

Non-perturbative input in the Parton Branching Evolution

L. Keersmaekers, F. Hautmann, A. Lelek, M. Mendizabal

September 29, 2022

Parton Branching equations

Parton branching equations for TMDs:

$$\begin{aligned}\tilde{\mathcal{A}}_a(x, k_\perp, \mu^2) &= \Delta_a(\mu^2) \tilde{\mathcal{A}}_a(x, k_\perp, \mu_0^2) + \sum_b \int \frac{d^2 \mu_\perp}{\pi \mu'^2} \frac{\Delta_a(\mu^2)}{\Delta_a(\mu'^2)} \Theta(\mu^2 - \mu'^2) \Theta(\mu'^2 - \mu_0^2) \times \\ &\quad \times \int_x^{z_M} dz P_{ab}^R(z) \tilde{\mathcal{A}}_b\left(\frac{x}{z}, k_\perp + (1-z)\mu_\perp, \mu'^2\right)\end{aligned}$$

- AO condition: $q_\perp^2 = (1-z)^2 \mu'^2$

Parton Branching equations

Parton branching equations for TMDs:

$$\begin{aligned}\tilde{\mathcal{A}}_a(x, k_\perp, \mu^2) &= \Delta_a(\mu^2) \tilde{\mathcal{A}}_a(x, k_\perp, \mu_0^2) + \sum_b \int \frac{d^2 \mu_\perp}{\pi \mu'^2} \frac{\Delta_a(\mu^2)}{\Delta_a(\mu'^2)} \Theta(\mu^2 - \mu'^2) \Theta(\mu'^2 - \mu_0^2) \times \\ &\quad \times \int_x^{z_M} dz P_{ab}^R(z) \tilde{\mathcal{A}}_b\left(\frac{x}{z}, k_\perp + (1-z)\mu_\perp, \mu'^2\right)\end{aligned}$$

- AO condition: $q_\perp^2 = (1-z)^2 \mu'^2$
- Resolution scale z_M : resolvable $z < z_M$ and non-resolvable $z > z_M$ branchings

Parton Branching equations

Parton branching equations for TMDs:

$$\begin{aligned}\tilde{\mathcal{A}}_a(x, k_\perp, \mu^2) &= \Delta_a(\mu^2) \tilde{\mathcal{A}}_a(x, k_\perp, \mu_0^2) + \sum_b \int \frac{d^2\mu_\perp}{\pi\mu'^2} \frac{\Delta_a(\mu^2)}{\Delta_a(\mu'^2)} \Theta(\mu^2 - \mu'^2) \Theta(\mu'^2 - \mu_0^2) \times \\ &\quad \times \int_x^{z_M} dz P_{ab}^R(z) \tilde{\mathcal{A}}_b\left(\frac{x}{z}, k_\perp + (1-z)\mu_\perp, \mu'^2\right)\end{aligned}$$

- AO condition: $q_\perp^2 = (1-z)^2\mu'^2$
- Resolution scale z_M : resolvable $z < z_M$ and non-resolvable $z > z_M$ branchings
 - Fixed z_M (set1/set2)
 - Dynamical $z_M = 1 - q_0/\mu'$
 q_0 smallest emitted transverse momentum

Parton Branching equations

Parton branching equations for TMDs:

$$\begin{aligned}\tilde{\mathcal{A}}_a(x, k_\perp, \mu^2) &= \Delta_a(\mu^2) \tilde{\mathcal{A}}_a(x, k_\perp, \mu_0^2) + \sum_b \int \frac{d^2\mu_\perp}{\pi\mu'^2} \frac{\Delta_a(\mu^2)}{\Delta_a(\mu'^2)} \Theta(\mu^2 - \mu'^2) \Theta(\mu'^2 - \mu_0^2) \times \\ &\quad \times \int_x^{z_M} dz P_{ab}^R(z) \tilde{\mathcal{A}}_b\left(\frac{x}{z}, k_\perp + (1-z)\mu_\perp, \mu'^2\right)\end{aligned}$$

