

Transverse Momentum Dependent Fragmentation Functions from recent BELLE data

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In collaboration with
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UNIVERSITÀ
DEGLI STUDI
DI TORINO



Based on :

M. Boglione, J. O. Gonzalez-Hernandez, and A. Simonelli (2022), 2206.08876.

Outlook

- Theoretical framework for $e^+e^- \rightarrow h X$
- Global fits
- Phenomenological analysis of recent BELLE data

Theoretical framework for $e^+e^- \rightarrow h X$

M. Boglione, A. Simonelli JHEP 02 (2021) 076

$$\begin{aligned}
 \frac{d\sigma^{\text{NLO, NLL}}}{dz_h dT dP_T^2} &= \\
 &= -\sigma_B \pi N_C \frac{\alpha_S(Q)}{4\pi} C_F \frac{3 + 8 \log(1 - T)}{1 - T} \exp \left\{ -\frac{\alpha_S(Q)}{4\pi} 3C_F (\log(1 - T))^2 \right\} \times \\
 &\times \sum_f e_f^2 \int \frac{d^2 \vec{b}_T}{(2\pi)^2} e^{i \frac{\vec{P}_T}{z_h} \cdot \vec{b}_T} \tilde{D}_{1, H/f}^{\text{NLL}}(z_h, b_T, Q, (1 - T) Q^2) \left[1 + \mathcal{O} \left(\frac{M_H^2}{Q^2} \right) \right] \\
 \\
 \tilde{D}_{1, H/f}(z, b_T; \mu, \zeta) &= \underbrace{\frac{1}{z^2} \sum_k \int_z^1 \frac{d\rho}{\rho} d_{H/k}(z/\rho, \mu_b) [\rho^2 \mathcal{C}_{k/f}(\rho, \alpha_S(\mu_b))]}_{\text{TMD at reference scale}} \times \\
 &\times \underbrace{\exp \left\{ \frac{1}{4} \tilde{K}(b_T^*; \mu_b) \log \frac{\zeta}{\mu_b^2} + \int_{\mu_b}^{\mu} \frac{d\mu'}{\mu'} \left[\gamma_D(\alpha_S(\mu'), 1) - \frac{1}{4} \gamma_K(\alpha_S(\mu')) \log \frac{\zeta}{\mu'^2} \right] \right\}}_{\text{Perturbative Sudakov Factor}} \\
 &\times \underbrace{(M_D)_{j, H}(z, b_T) \exp \left\{ -\frac{1}{4} g_K(b_T) \log \frac{z_h^2 \zeta}{M_H^2} \right\}}_{\text{Non-Perturbative content}}.
 \end{aligned}$$

Theoretical framework for $e^+e^- \rightarrow h X$

M. Boglione, A. Simonelli, Eur. Phys. J. C 81 (2021)

$$D = D^* \sqrt{M_S}$$



Usual definition
TMD FF

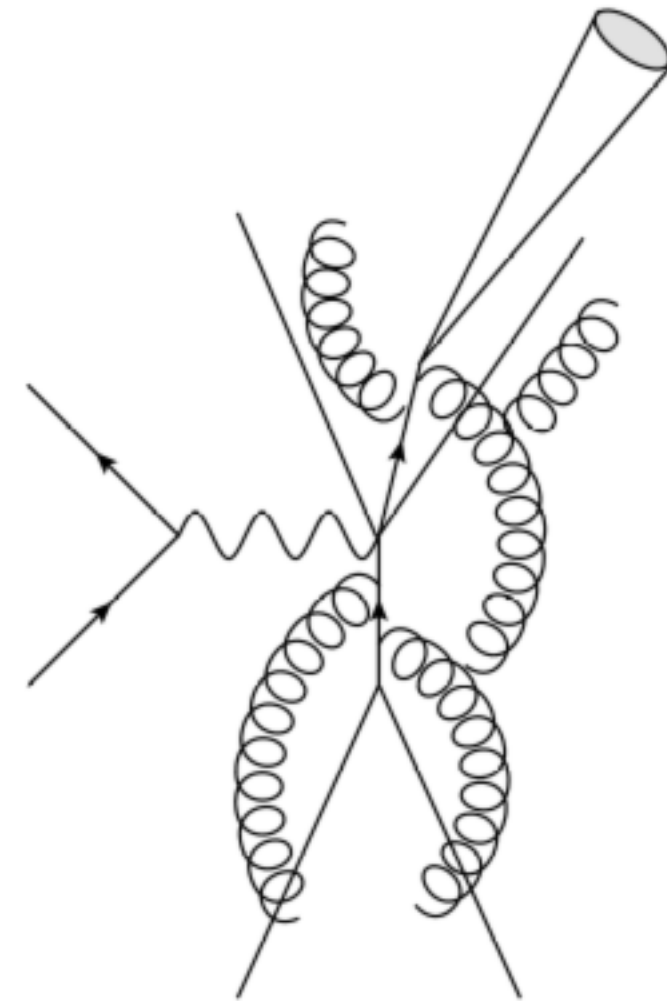


For this process
TMD FF



Soft non-
perturbative
Function

$$e^+e^- \rightarrow hX$$



$$\frac{d\sigma}{dP_T} = d\hat{\sigma} \otimes D^*(P_T)$$

Theoretical framework for $e^+e^- \rightarrow h X$

M. Boglione, A. Simonelli, Eur. Phys. J. C 81 (2021)

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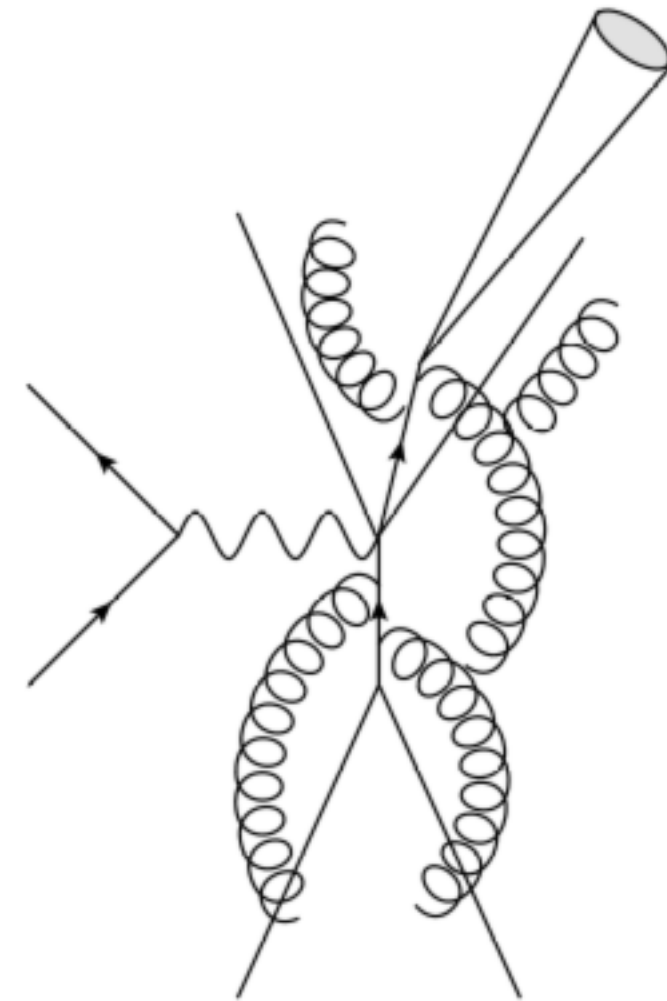


