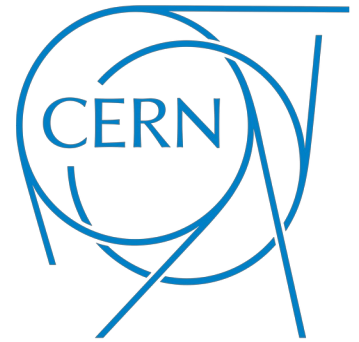


# Direct determination of the rapidity anomalous dimension. Experimental and phenomenological applications.



European Physical Society Conference on High Energy Physics 2023

21.08.2023

Armando Bermúdez Martínez

Based on the recent work:

Phys.Rev.D 106 (2022) 9, L091501

arXiv:2307.06704

<https://www.desy.de/f/students/2022/reports/David.Gutierrez.pdf>



European Physical Society

**Conference on High Energy Physics**

21-25 August 2023

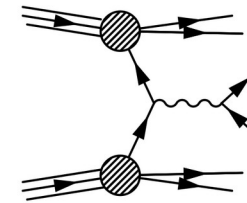


# From factorization to nucleon tomography

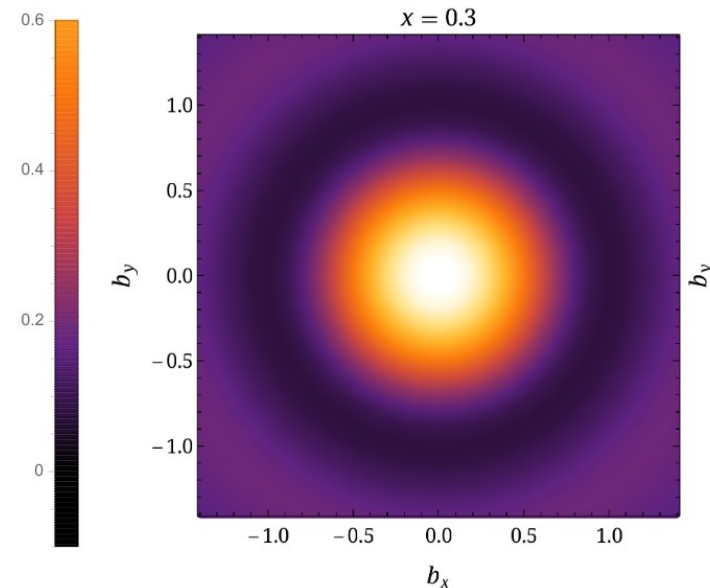
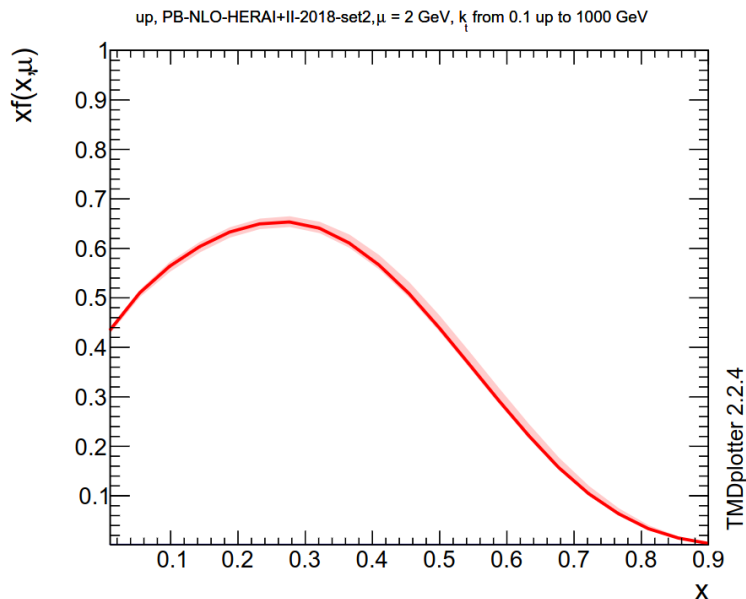
# From factorization to nucleon tomography

