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@ DESY theory workshop

# Dark Gauge Bosons in Neutron Stars

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(University of Padova)


*“Cooling of young neutron stars and dark gauge bosons,”* Deog Ki Hong (PNU), Chang Sub Shin (CNU), **SY**  
PRD 103 (2021) 12, 123031, arXiv:2012.05427

*“Dark gauge boson production from neutron stars via nucleon-nucleon bremsstrahlung,”* Chang Sub Shin (CNU), **SY**  
JHEP 02 (2022) 133, arXiv:2110.03362

Work in progress, Chang Sub Shin (CNU), **SY**

# Outline

**01 Introduction** • Effective theory of dark gauge boson (in a medium)

**02 Standard NS cooling** • SN1987A  
• (NS1987A)  
• Rapid cooling of Cas A  **Stellar cooling argument**  
- superfluidity?

**03  $\gamma'$  cooling** [Dark photon  $U(1)_{B-L}$ ] • NN-bremsstrahlung  
• (Compton pionic)  
• Cooper pairing

**04 Constraints on dark gauge bosons**

**05 Conclusion**

# ***Effective theory of Dark gauge boson***

- Chiral perturbation theory framework
  - Hadronic interactions
- Thermal field theory framework
  - Plasma effect : mixing & screening
- Two benchmarks : Dark photon &  $U(1)_{B-L}$

# Dark $U(1)$ gauge boson

- Effective  $\gamma'$  couplings at  $\mu > \Lambda_{\text{QCD}}$  - *perturbative QCD*

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} - \frac{1}{4}X_{\mu\nu}X^{\mu\nu} + \frac{\varepsilon}{2}F_{\mu\nu}X^{\mu\nu} - \frac{1}{2}m_{\gamma'}^2 X_\mu X^\mu + eA_\mu J_{\text{EM}}^\mu + e'A'_\mu J'^\mu$$

EM field strength      DP field strength      kinetic mixing       $\gamma'$  mass      EM current       $\gamma'$  current  
- quark  
- lepton

- Effective  $\gamma'$  couplings at  $\mu < \Lambda_{\text{QCD}}$  ?

✓ Chiral perturbation theory (ChPT) with hadronic resonances (e.g., nucleons and mesons)

- Effective  $\gamma'$  couplings in a medium?

✓ Thermal field theory → plasma effect

# Hadronic interactions

## ● Basic principle in ChPT?

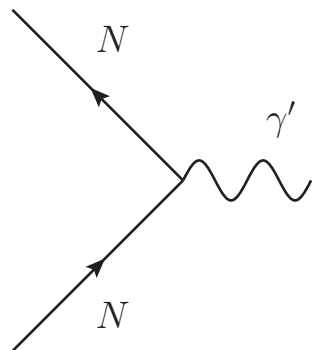
➡ Matching currents with the same symmetry properties

## ● $\gamma'$ couplings to hadrons from $e' A'_\mu J'^\mu$

$$J'_\mu = \sum_i q'_i \bar{\psi}_i \gamma_\mu \psi_i$$

[Kroll-Ruderman, 1954]

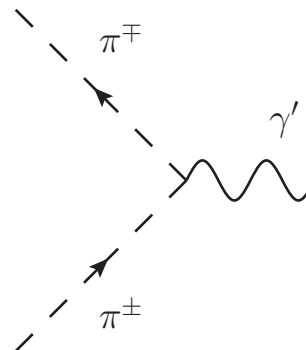
### • Nucleon currents



$$e' A'_\mu \sum_N q'_N \bar{N} \gamma^\mu N$$

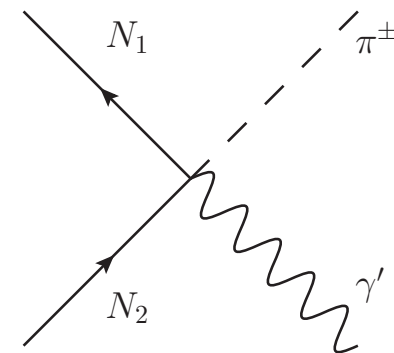
$$q'_p = 2q'_u + q'_d \quad q'_n = q'_u + 2q'_d$$

### • Pion currents



$$e' A'_\mu (q'_p - q'_n) i \pi^- \overleftrightarrow{\partial}^\mu \pi^+$$

### • N+Pion currents



$$-ie' A'_\mu (q'_p - q'_n) \frac{g_A}{f_\pi} [\pi^+ \bar{p} \gamma^\mu \gamma^5 n - \pi^- \bar{n} \gamma^\mu \gamma^5 p]$$

Isospin breaking

★ What is effective couplings in a dense medium?  
( $q'_p$  &  $q'_n$ )

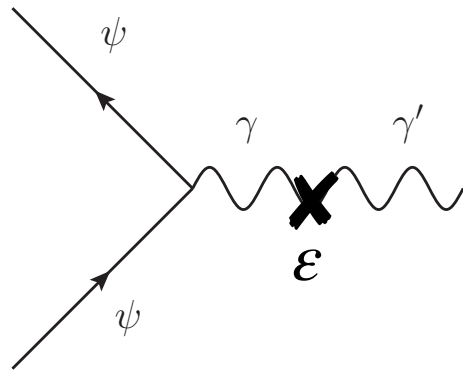
# Effective coupling of $\gamma'$ in medium

[E. Hardy, R. Lasenby, 16]

[C. S. Shin, D. Hong, SY, 20]

$$J_\psi^\mu = \bar{\psi} \gamma^\mu \psi$$

- Kinetic mixing ( $\epsilon$ )

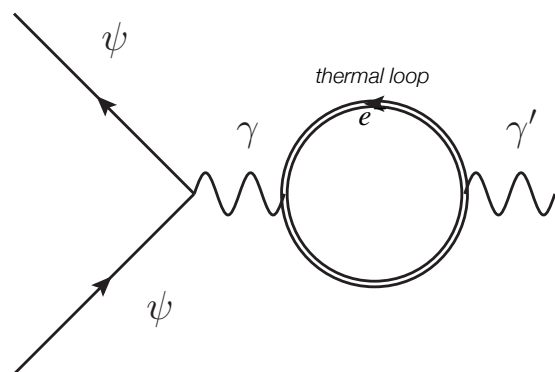


$$\mathcal{M} = e_\psi^\psi J_\psi^\mu \epsilon'_\mu$$

- $q$  : EM charge
- $q'$  : dark  $U(1)$  charge

$$e_{\text{eff}}^\psi = e' \left( q'_\psi + q'_e q_\psi \right) + \left( \epsilon e - q'_e e' \right) q_\psi \frac{m_{\gamma'}^2}{m_{\gamma'}^2 - \pi_{T,L}}$$

- Plasma mixing

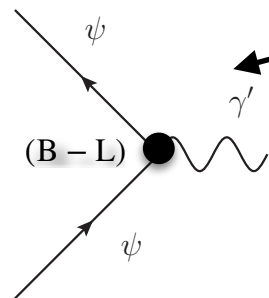


- If the gauge boson field basis coupling to  $\psi$  is  $(e\mathbf{q}_\psi \mathbf{A}_\mu + e'\mathbf{q}'_\psi \mathbf{A}'_\mu)$  is consistent with that of the **electron** (i.e., **plasmon**),  $(e\mathbf{q}_e \mathbf{A}_\mu + e'\mathbf{q}'_e \mathbf{A}'_\mu)$ ,
  - the first term vanishes, i.e.  $q'_\psi + q'_e q_\psi = 0$
  - plasma suppression ( $\propto \omega_{\text{pl}}^{-2}$ ) when  $\omega_{\text{pl}} > m_{\gamma'}$

# $U(1)_{B-L}$ gauge boson

$$\mathcal{L} = -\frac{1}{4} \underbrace{F_{\mu\nu} F^{\mu\nu}}_{\text{EM field strength}} - \frac{1}{4} \underbrace{X_{\mu\nu} X^{\mu\nu}}_{\text{DP field strength}} - \frac{1}{2} \underbrace{m_{\gamma'}^2 X_\mu X^\mu}_{\gamma' \text{ mass}} + \underbrace{e A_\mu J_{\text{EM}}^\mu}_{\text{EM current}} + \underbrace{e' A'_\mu J_{B-L}^\mu}_{\text{B-L current}}$$

⊙  $\gamma'$  emission



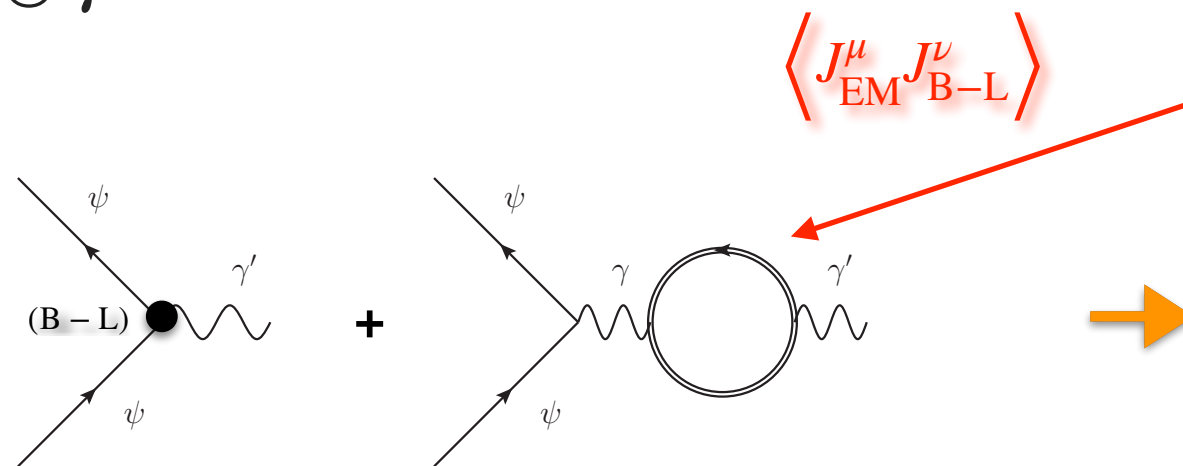
# $U(1)_{B-L}$ gauge boson

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} - \frac{1}{4}X_{\mu\nu}X^{\mu\nu} - \frac{1}{2}m_{\gamma'}^2 X_\mu X^\mu + eA_\mu J_{\text{EM}}^\mu + e'A'_\mu J_{B-L}^\mu$$

EM field strength
DP field strength
 $\gamma'$  mass
EM current
B-L current

plasma mixing

●  $\gamma'$  emission



	$e_{\text{eff}}^e$	$e_{\text{eff}}^p$	$e_{\text{eff}}^n$
$U(1)_{\text{DP}}$	$-\epsilon e \frac{m_{\gamma'}^2}{m_{\gamma'}^2 - \pi_{\text{T,L}}}$	$\epsilon e \frac{m_{\gamma'}^2}{m_{\gamma'}^2 - \pi_{\text{T,L}}}$	n/a
$U(1)_{B-L}$	$-e' \frac{m_{\gamma'}^2}{m_{\gamma'}^2 - \pi_{\text{T,L}}}$	$e' \frac{m_{\gamma'}^2}{m_{\gamma'}^2 - \pi_{\text{T,L}}}$	$e'$

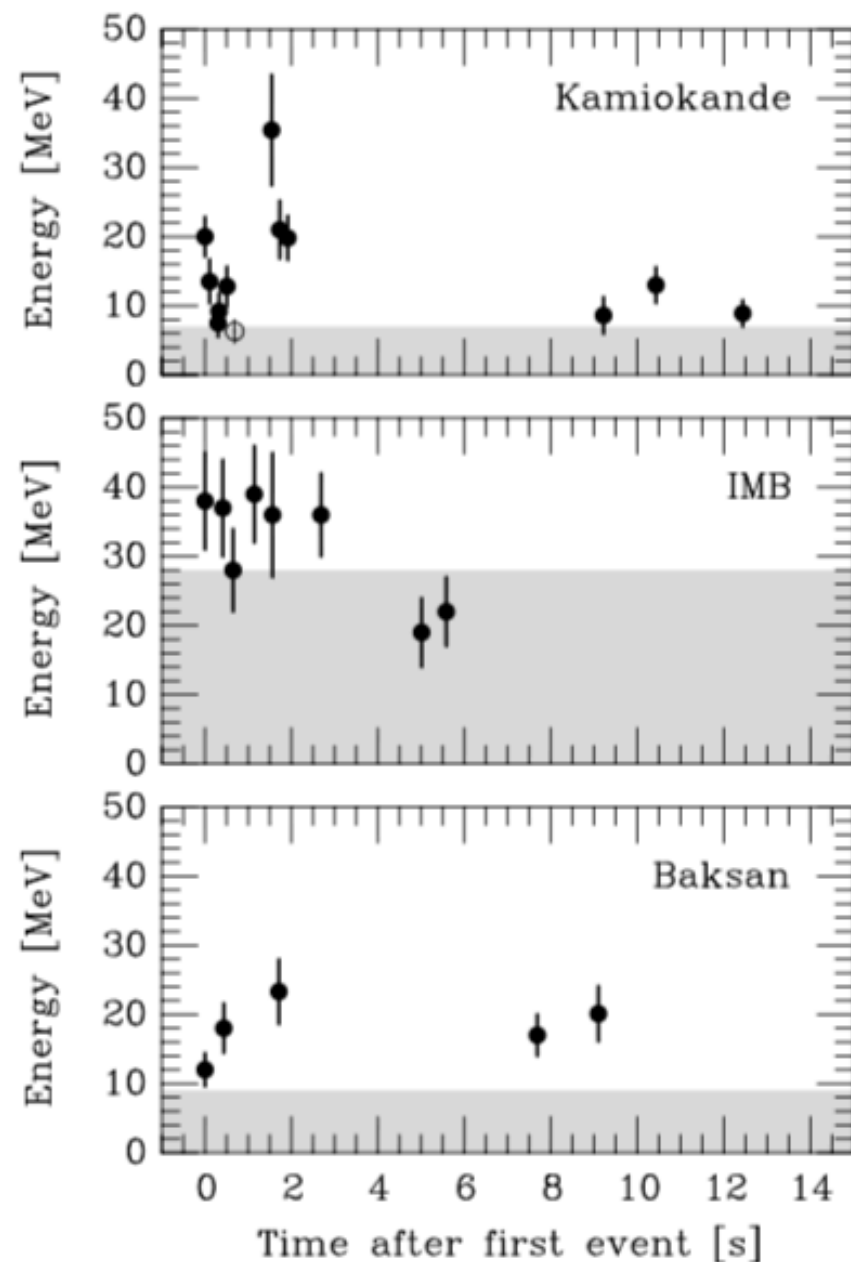
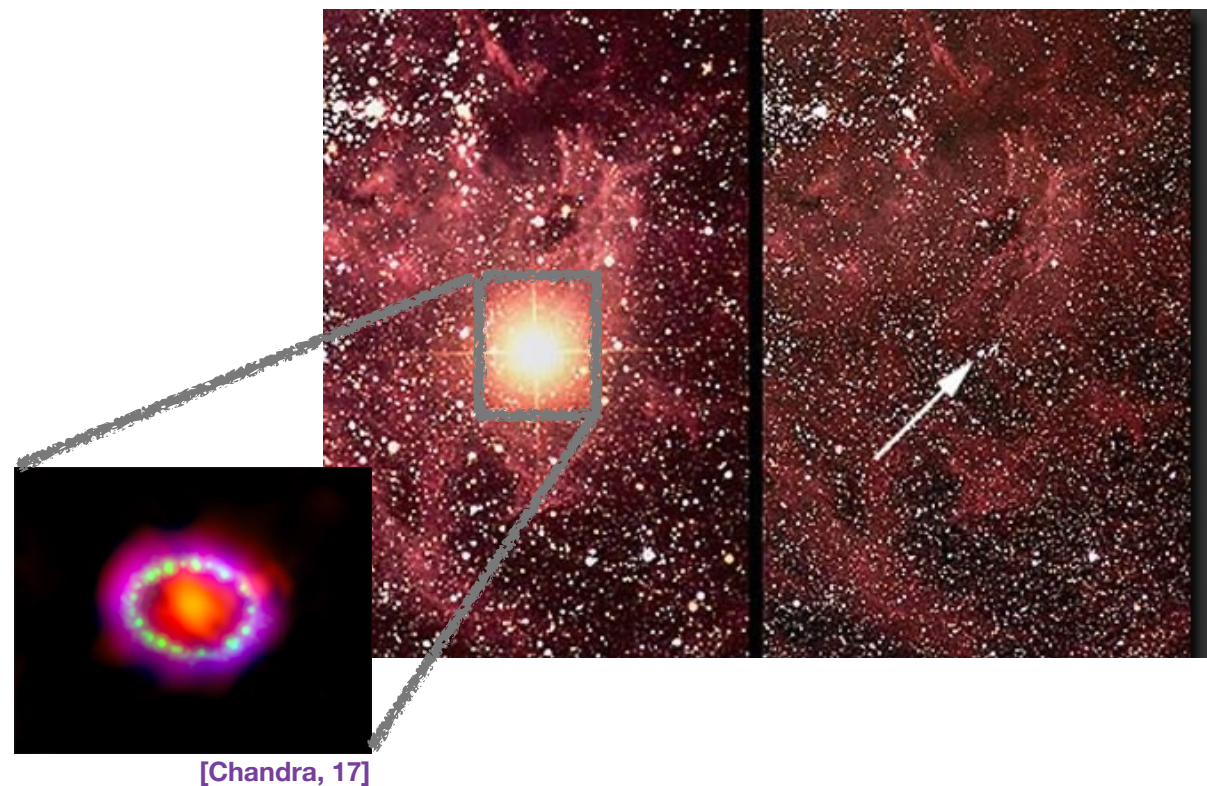
plasma mixing    **direct coupling**



# ***Core collapse Supernovae***

- Core collapse + the delayed neutrino mechanism
- Supernova neutrino signal of SN1987A
- Reduction of time duration by novel particle emission

# SN1987A



- The only direct observation of neutrinos occurred on 23 February 1987 when the blue supergiant Sanduleak–69 202 in the Large Magellanic Cloud exploded
- Time duration of SN neutrino flux (  $\sim 10$  sec)

➡  $L_x \lesssim 3 \times 10^{52} \text{ erg s}^{-1}$

[Raffelt, 96]

# Nucleon-Nucleon Bremsstrahlung

## ● Electrons in highly degenerate and relativistic limit

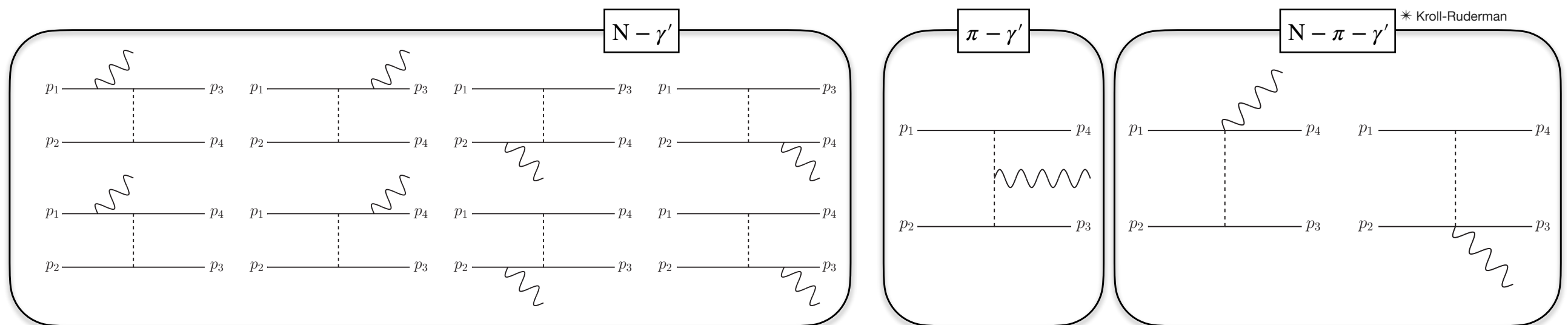
- $p_F \sim 1 \text{ fm}^{-1} \gg T \sim 10 \text{ MeV}$
- Pauli blocking makes processes associated with electrons (e.g.,  $e + \gamma \rightarrow e + \gamma'$ ) sub-dominant

## ● Nucleons in approximately non-degenerate and non-relativistic limit

- N-N bremsstrahlung with mediators associated with strong interaction

$N\text{-}\pi^-$  scatterings @ **Back up**

## ● Diagrams in one-pion-exchange approximation



# Multipole radiation

● Velocity as a good order parameter in expansion for scatterings

● Dipole radiation

- Leading order
- $\mathcal{M}_{\text{dipole}} \propto (e_{\text{eff}}^{N_1} - e_{\text{eff}}^{N_2})$  of isospin breaking combination
- Center of charges  $\neq$  Center of masses