Same
constraints
to collinear FF

$$g_K(b_T) = \tilde{K}(b_T^*; \mu) - \tilde{K}(b_T; \mu)$$

Same function for non-perturbative evolution

$$e^+e^- \rightarrow hX$$



$$\frac{d\sigma}{dP_T} = d\hat{\sigma} \otimes D^*(P_T)$$

Theoretical framework for $e^+e^- \rightarrow h X$

M. Boglione, A. Simonelli, Eur. Phys. J. C 81 (2021)

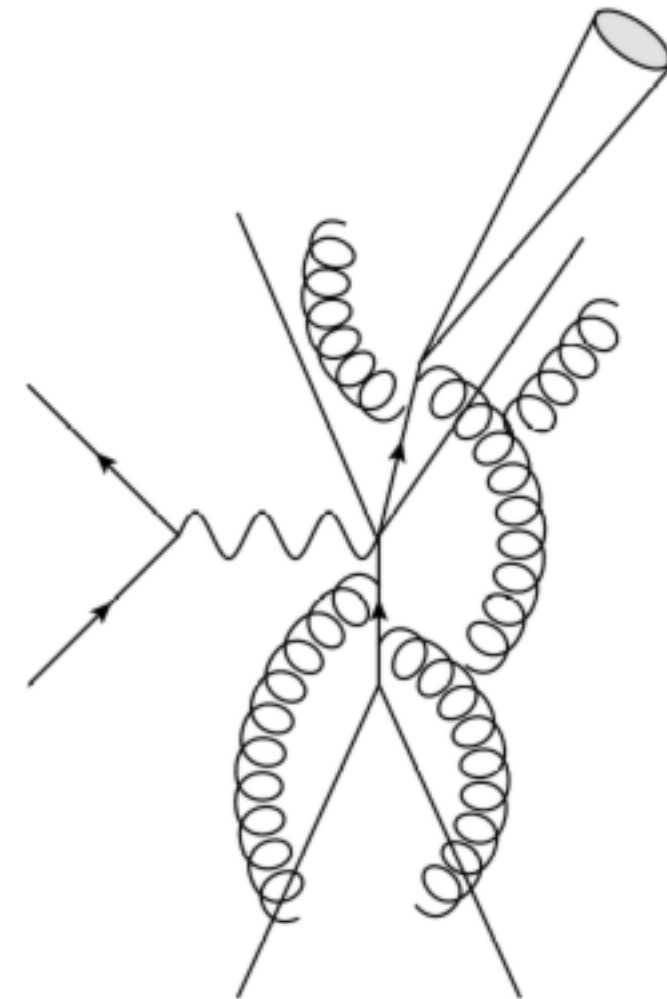
$$D = D^* \sqrt{M_S}$$



What is the effect of the collinear FFs (and PDFs in general) ?

Large- b_T behaviour of g_K ?

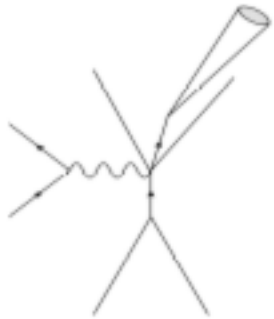
$$e^+e^- \rightarrow hX$$



$$\frac{d\sigma}{dP_T} = d\hat{\sigma} \otimes D^*(P_T)$$

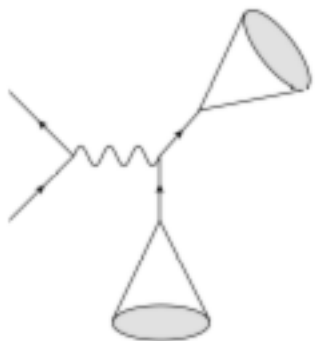
Global Fits

Possible roadmap



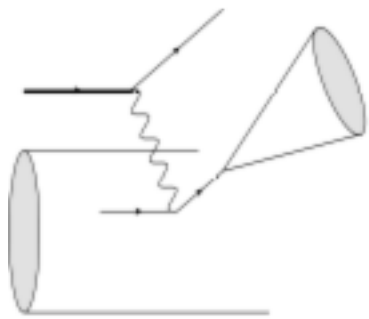
1.

Extraction of the unpolarized TMD FF, D^* , for charged pions from BELLE data (using factorization definition)



2.

Two non-perturbative functions:
 D^* , known from step 1
Soft Model M_S



3. SIDIS

Three non-perturbative functions in the cross section
 D^* , known from step 1.
Soft Model M_S , known from step 2.
Extraction of the TMD PDF, F^* (in the factorization definition, $F^* \neq F$).

(Global) Fits

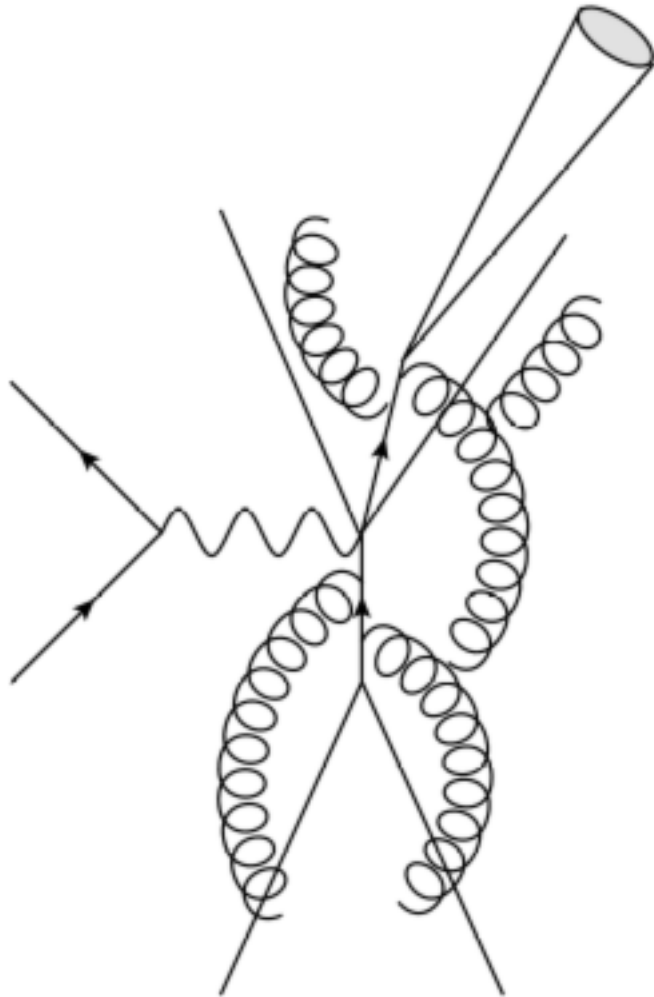
Must consider:

- Which collinear functions are more appropriate?
- Which regions in bT are being mapped by extractions.
- Constraints of bT -behaviour for TMDs.
- Physical pictures/theoretical arguments /models (not parametrizations)
- Non perturbative evolution (gK) should be consistent with SIDIS, DY, $e+e^-$ two-hadron production.

Phenomenological analysis of recent BELLE data

BELLE data overview

$$e^+e^- \rightarrow hX \quad (\text{Charged pions})$$



Binned in PT, zh and T (thrust)

$$Q=10.6 \text{ GeV}$$

$$0.06 < PT < 2.5 \text{ GeV}$$

$$0.125 < zh < 0.975 \text{ (18 bins)}$$

$$0.6 < T < 0.975. \text{ (6 bins)}$$

$$PT/zh < 0.15 Q$$

For our analysis

$$0.375 < zh < 0.725. \text{ (8 bins)}$$

$$0.750 < T < 0.875. \text{ (3 bins)}$$

zh-dependence and choice of collinear FFs

- We compare results obtained with NNFFnlo and JAM20nlo

$$\begin{aligned}
 \tilde{D}_{1,H/f}(z, b_T; \mu, \zeta) = & \underbrace{\frac{1}{z^2} \sum_k \int_z^1 \frac{d\rho}{\rho} d_{H/k}(z/\rho, \mu_b) [\rho^2 \mathcal{C}_{k/f}(\rho, \alpha_S(\mu_b))]}_{\text{TMD at reference scale}} \times \\
 & \underbrace{\exp \left\{ \frac{1}{4} \tilde{K}(b_T^*; \mu_b) \log \frac{\zeta}{\mu_b^2} + \int_{\mu_b}^{\mu} \frac{d\mu'}{\mu'} \left[\gamma_D(\alpha_S(\mu'), 1) - \frac{1}{4} \gamma_K(\alpha_S(\mu')) \log \frac{\zeta}{\mu'^2} \right] \right\}}_{\text{Perturbative Sudakov Factor}} \\
 & \underbrace{\times (M_D)_{j,H}(z, b_T) \exp \left\{ -\frac{1}{4} g_K(b_T) \log \frac{z_h^2 \zeta}{M_H^2} \right\}}_{\text{Non-Perturbative content}}.
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zh-dependence and choice of collinear FFs

