

Charmless multibody B decays at LHCb

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Conditions for CP violation in decay

Several other types of CP violation: in mixing, mixing-induced

In flavour-specific B decays, presence of multiple amplitudes may lead to (direct) CP violation in decay

$$A(B \rightarrow f) = \sum_i |A_i| e^{i(\delta_i + \phi_i)}$$

$$\bar{A}(\bar{B} \rightarrow \bar{f}) = \sum_i |A_i| e^{i(\delta_i - \phi_i)}$$

Strong phase (δ) invariant under CP

Weak phase (ϕ) changes sign under CP

$$\mathcal{A}_{CP}(B \rightarrow f) \equiv \frac{|\bar{A}|^2 - |A|^2}{|\bar{A}|^2 + |A|^2} \propto \sum_{i,j} |A_i| |A_j| \sin(\delta_i - \delta_j) \sin(\phi_i - \phi_j)$$

3 conditions required for CP violation in decay

At least 2 competitive amplitudes

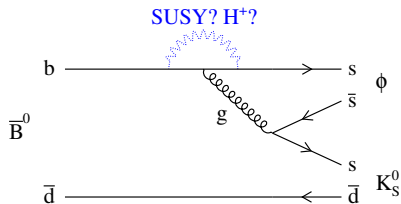
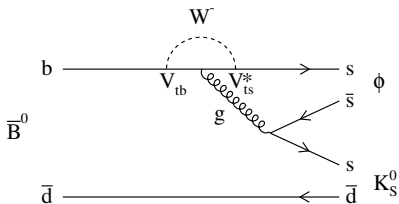
Non-zero strong phase difference, $\delta_i - \delta_j \neq 0$

Non-zero weak phase difference, $\phi_i - \phi_j \neq 0$

Weak phase comes from different CKM phases in each amplitude

Strong phase structure from decay amplitude in multibody phase space

2nd-order loop diagrams potentially sensitive to New Physics



Unknown heavy particle could be present in the loop

May carry a new CP violating phase

Short time-scale in loop the key to accessing arbitrarily high energies

Heisenberg Uncertainty Principle: $\Delta E \Delta t \geq \hbar/2$

Can reach higher mass scales than possible with direct searches at LHC

Mass scale of New Physics through loop processes a matter of statistics rather than CMS energy

Deviation of measured CP violating phase and Standard Model expectation a clear signature of New Physics

1. Charmless 3-body B decays

- Impact of strong phase motion on CP violation structure
- Binned analysis of $B^+ \rightarrow h^+ h'^+ h'^-$

2. Charmless baryonic B decays

- Triple-product asymmetries
- $B^0 \rightarrow p\bar{p}K^+\pi^-$

Short-distance contributions

Direct CP violation more complicated in $B \rightarrow 3h$ decay channels compared to 2-body decays

There are at least 4 possible sources of strong phase

1. Short-distance contributions (quark level)

BSS mechanism, PRL **43** (1979) 242

Tree contribution (a)

Penguin diagram (b) contains 3 quark generations in loop

From S -matrix unitarity, CPT conservation requires absorptive amplitude

Quark rescattering (c) generates a phase

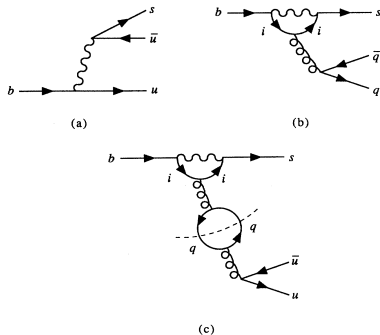
If gluon in penguin is on-shell

Momentum transfer $q^2 > 4m_i^2$
where $i = u, c$

Imaginary part depends on quark masses

CP violation in 2-body processes caused by this effect

eg. $B^0 \rightarrow K^+\pi^-$



PRD **43** (1991) 2909

Remaining sources more associated with multibody decays

Long-distance contributions (hadronic level)

