

PARISI-SOURLAS SUPERSYMMETRY

in RFIM and branched polymers

Apratim Kaviraj
DESY

Based on
1912.01617 , 2009.10087, 2112.06942, 2203.12629 w/ S. Rychkov & E. Trevisani



Invitation

[1912.01617] AK, Rychkov, Trevisani, JHEP

[2009.10087] AK, Rychkov, Trevisani, JHEP

[2112.06942] AK, Rychkov, Trevisani, PRL

[2203.12629] AK, Trevisani, JHEP

A class of disordered (impure) QFTs near criticality has a mysterious hidden supersymmetry (Parisi-Sourlas SUSY).

The SUSY results in the phenomenon of “dimensional reduction” of the disordered CFT (Parisi-Sourlas conjecture).

Some findings from experiments/numerics in this model have not been explained till date.

Main results:

- Systematic set-up of RG flow, understanding of the SUSY and its emergence
- Mechanism of how SUSY “breaks” and consistent explanation of all experimental/numerical findings.

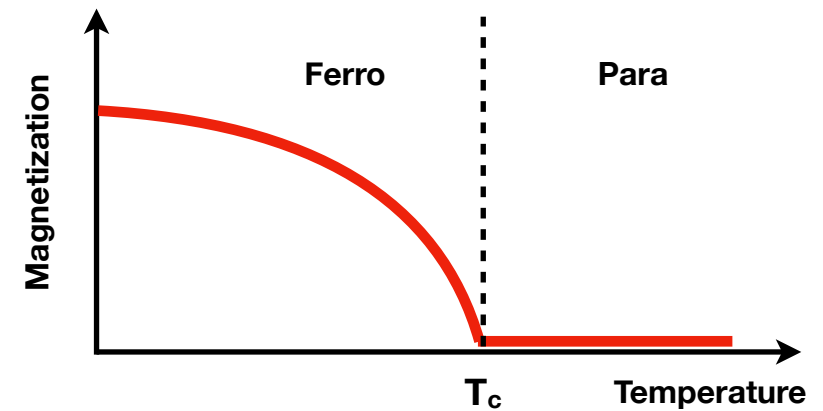
Disordered model	Dim reduced model	Dim. red. (Numerical finding)	Our explanation
RFIM	Ising	$6 \leq d < 4$	✓
Branched poly	Lee-Yang	$8 \leq d < 2$	✓

**Resolution of a
46 year old problem !**

RG flow and critical points

Conformal Field Theories (CFTs) are scale invariant theories, important in most aspects of theoretical physics (String Theory, AdS/CFT, Statistical Physics, etc).

CFTs describe critical behavior of systems, which is independent of microscopic details.



RG flow and critical points

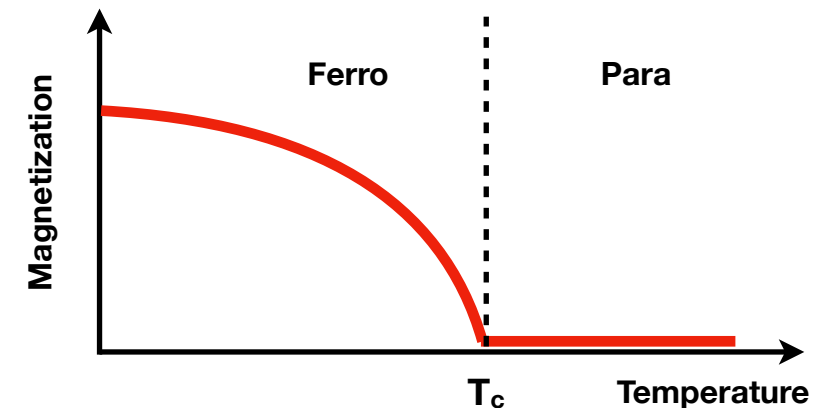
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A simple example: Ising model

$$H(J) = - \sum_{\langle ij \rangle} J s_i s_j$$

↑	↓	↑	↑	↓	↑
↑	↓	↑	↓	↑	↓
↑	↑	↓	↑	↓	↑



RG flow and critical points

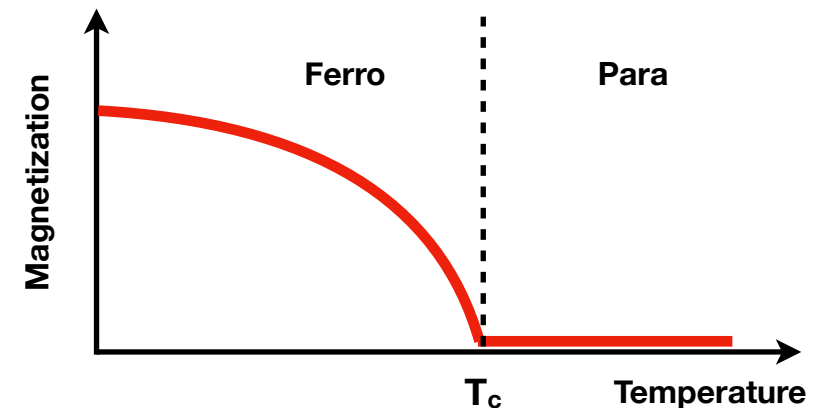
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RG flow of Ising model is often described by a scalar quantum field theory ($Z_{\phi^4} = Z_{\text{Ising}}$)

$$S[\phi] = \int d^d x \left[(\partial\phi)^2 + m^2 \phi^2 + \lambda \phi^4 \right]$$

RG flow and critical points

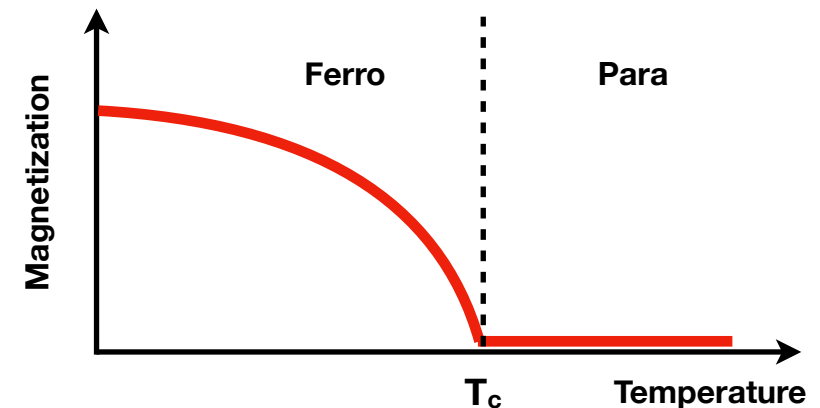
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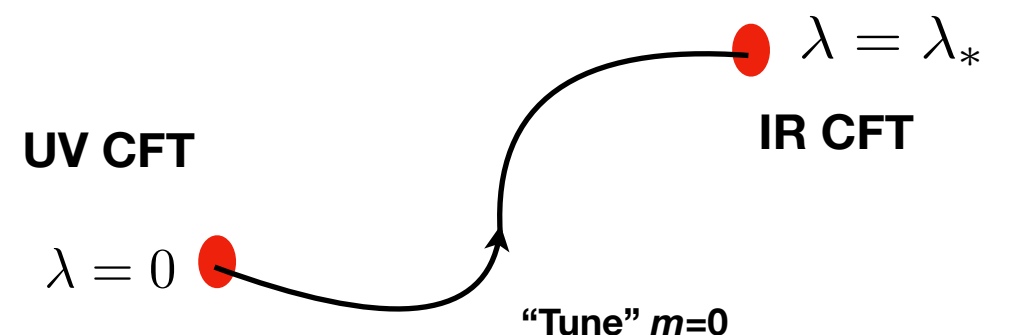
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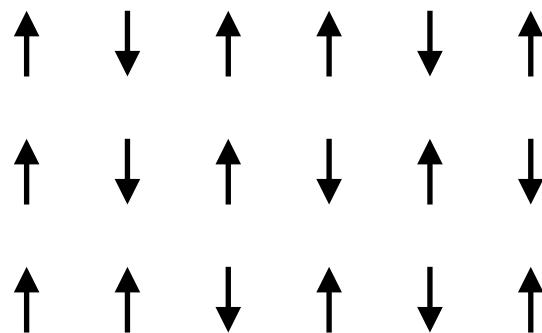
Beta function $\beta_\lambda = -\epsilon\lambda + \frac{3\lambda^2}{16\pi^2}$