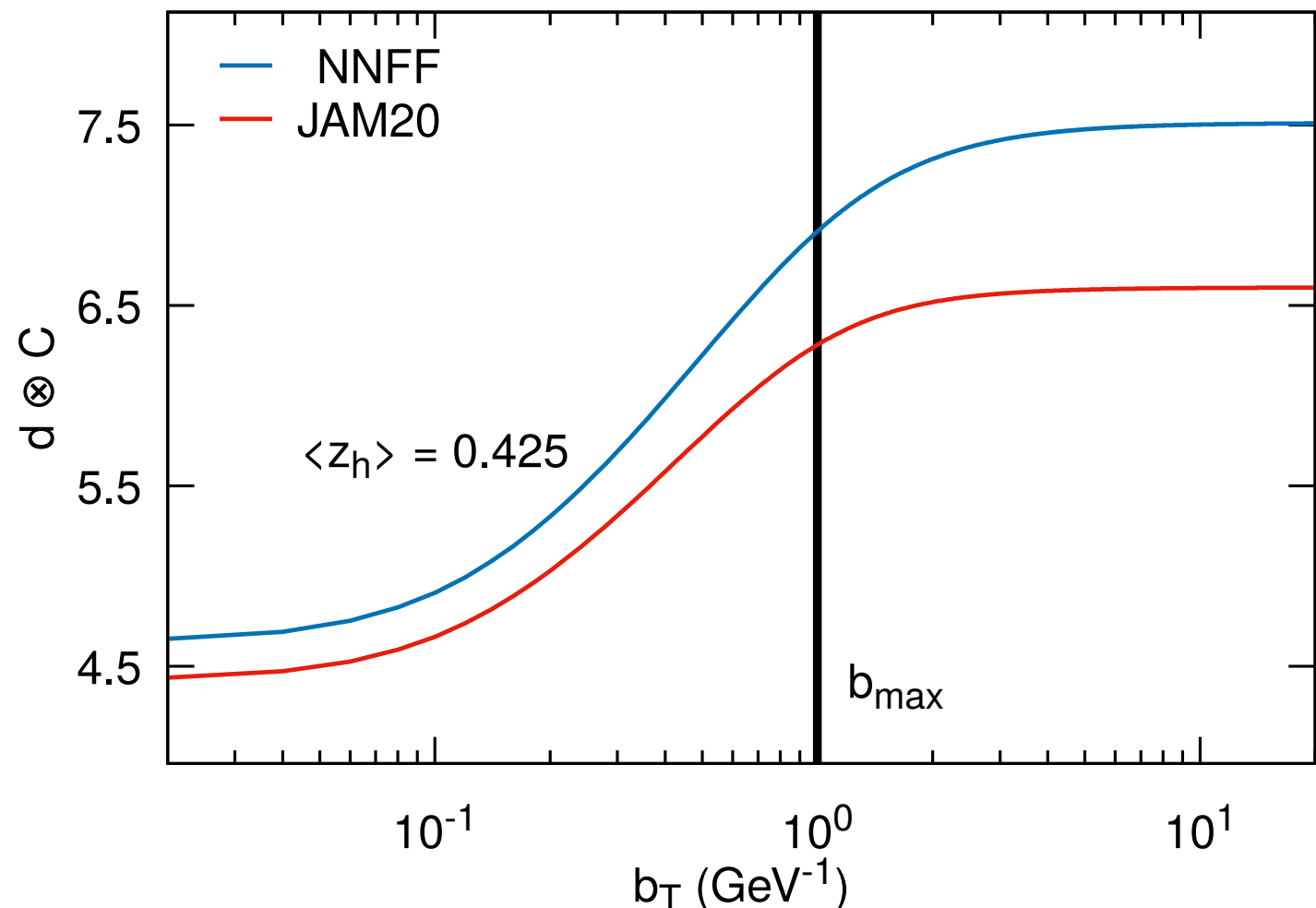
- We compare results obtained with NNFFnlo and JAM20nlo

$$\frac{1}{z^2} \sum_k \int_z^1 \frac{d\rho}{\rho} d_{H/k}(z/\rho, \mu_b) [\rho^2 \mathcal{C}_{k/f}(\rho, \alpha_S(\mu_b))]$$

$$\mu_b = \frac{2e^{-\gamma_E}}{b_T^*(b_T)}$$

$$\vec{b}_T^*(b_T) = \frac{\vec{b}_T}{\sqrt{1 + b_T^2/b_{\max}^2}}$$

$$\vec{b}_T^*(b_c(b_T)) = \vec{b}_T^* \left(\sqrt{b_T^2 + b_{\min}^2} \right)$$



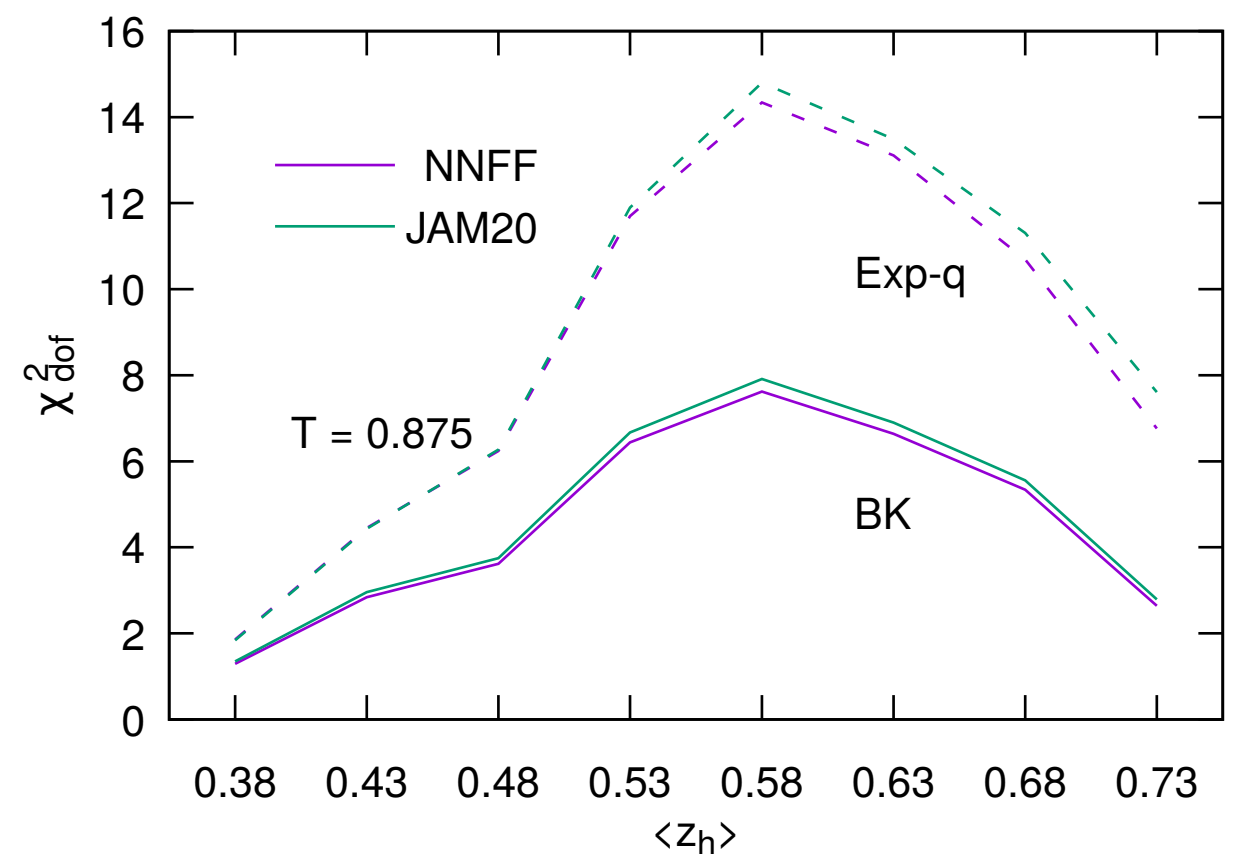
zh-dependence and choice of collinear FFs

- We compare results obtained with NNFFnlo and JAM20nlo

Nomenclature	M_D -model	parameters
z_h -independent models		
1) Exponential-q	$e^{-(M_0 b_T)^q}$	M_0, q
2) Bessel-K	$\frac{2^{2-p} (b_T M_0)^{p-1}}{\Gamma(p-1)} K_{p-1}(b_T M_0)$	M_0, p