2. Breit-Wigner phase

Propagator represents intermediate resonance states

$$T_R^{\text{BW}}(s) = \frac{1}{m_R^2 - s - im_R\Gamma_R(s)}$$

Phase varies across mass-squared, s

3. Relative CP -even phase in the isobar model

$$A(B \rightarrow f) = \sum_i |A_i| e^{i(\delta_i + \phi_i)}$$

$$\bar{A}(\bar{B} \rightarrow \bar{f}) = \sum_i |\bar{A}_i| e^{i(\delta_i - \phi_i)}$$

Related to final state interactions between different resonances

Manifestation of CP violation

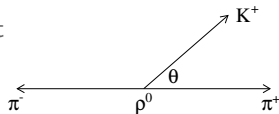
Each source of strong phase leaves a unique signature in phase space
 Illustrate with series of examples

Consider $B^\pm \rightarrow K^\pm \pi^+ \pi^-$ with only 2 isobars

$B^\pm \rightarrow \rho^0 K^\pm$ and flat non-resonant (NR) component

ρ^0 lineshape a Breit-Wigner, T_ρ^{BW}

ρ^0 is a vector resonance, so angular distribution follows $\cos \theta$



$$B^+ : A_+ = |a_+^\rho| e^{i\delta_+^\rho} T_\rho^{\text{BW}} \cos \theta + |a_+^{\text{NR}}| e^{i\delta_+^{\text{NR}}}$$

$$B^- : A_- = |a_-^\rho| e^{i\delta_-^\rho} T_\rho^{\text{BW}} \cos \theta + |a_-^{\text{NR}}| e^{i\delta_-^{\text{NR}}}$$

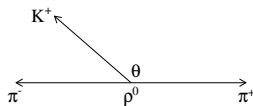
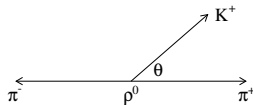
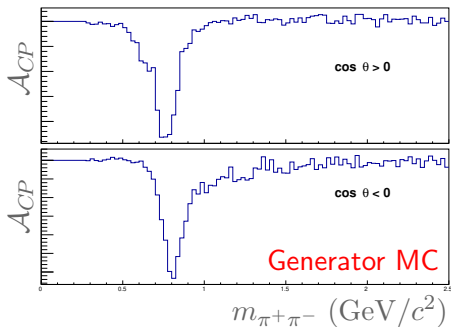
$$\begin{aligned} \mathcal{A}_{CP} &\propto |A_-|^2 - |A_+|^2 \\ &\propto (|a_-^\rho|^2 - |a_+^\rho|^2) |T_\rho^{\text{BW}}|^2 \cos^2 \theta \dots \\ &\quad - 2(m_\rho^2 - s) |T_\rho^{\text{BW}}|^2 \cos \theta \dots \\ &\quad + 2m_\rho \Gamma_\rho |T_\rho^{\text{BW}}|^2 \cos \theta \dots \end{aligned}$$

Short-distance effects

$$\begin{aligned}
 \mathcal{A}_{CP} \propto & (|a_-^\rho|^2 - |a_+^\rho|^2) |T_\rho^{\text{BW}}|^2 \cos^2 \theta \dots \\
 & -2(m_\rho^2 - s) |T_\rho^{\text{BW}}|^2 \cos \theta \dots \\
 & +2m_\rho \Gamma_\rho |T_\rho^{\text{BW}}|^2 \cos \theta \dots
 \end{aligned}$$

Only depends on ρ resonance

Maximum difference at ρ pole, quadratic in helicity



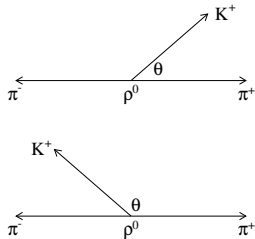
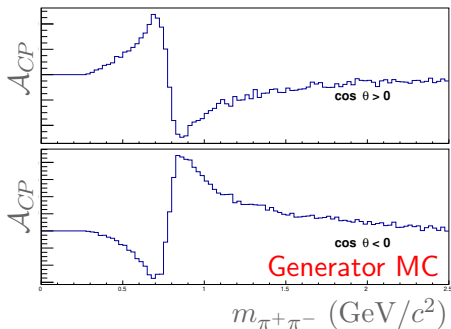
Only short-distance effects can create $|a_+^\rho| \neq |a_-^\rho|$