- AO condition: $q_\perp^2 = (1-z)^2\mu'^2$
- Resolution scale z_M : resolvable $z < z_M$ and non-resolvable $z > z_M$ branchings
 - Fixed z_M (set1/set2)
 - Dynamical $z_M = 1 - q_0/\mu'$
 q_0 smallest emitted transverse momentum
- Implicit in $P_{ab}^R(z)$ and $\Delta_a(\mu^2)$: $\alpha_s(\mu')$ (set1)/ $\alpha_s(q_\perp)$ (set2)

- Starting distribution:

$$\tilde{A}_a(x, k_{\perp,0}, \mu_0^2) = \tilde{f}_a(x, \mu_0^2) \cdot \frac{1}{q_s^2 \pi} \exp\left(-\frac{k_{\perp,0}^2}{q_s^2}\right)$$

- Starting distribution:

$$\tilde{A}_a(x, k_{\perp,0}, \mu_0^2) = \tilde{f}_a(x, \mu_0^2) \cdot \frac{1}{q_s^2 \pi} \exp\left(-\frac{k_{\perp,0}^2}{q_s^2}\right)$$

Collinear starting distribution $\tilde{f}_a(x, \mu_0^2)$ not focus of this talk
Study influence of intrinsic $k_{\perp,0}$, distributed according to a
Gaussian with width q_s

The non-perturbative input

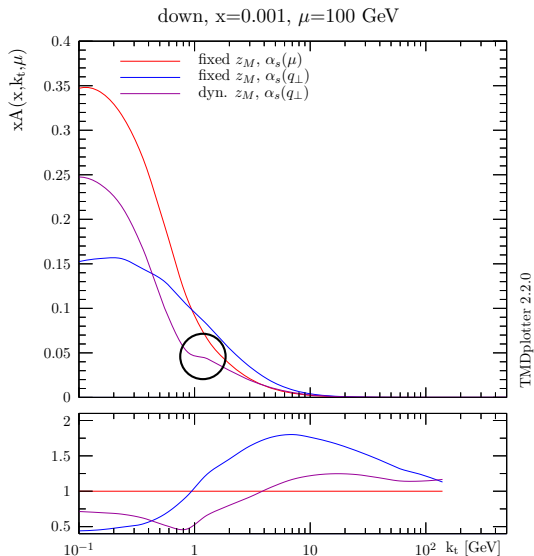
- Starting distribution:

$$\tilde{A}_a(x, k_{\perp,0}, \mu_0^2) = \tilde{f}_a(x, \mu_0^2) \cdot \frac{1}{q_s^2 \pi} \exp\left(-\frac{k_{\perp,0}^2}{q_s^2}\right)$$

Collinear starting distribution $\tilde{f}_a(x, \mu_0^2)$ not focus of this talk
Study influence of intrinsic $k_{\perp,0}$, distributed according to a Gaussian with width q_s

- Resolution scale z_M
Dynamical $z_M = 1 - q_0/\mu'$

TMD with dynamical resolution scale



TMD with dynamical $z_M \rightarrow$ bump around $k_\perp = q_0 = 1$ GeV

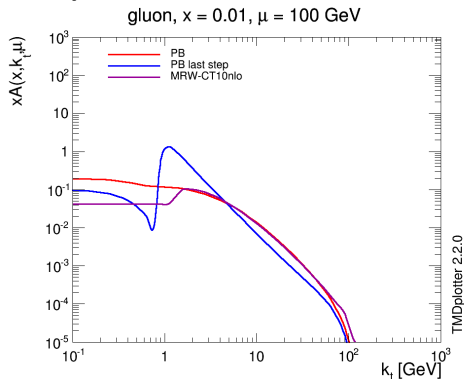
Single emission evolution

Standard PB: $\mathbf{k} = \mathbf{k}_0 - \sum_{i=1}^n \mathbf{q}_i$

PB last step: toy model with $\mathbf{k} = \mathbf{k}_0 - \mathbf{q}_n$

$q_{\perp} > q_0 \rightarrow$ part from evolution starts to accumulate around q_0