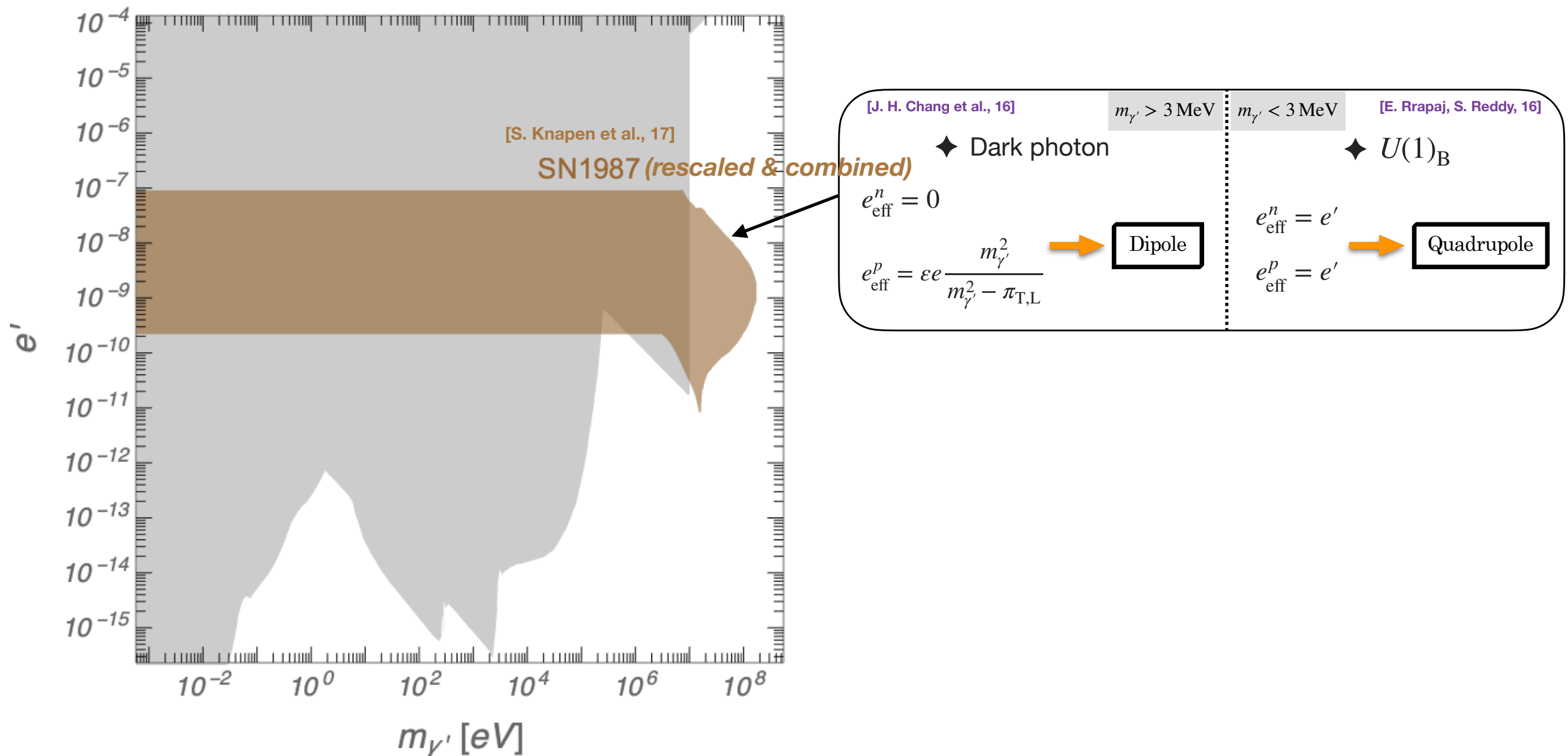


only for n-p bremsstrahlung

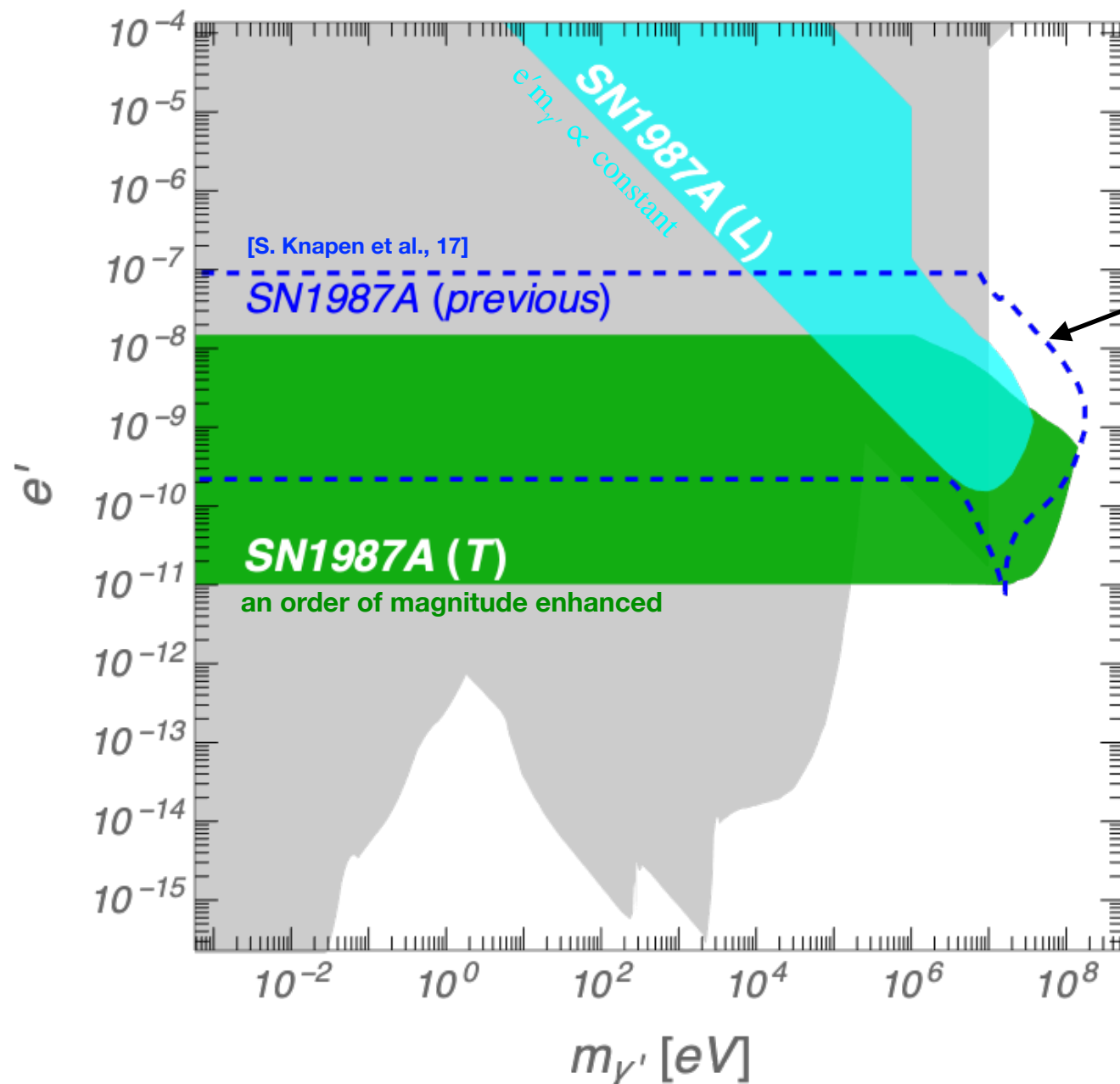
● Quadrupole radiation

- Next-leading order
- $\mathcal{M}_{\text{quadrupole}}$  as  $\left| \frac{\mathcal{M}_{\text{quadrupole}}}{\mathcal{M}_{\text{dipole}}} \right| \sim v_N$
- Center of charges = Center of masses

# SN1987A bound on $U(1)_{B-L}$



# SN1987A bound on $U(1)_{B-L}$



[J. H. Chang et al., 16]	$m_{\gamma'} > 3 \text{ MeV}$	$m_{\gamma'} < 3 \text{ MeV}$	[E. Rrapaj, S. Reddy, 16]
◆ Dark photon			◆ $U(1)_B$
$e_{\text{eff}}^n = 0$			$e_{\text{eff}}^n = e'$
$e_{\text{eff}}^p = \epsilon e \frac{m_{\gamma'}^2}{m_{\gamma'}^2 - \pi_{T,L}}$	→ Dipole		$e_{\text{eff}}^p = e'$
			→ Quadrupole

◆  $U(1)_{B-L}$

$e_{\text{eff}}^n = e'$

$e_{\text{eff}}^p = e' \frac{m_{\gamma'}^2}{m_{\gamma'}^2 - \pi_{T,L}}$  → Dipole

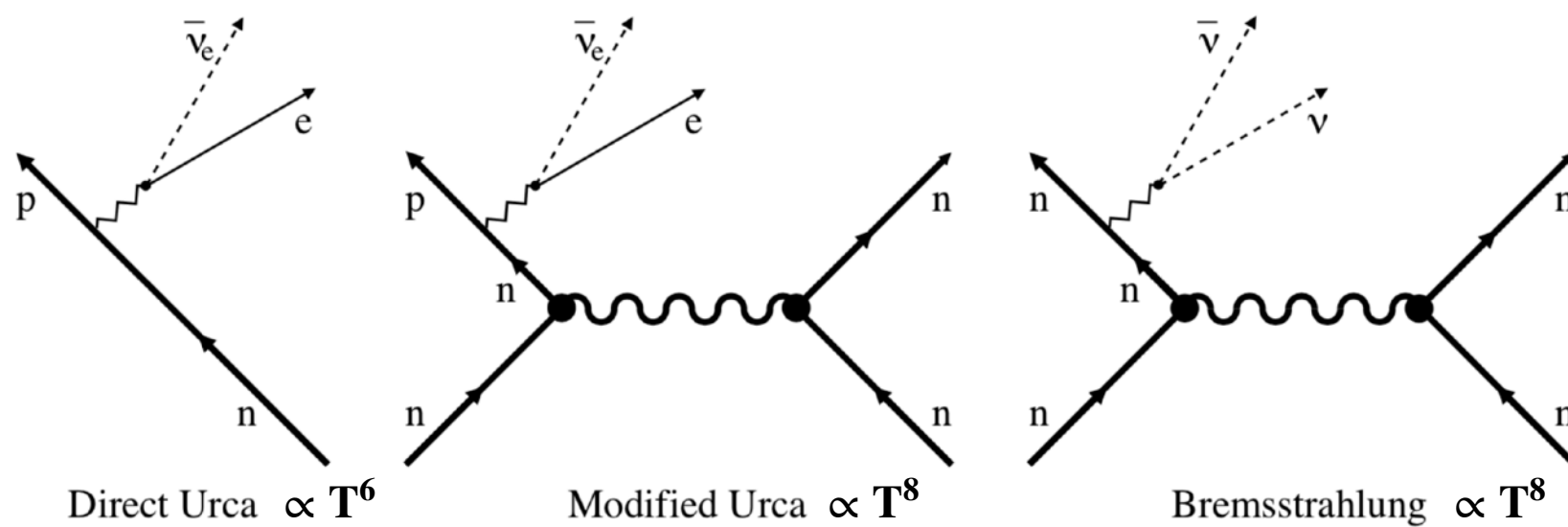
⚡ plasma screening

# *Neutron star cooling*

- Minimal cooling scenario
- Observation of rapid cooling at early era (CAS-A)
  - evidence of superfluidity?
- Constraints on novel particles

# Young NS cooling

- Neutrino (volume, mostly from core) emission

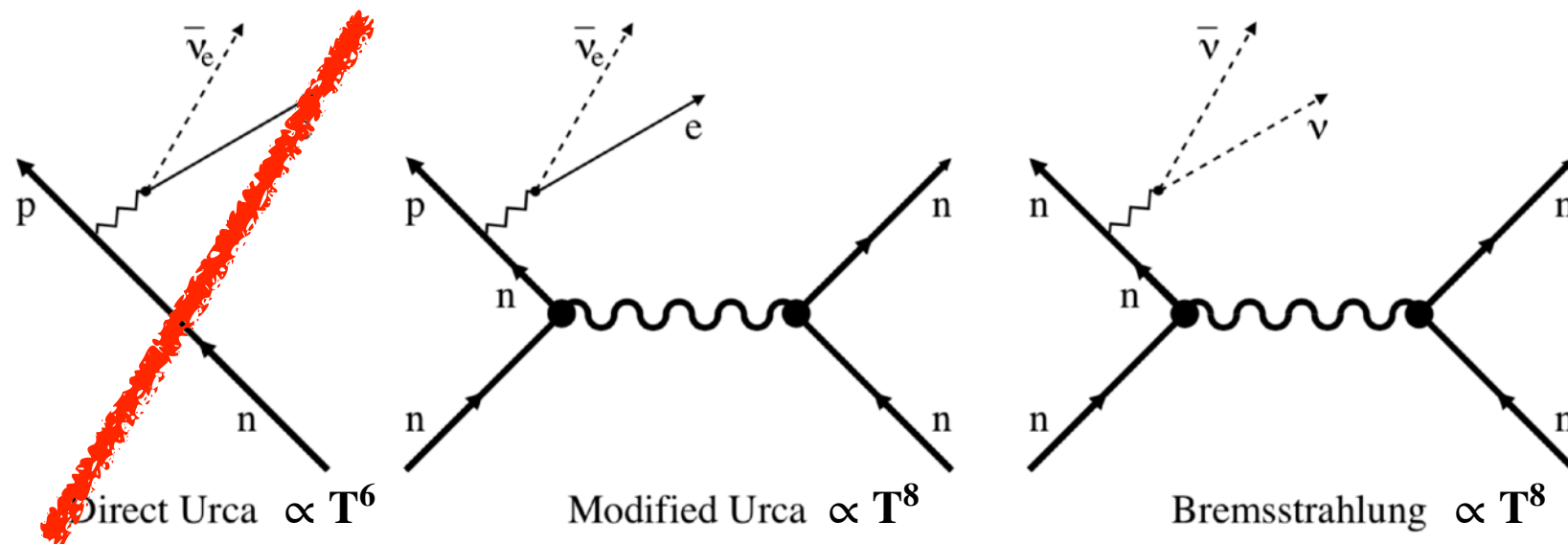


- Photon (surface) emission



# Young NS cooling

- Neutrino (volume, mostly from core) emission



only when  $> 2M_{\odot}$

© Energy balance equation

$$C \frac{dT}{dt} = -L_{\nu} - \cancel{L_{\gamma}} + \cancel{\dot{H}}$$

$\xrightarrow{C \propto T}$

$$\frac{\Delta T}{T} \sim -\frac{1}{6} \frac{dt}{t}$$

~~• Photon (surface) emission~~ dominant when  $> 10^5$  year

# Cassiopeia A

[Chandra]

- Exploded 340 years ago
- Carbon (light element) atmosphere
- Rapid cooling of surface temperatures derived from X-ray observations:

[M. J. P. Wijngaarden et al, 2019]

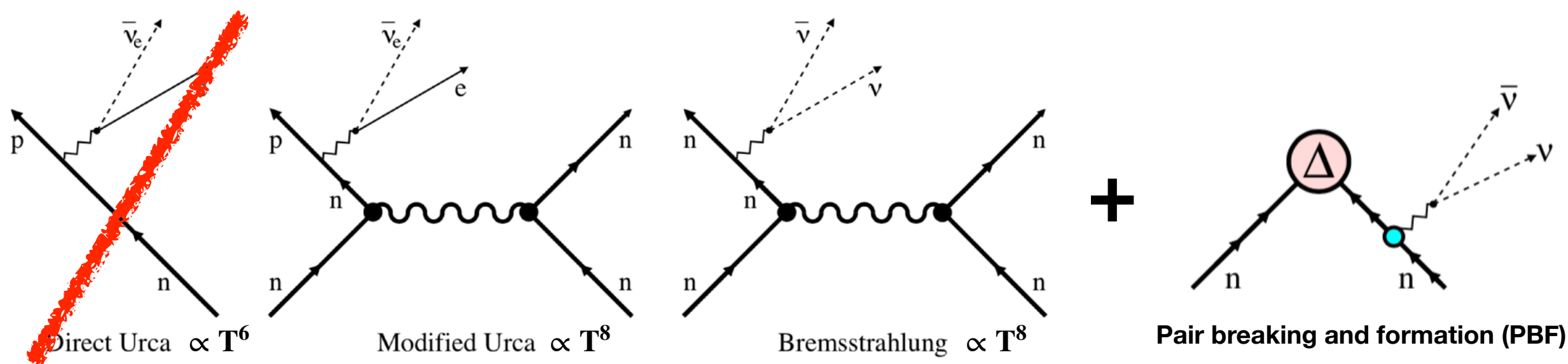
**2~3% for 10 yr decline rate**

expectation w/o superfluidity

$$\gg \frac{\Delta T}{T} \sim -\frac{1}{6} \frac{dt}{t} \sim 0.5 \%$$

# Young NS cooling

- Neutrino (volume) emission



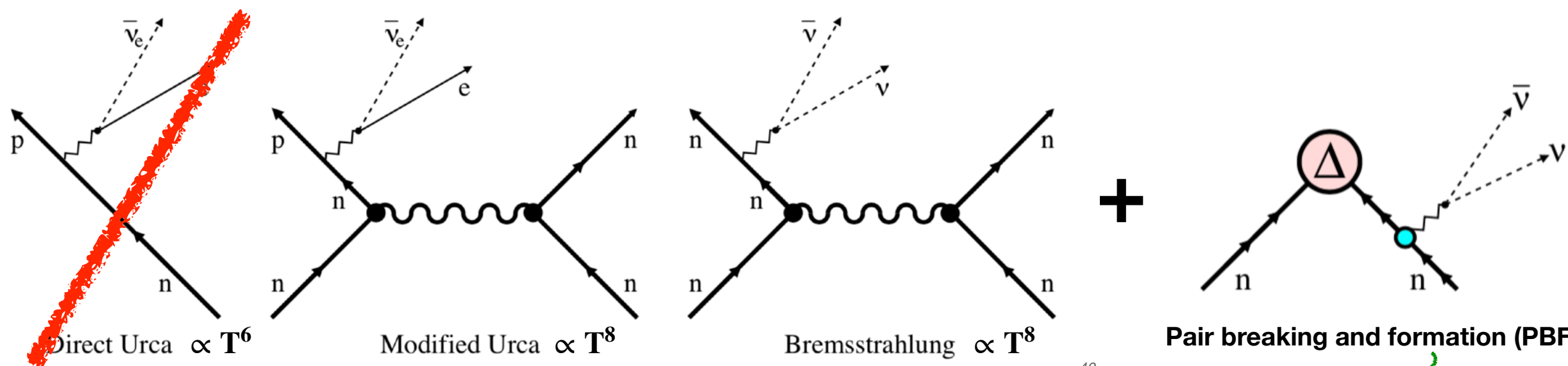
only when  $> 2M_{\odot}$

~~• Photon (surface) emission~~

dominant when  $> 10^5$  year

# Young NS cooling

- Neutrino (volume) emission

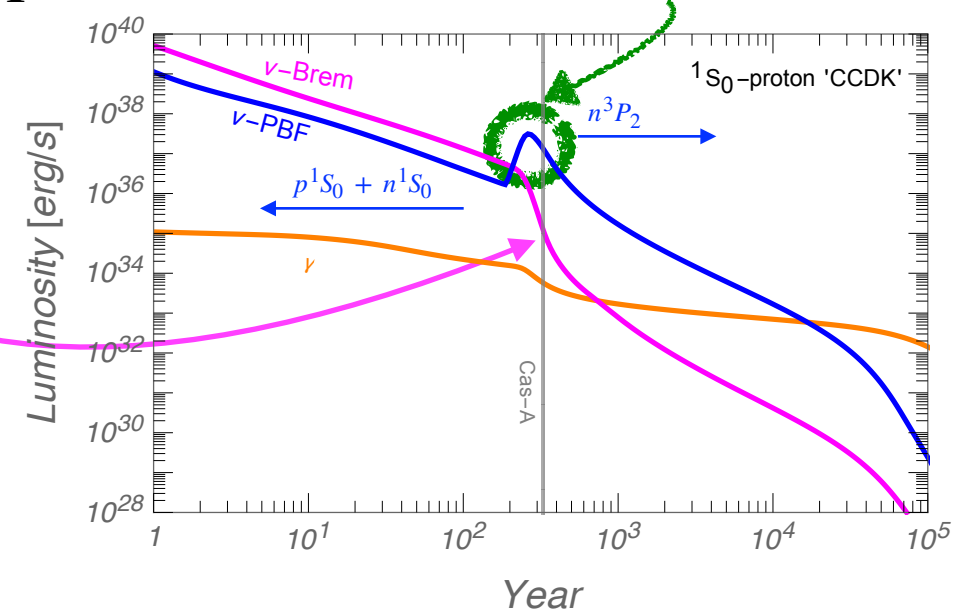


only when  $> 2M_{\odot}$

Suppressed at  $T < T_c$

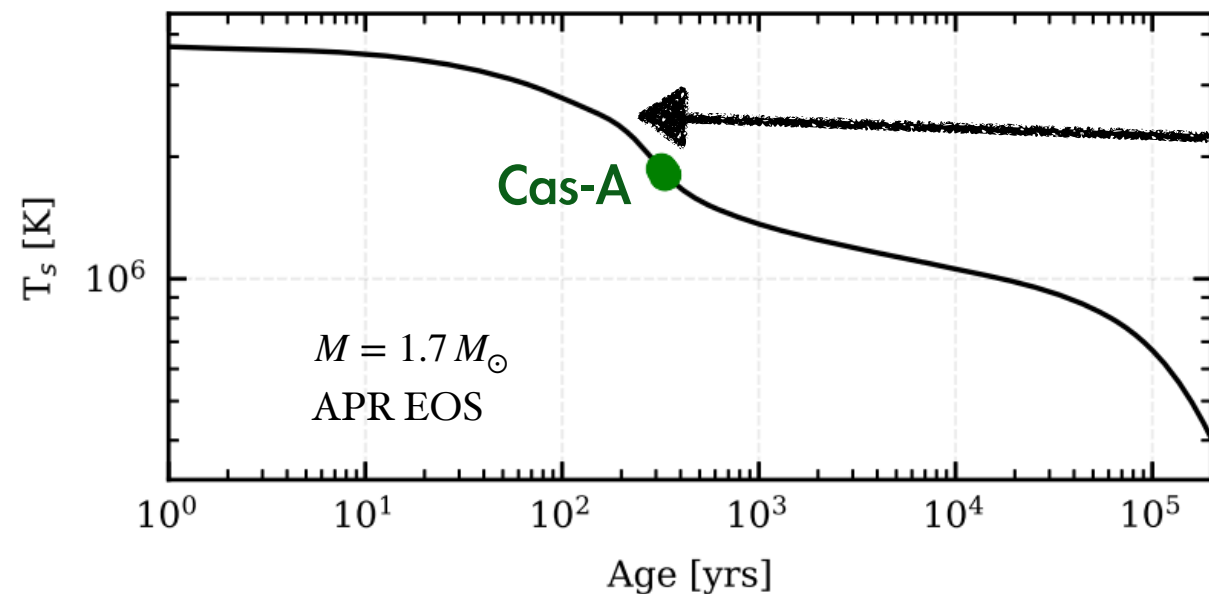
~~Photon (surface) emission~~

dominant when  $> 10^5$  year





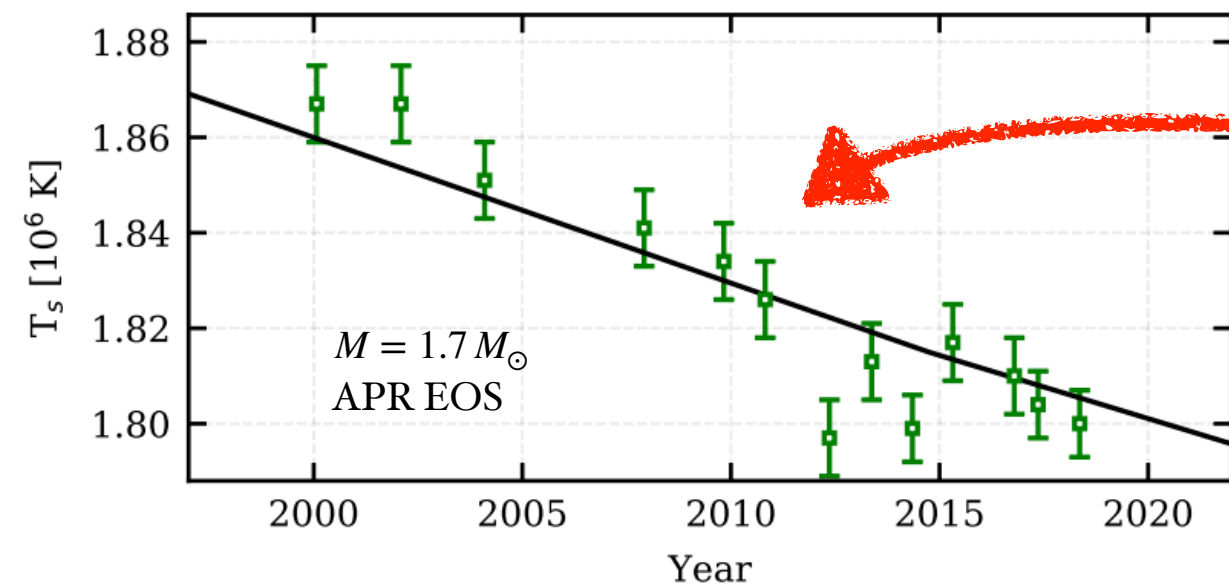
# Evidence for superfluidity?



