

Aligned Two Higgs Doublet Model and the Global Fits

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In Collaboration With: **Victor Miralles** and **Antonio Pich**

ArXiv: [2307.15419](https://arxiv.org/abs/2307.15419)

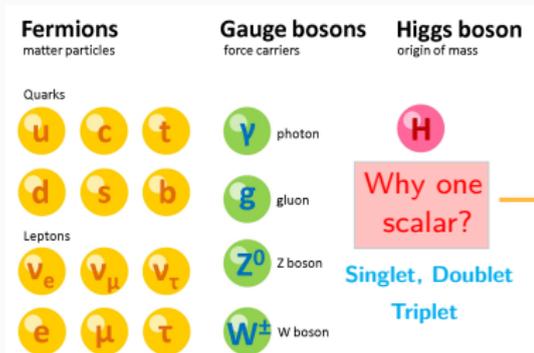
Updated version of Eberhardt, Peñuelas, Pich JHEP 05 (2021) 005.

IFIC (CSIC – Universitat de València)



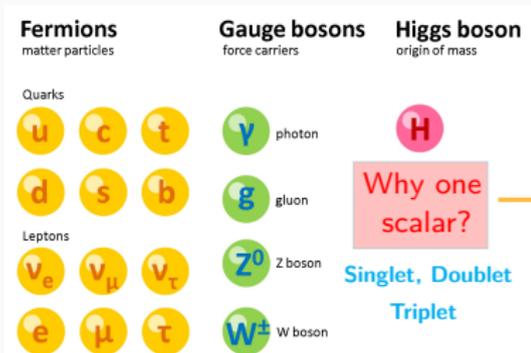
EPS-HEP2023, Universität Hamburg
23rd August, 2023

Motivation



$$\rho = \frac{\sum_i c_i \langle \phi_i^0 \rangle^2 [T_i(T_i+1) - Y_i^2]}{2 \sum_i \langle \phi_i^0 \rangle^2 Y_i^2} \rightarrow 2\text{HDM}$$

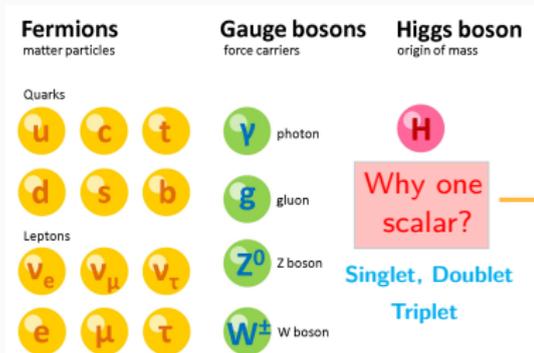
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Prospects: New sources of CP violation, Axion-like phenomenology, Dark matter aspects, Electroweak Baryogenesis, Stability of scalar potential till Planck scale, EFT for SUSY, etc.

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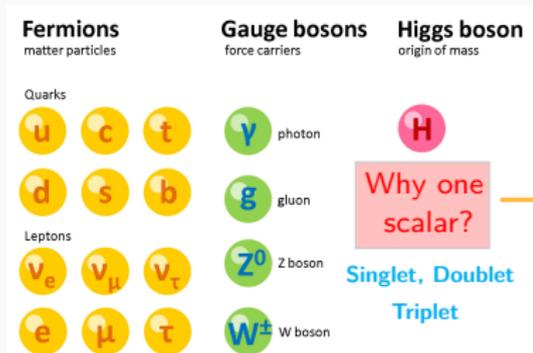


$$\rho = \frac{\sum_i c_i \langle \phi_i^0 \rangle^2 [\mathcal{T}_i(\mathcal{T}_i+1) - \mathcal{Y}_i^2]}{2 \sum_i \langle \phi_i^0 \rangle^2 \mathcal{Y}_i^2} \rightarrow 2\text{HDM}$$

► **Prospects:** New sources of CP violation, Axion-like phenomenology, Dark matter aspects, Electroweak Baryogenesis, Stability of scalar potential till Planck scale, EFT for SUSY, etc.

► **Problems:** FCNC

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► **Problems:** FCNC

► **Solutions:** 1) Additional \mathcal{Z}_2 symmetry, 2) **A2HDM**

Pich, Tuzon PRD 80 (2009) 091702; Ferreira, Lavoura, Silva PLB 688 (2010) 341; Jung, Pich, Tuzon JHEP 11 (2010) 003; Braeuninger, Ibarra, Simonetto PLB 692 (2010) 189; Bijnens, Lu, Rathsman JHEP 05 (2012) 118; Li, Lu, Pich JHEP 06 (2014) 022; Abbas, et al. JHEP 06 (2015) 005; Botella, et al. EPJC 75 (2015) 286; Gori, Haber, Santos JHEP 06 (2017) 110; Kanemura, Mondal, Yagyu JHEP 02 (2023) 237; etc...

