

# Quantum Field Theory Concepts

Hitoshi Murayama (Berkeley, Kavli IPMU)  
DESY Theory Workshop, September 30, 2022



# HIGGS, FLAVOR AND BELOW

DESY Theory Workshop

27 – 30 September 2022 at DESY Hamburg, Germany

## Plenary Talks

**P. Agrawal** (Oxford U.)

**A. El-Khadra** (UIUC)

**A. Greljo** (Bern U.)

**S. Hansmann-Menzemer** (U. Heidelberg)

**G. Heinrich** (KIT)

**S. Höche** (Fermilab)

**M. Kado** (La Sapienza)

**F. Kahlhoefer** (KIT)

**G. Landsberg** (Brown U.)

**J. Lindert** (U. of Sussex)

**M. Mangano** (CERN)

**P. Monni** (CERN)

**H. Murayama** (IPMU & Berkeley)

**F. Riva** (Geneva U.)

**Y. Shadmi** (Technion)

**D. Shih** (Rutgers U.)

**A. Vladimirov** (UC Madrid)

**K. Vos** (U. of Maastricht)

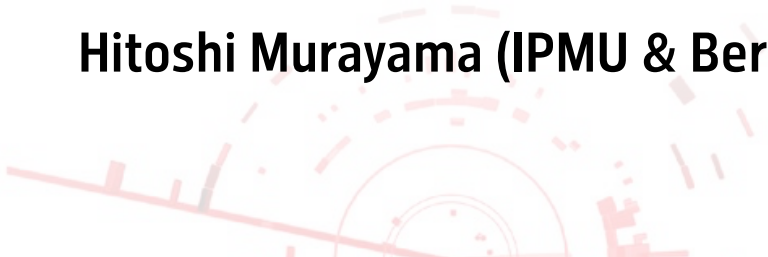
**G. Weiglein** (DESY)

**M. Worek** (RWTH Aachen)

## DESY Heinrich-Hertz-Lecture

Hitoshi Murayama (IPMU & Berkeley)

## 28 September 2022





# references

vector-like  
chiral

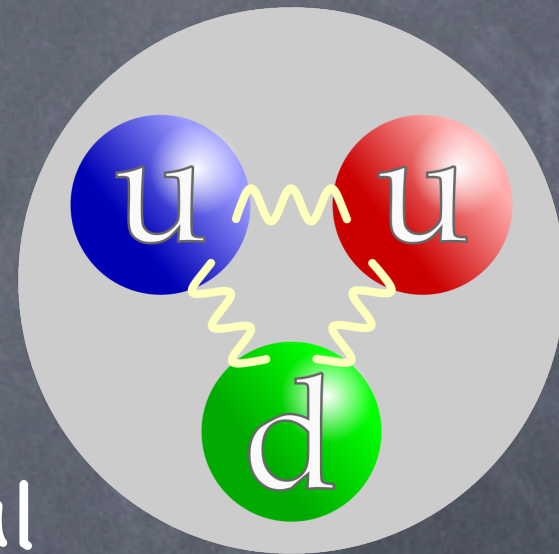
condensed matter

- HM, 2104.01179, PRL
- Csáki, HM, Telem, 2104.10171, PRD
- Csáki, HM, Telem, 2105.03444, PRD
- Csáki, Gomes, HM, Telem, 2106.10288, PRL
- Csáki, Gomes, HM, Telem, 2107.02813, PRD
- HM, Noether, Varier, 2111.09690
- Leedom, HM, Suter, in preparation
- Csáki, Gomes, HM, Noether, Telem, Varier, in preparation
- Kondo, HM, Sylber, 2209.09287
- HM and Wang, in preparation



# I felt cheated by QCD

- When we first learn about quarks, we get told
  - Here, look at this proton.
  - There are actually colorful, beautiful quarks inside.
  - But you can never see them, you can never take them out.
  - Nonetheless, believe, they are in there!
  - Who would buy that? Internet Scam?





Dear friend,

I am Andre Ouedraogo, a banker by profession from Burkina Faso in West Africa and currently holding the post of Director Auditing and Accounting unit of the bank. It's my urgent need for a foreign partner that made me to contact you for this business. I have the opportunity of transferring the left over funds (\$11.5 million) of one of my bank clients who died along with his entire family on 31 July 2000 in a plane crash. You can confirm the genuineness of the deceased death by clicking on this website.  
<http://news.bbc.co.uk/1/hi/world/europe/859479.stm>

I need a foreign partner who will support me because i can not claim this money alone without a foreign partner since the deceased client (the owner of the fund) was a foreigner.

This fund (\$11.5 million) will be shared between us in the ratio of 60/40. I agreed that 40% of this money will be for you as a respect to the provision of a foreign account while 60% will be for me and I want to assure you that this transaction is absolutely legal and risk free since i work in this bank and i have all the necessary information that might be needed. Before we proceed, i would like to know your ability to handle this over there in your country.

Please tell me more about the political/economic stability/monetary policy of your country. I need to know all these because i don't want to have problem with the Government of your country.

Kindly update me with the following information because i want to know you more before we proceed on this transaction. Hope you will understand the importance of this request.

- 1. Your full name.....
- 2. Your age/sex .....
- 3. your occupation .....
- 4. Your residential address .....
- 5. Your nationality .....
- 6. Your private phone number .....
- 7. Your fax number .....

I will be waiting for your response.

Thanks for your understanding.

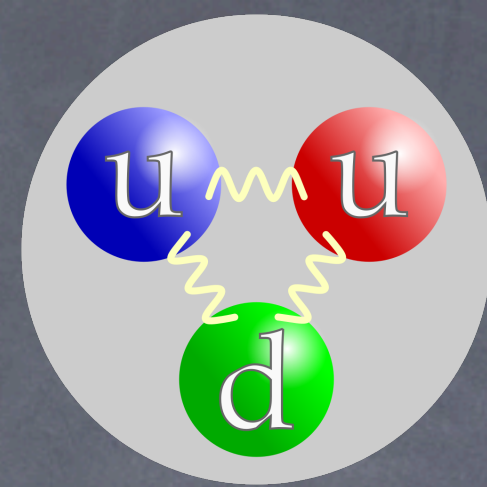
Have a great day.

Yours.