Long-distance effects

$$\begin{aligned}
 \mathcal{A}_{CP} &\propto (|a_-^\rho|^2 - |a_+^\rho|^2) |T_\rho^{BW}|^2 \cos^2 \theta \dots \\
 &\quad - 2(m_\rho^2 - s) |T_\rho^{BW}|^2 \cos \theta \dots \\
 &\quad + 2m_\rho \Gamma_\rho |T_\rho^{BW}|^2 \cos \theta \dots
 \end{aligned}$$

Interference term from real part of Breit-Wigner

Zero at ρ pole, linear in helicity



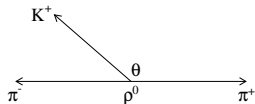
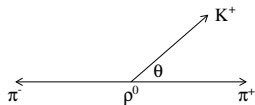
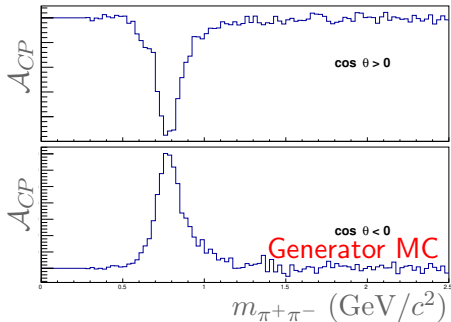
Caused by long-distance effects from final state interactions

Long-distance effects

$$\begin{aligned}
 \mathcal{A}_{CP} &\propto (|a_-^\rho|^2 - |a_+^\rho|^2) |T_\rho^{\text{BW}}|^2 \cos^2 \theta \dots \\
 &\quad - 2(m_\rho^2 - s) |T_\rho^{\text{BW}}|^2 \cos \theta \dots \\
 &\quad + 2m_\rho \Gamma_\rho |T_\rho^{\text{BW}}|^2 \cos \theta \dots
 \end{aligned}$$

Interference term from imaginary part of Breit-Wigner

Maximum at ρ pole, linear in helicity



Caused by long distance effects from Breit-Wigner phase and final state interactions

Last source of strong phase

4. Final state $KK \leftrightarrow \pi\pi$ rescattering

Can occur between decay channels with the same flavour quantum numbers

eg. $B^\pm \rightarrow K^\pm K^+ K^-$ and $B^\pm \rightarrow K^\pm \pi^+ \pi^-$

CPT conservation constrains hadron rescattering

For given quantum numbers, sum of partial widths equal for charge-conjugate decays

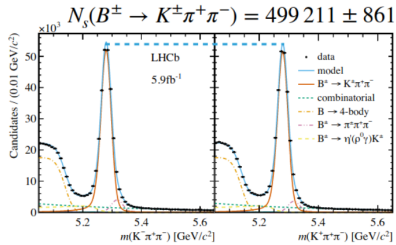
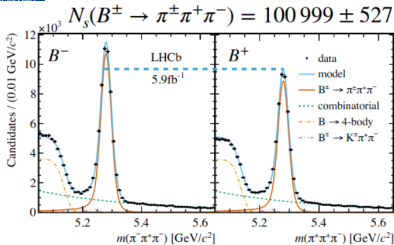
$KK \leftrightarrow \pi\pi$ rescattering generates a strong phase

Look into rescattering region

If rescattering phase in one decay channel generates direct CP violation in this region

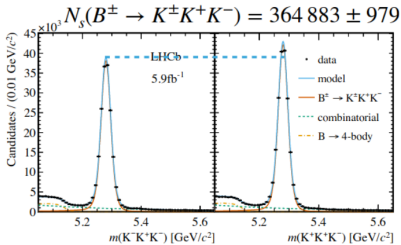
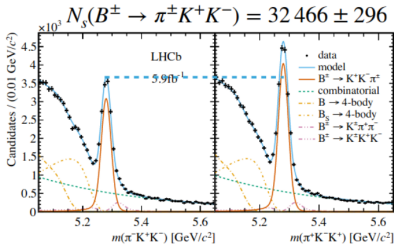
Rescattering phase should generate opposite sign direct CP violation in partner decay channel

