



p-Pb $\sqrt{s_{NN}}$ = 5.02 TeV

p+Pb, $\sqrt{s_{NN}} = 5.02 \text{ TeV}$, 28 nb⁻¹

→ v₂{2} template fit

-- v₂{2} peripheral subtraction

-- v₂{4} three-subevent method

 $\langle N_{ch} \rangle$

0.3<p_<3 GeV *ATLAS*

N_{ch} for 0.3<p₋<3 GeV

ALI-PREL-503267

Improved template f

Collective effects in PYTHIA 8 simulations of pp and p-Pb collisions

Executive Agency for Higher Education, Research, Development and Innovation Funding

C. Brandibur, A. Danu, A. Dobrin, A. Manea Institute of Space Science (RO)

Introduction

- Collective effects in small collision systems
 - Final state effects (Hydrodynamics, Parton escape)
 - Initial state effects (Color Glass Condensate Glasma, Color-field domains)

How to disentangle initial vs final state effects?

Scalar product (SP) method [4]

$$v_n\{SP\} = \frac{\langle \langle \mathbf{u}_{n,k} \mathbf{Q}_n^* / M \rangle \rangle}{\sqrt{\langle \mathbf{Q}_n^{*a} \mathbf{Q}_n^{*b} / (M^a M^b) \rangle}}$$

Particle of Interest $u_{n,x} = \cos(n \varphi)$ $u_{n,y} = \sin(n \varphi)$

Reference Particles $Q_{n,x} = \sum_{i} \cos(n \varphi_i)$ $Q_{n,y} = \sum_{i} \sin(n \, \phi_i)$



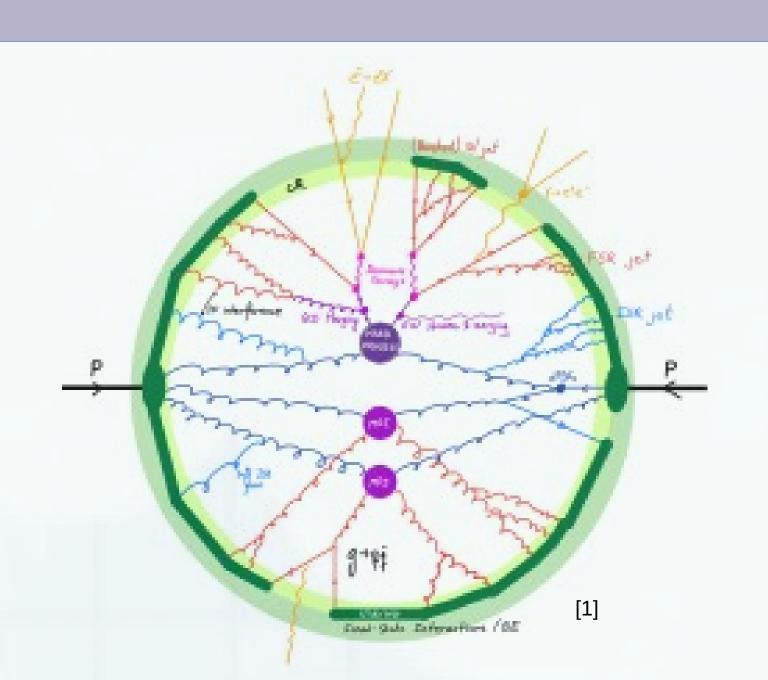
→ 2- and 4-particle azimuthal correlations for an event

$$\langle 2 \rangle \equiv \langle \cos(n(\varphi_i - \varphi_j)) \rangle, i \neq j$$
$$\langle 4 \rangle \equiv \langle \cos(n(\varphi_i + \varphi_j - \varphi_k - \varphi_l)) \rangle, i \neq j \neq k \neq l$$

→ Averaging over all events → 2nd and 4th order cumulants

$$c_n\{2\} = \langle \langle 2 \rangle \rangle = v_n^2$$

$$c_n\{4\} = \langle \langle 4 \rangle \rangle - 2 \langle \langle 2 \rangle \rangle^2 = -v_n^4$$