Only one branching $\rightarrow$ large bump



Single emission evolution

Standard PB: $\mathbf{k} = \mathbf{k}_0 - \sum_{i=1}^n \mathbf{q}_i$

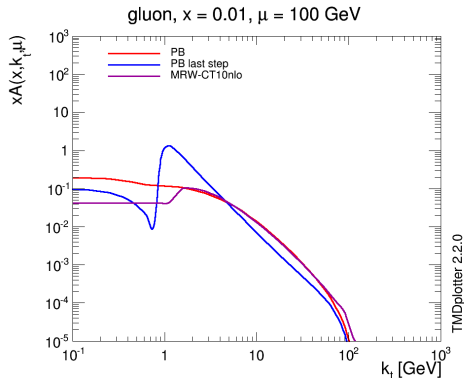
PB last step: toy model with $\mathbf{k} = \mathbf{k}_0 - \mathbf{q}_n$

$q_{\perp} > q_0 \rightarrow$ part from evolution starts to accumulate around q_0

Only one branching $\rightarrow$ large bump

Vector sum $\rightarrow$ with multiple branchings we can reach $k_{\perp} < q_0$

$\rightarrow$ smooth out bump



Single emission evolution

Standard PB: $\mathbf{k} = \mathbf{k}_0 - \sum_{i=1}^n \mathbf{q}_i$

PB last step: toy model with $\mathbf{k} = \mathbf{k}_0 - \mathbf{q}_n$

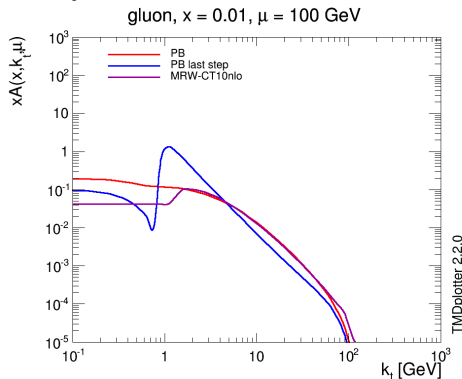
$q_{\perp} > q_0 \rightarrow$ part from evolution starts to accumulate around q_0

Only one branching $\rightarrow$ large bump

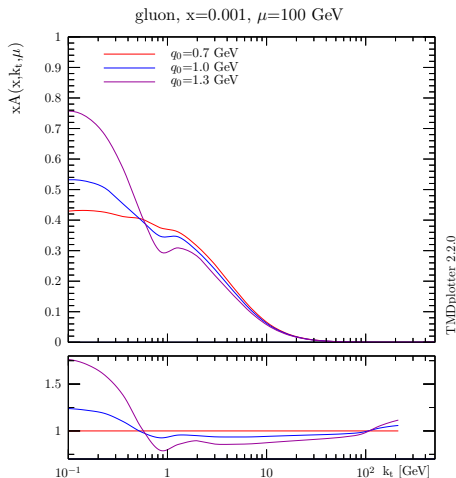
Vector sum $\rightarrow$ with multiple branchings we can reach $k_{\perp} < q_0$
 $\rightarrow$ smooth out bump

Often not enough branchings to completely smooth out, small bump left over

(Not visible here because of log-scale)



Different values of q_0



$$q_s = 0.5 \text{ GeV}$$

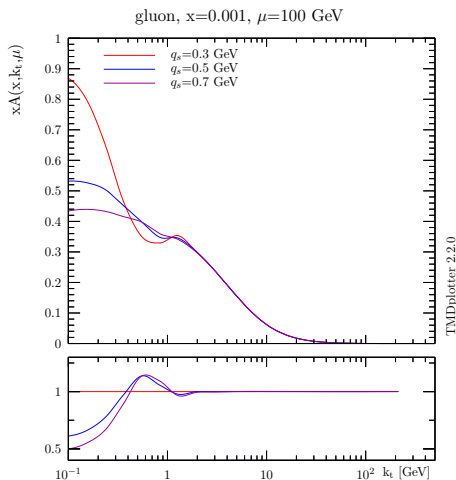
Bump smoother with small values of q_0 :

- $q_\perp > q_0$, with smaller q_0 more branchings are allowed
- More overlap between peak of intrinsic $k_\perp$ and peak around q_0

Different values of q_s

$$q_0 = 1 \text{ GeV}$$

- Smoother with q_s close to q_0 (more overlap peaks)
- **Almost no effect on $k_{\perp}$ -tail** (from ± 2 GeV on)