$$T = T_c \sim 5.9 \times 10^8 \text{ K for } {}^3P_2$$



Rapid cooling by **PBF**



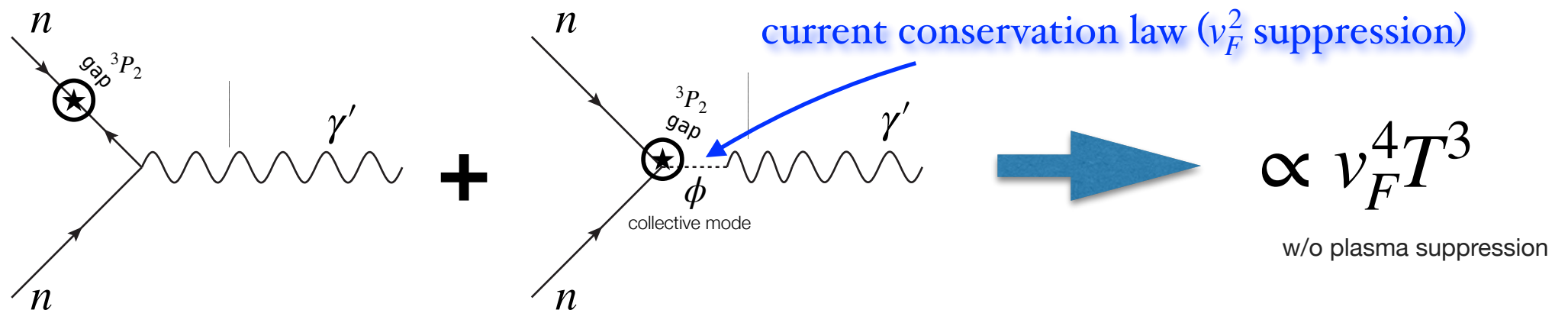
Minimal cooling scenarios with  
superfluidity fit well to the data

[M. J. P. Wijngaarden et al, 2019]

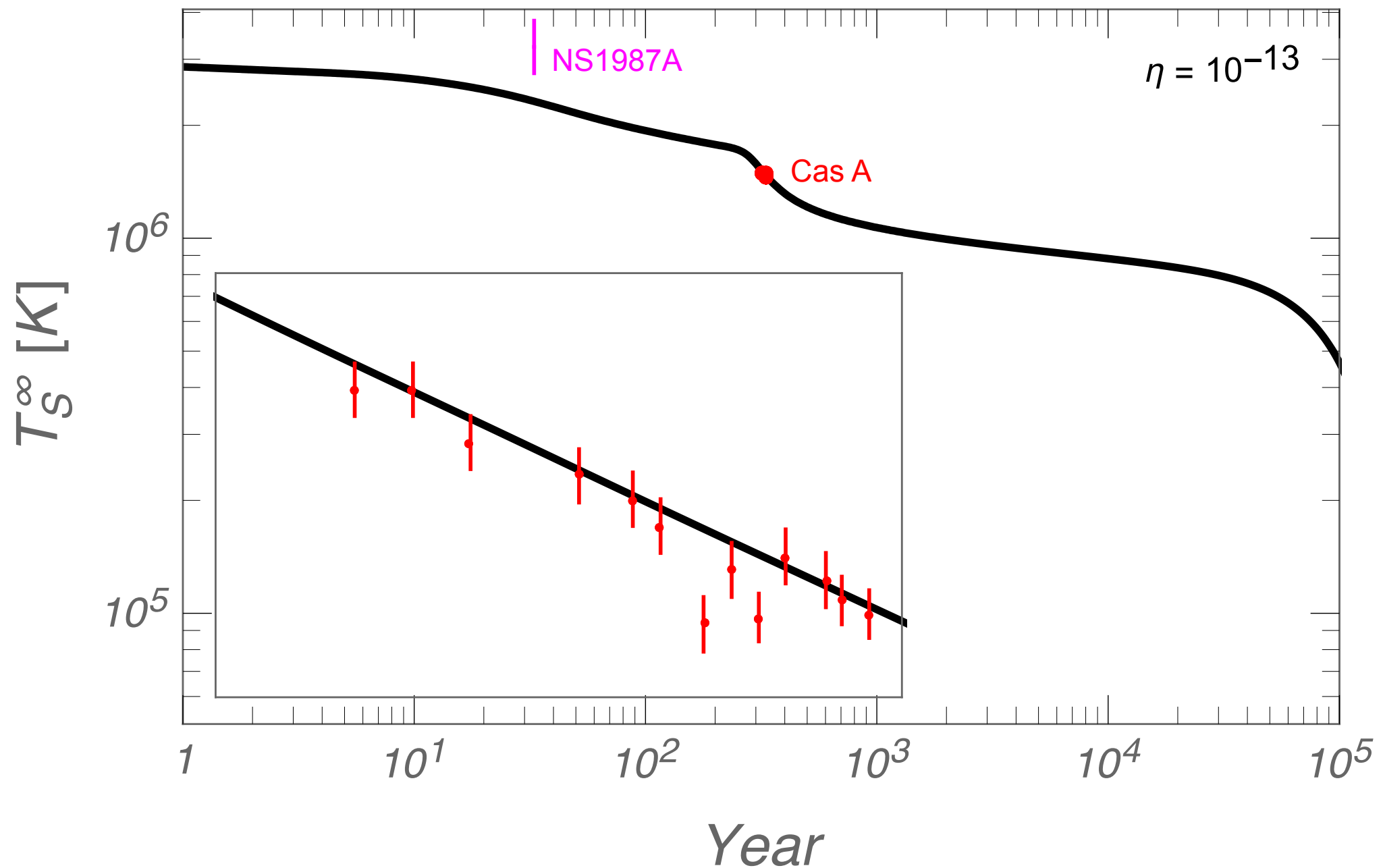
# $U(1)_{B-L}$ gauge boson in NS

- Emission processes through the nucleon current :
  - nn- or np-bremsstrahlung :  $nn (np) \rightarrow nn (np) + \gamma'$  (dominant for very young NS)
  - Neutron singlet PBF in crust
  - Neutron triplet PBF in core  $\rightarrow$  **Dominant!!** (after  $\sim 200$  yr)

