

# Lepton and Neutron EDM as Probe of General Two Higgs Doublet Model

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# Outline

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# Motivation

- Baryogenesis: CP violation (CPV) beyond Kobayashi-Masakawa
- Electric dipole moments (EDMs): Precision measurements that put constraints on CPV

# g2HDM

The *general* two Higgs doublet model (g2HDM)

- Add a second Higgs doublet, natural extension of SM
- No imposed  $\mathbb{Z}_2$  symmetry constraint (hence *general*)
- Flavor-changing neutral Higgs couplings controlled by flavor hierarchies and alignment ( $c_\gamma \equiv \cos \gamma = \cos(\beta - \alpha)$  is small)

Lagrangian:

$$\begin{aligned}\mathcal{L} = & -\frac{1}{\sqrt{2}} \sum_{f=u,d,\ell} \bar{f}_i \left[ \left( -\lambda_i^f \delta_{ij} s_\gamma + \rho_{ij}^f c_\gamma \right) h \right. \\ & + \left( \lambda_i^f \delta_{ij} c_\gamma + \rho_{ij}^f s_\gamma \right) H - i \operatorname{sgn}(Q_f) \rho_{ij}^f A \Big] R f_j \\ & - \bar{u}_i \left[ (V \rho^d)_{ij} R - (\rho^{u\dagger} V)_{ij} L \right] d_j H^+ \\ & - \bar{\nu}_i \rho_{ij}^L R \ell_j H^+ + \text{h.c.},\end{aligned}$$

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Takeaway

$c_\gamma$  small (alignment)

- One  $\rho$  matrix per fermion family ( $u, d, \ell$ )
- SM Yukawas control SM-like Higgs,  $\rho$  controls exotic Higgses

Important parameters

- $\mathcal{O}(1)$   $\rho_{tt}$  can drive baryogenesis through<sup>1</sup>  $\lambda_t \operatorname{Im} \rho_{tt}$
- $\rho_{tt}$  along with  $\rho_{ff}$  drive EDM for fermion  $f$

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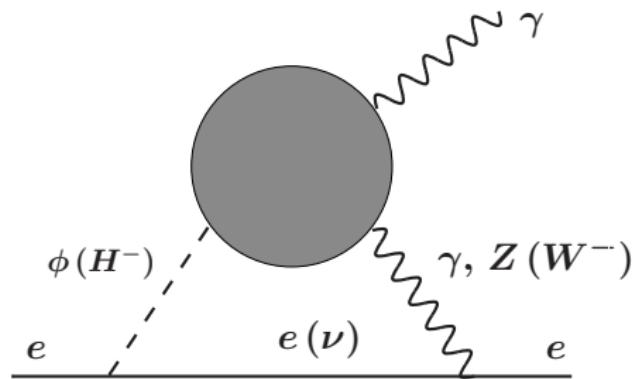
<sup>1</sup>K. Fuyuto, W.-S. Hou, E. Senaha, PLB **776**(2018)402, [1705.05034]

# g2HDM and EDMs

Interaction term for EDMs:

$$-\frac{i}{2} d_f (\bar{f} \sigma^{\mu\nu} \gamma_5 f) F_{\mu\nu}$$

Main drive for EDMs in g2HDM is two-loop Barr-Zee diagram<sup>2</sup>.



<sup>2</sup>S.M. Barr and A. Zee, PRL**65**(1990)21

# Electron EDM

Current bound for electron EDM (eEDM)

- ACME (2018)<sup>3</sup>:  $|d_e| < 1.1 \times 10^{-29} \text{ e cm}$
- JILA (2023)<sup>4</sup>:  $|d_e| < 4.1 \times 10^{-30} \text{ e cm}$

Very high precision from “tabletop experiments” !

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<sup>3</sup>V. Andreev et al. [ACME], Nature**562**(2018)355

<sup>4</sup>T.S. Roussy et al., Science**381**(2023)46

# Electron EDM

Large  $\rho_{tt}$  needed for baryogenesis, but would also drive large EDMs.