- ▶ Modern factorization theorems separate the hadron structure from the low distance hard scattering

$$\frac{d\sigma}{dp_T} = \sigma_0 \int \frac{d^2b}{(2\pi)^2} e^{i(bp_T)} C\left(\frac{Q}{\mu}\right) F_1(x_1, b; \mu, \zeta) F_2(x_2, b; \mu, \zeta)$$



- ▶ More complete description, going **beyond the simplest 1D parton structure**



# From factorization to nucleon tomography

- ▶ A new function emerges and it dictates the evolution of the parton distributions:

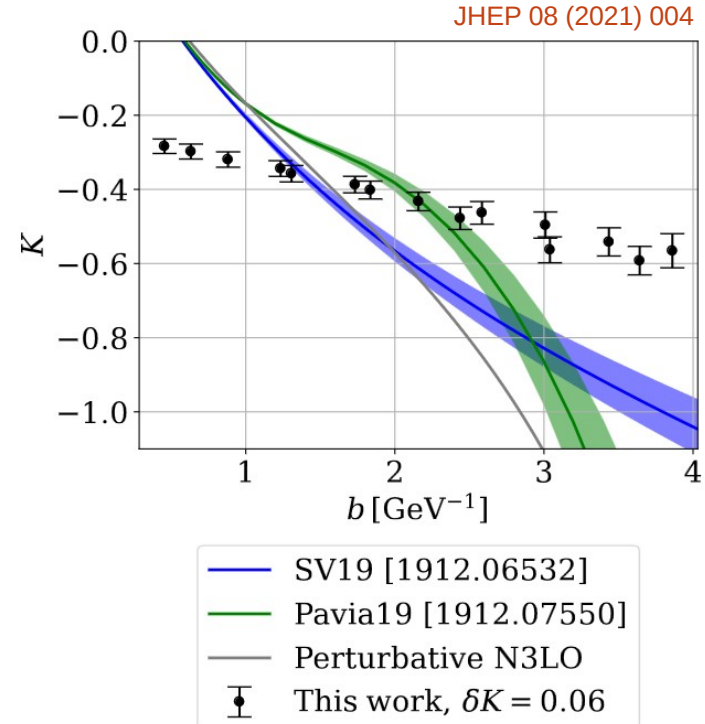
$$\frac{df_{q,h}(x, b; \mu, \zeta)}{d \ln \mu^2} = \frac{\gamma_F(\mu, \zeta)}{2} f_{q,h}(x, b; \mu, \zeta),$$

$$\frac{df_{q,h}(x, b; \mu, \zeta)}{d \ln \zeta} = -\mathcal{D}(b, \mu) f_{q,h}(x, b; \mu, \zeta).$$

- ▶ It is a self-contained object with new non-perturbative information

Phys. Rev. Lett. 125, 192002 (2020)

- ▶ RAD has been studied extensively
- ▶ Yet, only QCD function which is largely unknown



# From factorization to nucleon tomography

- ▶ A new function emerges and it dictates the evolution of the parton distributions:

$$\frac{df_{q,h}(x, b; \mu, \zeta)}{d \ln \mu^2} = \frac{\gamma_F(\mu, \zeta)}{2} f_{q,h}(x, b; \mu, \zeta),$$

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- ▶ It is a self-contained object with new non-perturbative information

Phys. Rev. Lett. 125, 192002 (2020)

- ▶ A direct measurement of RAD would imply:

- Stringent test of factorization and universality of the 3D structure
- Higher precision imaging of hadrons
- Higher precision for measurements, e.g W mass
- Input to probe parton spin-orbit correlations

...

# Novel method to determine RAD

Phys.Rev.D 106 (2022) 9, L091501

# Novel method to determine RAD

- ▶ Start from the factorization formula:

$$\frac{d\sigma}{dQ^2 dy dq_T^2} = \frac{2\pi \alpha_{\text{em}}^2(Q)}{9 s Q^2} W_{f_1 f_2}(x_1, x_2, Q, q_T),$$

- ▶ Apply the inverse Hankel transform:

$$\Sigma(y, Q; b) = \frac{2\pi \alpha_{\text{em}}^2(Q)}{9 s Q^2} |C_V(Q)|^2 \sum_q e_q^2 f_{q_1}(x_1, b; Q, Q^2) f_{q_2}(x_2, b; Q, Q^2)$$