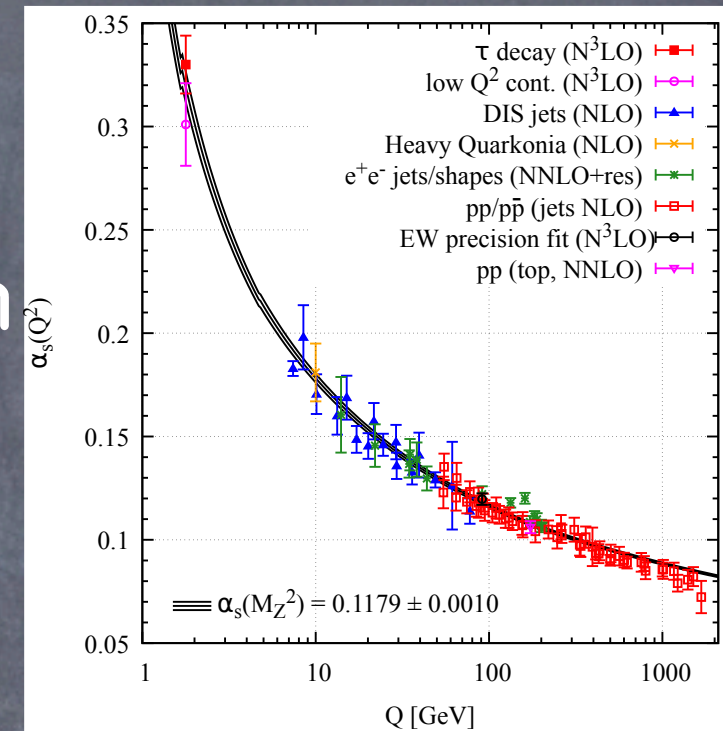
Andre Ouedraogo



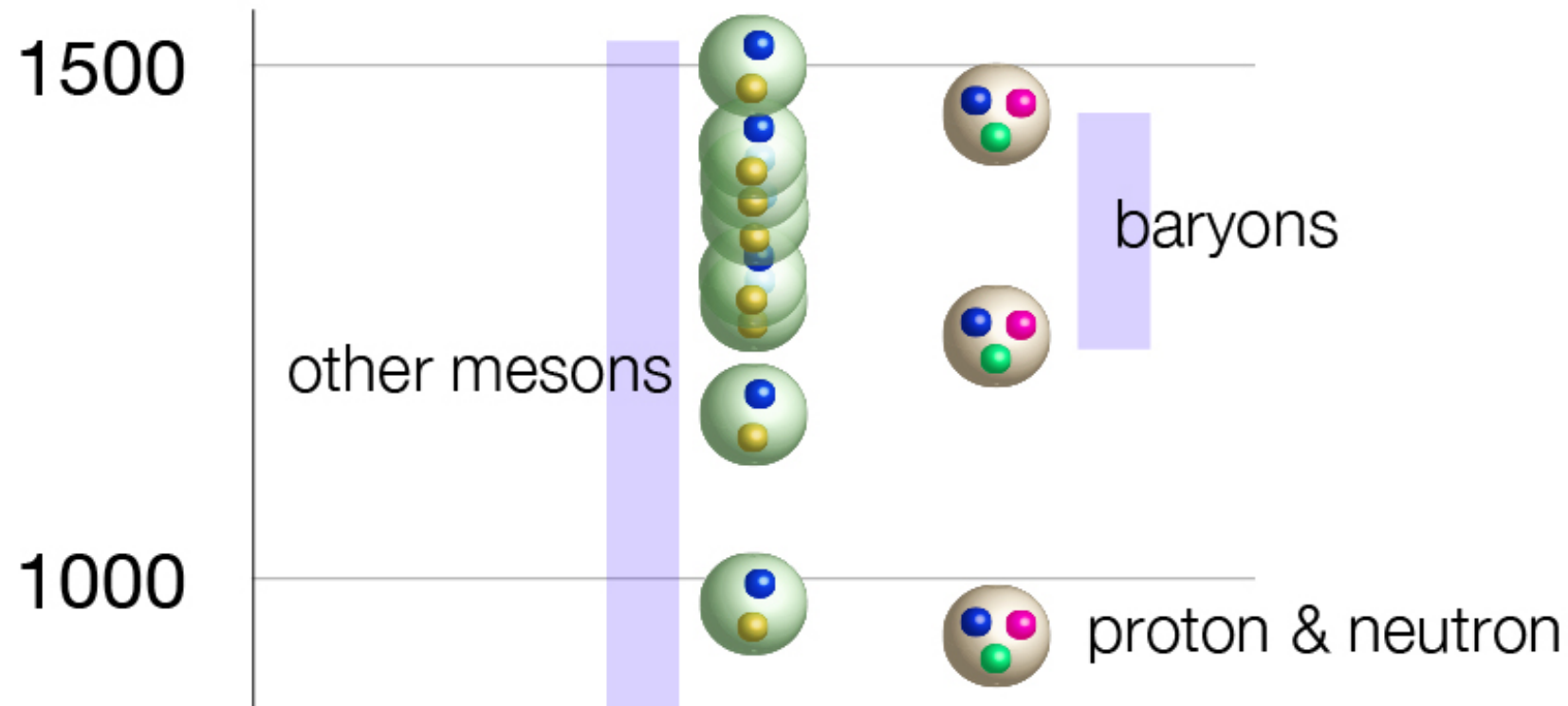
# Can we solve QCD?



- When we first learn about quarks, we get told we can never see them
  - Internet Scam?
  - Confinement!
  - $\beta(g^2) < 0$  and asymptotic freedom
  - only suggestive, doesn't prove confinement
- Another puzzle: proton and pion are made of the same quarks
  - why pion  $\approx$  massless  $\ll$  proton?
- very mysterious!
- Very little known about chiral gauge theories







$\hbar$

$$\frac{\hbar}{m_{\pi}c}$$

electric force

strong force

strong force

electric force

electric force

strong force

strong force

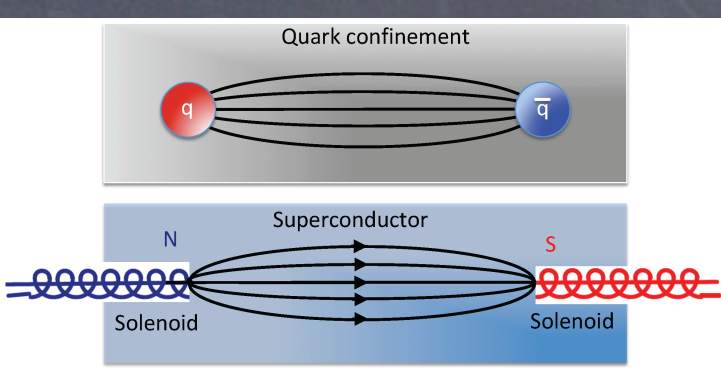
electric force

range  $\approx$

If pions are heavy

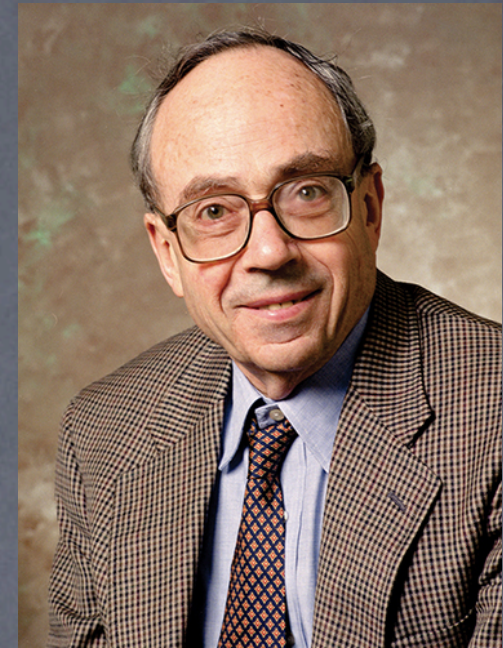
With real light-weight pions





# Feeling better

- 👁 Qualitative picture makes us feel better
- 👁 **Confinement**
  - 👁 dual Meißner effect (Mandelstam)
  - 👁 **assume** monopole condensation
  - 👁 quarks confined by electric flux tube
- 👁 **Chiral symmetry breaking** (Nambu)
  - 👁 massless QCD invariant under  $SU(N_f)_L \times SU(N_f)_R \times U(1)_B$
  - 👁 **assume** broken to  $SU(N_f)_V \times U(1)_B$
  - 👁 pion = Nambu-Goldstone boson = massless
- 👁 Are there really monopoles?
- 👁 How do we know  $\langle \bar{q}q \rangle \neq 0$ ?