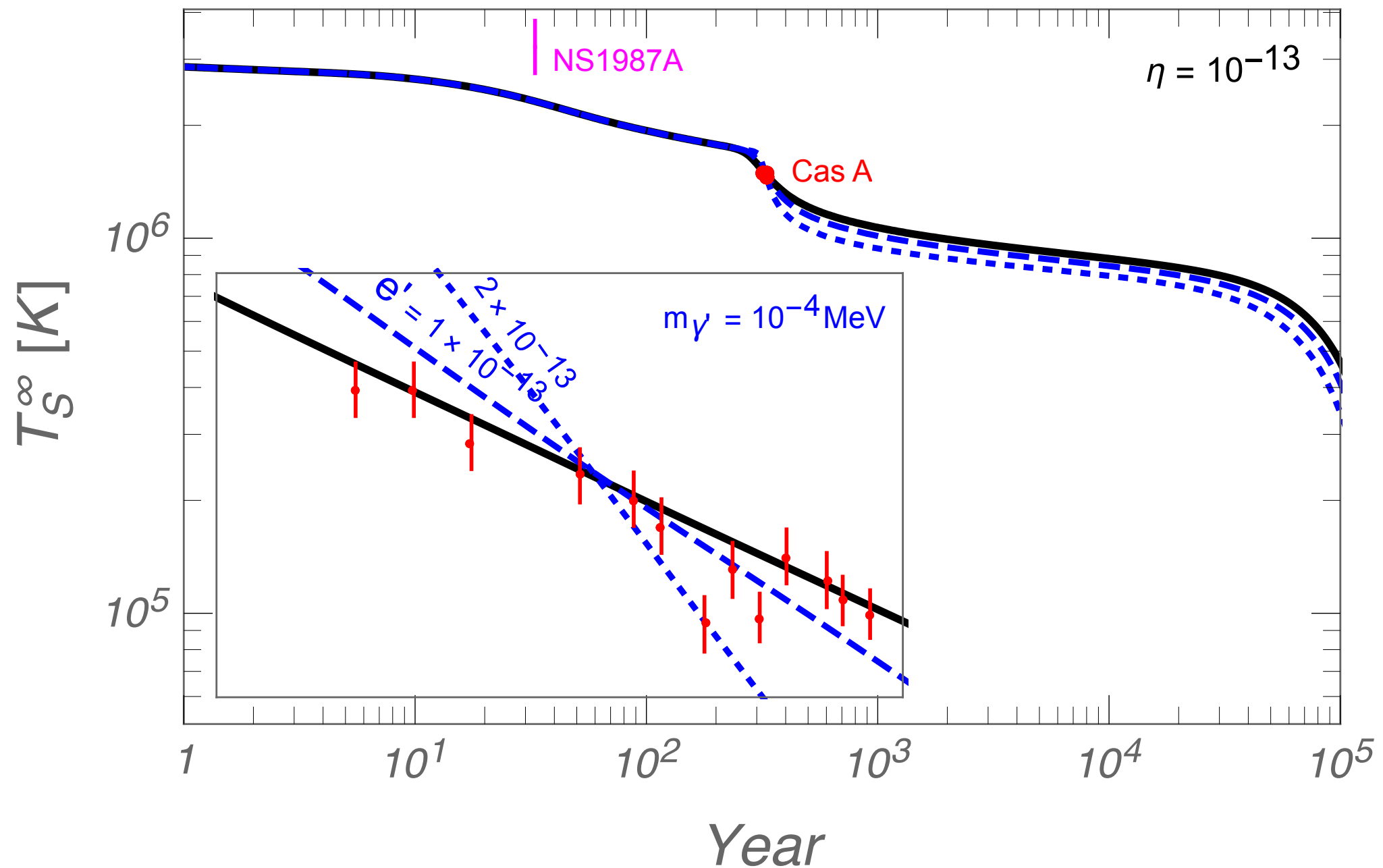
## ⊙ Neutron triplet **PBF** in core



# $U(1)_{B-L}$ gauge boson cooling

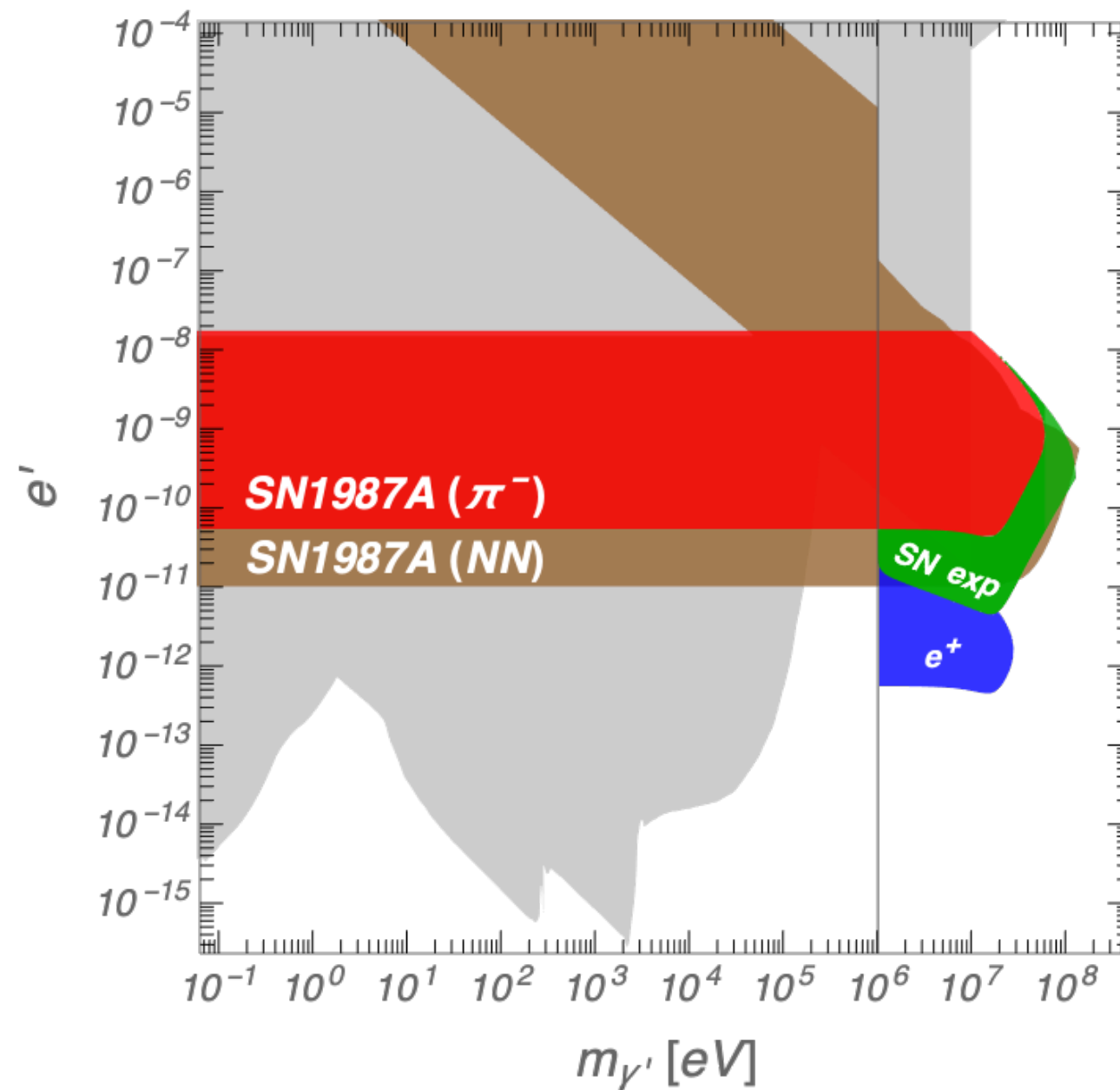


# $U(1)_{B-L}$ gauge boson cooling

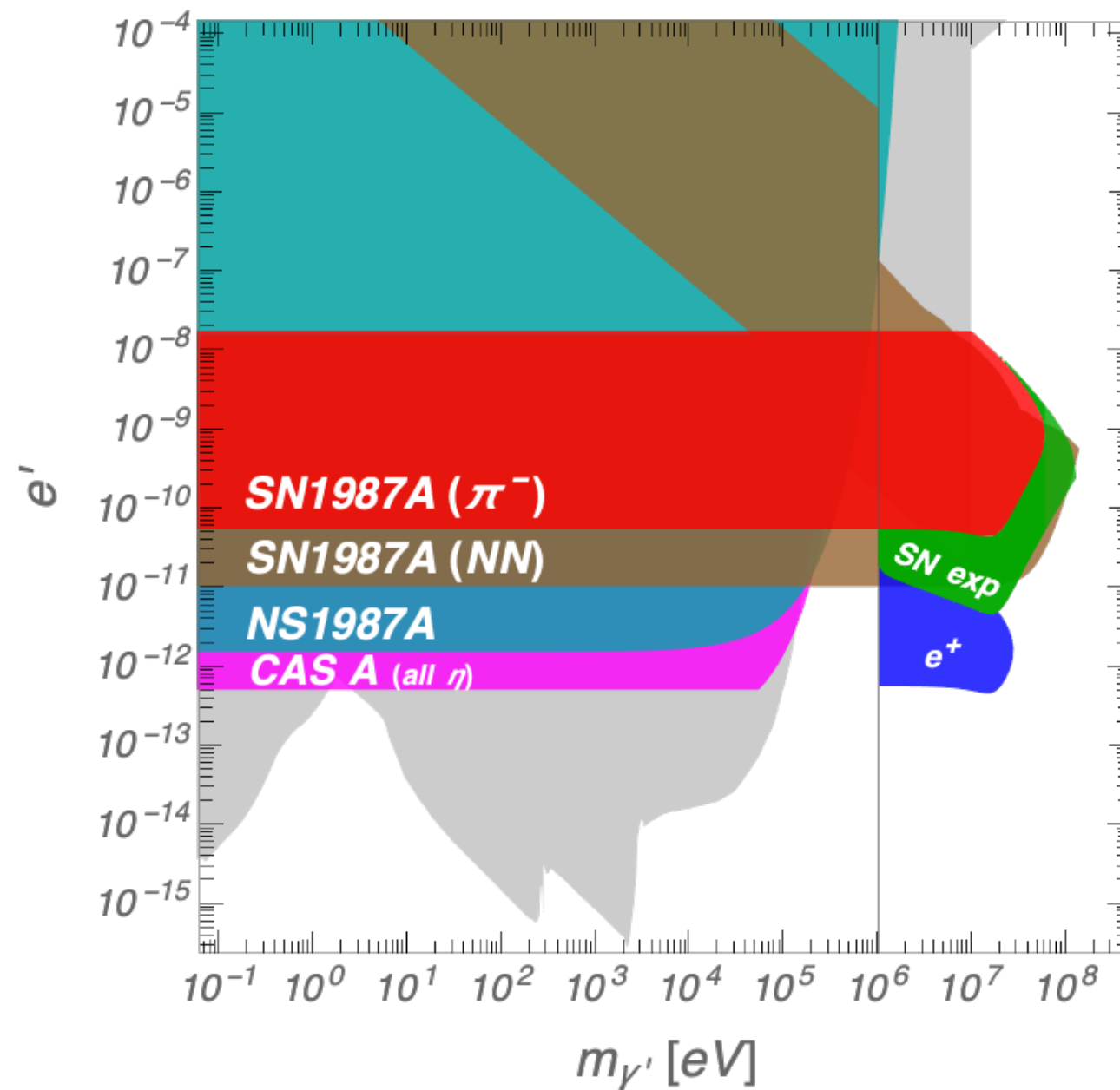




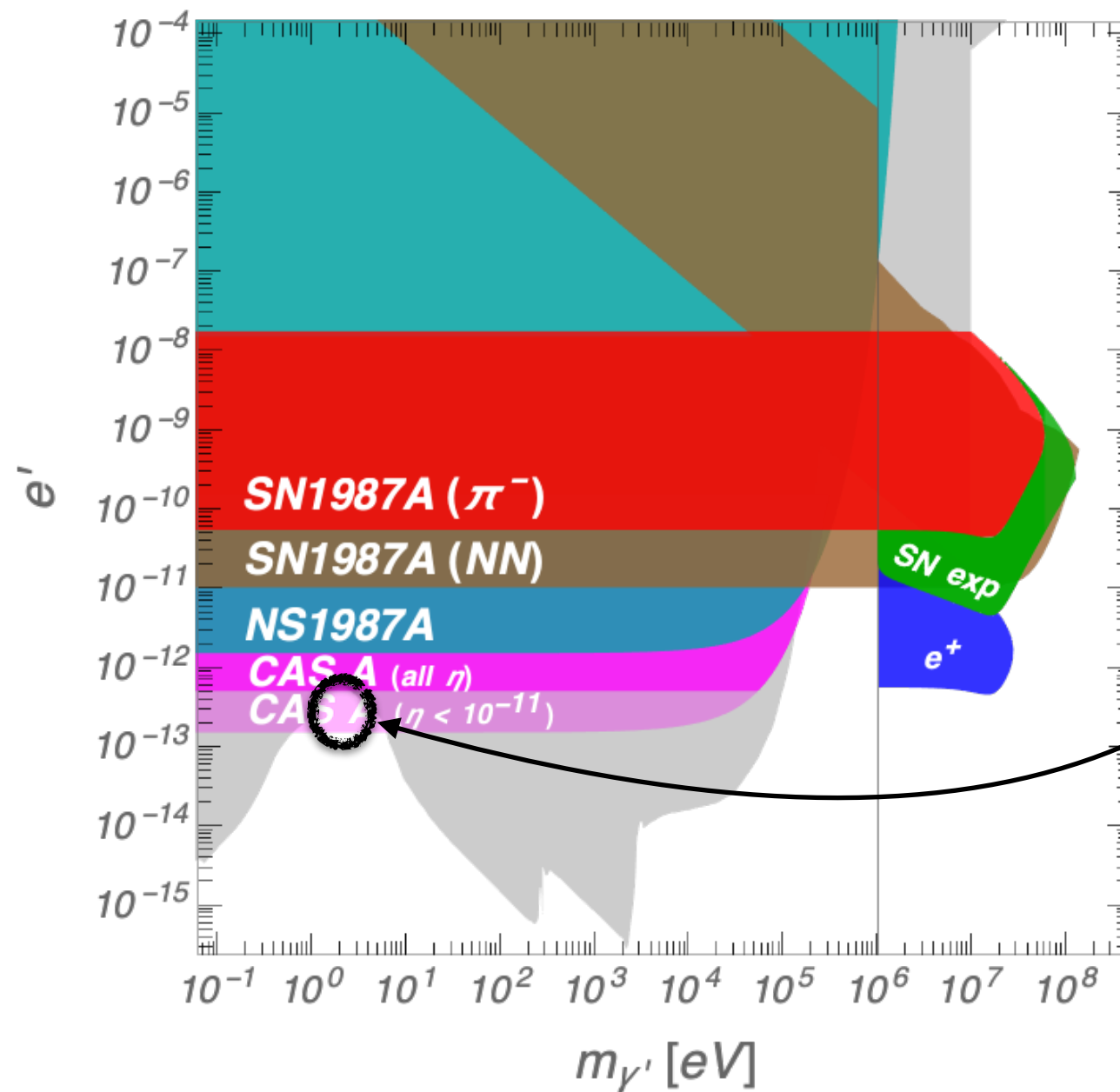
# $U(1)_{B-L}$ gauge boson constraints



# $U(1)_{B-L}$ gauge boson constraints

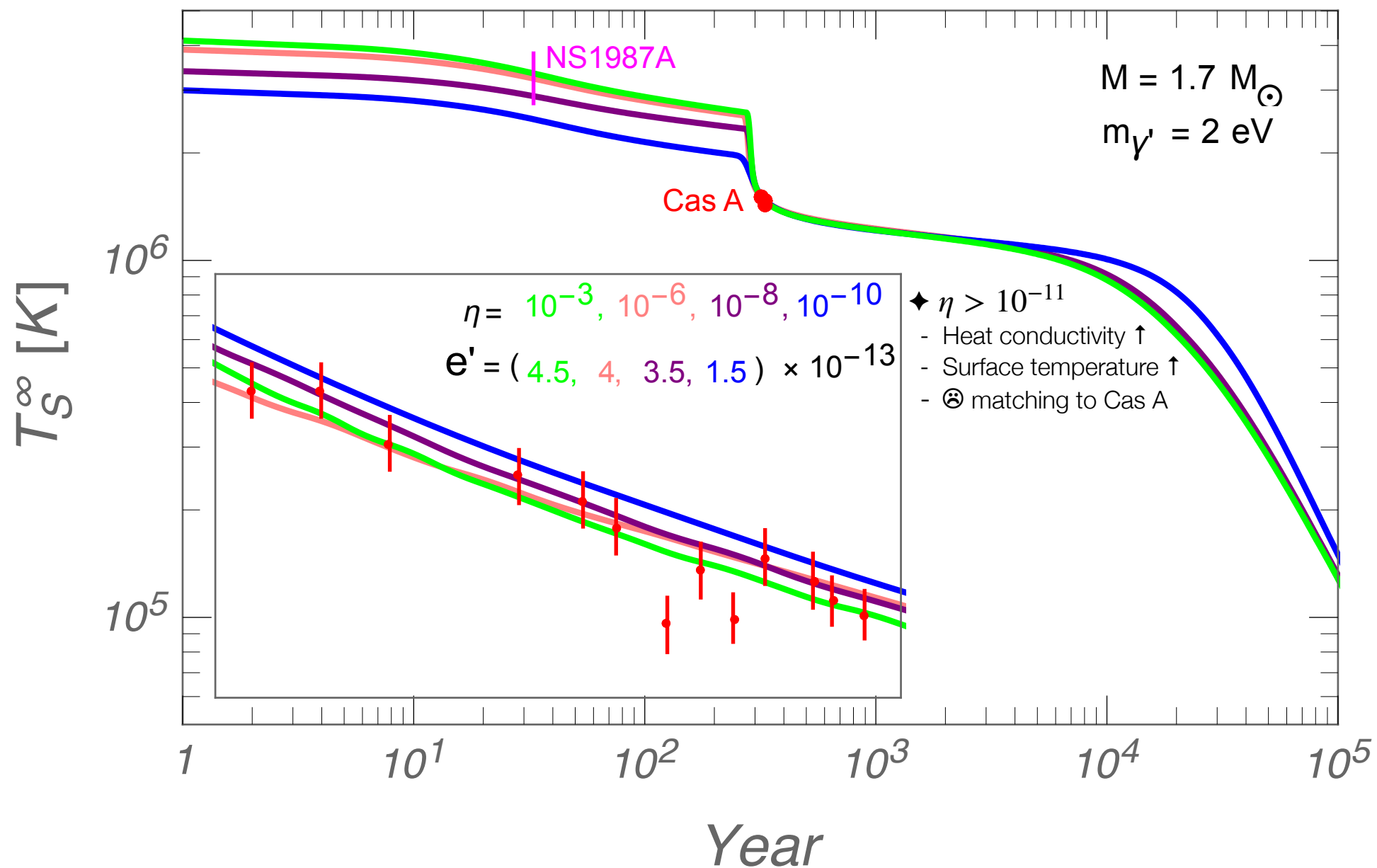


# $U(1)_{B-L}$ gauge boson constraints



Implications?

# $U(1)_{B-L}$ gauge boson implications



# Conclusion

- Neutron stars are attractive objects to probe fundamental interactions
- Stellar cooling argument on young neutron star cooling
- 1. Revisited SN1987A bound on  $U(1)_{B-L}$  gauge bosons
  - ChPT & Plasma effect
  - NN-brem (multipole rad)
  - Pion-induced

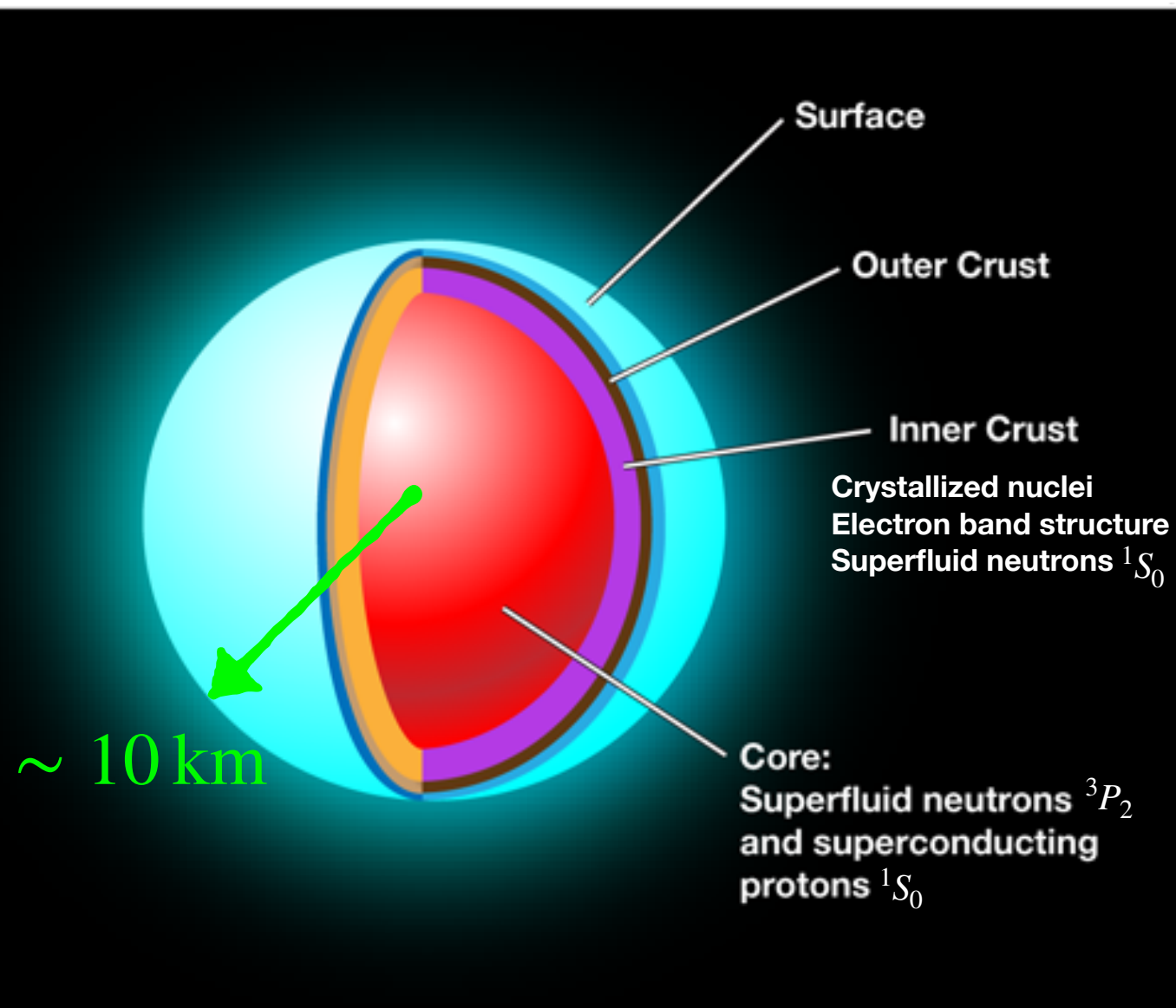
***An order of magnitude enhanced!***

- 2. Direct evidence of superfluidity from the rapid cooling of Cas A → minimal scenario with superfluidity fits well to datas (incl. NS1987A?)

***New stringent constraints on dark gauge bosons!***

# Back up

# Neutron stars



- Remnant core of a massive star after a type-II supernovae explosion
  - Neutrino signal for a few seconds (SN1987A)
  - Lack of prompt gamma-ray signal
- Young neutron star cooling
  - Neutrino (volume) emission
  - Photon (surface) emission  $> 10^{4-5}$  yr
  - Superfluidity?
- Etc.
  - Strong magnetic field
  - Gravitational source?

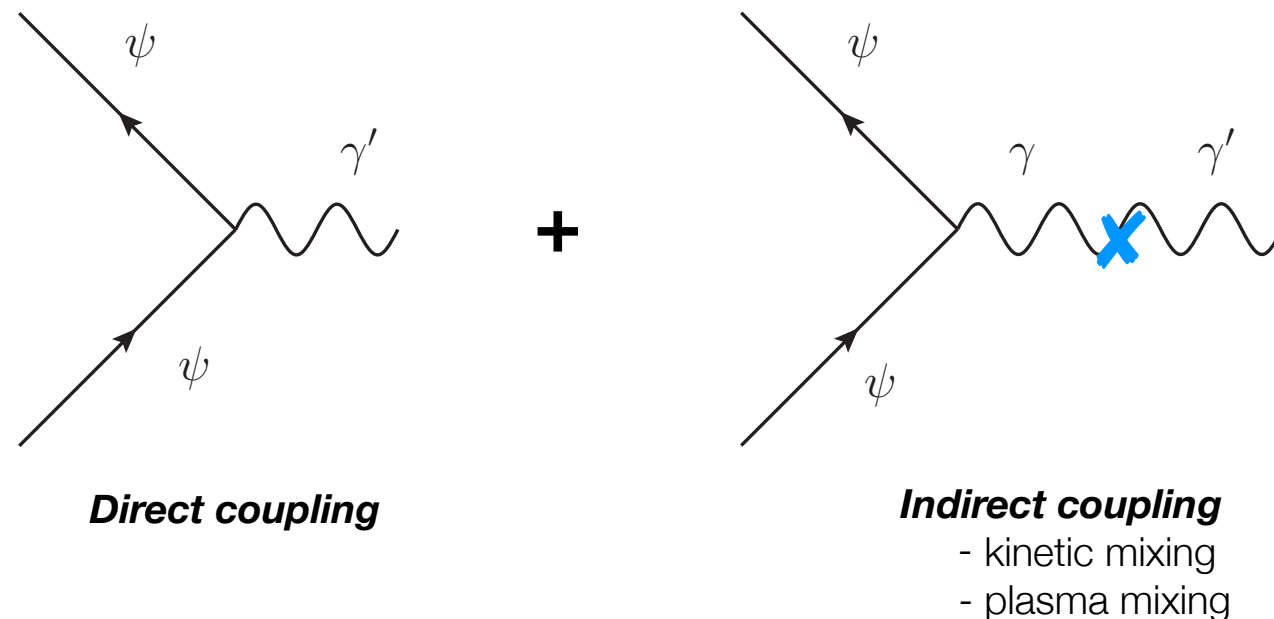
# In medium (+ kinetic) mixing

## Effective $\gamma'$ couplings within ChPT framework

$$\mathcal{L} = \underbrace{-\frac{1}{4}F_{\mu\nu}F^{\mu\nu}}_{\text{EM field strength}} - \underbrace{\frac{1}{4}X_{\mu\nu}X^{\mu\nu}}_{\text{DP field strength}} + \underbrace{\frac{\varepsilon}{2}F_{\mu\nu}X^{\mu\nu}}_{\text{kinetic mixing}} - \underbrace{\frac{1}{2}m_{\gamma'}^2 X_\mu X^\mu}_{\gamma' \text{ mass}} + \underbrace{eA_\mu J_{\text{EM}}^\mu}_{\text{EM current}} + \underbrace{e'A'_\mu J'^\mu}_{\gamma' \text{ current}}$$

- hadron (ChPT)  
- lepton

## Effective $\gamma'$ couplings: direct + indirect (mixing)

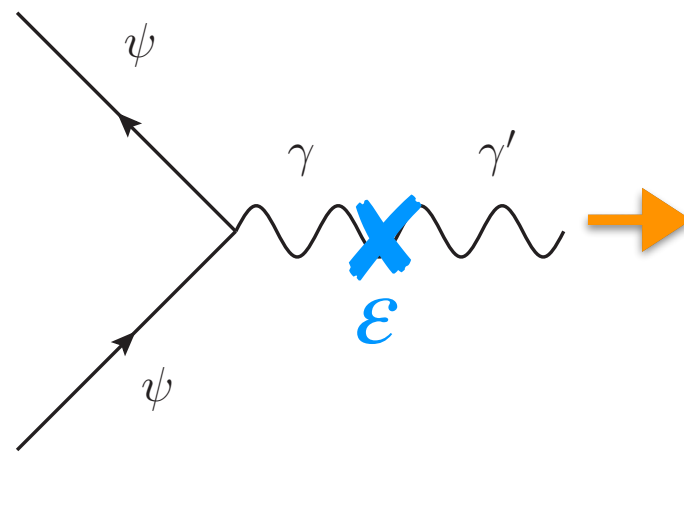




# Dark photon (vector portal)

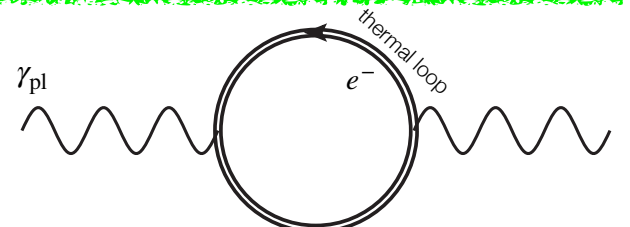
$$\mathcal{L}_{\text{eff}} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} - \frac{1}{4}X_{\mu\nu}X^{\mu\nu} + \frac{\epsilon}{2}F_{\mu\nu}X^{\mu\nu} - \frac{1}{2}\overset{\langle J_{\text{EM}}^\mu J_{\text{EM}}^\nu \rangle}{\Pi^{\mu\nu}}A_\mu A_\nu - \frac{1}{2}m_{\gamma'}^2 X_\mu X^\mu$$

EM field strength
DP field strength
kinetic mixing
plasmon
DP mass



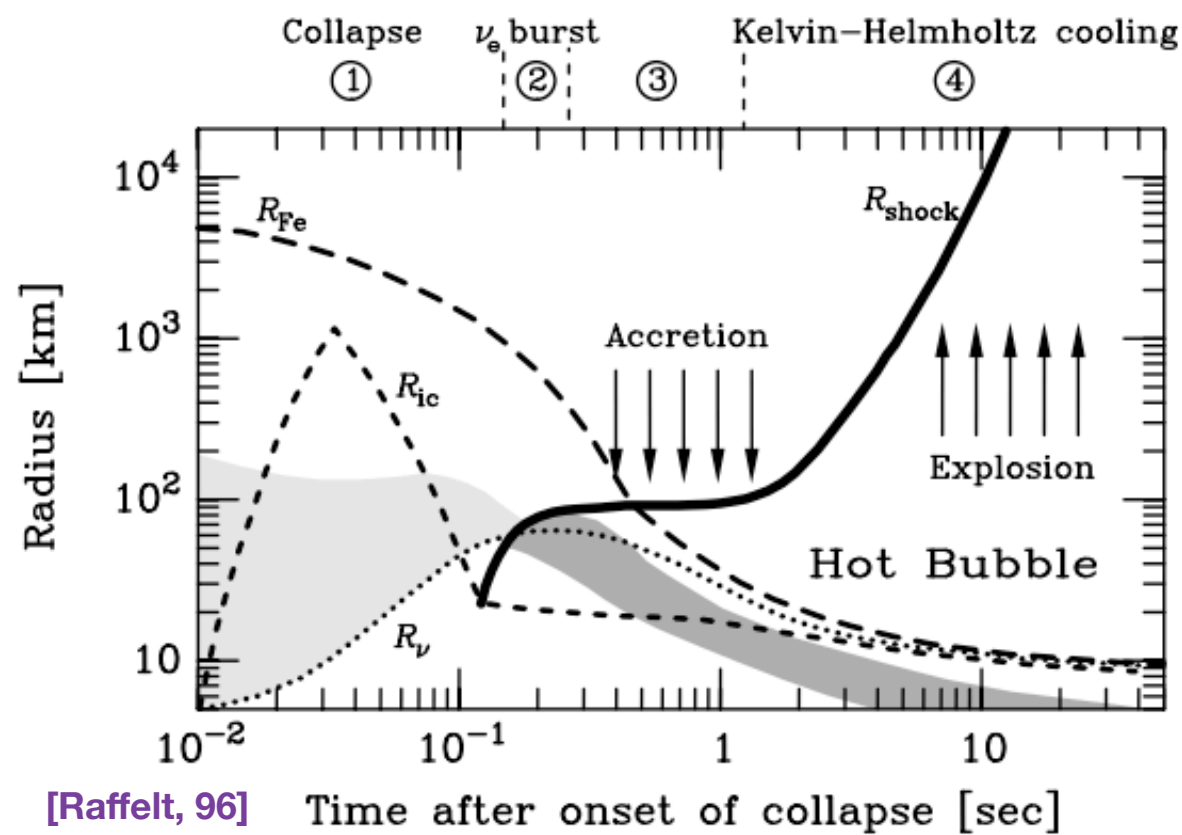
$$\mathcal{M} = \underbrace{\frac{\epsilon e m_{\gamma'}^2}{m_{\gamma'}^2 - \pi_{\text{T,L}}}}_{e_{\text{eff}}^\psi} \left[ J_{\text{EM}}^\mu \right] \epsilon'_\mu$$

\*medium effect



$$\pi_{\text{T,L}} = \begin{cases} \text{transverse} & \omega_{\text{pl}}^2 \left( 1 + \frac{1}{2} G(v_*^2 k^2 / \omega^2) \right) \\ \text{longitudinal} & \omega_{\text{pl}}^2 \frac{m_{\gamma'}^2}{\omega^2} \frac{1 - G(v_*^2 k^2 / \omega^2)}{1 - v_*^2 k^2 / \omega^2} \end{cases}$$

# Raffelt's criterion for SN1987A



- The “**delayed explosion scenario**” ( $D_{\nu}M$ )
  - Thermal relaxation of the proto-neutron star as the source of neutrino emission for  $> 5\text{sec}$
  - the energy deposition of the neutrino flow causes the explosion ( $\mathcal{O}(1)\%$  of total neutrino bulk energy)

➡  $L_x \lesssim 3 \times 10^{52} \text{ erg s}^{-1}$

# SN1987A bound

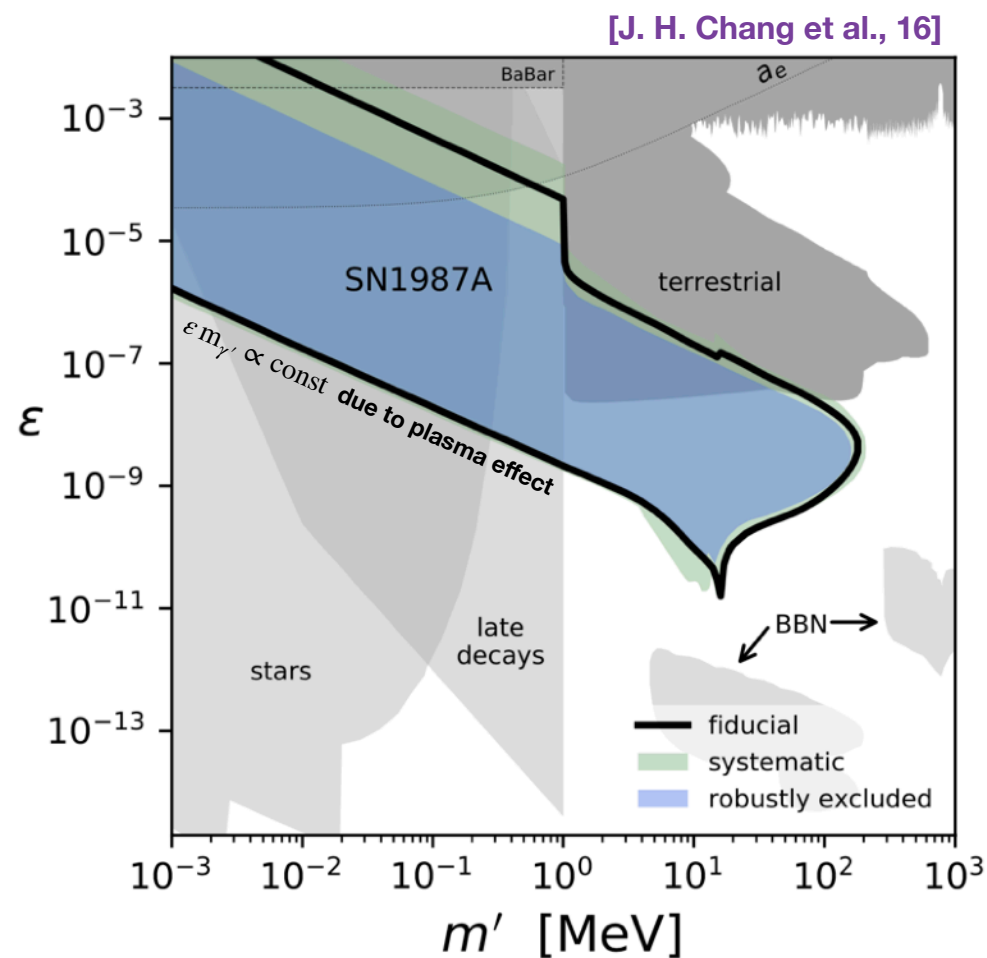
## ◆ Dark photon

$$e_{\text{eff}}^n = 0$$

$$e_{\text{eff}}^p = \epsilon e \frac{m_{\gamma'}^2}{m_{\gamma'}^2 - \pi_{T,L}}$$