- ▶ Evolve the parton distribution to a reference scale:

$$\Sigma(y, Q; b) = \frac{2\pi \alpha_{\text{em}}^2(Q)}{9 s Q^2} |C_V(Q)|^2 e^{2\Delta(b; Q \rightarrow (\mu_0, \zeta_0))} \sum_q e_q^2 f_{q_1}(x_1, b; \mu_0, \zeta_0) f_{q_2}(x_2, b; \mu_0, \zeta_0)$$

the target
disposable

- ▶ Build ratios of the cross sections at different scales:

$$\frac{\Sigma_1}{\Sigma_2} = \frac{\frac{\alpha_{\text{em}}^2(Q_1)}{s_1 Q_1^2} |C_V(Q_1)|^2 e^{2\Delta(b; Q_1 \rightarrow (\mu_0, \zeta_0))}}{\frac{\alpha_{\text{em}}^2(Q_2)}{s_2 Q_2^2} |C_V(Q_2)|^2 e^{2\Delta(b; Q_2 \rightarrow (\mu_0, \zeta_0))}} \frac{\sum e_q^2 f_{q_1}(x_1, b; \mu_0, \zeta_0) f_{q_2}(x_2, b; \mu_0, \zeta_0)}{\sum e_q^2 f_{q_1}(x_1, b; \mu_0, \zeta_0) f_{q_2}(x_2, b; \mu_0, \zeta_0)}$$