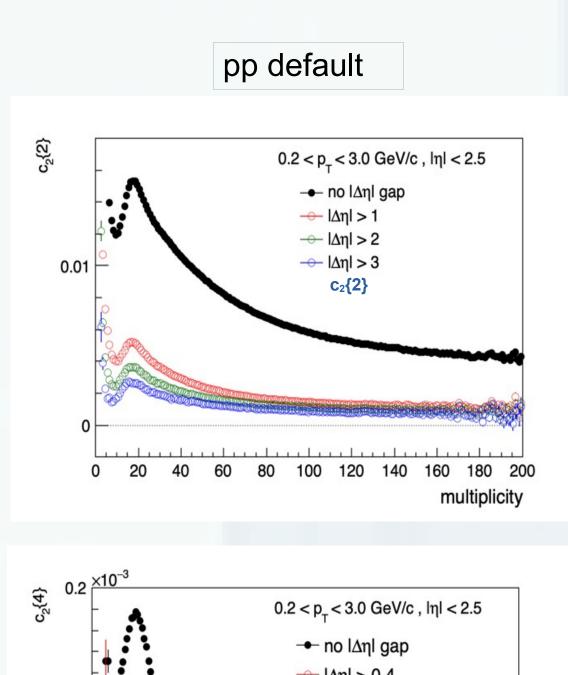


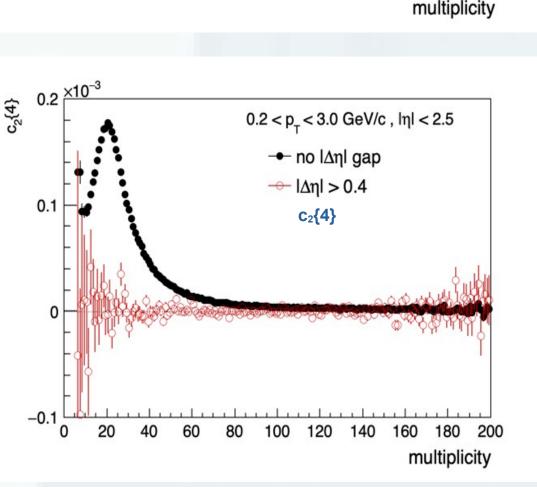
Microscopic model: PYTHIA8^[2]

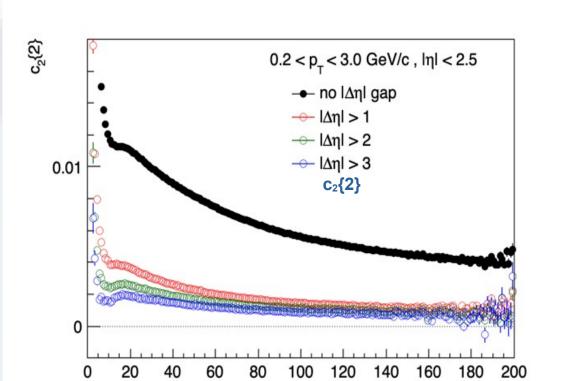
- QCD strings with LUND fragmentation
- Collective effects from new processes

· Color reconnection, rope hadronization, ...