$B^+ \rightarrow h^+ h'^+ h'^-$ at LHCb



$$A_{CP} = +0.080 \pm 0.004 \pm 0.003 \pm 0.003 \quad (14.1\sigma)$$

$$A_{CP} = +0.011 \pm 0.002 \pm 0.003 \pm 0.003 \quad (2.4\sigma)$$

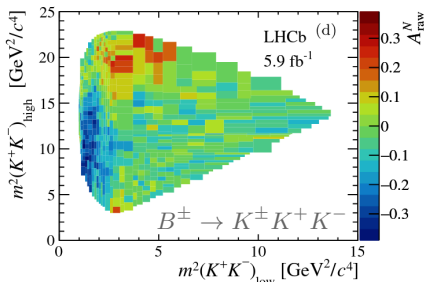
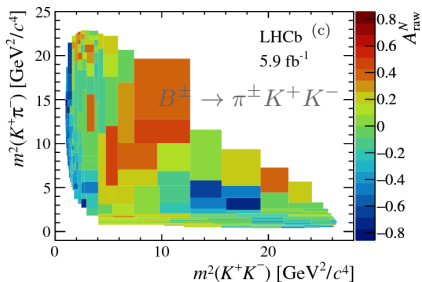
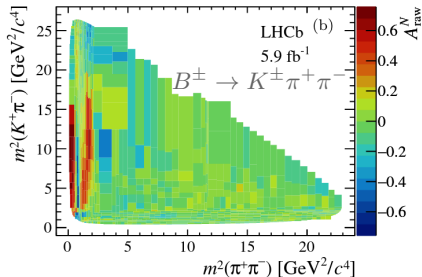
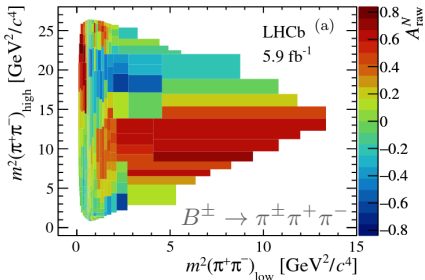


$$A_{CP} = -0.114 \pm 0.007 \pm 0.003 \pm 0.003 \quad (13.6\sigma)$$

$$A_{CP} = -0.037 \pm 0.002 \pm 0.002 \pm 0.003 \quad (8.5\sigma)$$

Phys. Rev. **D108** (2023) 012013, Phys. Rev. **D108** (2023) 012008

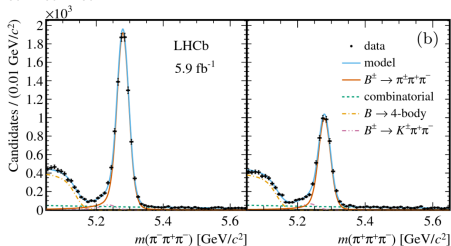
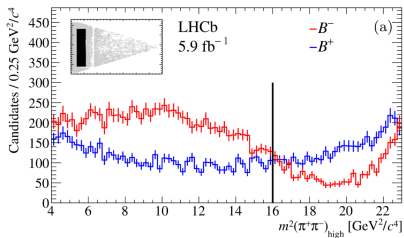
$B^+ \rightarrow h^+ h'^+ h'^-$ at LHCb



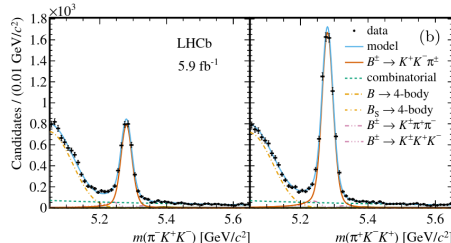
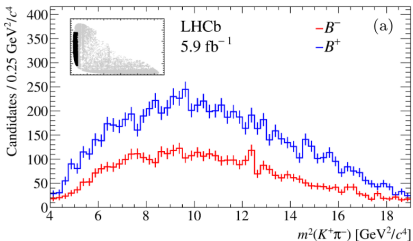
Phys. Rev. **D108** (2023) 012013, Phys. Rev. **D108** (2023) 012008