**Proxy models: performed fits
at fixed $T=0.875$.**

**One INDEPENDENT fit for each
zh-bin in the range
 $0.375 < z_h < 0.725$.**

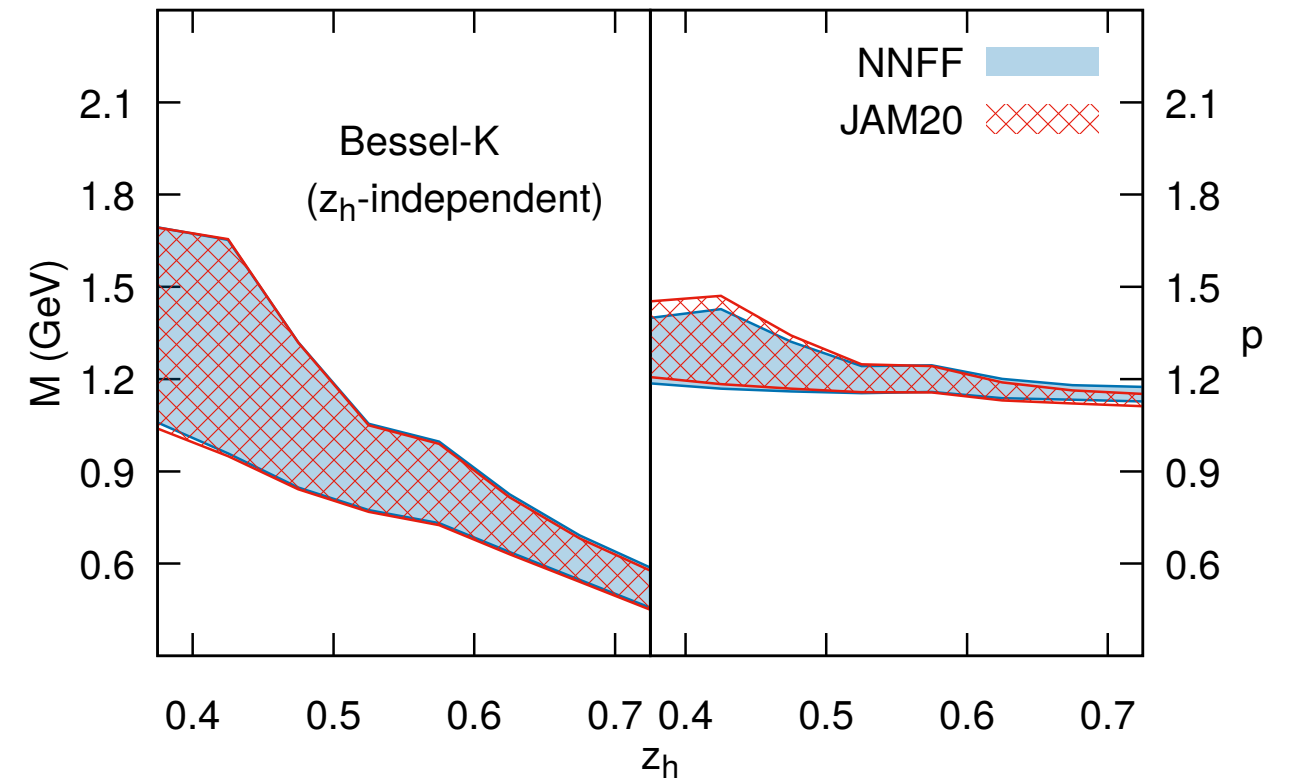


**JAM20 & NNFF perform
equally.**

zh-dependence and choice of collinear FFs

Nomenclature	M_D -model	parameters
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**Proxy models: performed fits
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One INDEPENDENT fit for each
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**Stronger z_h -dependence in *dimensionfull*
parameter**

Model for PT-dependence

Example: preliminary fits at fixed T

$\chi^2_{\text{d.o.f.}}$ (fixed- T fits)			
M_{D} model	$T = 0.750$	0.825	0.875
I	1.2	0.38	1.02
IG	1.46	0.47	1.51

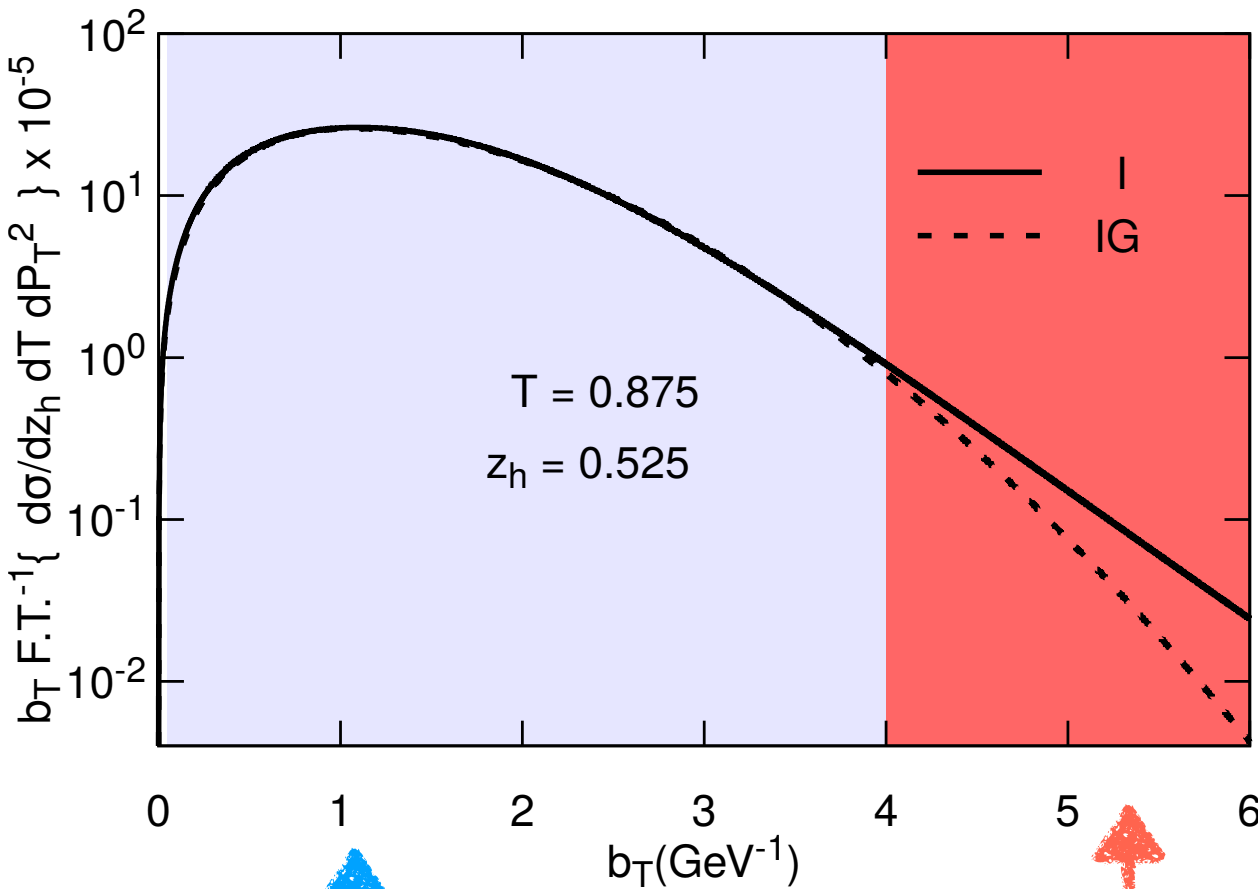
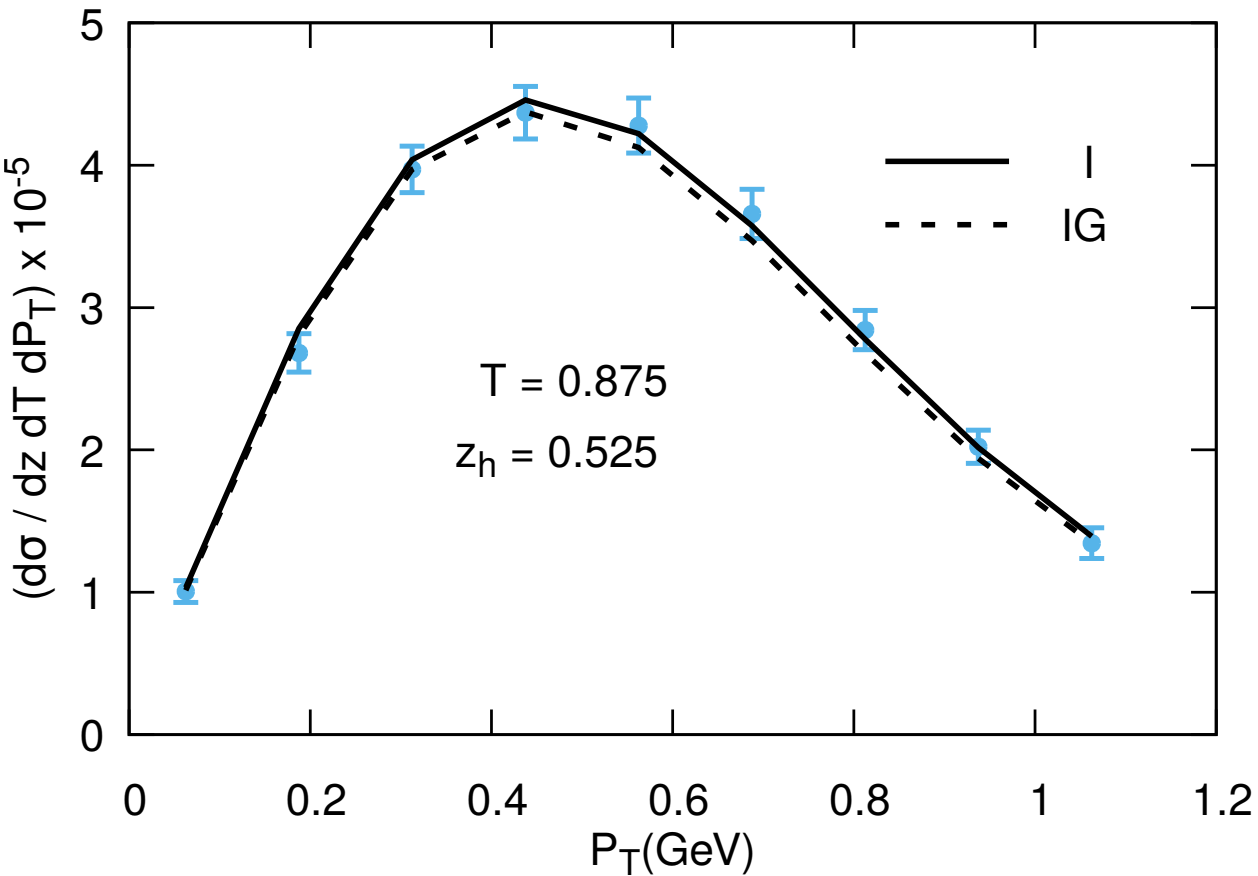
$M_{\text{D}} = \frac{2^{2-p} (b_{\text{T}} M_0)^{p-1}}{\Gamma(p-1)} K_{p-1}(b_{\text{T}} M_0) \times F(b_{\text{T}}, z_h)$ $M_z = -M_1 \log(z_h)$		
ID	F -model	parameters
I	$F = \left(\frac{1 + \log(1 + (b_{\text{T}} M_z)^2)}{1 + ((b_{\text{T}} M_z)^2)} \right)^q$	$M_0, M_1, p, q = 8$
IG	$F = \exp((M_g b_{\text{T}})^2 \log(z_h))$	M_0, M_g, p