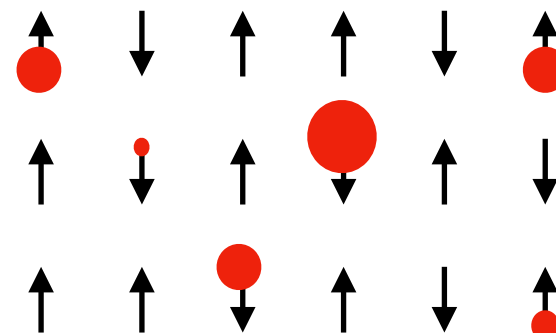
Pure vs impure CFTs

At large correlation lengths (near critical point) the effect of impurities become prominent.

Quenched disorder: When impurities are frozen - not in thermal equilibrium with system.



Pure



Impure

Random field models

Definition: $S[h] = \int d^d x \left[(\partial\phi)^2 + V(\phi) + h(x)\phi(x) \right]$

Impurity distribution:
(Gaussian) $P(h) = \exp \left[-\frac{1}{2H} \int d^d x h^2 \right]$

H = strength of disorder.

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Disorder averaging (quenched)

$$\overline{\langle \phi(x_1)\phi(x_2)\phi(x_3)\cdots \rangle} = \int \mathcal{D}h(x) \langle \phi(x_1)\phi(x_2)\cdots \rangle_h P[h(x)]$$

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Mostly studied:

$$V(\phi) \sim \phi^4$$

Random field Ising Model (RFIM)

$$V(\phi) \sim \phi^3$$

Random field ϕ^3 Model

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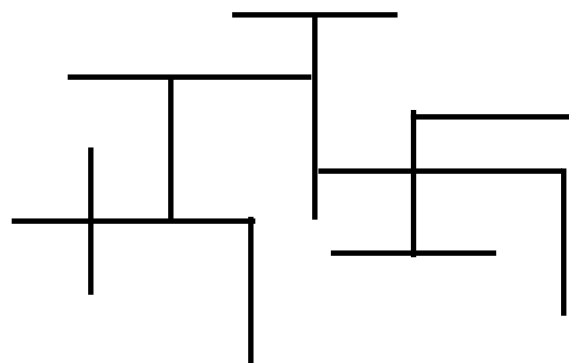
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Random field ϕ^3 Model



Lattice model



Branched polymers

Random field models

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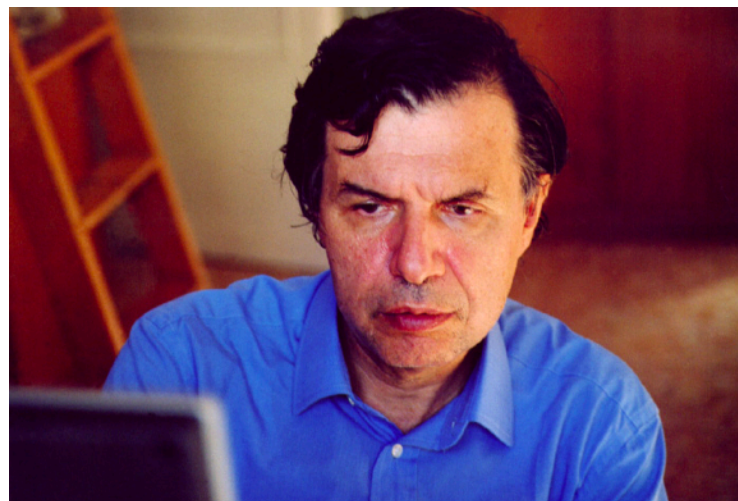
Disorder-averaged correlators have a new critical point !

**What CFT describes
the critical point of random field models?**

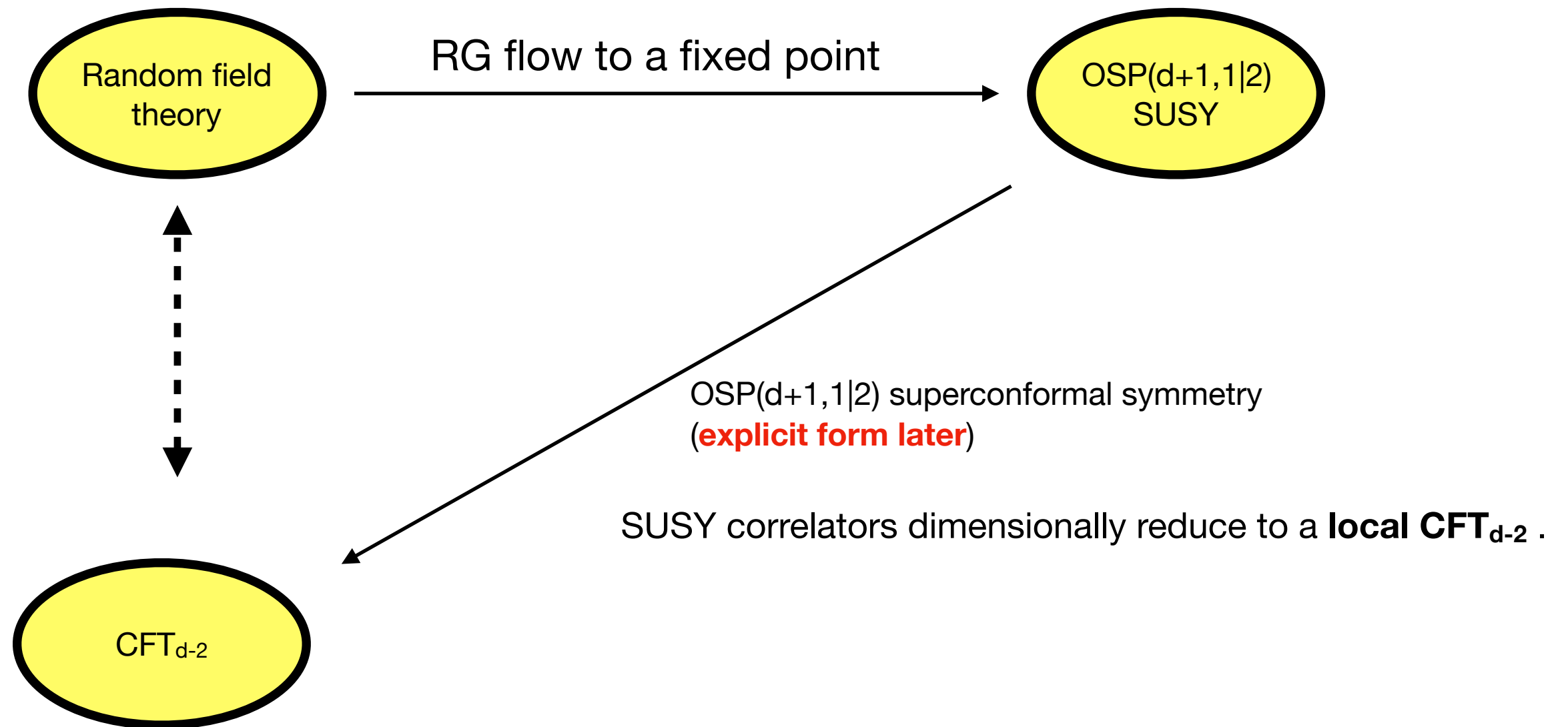
Parisi-Sourlas conjecture

[Parisi, Sourlas, 1979]

