f_1} + C_2^R C_3^R m_{f_2} + C_2^R C_3^L m_{f_3}) + C_1^R(C_2^R C_3^L m_{f_1} + C_2^L C_3^L m_{f_2} + \\
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 &(C_1^L C_2^R C_3^L + C_1^R C_2^L C_3^R)m_{f_3})) + \mathbf{C2}(p_2^2, p_3^2, p_1^2, m_1^2, m_3^2, m_2^2)((p_1^2 + p_2^2 - \\
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 \end{aligned}$$

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 \end{aligned}$$

- > insert concrete BSM model
- > evaluate with the help of COLLIER [\[Denner et al. '16\]](#)
- > public code anyH3 [\[Bahl, Braathen, MG, Weiglein tbp\]](#)

Another ingredient: Renormalisation of λ_{hhh}

$$\delta\lambda_{hhh}^{\text{CT}} = \text{---} \bigcirc \text{---} = ?$$

- > one-loop \rightarrow renormalisation of all parameters entering $\lambda_{hhh}^{(0),\text{BSM}}$
- > reminder: $\lambda_{hhh}^{(0),\text{SM}} = 3 \frac{m_h^2}{v^2}$
- > \rightarrow generalization (reminder: quasi-alignment!):

$$\lambda_{hhh}^{(0),\text{BSM}} \equiv \lambda_{hhh}^{(0),\text{BSM}} \left(\underbrace{v^{\text{SM}}, m_h^{\text{SM}}}_{\text{SM Higgs sector}}, m_{H_i}^{\text{BSM}}, v_i^{\text{BSM}}, \dots \right)$$

\swarrow
 \swarrow

BSM masses
vevs

couplings etc. (that can't be expressed in terms of masses)

- > user's choice:
 - SM sector: fully OS **or** $\overline{\text{MS}}/\overline{\text{DR}}$
 - BSM masses: OS **or** $\overline{\text{MS}}/\overline{\text{DR}}$
 - additional couplings/vevs: most likely $\overline{\text{MS}}$ **but also custom ren. conditions possible!**

$$\delta\lambda_{hhh}^{\text{CT}} = \sum_x \frac{\partial \lambda_{hhh}^{(0),\text{BSM}}}{\partial x} \delta x, \quad x = (m_h^{\text{SM}})^{\text{OS}/\overline{\text{MS}}}, (v^{\text{SM}})^{\text{OS}/\overline{\text{MS}}}, (m_{H_i}^{\text{BSM}})^{\text{OS}/\overline{\text{MS}}}, (\dots)^{\overline{\text{MS}}/\text{custom}}$$

(Default) Renormalization choice of $(v^{\text{SM}})^{\text{OS}}$ and $(m_i^2)^{\text{OS}}$

> $v^{\text{OS}} \equiv \frac{2M_W^{\text{OS}}}{e} \sqrt{1 - \frac{M_W^{2\text{OS}}}{M_Z^{2\text{OS}}}}$ with

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All bosonic one- & two-point functions and their derivatives for general QFTs are required for flexible OS renormalisation.

Feature list (so far) of anyH3

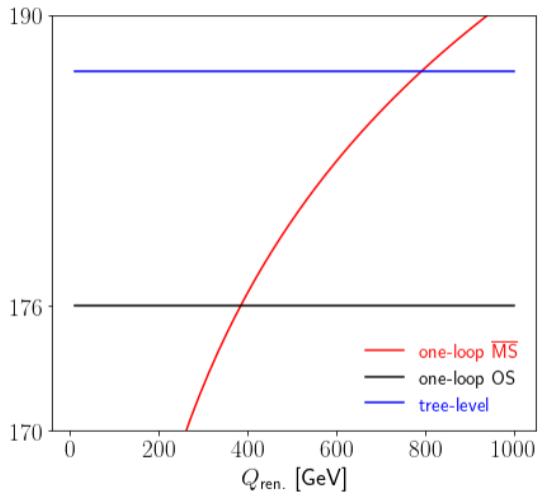
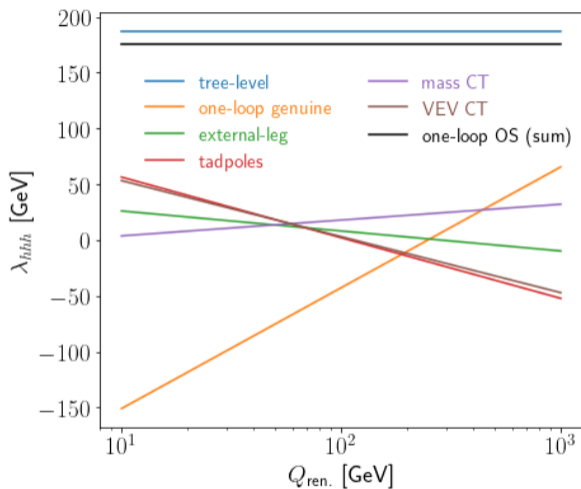
- > import/convert arbitrary UFO models
- > definition of renormalisation schemes

```
# schemes.yml
renormalization_schemes:
  OS:
    mass_counterterms:
      h1: OS
      h2: OS
    VEV_counterterm: OS
  MS:
    mass_counterterms:
      h1: MS
      h2: MS
    VEV_counterterm: MS
```

- > optional: full p^2 dependence
 - > numerical / analytical / \LaTeX outputs
 - > restrict to certain topologies
 - > restrict to certain particles in the loop
 - > python-library with command-line- and Mathematica-interface
- ```
from anyBSM import anyH3
myfancymodel = anyH3(
 'path/to/UFO/model',
 scheme = 'OS')
result = myfancymodel.lambdahhh()
```
- > ...