zh-dependence
Factor

Model for PT-dependence

Very large bT behaviour
might not be resolved by data

$\chi^2_{\text{d.o.f.}}$ (fixed- T fits)			
M_D model	$T = 0.750$	0.825	0.875
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MD $\rightarrow$ 1 for
 $b_T \rightarrow 0$

Different asymptotic
behaviour

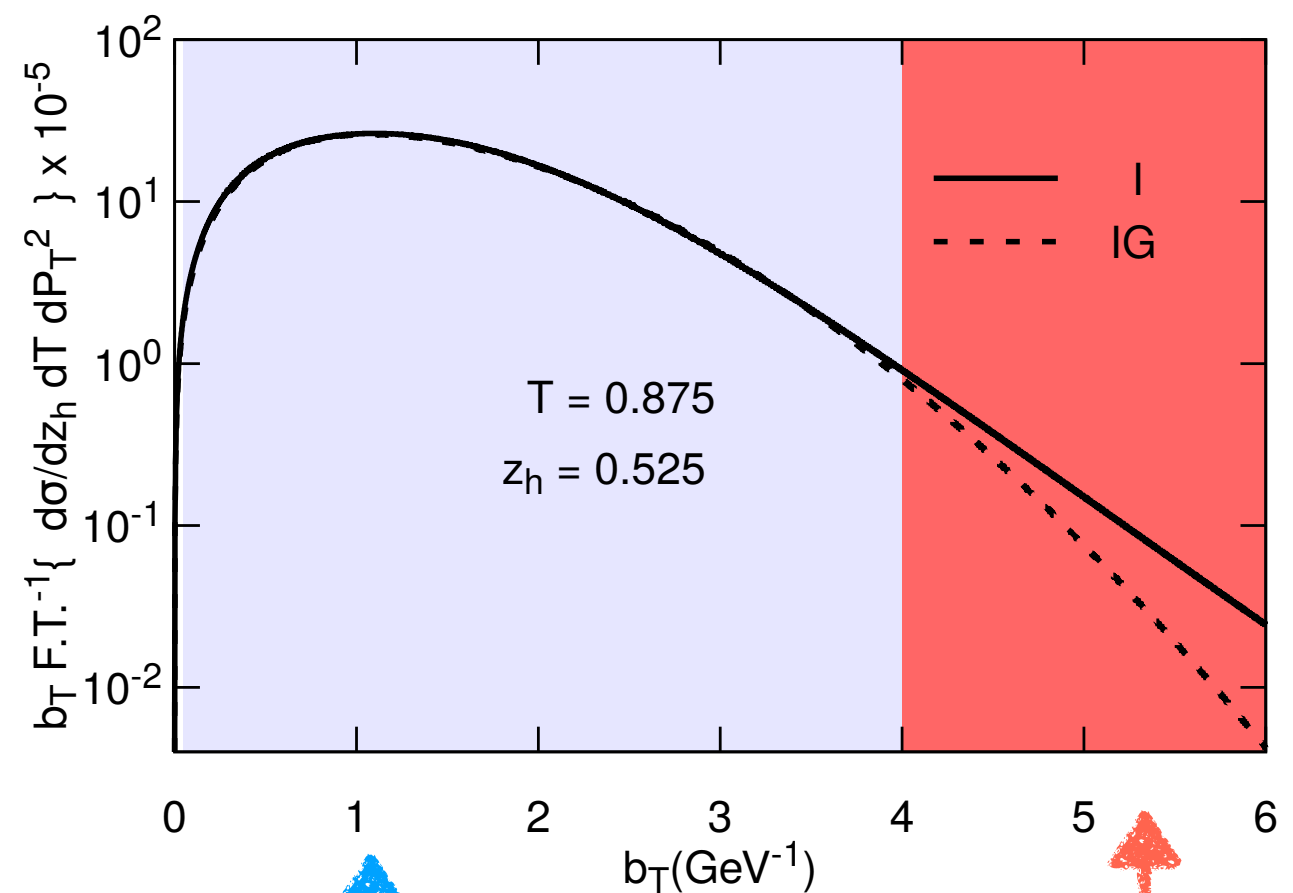
Model for PT-dependence

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Very large b_T behaviour
might not be resolved by data

At smaller energies
(say, COMPASS)
one has sensitivity to
larger values of b_T .

Need to constraint
very large b_T region.



MD $\rightarrow 1$ for
 $b_T \rightarrow 0$

Different asymptotic
behaviour

Model for PT-dependence

Which asymptotic behaviour?