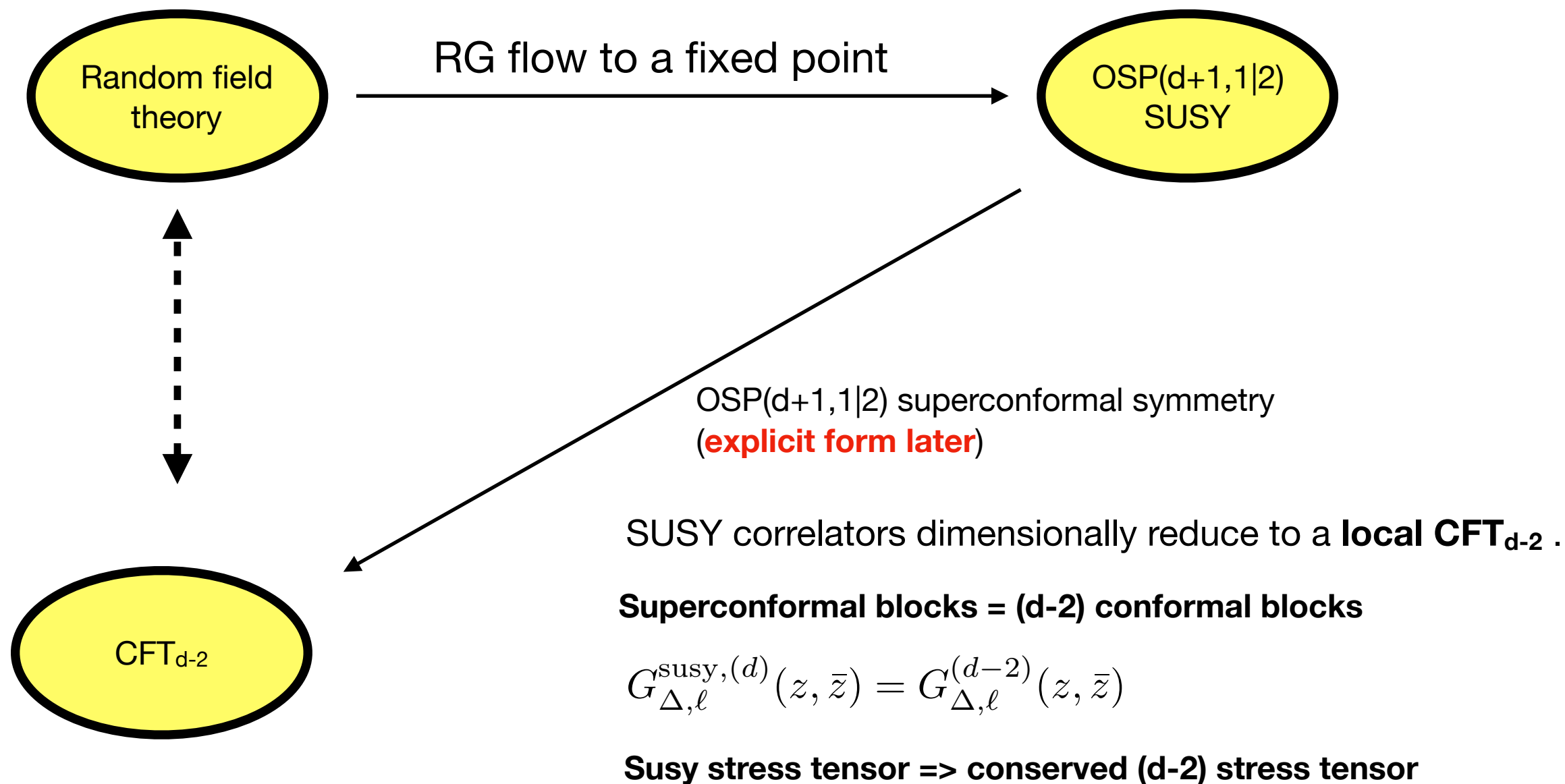
1. **Parisi-Sourlas SUSY:** The fixed point of a random field theory is described by a supersymmetric CFT.
2. **Dimensional reduction:** Observables in the SUSY CFT are same as a $(d-2)$ -dimensional CFT.



Parisi-Sourlas conjecture



Parisi-Sourlas conjecture



It does not always work !!

Numerical check of dimensional reduction (Monte Carlo)

[Fytas, Martin-Mayor, Parisi, Picco, Sourlas 2013-2019]

$$V(\phi) = \lambda \phi^4$$

Lattice model (Disordered CFT) d dimension	Lattice model (Pure CFT) d-2 dimension	Dimensional reduction
6d free theory	4d free theory	✓
5d RFIM	3d Ising	✓
4d RFIM	2d Ising	✗
3d RFIM	No Ising CFT	✗

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$V(\phi) = g \phi^3$

8d free theory	6d free theory	✓
7d Branched polymer	5d Lee-Yang	✓
6d Branched polymer	4d Lee-Yang	✓
5d Branched polymer	3d Lee-Yang	✓
4d Branched polymer	2d Lee-Yang	✓
3d Branched polymer	1d Lee-Yang	✓

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$V(\phi) = \lambda\phi^4$	6d free theory	4d free theory	✓
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	4d RFIM	2d Ising	✗

How does the Parisi-Sourlas conjecture work/fail?

$V(\phi) = g\phi^3$	6d Branched polymer	4d Lee-Yang	✓
	5d Branched polymer	3d Lee-Yang	✓
	4d Branched polymer	2d Lee-Yang	✓
	3d Branched polymer	1d Lee-Yang	✓

Main question: how does SUSY emerge in the RG flow?

To answer: We need a well-defined QFT

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We are interested in disorder-averaged correlation functions

$$\overline{\langle \phi(x_1) \phi(x_2) \cdots \rangle}$$

come from

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We are interested in disorder-averaged correlation functions

$$\overline{\langle \phi(x_1) \phi(x_2) \cdots \rangle} \quad \text{come from} \quad \overline{\log Z(h)}$$

$$\overline{\log Z(h)} = \lim_{n \rightarrow 0} \frac{\overline{Z^n} - 1}{n}$$

Enter the **Replica Method**

Replica Method

$$\overline{Z^n} = \int \prod_{i=1}^n \mathcal{D}\phi_i e^{-S_r}$$

$$S_r = \int d^d x \left[\frac{1}{2} \sum_{i=1}^n (\partial \phi_i)^2 - H \left(\sum_{i=1}^n \phi_i \right)^2 + \sum_{i=1}^n V(\phi_i) \right]$$

This action at $n \rightarrow 0$

is a QFT that describes random field models

Replica action in Cardy variables

$$\sum_i (\partial\phi_i)^2 + H(\sum_i \phi_i)^2 + \sum_i (m^2\phi_i^2 + \lambda\phi_i^4)$$

$$\phi_1 = \varphi + \frac{\omega}{2} \quad \phi_i = \varphi - \frac{\omega}{2} + \chi_i \quad [i = 2, \dots, n]$$

(Gives good scaling operators)

[Cardy, 1985]

$$= \partial\varphi\partial\omega - H\omega^2 + \frac{1}{2} \sum_{i=2}^n (\partial\chi_i)^2 + m^2(\varphi\omega + \frac{\chi_i^2}{2}) + \lambda(4\varphi^3\omega + 6\varphi^2\chi_i^2) \\ + O(n) + \text{irrelevant}$$