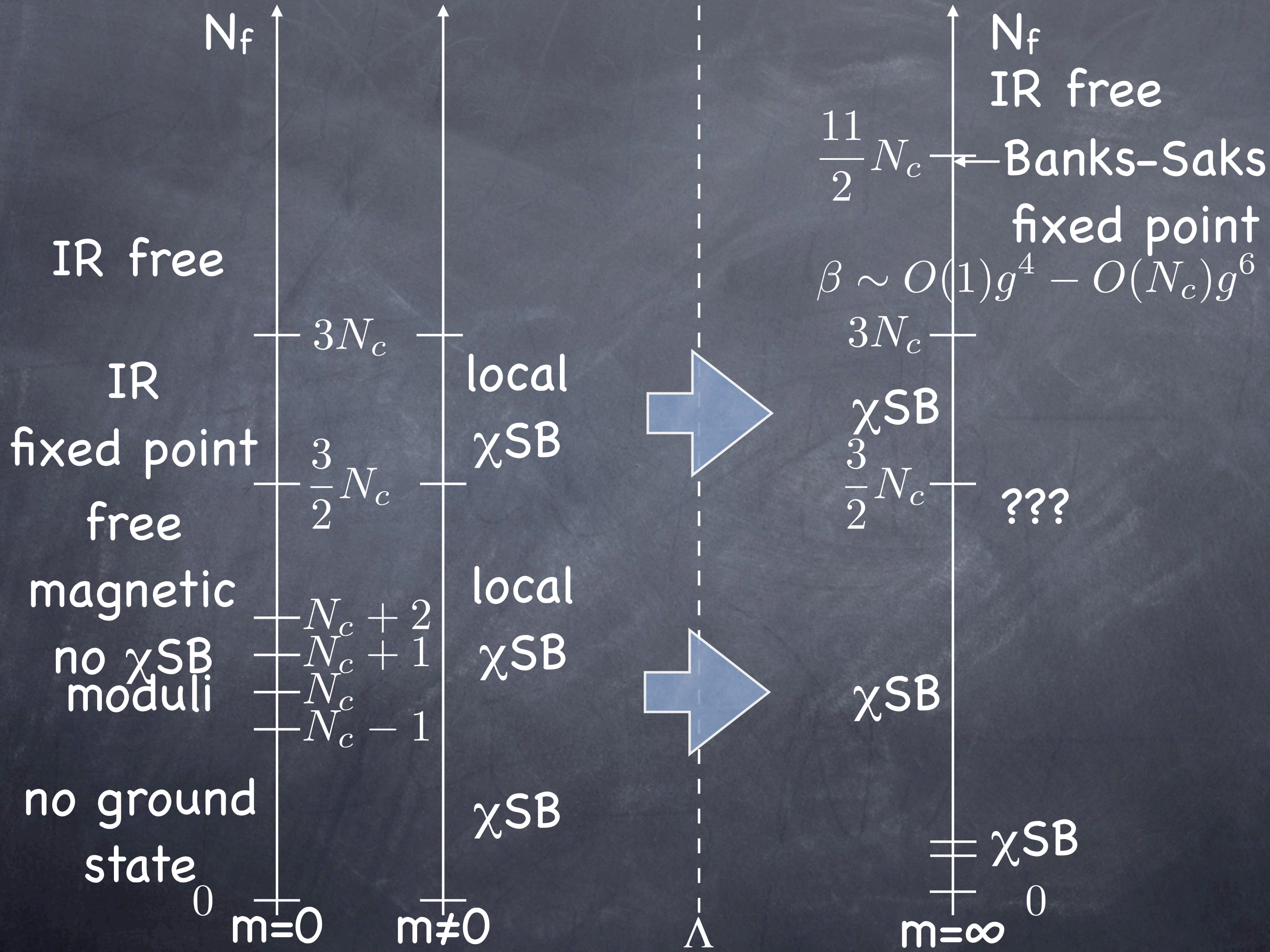




# Feeling even better but not there yet

- Progress in understanding QCD
- Confinement** (Seiberg-Witten)
  - N=2 SYM has Coulomb branch  $u = \text{Tr} \Phi^2$
  - singularities = massless monopole/dyon
  - N=1 perturbation  $W = \mu u - (u - \Lambda^2) M^+ M^-$
  - $M^+ = M^- = \sqrt{\mu} \neq 0$ : monopole condensation!
  - can further perturb to N=0 with  $m_\lambda \neq 0$
- Chiral symmetry breaking**
  - N=2 doesn't have  $\chi S$   $W = \sqrt{2} \tilde{Q}_i \Phi Q^i$
  - N=1 (Seiberg) has too unusual phases







# Main message

- Supersymmetric QCD is “solved” exactly by Seiberg in the 90s
  - but far removed from the real world
- Adding small SUSY breaking via anomaly mediation still allows for exact solution
- derive non-perturbative behavior analytically
  - chiral symmetry breaking
  - monopole condensation and confinement
- can solve chiral gauge theories exactly, too!
- Sometimes phase transitions, but local minima are still useful



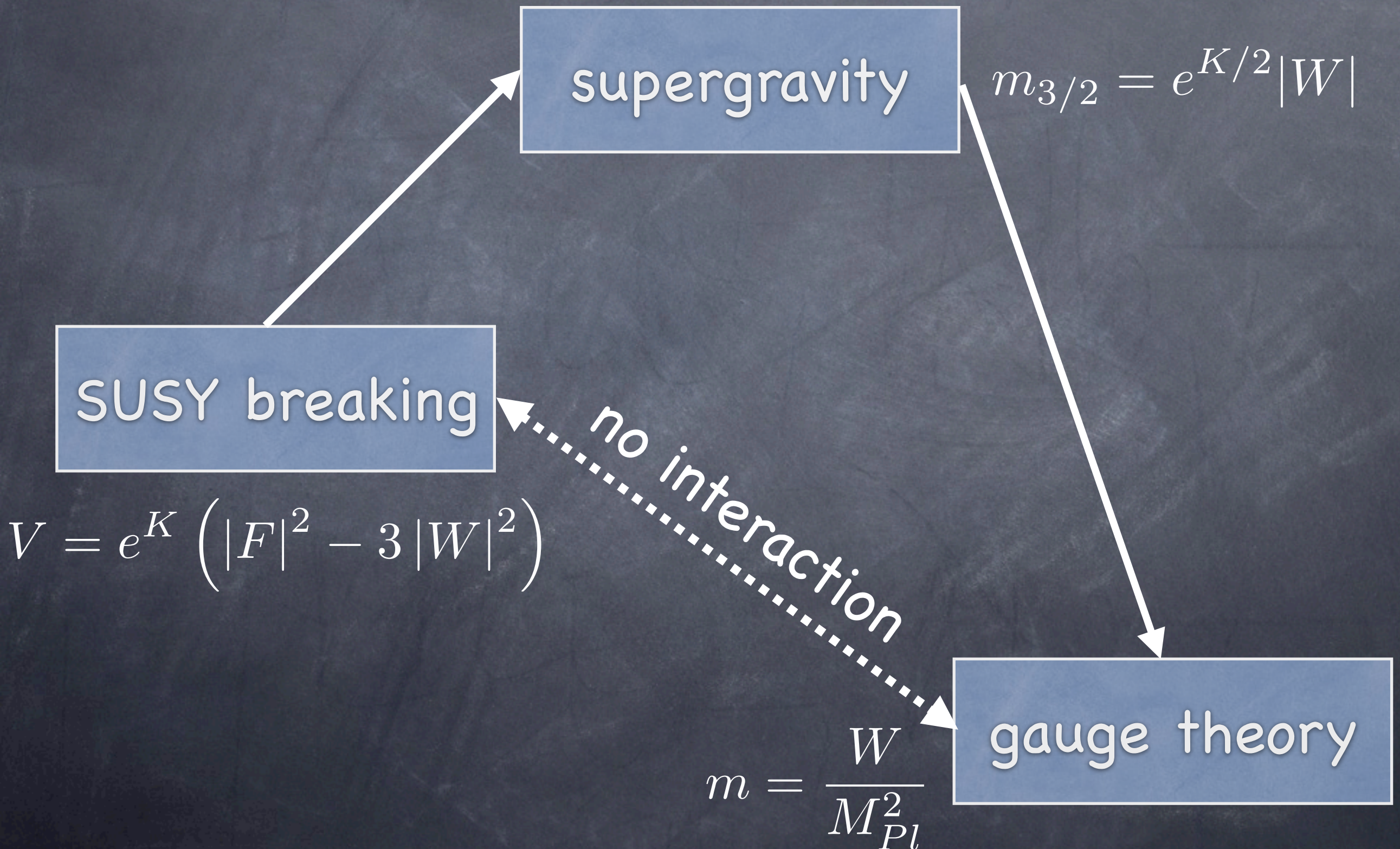
# Our Needs

- We'd like to connect N=1 SUSY results by Seiberg to non-SUSY gauge theories
  - decouple gauginos and squarks!
  - SUSY breaking  $m_\lambda$  and  $m_{\tilde{Q}}$
- But we need to deal with composites such as mesons and baryons
- SUSY breaking effects on composites may be non-trivial
- Anomaly mediation!

Randall, Sundrum; Giudice, Luty, HM, Rattazzi (1998)



# Sequestering





# AMSB Summary

- Tree-level piece on dimensionful parameters

$$V_{\text{AMSB}} = -m \left( \phi \frac{\partial W}{\partial \phi} - 3W \right)$$

- loop-level piece from running

$$M_i = -\frac{\beta_i(g^2)}{2g_i^2} m_{3/2}, \quad m_i^2 = -\frac{\dot{\gamma}_i}{4} m_{3/2}^2, \quad A_{ijk} = -\frac{1}{2}(\gamma_i + \gamma_j + \gamma_k) m_{3/2}$$

- determined only by physics at the energy scale of interest
- UV insensitivity!