Scalar Potential

$$\phi_a : \langle 0 | \phi_a^T | 0 \rangle = (0, v_a e^{i\theta_a}) \quad a \in \{1, 2\}$$

Global $SU(2) \implies$ "Higgs basis"

$$\Phi_a : \quad \Phi_1 = \frac{1}{\sqrt{2}} \begin{pmatrix} \sqrt{2} G^+ \\ S_1 + v + i G^0 \end{pmatrix}, \quad \Phi_2 = \frac{1}{\sqrt{2}} \begin{pmatrix} \sqrt{2} H^+ \\ S_2 + i S_3 \end{pmatrix}$$

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Scalar Potential:

$$V = \mu_1 \Phi_1^\dagger \Phi_1 + \mu_2 \Phi_2^\dagger \Phi_2 + [\mu_3 \Phi_1^\dagger \Phi_2 + h.c.] + \frac{\lambda_1}{2} (\Phi_1^\dagger \Phi_1)^2 + \frac{\lambda_2}{2} (\Phi_2^\dagger \Phi_2)^2 + \lambda_3 (\Phi_1^\dagger \Phi_1)(\Phi_2^\dagger \Phi_2) \\ + \lambda_4 (\Phi_1^\dagger \Phi_2)(\Phi_2^\dagger \Phi_1) + \left[\left(\frac{\lambda_5}{2} \Phi_1^\dagger \Phi_2 + \lambda_6 \Phi_1^\dagger \Phi_1 + \lambda_7 \Phi_2^\dagger \Phi_2 \right) (\Phi_1^\dagger \Phi_2) + h.c. \right].$$

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Gauge-Higgs Coupling:

$$g_{hVV} = \cos \tilde{\alpha} g_{hVV}^{SM}, \quad g_{HVV} = -\sin \tilde{\alpha} g_{hVV}^{SM}, \quad g_{AVV} = 0, \quad VV \equiv (W^+ W^-, ZZ)$$

Independent parameters for scalar potential

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* **Masses:**

$$M_{H^\pm}^2 = \mu_2 + \frac{\lambda_3}{2} v^2, \quad M_{h,H}^2 = \frac{1}{2} (\Sigma \mp \Delta), \quad M_A^2 = M_{H^\pm}^2 + \frac{v^2}{2} (\lambda_4 - \lambda_5),$$

with $\Sigma = M_{H^\pm}^2 + \left(\lambda_1 + \frac{\lambda_4}{2} + \frac{\lambda_5}{2} \right) v^2$ and $\Delta = \sqrt{(\Sigma - 2\lambda_1 v^2)^2 + 4\lambda_6^2 v^4}$.

* **Mixing angle:** $\tan \tilde{\alpha} = \frac{M_h^2 - v^2 \lambda_1}{v^2 \lambda_6} = \frac{v^2 \lambda_6}{v^2 \lambda_1 - M_H^2}$.

* **Parameter set:** $v, M_h, M_{H^\pm}, M_H, M_A, \tilde{\alpha}, \lambda_2, \lambda_3$ and λ_7

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* **Parameter set:** ~~μ_1, μ_3~~ , $M_{H^\pm}, M_H, M_A, \tilde{\alpha}, \lambda_2, \lambda_3$ and $\lambda_7 \implies$ 7 parameters.

Fermionic interaction

✿ Yukawa interaction:

$$\begin{aligned}
 -\mathcal{L}_Y &= \left(1 + \frac{S_1}{v}\right) \left\{ \bar{u}_L M_u u_R + \bar{d}_L M_d d_R + \bar{\ell}_L M_\ell \ell_R \right\} \\
 &+ \frac{1}{v} (S_2 + iS_3) \left\{ \bar{u}_L Y_u u_R + \bar{d}_L Y_d d_R + \bar{\ell}_L Y_\ell \ell_R \right\} \\
 &+ \frac{\sqrt{2}}{v} H^+ \left\{ \bar{u}_L V Y_d d_R - \bar{u}_R Y_u^\dagger V d_L + \bar{\nu}_L Y_\ell \ell_R \right\} + \text{h.c.},
 \end{aligned}$$

✿ Alignment:

$$Y_u = \varsigma_u^* M_u \quad \text{and} \quad Y_{d,\ell} = \varsigma_{d,\ell} M_{d,\ell},$$

$$-\mathcal{L}_Y = \sum_{i,f} \left(\frac{y_f^{\varphi_i^0}}{v} \right) \varphi_i^0 \left[\bar{f} M_f \mathcal{P}_R f \right] + \left(\frac{\sqrt{2}}{v} \right) H^+ \left[\bar{u} \left\{ \varsigma_d V M_d \mathcal{P}_R - \varsigma_u M_u^\dagger V \mathcal{P}_L \right\} d + \varsigma_\ell \bar{\nu} M_\ell \mathcal{P}_R \ell \right] + \text{h.c.}$$

$$\begin{aligned}
 y_u^H &= -\sin \tilde{\alpha} + \varsigma_u^* \cos \tilde{\alpha}, & y_u^h &= \cos \tilde{\alpha} + \varsigma_u^* \sin \tilde{\alpha}, & y_u^A &= -i\varsigma_u^*, \\
 y_{d,\ell}^H &= -\sin \tilde{\alpha} + \varsigma_{d,\ell} \cos \tilde{\alpha}, & y_{d,\ell}^h &= \cos \tilde{\alpha} + \varsigma_{d,\ell} \sin \tilde{\alpha}, & y_{d,\ell}^A &= i\varsigma_{d,\ell}.
 \end{aligned}$$