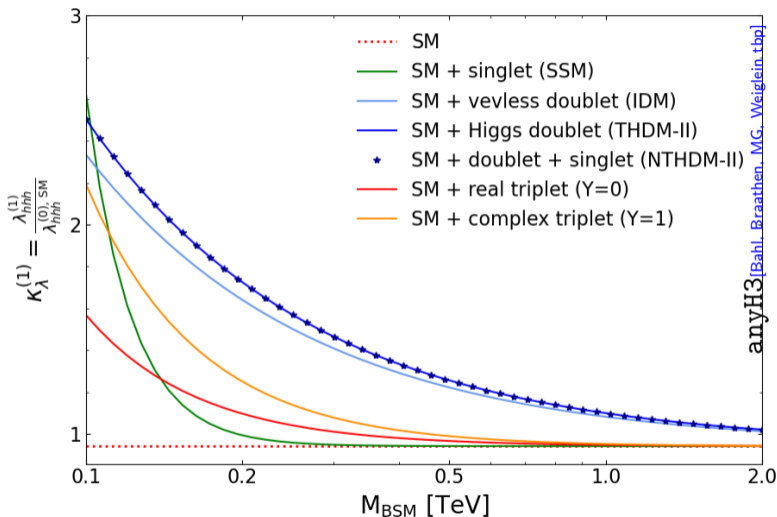
# Results

# Simplest model to consider for cross-check: SM



Leading two-loop  $\mathcal{O}(\alpha_t(\alpha_t + \alpha_s))^{\text{OS}}$ :  $\mathcal{O}(+1.4\%)$  [Senaha '18] [Braathen et al. '19]

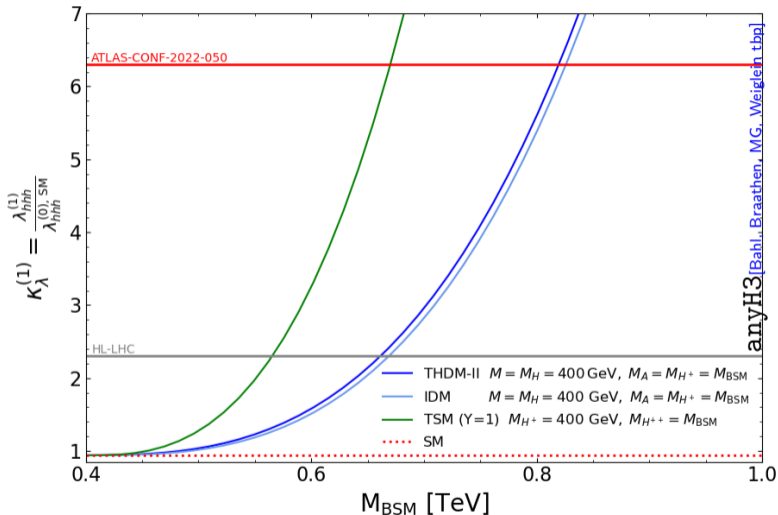
# Decoupling in the alignment limit



- > alignment means  $\kappa_\lambda^{\text{tree-level}} = 1$
- > recover SM result for  $M_{\text{BSM}} \rightarrow \infty$
- > many models built-in and cross-checked
- > easy to implement new models (UFO)



# Non-decoupling in the alignment limit



- > mass splitting within the same multiplet
- > induces large couplings for  $M_{BSM} \rightarrow \infty$
- > corrections large-enough to exclude otherwise unconstrained parameter space
- > (see also talk by Johannes Braathen before)

## Summary

- > developed computer code anyH3 (anyBSM) for  $\lambda_{hhh}$  in arbitrary ren. QFTs
  - at the full one-loop order
  - with arbitrary choice of renormalization schemes
- > uses UFO input (generate with SARAH, FeynRules or use a custom one)
- > analytical results; fast numerical results; SM:  $\mathcal{O}(0.2\text{ s})$ , MSSM:  $\mathcal{O}(0.5\text{ s})$
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Future todos:

- > publish
- > more models / cross-checks
- > go beyond one-loop
- > non-SM self-couplings (e.g.  $\kappa_{\lambda_{Hhh}}$ )
- >  $\kappa_t$  and  $\kappa_{tt}$