P. Schweitzer, M. Strikman, and C. Weiss, JHEP 1301, 163 (2013)

J. Collins and T. Rogers, Phys. Rev. D91 (2015) 074020, [1412.3820].

$$\text{TMDs} \rightarrow \frac{1}{b_T^\alpha} e^{-mb_T} \quad \text{as} \quad b_T \rightarrow \infty.$$

We consider $\log(M_D) \underset{b_T \rightarrow \infty}{\sim} -C b_T + o(b_T)$

Non-

perturbative
function



Mass parameter

Hypotheses for gK

We consider $\log(M_D) \underset{b_T \rightarrow \infty}{\sim} -C b_T + o(b_T)$

**TMD
model**

$$\tilde{D}(b_T, \zeta) = \tilde{D}(b_T, \zeta_0) \exp \left\{ -\frac{g_K}{4} \log \left(\frac{\zeta}{\zeta_0} \right) \right\} (\dots)$$

In general

Hypotheses for gK

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In general

$$\log \left(\tilde{D}(b_T, \zeta) \right) \underset{b_T \rightarrow \infty}{\sim} -C b_T - \underbrace{\frac{g_K^{\text{large } b_T}}{4} \log \left(\frac{\zeta}{\zeta_0} \right)}_{g_K(b_T) = o(b_T)} + o(b_T)$$

**With TMD
model**

$$g_K(b_T) = o(b_T)$$

Hypotheses for gK

**Asymptotic behaviour of TMD
preserved under evolution**



$$\tilde{D}(b_T, \zeta) = \tilde{D}(b_T, \zeta_0) \exp \left\{ -\frac{g_K}{4} \log \left(\frac{\zeta}{\zeta_0} \right) \right\} (\dots)$$

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**With TMD
model**

$$g_K(b_T) = o(b_T)$$

Models for M_D

$$M_D = \frac{2^{2-p}(b_T M_0)^{p-1}}{\Gamma(p-1)} K_{p-1}(b_T M_0) \times F(b_T, z_h)$$

We consider two models; power law as starting point

ID	M_D model	parameters
I	$F = \left(\frac{1 + \log(1 + (b_T M_z)^2)}{1 + (b_T M_z)^2} \right)^q$ $M_z = -M_1 \log(z_h)$	M_0, M_1 $p = 1.51, q = 8$
II	$F = 1$ $M_z = M_h \frac{1}{z f(z)^2} \sqrt{\frac{3}{1 - f(z)}}$ $p_z = 1 + \frac{3}{2} \frac{f(z)}{1 - f(z)}$ $f(z) = 1 - (1 - z)^\beta, \quad \beta = \frac{1 - z_0}{z_0}$	z_0

Two different approaches to model z_h -dependence

Models for gK

g_K model		
A	$g_K = \log (1 + (b_T M_K)^{p_K})$	M_K, p_K
B	$g_K = M_K b_T^{(1-2p_K)}$	M_K, p_K

**I focus on gK since it can be
compared to other extractions.**

**Recall that the TMDFF in this case is differs
from the usual definition by a soft factor.**

Goodness of Fit

$q_T/Q < 0.15$ (pts = 168)		
	IA	IB
$\chi^2_{\text{d.o.f.}}$	1.25	1.19
$M_0(\text{GeV})$	$0.300^{+0.075}_{-0.062}$	$0.003^{+0.089}_{-0.003}$
$M_1(\text{GeV})$	$0.522^{+0.037}_{-0.041}$	$0.520^{+0.027}_{-0.040}$
p^*	1.51	1.51
q^*	8	8
$M_K(\text{GeV})$	$1.305^{+0.139}_{-0.146}$	$0.904^{+0.037}_{-0.086}$
p_K^*	0.609	0.229

$q_T/Q < 0.15$ (pts = 168)		
	IIA	IIB
$\chi^2_{\text{d.o.f.}}$	1.35	1.33
z_0	$0.574^{+0.039}_{-0.041}$	$0.556^{+0.047}_{-0.051}$
$M_K(\text{GeV})$	$1.633^{+0.103}_{-0.105}$	$0.687^{+0.114}_{-0.171}$
p_k	$0.588^{+0.127}_{-0.141}$	$0.293^{+0.047}_{-0.038}$

**Comparison of all
models similar
(within error bands)**

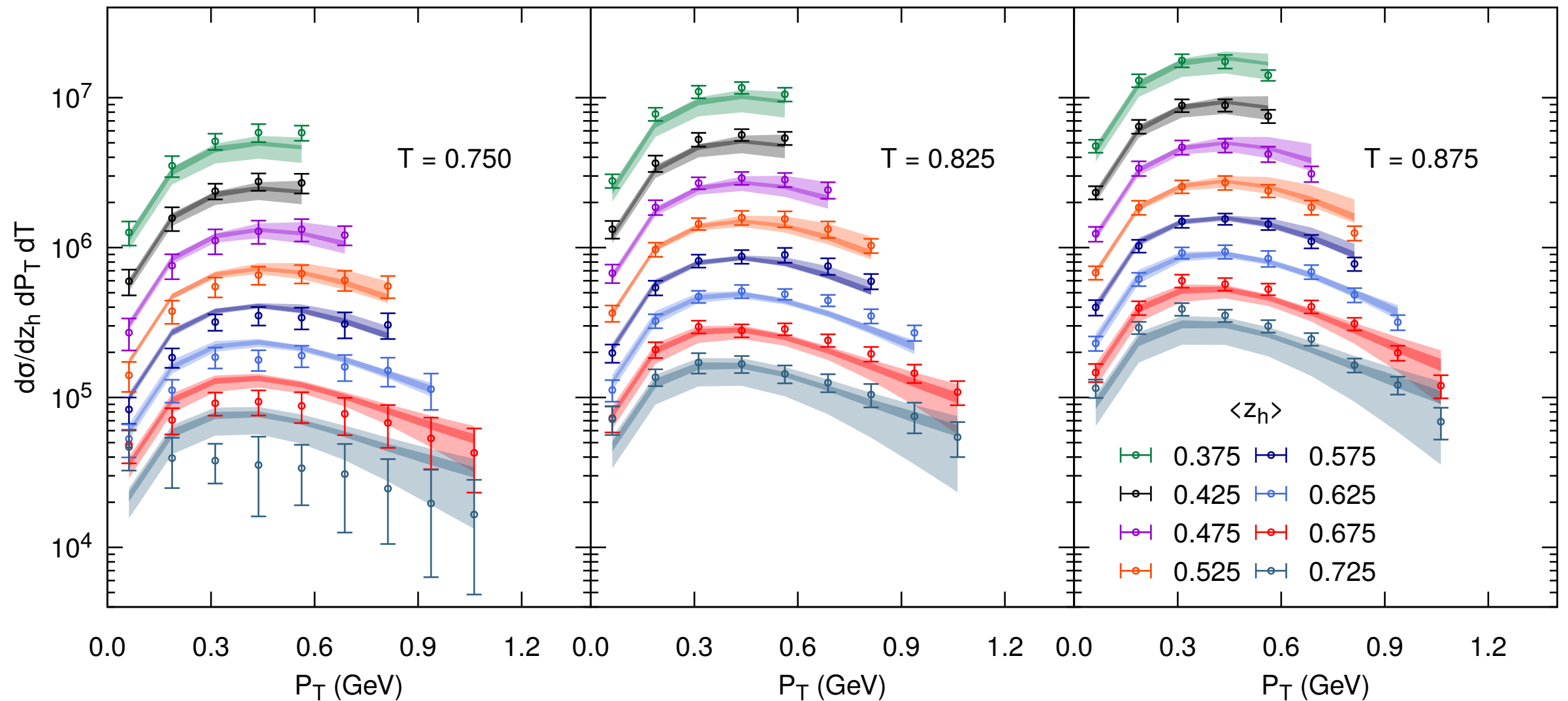
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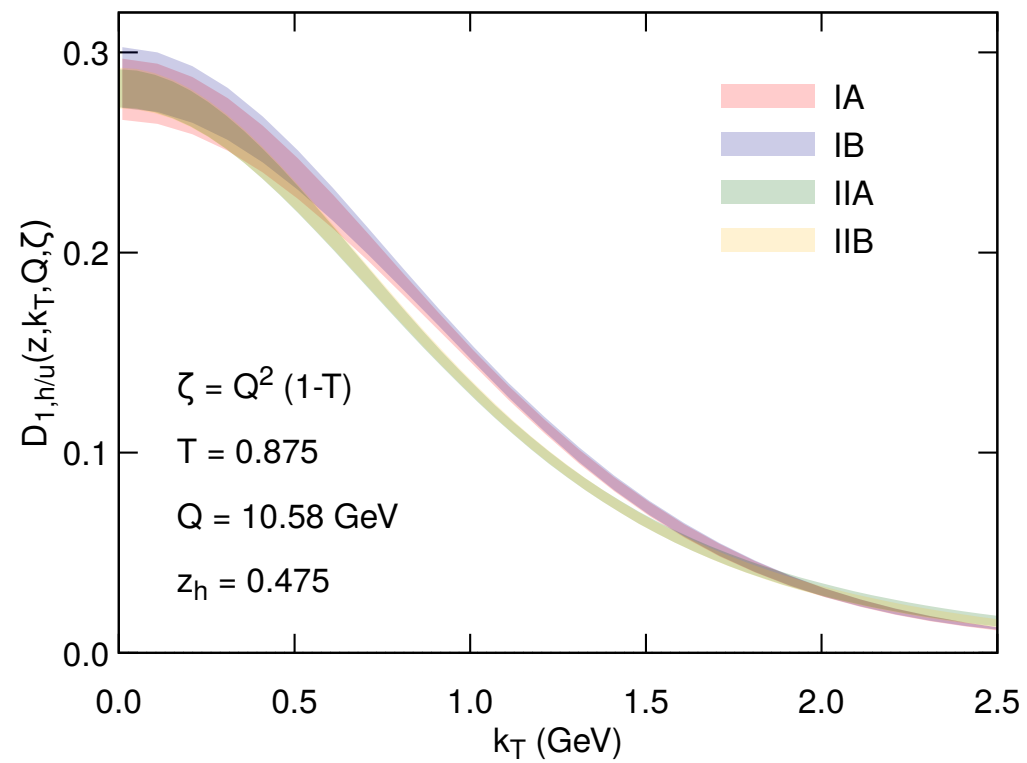
Comparison to data