Different values of q_s

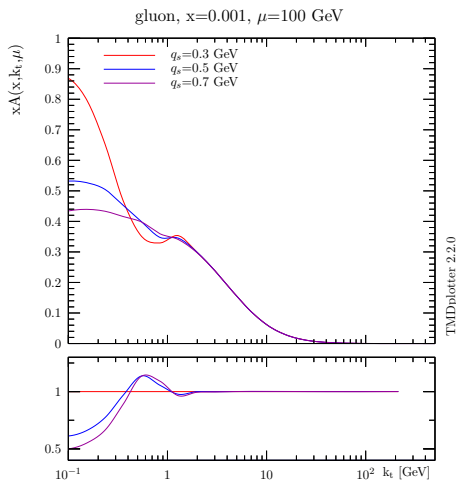
$$q_0 = 1 \text{ GeV}$$

- Smoother with q_s close to q_0 (more overlap peaks)
- **Almost no effect on $k_{\perp}$ -tail** (from ± 2 GeV on)

Vector sum

$$\mathbf{k} = \mathbf{k}_0 - \sum_{i=1}^n \mathbf{q}_i:$$

Random angle between intrinsic $k_{\perp,0}$ and part from evolution $\rightarrow$ only small influence from $k_{\perp,0}$ especially when $k_{\perp}$ from evolution is smooth



Different values of q_s

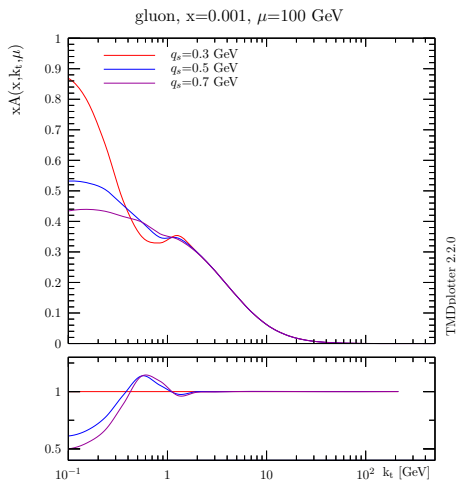
$$q_0 = 1 \text{ GeV}$$

- Smoother with q_s close to q_0 (more overlap peaks)
- **Almost no effect on $k_\perp$ -tail** (from ± 2 GeV on)

Vector sum

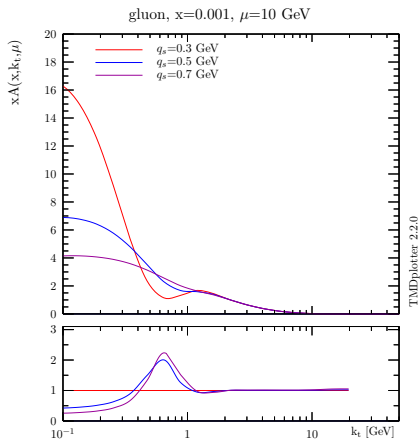
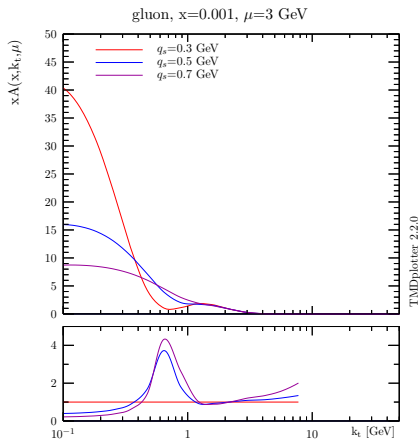
$$\mathbf{k} = \mathbf{k}_0 - \sum_{i=1}^n \mathbf{q}_i$$

Random angle between intrinsic $k_{\perp,0}$ and part from evolution $\rightarrow$ only small influence from $k_{\perp,0}$ especially when $k_\perp$ from evolution is smooth



Since $k_\perp$ -tail unaffected and $f_a = \int dk_\perp^2 A_a$ independent of q_s , total amount of partons below 2GeV unaffected by q_s