Type I: $\varsigma_u = \varsigma_d = \varsigma_\ell = \cot \beta$, **Type II:** $\varsigma_u = -\frac{1}{\varsigma_d} = -\frac{1}{\varsigma_\ell} = \cot \beta$, **Inert:** $\varsigma_u = \varsigma_d = \varsigma_\ell = 0$,

Type X: $\varsigma_u = \varsigma_d = -\frac{1}{\varsigma_\ell} = \cot \beta$ and **Type Y:** $\varsigma_u = -\frac{1}{\varsigma_d} = \varsigma_\ell = \cot \beta$.

Package: HEPfit (Bayesian approach)

| Priors | | | |
|--|--|--|--------------------------------------|
| $M_{H^\pm} \subset [0.125, 1.0 (1.5)] \text{ TeV}$ | $M_H \subset [0.125, 1.0 (1.5)] \text{ TeV}$ | $M_A \subset [0.125, 1.0 (1.5)] \text{ TeV}$ | |
| $\lambda_2 \subset [0, 11]$ | $\lambda_3 \subset [-3, 17]$ | $\lambda_7 \subset [-5, 5]$ | |
| $\tilde{\alpha} \subset [-0.16, 0.16]$ | $\varsigma_u \subset [-1.5, 1.5]$ | $\varsigma_d \subset [-50, 50]$ | $\varsigma_\ell \subset [-100, 100]$ |

- ⊠ Linear prior on masses.
- ⊠ Ranges for quartic couplings and alignment parameters are chosen from theoretical constraints.
- ⊠ Range of $\tilde{\alpha}$ is chosen to incorporate the 5σ region.
- ⊠ Taken the experimental values of CKM matrix element carefully and fitted the Wolfenstein parameters.

Theoretical Constraints: Stability

- **Bounded from below:**

$$V = -M_\mu r^\mu + \frac{1}{2} \Lambda^\mu{}_\nu r_\mu r^\nu, \text{ where,}$$

$$M_\mu = \left(-\frac{\mu_1 + \mu_2}{2}, -\operatorname{Re} \mu_3, \operatorname{Im} \mu_3, -\frac{\mu_1 - \mu_2}{2} \right),$$

$$r^\mu = \left(|\Phi_1|^2 + |\Phi_2|^2, 2\operatorname{Re}(\Phi_1^\dagger \Phi_2), 2\operatorname{Im}(\Phi_1^\dagger \Phi_2), |\Phi_1|^2 - |\Phi_2|^2 \right),$$

$$\Lambda^\mu{}_\nu = \frac{1}{2} \begin{pmatrix} \frac{1}{2}(\lambda_1 + \lambda_2) + \lambda_3 & \operatorname{Re}(\lambda_6 + \lambda_7) & -\operatorname{Im}(\lambda_6 + \lambda_7) & \frac{1}{2}(\lambda_1 - \lambda_2) \\ -\operatorname{Re}(\lambda_6 + \lambda_7) & -\lambda_4 - \operatorname{Re}\lambda_5 & \operatorname{Im}\lambda_5 & -\operatorname{Re}(\lambda_6 - \lambda_7) \\ \operatorname{Im}(\lambda_6 + \lambda_7) & \operatorname{Im}\lambda_5 & -\lambda_4 + \operatorname{Re}\lambda_5 & \operatorname{Im}(\lambda_6 - \lambda_7) \\ -\frac{1}{2}(\lambda_1 - \lambda_2) & -\operatorname{Re}(\lambda_6 - \lambda_7) & \operatorname{Im}(\lambda_6 - \lambda_7) & -\frac{1}{2}(\lambda_1 + \lambda_2) + \lambda_3 \end{pmatrix}.$$

- ✓ The **necessary & sufficient** conditions for bounded below potential :

1. All the eigenvalues ($\Lambda_{0,1,2,3}$) of $\Lambda^\mu{}_\nu$ are real.
2. $\Lambda_0 > 0$ with $\Lambda_0 > \Lambda_i \forall i \in \{1, 2, 3\}$.

For several necessary conditions: [H. Bahl, et al., JHEP 03 \(2023\) 165](#)

- **Absolute stability:**

$$D = \operatorname{Det}[\xi \mathbb{I}_4 - \Lambda^\mu{}_\nu] = -\prod_{k=0}^3 (\xi - \Lambda_k) \quad \text{with} \quad \xi = \frac{m_{H^\pm}^2}{v^2}.$$

- ✓ The conditions for global minimum : 1) $D > 0$, or 2) $D < 0$ with $\xi > \Lambda_0$.