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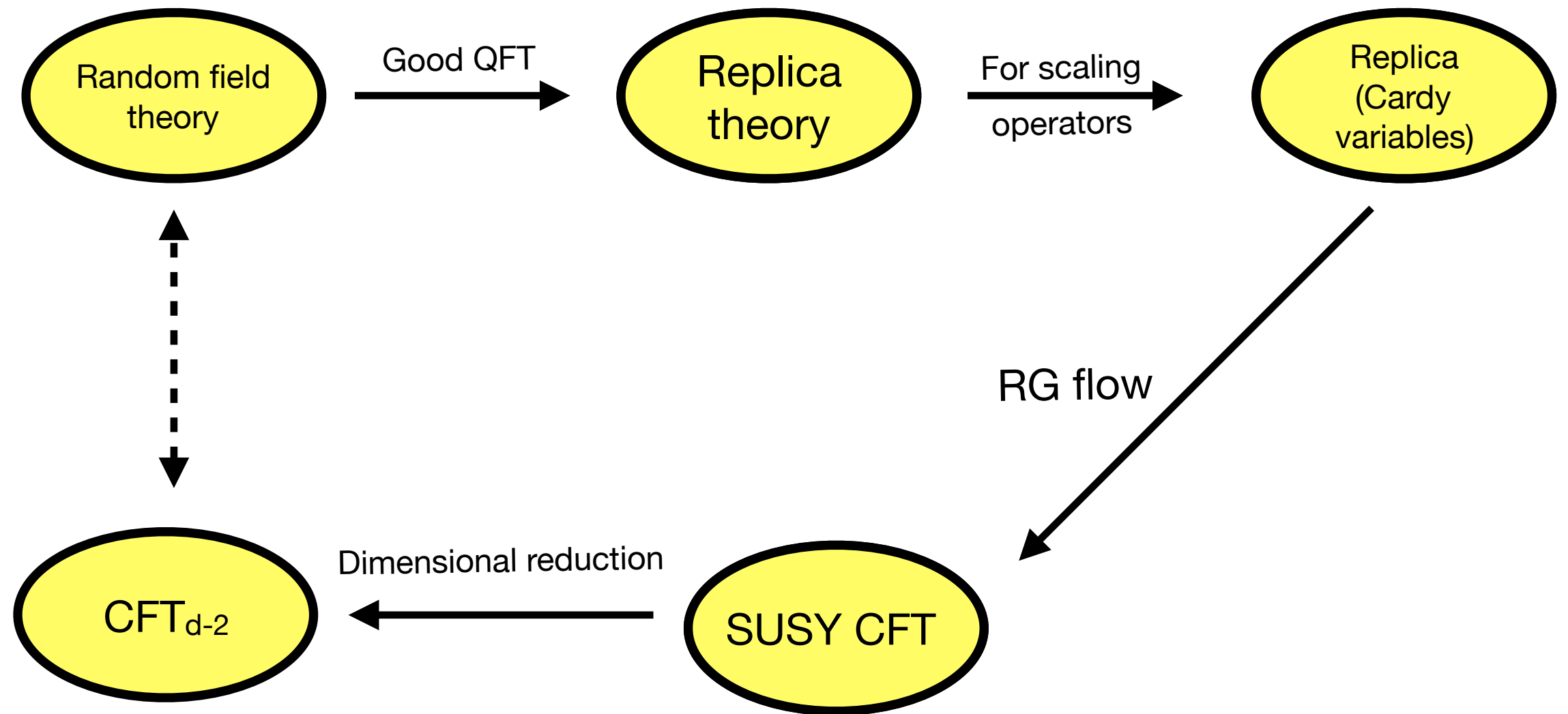
[Cardy, 1985]

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$$\chi_i \xrightarrow{n \rightarrow 0} \psi, \bar{\psi} \text{ (fermions)}$$

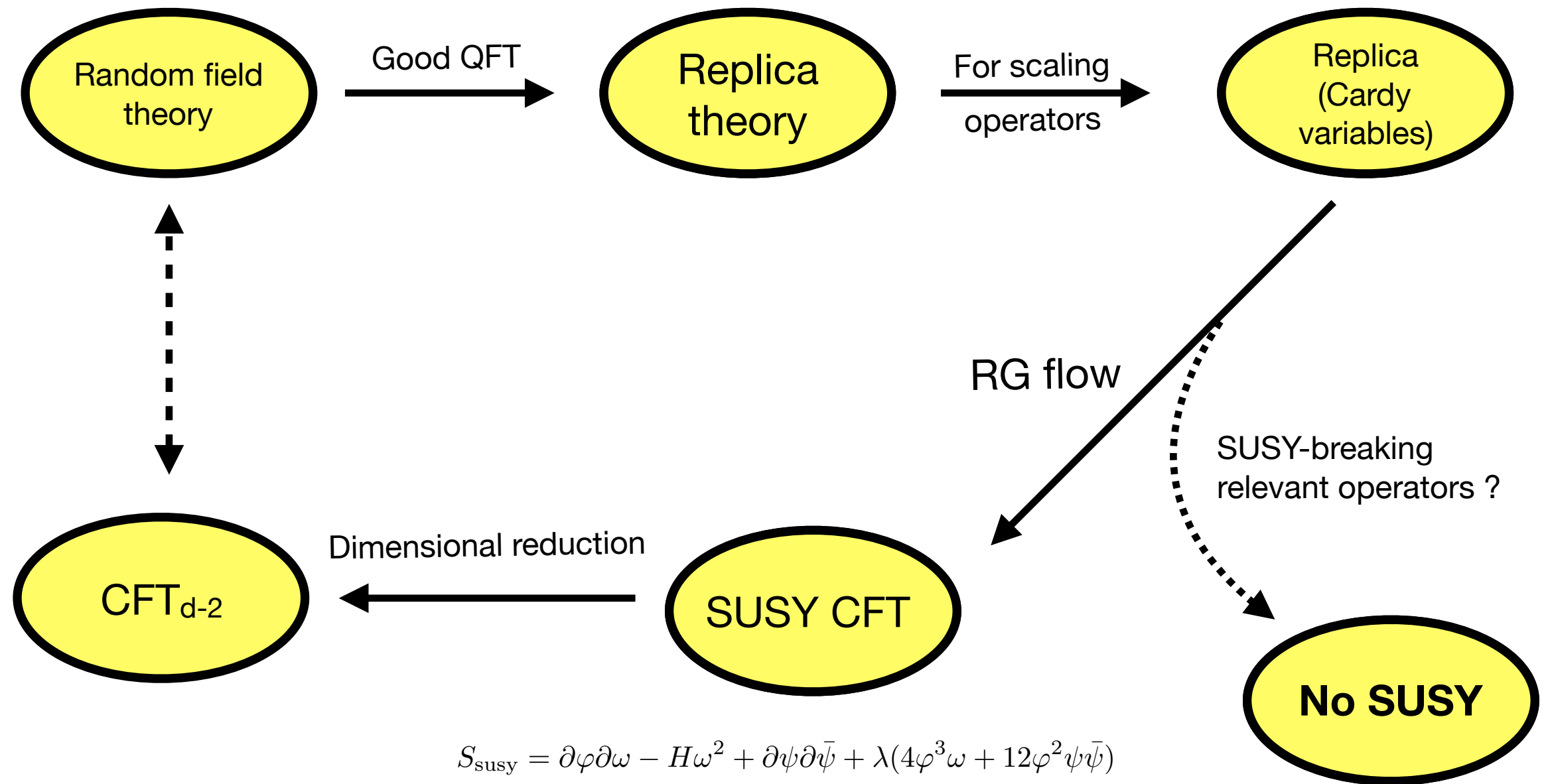
→ **OSp supersymmetric theory**

Emergence of SUSY

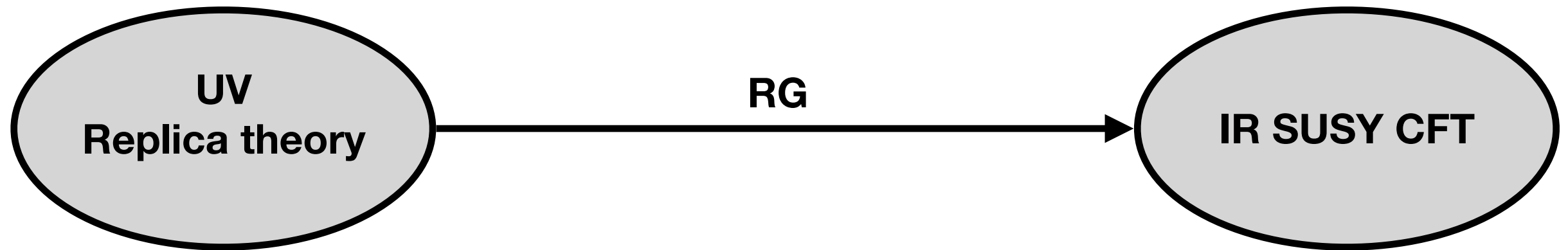


$$S_{\text{susy}} = \partial\varphi\partial\omega - H\omega^2 + \partial\psi\partial\bar{\psi} + \lambda(4\varphi^3\omega + 12\varphi^2\psi\bar{\psi})$$

Emergence of SUSY



A paradox!!