A previous study by Fuyuto, Senaha, and Hou<sup>5</sup> proposed a “cancellation ansatz” between  $\rho_{ee}$  and  $\rho_{tt}$  to evade eEDM bounds.

$$\text{Re } \rho_{ee} = -r \frac{\lambda_e}{\lambda_t} \text{Re } \rho_{tt}, \quad \text{Im } \rho_{ee} = +r \frac{\lambda_e}{\lambda_t} \text{Im } \rho_{tt} \quad (1)$$

which gives both a flavor hierarchy  $|\rho_{ee}|/|\rho_{tt}| \sim \lambda_e/\lambda_t$  that reflects SM, as well as a phase lock.  $r$  depends on loop functions.

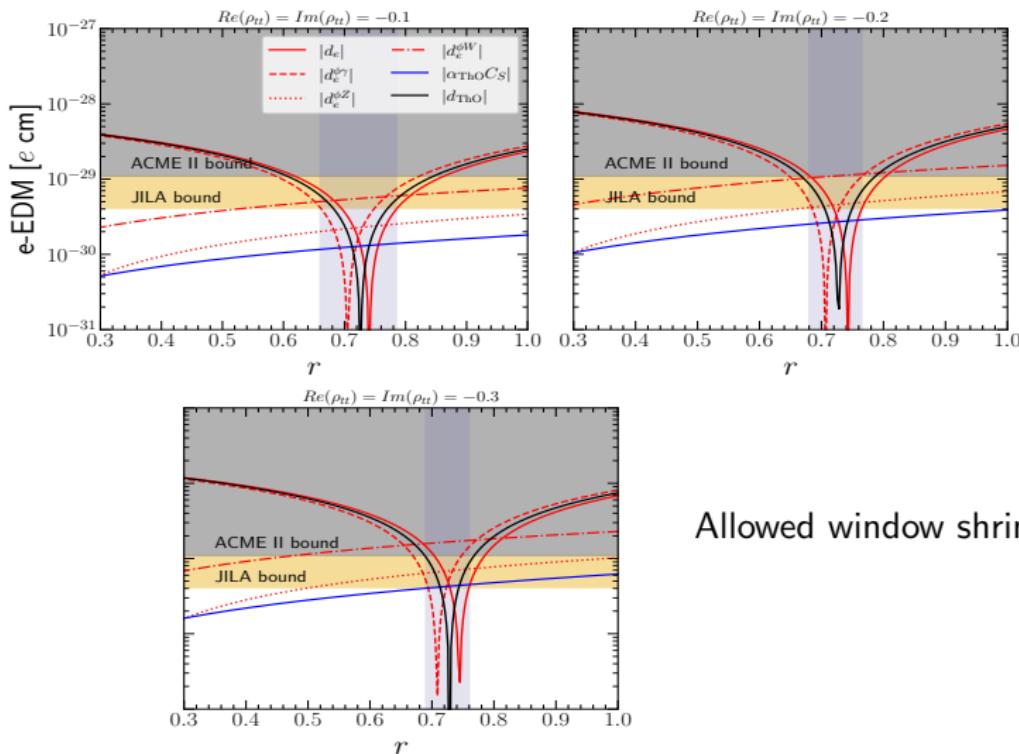
For sake of numerical illustration of the flavor hierarchy, we extend this ansatz to all  $\rho_{ff}$  except top.

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<sup>5</sup>K. Fuyuto, W.-S. Hou, E. Senaha, PRD**101**(2020)011901(R), [1910.12404] 

# Electron EDM

We explore larger range of  $|\rho_{tt}|$ . ( $c_\gamma = 0.1$ ,  $m_{H,A,H^\pm} = 500 \text{ GeV}$ )



Allowed window shrinks.

# Echo in Neutron EDM?

Current bound for Neutron EDM (nEDM)

- PSI (2020)<sup>6</sup>:  $|d_n| < 1.8 \times 10^{-26} \text{ e cm}$

Precision poorer than eEDM.