Theoretical Constraints: Perturbativity

- Tree-level partial-wave amplitudes:

$$(a_0)_{i,f} = \frac{1}{16\pi s} \int_{-s}^0 dt \mathcal{M}_{i \rightarrow f}(s, t)$$

- Eigenvalue in the j th partial wave:

$$(a_j^0)^2 \leq \frac{1}{4}$$

- There are fourteen neutral, eight single-charged and three doubly-charged two-body scalar $2 \rightarrow 2$ scattering states. $\implies a_0^{(0,+,++)} = \oplus \{X_{Y,\sigma}\}$

$$X_{(1,0)} = \lambda_3 - \lambda_4,$$

$$X_{(1,1)} = \begin{pmatrix} \lambda_1 & \lambda_5 & \sqrt{2}\lambda_6 \\ \lambda_5^* & \lambda_2 & \sqrt{2}\lambda_7^* \\ \sqrt{2}\lambda_6^* & \sqrt{2}\lambda_7^* & \lambda_3 + \lambda_4 \end{pmatrix}, \quad X_{(0,1)} = \begin{pmatrix} \lambda_1 & \lambda_4 & \lambda_6 & \lambda_6^* \\ \lambda_4 & \lambda_2 & \lambda_7 & \lambda_7^* \\ \lambda_6^* & \lambda_7^* & \lambda_3 & \lambda_5^* \\ \lambda_6 & \lambda_7 & \lambda_5 & \lambda_3 \end{pmatrix},$$

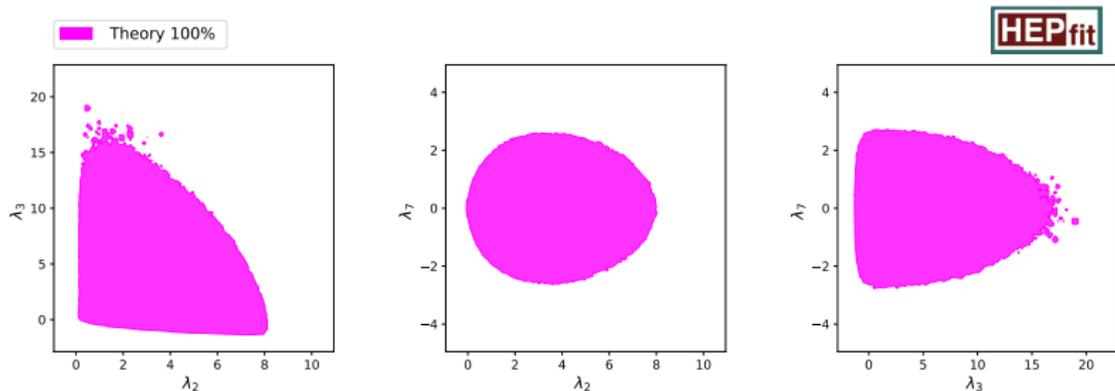
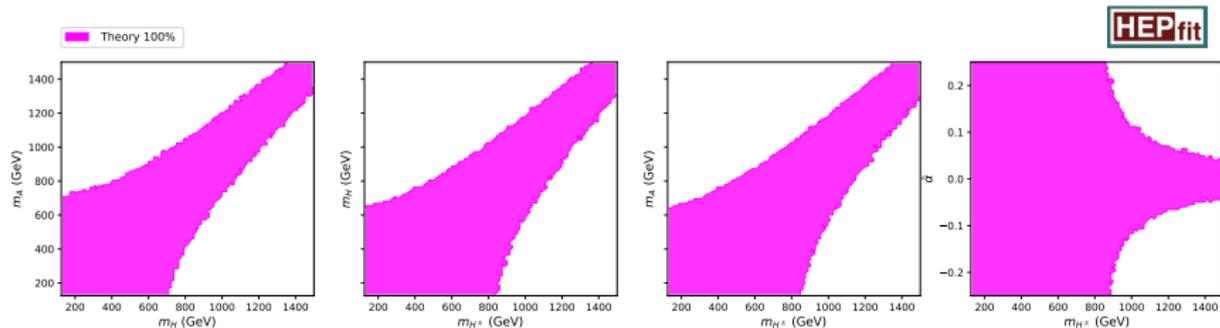
$$X_{(0,0)} = \begin{pmatrix} 3\lambda_1 & 2\lambda_3 + \lambda_4 & 3\lambda_6 & 3\lambda_6^* \\ 2\lambda_3 + \lambda_4 & 3\lambda_2 & 3\lambda_7 & 3\lambda_7^* \\ 3\lambda_6^* & 3\lambda_7^* & \lambda_3 + 2\lambda_4 & 3\lambda_5^* \\ 3\lambda_6 & 3\lambda_7 & 3\lambda_5 & \lambda_3 + 2\lambda_4 \end{pmatrix}.$$

- Eigenvalue: $|e_i| < 8\pi$.