Dipole radiation



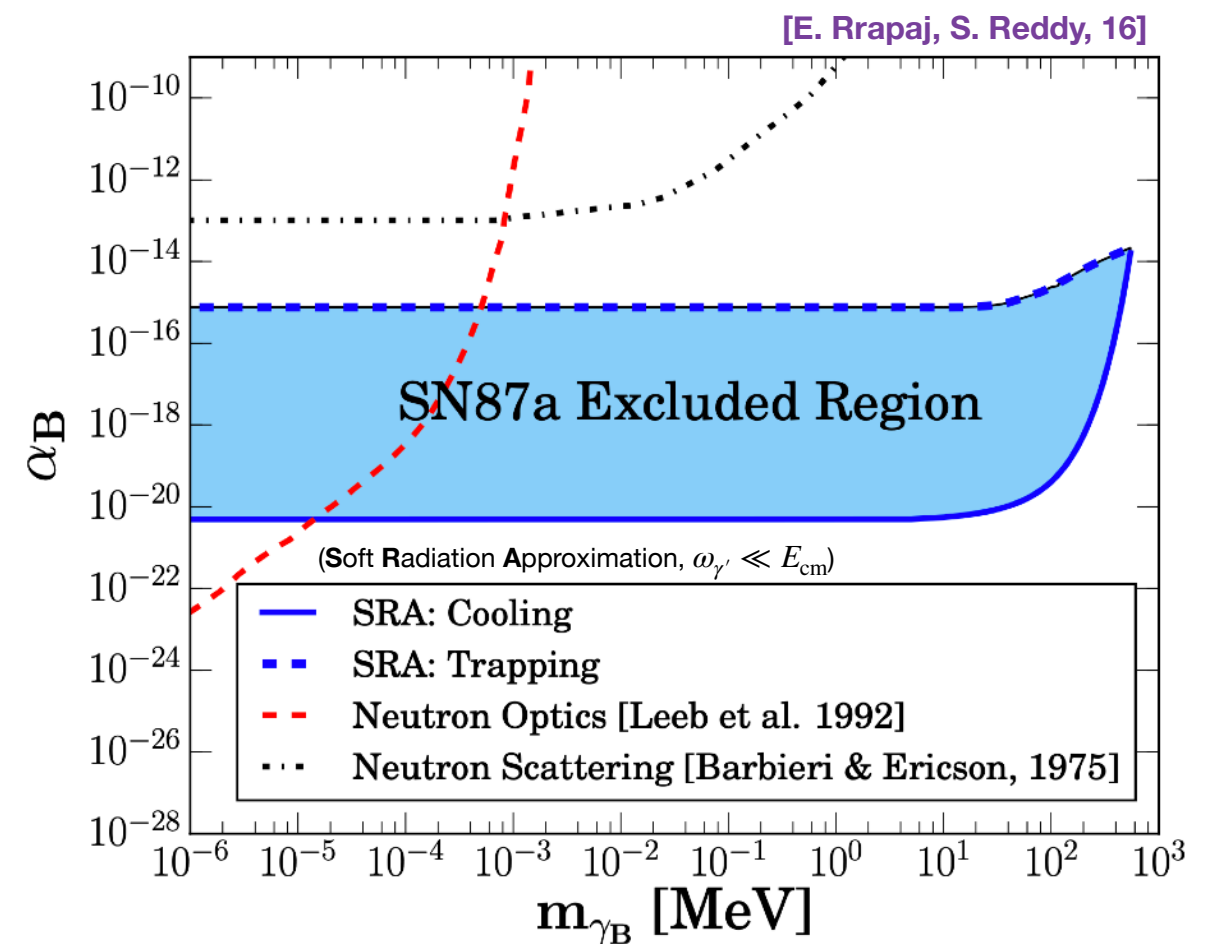
## ◆ $U(1)_B$

$$e_{\text{eff}}^n = 1$$

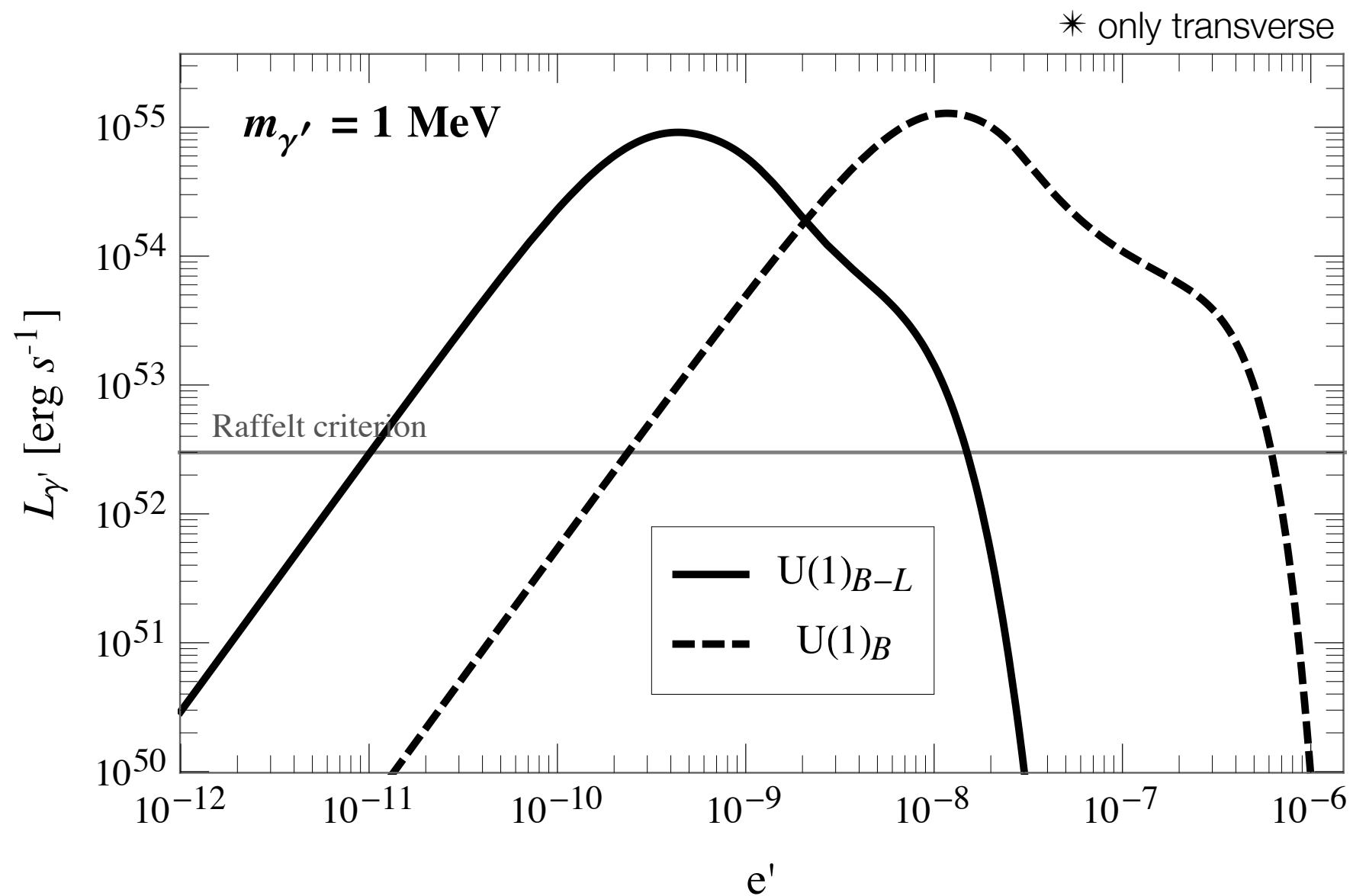
$$e_{\text{eff}}^p = 1$$



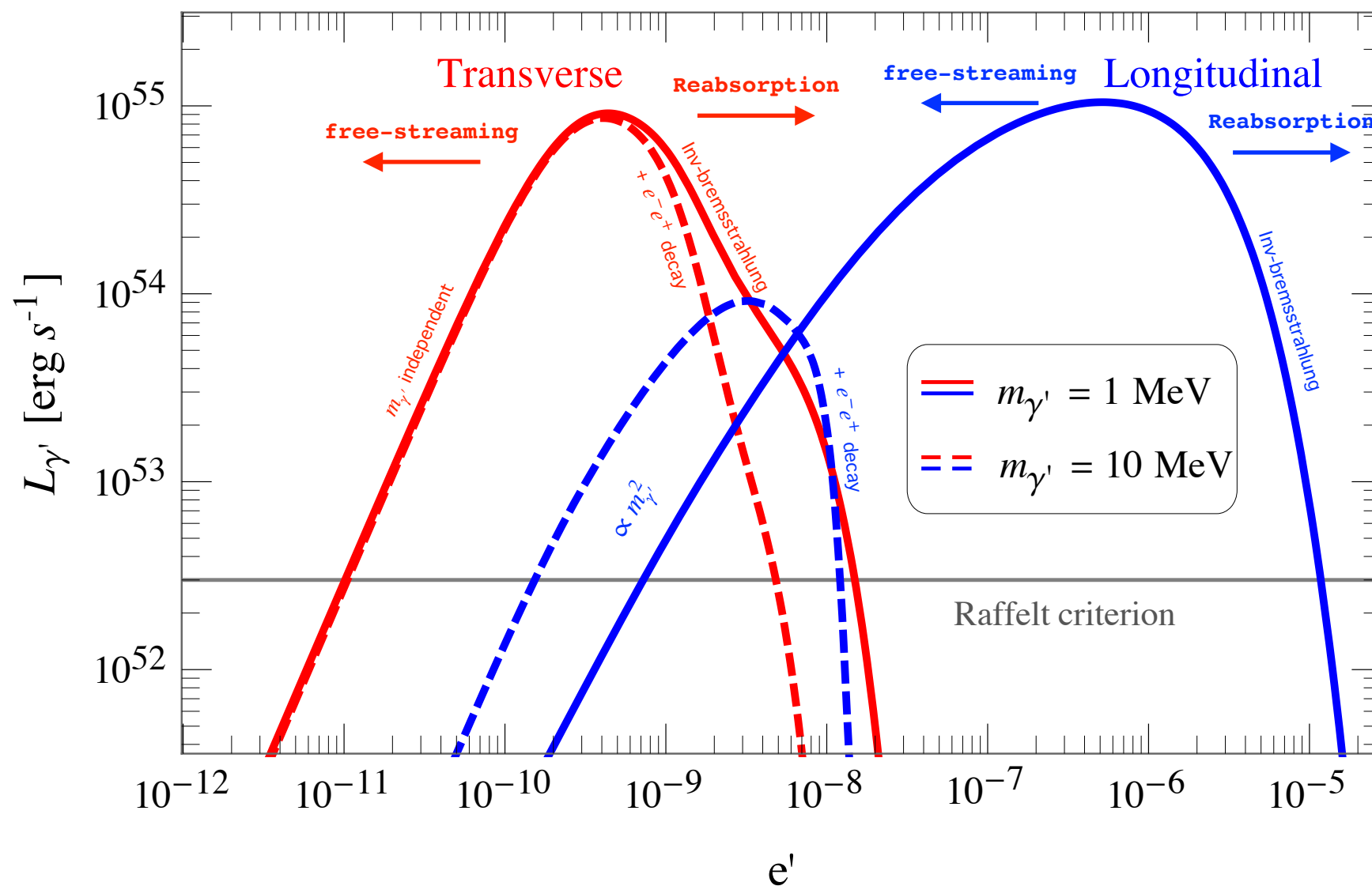
Quadrupole radiation



# $U(1)_{B-L}$ vs $U(1)_B$ in SN1987A



# $U(1)_{B-L}$ luminosity in SN1987A



# *Dark gauge boson via Pion-induced reaction*

- Thermal  $\pi^-$  abundance
- Dominant production during supernovae @  $\omega \sim m_\pi$
- Constraints from
  - i) SN1987A (cooling)
  - ii) SN explosion energy
  - iii) galactic  $e^+$  injection
  - iv) Absence of prompt SN  $\gamma$ -ray

# Thermal $\pi^-$ -induced production

- Large  $\mu_n - \mu_p$  & temperature (  $\sim 30 \text{ MeV}$  )

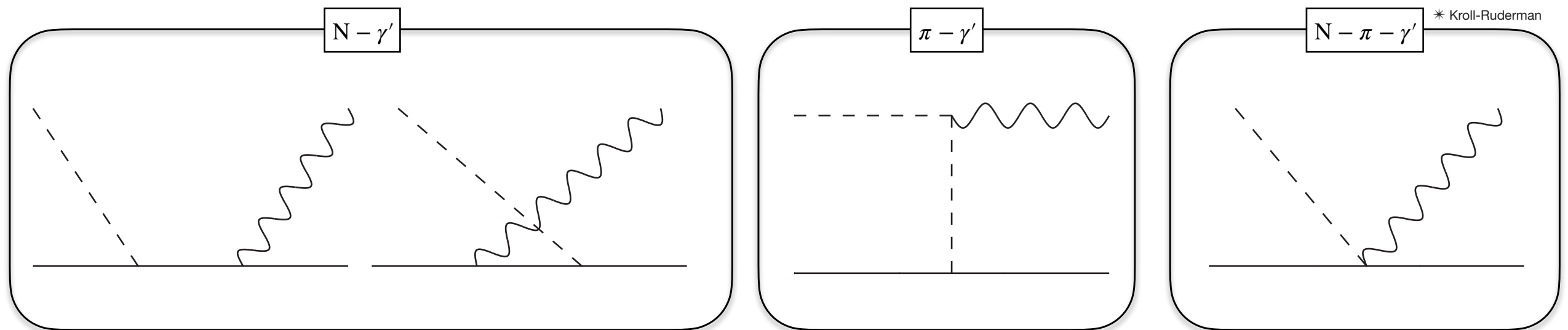
- source for negatively charged particles (e.g.,  $e^-$ ,  $\mu^-$ ,  $\pi^-$ )
- thermal pion abundance with  $n_{\pi^-}/n_B \sim \mathcal{O}(1)\%$  [B. Fore, S. Reddy, 19]

- Strong interactions with nucleons

- modification of thermodynamics
- effective dispersion relation of  $\pi^-$

$\Rightarrow$  **Viral expansion**

- Diagrams of Compton-like processes



# Phenomenological implications?

## ● Cooling argument on SN1987A

- comparable constraints to the results based on NN-bremsstrahlung
- If absorbed within  $r_{\text{far}} \sim 10^{2-3}$  km, no efficient cooling
- $L_{\gamma'} e^{-(m_{\gamma'}/\omega_{\gamma'}) \Gamma_{\gamma'} v_{\gamma'} r_{\text{far}}} < 10^{19} \text{ erg g}^{-1} \text{ s}^{-1}$  “Raffelt criterion”

## ● High energy galactic $e^+$ injection [W. DeRocco, et al., 19] [F. Calore, et al., 21]

- No excess of the galactic gamma ray flux  $\Rightarrow$  constrain new  $e^+$  sources [J. F. Beacom, et al., 06]  
[P. Sizun, et al., 06]
- $N_{\gamma'} e^{-(m_{\gamma'}/\omega_{\gamma'}) \Gamma_{\gamma'} v_{\gamma'} r_{\text{esc}}} < 10^{52} e^+$  for  $\omega_{e^+} = 100 \text{ MeV}$  &  $r_{\text{esc}} \sim 10^9 \text{ km}$  (type-II SN)

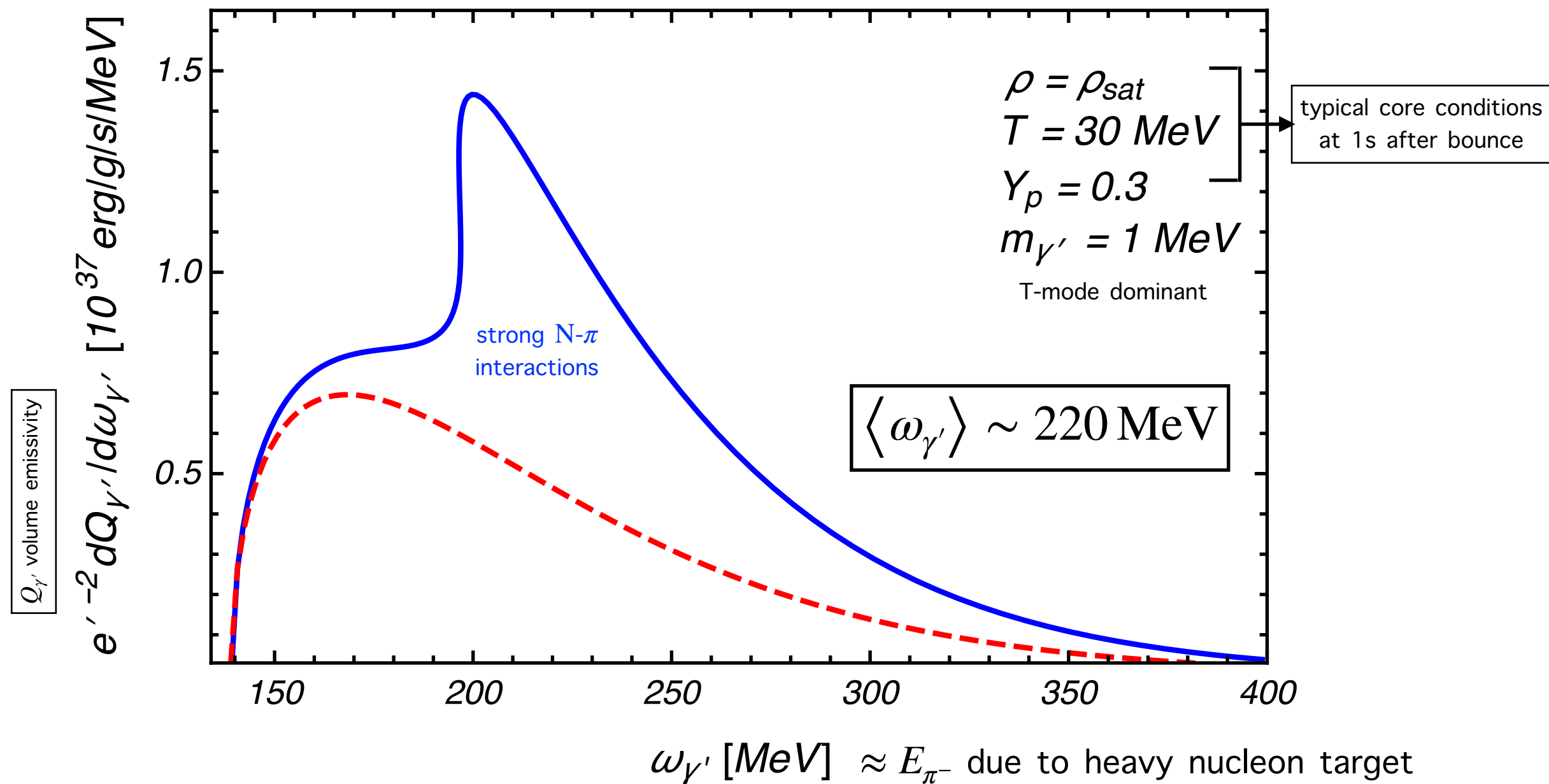
## ● Trigger of supernova explosion [S. W. Falk et al., 78] [A. Sung, et al., 19] [A. Caputo et al., 22]

- trapping or reabsorption within  $r_{\text{far}} < r < r_{\text{esc}}$  can transfer energy
- inferred SN explosion energy  $< 10^{51} \text{ erg} \Rightarrow$  constrain new energy transfer sources [K. Nomoto, et al., 13]  
[S. W. Bruenn, et al., 14]
- $L_{\gamma'} \left( e^{-(m_{\gamma'}/\omega_{\gamma'}) \Gamma_{\gamma'} v_{\gamma'} r_{\text{far}}} - e^{-(m_{\gamma'}/\omega_{\gamma'}) \Gamma_{\gamma'} v_{\gamma'} r_{\text{esc}}} \right) < 10^{17} \text{ erg g}^{-1} \text{ s}^{-1}$