Dark bands:
statistical uncertainties
from extraction

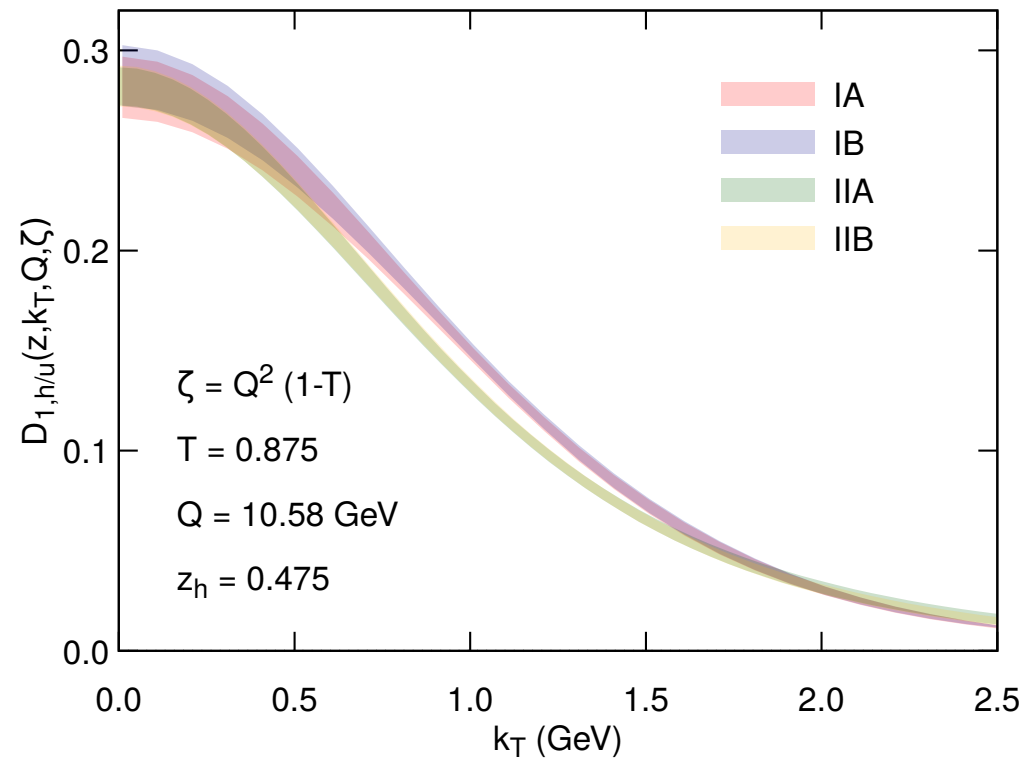
Light bands:
rough estimate of
uncertainty due to collinear
function uncertainties
(probably an overestimation)

Extracted functions & correlations

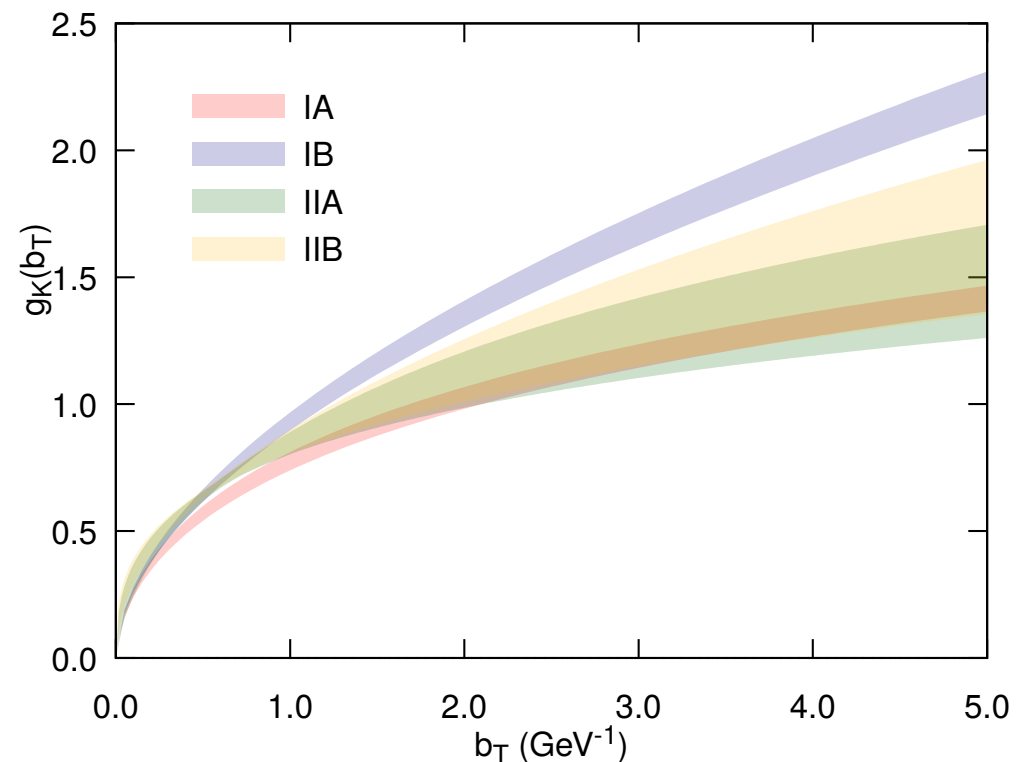
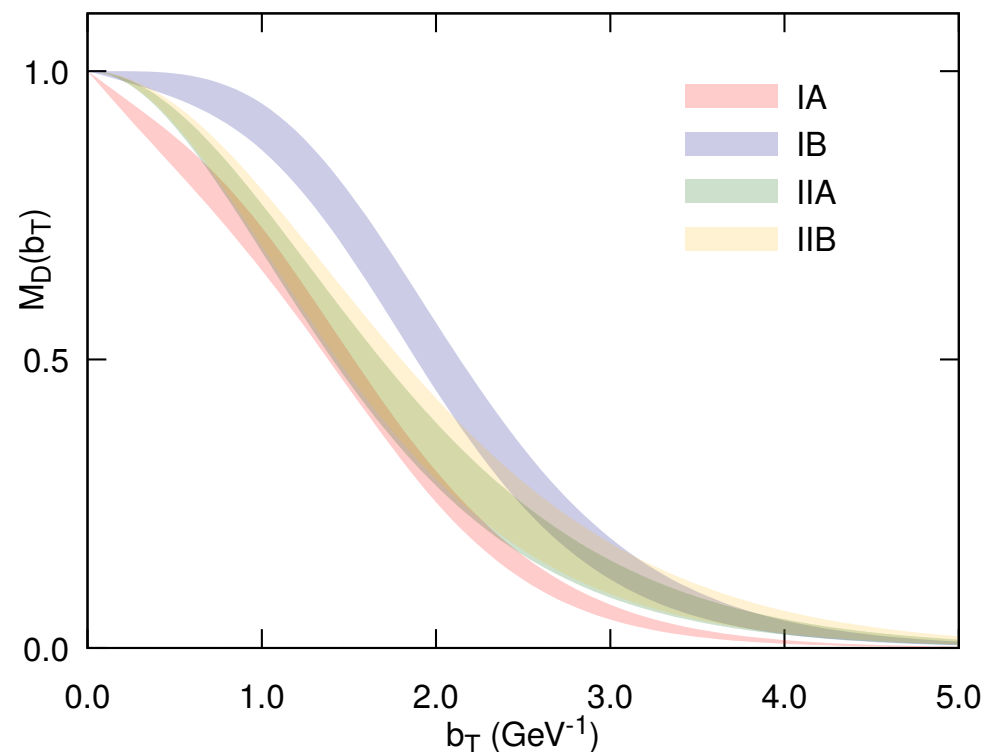