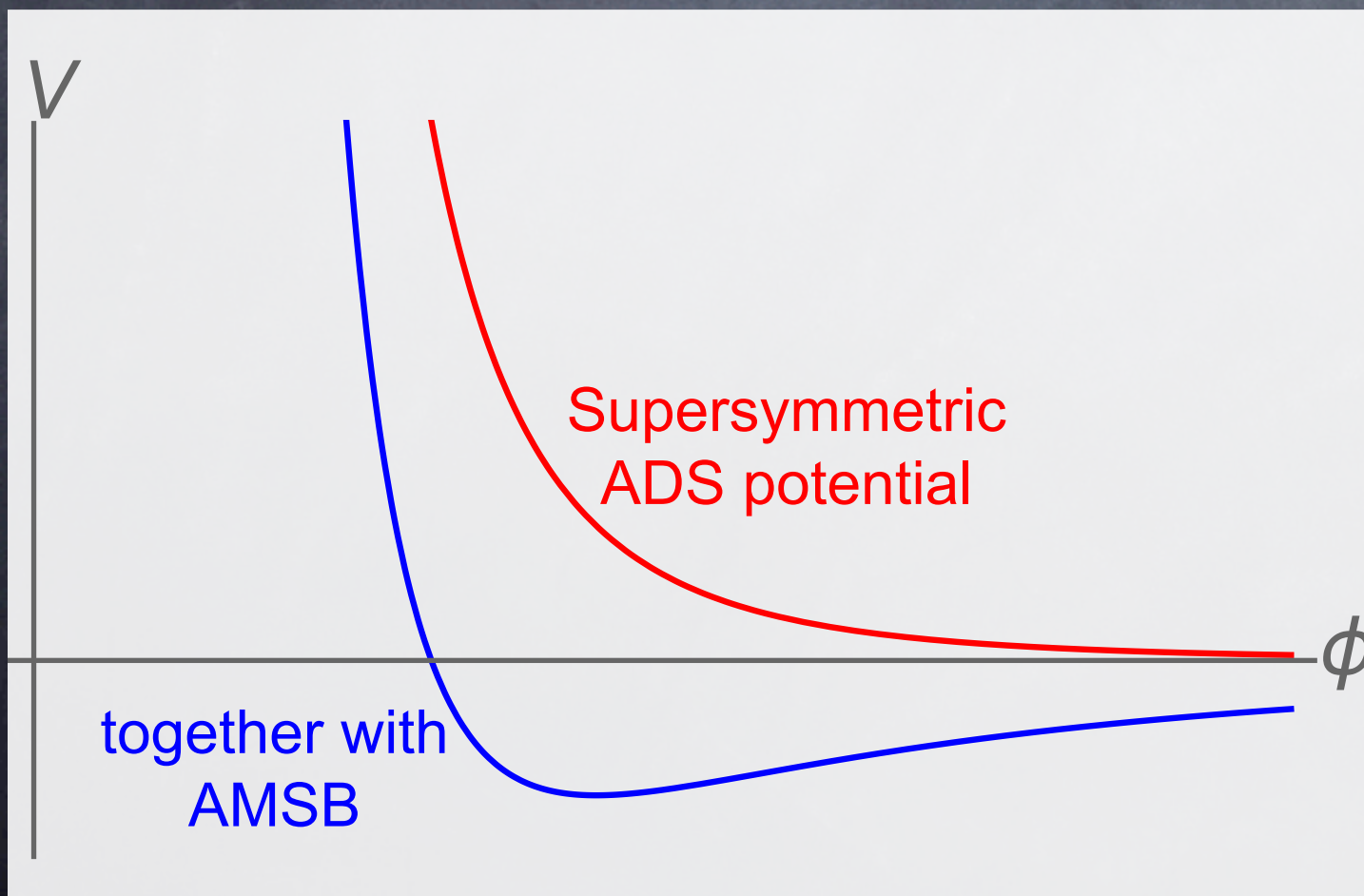


$$N_f < N_c$$

run-away superpotential for  $M^{ij} = \tilde{Q}^i Q^j$

$$W = (N_c - N_f) \left( \frac{\Lambda^{3N_c - N_f}}{\det M} \right)^{1/(N_c - N_f)} \quad M^{ij} = \delta^{ij} \phi^2$$

$$V = \left| 2N_f \frac{1}{\phi} \left( \frac{\Lambda^{3N_c - N_f}}{\phi^{2N_f}} \right)^{1/(N_c - N_f)} \right|^2 - (3N_c - N_f)m \left( \frac{\Lambda^{3N_c - N_f}}{\phi^{2N_f}} \right)^{1/(N_c - N_f)} + c.c.$$



$$M_{ij} = \Lambda^2 \left( \frac{4N_f(N_c + N_f)}{3N_c - N_f} \frac{\Lambda}{m} \right)^{(N_c - N_f)/N_c} \delta_{ij}$$

$$SU(N_f)_L \times SU(N_f)_R \rightarrow SU(N_f)_V$$

$\chi$ SB! Proving Nambu  
mesino loop  $\rightarrow$  WZW term

$N_f=1$  special  
no NGB, gapped



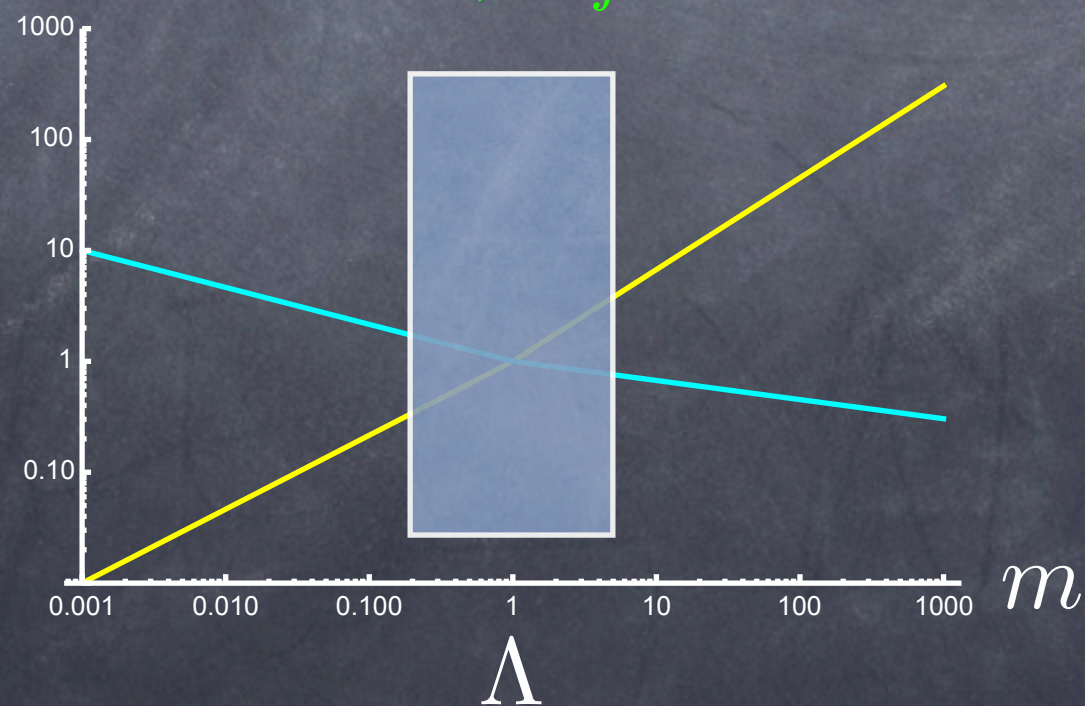
# fermion bilinear

$$M^{ij} = \tilde{q}_L^{i*} \tilde{q}_R^j + \theta^2 \bar{q}_L^i q_R^j$$

$$\tilde{q}_R^{i*} \tilde{q}_L^j \sim \left( m^{N_f - N_c} \Lambda^{3N_c - N_f} \right)^{1/N_c}$$

$$\bar{q}_R^i q_L^j \sim m \left( m^{N_f - N_c} \Lambda^{3N_c - N_f} \right)^{1/N_c}$$