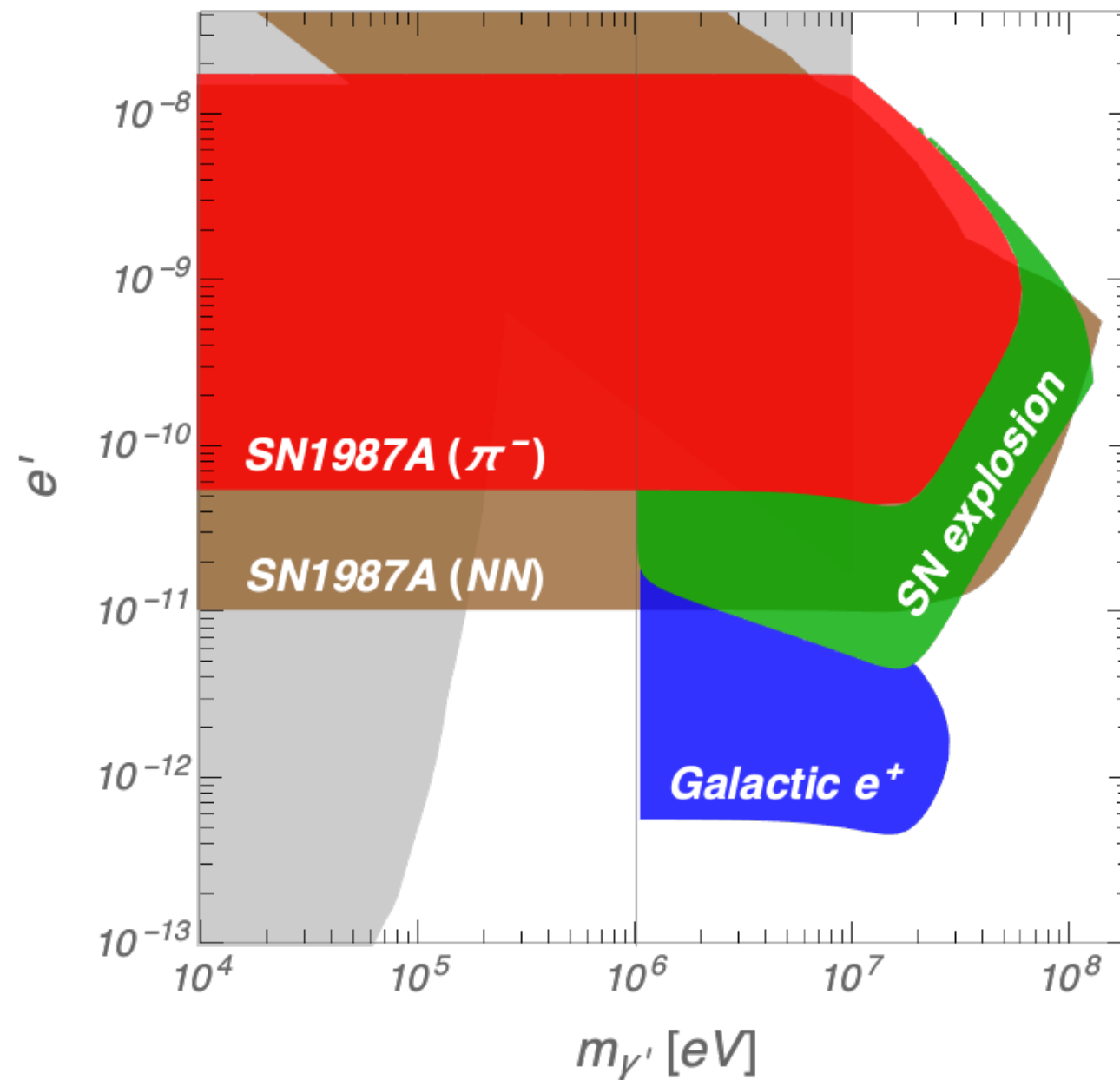
$m_{\gamma'} > 2m_e$



# $\pi^-$ -induced $\gamma'$ spectrum

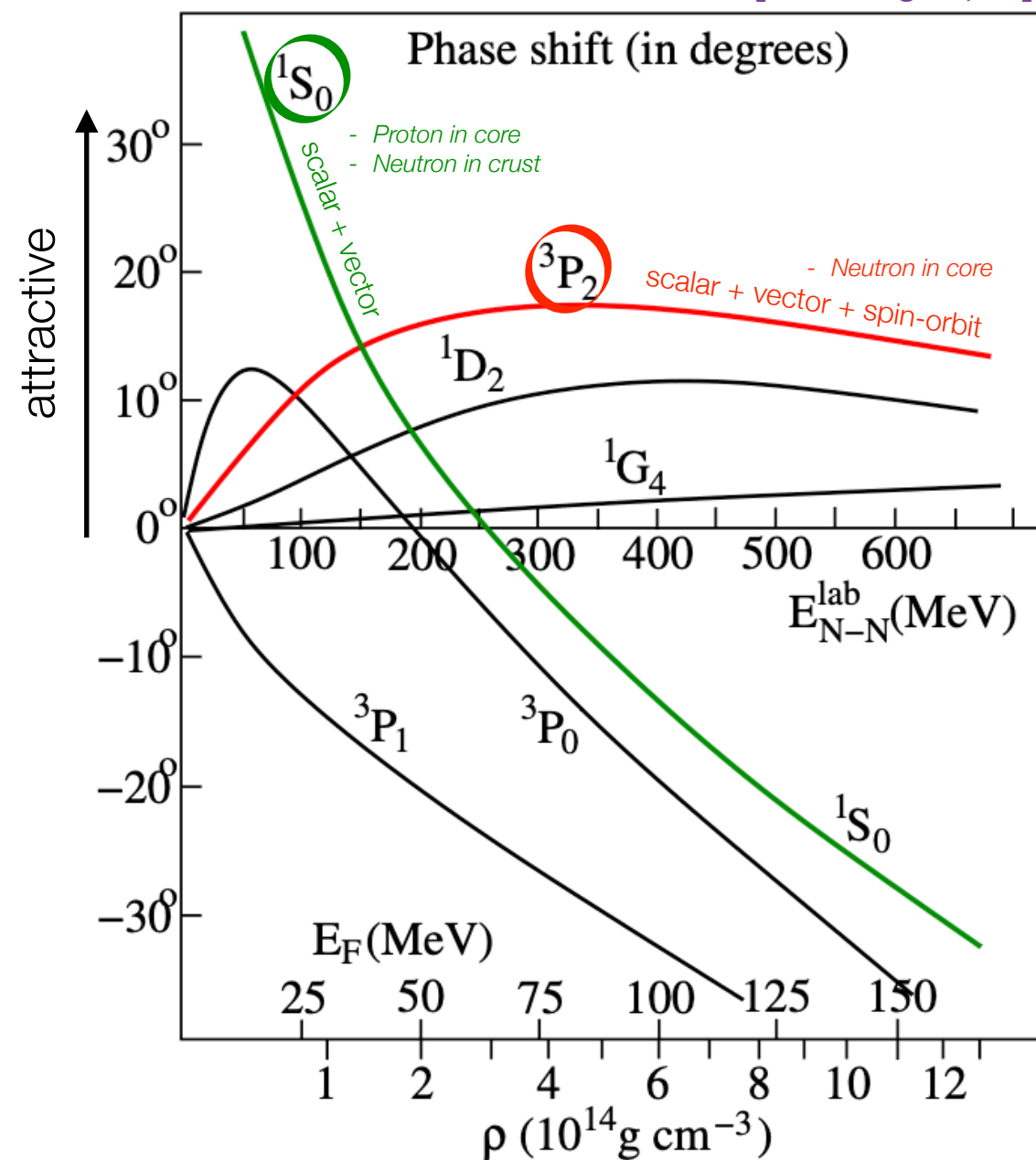


# $U(1)_{B-L}$ gauge boson constraints

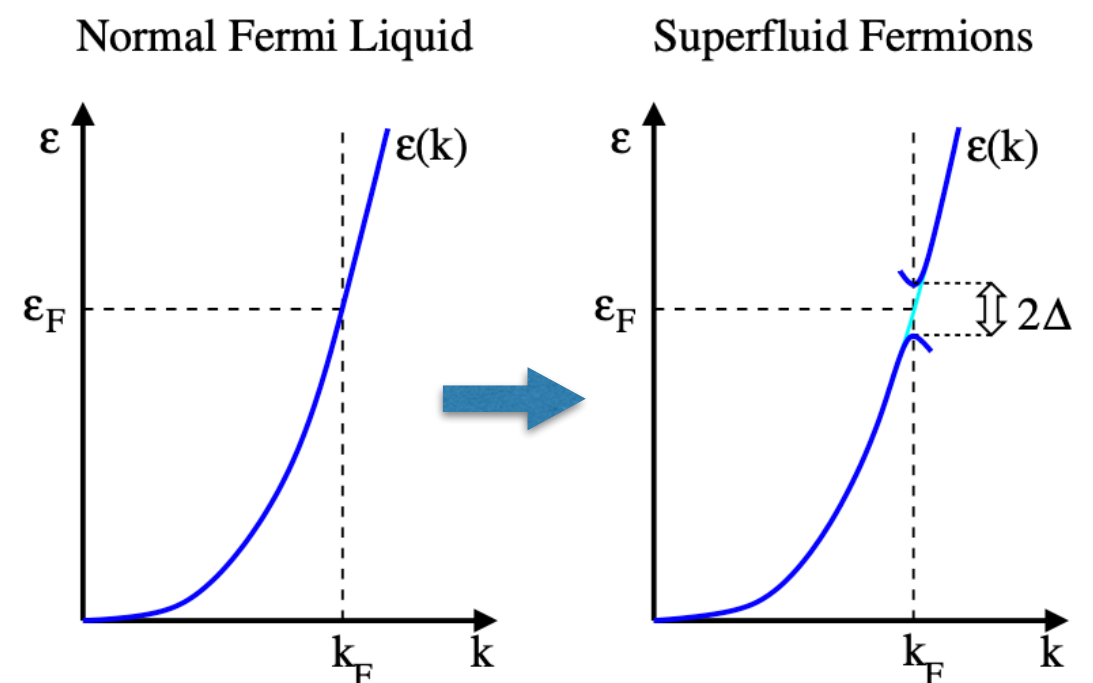


# Superfluidity of nucleon

[R. Tamagaki, 70]

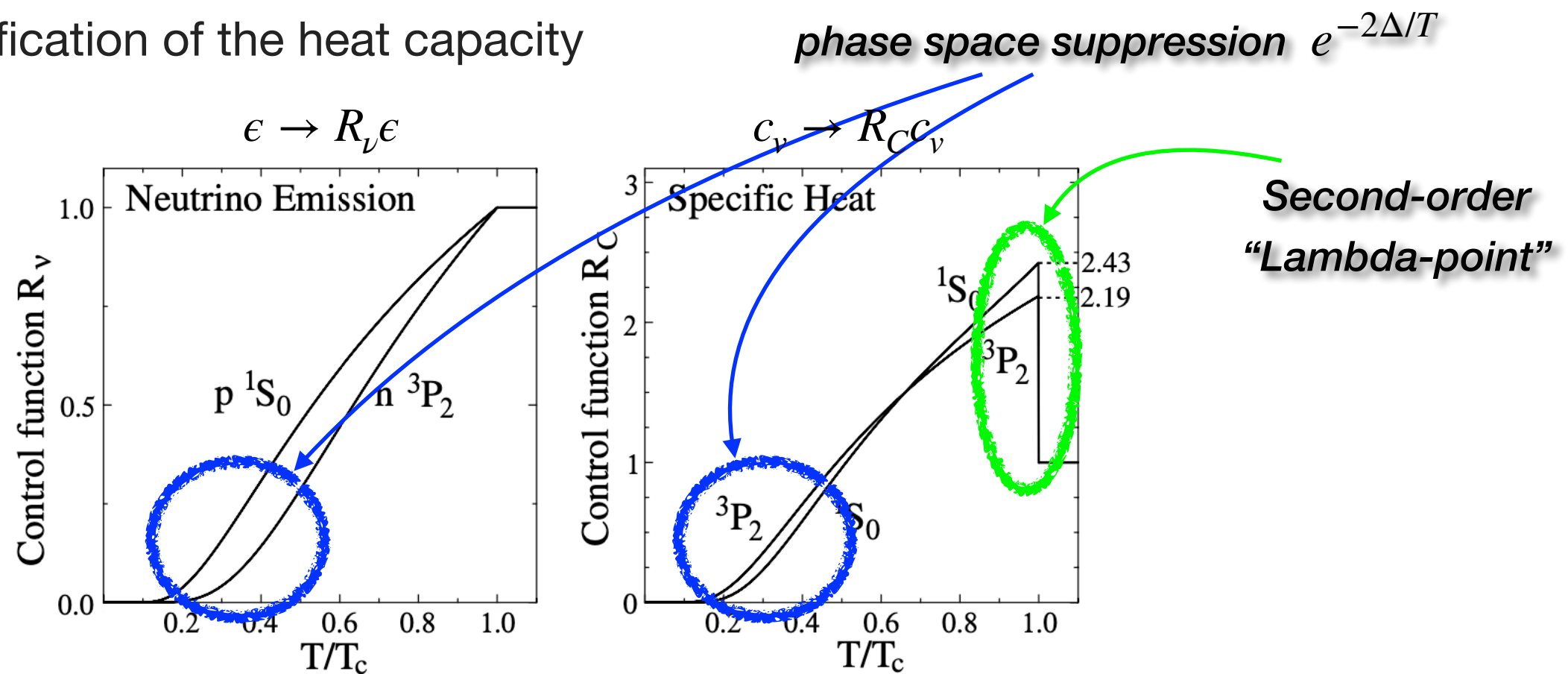


- N-N interaction by exchange of meson
  - Scalar (attractive)
  - Vector (repulsive)
  - Spin-orbit (attractive or repulsive)
- Below  $T_c$ , nucleons form the pairs
- Appearance of a gap in the momentum spectrum



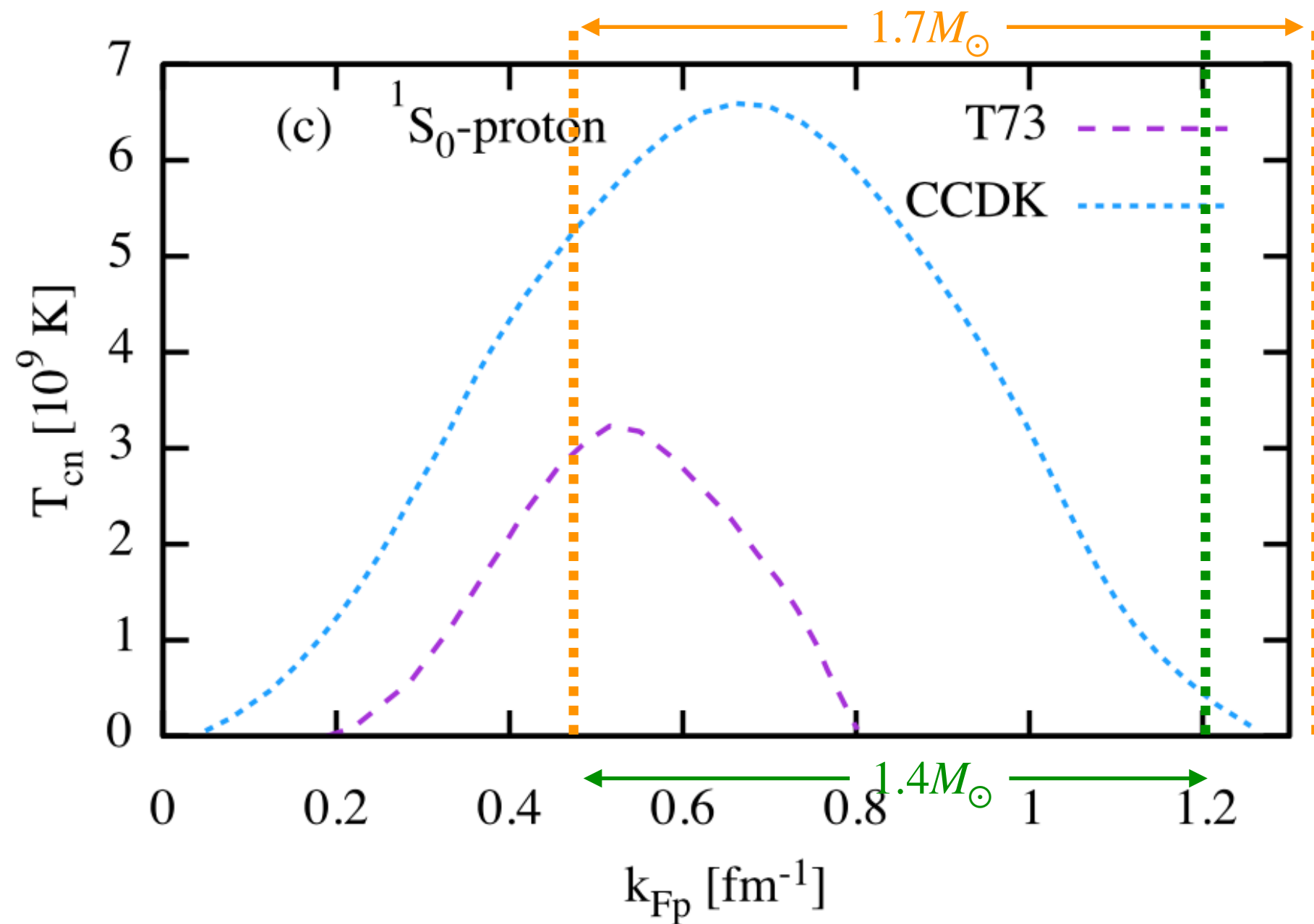
# Effect of pairing for NS cooling

1. Suppression of the emissivity
2. Modification of the heat capacity

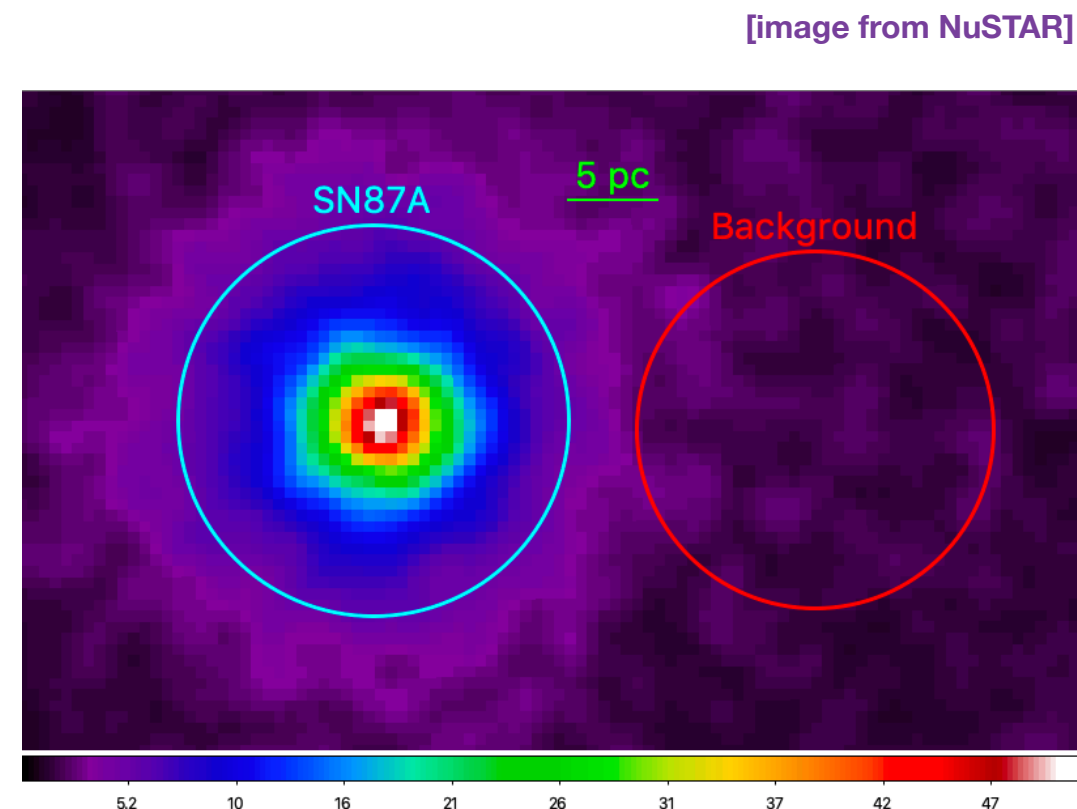
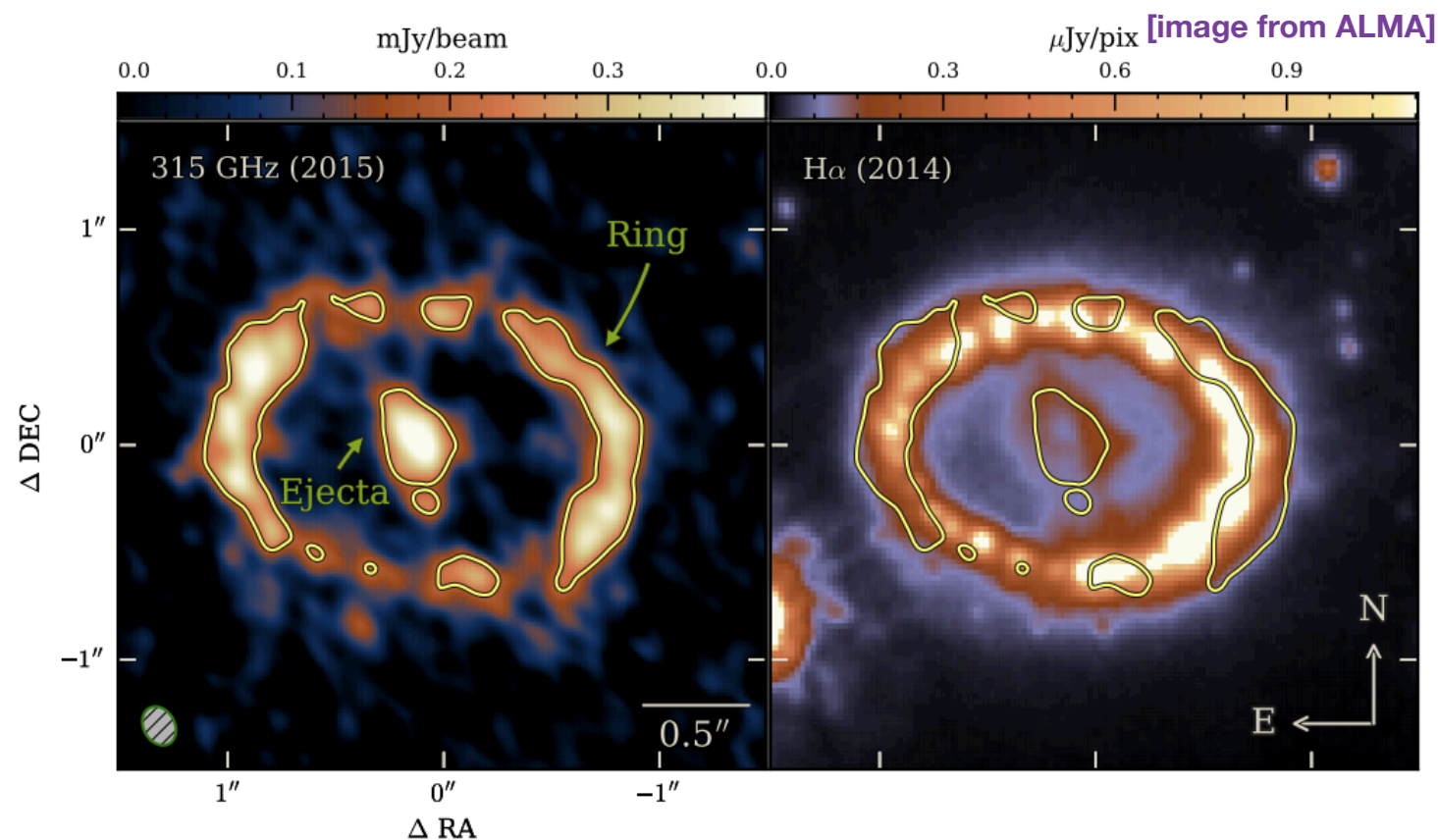