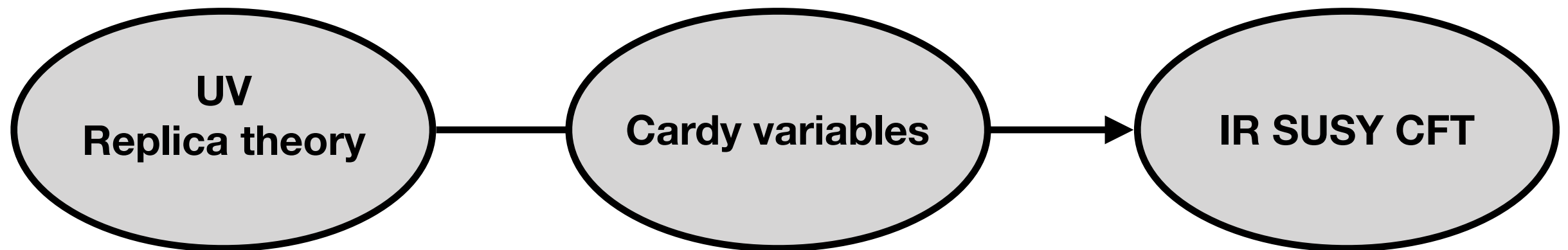
Permutation symmetry
of n replicas
(S_n symmetry)



No S_n symmetry

But RG should preserve all symmetries !!

Resolution



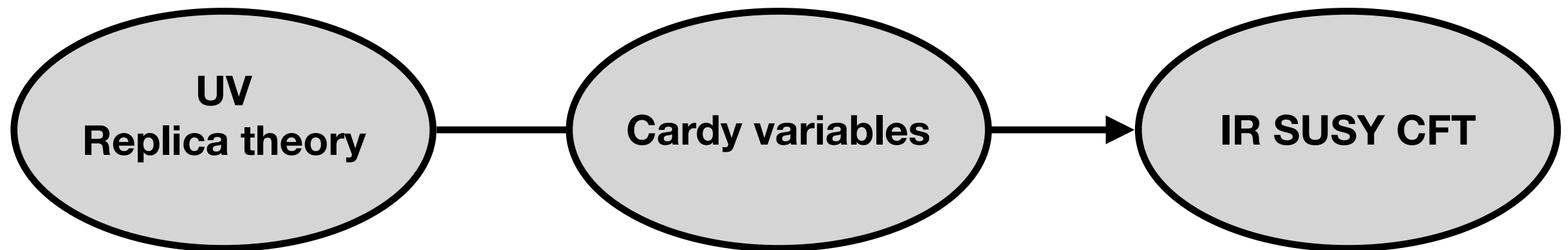
**Obscures S_n symmetry
to get scaling operators.**

$$[\varphi] = \frac{d}{2} - 2, [\chi_i] = \frac{d}{2} - 1, [\omega] = \frac{d}{2}$$

$$\sum_{i=1}^n \phi_i^4 : \stackrel{\text{Cardy}}{=} 4\varphi^3\omega + 6\varphi^2\chi_i^2 + \frac{1}{H}(\chi_i^4 - 6\varphi\chi_i^2\omega) + \frac{1}{H^2}\varphi\omega^3 + \dots$$

Relevant
Irrelevant
More irrelevant

Resolution

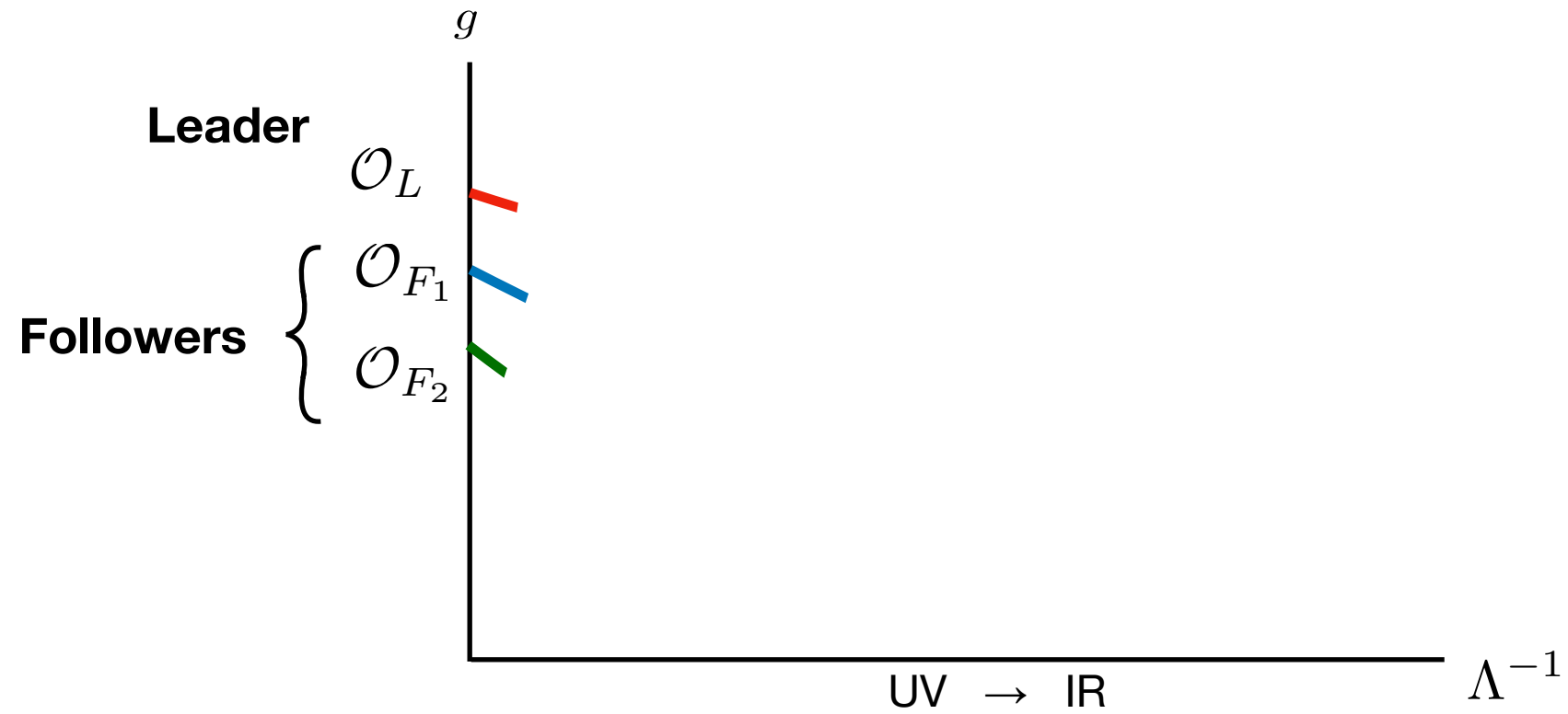


Obscures S_n symmetry
to get scaling operators.