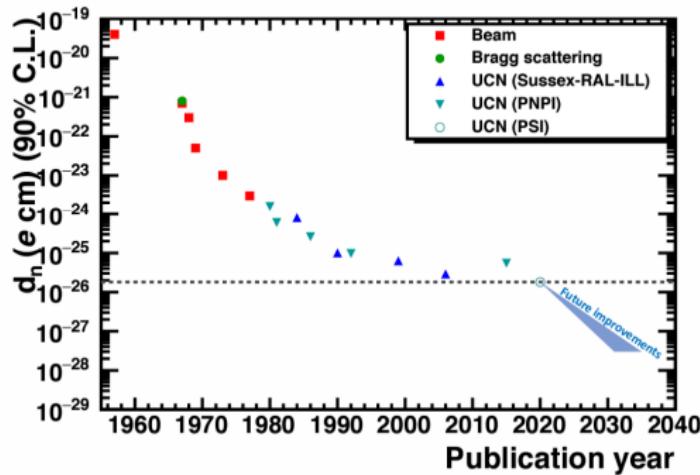


Figure from Snowmass report<sup>7</sup>

<sup>6</sup>C. Abel et al. [nEDM], PRL **124**(2020)081803, [2001.11966]

<sup>7</sup>R. Alarcon et al. [2203.08103]

# Neutron EDM

Two additional contributions:

Chromo-EDM of fermions

$$-\frac{ig_s}{2} \tilde{d}_f (\bar{f} \sigma^{\mu\nu} T^a \gamma_5 f) G_{\mu\nu}^a$$

and

Weinberg term for gluons

$$-\frac{1}{3} C_W f^{abc} G_{\mu\sigma}^a G_{\nu}^{b,\sigma} \tilde{G}^{c,\mu\nu}$$

We use the following formula<sup>8</sup>

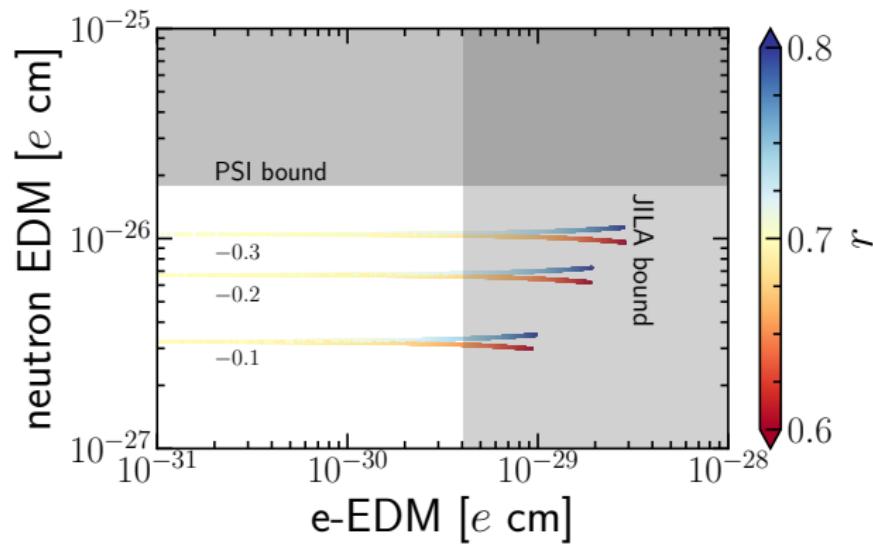
$$d_n = -0.20 d_u + 0.78 d_d + e (0.29 \tilde{d}_u + 0.59 \tilde{d}_d) + e 23 \text{ MeV } C_W$$

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<sup>8</sup> Hisano, Kobayashi, Kuramoto, Kuwahara, JHEP11(2015)085, [1507.05836] 

# Neutron EDM, even $|\rho_{tt}| \sim 0.42$ not ruled out!

We show combined results for eEDM and nEDM.



Note:  $r \approx 0.7$  is a good loop function value.