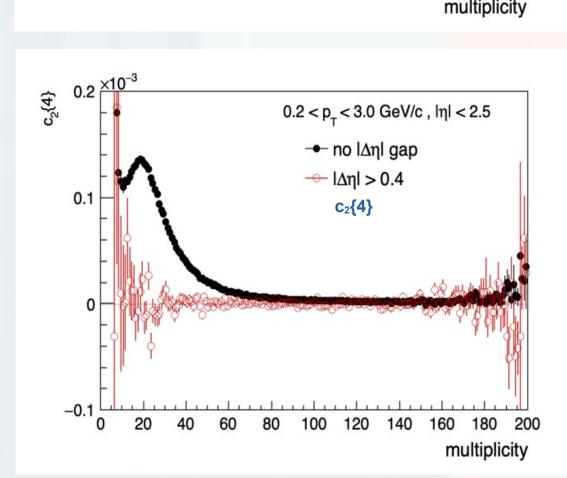
Results: pp @ 13.6 TeV and p-Pb @ 5.02





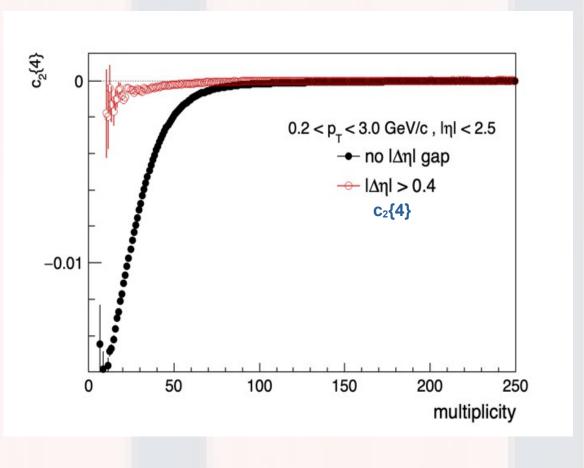


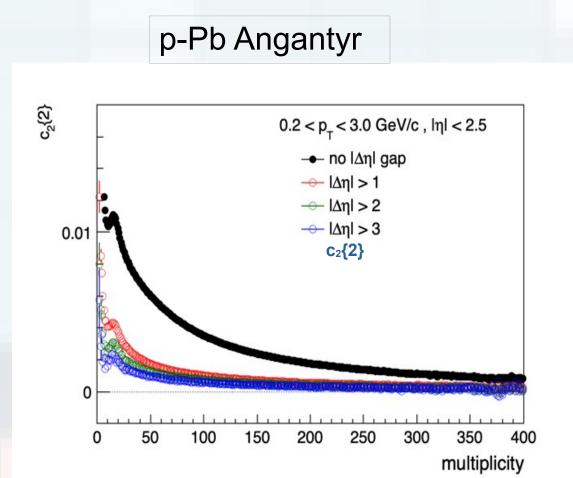
pp rope hadronization

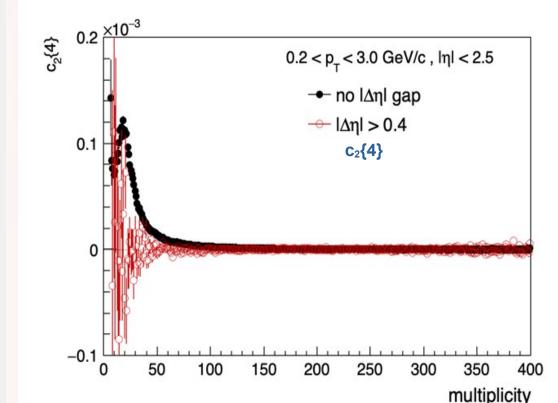


$0.2 < p_{_{\rm T}} < 3.0$ GeV/c , $|\eta| < 2.5$ - no l∆ηl gap - $|\Delta \eta| > 1$ \rightarrow $|\Delta \eta| > 2$ - |Δη| > 3 200