$$N_c = 3, N_f = 2$$

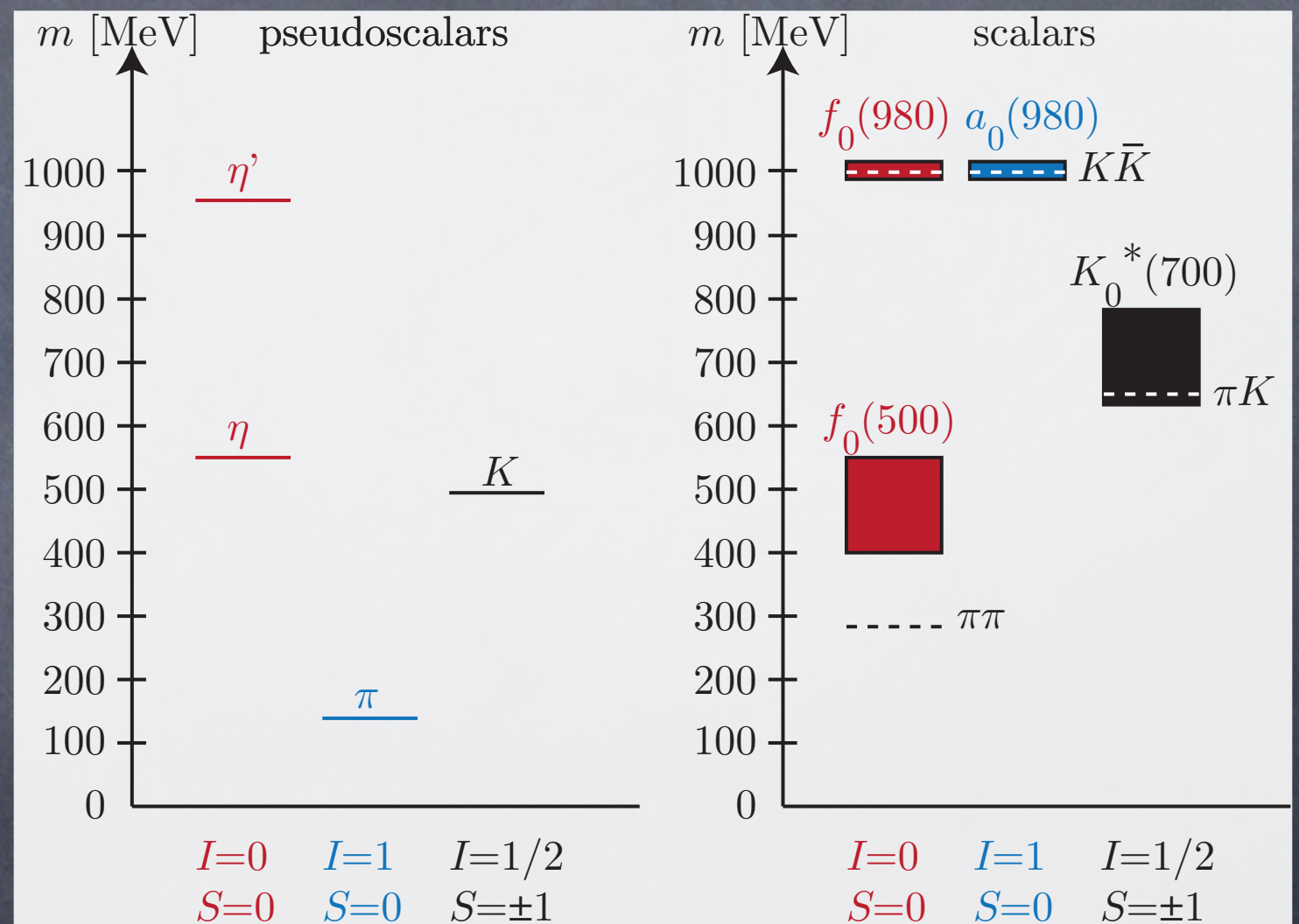




# Light spectrum



— fermions





# confinement vs screening

- We've derived  $\chi$ SB in  $SU(N_c)$  QCD
  - it has no confinement
  - massless quarks in the fundamental rep can screen any color charges
  - Wilson loop is perimeter law
- $SO(N_c)$  QCD with quarks in vector rep
  - cannot screen  $Z_2$  center (e.g. spinor rep)
  - rigorous definition of confinement
  - can we see an interplay with  $\chi$ SB?



$$N_f = N_c - 2$$

$$Q = \begin{pmatrix} v & \dots & 0 \\ \vdots & \ddots & \vdots \\ 0 & \dots & v \\ 0 & \dots & 0 \\ 0 & \dots & 0 \end{pmatrix}$$

• for  $M_{ij} = Q_i Q_j \neq 0$  with rank  $M = N_f$ ,  
 $SO(N_c)$  is broken to  $SO(2)$

• 't Hooft Polyakov monopoles!

• Coulomb branch  $u = \det M$

• two singularities

•  $u = \det M = 0$

• dyons:  $q_i^\pm$

$$V \approx - \left( \frac{\lambda^2}{16\pi^2} \right)^4 m^4$$

$$W = \frac{1}{\mu} M^{ij} q_i^+ q_j^-$$

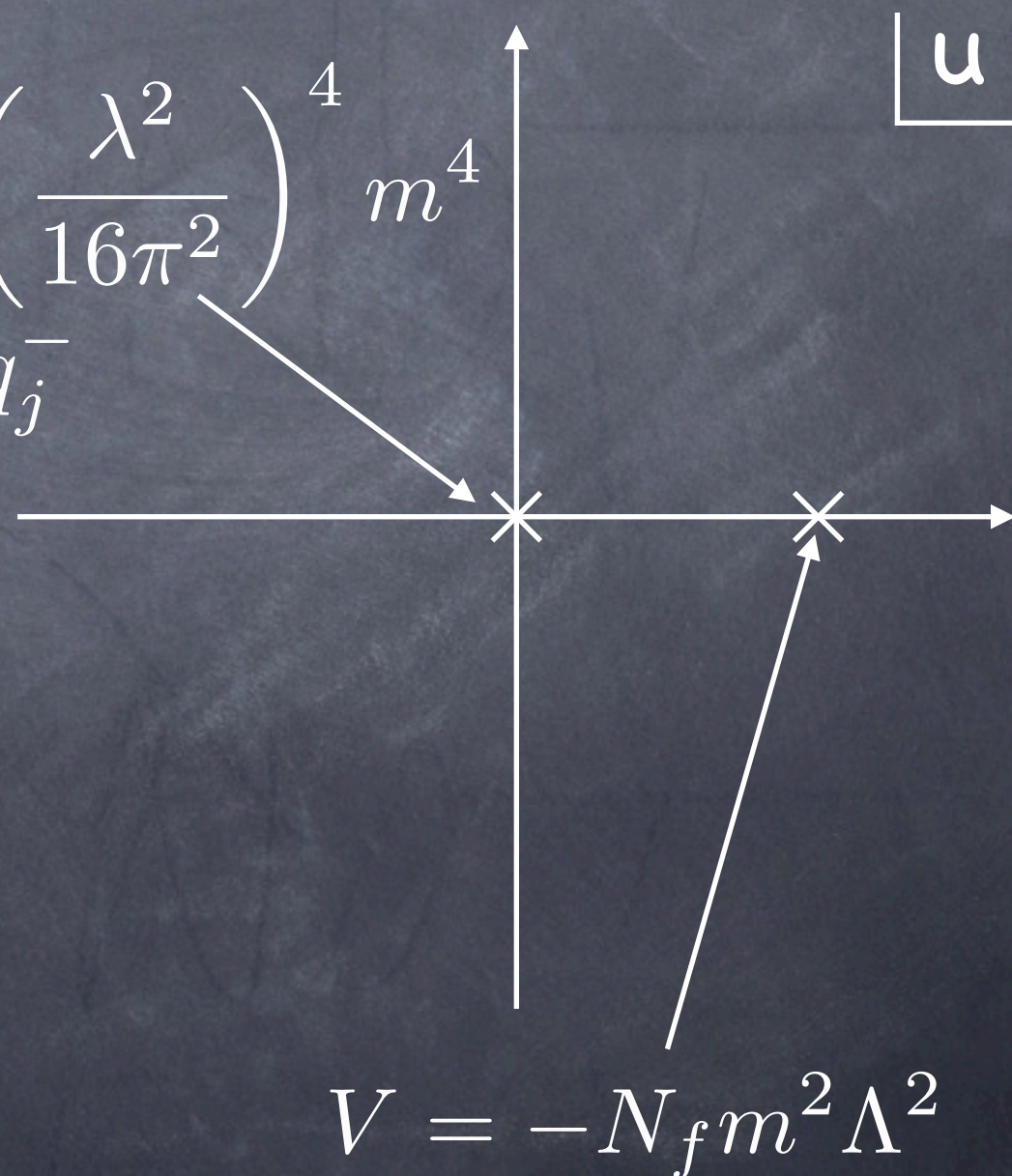
•  $u = \det M = \Lambda^{2N_f}$

• monopoles:

$$W = (u - \Lambda^{2N_f}) E^+ E^-$$

$$|E^\pm| = (m\Lambda)^{1/2}$$

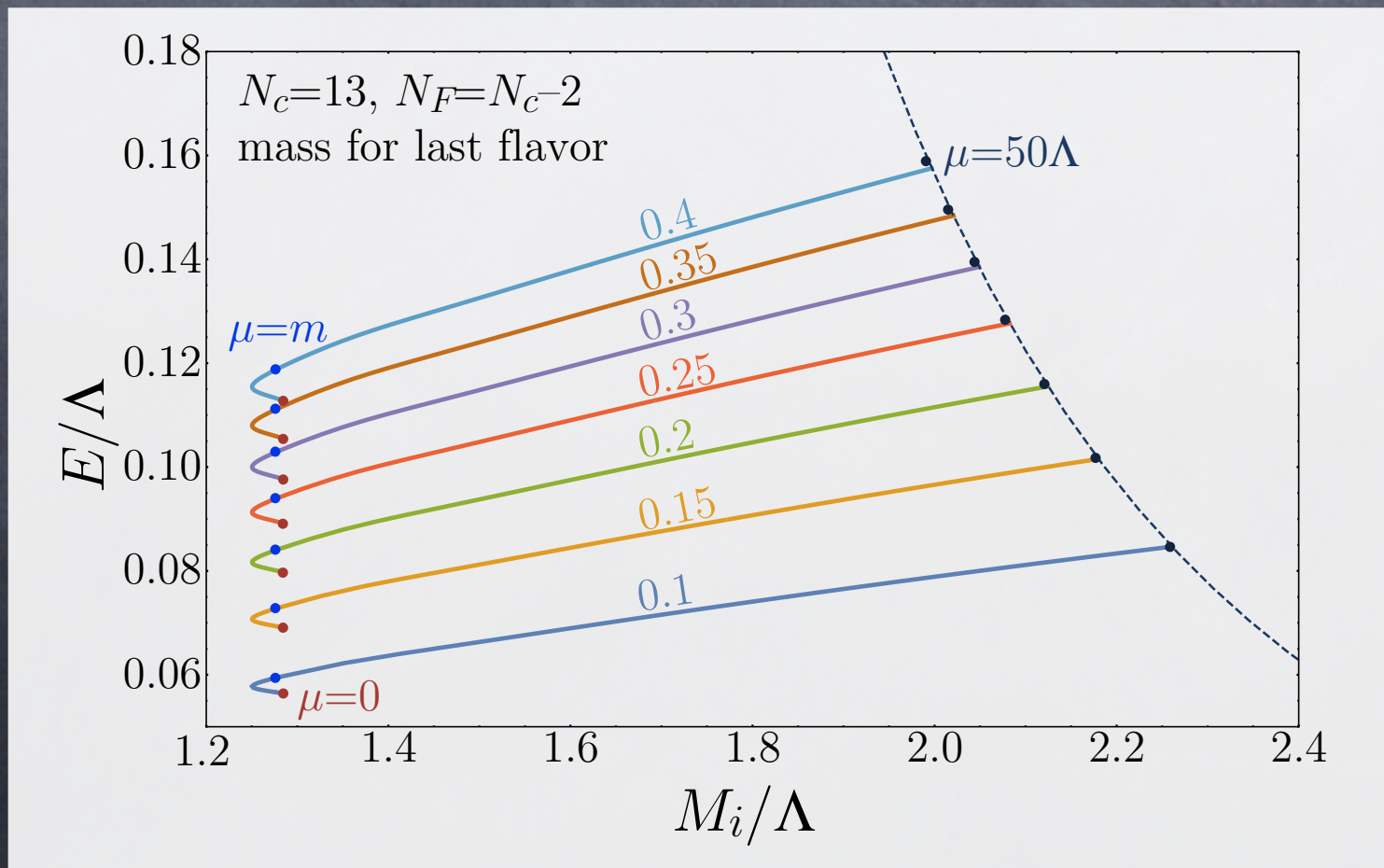
• both monopoles and meson  
 condense!