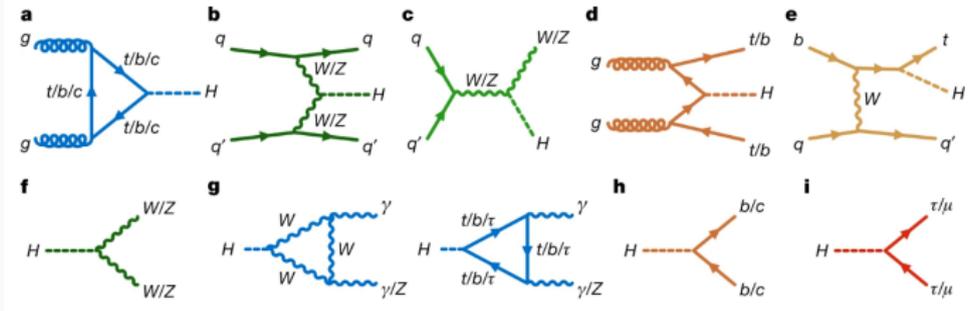
Ginzburg, Ivanov, PRD 72 (2005) 115010; H. Bahl, et al., JHEP 03 (2023) 165

- Charged scalar coupling to fermions: $\sqrt{2} |\zeta_f| m_f / v < 1$

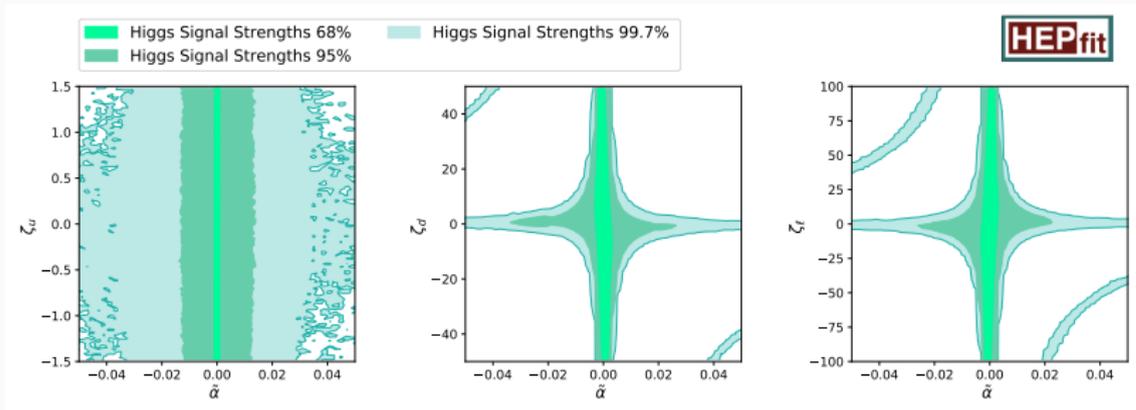
Theoretical Constraints: Fit Results



Signal Strength

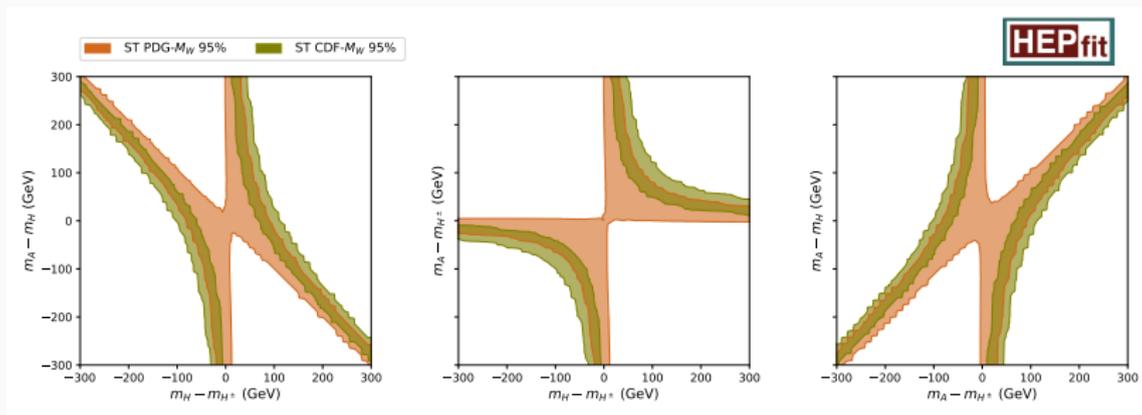


ATLAS, Nature 607(2022) 52–59



Data: ATLAS and CMS

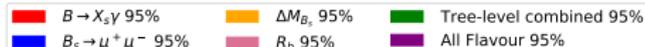
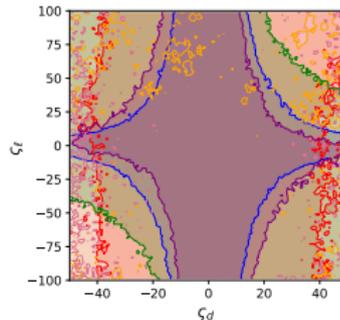
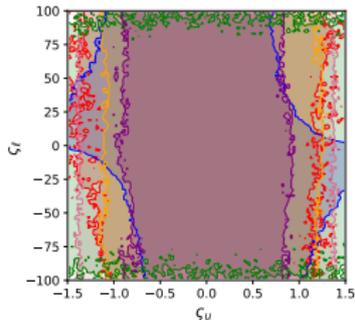
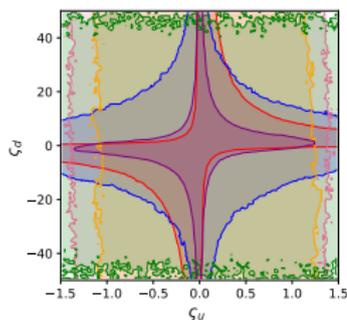
- ♣ M_W (PDG): [PDG, PTEP 2022 \(2022\) 083C01](#)
- ♣ M_W (CDF): [CDF, Science 376 \(2022\) 170](#); [Blas, et al. PRL 129 \(2022\) 271801](#)



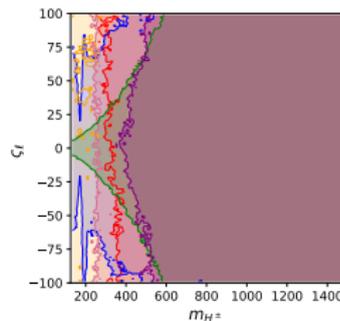
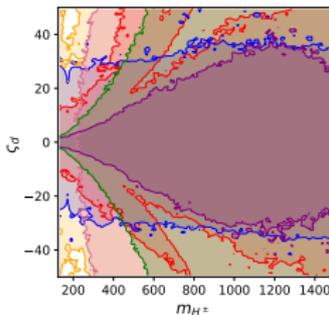
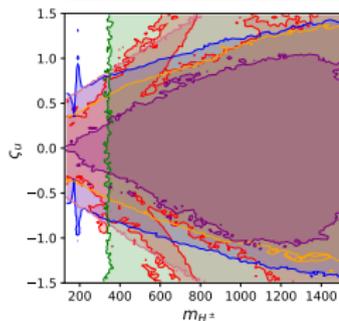
Flavour Observables