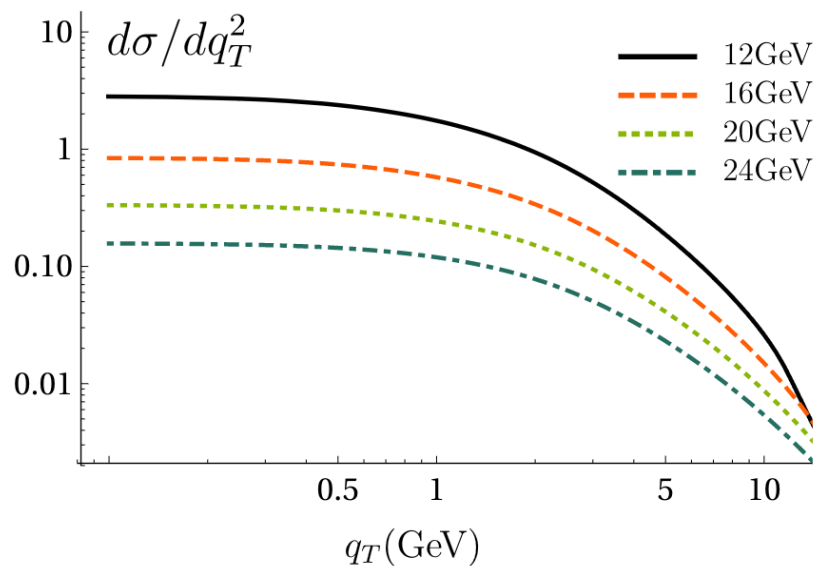
# Applying the method to simulated data

Phys.Rev.D 106 (2022) 9, L091501

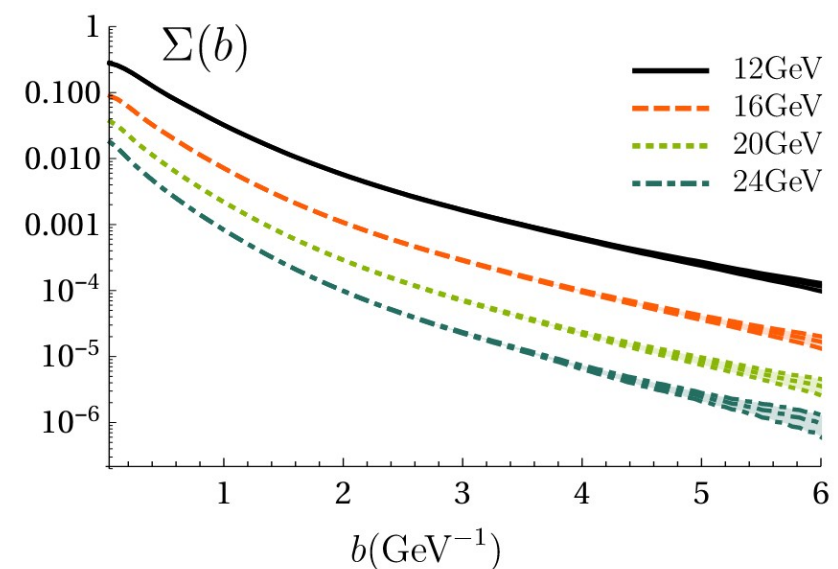
# Applying the method to simulated data

Phys.Rev.D 106 (2022) 9, L091501

- ▶ Master formula can be used with data, provided:
  - small  $q_T$  and  $Q$  bin sizes
  - choices of  $y$ ,  $Q$  and center-of-mass energy ensure same  $x$  range
  - $Q$  below  $Z$  peak
- ▶ Simulation using the CASCADE MC generator:

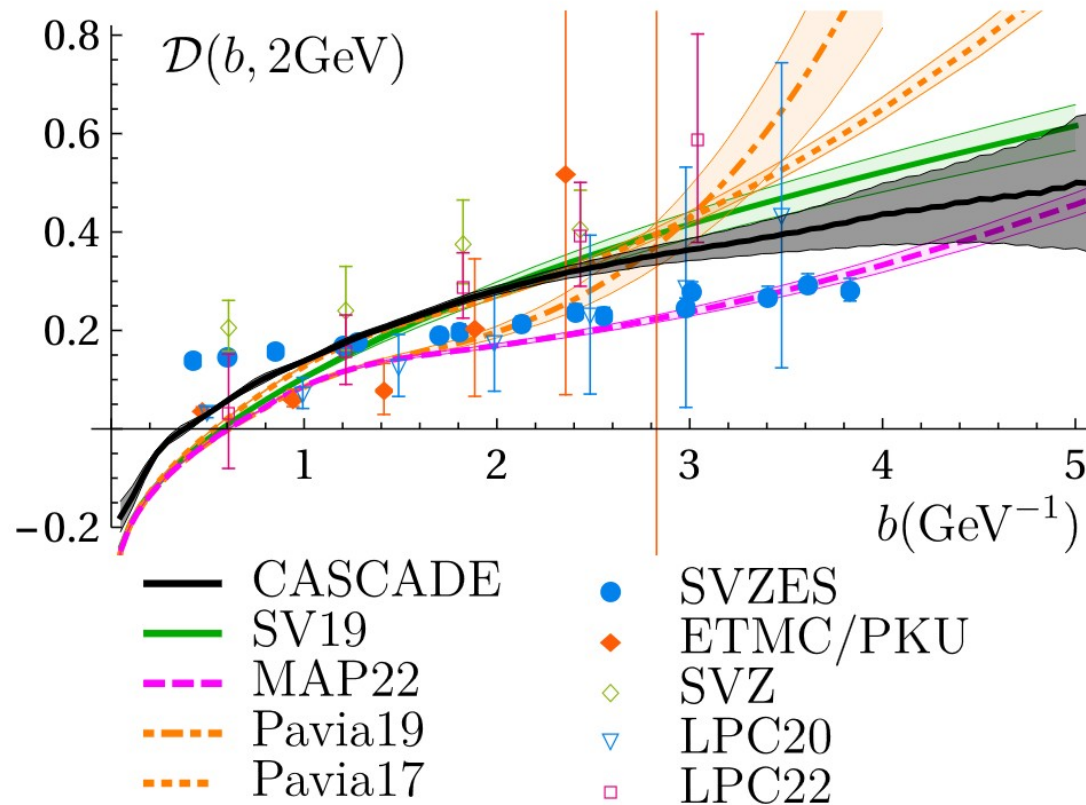


→  
Inverse  
Hankel



# Applying the method to simulated data

- ▶ All properties of RAD, like universality, are observed for the PB approach
- ▶ This non-trivially supports both factorization and PB approaches, sets the stage for a comparison
- ▶ The method can be applied to the experimental data!