3. Triggering of the “pair breaking and formation” (PBF) emission

# Proton singlet pairing models



# Remnant of SN1987A?

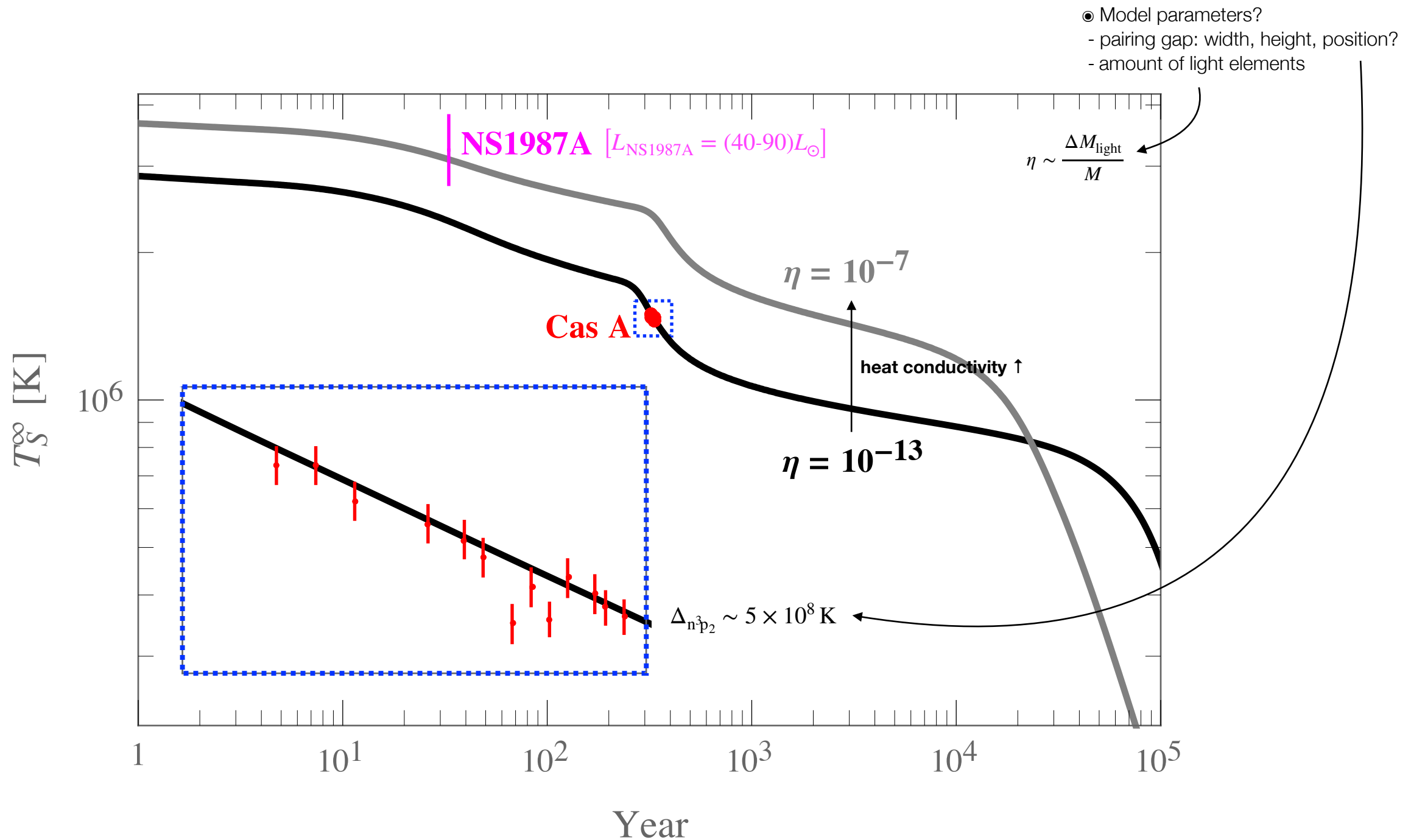


- Presence of a compact source in the remnant is indicated.
  - heated dust clump (at the predicted position) + the hard X-ray spectrum
  - a neutron star rather than a black hole is likely to have formed

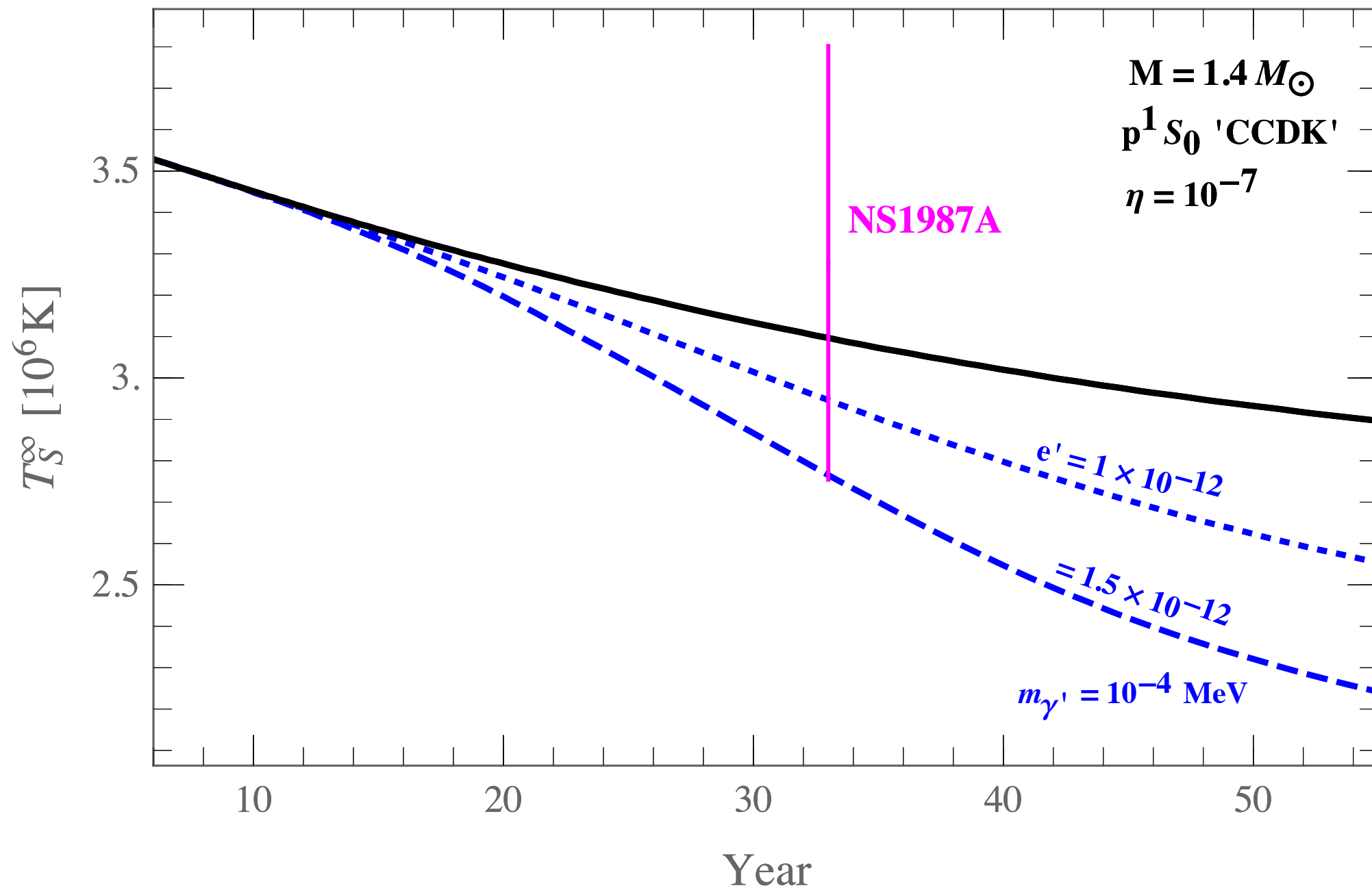


**'NS1987A'**

# Minimal cooling scenario



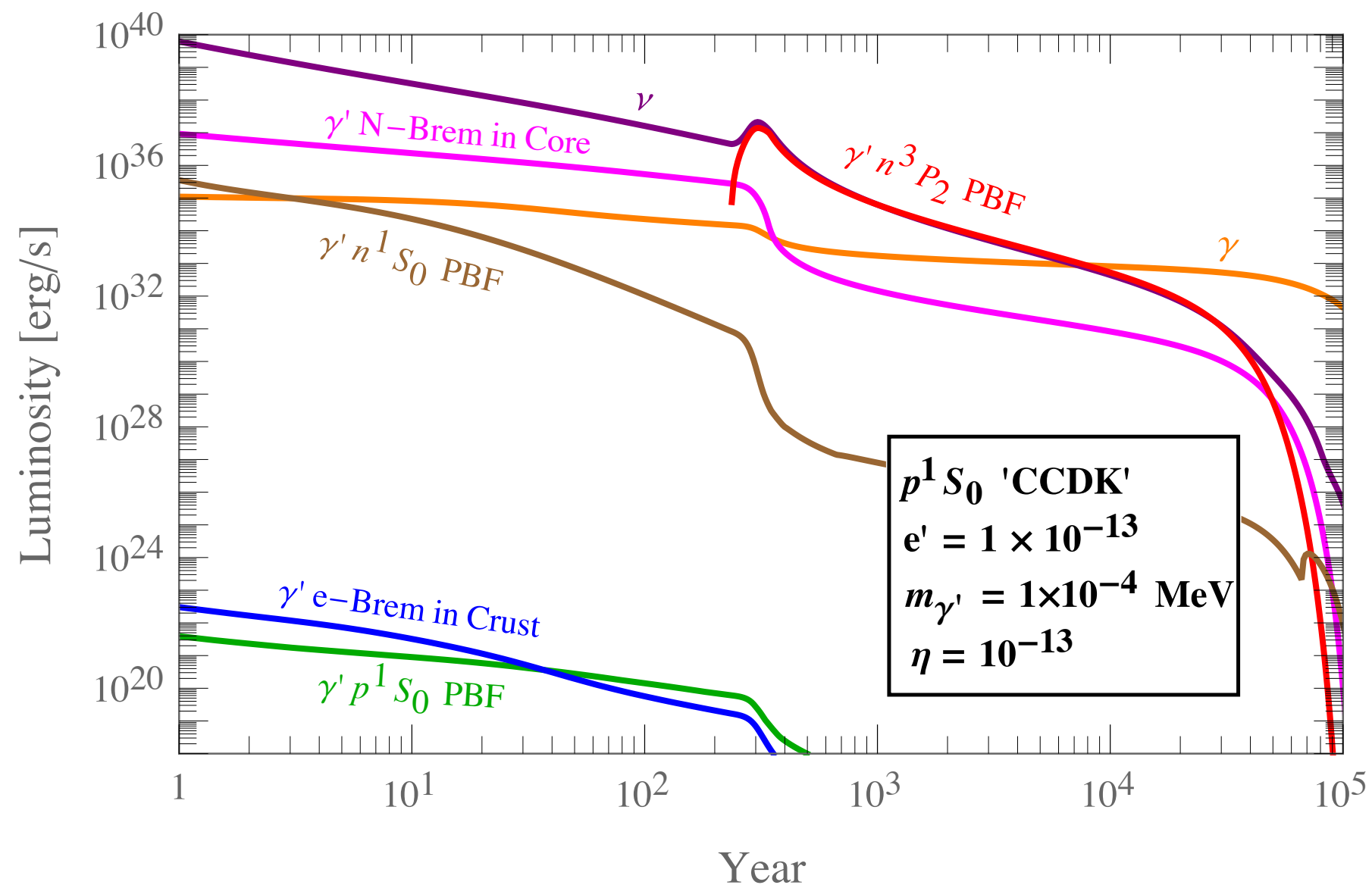
# $U(1)_{B-L}$ gauge boson cooling





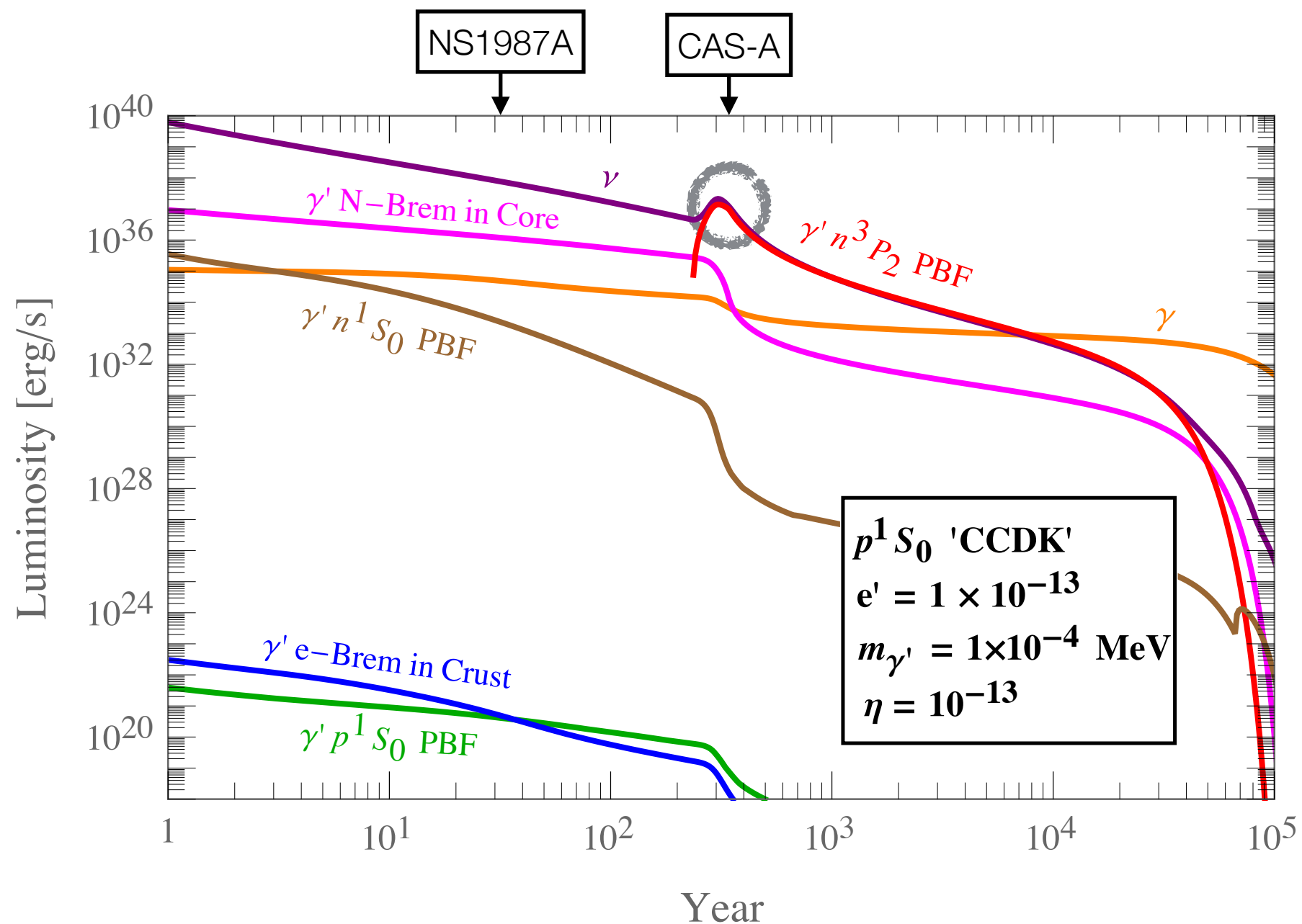
# Luminosity of $U(1)_{B-L}$ (vs. $\nu$ )

\* Stellar cooling argument



# Luminosity of $U(1)_{B-L}$ (vs. $\nu$ )

\* Stellar cooling argument

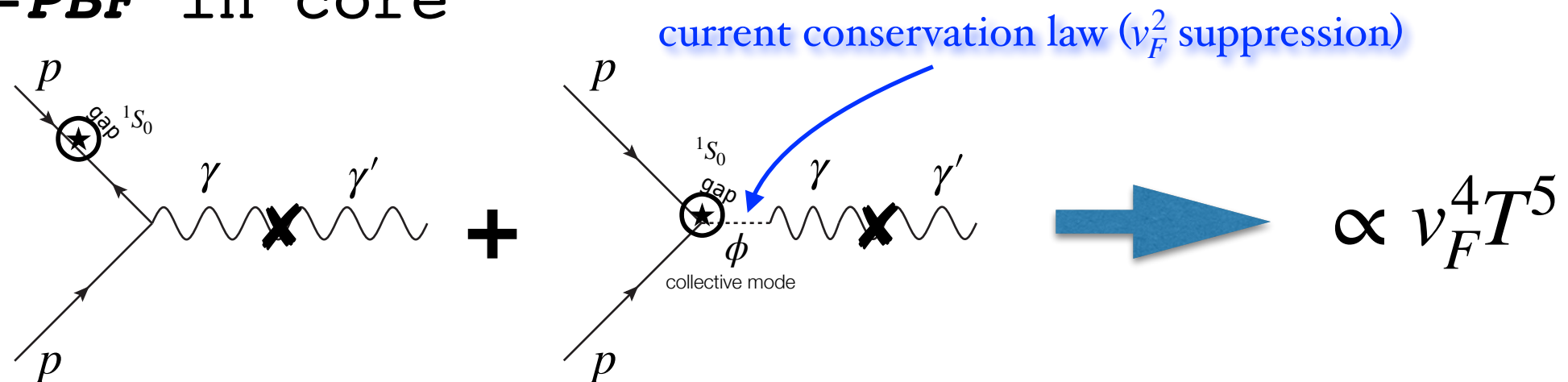


# Dark photon in NS

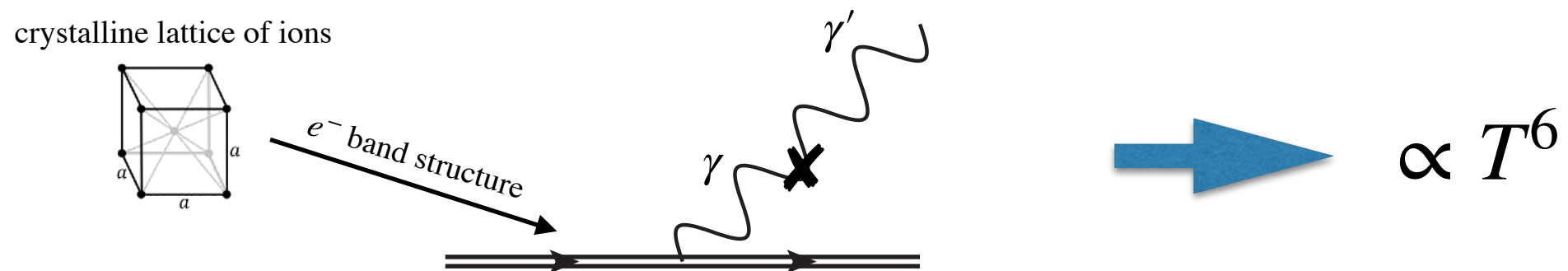
© **Longitudinal** emission is dominant

$$\epsilon_L \propto \epsilon^2 m_{\gamma'}^2 T^2 \gg \epsilon_T \propto \epsilon^2 m_{\gamma'}^4$$