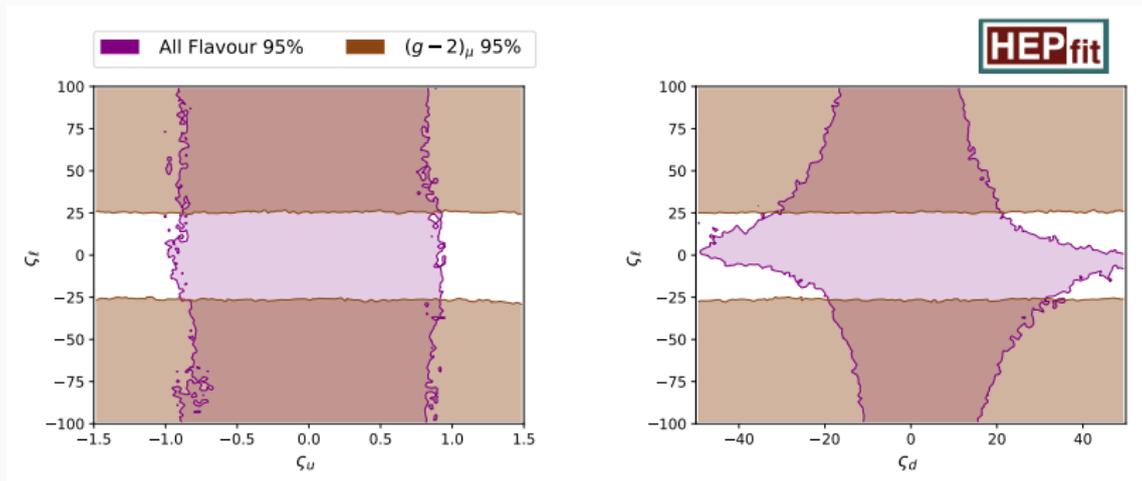
HEPfit



HEPfit

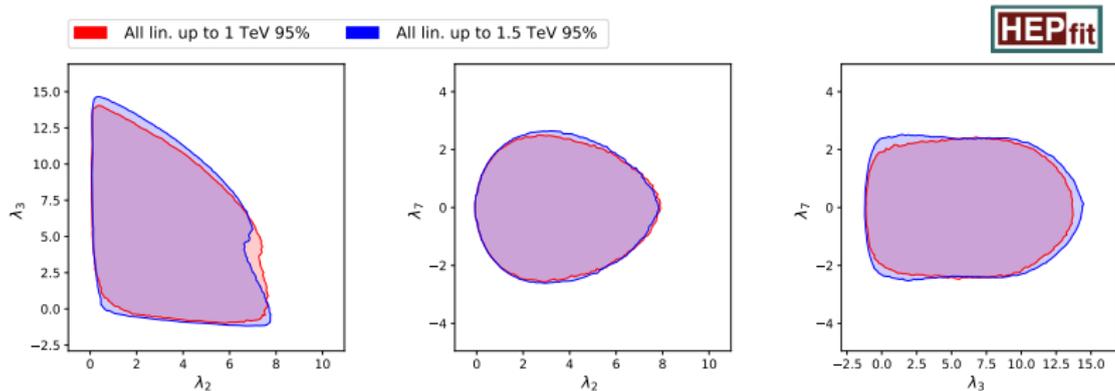
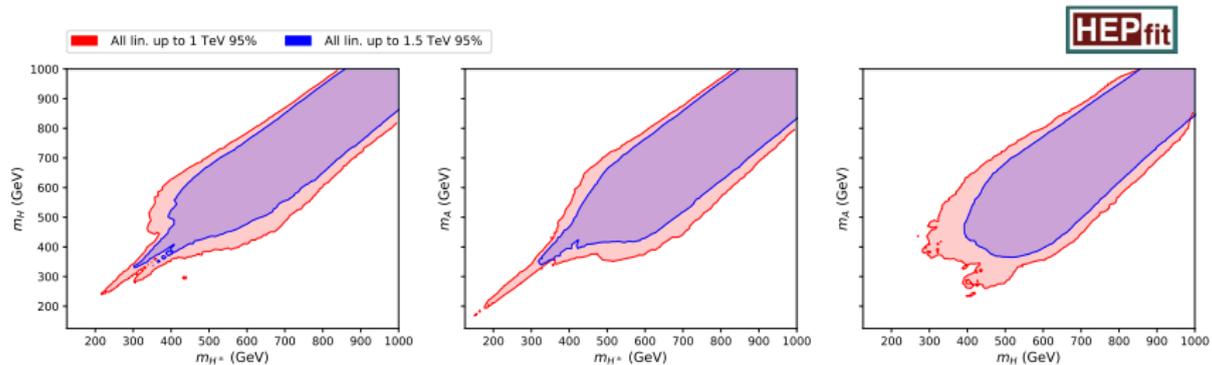


Flavour Observables: muon g-2



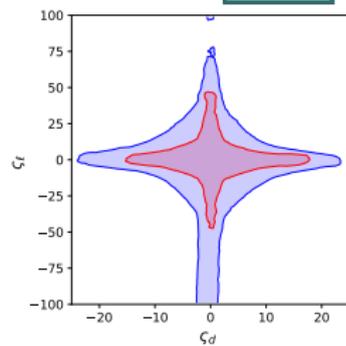