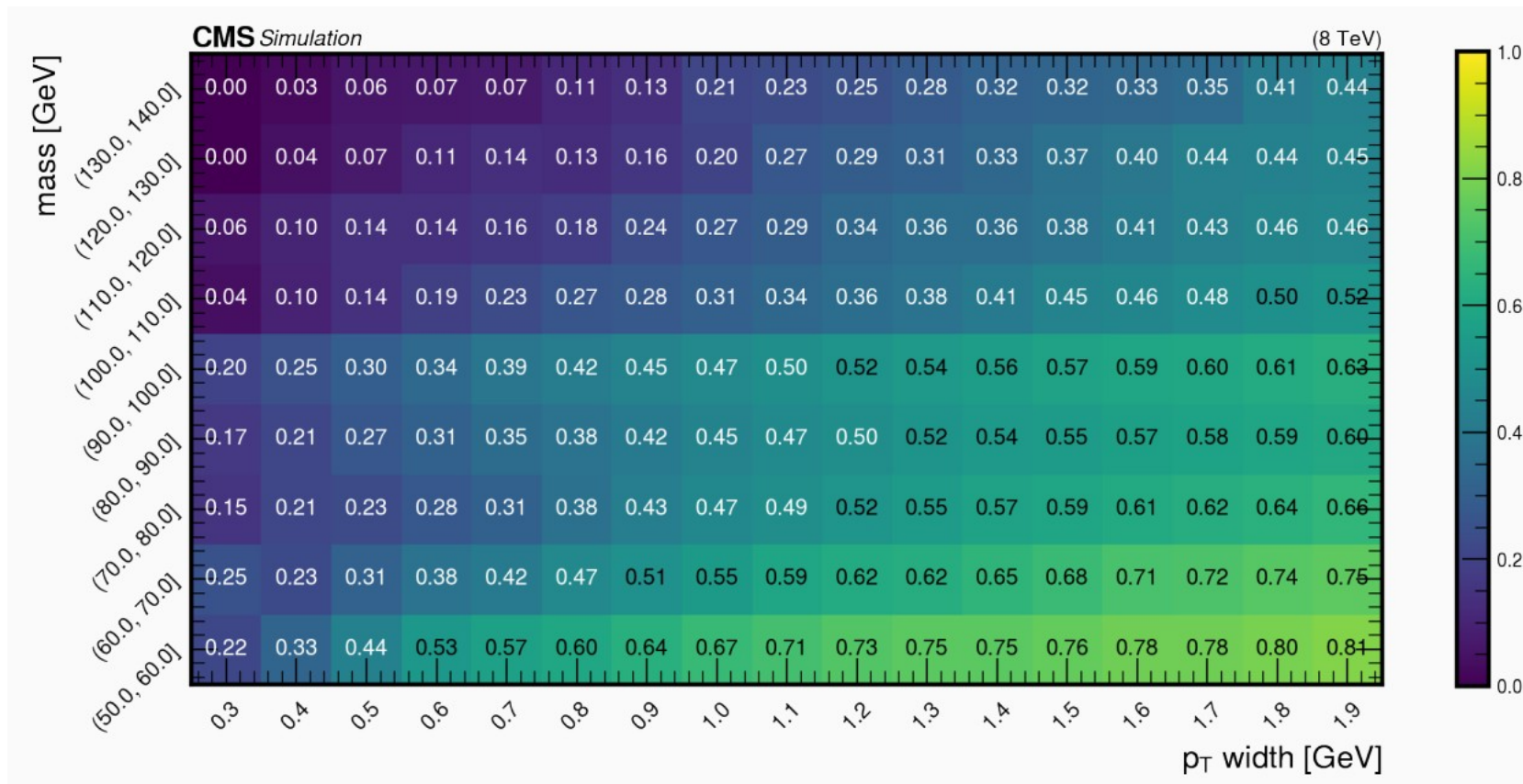
# Applying the method to **experimental** data