$B^+ \rightarrow \pi^+ h'^+ h'^-$ rescattering region

$$B^+ \rightarrow \pi^+ \pi^+ \pi^-$$



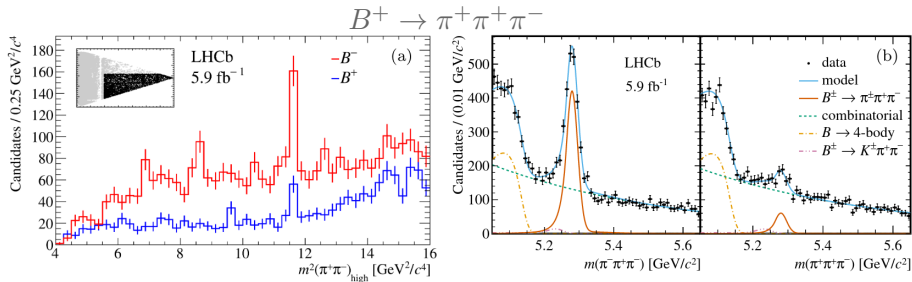
$$B^+ \rightarrow \pi^+ K^+ K^-$$



Phys. Rev. **D108** (2023) 012013, Phys. Rev. **D108** (2023) 012008

Clear opposite sign CP asymmetry in CPT -coupled $KK \leftrightarrow \pi\pi$ channels

$B^+ \rightarrow \pi^+ h'^+ h'^-$ charmonium region



Phys. Rev. **D108** (2023) 012013, Phys. Rev. **D108** (2023) 012008

Large amount of CP violation observed

Double-charm rescattering mechanism proposed

I. Bediaga, T. Frederico and P.C. Magalhães, PLB **806** (2020) 135490

Clear CP violation involving $\chi_{c0}(1P)$

Short-distance CP violation not expected

Large long-distance CP violation through interference predicted long ago

G. Eilam, M. Gronau and R.R. Mendel, PRL **74** (1995) 4984

1. Charmless 3-body B decays

- Impact of strong phase motion on CP violation structure
- Binned analysis of $B^+ \rightarrow h^+ h'^+ h'^-$

2. Charmless baryonic B decays

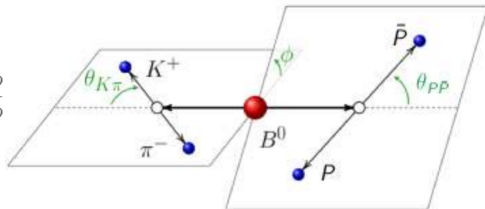
- Triple-product asymmetries
- $B^0 \rightarrow p\bar{p}K^+\pi^-$

Rich resonant structure in 4-body decays conducive to CP violation
 Amplitude analysis ambitious, begin with model-independent approach
 Probe CP violation with triple-product asymmetry measurements

P -odd triple products

$$B: C_{\hat{T}} = \vec{p}_{K^+} \cdot (\vec{p}_{\pi^-} \times \vec{p}_P) \propto \sin \phi$$

$$\bar{B}: \bar{C}_{\hat{T}} = \vec{p}_{K^-} \cdot (\vec{p}_{\pi^+} \times \vec{p}_{\bar{P}}) \propto \sin \phi$$



P -odd asymmetries of \hat{T} operator

$$B: A_{\hat{T}} = \frac{N(C_{\hat{T}} > 0) - N(C_{\hat{T}} < 0)}{N(C_{\hat{T}} > 0) + N(C_{\hat{T}} < 0)}$$

$$\bar{B}: \bar{A}_{\hat{T}} = \frac{\bar{N}(-\bar{C}_{\hat{T}} > 0) - \bar{N}(-\bar{C}_{\hat{T}} < 0)}{\bar{N}(-\bar{C}_{\hat{T}} > 0) + \bar{N}(-\bar{C}_{\hat{T}} < 0)}$$