TMDFF:
contains both functions
extracted: MD and gK

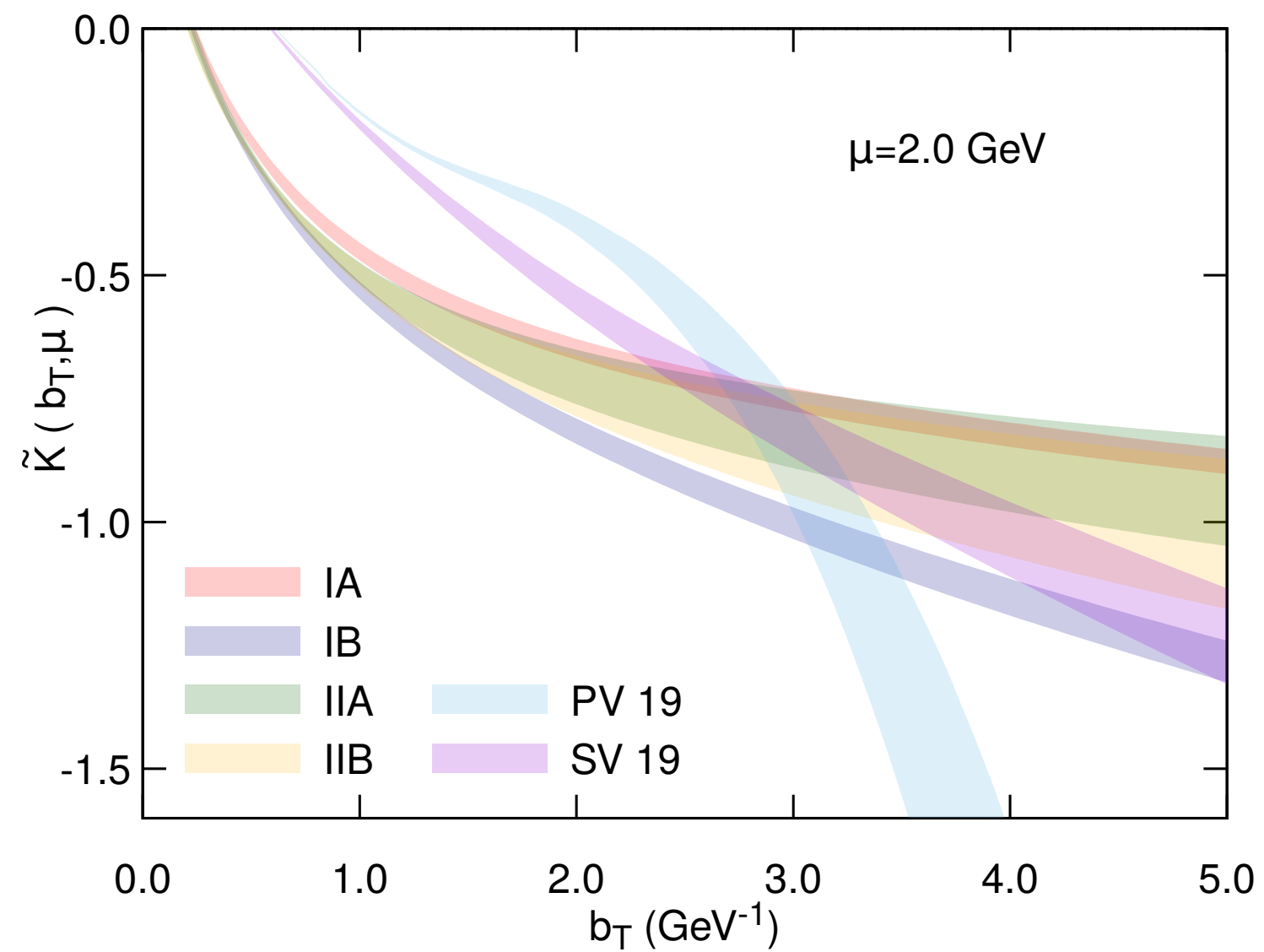
Extracted functions & correlations



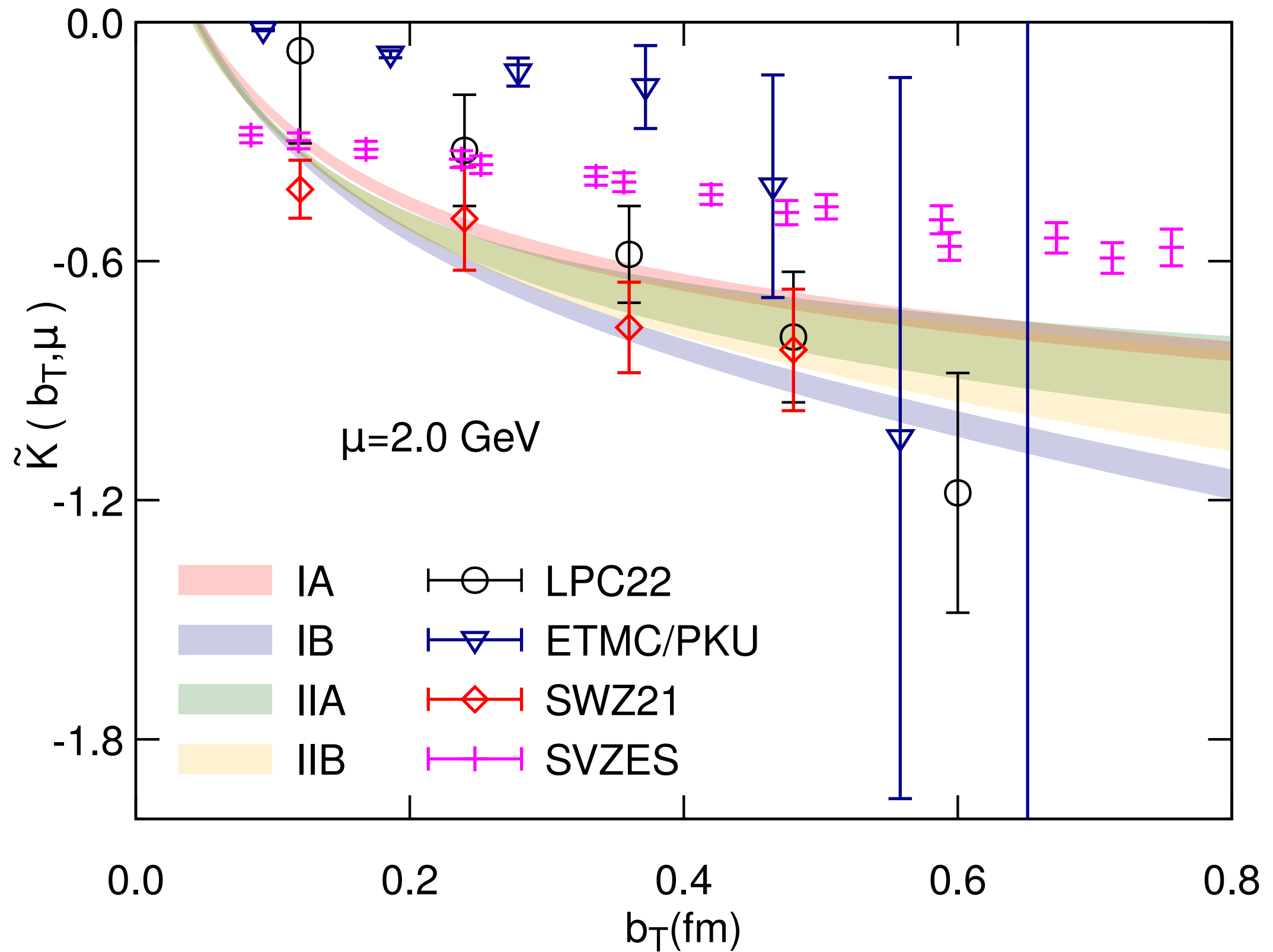
**MD and gK differences
beyond statistical error bands.
Can be thought as a type of
theoretical error (model bias) .**



Comparison to other extractions



Comparison to lattice calculations



Thanks.