<https://www.desy.de/f/students/2022/reports/David.Gutierrez.pdf>

# Applying the method to experimental data

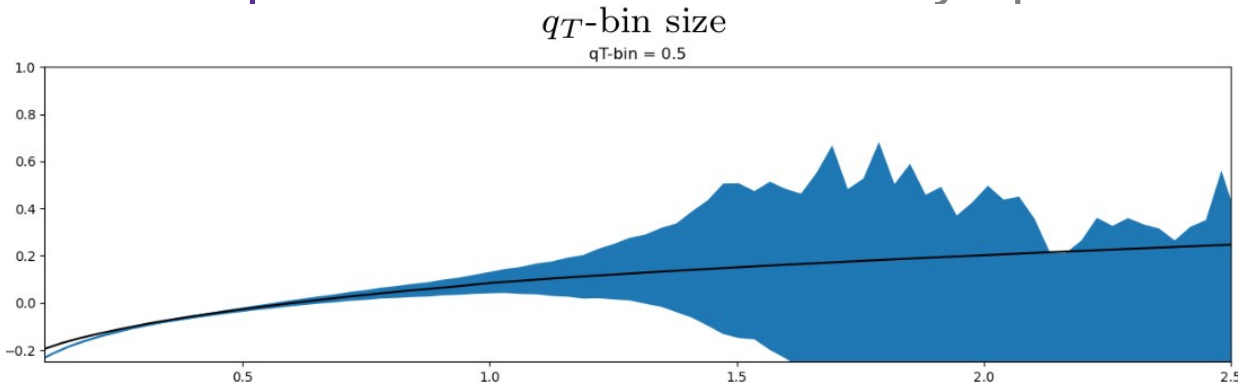
<https://www.desy.de/f/students/2022/reports/David.Gutierrez.pdf>

- ▶ CMS provides excellent muon capabilities
- ▶ High quality data at 7, 8, 13, 13.5 TeV
- ▶ Feasibility studies on the di-muon resolution show promising results:

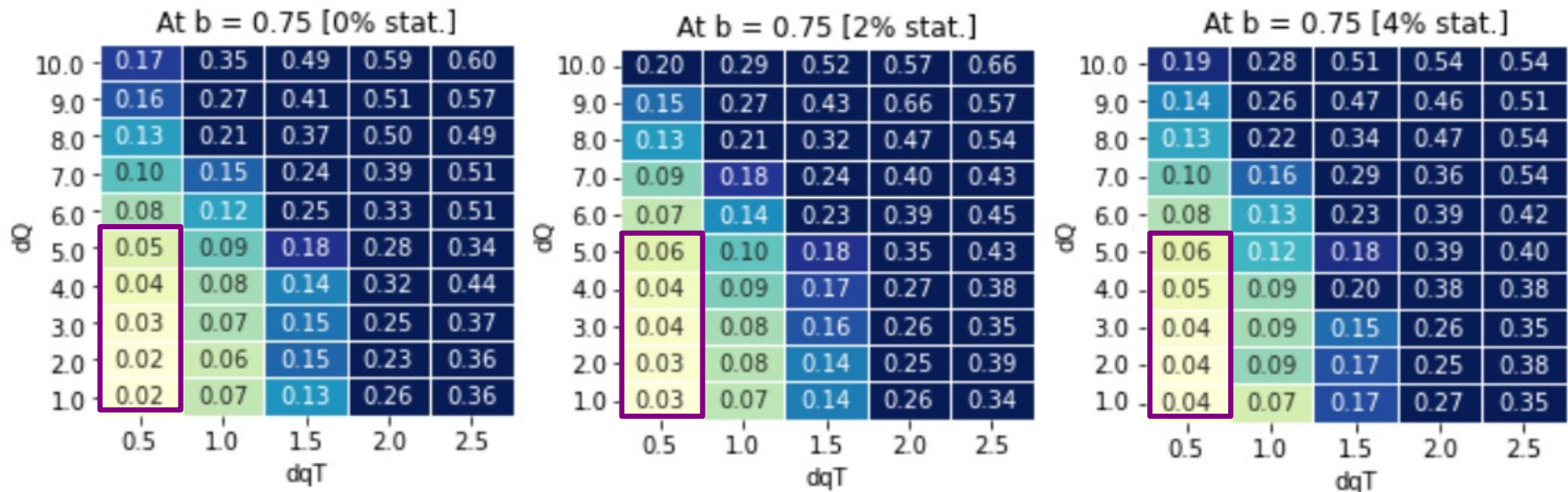


# Applying the method to **experimental data**

- ▶ As an example  $Q_1, Q_2 = 28, 46$  GeV
- ▶ **Small  $q_T$  bin size** ensure sensitivity up to around  $b = 1.5$