$$\mathcal{O}_{S_n \text{ singlet}} \stackrel{\text{Cardy}}{=} \mathcal{O}_L + \mathcal{O}_{F_1} + \mathcal{O}_{F_2} + \dots$$

same RG correction (anomalous dimension)
but different overall scaling

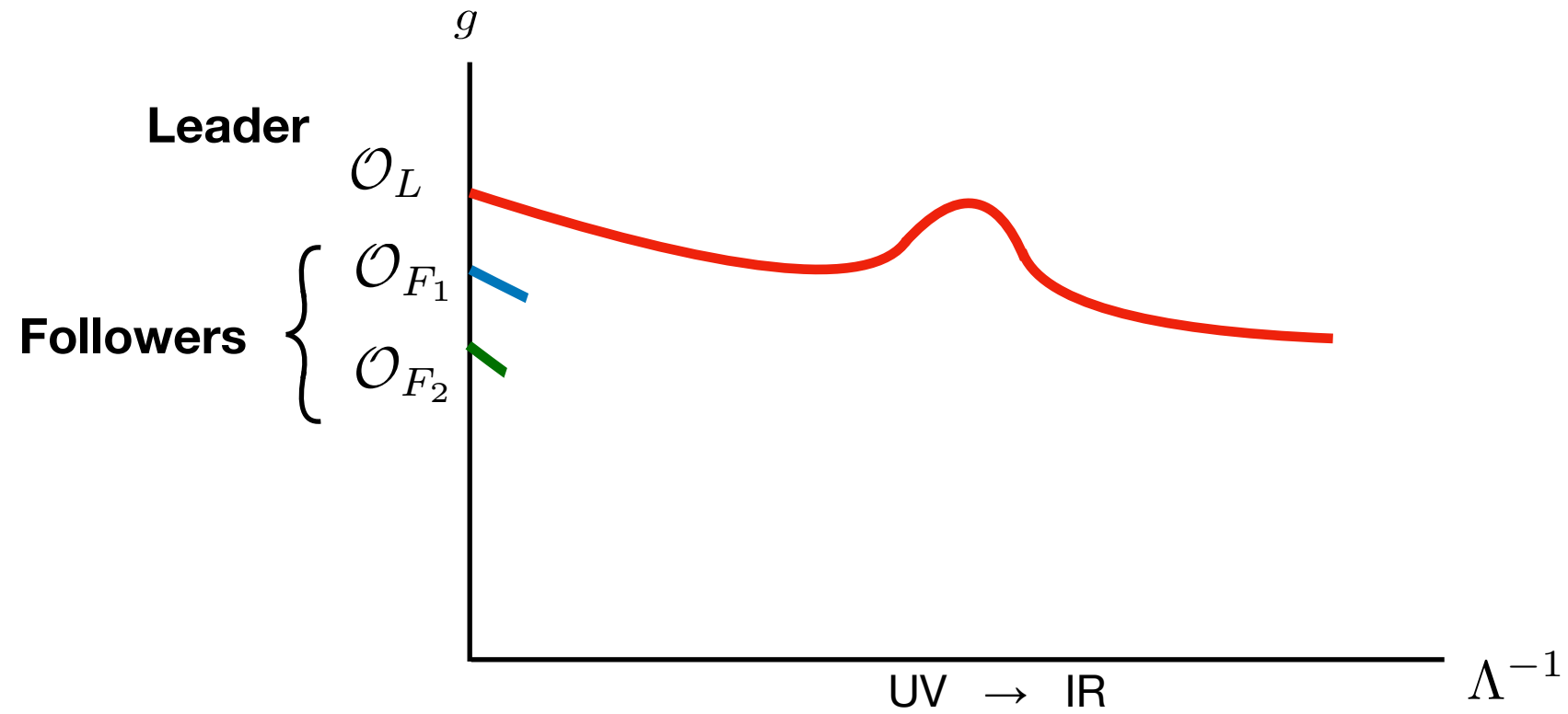
RG evolution of leader and followers



[2009.10087] AK, Rychkov, Trevisani

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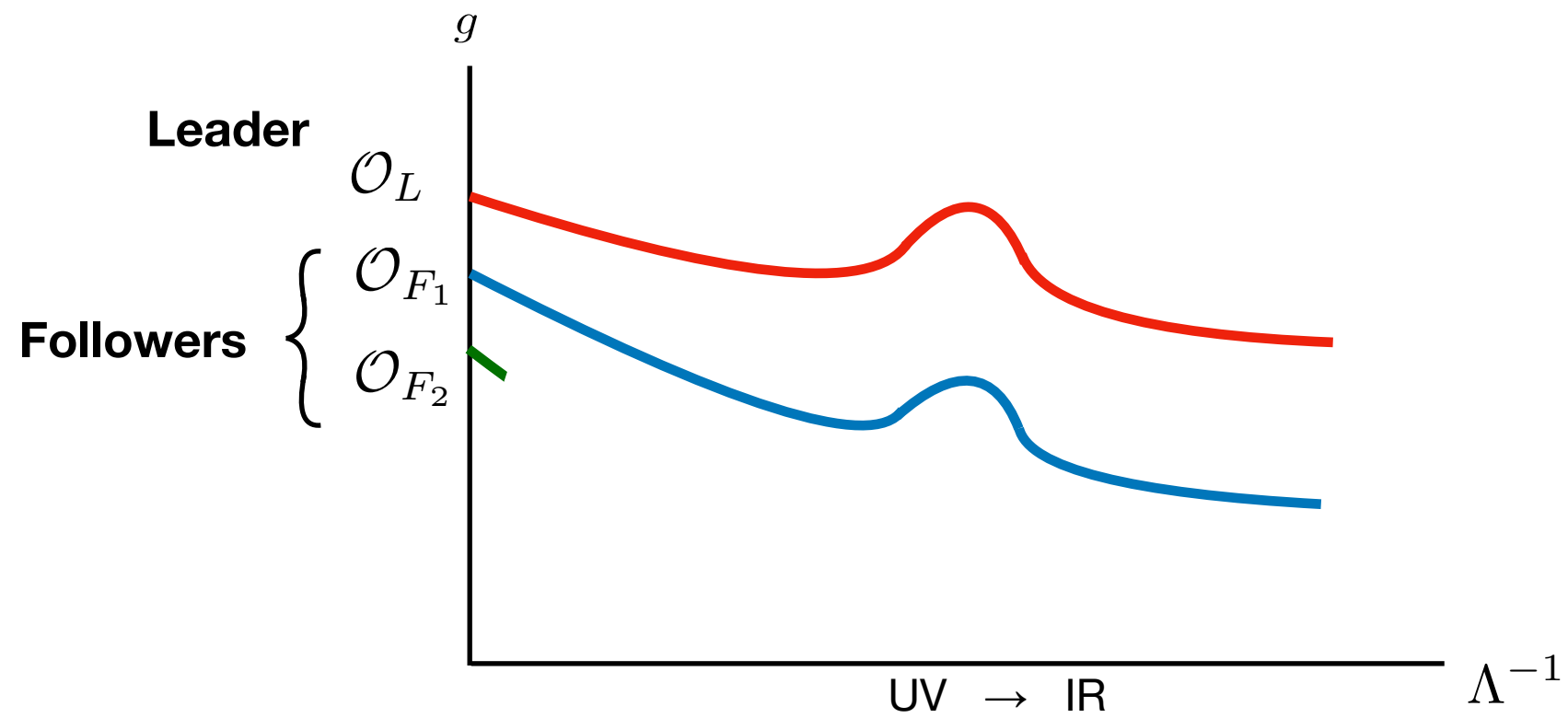
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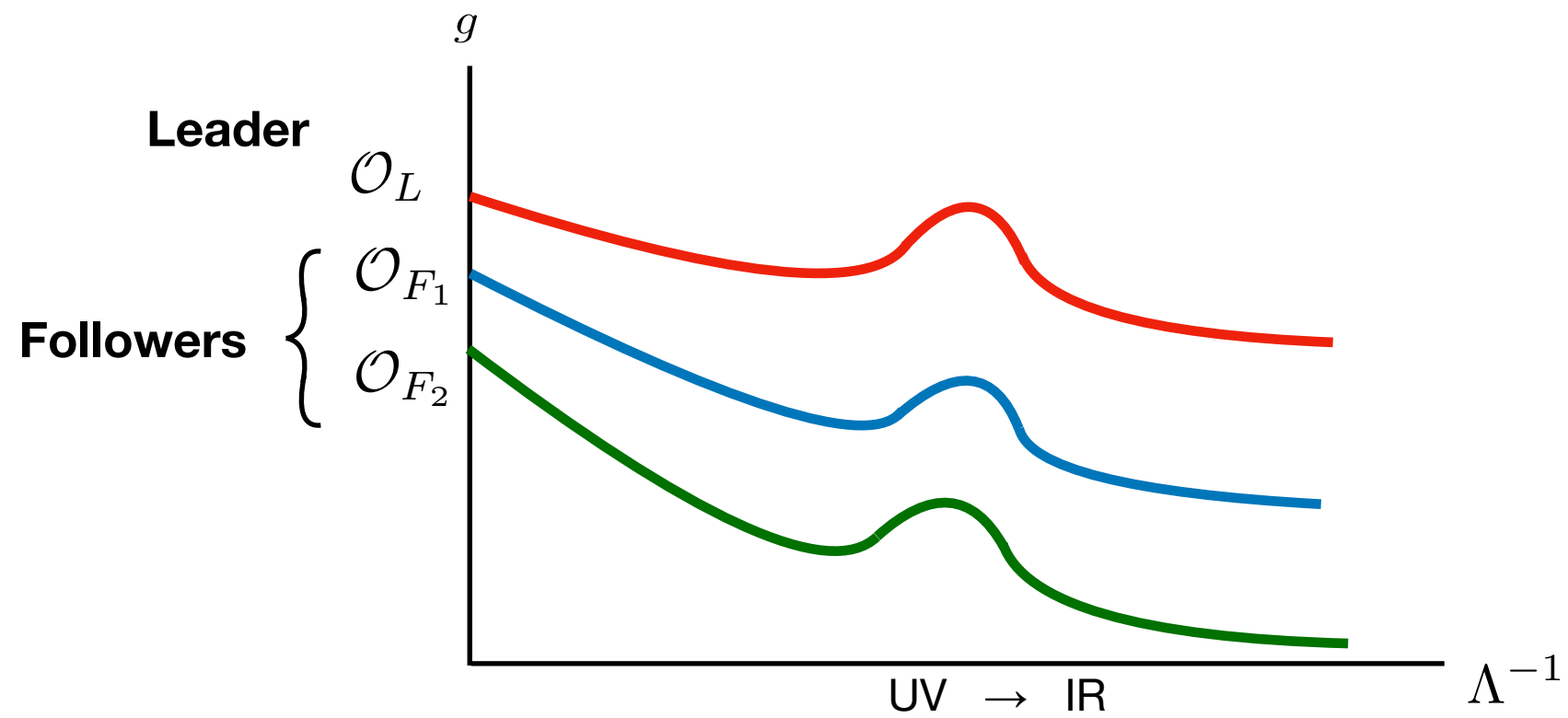