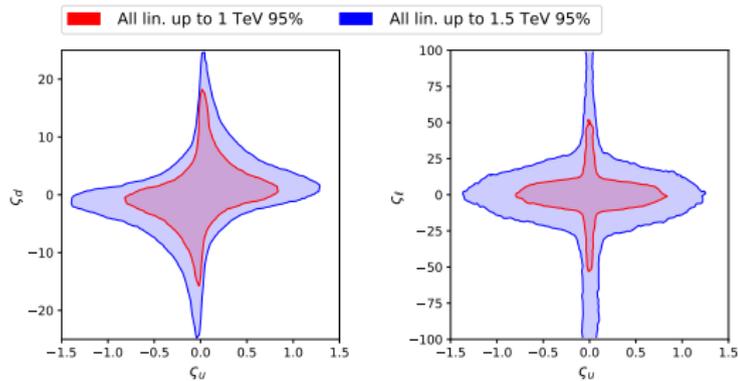
- Theoretical Constraints
- EWPO
- Flavour Observables
- Higgs Signal Strength
- Direct Searches

Global Fits

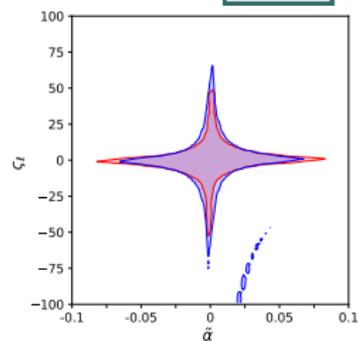
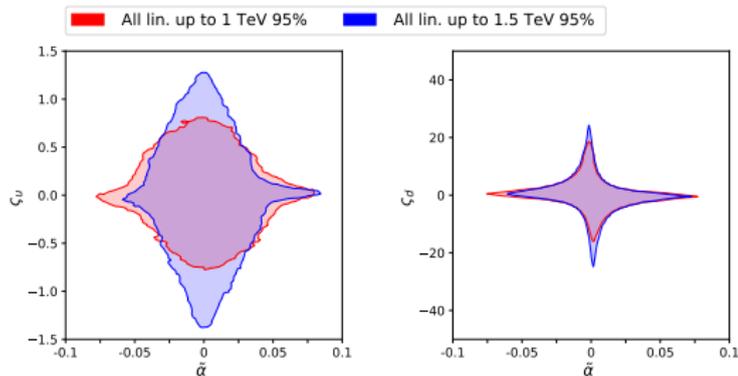