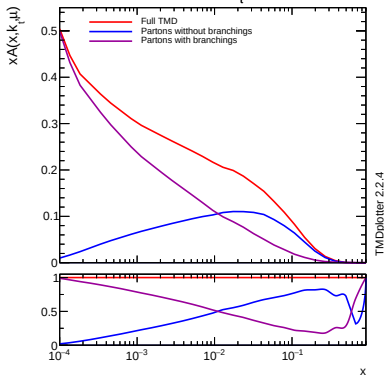
Different values of q_s at low scales



Contributions from partons **with** or **without** branchings

Low $k_{\perp}$ region

gluon, $\mu = 100$ GeV, $k_t = 0.1$ GeV



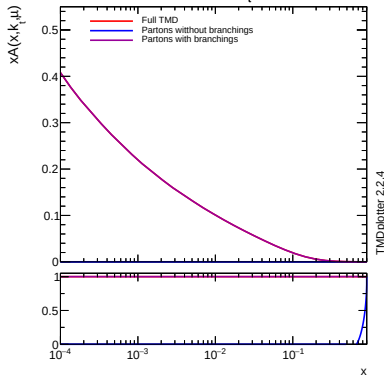
At small x : mostly partons that had branchings

At large x : mostly partons that have not branched

Effects from intrinsic $k_{\perp}$ at small $k_{\perp}$ can be important in a large region of x

$k_{\perp} = 2$ GeV

gluon, $\mu = 100$ GeV, $k_t = 2$ GeV



Dominated in whole x -region by partons that had branchings

Intrinsic $k_\perp$ important in TMD distribution for $k_\perp < 2$ GeV, but total amount of partons with $k_\perp < 2$ GeV constant:

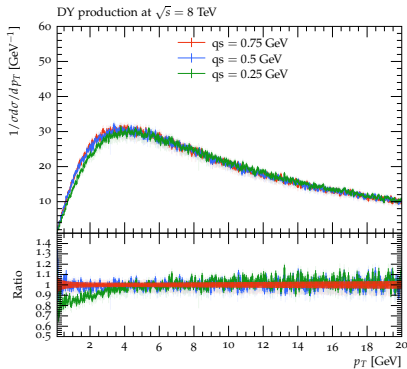
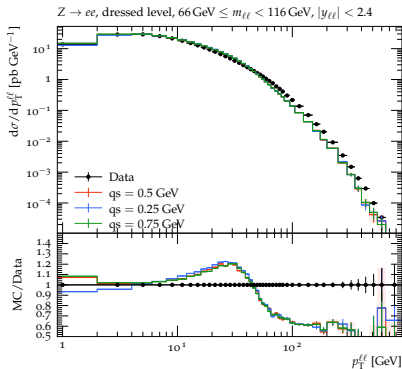
- Effects of intrinsic $k_\perp$ in $p_\perp$ -data might be difficult to see when bin sizes are larger than 2 GeV
- We think data will be most sensitive to intrinsic $k_\perp$ in small $p_\perp$ region:

Only contributions from two partons with

- $k_{\perp,1}$ and $k_{\perp,2}$ both small
- $k_{\perp,1} \approx k_{\perp,2}$ AND $\phi \approx 180^\circ$

Drell Yan with $q_0=1$

Colours do not correspond on both figures



DY at LHC (Atlas), right: fine binning (temporary result)

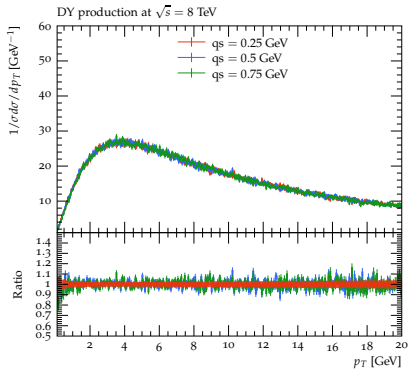
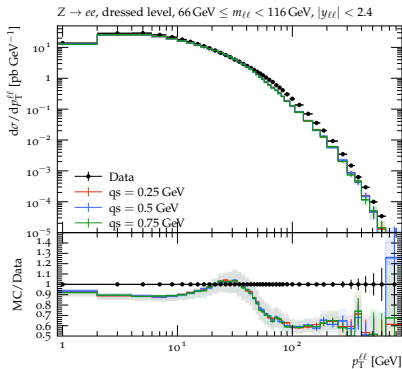
With $q_s=0.25$ GeV (blue left, green right) we can see the effect of q_s at DY spectrum at LHC