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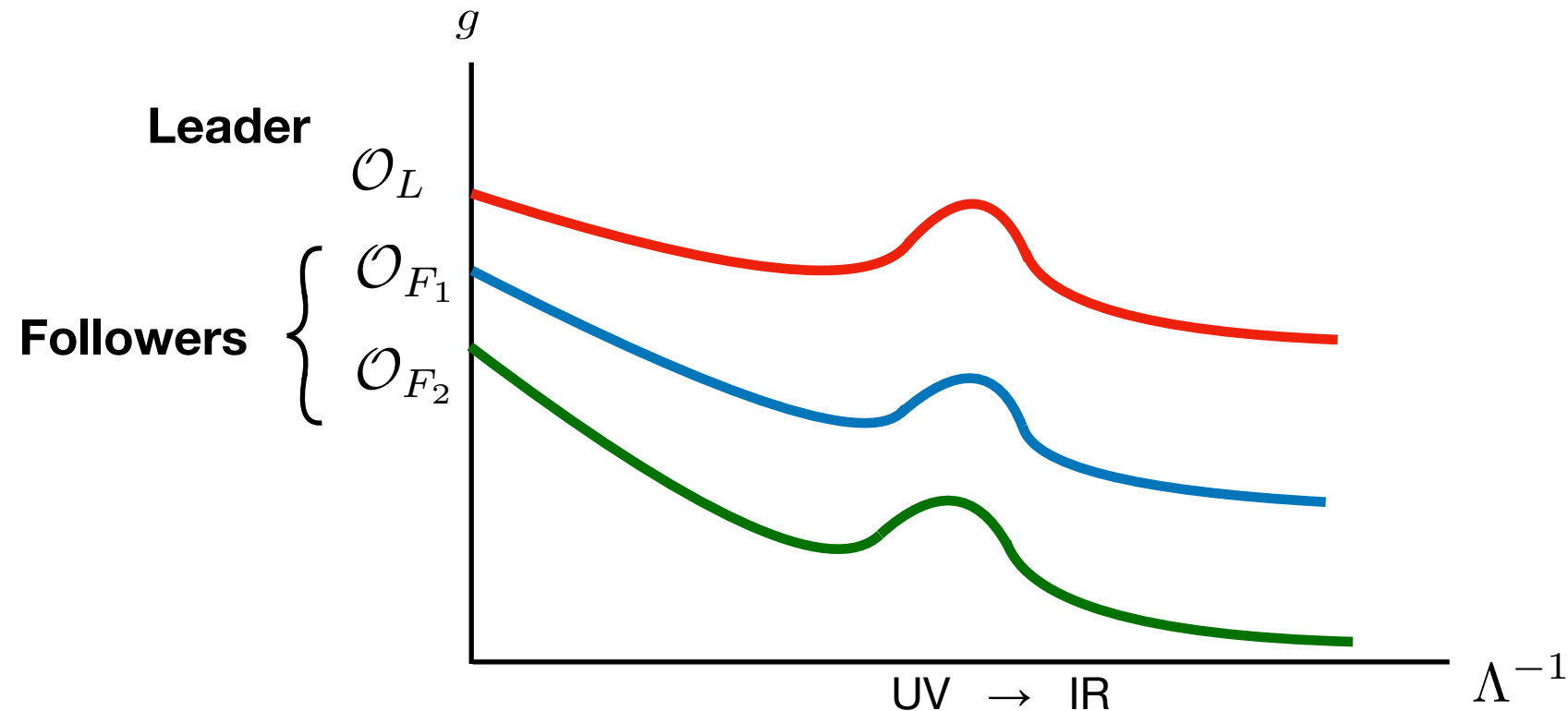
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RG evolution of leader and followers



**The leaders may flow to (an emergent) SUSY CFT
....or may not !!**

**We classified and computed anomalous dimensions of
low lying leaders in perturbation theory**

[2009.10087] AK, Rychkov, Trevisani

[2112.06942] AK, Rychkov, Trevisani

Findings

For $V(\phi) = \lambda\phi^4$ we found SUSY breaking leader operators become relevant in $d \leq d_c \in [4.2, 4.5]$ (unstable fixed point)