# Neutron EDM

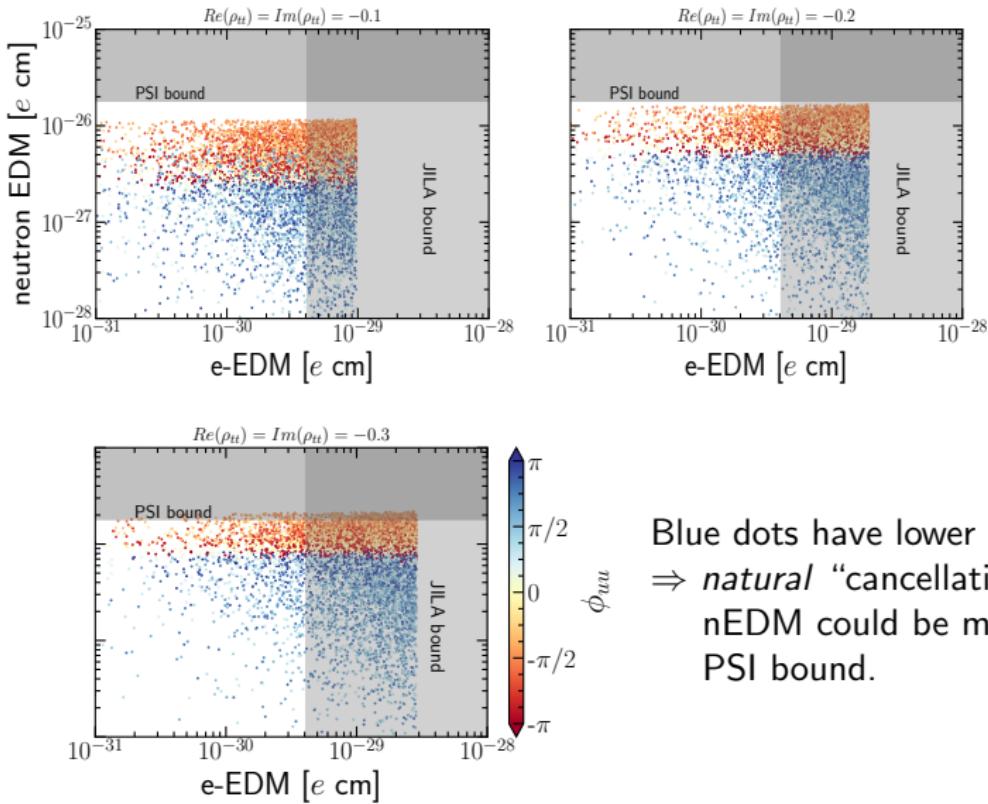
Since the ansatz is just numerical illustration of flavor hierarchy, we further explore<sup>9</sup>  $\rho_{uu} \sim \mathcal{O}(\lambda_u)$  by varying

$$|\rho_{uu}| \in [0.3\lambda_u, 3\lambda_u], \quad \arg \rho_{uu} \in [-\pi, \pi]$$

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<sup>9</sup>W.-S. Hou, G. Kumar, PRD**102**(2020)115017, [2008.08469]

# Neutron EDM, with $|\rho_{uu}| \in [0.3\lambda_u, 3\lambda_u]$ , $\arg \rho_{uu} \in [-\pi, \pi]$



Blue dots have lower nEDM  
⇒ *natural* “cancellation mechanism”  
nEDM could be much below  
PSI bound.

# Discussion and Summary

- Future sensitivity:
  - n2EDM at PSI ( $\sim 10^{-27}$ )
  - Spallation Neutron Source at ORNL ( $\sim 10^{-28}$ )can probe most of our projected nEDM range.
- Prospect of Discovery:  
If g2HDM is indeed behind electroweak baryogenesis, then combined eEDM-nEDM are poised for **Discovery**.
- Note on Heavy Higgs masses: we have taken  $H, A, H^\pm$  degenerate at 500 GeV; if nondegenerate, need to face electroweak precision constraints.