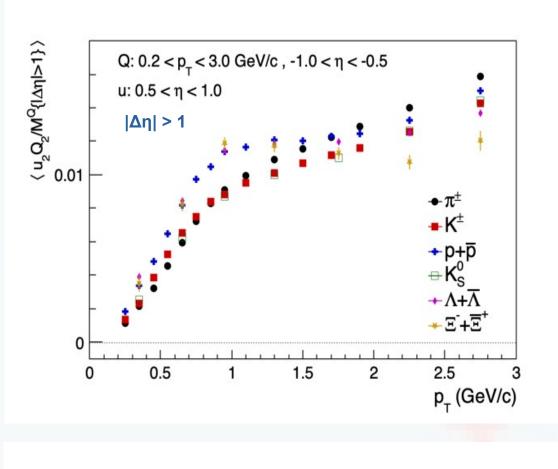
pp hardQCD

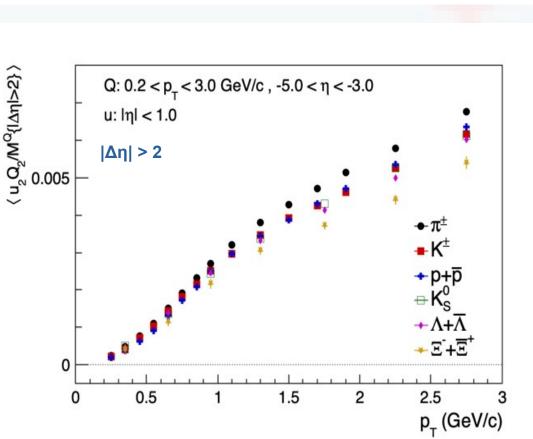




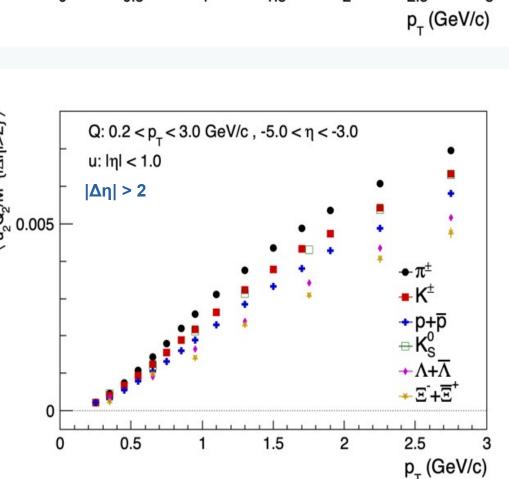


- $c_2\{2\} > 0$ and $c_2\{4\} \sim 0$ at high multiplicities
- → Small dependence on |∆η| gap for $c_2\{2\}$
- \rightarrow c₂{4} \sim 0 \rightarrow expected for Gaussian fluctuations
- Differences between pp default, hardQCD and rope hadronizaton
- Similar trends in pp and p-Pb default

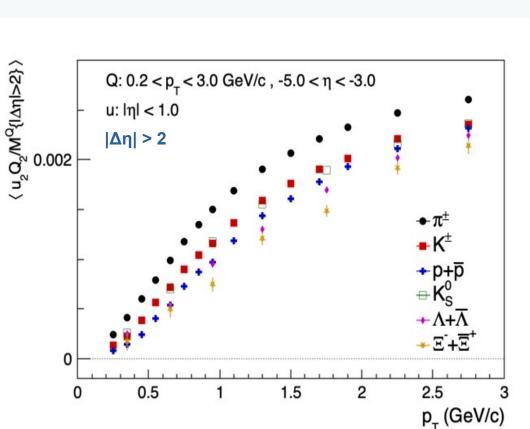




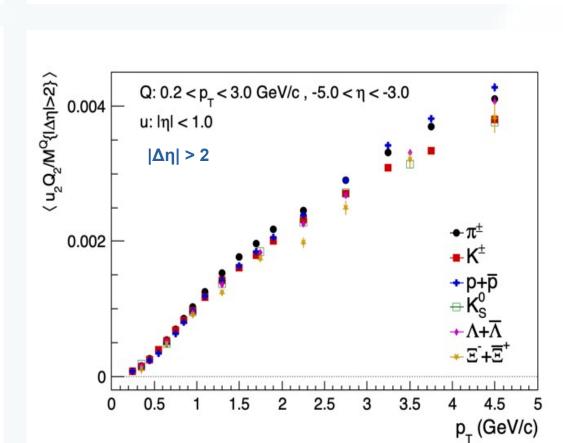
Q: $0.2 < p_{_T} < 3.0 \text{ GeV/c}$, $-1.0 < \eta < -0.5$ u: $0.5 < \eta < 1.0$ $|\Delta \eta| > 1$ o² 0.01 $+\Lambda + \overline{\Lambda}$ p_{_} (GeV/c)



Q: $0.2 < p_T < 3.0 \text{ GeV/c}$, $-1.0 < \eta < -0.5$ 0.04 $u: 0.5 < \eta < 1.0$ $|\Delta n| > 1$ p_T (GeV/c)



Q: $0.2 < p_{_T} < 3.0 \text{ GeV/c}$, $-1.0 < \eta < -0.5$ $u: 0.5 < \eta < 1.0$ 0.5 1 1.5 2 2.5 3 3.5 4 4.5 5 p_T (GeV/c)



- Differences between pp default, hardQCD and rope hadronizaton
- Similar trends in pp and p-Pb default
- Small mass ordering for large |Δη| gap
- → More pronounced for rope hadronization and hardQCD
- Crossing between proton and pion u_2Q_2 for large $|\Delta\eta|$ gap in p-Pb
- → No particle type grouping

Summary

- $c_2\{2\}$ decreasing with increasing multiplicity and $|\Delta\eta|$ gap
- → Small dependence on |Δη| gap
- $c_2{4} \sim 0$ at high multiplicities
- → Expected for Gaussian fluctuations
- Mass ordering for u_2Q_2 when a large $|\Delta\eta|$ gap is employed
- → Crossing between pions and protons u₂Q₂ in PYTHIA 8 Angantyr p-Pb simulations
- → No particle type grouping

References

[1] figure by P. Skands

[2] C. Bierlich et al., arXiv: 2203.11601

[3] ATLAS, PRC 97 (2018) 024904

[4] A. Bilandzic, PRC 83 (2011) 044913 [5] S. Voloshin et al., arXiv:0809.2949