# Backup

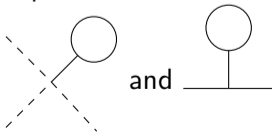
## Tadpole contributions to $\lambda_{hhh}$

- > In the SM: once  $\lambda_{hhh}$  is expressed in terms of *physical* input parameters, its result is independent of the treatment of the tadpoles:

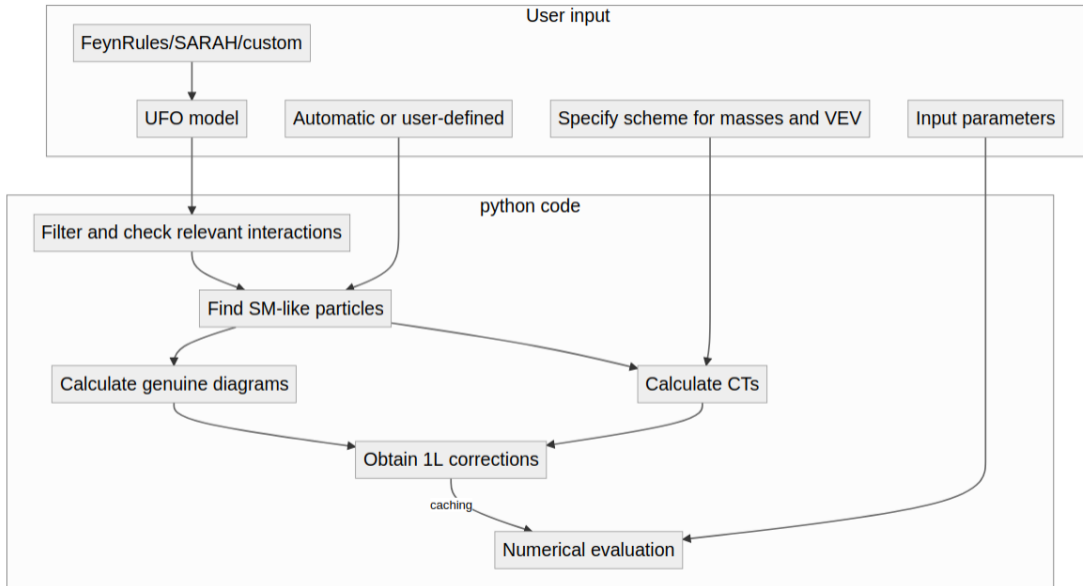
$$\delta^{(1)}\lambda_{hhh} \supset -\frac{3}{v^2}\delta^{(1)}t_{\text{finite}}$$

- > However: UFO models do (often) **not** contain the explicit dependence on the tree-level/one-loop/one-loop-CT tadpoles.
- > Thus: we choose the Fleischer-Jegerlehner treatment  $t^{\text{tree-level}} = 0$  and renormalize  $\delta^{(1)}t^{\text{CT}}|_{\text{finite}} = 0$  in the  $\overline{\text{MS}}$  scheme per default (can also turn-off automatic tadpoles and implement own scheme).

→ only need to take into account tadpole contributions to all two- and three-point functions:



# Workflow



# W Mass

- > start with HO corrections to muon decay:  $M_W^2 \left(1 - \frac{M_W^2}{M_Z^2}\right) = \frac{\pi\alpha_{em}}{\sqrt{2}G_F} [1 + \Delta r]$
- > and solve for:  $M_W^2 = M_Z^2 \left[ \frac{1}{2} + \sqrt{\frac{1}{4} - \frac{\pi\alpha_{em}}{\sqrt{2}G_F M_Z^2} (1 + \Delta r)} \right]$
- > with:  $\Delta r^{(1)} = 2\delta^{(1)} e + \frac{\Pi_W^{(1),T}(0) - \delta^{(1)} M_W^2}{M_W^2} - \frac{\delta^{(1)} \sin^2 \theta_W}{\sin^2 \theta_W} + \delta_{\text{vertex+box}}$
- > and:  $\frac{\delta^{(1)} \sin^2 \theta_W}{\sin^2 \theta_W} = \frac{\cos^2 \theta_W}{\sin^2 \theta_W} \left( \frac{\Pi_W^{(1),T}(M_W^2)}{M_W^2} - \frac{\Pi_W^{(1),T}(M_W^2)}{M_W^2} \right)$

It's all there but:

- >  $\delta_{\text{vertex+box}}^{\text{SM}} = -\frac{2 \text{sign}(\sin \theta_W)}{\cos \theta_W \sin \theta_W M_Z^2} \Pi_{Z\gamma}(p^2 = 0) + \frac{\alpha_{QED}}{4\pi \sin^2 \theta_W} \left( 6 + \frac{7-4 \sin^2 \theta_W}{2 \sin^2 \theta_W} \right) \log(\cos^2 \theta_W)$
- >  $\delta_{\text{vertex+box}}^{\text{BSM}} = \text{needs to be implemented}$

However:

- > in many models  $\Delta r \supset \frac{\delta \sin^2 \theta_W}{\sin \theta_W} \approx \delta\rho$  is the dominant effect!