$$N_f < N_c - 2$$

- add mass  $m_q$  to some of the quarks
- can show monopole VEVs persist  $m_q \rightarrow \infty$
- demonstration of confinement and chiral symmetry breaking for all  $N_f \leq N_c - 2$





$$N_f = N_c + 1$$

“Confinement without  $\chi$ SB”

$$W = \frac{\det M - \tilde{B} M B}{\Lambda^{2N_c - 1}}$$

$$W = \lambda \frac{\det M}{\Lambda^{N_f - 3}} - \kappa \tilde{B} M B$$

$$B_i = \epsilon_{ij_1 \dots j_{N_c}} Q^{j_1} \dots Q^{j_{N_c}}$$

$$V = N_f \lambda^2 \frac{|\phi|^{2N_f - 2}}{\Lambda^{2N_f - 6}} - \lambda(N_f - 3)m\phi^{N_f} + c.c.$$

$$\phi = \kappa^{-1} \Lambda \left( \frac{2N_f - 3N_c}{N_c} \frac{m}{\Lambda} \right)^{(N_f - N_c)/(2N_c - N_f)} \ll \Lambda$$

Supersymmetric  
higher power potential

together with  
AMSB

$$SU(N_f)_L \times SU(N_f)_R \rightarrow SU(N_f)_V$$

$\chi$ SB!

massless pions

Deeper minimum  
with  $U(1)_B$  breaking

Andrea Luzio, Ling-Xiao Xu

2202.01239

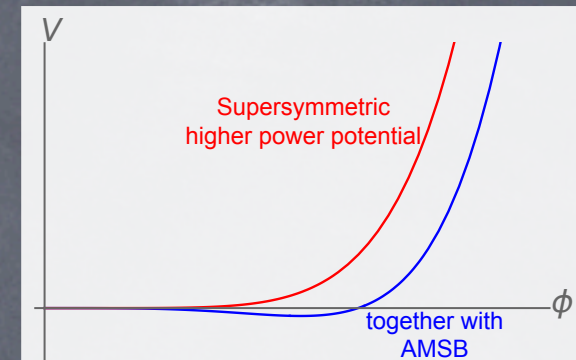
Csaki, Gomes, HM, Noether,  
Telem, Varier, in preparation



$$N_c + 2 \leq N_f < 3N_c/2$$

- “magnetic” IR-free  $SU(N_f - N_c)$  theory

$$W = \frac{1}{\mu} M^{ij} q_i \tilde{q}_j \rightarrow \lambda \tilde{M}^{ij} q_i \tilde{q}_j$$



- go along the meson direction with rank  $M = N_f$
- integrate out dual quarks with  $M^{ij} = \phi \delta^{ij}$
- pure  $SU(N_f - N_c)$  YM forms gaugino condensate

$$W = (N_f - N_c) \left( \frac{\kappa^{N_f} \det M}{\Lambda^{3N_c - 2N_f}} \right)^{1/(N_f - N_c)}$$

$$V = N_f \Lambda^4 \left| \frac{\kappa \phi}{\Lambda} \right|^{2N_c/(N_f - N_c)} - (2N_f - 3N_c) m \Lambda^3 \left( \frac{\kappa \phi}{\Lambda} \right)^{N_f/(N_f - N_c)} + c.c.$$

$$\phi = \kappa^{-1} \Lambda \left( \frac{2N_f - 3N_c}{N_c} \frac{m}{\Lambda} \right)^{(N_f - N_c)/(2N_c - N_f)} \ll \Lambda$$

- $SU(N_f)_L \times SU(N_f)_R \rightarrow SU(N_f)_V$   $V \approx -\Lambda^4 \left( \frac{m}{\Lambda} \right)^{2N_c/(2N_c - N_f)}$
- deeper minimum with  $U(1)_B$  breaking

Csaki, Gomes, HM, Noether, Telem, Varier, in preparation



$$3N_c/2 < N_f < 3N_c$$

- integrate  $q$  out like in free magnetic phase

$$W = (N_f - N_c) \left( \frac{\kappa^{N_f} \det M}{\Lambda^{3N_c - 2N_f}} \right)^{1/(N_f - N_c)}$$

- No stable minimum for  $N_f > 2N_c$ ?

- Actually remember  $Z_M(\mu) = \left( \frac{\mu}{\Lambda} \right)^{(6N_c - 4N_f)/N_f}$

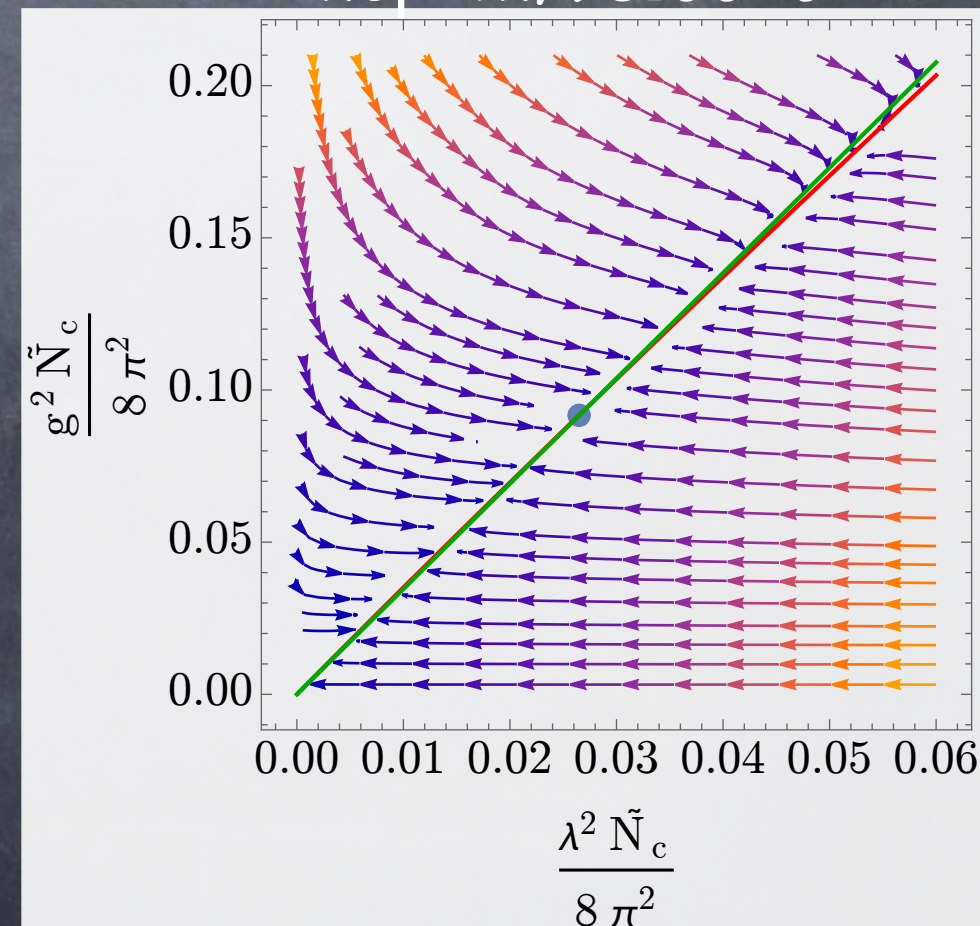
- Effective potential

$$V = Z_M(\phi)^{-1} \left| \frac{\partial W}{\partial M} \right|^2 + (N_f - 3)mW$$