- ▶ Adding statistical and  $dQ$  uncertainties:



- ▶ **Statistical uncertainty mild**, main uncertainty from  $q_T$  binning

# Applying the method to transform PB TMDs to CSS

# Applying the method to transform PB TMDs to CSS

▶ This is a long standing problem

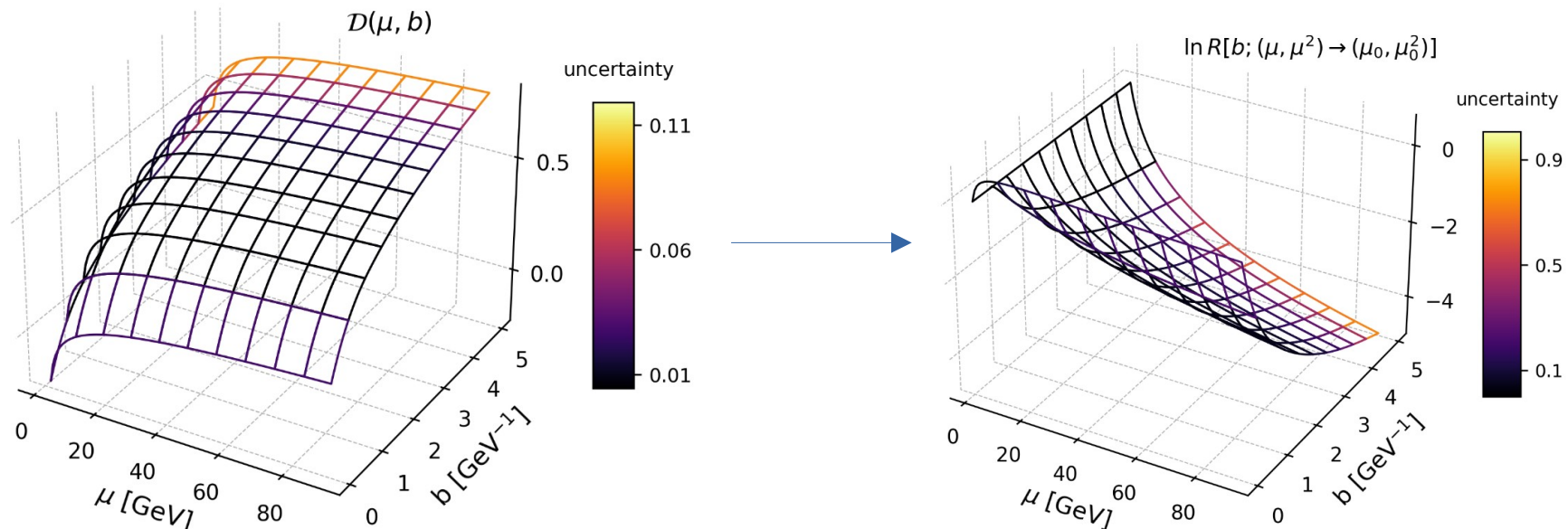
▶ Evolution of a TMD can be expressed as:

▶ Evolution factor

$$F(x, b; \mu, \zeta) = R[b; (\mu, \zeta) \rightarrow (\mu_0, \zeta_0)] F(x, b)$$

▶  $R[b; (\mu, \mu^2) \rightarrow (\mu_0, \mu_0^2)] = \exp \left\{ - \int_{\mu_0}^{\mu} \frac{d\mu'}{\mu'} (2\mathcal{D}(\mu', b) + \gamma_V(\mu')) \right\}$

▶ We use the method to determine RAD from DY in CASCADE:





# Applying the method to transform PB TMDs to CSS

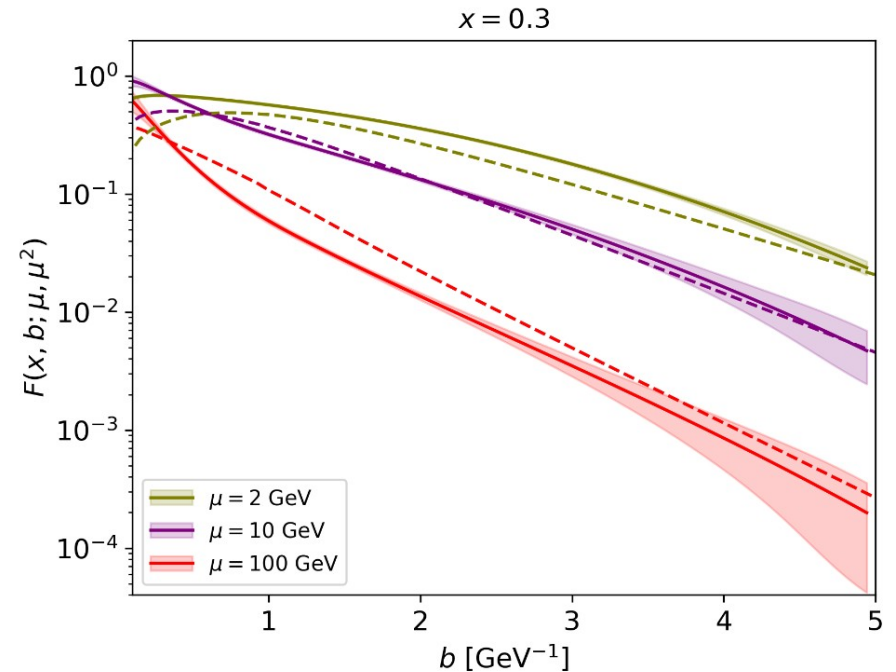
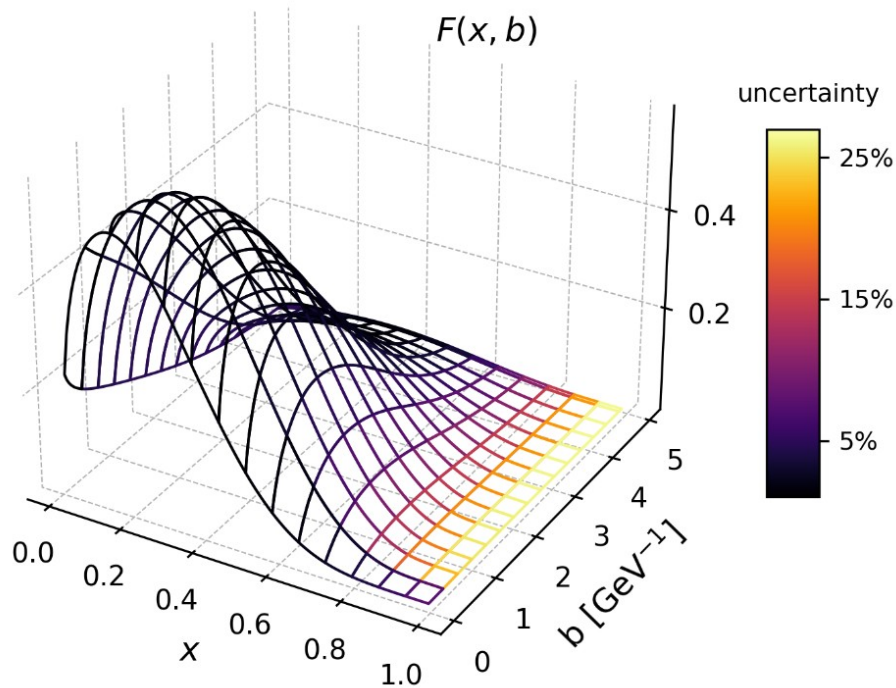
- ▶ This is a long standing problem

- ▶ Evolution of a TMD can be expressed as:

$$F(x, b; \mu, \zeta) = R[b; (\mu, \zeta) \rightarrow (\mu_0, \zeta_0)] F(x, b)$$

▶ Evolution factor

- ▶ Comparing PB TMD set2 to MAPP22



# Summary and conclusions

- ▶ Determination of RAD would be a **stringent test of factorization** and can have a deep impact on hadron 3D imaging
- ▶ Novel method to determine RAD was introduced
- ▶ Its application to simulated data from PB approach has **solved long standing problem of comparison** between factorization and PB
- ▶ **Feasibility studies** using CMS full simulated public data have shown promising results