Differences between $q_s=0.5$ and $q_s=0.75$ are small

Largest differences at small $p_{\perp} < 4$ GeV, but also differences at higher $p_{\perp} \sim 20$ GeV

Drell Yan with $q_0=0.5$

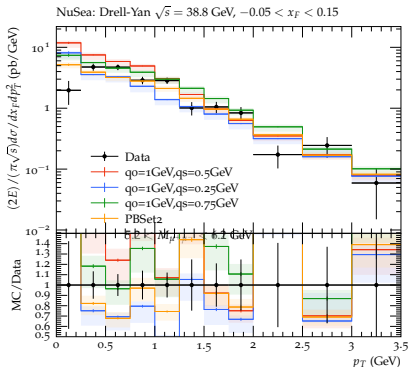
With $q_0=0.5$ more branchings $\rightarrow$ less sensitivity to q_s



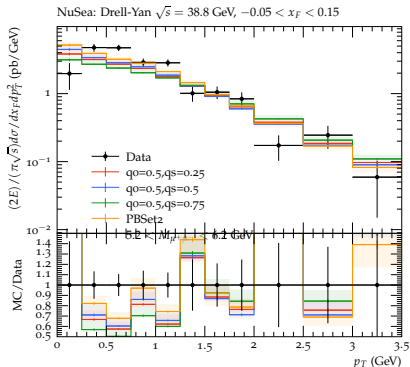
We need higher statistics to see whether we can find some differences with finer binning

Low mass DY

$q_0 = 1$



$q_0 = 0.5$

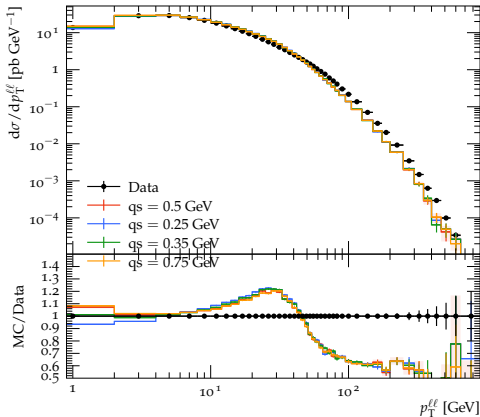


Low mass DY more sensitive to q_s :
Small binnings and less branchings

Set of parameters

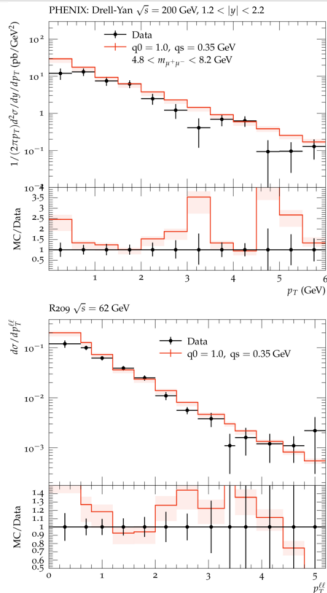
$q_0 = 1 \text{ GeV}$, $q_s = 0.35 \text{ GeV}$ (below green)

$Z \rightarrow e\bar{e}$, dressed level, $66 \text{ GeV} \leq m_{\ell\bar{\ell}} < 116 \text{ GeV}$, $|y_{\ell\bar{\ell}}| < 2.4$



A set of parameters might describe one experiment best, but might describe another experiment worse

$\Rightarrow$ Further studies needed for best parameters



- The distribution at low $k_{\perp}$ comes from both intrinsic $k_{\perp}$ and the effect of multiple branchings in the evolution
- Intrinsic $k_{\perp}$ affects only small $k_{\perp}$ -region (below 2 GeV)
- In certain cases, low $p_{\perp}$ -region of DY at LHC can have some sensitivity to intrinsic $k_{\perp}$, but more sensitivity in low mass DY
- Further studies needed for values of q_0 , q_s