- Well-defined local minimum

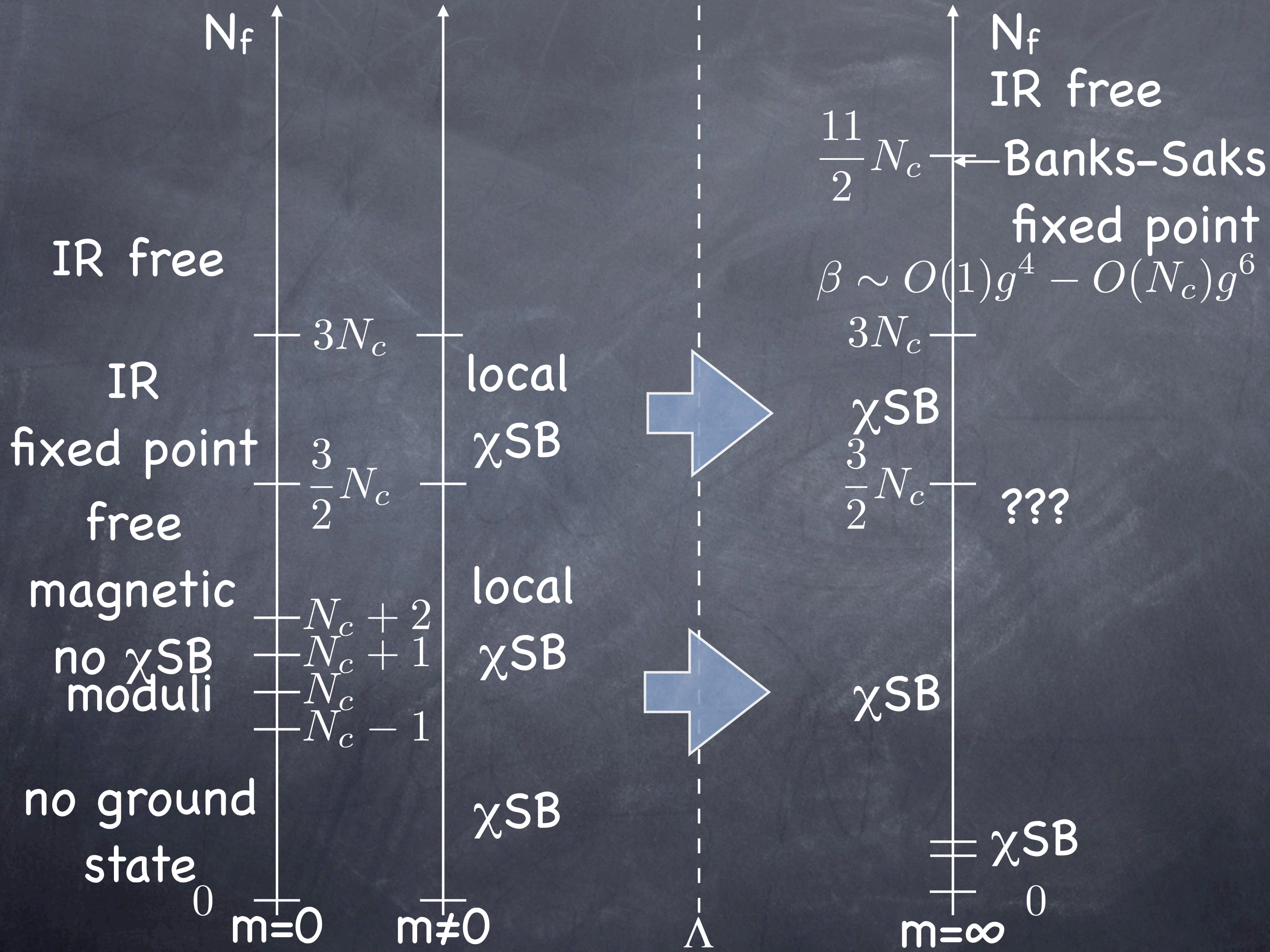
- $\chi$ SB up to  $N_f < 3N_c$

Friedland, de Gouvêa, HM  
hep-th/9810020



HM, Bea Noether, Digvijay Roy Varier  
arxiv:2111.09690







# Chiral gauge theories

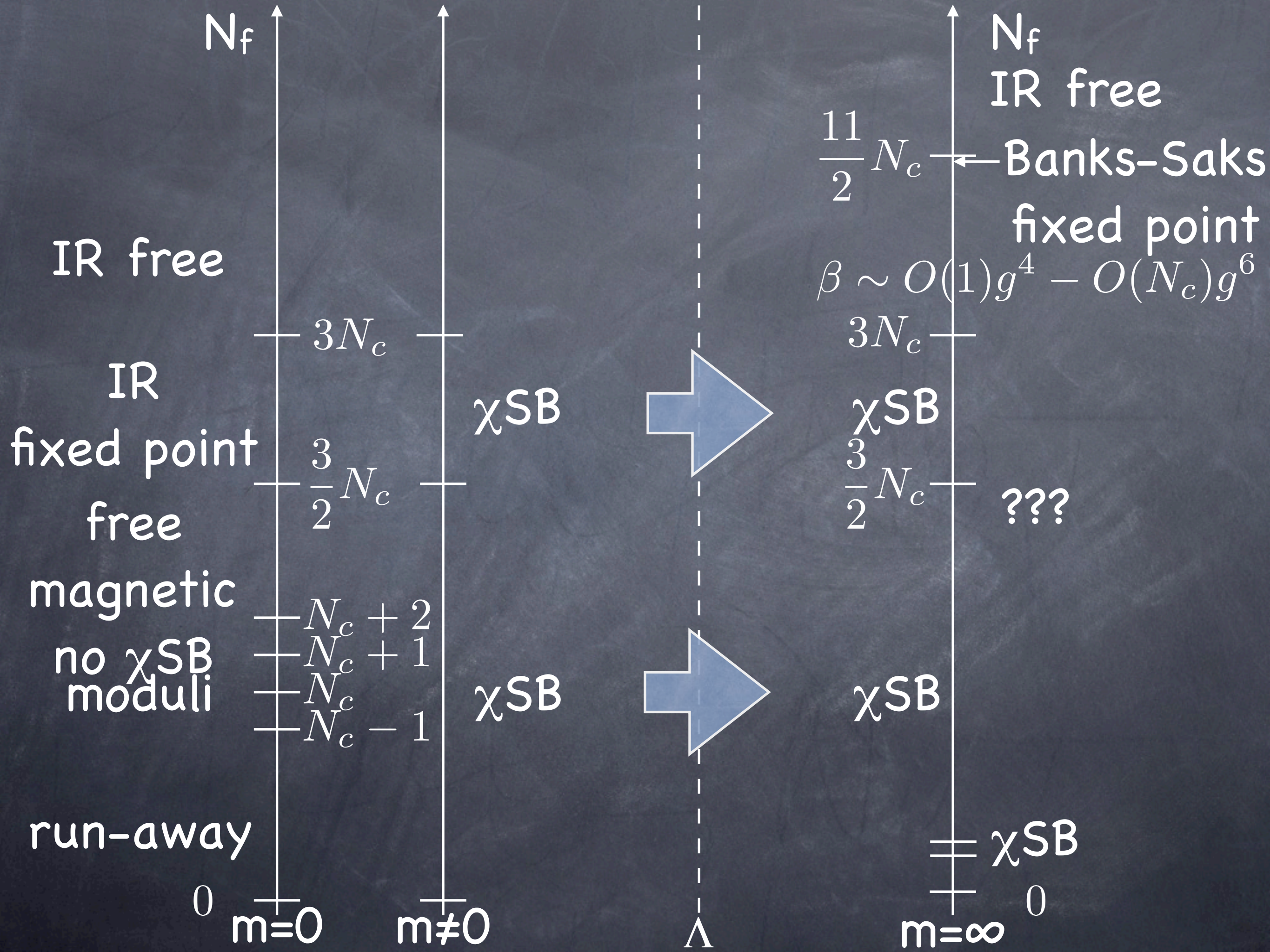
- It has been very difficult to formulate them on lattice because of fermion doubling
- Recent progress: domain wall fermions and "overlap lattice Dirac operator"
- Still challenging numerical problems
- Never simulated on lattice
- Only hand-waving ideas called "tumbling"
  - Theory breaks itself due to fermion bilinear condensates in the MAC



# $SO(10)$ with $N_f$ 16's

- Yoshio Kikukawa: most likely chiral gauge theory to be simulated on lattice
- $SO(10)$ : smallest anomaly-free group with complex representations
- 16: smallest complex representation
- $N_f=1$ : dynamical SUSY breaking, gapped
- $N_f=2$ : gapped,  $SU(2)$  unbroken, not tumbling
- $N_f=3$ :  $SU(3)$  broken dynamically to  $SO(3)$
- $N_f=4$ :  $SU(4)=SO(6)$  broken dynamically to  $SO(3)\times SO(3)$  or  $SO(5)$
- No massless composite fermions