Numerical findings of SUSY or dimensional reduction

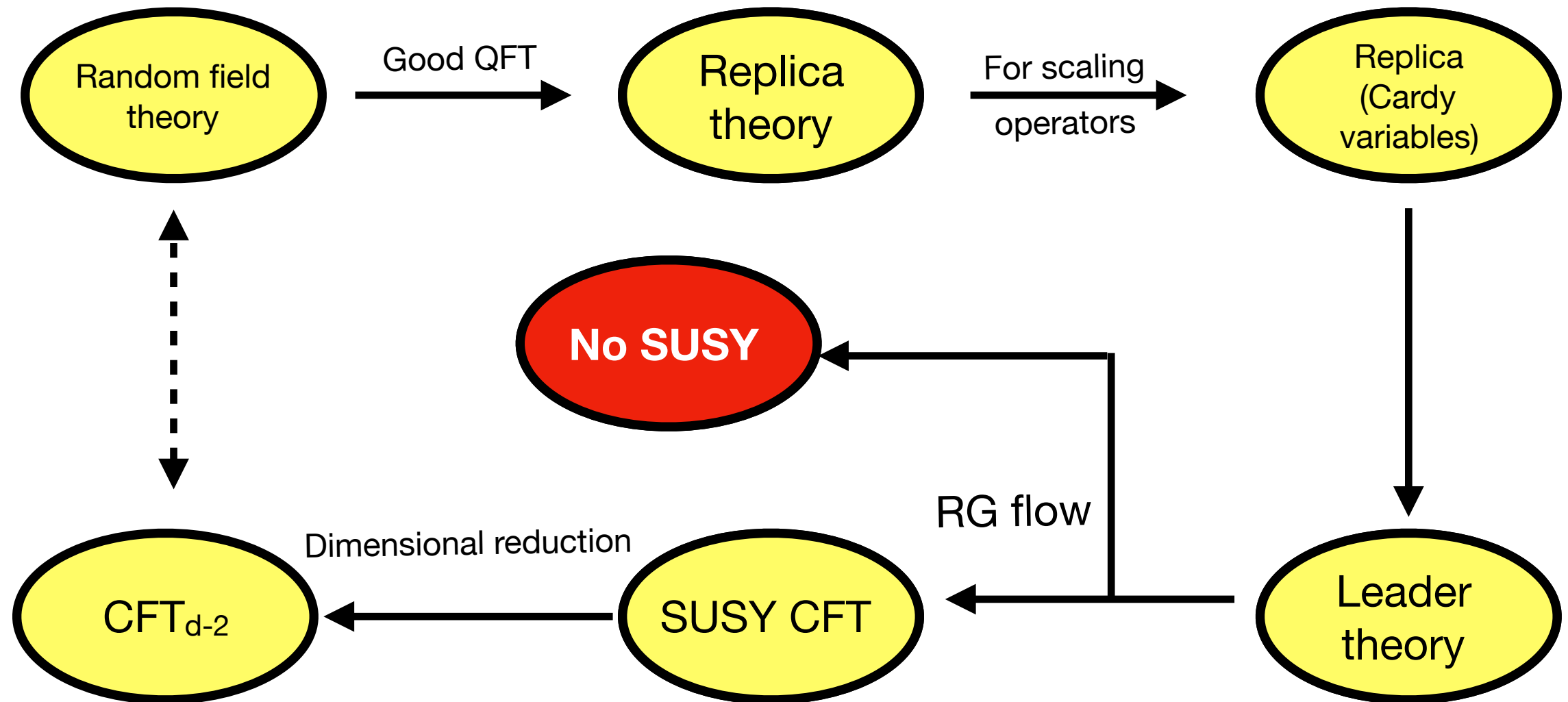
$d=6$	✓
$d=5$	✓
$d=4$	✗
$d=3$	✗

$d=8$	✓
$d=7$	✓
$d=6$	✓
$d=5$	✓
$d=4$	✓
$d=3$	✓

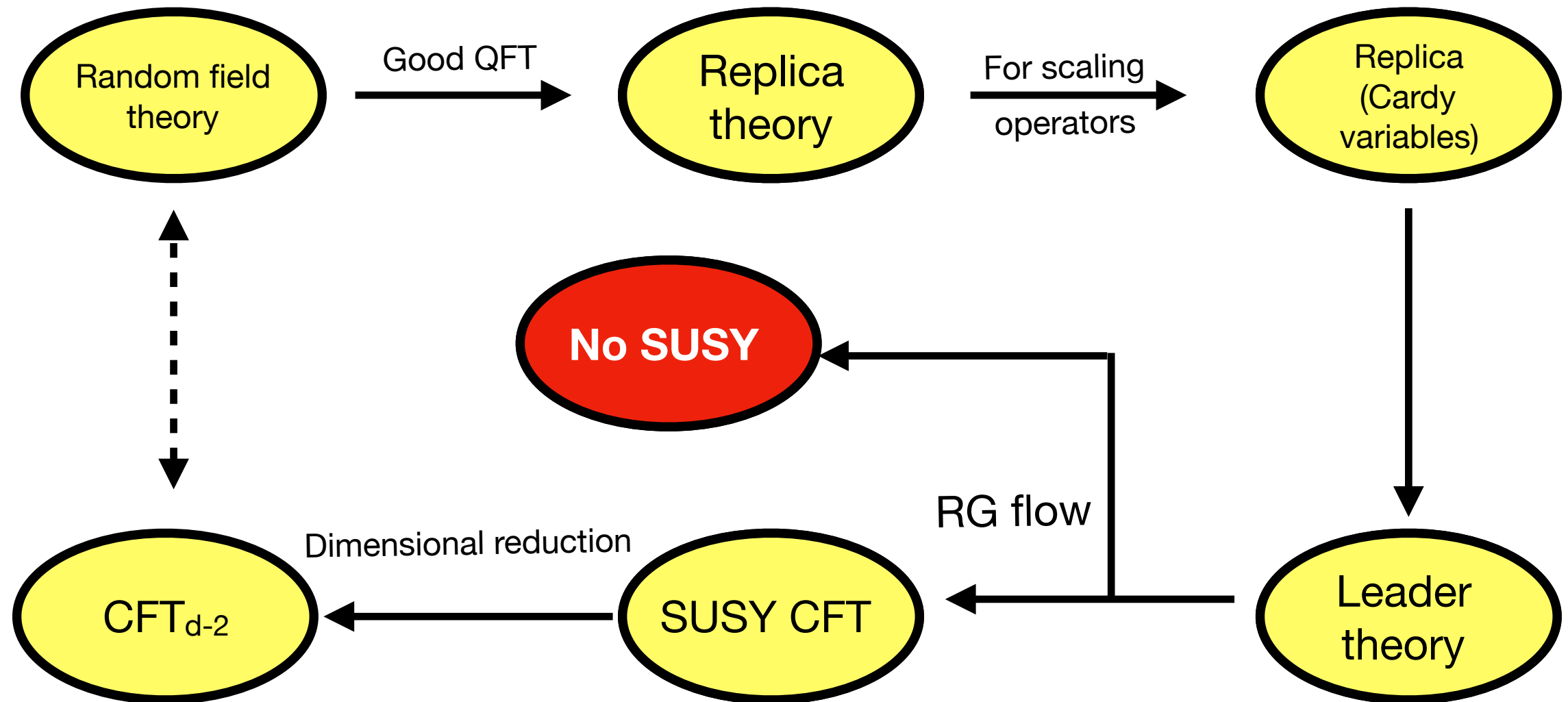
For $V(\phi) = g\phi^3$ we found no SUSY breaking leader operators becoming relevant.

[2009.10087] AK, Rychkov, Trevisani
[2112.06942] AK, Rychkov, Trevisani
[2203.12629] AK, Trevisani
(JHEP, PRL)

Summary



Summary



Future directions

- Tuning to an unstable SUSY fixed point in $d=4$ RFIM (confirmation of our scenario).
- Conformal bootstrapping the RFIM (Difficulty: taking limits)
- Unexplored features of logarithmic CFTs.
- Understanding the implications SUSY and dimensional reduction.

Thanks