Global Fits

HEPfit



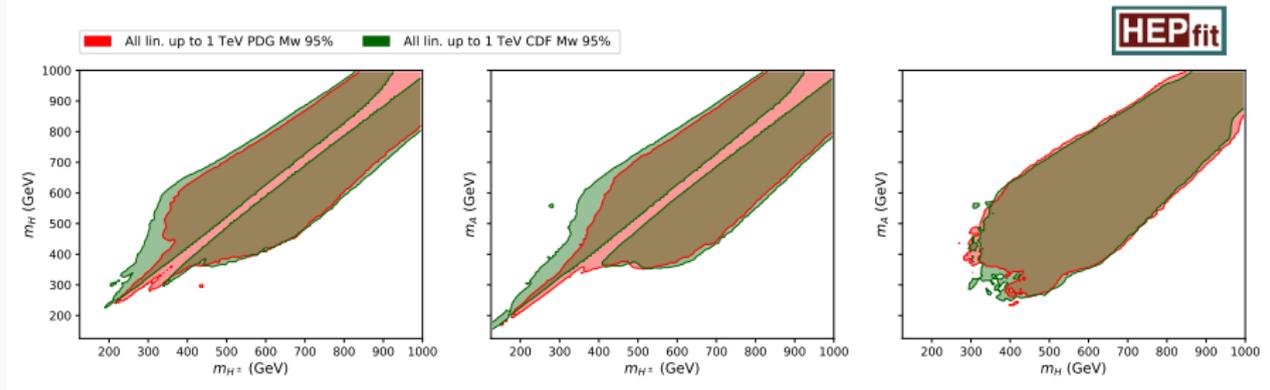
HEPfit



Fit Results

| Marginalised Individual results | | | |
|---|-----------------------------------|-----------------------------------|-----------------------------------|
| <i>Masses up to 1 TeV</i> | | | |
| $M_{H^\pm} \geq 390 \text{ GeV (95\%)}$ | $M_H \geq 410 \text{ GeV (95\%)}$ | $M_A \geq 370 \text{ GeV (95\%)}$ | |
| $\lambda_2: 3.2 \pm 1.9$ | $\lambda_3: 5.9 \pm 3.5$ | $\lambda_7: 0.0 \pm 1.1$ | |
| $\tilde{\alpha}: (0.05 \pm 21.0) \cdot 10^{-3}$ | $\varsigma_u: 0.006 \pm 0.257$ | $\varsigma_d: 0.12 \pm 4.12$ | $\varsigma_\ell: -0.39 \pm 11.69$ |
| <i>Masses up to 1.5 TeV</i> | | | |
| $M_{H^\pm} \geq 480 \text{ GeV (95\%)}$ | $M_H \geq 490 \text{ GeV (95\%)}$ | $M_A \geq 480 \text{ GeV (95\%)}$ | |
| $\lambda_2: 3.2 \pm 1.9$ | $\lambda_3: 5.9 \pm 3.8$ | $\lambda_7: 0.0 \pm 1.2$ | |
| $\tilde{\alpha}: (0.8 \pm 16.8) \cdot 10^{-3}$ | $\varsigma_u: -0.011 \pm 0.407$ | $\varsigma_d: -0.096 \pm 6.22$ | $\varsigma_\ell: -1.18 \pm 17.54$ |

CDF measurement of M_W



Summary

- ✍ We fitted total 10 parameters: 3 masses, 3 quartic couplings, 3 alignment parameters and one angle.
- ✍ The lower limit on masses are constrained from direct searches and flavour observables.
- ✍ Mass splittings are constrained by EWPO and theoretical bounds.
- ✍ Quartic couplings are bounded by theoretical constraints.
- ✍ Alignment parameters are constrained by direct searches and flavour observables.
- ✍ Mixing angle is restricted by Higgs signal strength.
- ✍ We also studied the effects of $(g - 2)_\mu$ and CDF measurement of M_W , but did not include into the fit. However, parameter-space of A2HDM has the flexibility to accommodate both.



Back up

V_{ud} : $0^+ \rightarrow 0^+$ nuclear β decay

V_{ud} : Semi-leptonic decays of kaons to electrons (not muons). Cabibbo Anomaly \rightarrow
Double the error.

V_{cd} : Semileptonic (Not the leptonic) decays of D and neutrino scattering.

V_{cs} : Semileptonic (Not the leptonic) decays of D_s

V_{ub} : PDG

V_{cb} : PDG

V_{td}/V_{ts} : PDG; The ratio kills the NP contribution in $B - \bar{B}$ mixing.