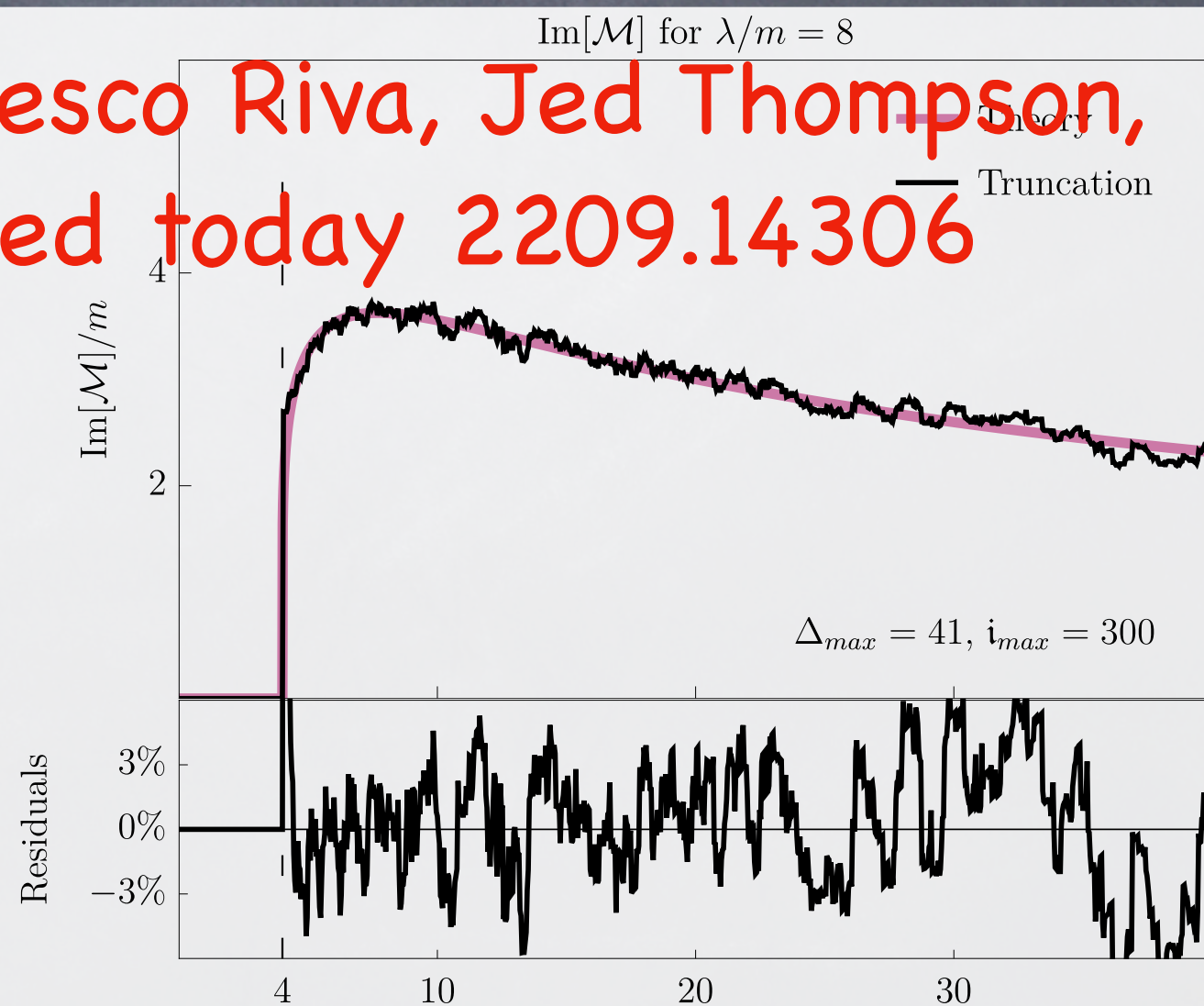
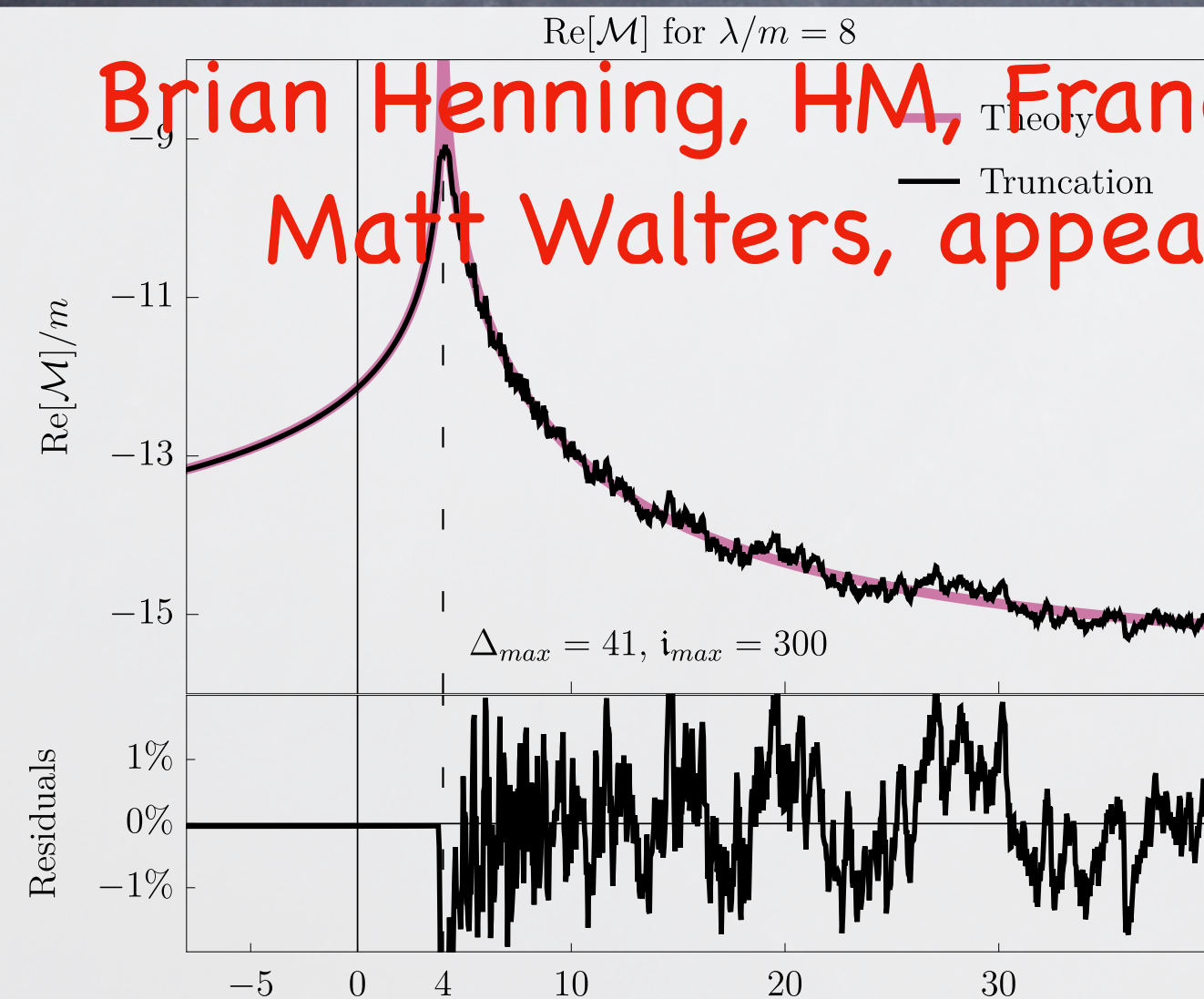
# Conclusion

- N=1 SUSY + AMSB: a great tool to study non-SUSY gauge theories
- vector-like theories
  - Can show  $\langle \bar{q}q \rangle \neq 0$
  - Monopole condensation for  $SO(N_c)$
  - For  $N_f > N_c$ , local minima seem useful
- Chiral gauge theories can also be solved
  - Concrete predictions for  $SO(10)$  with 16's
  - Also for  $SU(N_c)$  with  $A+Fbar$ ,  $S+Fbar$
  - Need to understand general symmetry breaking patterns with more examples



# Hamiltonian truncation

- One can approximate QFT with a finite-dimensional Hilbert space w/o Wick rotation
- Can study scattering problems directly
- $O(N)$  sigma model in 2+1 dim with non-perturbative Wilson-Fischer fixed point



Brian Henning, HM, Francesco Riva, Jed Thompson,  
Matt Walters, appeared today 2209.14306



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Hitoshi Murayama (IPMU & Berkeley)

## 28 September 2022