P -odd observable

$$a_P^{\hat{T}\text{-odd}} = \frac{1}{2}(A_{\hat{T}} + \bar{A}_{\hat{T}})$$

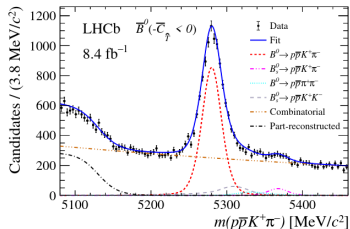
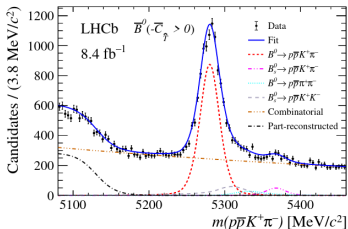
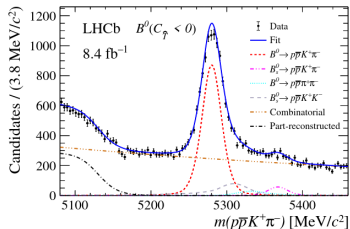
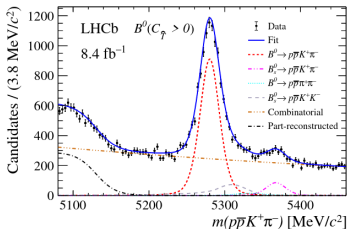
CP -odd observable

$$a_{CP}^{\hat{T}\text{-odd}} = \frac{1}{2}(A_{\hat{T}} - \bar{A}_{\hat{T}})$$

Sensitive to interference between P -even and P -odd amplitudes

$B^0 \rightarrow p\bar{p}K^+\pi^-$ at LHCb

Phys. Rev. **D108** (2023) 032007



$$a_P^{\hat{T}^{\text{odd}}} = [1.49 \pm 0.85 \text{ (stat)} \pm 0.08 \text{ (syst)}]\%$$

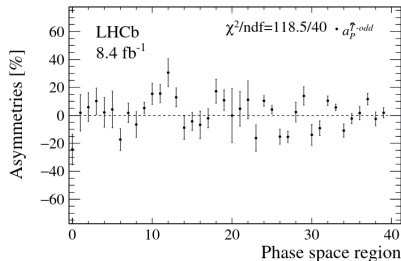
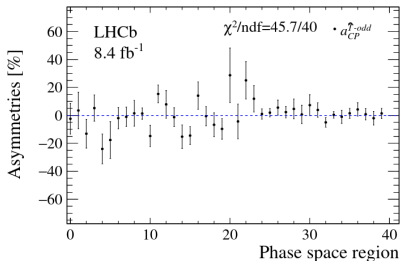
$$a_{CP}^{\hat{T}^{\text{odd}}} = [0.51 \pm 0.85 \text{ (stat)} \pm 0.08 \text{ (syst)}]\%$$

$B^0 \rightarrow p\bar{p}K^+\pi^-$ at LHCb

Can enhance sensitivity by dividing phase space into bins

Based on $K^+\pi^-$ mass and helicity angles

Phys. Rev. **D108** (2023) 032007



Control sample $B^0 \rightarrow p\bar{p}\bar{D}^0[K^+\pi^-]$ estimates detection asymmetry

Significant P violation in low $p\bar{p}$ and $K^*(892)^0$ mass regions

P conservation rejected at $\sim 6\sigma$

CP conserved within $\sim 1\sigma$

Theoretical prediction up to 20% level not excluded

Int. J. Mod. Phys. **23** (2008) 3290, Phys. Rev. Lett. **98** (2007) 011801,

Eur. Phys. J. **C80** (2020) 565

Multibody decays provide excellent environment for CP violation studies

- Large-valued measurements more accessible

- Diverse structures provide another view to strong phase motion

Recent results from LHCb

Model-independent study of 3-body charmless hadronic B^\pm decays

- Attention should turn to modelling rescattering regions

- CP violation driven by long-distance $\chi_{c0}(1P)$ interactions foreseen

Triple-product asymmetries in $B^0 \rightarrow p\bar{p}K^+\pi^-$ decays

- World-first measurement

- Still no observation of CP violation in baryonic B decays at this time

- Predicted up to the 20% level