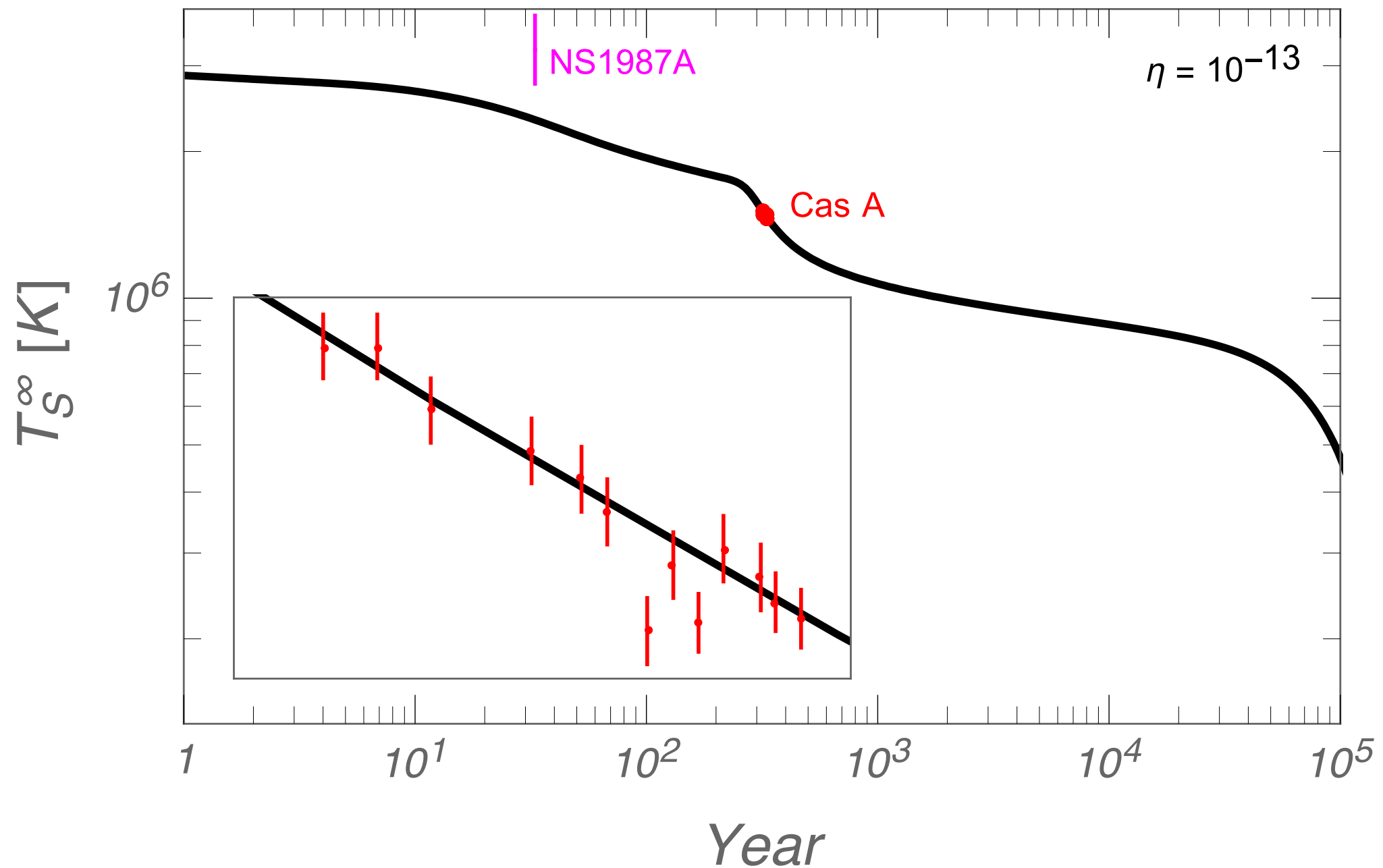
## 1. p-**PBF** in core



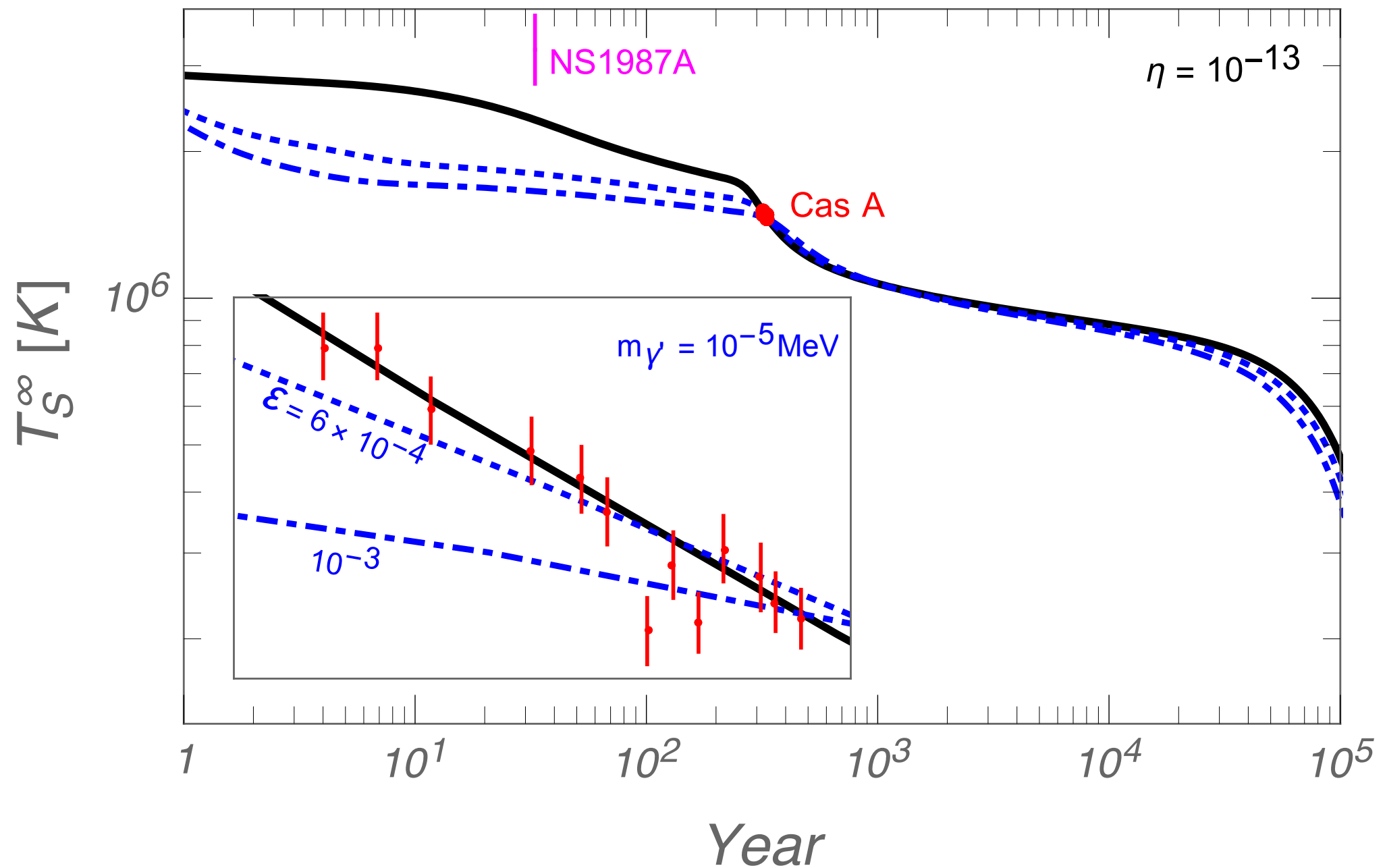
## 2. **e-Bremsstrahlung** in crust



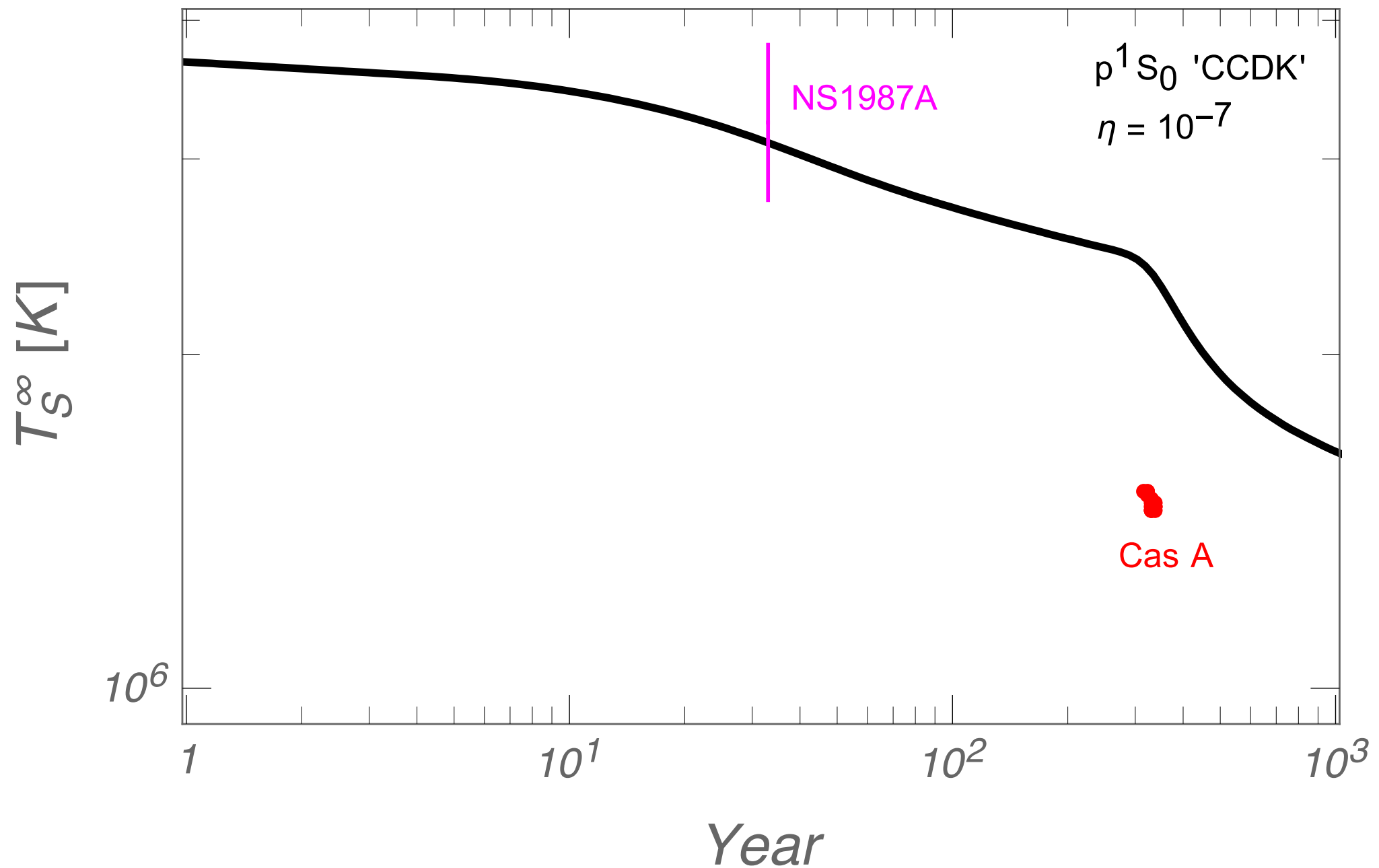
# Dark photon cooling



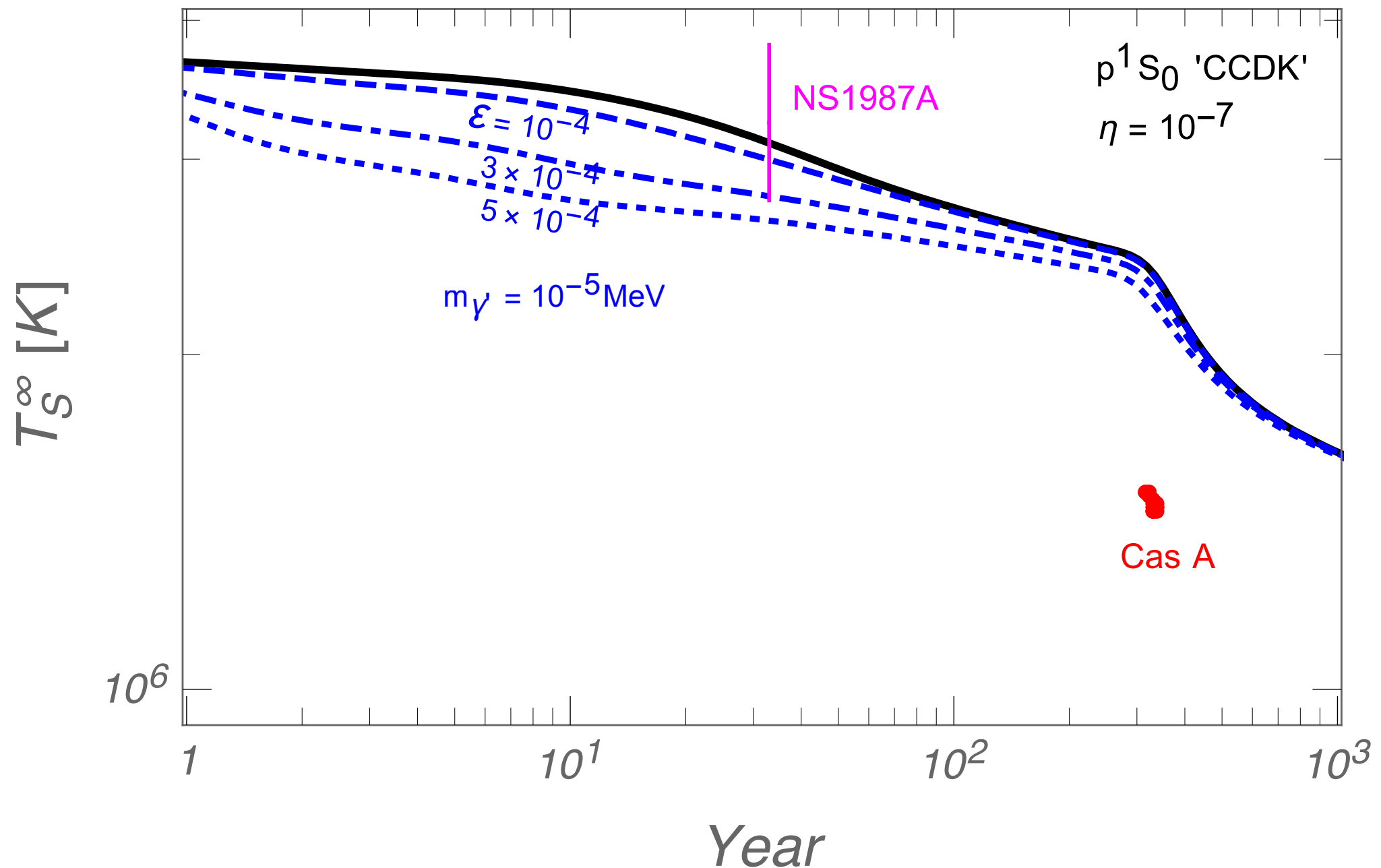
# Dark photon cooling



# Dark photon cooling

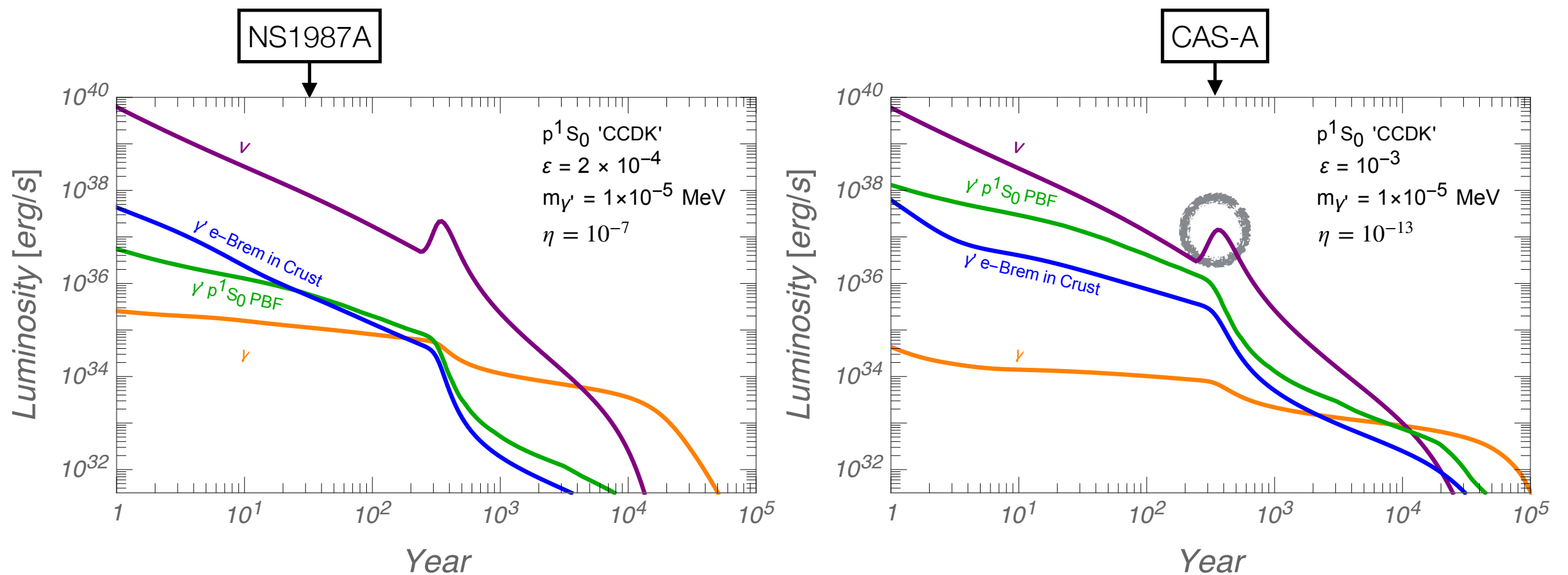


# Dark photon cooling



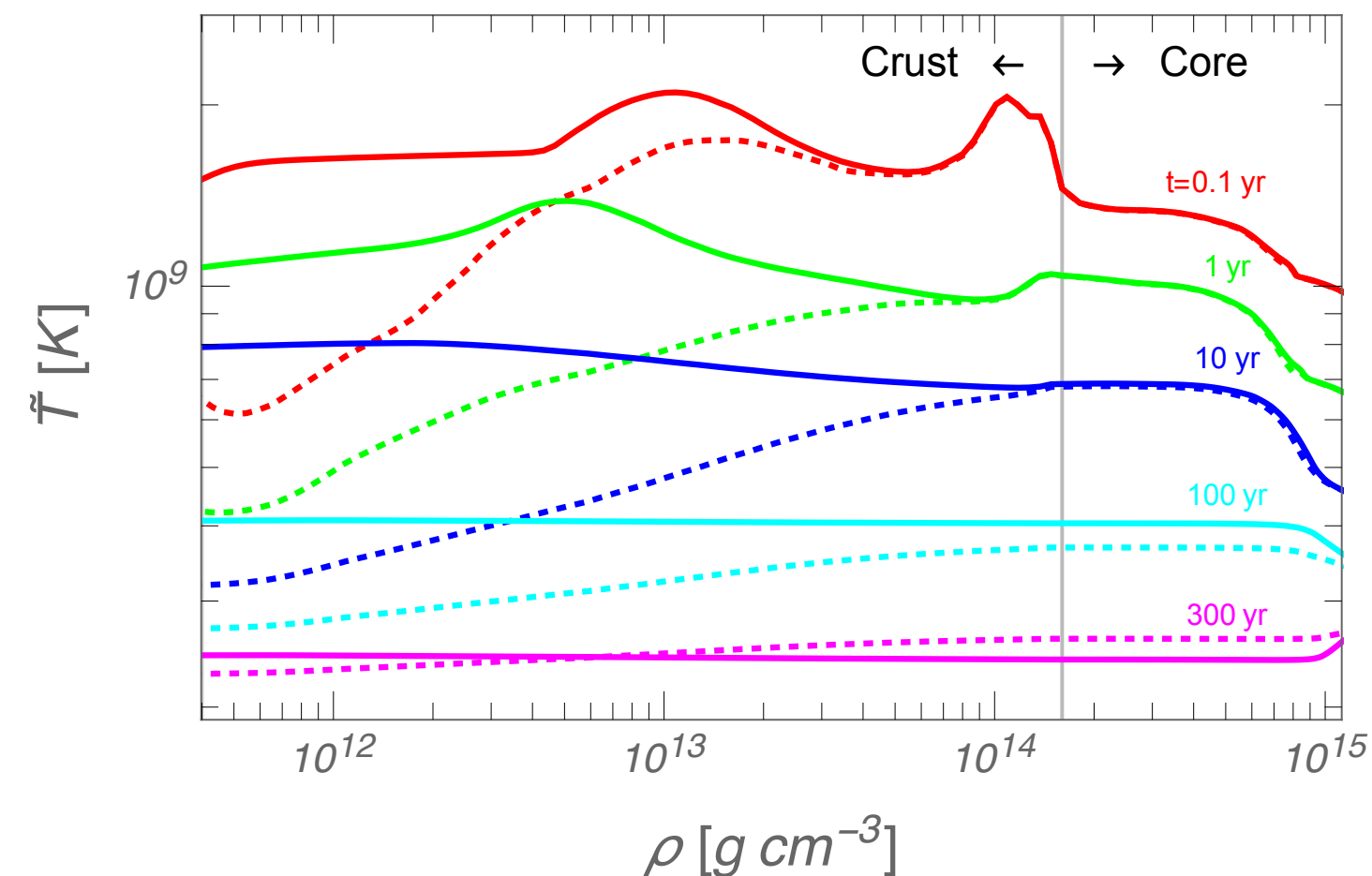
# Luminosity of dark photon (vs. $\nu$ )

\* Stellar cooling argument



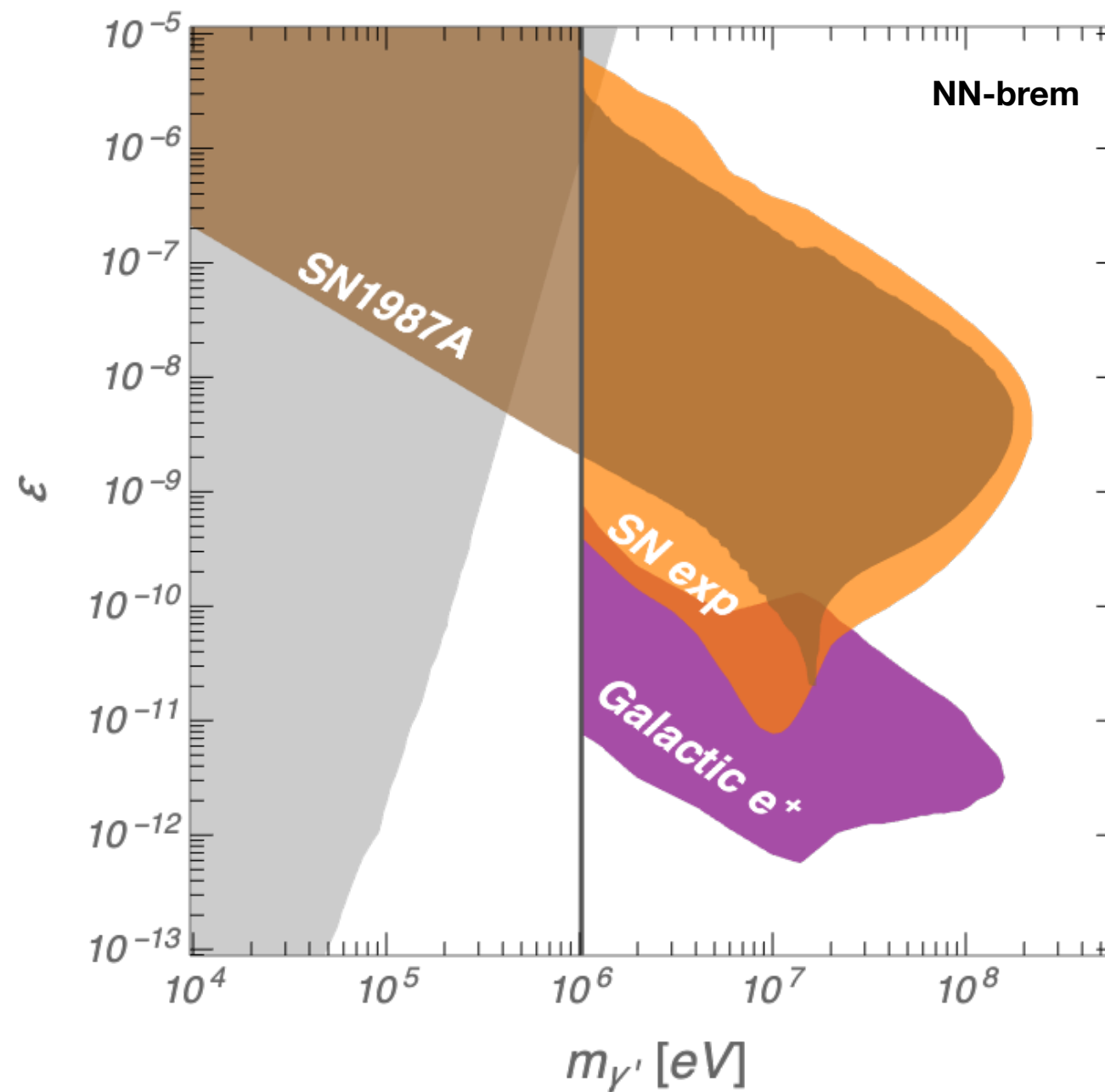


# Temperature profiles



- Before thermal relaxation ( $\sim 20$  years), the surface temperature is mainly determined by the crust
- Due to small neutrino emission at the crust, the dark photon cooling could be efficient before thermal relaxation
- $m_\gamma^2 \propto n_e \propto \rho$ , thus plasma suppression becomes weaker at outer part

# Dark photon constraints



# Dark photon constraints

